

Patterns of Flavour Violation in a RS Model with Custodial Protection

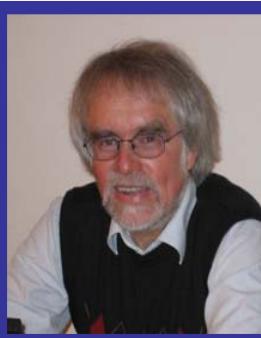
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CERN, March 16th, 2009

Mixing, \mathcal{CP} in RS model [hep-ph/0809.1073]



M. Blanke



AJB



B. Duling



S. Gori



A. Weiler

(50)

Rare K and B Decays in RS model [hep-ph/0812.3803]



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Electroweak and Flavour Structure [hep-ph/0903]



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Impact of KK Fermions on SM fermion couplings [hep-ph/0903]



(31)

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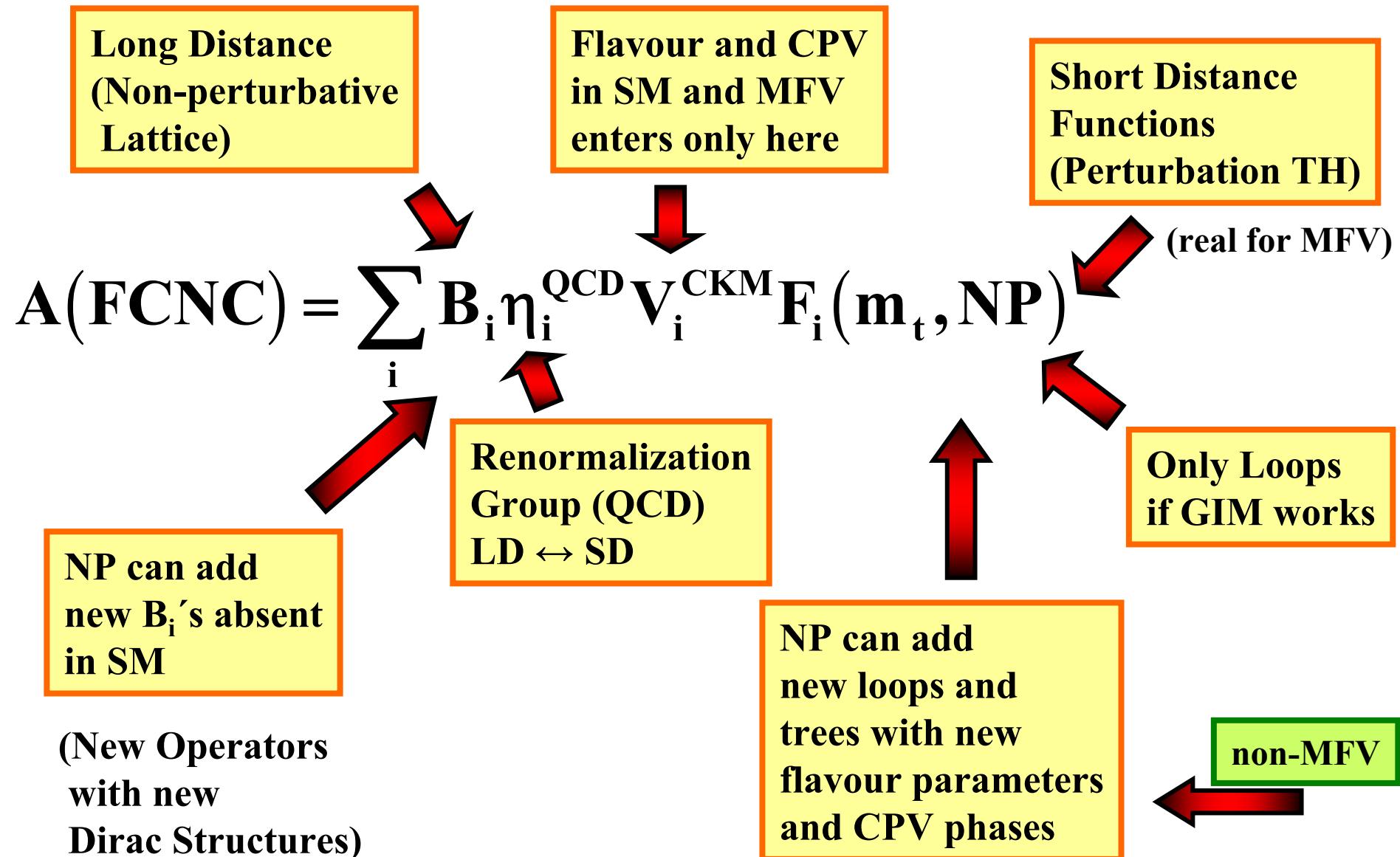
Strategy for the next 19 Minutes

- 1. Going beyond MFV**
- 2. RS Framework**
- 3. A RS-Model with Custodial Protection**
- 4. Patterns of Flavour Violation in 3 Steps**
- 5. Selected Numerical Results**
- 6. Final Messages**

1.

Going beyond MFV

Basic Structure of FCNC Amplitudes



Possible Dirac Structures in $K^0 - \bar{K}^0$ and $B_{d,s}^0 - \bar{B}_{d,s}^0$

SM:

$$\gamma_\mu (1 - \gamma_5) \otimes \gamma^\mu (1 - \gamma_5)$$

LHT

**Strong
enhancements**



Beyond SM:

$$\gamma_\mu (1 - \gamma_5) \otimes \gamma^\mu (1 + \gamma_5)$$

$$\star (1 - \gamma_5) \otimes (1 + \gamma_5)$$

$$(1 - \gamma_5) \otimes (1 - \gamma_5)$$

$$\sigma_{\mu\nu} (1 - \gamma_5) \otimes \sigma^{\mu\nu} (1 - \gamma_5)$$

} RS } SUSY

MSSM with large $\tan\beta$

General Supersymmetric Models

Models with complicated Higgs System; RS Models

NLO $\left[\eta_{\text{QCD}}^i \right]^{\text{New}}$: Ciuchini, Franco, Lubicz,
Martinelli, Scimemi, Silvestrini
AJB, Misiak, Urban, Jäger

Enhancements of Q_{LR} versus Q_{LL} in $\Delta F=2$ Transitions

$$\mu \approx 3\text{TeV} \xrightarrow{\text{RG}} \begin{array}{l} \mu_B \approx 5\text{GeV} \\ \mu_K \approx 2\text{GeV} \end{array}$$

	Wilson Coefficient (RG Enhancement)	Hadronic Matrix Element (Chiral Enhancement)	Total
$K^0 - \bar{K}^0$	~ 7	20	140
$B_{d,s}^0 - \bar{B}_{d,s}^0$	~ 4.3	1.5	6.5

Comparison of Beyond-MFV Scenarios

Scenario	New Flavour and CP Violation	New Operators	FCNC at Tree Level
LHT	★		
SUSY	★	★	
RS	★	★	★

(non-universalities in gauge couplings implied by the manner CKM and mass hierarchies are explained)

General Structure of New Physics Contributions

SM : $\lambda_t^{(K)} = V_{ts}^* V_{td}$ $\lambda_t^{(d)} = V_{tb}^* V_{td}$ $\lambda_t^{(s)} = V_{tb}^* V_{ts}$

Amplitudes

: $\lambda_t^{(i)} X_{\text{SM}}(m_t)$
 $v\bar{v}$ in the final state $\lambda_t^{(i)} Y_{\text{SM}}(m_t)$
 $\mu^+ \mu^-$ in the final state

**Universality
of short
distance
functions**

$i = K, B_d, B_s$

LHT

:

$$X_i = X_{\text{MFV}} + \frac{1}{\lambda_t^{(i)}} \xi_i \bar{X} \equiv |X_i| e^{i\theta_X^i}$$

real complex

RS

**Breakdown
of
Universality**

$$Y_i = Y_{\text{MFV}} + \frac{1}{\lambda_t^{(i)}} \xi_i \bar{Y} \equiv |Y_i| e^{i\theta_Y^i}$$

**New
Flavour
and CP
in ξ_i**

Non-MFV

Natural Expectations

$$X_i = X_{\text{MFV}} + \frac{1}{\lambda_t^{(i)}} \xi_i \bar{X} \equiv |X_i| e^{i\theta_X^i}$$

$i = K, B_d, B_s$

(similarly for Y_i) Non-MFV

$$\frac{1}{\lambda_t^{(K)}} \approx 2 \cdot 10^3$$

$$\frac{1}{\lambda_t^{(d)}} \approx 100$$

$$\frac{1}{\lambda_t^{(s)}} \approx 25$$

$$\left\{ \begin{array}{l} \text{Natural} \\ \text{size} \\ \text{of NP} \\ \text{contributions} \end{array} \right\} : K \gg B_d > B_s$$

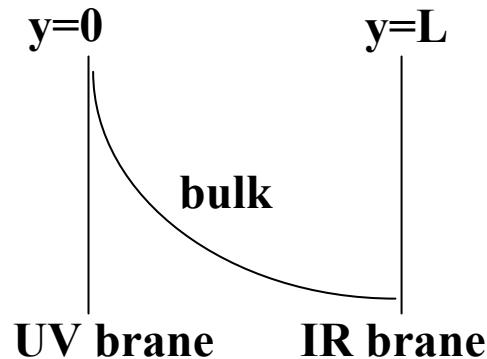
But can be modified for
special structures of ξ_i

2.

Randall-Sundrum Framework (Express Summary)

5D spacetime with warped metric:

$$ds^2 = e^{-2ky} \eta_{\mu\nu} dx^\mu - dy^2 \quad 0 \leq y \leq L$$



- fermions and gauge bosons live in the bulk
- Higgs localised on IR brane

(Chang, Okada et al.
Grossman, Neubert
Gherghetta, Pomarol)

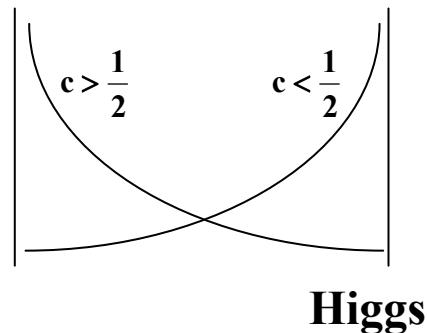
- energy scales suppressed by warp factor e^{-ky}
natural solution to the gauge hierarchy problem.
- Kaluza-Klein (KK) excitations of both SM fermions and gauge bosons live close to the IR brane.

