

Conformal Neutrinos

with M. Quirós, 0901.0006

Gero von Gersdorff

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Outline

Introduction

Conformal (Dirac) Neutrinos

The model

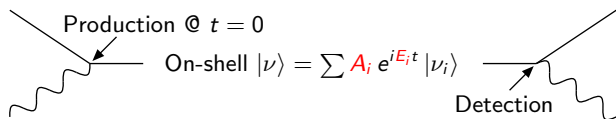
Flavour Changing Processes

Higgs Decay

Conformal Majorana Neutrinos

Neutrino Masses - Experimental Results

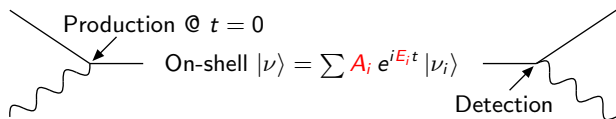
- ▶ Neutrino masses measured via oscillations



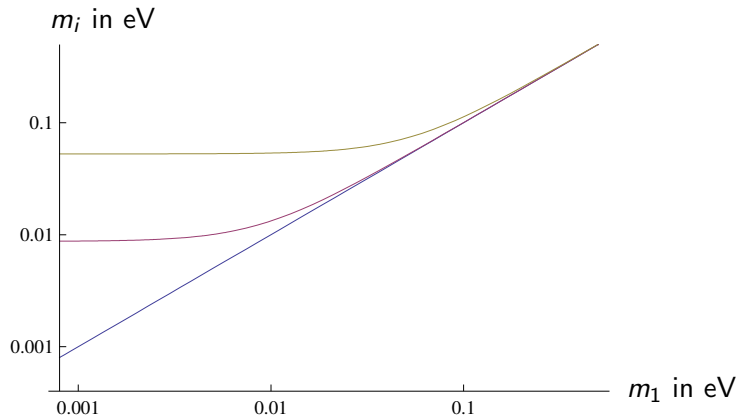
- ▶ **Amplitude** gives **Mixing angles** ... **Frequency** gives Δm^2
 - ▶ $m_2^2 - m_1^2 = 7.6 \cdot 10^{-5} \text{ eV}^2$ [KAMLAND]
 - ▶ $|m_3^2 - m_2^2| = 2.7 \cdot 10^{-3} \text{ eV}^2$ [MINOS]

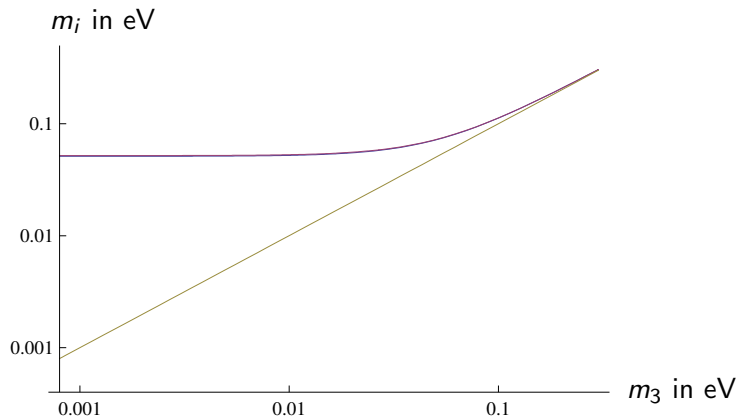
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- ▶ Still leaves one free continuous parameter and one sign
 - ▶ $m_3 > m_2 > m_1 \Rightarrow$ regular "hierarchy"
 - ▶ $m_2 > m_1 > m_3 \Rightarrow$ inverted "hierarchy"

Regular "Hierarchy": $m_3 > m_2 > m_1$ 

Inverted "Hierarchy": $m_2 > m_1 > m_3$ 

Neutrino Masses - Theory

- ▶ Fermion masses are **multiplicatively** renormalized
- ▶ Still, scale $m_\nu \lesssim 1$ eV clearly requires some explanation
- ▶ Creating m_ν with Yukawa couplings

$$\mathcal{L}^Y = h_\nu H \ell_L \bar{n}_R \quad \implies \quad h_\nu = \frac{m_\nu}{v} = 10^{-12}$$

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- ▶ OR add **large mass term** for RH neutrino

$$m_R^M (n_R)^2 \xrightarrow{\text{integr. out } n_R} \frac{h_\nu^2}{m_R^M} (H \ell_L)^2 \xrightarrow{EWSB} \frac{v^2 h_\nu^2}{m_R^M} (n_L)^2$$

- ▶ **See Saw** mechanism: $h_\nu = \mathcal{O}(1)$ for $m_R^M \approx 10^{14}$ GeV.

