Heavy Quarks above the top

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Interplay of Collider and Flavour Physics 2<sup>nd</sup> meeting, CERN March 16, 2009

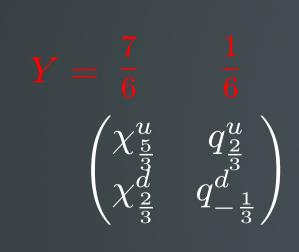
Atre, Carena, Han, J.S., Phys. Rev. D (in press) [arXiv:0806.966]

## The Big Question

- How is electroweak symmetry breaking realized in nature?
  - The answer is around the corner!!
  - A very good candidate is a composite Higgs
    - Compositeness can explain the "large" Hierarchy
    - PNGB nature can explain the "little" Hierarchy
  - Calculable models from warped extra dimensions
    - EWSB, EWPT, Flavour "reasonably" under control (difficult to escape few % fine-tuning in minimal models)

## Realistic Composite Higgs

- A composite Higgs suggests a composite top
  - Large corrections to T and Zbb unless protected
  - A natural framework can be obtained with custodial symmetry and fermions in bidoublets



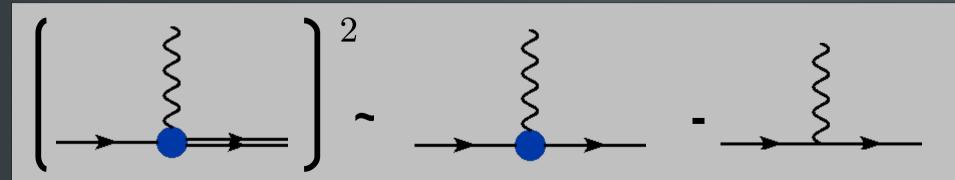
Agashe, Delgado, May, Sundrum 03 Agashe, Contino, Da Rold, Pomarol 06

 $\begin{pmatrix} \chi^{u}_{\frac{5}{3}} & q^{u}_{\frac{2}{3}} \\ \chi^{d}_{2} & q^{d}_{1} \end{pmatrix} \sim (2,2) \text{ under } SU(2)_{L} \times SU$  $(2)_{R_{1}}$ 

- New vector-like quarks that mix with SM quarks can be produced:
  - In pairs: QCD production
    - Iarge (but strong suppression at high masses)
    - model independent, no low energy constraints



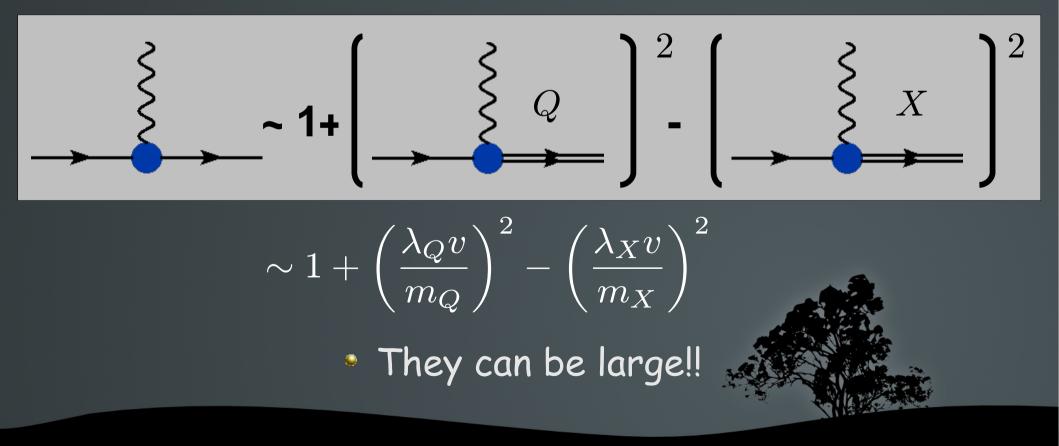
- New vector-like quarks that mix with SM quarks can be produced:
  - Singly: EW production
    - smaller for low masses, larger for heavy quarks
    - model dependent, production cross section constrained by EWPT and flavour experiments