Fermion Localisation and Yukawa Couplings

SM fermion (zero mode) shape function depends strongly on bulk mass parameter characteristic for a given fermion:

$$f^{(0)}(y, c) \propto e^{\left(\frac{1}{2}-c\right)y}$$

UV brane IR brane



$c > \frac{1}{2}$: localisation near UV brane

$c < \frac{1}{2}$: localisation near IR brane

effective 4D Yukawa couplings:

$$(Y_{u,d})_{ij} = (\lambda_{u,d})_{ij} f_i^Q f_j^u$$

- $\lambda_{u,d} \sim 0(1)$ anarchic complex 3×3 matrices $\equiv Y_{5D}$
- hierarchical structure of quark masses and CKM parameters can be naturally generated by exponential suppression of $f^{Q,u,d}$ at IR brane.

Bulk Profiles of SM Gauge Bosons

- Gluons and Photon : flat (protection by Gauge symmetry)
- W^\pm, Z : flat before EWSB
but

distorted near the IR brane after EWSB $\propto \left(\frac{v^2}{M_{KK}^2} \right)$

Equivalently : Mixing of KK gauge bosons with W^\pm, Z in the process of EWSB modifies the couplings of mass eigenstates W^\pm, Z

- Recall : All KK gauge bosons live close to the IR brane

All KK fermions live close to the IR brane

First Implications for Phenomenology

1.

Gauge-Fermion
Interactions:
Overlaps of shape
functions



Non-universalities
in
Gauge Couplings

(in flavour)

of $\left\{ \begin{array}{l} \text{KK-gauge bosons} \\ W^\pm, Z \end{array} \right\}$
to $\{\text{SM fermions}\}$



2.

Impact on
Electroweak Precision
Observables

$SU(2)_L \otimes U(1)_Y$

S parameter : $M_{KK} \geq (2-3) \text{ TeV}$
T parameter: $M_{KK} \geq 10 \text{ TeV}$

Agashe, Delgado, May, Sundrum (2003)
Csaki, Grojean, Pilo, Terning (2003)

Also problems with $Z b_L \bar{b}_L$

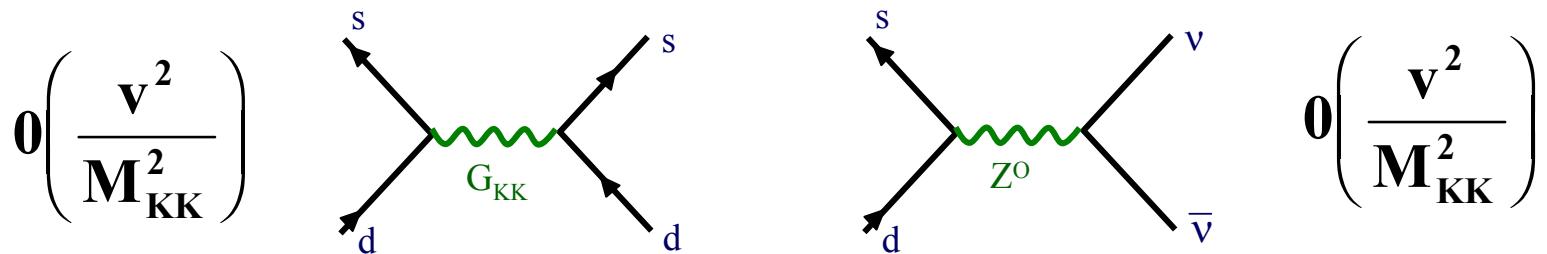
3.

Tree Level FCNC mediated by KK gauge bosons
and Z (breakdown of standard GIM mechanism)

$$\mathbf{d} \equiv \begin{pmatrix} \mathbf{d} \\ \mathbf{s} \\ \mathbf{b} \end{pmatrix}$$

$$\bar{\mathbf{d}} \mathbf{D}_L^+ \begin{pmatrix} \mathbf{a} \\ \mathbf{b} \\ \mathbf{c} \end{pmatrix} \mathbf{D}_L \gamma_\mu \mathbf{Z}^\mu \mathbf{d} \neq \bar{\mathbf{d}} \gamma_\mu \mathbf{Z}^\mu \mathbf{d}$$

(non-universality)



But RS-GIM helps in avoiding disaster.

Gherghetta, Pomarol
Huber, Shafi
Agashe, Soni, Perez

4.

Mixing of KK fermions with SM fermions and
mixing of KK gauge bosons with SM gauge bosons



Breakdown of Unitarity of the CKM matrix

5.

$\left\{ \text{Tree level exchanges of } G_{KK} \text{ and } Z \right\} \rightarrow \left\{ \text{Contributions of new operators. In particular } Q_{LR} \text{ operators in addition to } Q_{LL}, Q_{RR} \right\}$

6.

$\left\{ \text{The presence of three } 3 \times 3 \text{ hermitian bulk matrices } c^q, c^u, c^d \text{ in addition to usual Yukawa couplings} \right\} \rightarrow \left\{ \begin{array}{l} \text{New flavour and CP violating parameters:} \\ \left\{ \begin{array}{ll} 3 * 6 = 18 & \text{real} \\ 3 * 3 = 9 & \text{phases} \end{array} \right. \end{array} \right\}$

Non-MFV

3.

A RS Model with Custodial Protection

$$\mathbf{SU(3)_C \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_X \otimes P_{LR}}$$

Gauge Group in the Bulk

$$P_{LR} : SU(2)_L \leftrightarrow SU(2)_R$$

P_{LR} symmetric fermion representations

What is protected in this Model?

(up to small symmetry breaking due to UV boundry conditions)

A.

T-Parameter

Agashe, Delgado, May, Sundrum (0308036)
Csaki, Grojean, Pilo, Terning (0308038)



B.

$$Z\bar{b}_L b_L$$

Agashe, Contino, Rold, Pomarol (0605341)

C.

$$Z\bar{d}_L^i d_L^j$$

Blanke, AJB, Duling, Gori, Weiler (0809.1073)
Blanke, AJB, Duling, Gemmler, Gori (0812.3803)



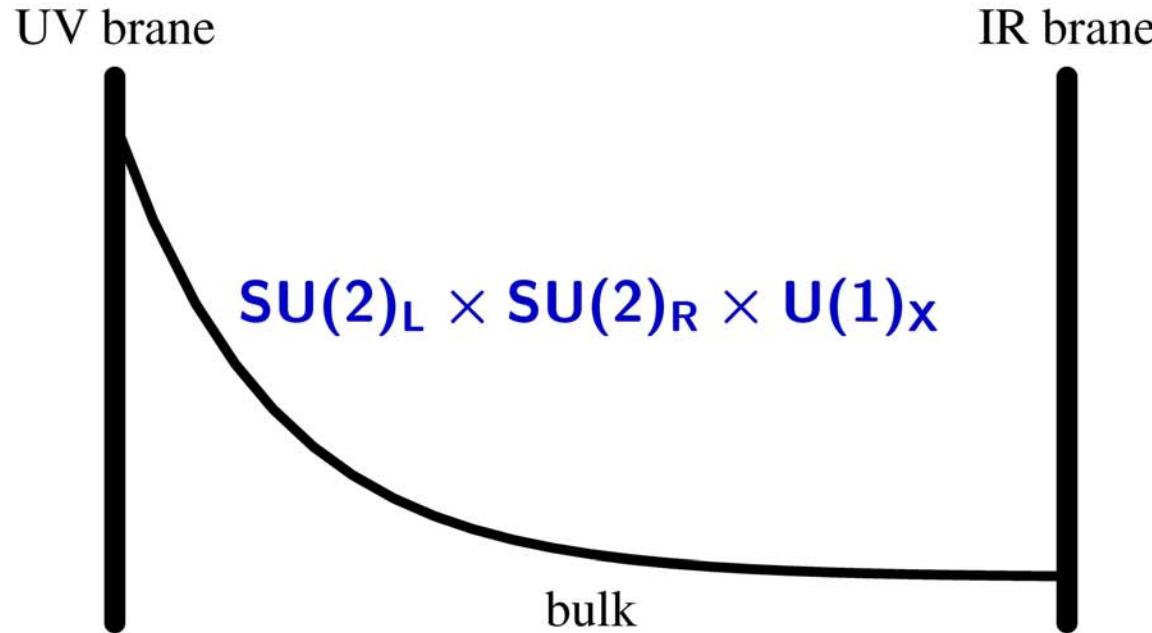
D.

$$Z\bar{u}_R^i u_R^j$$

AJB, Duling, Gori (0903xxx)

But: $Z\bar{d}_R^i d_R^j$, $Z\bar{u}_L^i u_L^j$, $W^+ \bar{u}_L^i d_L^j$ not protected

A Realistic Model in the Reach of the LHC



$$SU(2)_R \times U(1)_X \rightarrow U(1)_Y$$

by boundary conditions

+ $(L \leftrightarrow R)$ -symmetric fermion representations

$$SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$$

by Higgs VEV

low energy theory: $SU(2)_L \times U(1)_Y \rightarrow U(1)_{\text{em}}$

Particle Content of the Model

Albrecht, Blanke, AJB, Duling, Gemmler (0903xxx)

Gauge sector

$$W^\pm, \quad W_H^\pm, \quad W'^\pm \\ Z^0, \quad Z_H, \quad Z'$$

KK

$$A, \quad A^{(\prime)} \\ G^a, \quad G^{a(\prime)}$$

KK

Quark sector
(i=1,2,3)

$$(2,2) = \begin{pmatrix} \chi^{u_i} (-+)_{5/3} & q^{u_i} (++)_{2/3} \\ \chi^{d_i} (-+)_{2/3} & q^{d_i} (++)_{1/3} \end{pmatrix}_L$$

$SU(2)_L \otimes SU(2)_R$

$$(3,1) = \begin{pmatrix} \psi'^i (-+)_{5/3} \\ U'^i (-+)_{2/3} \\ D'^i (-+)_{-1/3} \end{pmatrix}_R \oplus \begin{pmatrix} \psi''^i (-+)_{5/3} \\ U''^i (-+)_{2/3} \\ D^i (++)_{-1/3} \end{pmatrix}_R = (1,3)$$

+
states of
opposite
chirality

$Q=5/3$
Fermions!