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- ▶ Assume RH neutrinos n_R belong to (strongly coupled) **conformal sector**
- ▶ Allow for **large** anomalous dimension $1/2 \lesssim \gamma < 1$
- ▶ Impose **lepton number conservation** (can be relaxed)
- ▶ Lagrangian

$$\mathcal{L} = \mathcal{L}_{CFT}(n_R, \dots) + \mathcal{L}_{SM} + \Lambda^{-\gamma} H \bar{n}_R \ell_L$$

Breaking of Scale Invariance

- ▶ EWSB induces a Dirac mass term

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Conclusion

The scale of conformal breaking defines the neutrino mass scale.
Parametrically suppressed wrt electroweak scale.

The RG Flow

- ▶ Use rescaled fields to display running Yukawa/mass

$$n_R = \mu^\gamma \nu_R \quad h(\mu) = \left(\frac{\mu}{\Lambda}\right)^\gamma, \quad m^D(\mu) = \left(\frac{\mu}{\Lambda}\right)^\gamma v$$

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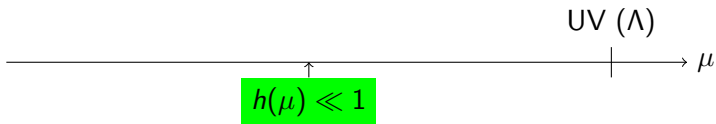
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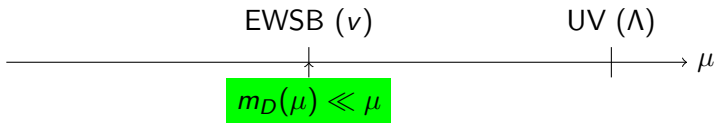
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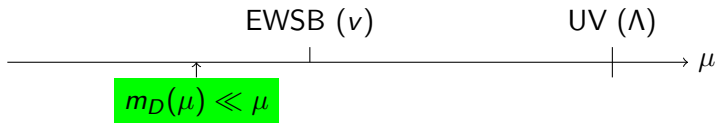
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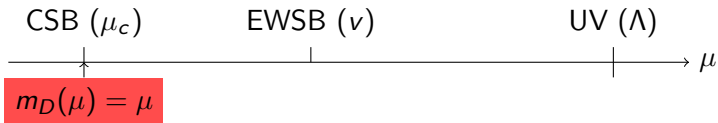
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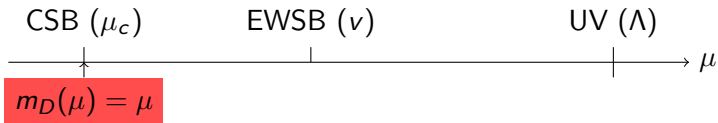
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- ▶ Yields previous result

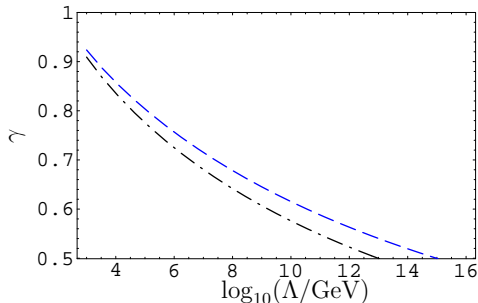
$$\mu_c = m^D = v \left(\frac{v}{\Lambda}\right)^{\frac{\gamma}{1-\gamma}}$$

UV-scale vs anomalous dimension

- ▶ For fixed neutrino mass scale, Λ and γ are related

$$m^D = v \left(\frac{v}{\Lambda} \right)^{\frac{\gamma}{1-\gamma}}$$

$$\left\{ \begin{array}{l} 1 \text{ for } \gamma = 0.5 \\ 9 \text{ for } \gamma = 0.9 \end{array} \right.$$



Here: $m^D = 0.05 \dots 1 \text{ eV}$

Three Generations

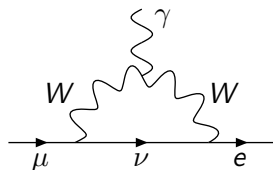
- ▶ Degenerate masses natural if anomalous dimensions equal
- ▶ Hierarchies can be achieved with **different anomalous dimensions**
- ▶ assume $\gamma_3 = \gamma_2 < \gamma_1$. CSB occurs at $\mu_c = m_3 = m_2$ but

$$m_1(\mu_c) = \mu_c \left(\frac{\mu_c}{\Lambda} \right)^{\gamma_1 - \gamma}$$

- ▶ Can be **suppressed** w.r.t. $m_{2,3}$
- ▶ Just two RH neutrinos with same γ give $m_3 = m_2$, $m_1 = 0$.
- ▶ Also differences in Yukawas at high scales, SM corrections, etc....