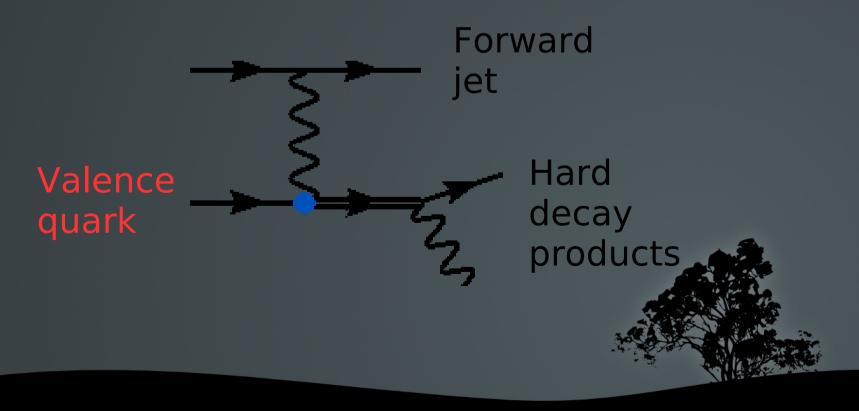
- Can we have large mixing with valence quarks without violating EWPT and Flavour? YES!!
  - Degenerate doublets of hypercharges 1/6 and 7/6 that mix only with u<sub>R</sub>, in the basis of diagonal up Yukawas

$$\begin{split} \mathcal{L} &= \mathcal{L}_{\rm K} - \left[ \lambda_{u}^{i} \bar{q}_{L}^{(0)i} \tilde{\varphi} u_{R}^{(0)i} + \lambda_{d}^{j} V_{ij} \bar{q}_{L}^{(0)i} \varphi d_{R}^{(0)j} \right. \\ &+ \lambda_{Q} \big( \bar{Q}_{L}^{(0)} \tilde{\varphi} + \bar{X}_{L}^{(0)} \varphi \big) u_{R}^{(0)} \\ &+ m_{Q} \big( \bar{Q}_{L}^{(0)} Q_{R}^{(0)} + \bar{X}_{L}^{(0)} X_{R}^{(0)} \big) + \text{h.c.} \end{split}$$

- How can it work?
  - The same cancellation that protects Zbb



- Why single production?
  - Large (unconstrained) coupling to valence quarks
  - Distinctive kinematics



- Model independent analysis
  - Assume two new quarks U, D (charge 2/3, -1/3)
  - Arbitrary couplings to valence quarks

$$\mathcal{L}_{\rm int} = \frac{g}{\sqrt{2}} W^+_{\mu} (\kappa_{uD} \bar{u}_R \gamma^{\mu} D_R + \kappa_{dU} \bar{d}_R \gamma^{\mu} U_R)$$

$$+\frac{g}{2c_W}Z_{\mu}(\kappa_{uU}\bar{u}_R\gamma^{\mu}U_R+\kappa_{dD}\bar{d}_R\gamma^{\mu}D_R)+\text{h.c.}$$

$$\kappa_{qQ} = \tilde{\kappa}_{qQ} \frac{v}{m_Q}$$



• Analysis classified by final state

 $l^{\pm} \not E_T + jj \qquad (W \to l\nu)$   $l^{+}l^{-} + jj \qquad (Z \to l^{+}l^{-})$   $\not E_T + jj \qquad (Z \to \bar{\nu}\nu)$   $(l = e, \mu \text{ here})$ 

- Model independent analysis
  - Parametrize cross section (narrow width approx.)

$$\sigma(pp \to q_1 q_2 f\bar{f}) \equiv S_Q^{CC(NC)} \sigma_{\text{prod}}^{CC(NC)} Br(V \to f\bar{f}),$$

Encode model dependence

Stripped off couplings (depends on kinematics)



- Model independent analysis
  - Parametrize cross section (narrow width approx.)

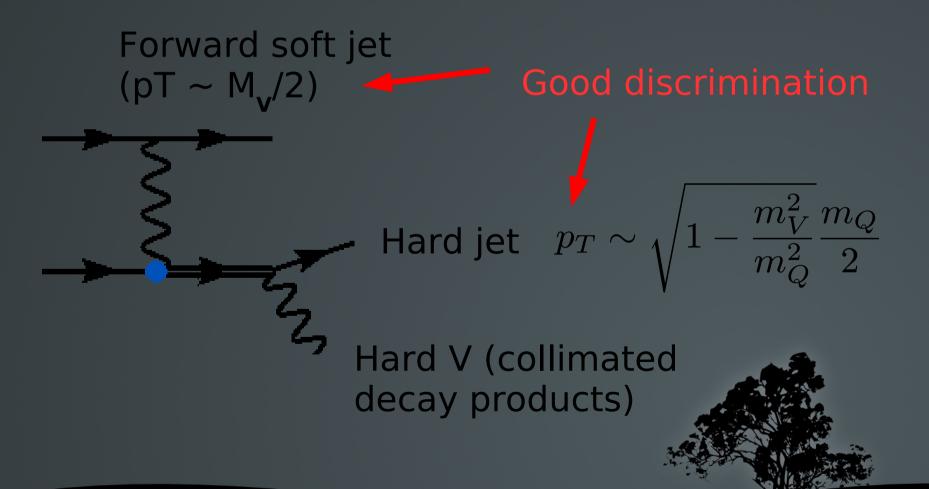
$$\sigma(pp \to q_1 q_2 f\bar{f}) \equiv S_Q^{CC(NC)} \sigma_{\text{prod}}^{CC(NC)} Br(V \to f\bar{f}),$$