(Feynman rules worked out for SM and n=1 KK modes)

Higgs and Mass Matrices

$$H = \begin{pmatrix} \pi^+ / \sqrt{2} & -(\mathbf{h}^0 - i\pi^0) / 2 \\ (\mathbf{h}^0 + i\pi^0) / 2 & \pi^- / \sqrt{2} \end{pmatrix}$$

For SM and n=1 KK modes with flavour i=1,2,3

$$M(5/3) = 9 \times 9$$

$$M(2/3) = 18 \times 18$$

$$M(-1/3) = 12 \times 12$$

a) Numerical Diagonalization

(BB D GW)

b) Analytic Reduction to 3 x 3 by means of effective Lagrangians
(integrating out of KK modes) (AJB, Duling, Gori)

4.

Patterns of Flavour Violation in 3 Steps

First look at $\Delta F = 2$

: Burdman; Agashe, Perez, Soni

First more sophisticated analysis

: Csaki, Falkowski, Weiler (0804.1954)

Application of model-independent results of Utfit group to RS-type models.

**Hierarchy of fermion masses and weak mixings solely due to geometry
 Y_{5D} anarchic and $O(1)$**

KK-Gluon

Contribution to ϵ_k

$M_{KK} \gtrsim 21 \text{ TeV}$

Step 1

: ($\Delta F = 2$ Processes)

(Blanke, AJB, Duling, Gori, Weiler (0809.1073))

- A. Full RG analysis at the NLO level: (using AJB, Misiak, Urban; Jäger) (2000)**

$$Q_1^{VLL} = (\bar{s}\gamma_\mu P_L d)(\bar{s}\gamma^\mu P_L d) \quad Q_1^{LR} = (\bar{s}\gamma_\mu P_L d)(\bar{s}\gamma^\mu P_R d)$$

$$Q_1^{VRR} = (\bar{s}\gamma_\mu P_R d)(\bar{s}\gamma^\mu P_R d) \quad Q_2^{LR} = (\bar{s}P_L d)(\bar{s}P_R d)$$

(For $K^0 - \bar{K}^0$, $B_d^0 - \bar{B}_d^0$ and $B_s^0 - \bar{B}_s^0$ systems)

- B. Inclusion of the contributions of all gauge bosons:**

$(G_{KK}, A_{KK}, Z, Z_H, Z')$ (Protection of Z and Z' pointed out)

$$Z \bar{d}_d^i d_L^j$$

$$Z' \bar{d}_L^i d_L^j$$

- C. Phenomenology of ε_K , ΔM_K , ΔM_S , ΔM_d , $S_{\psi Ks}$, $S_{\psi\phi}$, A_{SL}^q , $\Delta\Gamma_q$**

- D. Relation of RS flavour model to Froggatt-Nielsen (analytic formulae for masses and mixings)**

- E. Calculation of fine tuning (Barbieri + Giudice) of Yukawa couplings $\Delta_{BG}(\varepsilon_K)$ necessary to satisfy ε_K with $M_{KK} \sim 2\text{-}3 \text{ TeV}$**

Main Results of Step 1

- A. Confirmation of CFW analysis for anarchic 5D Yukawa's.
- B. Identifications of regions in parameter space with only modest fine-tuning of Y_{5D} which satisfy all $\Delta F = 2$ constraints, agree with quark masses and mixings and electroweak constraints for $M_{KK} \sim 2\text{-}3 \text{ TeV}$.
- C. Pattern of NP contributions
 - : $\epsilon_K, \Delta M_K$: dominated by Q_2^{LR} and G_{KK}
 - $\Delta M_d, \Delta M_s, S_{\psi K_s}, S_{\psi\phi}$: Competition between Q_1^{VLL} and Q_2^{LR}
 $(Z_H \text{ and } G_{KK} \text{ dominate})$
- D. $S_{\psi\phi}$ asymmetry can be by order of magnitude larger than $(S_{\psi\phi})_{SM}$.

Step 2

: (Rare K and B Decays) ($\Delta F=1$)

(Blanke, AJB, Duling, Gemmeler, Gori (0812.3803))

- A. Calculation of $K^+ \rightarrow \pi^+ \nu \bar{\nu}$, $K_L \rightarrow \pi^0 \nu \bar{\nu}$, $K_L \rightarrow \pi^0 l^+ l^-$,
 $K_L \rightarrow \mu^+ \mu^-$, $B_{s,d} \rightarrow \mu^+ \mu^-$, $B \rightarrow K \nu \bar{\nu}$,
 $B \rightarrow K^* \nu \bar{\nu}$, $B \rightarrow X_{s,d} \nu \bar{\nu}$

For all allowed regions of parameters from Step 1 with

$$\Delta_{BG}(\varepsilon_K) \leq 20$$

- B. Dominance of tree level Z-exchanges but through its right-handed couplings.

- C. Study of correlations between various $\Delta F=1$ branching ratios and of $(\Delta F=1) \leftrightarrow (\Delta F=2)$ correlations.

Main Results of Step 2

- A. Enhancements of $\text{Br}(K_L \rightarrow \pi^0 v\bar{v})$ by a factor 5
 $\text{Br}(K^+ \rightarrow \pi^+ v\bar{v})$ by a factor 2
 $\text{Br}(K_L \rightarrow \pi^0 l^+ l^-)$ by a factor 1.5 } possible
(even simultaneously)
- B. Large Enhancements of $\text{Br}(K_L \rightarrow \mu^+ \mu^-)$ but not simultaneously with $K^+ \rightarrow \pi^+ v\bar{v}$.
- C. SM-like $\text{Br}(B_{s,d} \rightarrow \mu^+ \mu^-)$, $\text{Br}(B \rightarrow K v\bar{v})$
 $\text{Br}(B \rightarrow K^* v\bar{v})$, $\text{Br}(B \rightarrow X_{s,d} v\bar{v})$ } (10-20% effects)
- D. Simultaneous large effects in $S_{\psi\phi}$ and $K \rightarrow \pi v\bar{v}$ not possible.
- E. Non-Universality of NP effects and consequently "golden relations" of MFV can be strongly broken.

Step 3

: Impact of KK fermions
(Effective Lagrangian approach)

(AJB, Duling, Gori (0903.xxx))

1.

General formulae for corrections to
SM fermion $\leftrightarrow (W^\pm, Z, H)$ couplings
from mixing with KK fermions.

2.

Explicit demonstration that the custodial
protection of $Zd_L^i \bar{d}_L^j$ and $Zu_R^i \bar{u}_R^j$
couplings remains valid in the presence of
mixing with KK fermions (guaranteed by
 P_{LR} symmetric fermion representations)

3.

Calculations of KK corrections to
unprotected $Zd_R^i \bar{d}_R^j$ and $Zu_L^i \bar{u}_L^j$

4.

Study of the violation of the unitarity
of the CKM due to KK mixing

+

Comparison
with brute
force numerical
diagonalization
of 18×18 and
 12×12 matices

5.

Selected Numerical Results

Fine Tuning in $\Delta F=2$ Processes (BBDGW)

$$\Delta_{\text{BG}}(\mathbf{Q}) = \max_i \left| \frac{\mathbf{x}_i}{\mathbf{Q}} \frac{\partial \mathbf{Q}}{\partial \mathbf{x}_i} \right| \quad \text{Barbieri + Giudice}$$