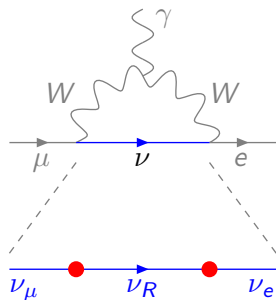
FCNC mediated by SM neutrinos

- ▶ Consider $\mu \rightarrow e \gamma$ in SM with Dirac Neutrinos
- ▶ Vanishes at tree level (GIM)
- ▶ Contribution at 1 loop



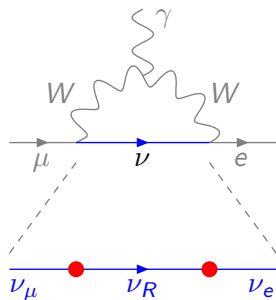
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- ▶ Need **two mass insertions**
- ▶ Branching scales as $B \sim \left(\frac{m^D}{M_W}\right)^4$.
- ▶ **Conformal Neutrinos**: Dirac masses grow with energy!



FCNC mediated by conformal neutrinos

- ▶ Result for conformal neutrinos:

$$B(\mu \rightarrow e\gamma) = \frac{3}{32} \frac{\alpha}{\pi} \left| \frac{\pi\gamma}{\sin(\pi\gamma)} \sum_i U_{ei} U_{\mu i}^* \left(\frac{m_i}{M_W} \right)^{2-2\gamma} \right|^2$$

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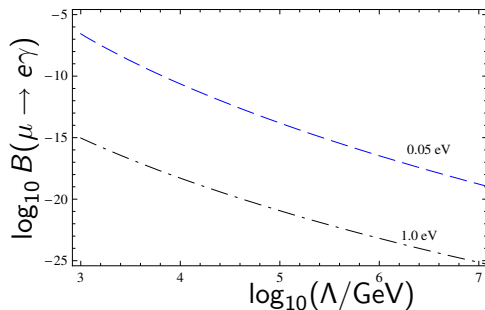
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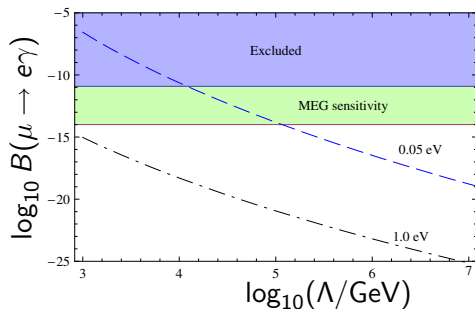
- ▶ $\gamma = 0$ gives SM result...
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- ▶ Prefactor not very important, except near $\gamma = 1$.
- ▶ **Parametric enhancement** for $\gamma > 0$ w.r.t SM.
- ▶ Can be attributed to **running of Dirac masses**.

Comparing to Experiment



- ▶ Contours for fixed ν masses
- ▶ (using best fits for U)
- ▶ Regular hierarchy

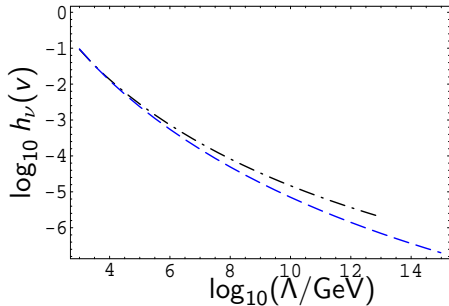
Comparing to Experiment



- ▶ Contours for fixed ν masses
- ▶ (using best fits for U)
- ▶ Regular hierarchy
- ▶ **Bound** (for 0.05 eV ν)
 $\Rightarrow \Lambda \gtrsim 10 \text{ TeV}$ ($\gamma \lesssim 0.86$)
- ▶ Future sensitivity
 $\Rightarrow \Lambda \sim 100 \text{ TeV}$ ($\gamma \sim 0.81$)

Coupling to Higgs

- ▶ Yukawa becomes stronger at larger scales
- ▶ Can we detect the running Yukawa?
- ▶ For $\Lambda = 10$ TeV Neutrino Yukawas are **comparable to c, τ** !



Yukawa coupling at the **weak scale**

The width of the Higgs

- ▶ Partial width of Higgs to $\nu\bar{\nu}$

$$\Gamma(H \rightarrow \nu\bar{\nu}) = h^2(m_H) \frac{m_H}{16\pi} \frac{2}{\Gamma(1-\gamma)\Gamma(3+\gamma)}$$

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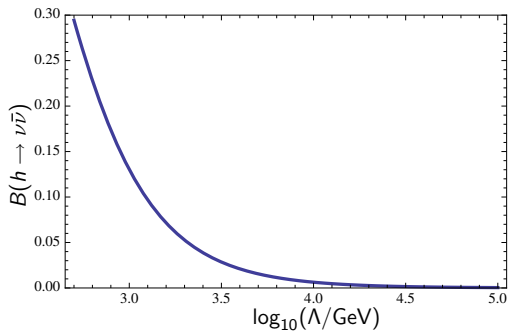
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- ▶ **Conformal neutrinos**: main effect is the enhanced Yukawa

$$h(m_H) \approx 10^{-6} \dots 10^{-2}$$

- ▶ Also some mild suppression for $\gamma = 1$

Higgs Branching in Neutrinos



- ▶ $m^D = 0.1$ eV.
- ▶ $m_H = 130$ GeV.