$$S_{D}^{CC} \equiv \left(\tilde{\kappa}_{uD}^{2} + \frac{\sigma_{\text{prod}}^{NC}}{\sigma_{\text{prod}}^{CC}} \tilde{\kappa}_{dD}^{2}\right) \quad Br(D \to qW)$$
$$S_{D}^{NC} \equiv \left(\tilde{\kappa}_{dD}^{2} + \frac{\sigma_{\text{prod}}^{CC}}{\sigma_{\text{prod}}^{NC}} \tilde{\kappa}_{uD}^{2}\right) \quad Br(D \to qZ)$$

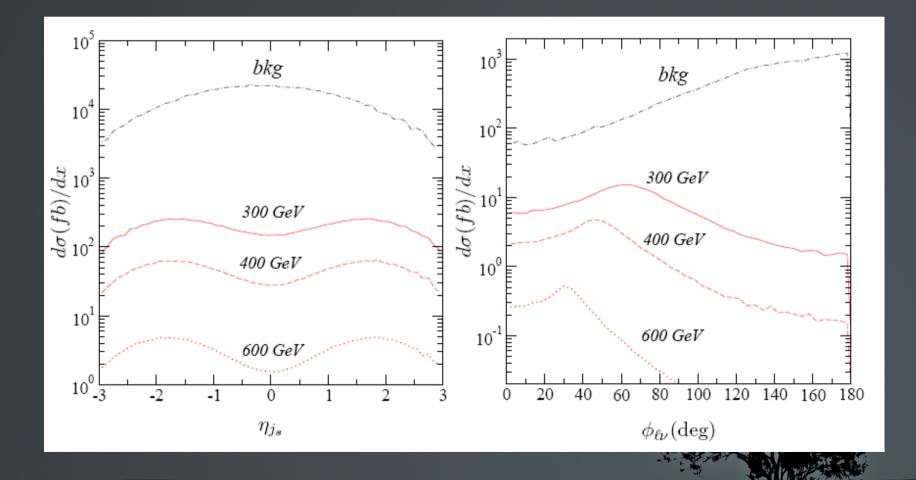
## **Tevatron Analysis**



Main kinematical features



#### Main kinematical features



- Cut based analysis

  - Improved cuts

 $\begin{aligned} p_T(j_h) &> \frac{1}{4} m_Q \quad p_T(W/Z) > \frac{1}{5} m_Q \\ 0.5 &< |\eta_{j_s}| < 3.0 \quad \Delta R(jj, jl) > 1.5, \ 0.8 \\ &+ \text{ cut on } \phi_{l\nu, ll} \end{aligned}$ 

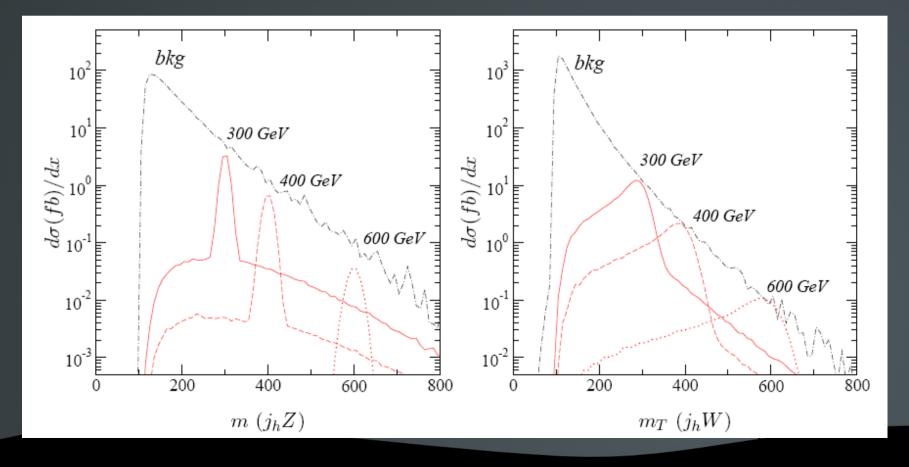
- Cut based analysis
  - Mass reconstruction