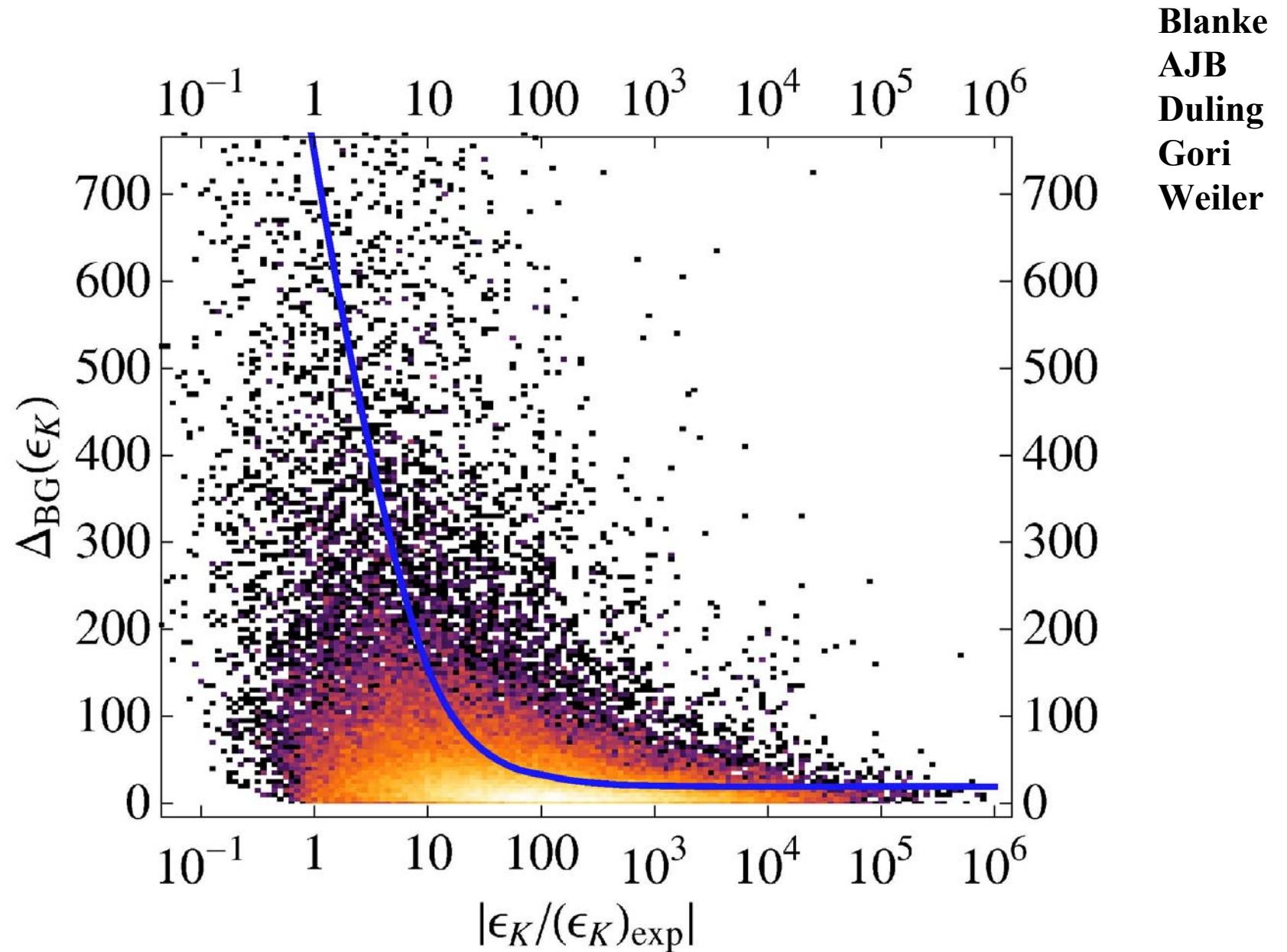
$M_{KK} \approx 2.5 \text{TeV}$

- ◆ Generically: $\varepsilon_K \cong 10^2 \varepsilon_K^{\text{exp}}$
- ◆ $\Delta_{\text{BG}}(\varepsilon_K)$ decreases with increasing ε_K
- ◆ Parameter sets with moderate $\Delta_{\text{BG}}(\varepsilon_K) \leq 20$ and $\varepsilon_K \approx \varepsilon_K^{\text{exp}}$ exist.

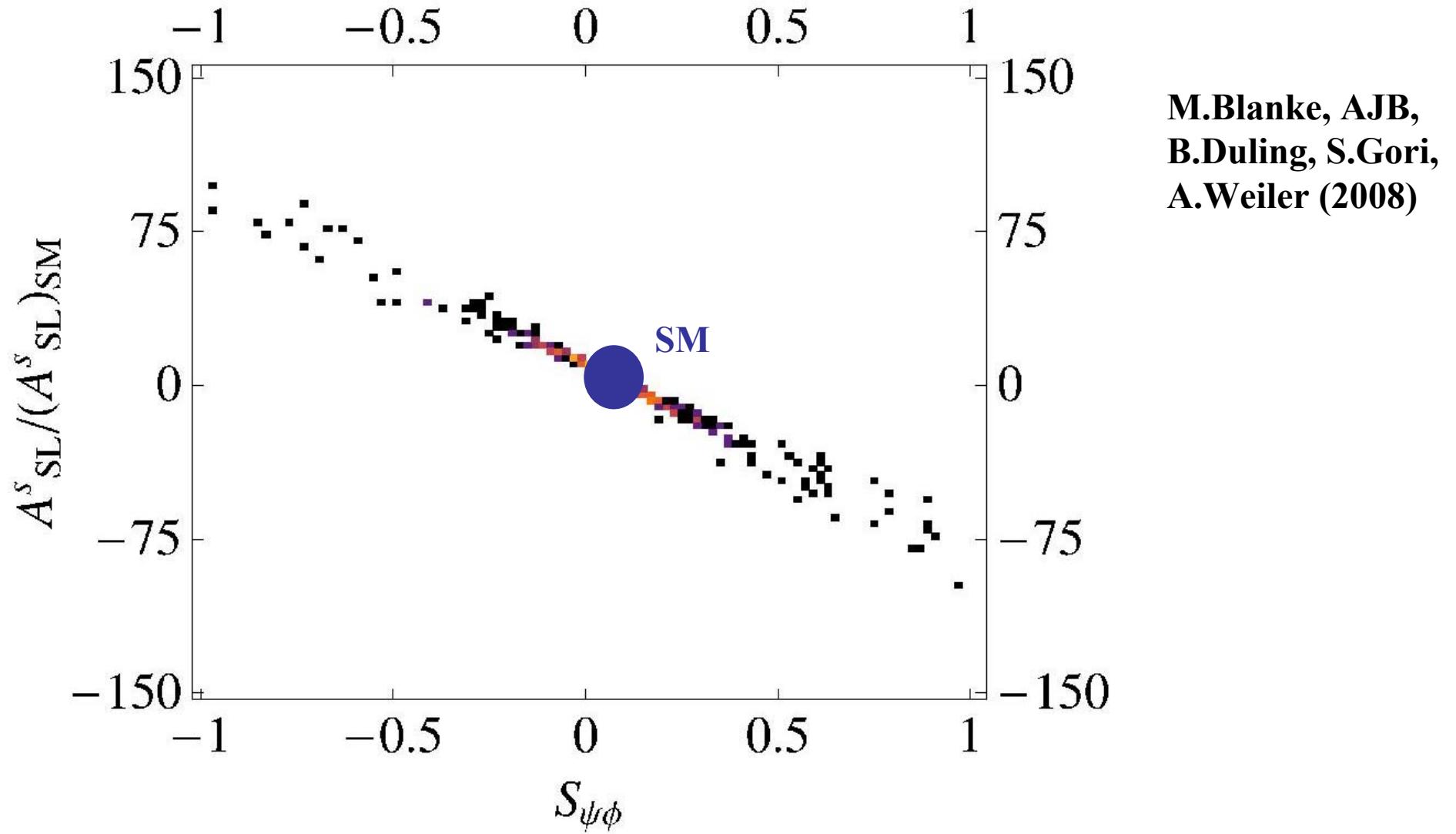
For ΔM_K and $\Delta B=2$ observables fine tuning much smaller.
Generically: $(\Delta_{\text{BG}} \lesssim 20)$ $(\Delta_{\text{BG}} \leq 5)$