Lepton Number Violation

- ▶ Neutrino Lagrangian reads

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- ▶ EWSB again creates LH Majorana mass, does not run, dominates at low energies...
- ▶ Like see saw, $m_L^M \sim c v^2/\Lambda$, but no heavy states!

An Inverse See-Saw

- ▶ Parametric dependence of mass matrix (1 flavour)

$$\begin{pmatrix} m_L^M & m^D \\ m^D & m_R^M \end{pmatrix} = v \begin{pmatrix} \frac{v}{\Lambda} & \left(\frac{v}{\Lambda}\right)^{\frac{\gamma}{1-\gamma}} \\ \left(\frac{v}{\Lambda}\right)^{\frac{\gamma}{1-\gamma}} & \left(\frac{v}{\Lambda}\right)^{\frac{3\gamma-1}{1-\gamma}} \end{pmatrix}$$

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- ▶ Disappears for $\gamma = 1/2$

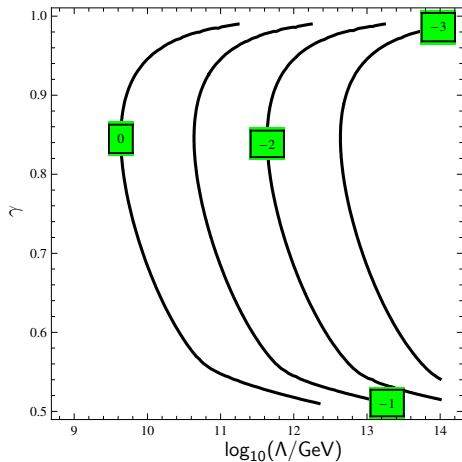
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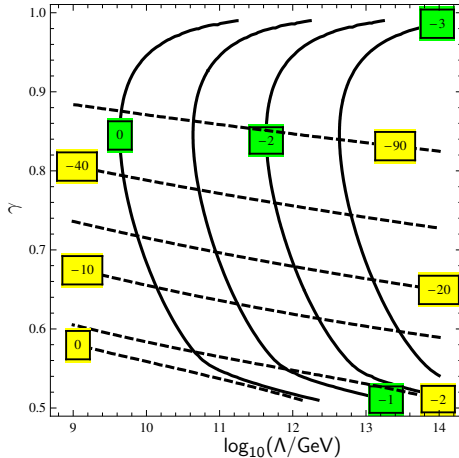
- ▶ Hierarchy $m_L^M \gg m^D \gg m_R^M$
- ▶ Disappears for $\gamma = 1/2$
- ▶ Leads to **inverse** See-saw mechanism (very light sterile states...)

The Masses



log (heavy mass / eV)

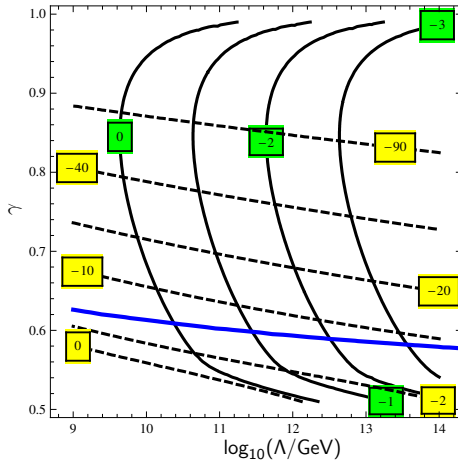
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The Masses



$\log(\text{heavy mass} / \text{eV})$

$\log(\text{light mass} / \text{eV})$

— $\text{Mixing} = 1\%$

- ▶ Conformal Neutrinos are an interesting method to obtain small neutrino masses
- ▶ RH neutrinos are conformal operators with large anomalous dimension
- ▶ Conformal Symmetry breaking naturally triggered by EWSB
- ▶ $\mu \rightarrow e\gamma$ and $H \rightarrow \nu\bar{\nu}$ strongly enhanced w.r.t. SM
- ▶ Conformal Majorana neutrinos also possible, similar to Seesaw, but no heavy states
- ▶ Lots of things to be done: cosmology, supersymmetry...