 $m_Q - 30 \text{ GeV} < M(j_h Z) < m_Q + 30 \text{ GeV}$ 

 $m_Q - \frac{1}{4}m_Q < M_T(j_h W/Z) < m_Q + 50 \text{ GeV}$ 

$$M_T^2 = \left(\sqrt{p_T^2}_{W,Z} + M_{W,Z}^2 + p_{Tj_h}\right)^2 - (\vec{p}_{TW,Z} + \vec{p}_{Tj_h})^2$$

- Cut based analysis
  - Mass reconstruction

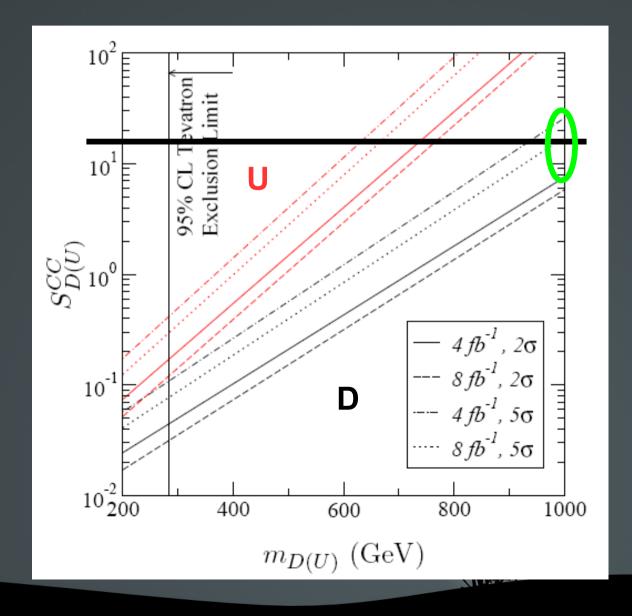


### Simulations:

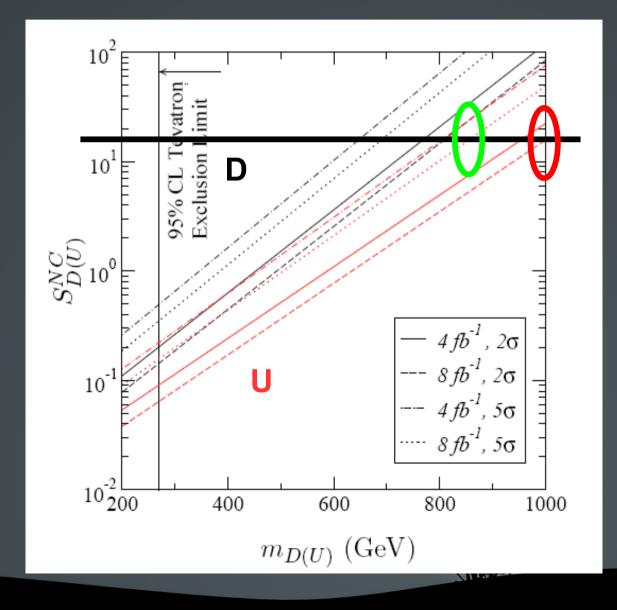
- We have performed a partonic analysis of signal and background with smeared momenta
- Cross-checked at selected points with pythia+PGS
  - $^{ullet}\lesssim7\%$  from hadronization/detector simulation
  - $\sim \lesssim 20\%$  from ISR/FSR + reducible backgrounds



- Results:
  - CC channel



- Results:
  - NC channel



### Conclusions and outlook

- New exciting collider and flavor phenomenology from composite Higgs models
- New quarks with strong mixing to valence quarks allowed by EWPT and Flavour data
- Great reach potential through single production at the Tevatron  $m_Q \lesssim ~{
  m TeV}$
- LHC analysis in progress



# Backup Slides



- Single production and decay
  - \* The spectrum is fixed for degenerate doublets in terms of  $\lambda_Q$  and  $m_Q$

| State             | $q^-$  | $q^+$               | $q^d$                                       | $\chi^u$                                    |
|-------------------|--|---------------------|---|---|
| Electric Charge   | 2/3  | 2/3                 | -1/3  | 5/3   |
| Coupling to $u_R$ |  |                     |   |   |
| CC                |  |                     | $\frac{-g}{\sqrt{2}}\frac{v}{m_Q}\lambda_Q$ | $\frac{-g}{\sqrt{2}}\frac{v}{m_Q}\lambda_Q$ |
| NC                | $\frac{-g}{\sqrt{2}c_W}\frac{v}{m_Q}\lambda_Q$ |                     | v = -;                                      | v –   |
| Yukawa            |  | $\sqrt{2}\lambda_Q$ |   |   |