RS-GIM
very
effective

Fine Tuning in ϵ_K

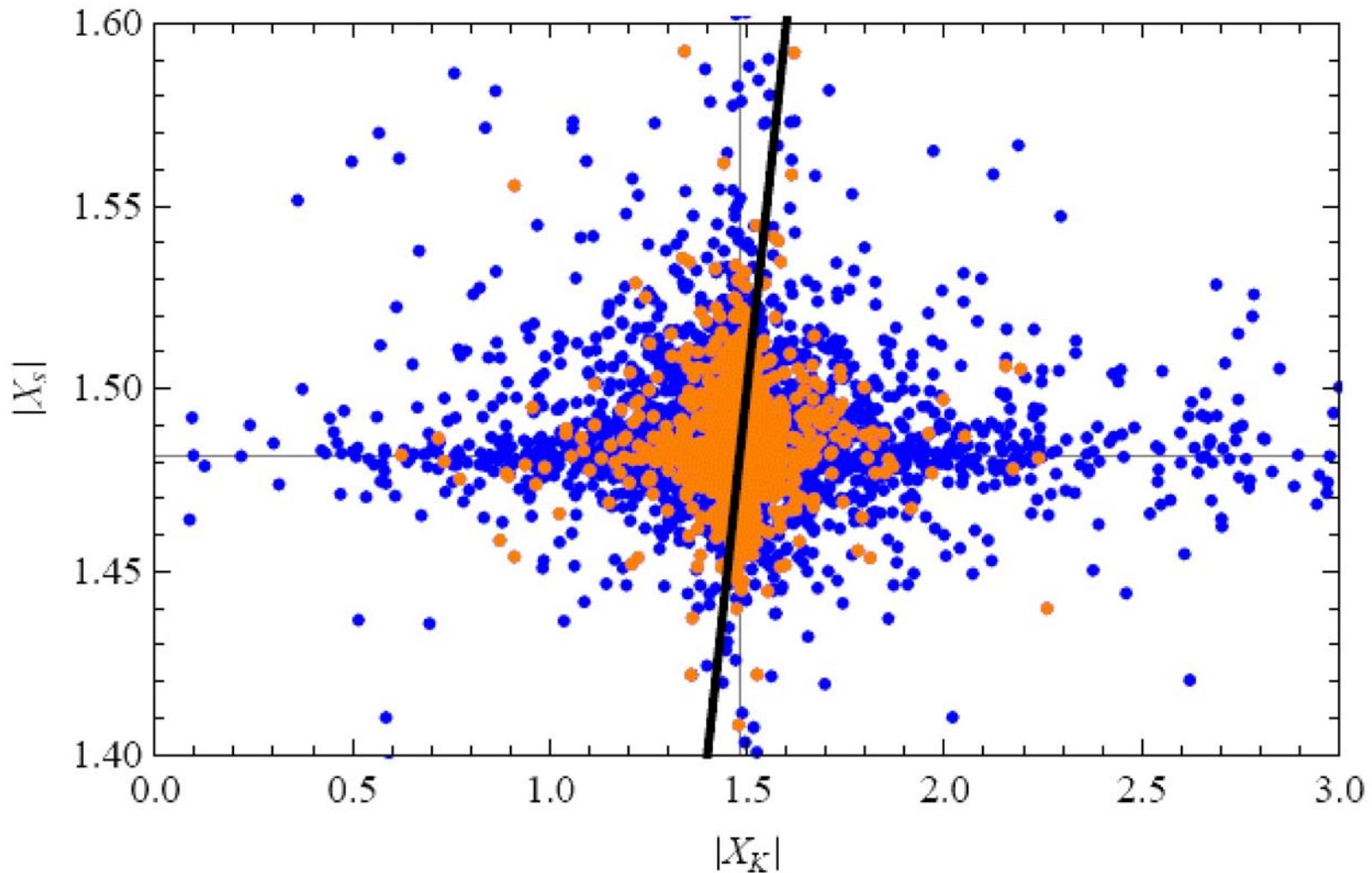


Correlation in Warped Extra Dimensions



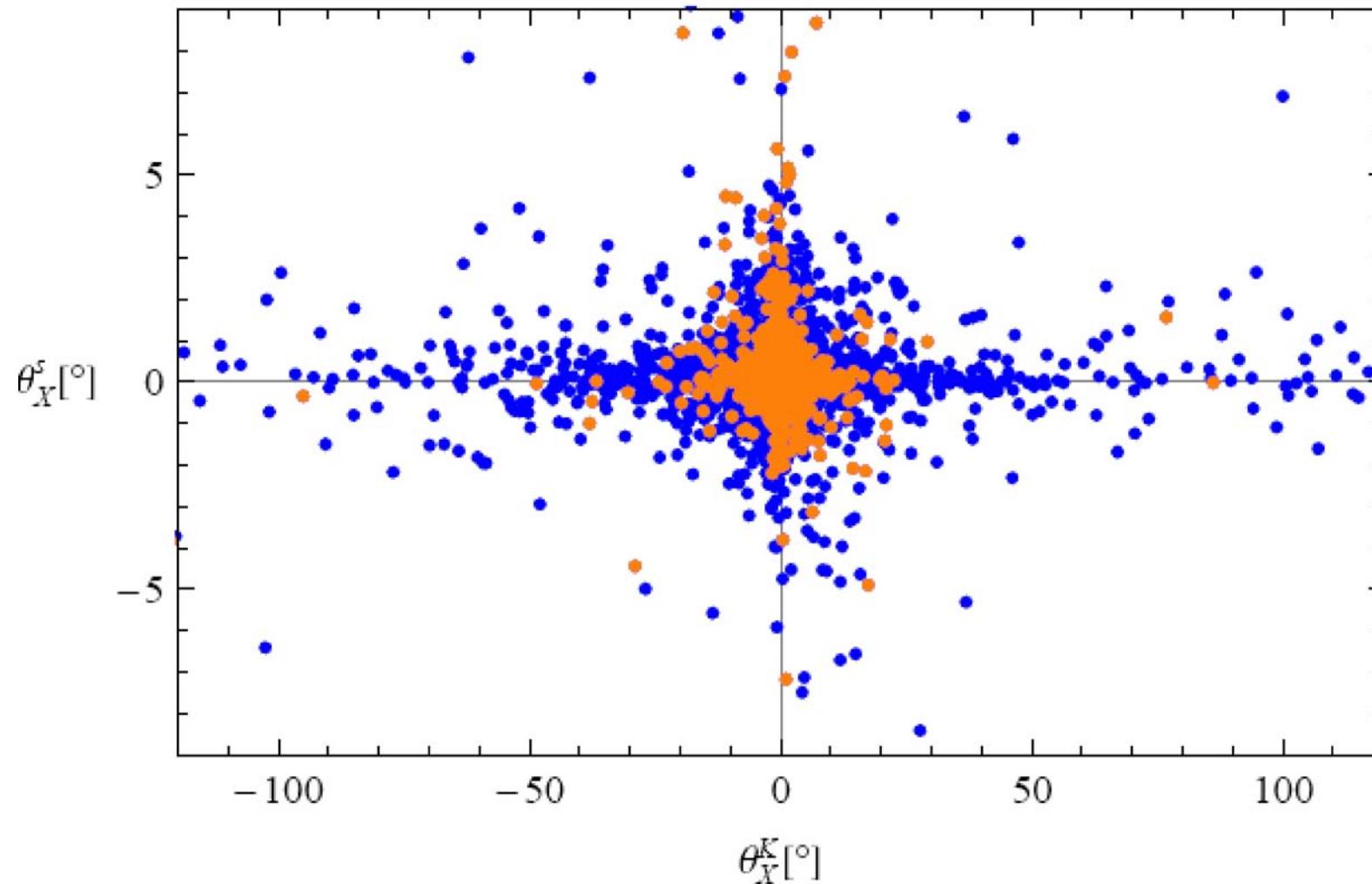
Breakdown of Universality in X

(NP effects much larger in K decays)



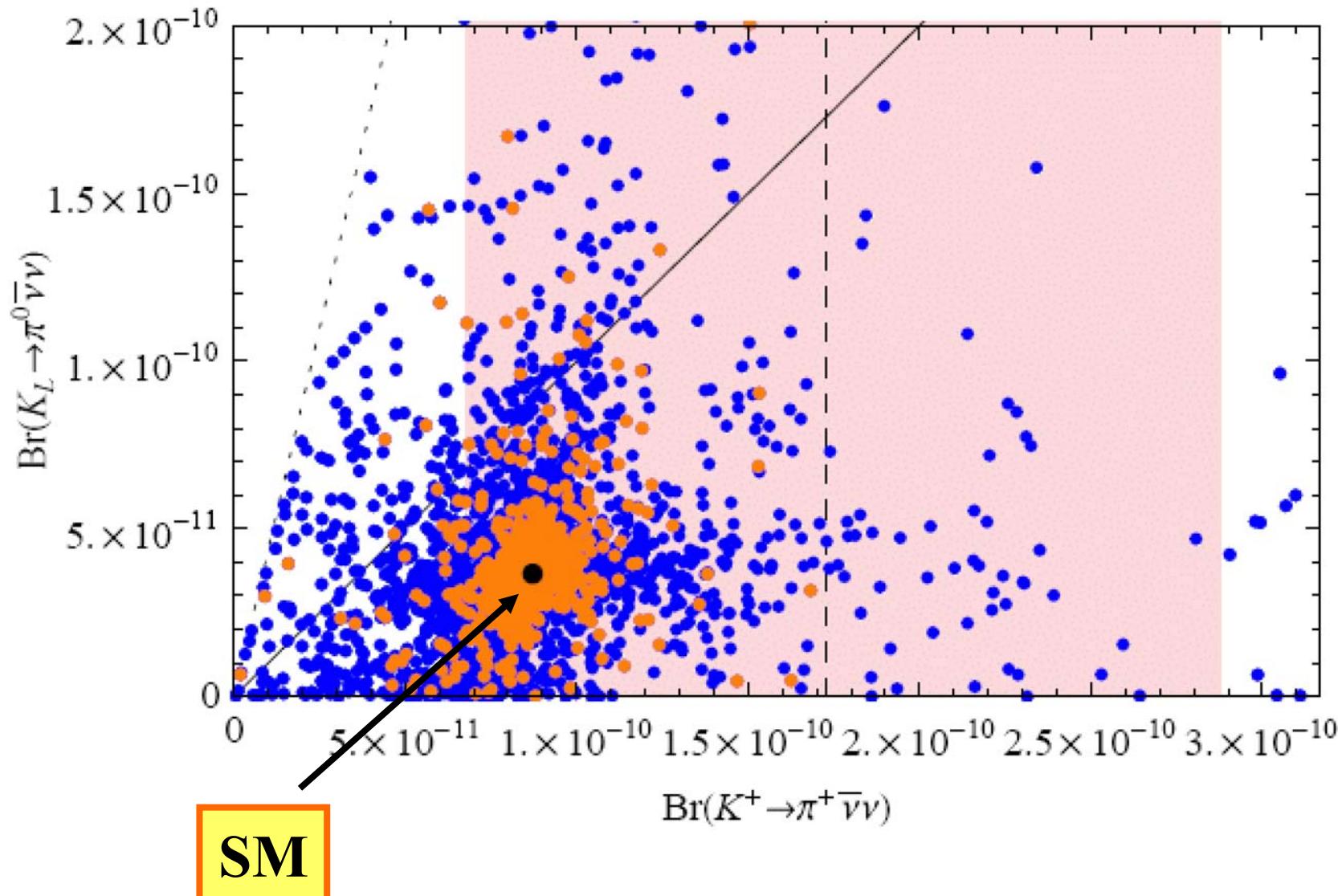
New CP-Violating Phases

(NP effects much larger in K decays)



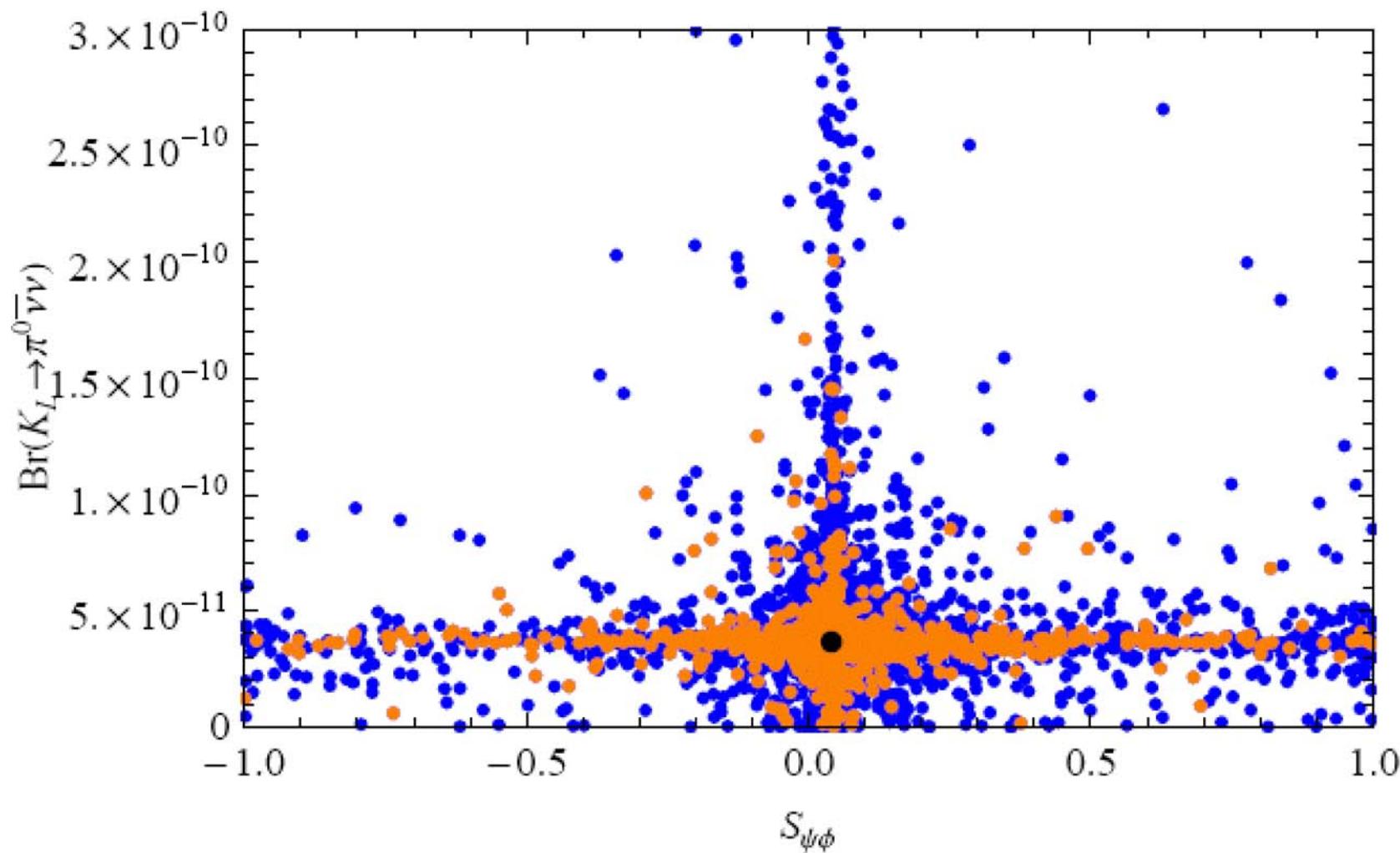
$K_L \rightarrow \pi^0 \bar{\nu} \nu$ vs. $K^+ \rightarrow \pi^+ \bar{\nu} \nu$

(Up to Factor 5 and 2 Enhancements)



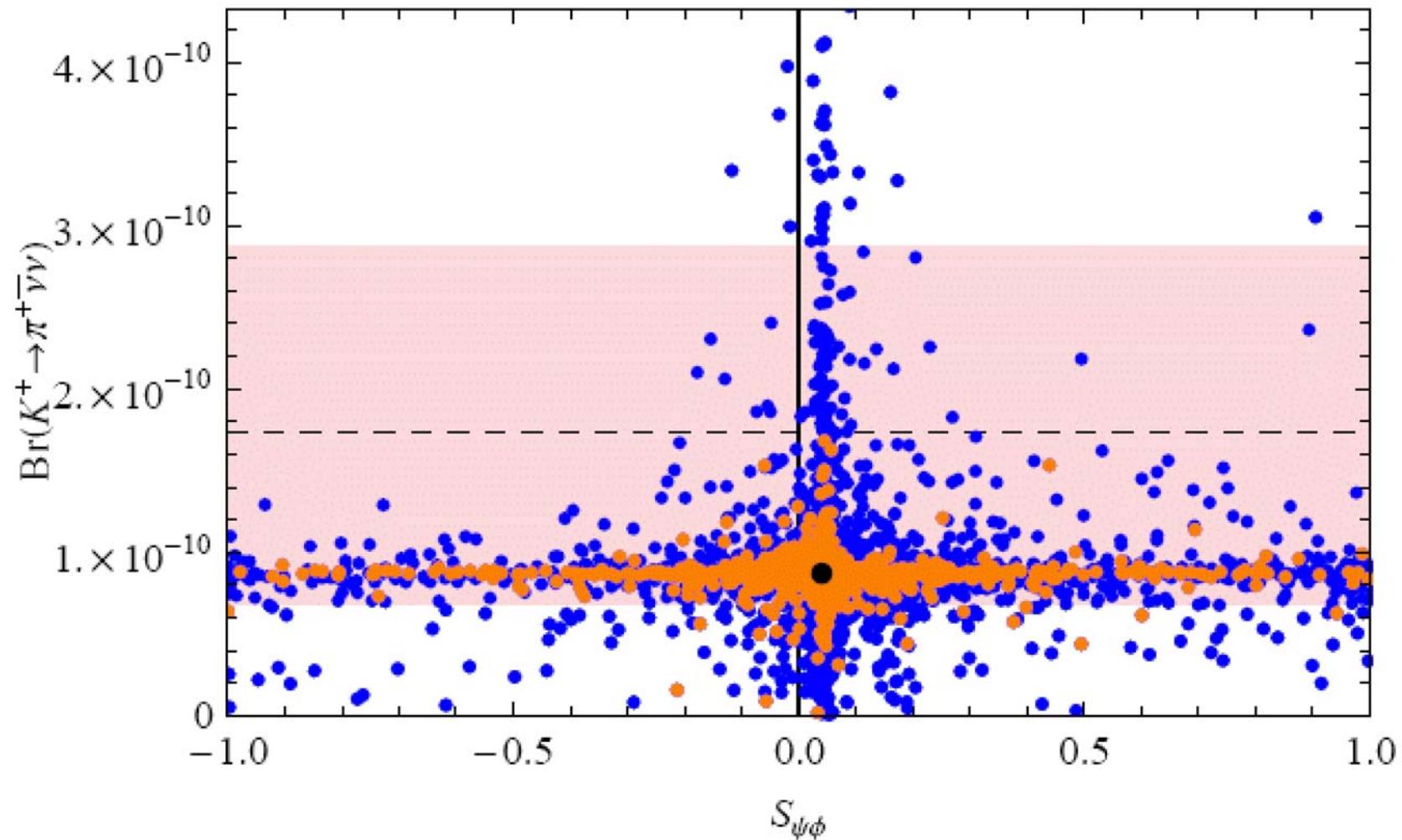
$$K_L \rightarrow \pi^0 \bar{\nu} \nu \text{ vs. } S_{\psi\phi}$$

(Simultaneous Large Enhancements unlikely)

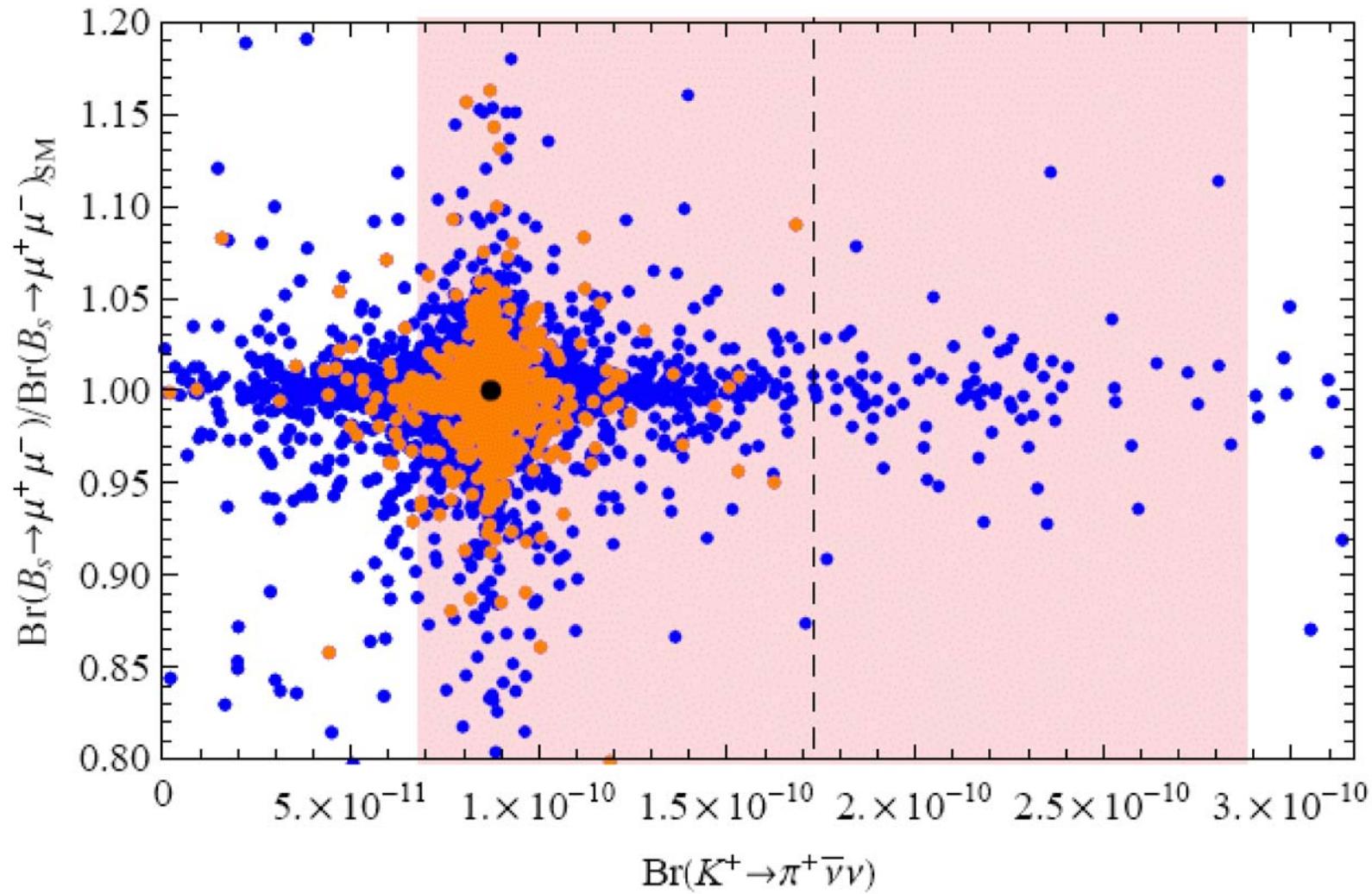


$$K^+ \rightarrow \pi^+ \nu\bar{\nu} \text{ vs. } S_{\psi\phi}$$

(Simultaneous Large Enhancements unlikely)