- How can it work?
  - The same cancellation that protects Zbb
  - Gauge and Yukawa SM couplings get corrections suppressed by the mass of the up quark

$$\begin{split} W^L_{ud_i} \approx V_{ud_i} \left[ 1 - \lambda_Q^2 \frac{v^2}{m_Q^2} \frac{m_u^2}{m_Q^2} \right] & \longrightarrow \sim 10^{-10} \\ W^R_{ud_i} = 0 \end{split}$$

Controls single production

#### Results:

 $m_Q = 400 \text{ GeV}$ 

 $S_Q^{CC} = 1$ 

 $\sigma (\text{in fb}^{-1})$ 

| channels                         | Basic cuts (10) | High $p_T$ (11) | $m_Q$ (12) |
|----------------------------------|-----------------|-----------------|------------|
| $D \to W^{\pm}q$                 | 270             | 190             | 160        |
| $U \to W^{\pm}q$                 | 49              | 35              | 29         |
| $W^{\pm} + 2j$                   | 79000           | 1200            | 280        |
| $W^{\pm}W^{\mp}(\rightarrow 2j)$ | 1500            | 15              | 1.4        |
| $W^{\pm}Z(\rightarrow 2j)$       | 230             | 4.7             | 0.52       |
| single top: $W^{\pm}b j$         | 330             | 10              | 2.9        |
| $t\bar{t}$ : fully leptonic      | 170(79)         | 2.0             | 0.40       |
| $t\bar{t}$ : semi-leptonic       | 600             | 0.19            | -          |

#### Results:

#### $m_Q = 400 \text{ GeV}$

 $S_Q^{CC} = 1$ 

 $\sigma (\text{in fb}^{-1})$ 

| channels                         | Basic cuts $(10)$ | High $p_T$ (11) | $m_Q$ (12) |
|----------------------------------|-------------------|-----------------|------------|
| $D \to Z(\to \ell \ell)q$        | 8.8               | 6.0             | 5.7        |
| $U \to Z (\to \ell \ell) q$      | 22                | 15              | 15         |
| $Z(\to \ell\ell) + 2j$           | 7000              | 120             | 14         |
| $Z(\to \ell\ell)W^{\pm}(\to 2j)$ | 60                | 0.65            | 0.08       |
| $Z(\to \ell\ell)Z(\to 2j)$       | 55                | 1.1             | 0.11       |
| $t\bar{t}$ : fully leptonic      | 160 (1.7)         | -               | -          |



#### Results:

| $m_Q = 400 \text{ Ge}$ | V |
|------------------------|---|
|------------------------|---|

 $S_Q^{CC} = 1$ 

 $\sigma (\text{in fb}^{-1})$ 

| channels                                   | Basic cuts (10) | High $p_T$ (11) | $m_Q$ (12) |
|--|-----------------|-----------------|------------|
| $D \to Z (\to \nu \nu) q$                  | 31              | 22              | 18         |
| $U \to Z (\to \nu \nu) q$                  | 79              | 56              | 46         |
| $Z(\rightarrow \nu \nu) + 2j$              | 28000           | 630             | 160        |
| $Z(\to \nu\nu)W^{\pm}(\to 2j)$             | 240             | 3.4             | 0.30       |
| $Z(\rightarrow \nu \nu) Z(\rightarrow 2j)$ | 220             | 6.1             | 0.76       |
| $t\bar{t}$ : fully leptonic                | 260(12)         | 1.5             | 0.89       |
| $t\bar{t}$ : semi-leptonic                 | 880 (290)       | 2.3             | 1.1        |



Reach:

| $\int \mathcal{L} dt$        | $4 \text{ fb}^{-1}$ |           | $8 \text{ fb}^{-1}$ |           |
|------------------------------|---------------------|-----------|---------------------|-----------|
| Sensitivity                  | $2\sigma$           | $5\sigma$ | $2\sigma$           | $5\sigma$ |
| $m_D$ for $S_D^{CC} = 1$ (2) | 720 (820)           | 580(670)  | 760 (860)           | 630 (710) |
| $m_U$ for $S_U^{CC} = 1$ (2) | 470 (530)           | 370 (440) | 490 (560)           | 400 (470) |
| $m_D$ for $S_D^{NC} = 1$ (2) | 450 (530)           | 350 (420) | 490 (570)           | 380 (470) |
| $m_U$ for $S_U^{NC} = 1$ (2) | $590 \ (680)$       | 460 (540) | 640 (730)           | 510 (590) |

TABLE V: Tevatron sensitivity for  $m_{D,U}$  (GeV).