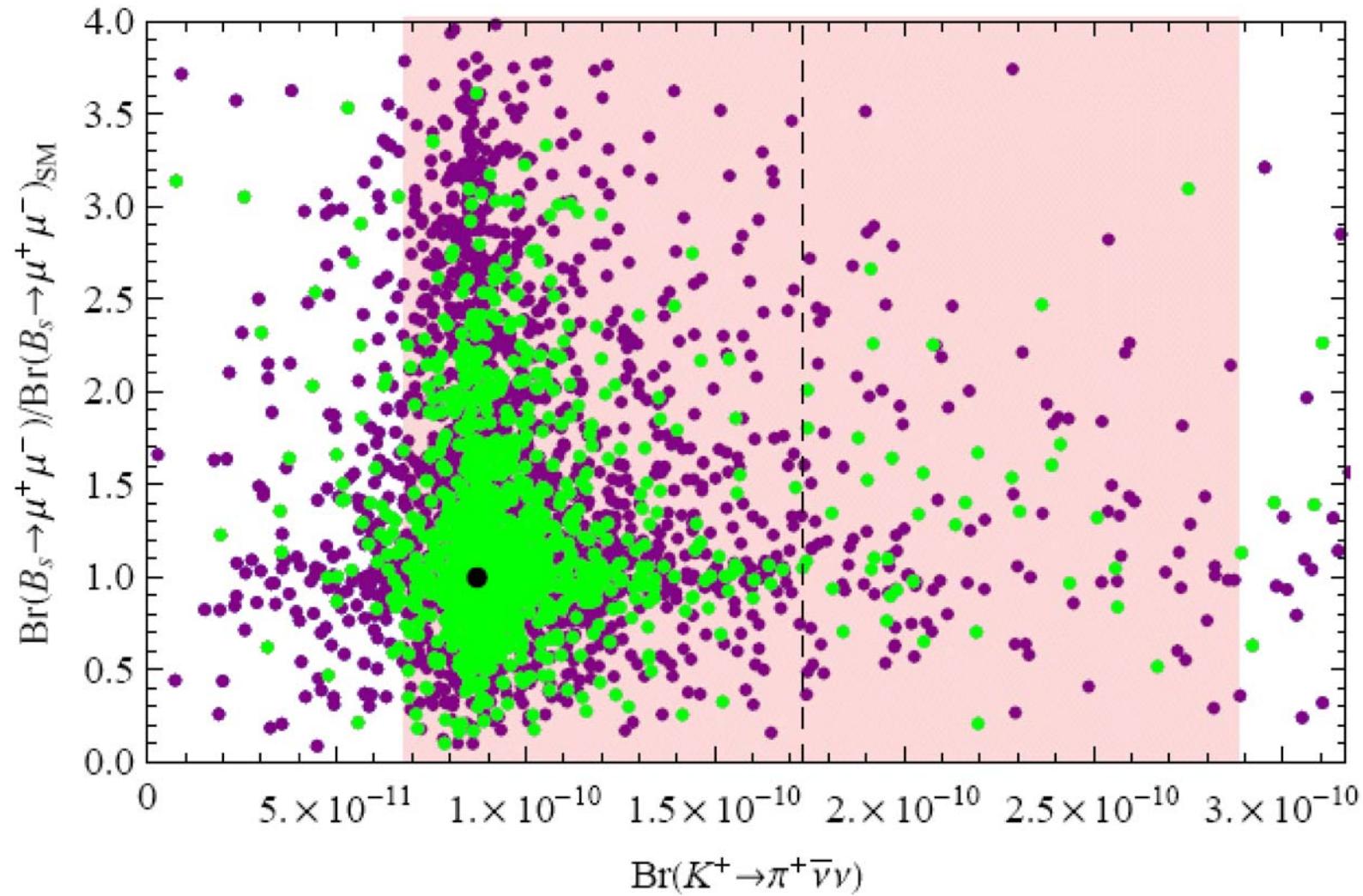


Very strong Protection for $B_s \rightarrow \mu^+ \mu^-$



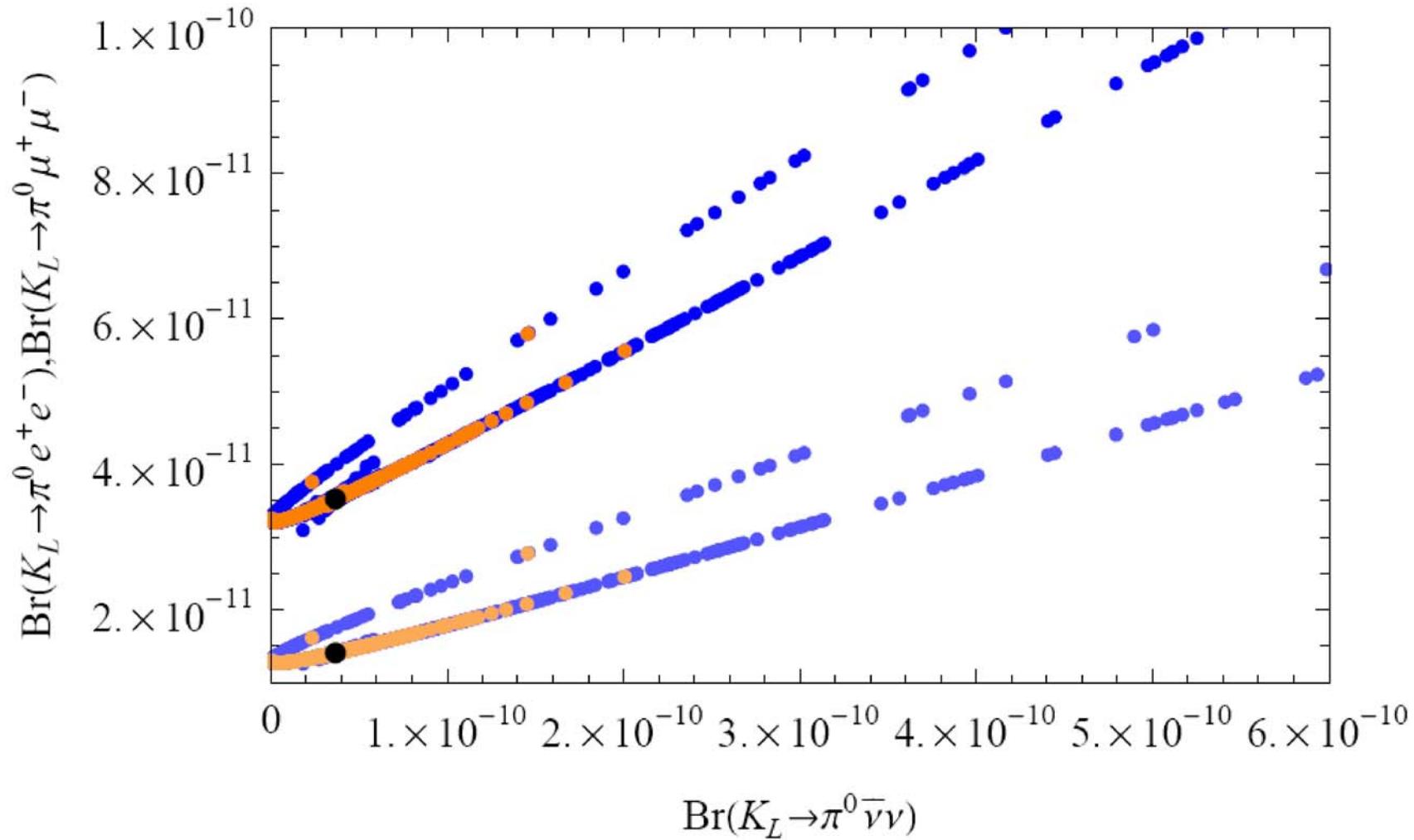
Removal of Protection

($\text{Br}(B_s \rightarrow \mu^+ \mu^-)$ enhanced up to 10^{-8})

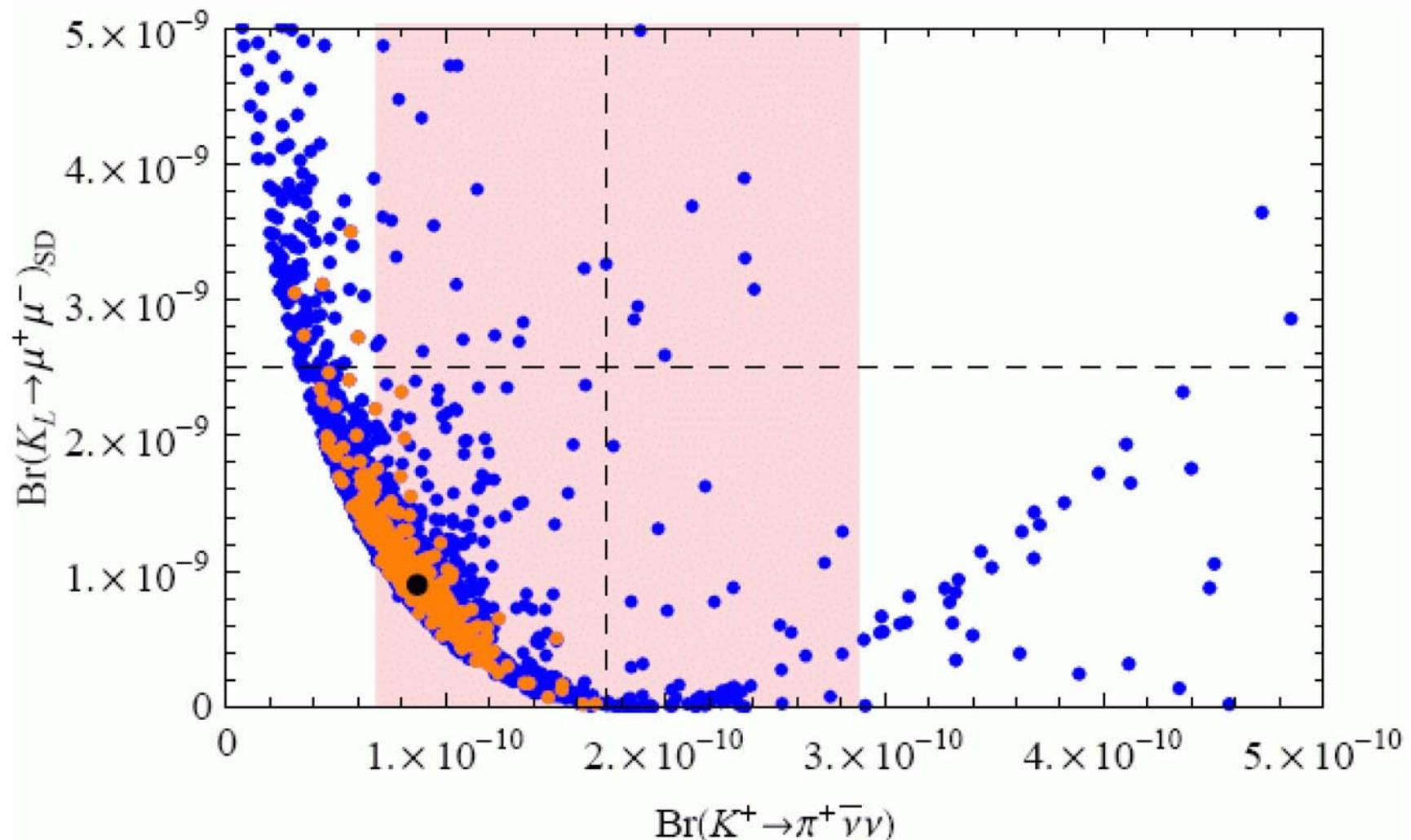


Strong Correlation between

$K_L \rightarrow \pi^0 l^+ l^-$ and $K_L \rightarrow \pi^0 \bar{\nu}\nu$



Correlation between $K_L \rightarrow \mu^+ \mu^-$ and $K^+ \rightarrow \pi^+ \bar{\nu}\nu$



Golden Relations of CMFV:

AJB (03)

$$\frac{Br(B_s \rightarrow \mu^+ \mu^-)}{Br(B_d \rightarrow \mu^+ \mu^-)} = \frac{\hat{B}_{B_d}}{\hat{B}_{B_s}} \frac{\tau(B_s)}{\tau(B_d)} \frac{\Delta M_s}{\Delta M_d} r$$

(CMFV)
r = 1

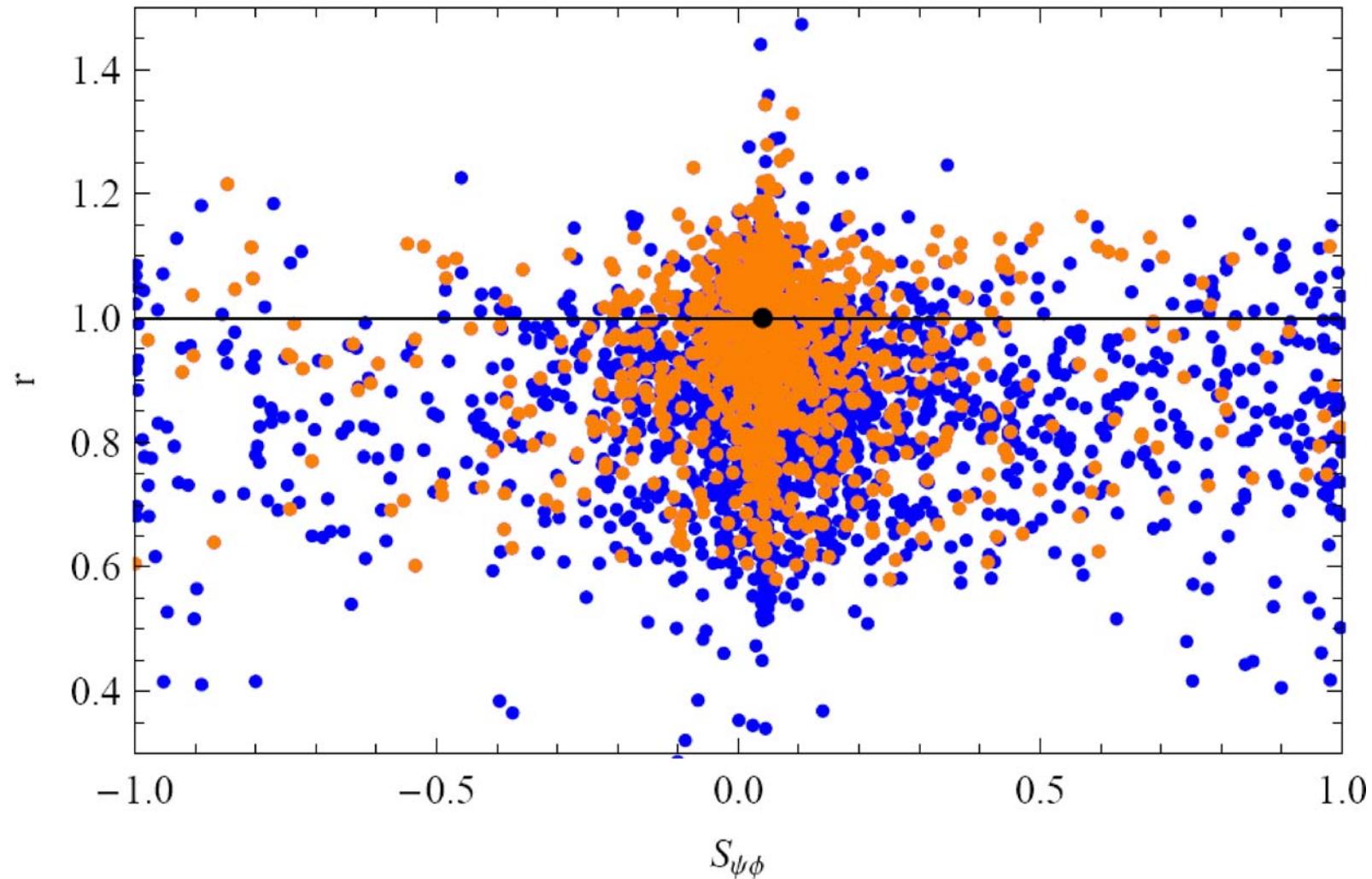
Buchalla
AJB (94)

$$(\sin 2\beta)_{B \rightarrow \psi K_S} = (\sin 2\beta)_{K \rightarrow \pi \nu \bar{\nu}}$$

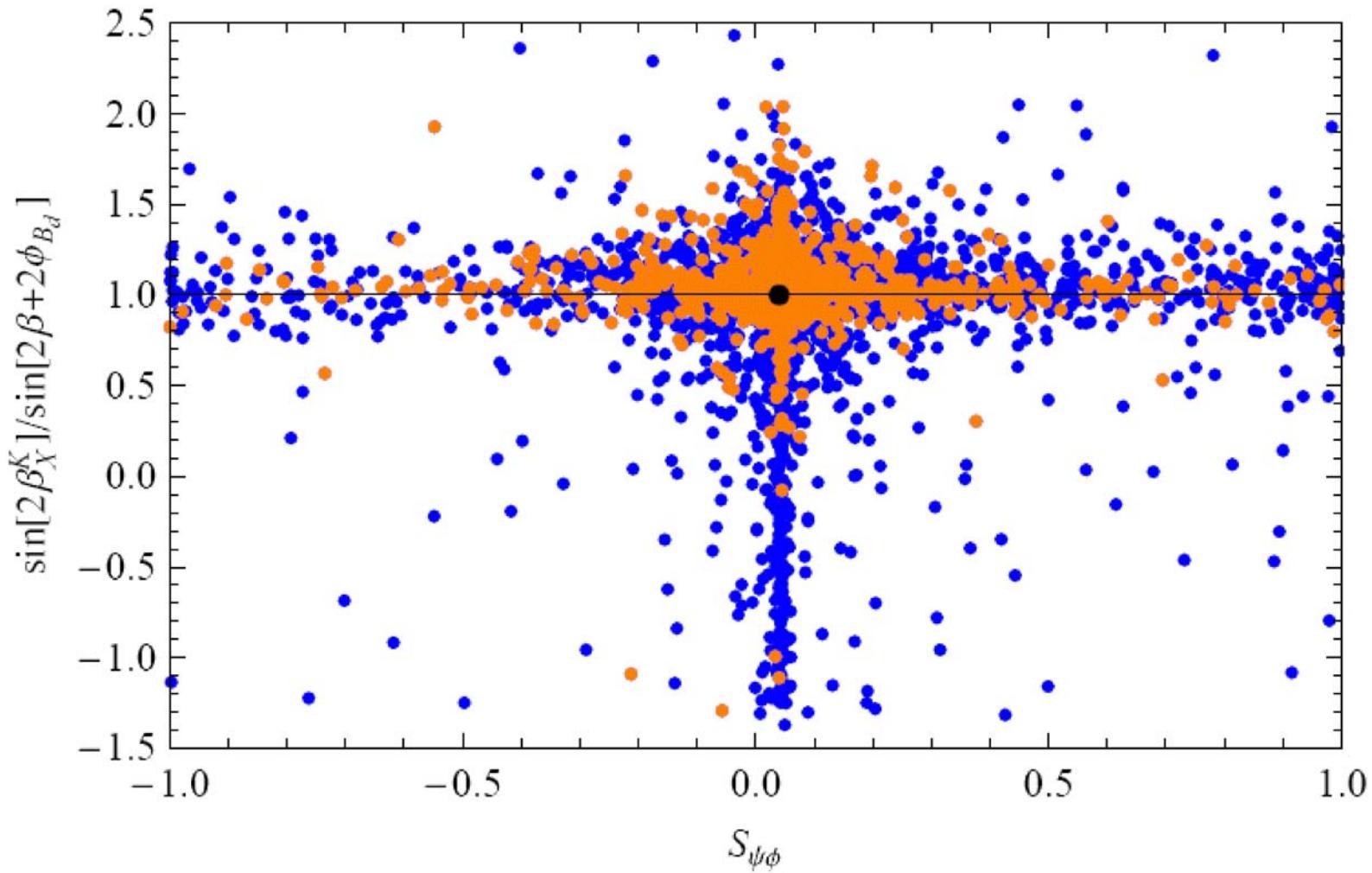
(MFV)

The violation of these model independent MFV (CMFV) relations would signal new flavour and CP-violating interactions (and/or new operators)

Violation of the Golden MFV Relation I



Violation of the Golden MFV Relation II



Clear Pattern of Flavour Violation

- 1.** ϵ_K can be made consistent with data for $M_{KK} \cong 2 - 3\text{TeV}$ with only moderate tuning of Y_{5D}
- 2.** $S_{\psi\phi}$ can be much larger than $(S_{\psi\phi})_{\text{SM}}$
- 3.** $K_L \rightarrow \pi^0 v\bar{v}$ and $K^+ \rightarrow \pi^+ v\bar{v}$ can be enhanced up to factors 5 and 2
- 4.** Rare B-decays SM-like
- 5.** Simultaneous enhancements of $S_{\psi\phi}$ and $K \rightarrow \pi v\bar{v}$ very unlikely
- 6.** Analysis of $B \rightarrow X_s \gamma$, $\mu \rightarrow e \gamma$, ... in progress

Backup

Number of new Flavour Parameters

(Quark Sector)

(physical)

Real

\mathcal{CP} Phases

SUSY

36

27

(R-parity)

LHT

7

3

some
sensitivity
to UV

RS

18

9

SM

9

1