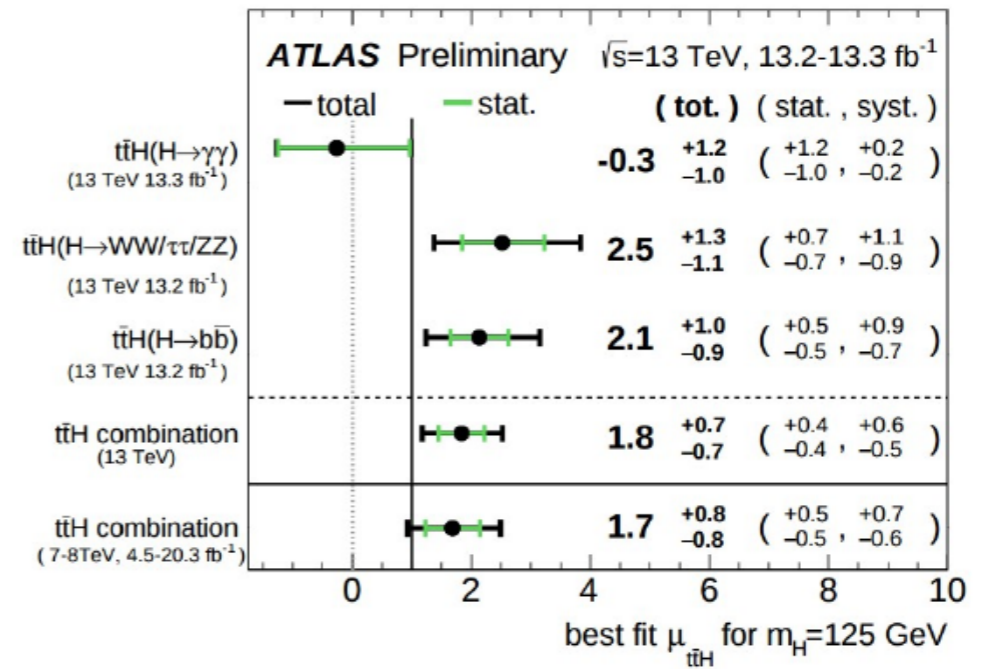
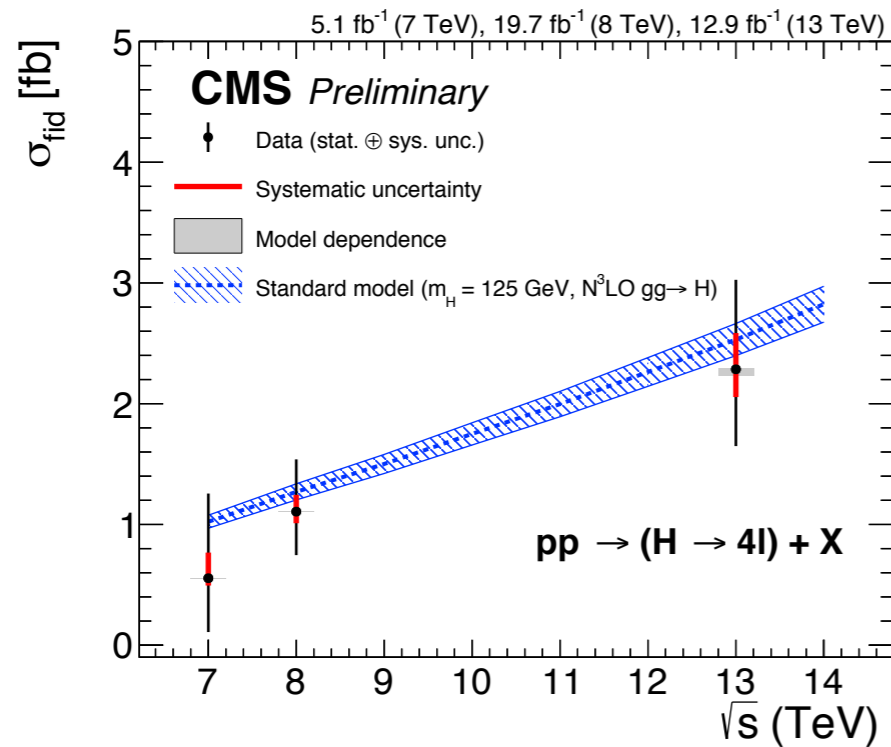
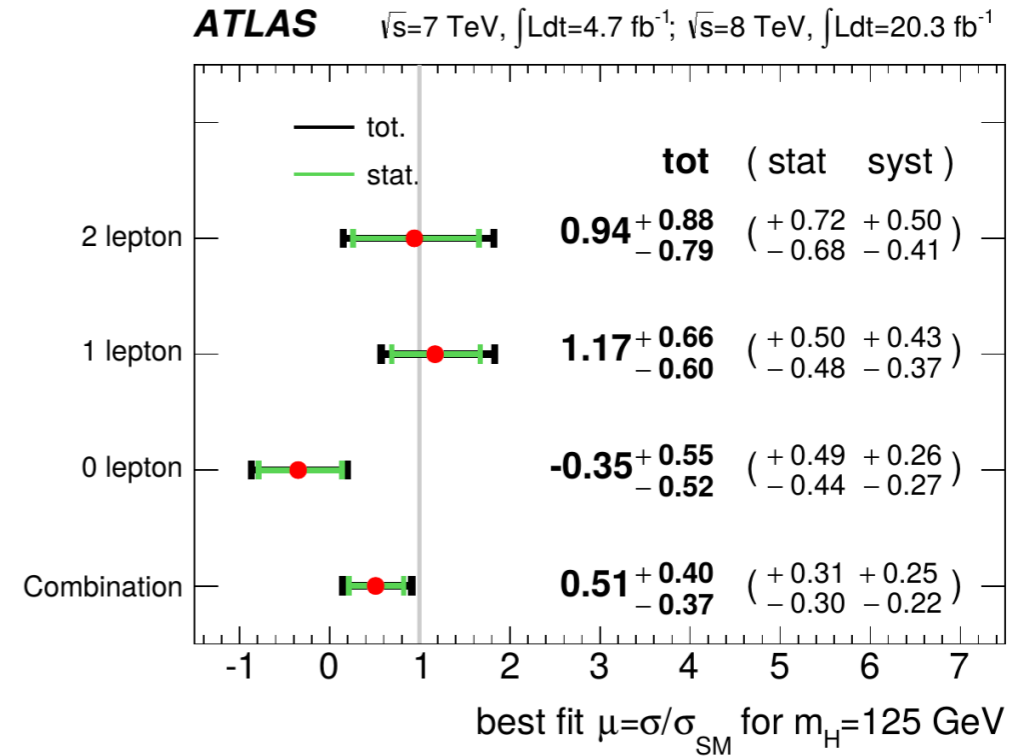
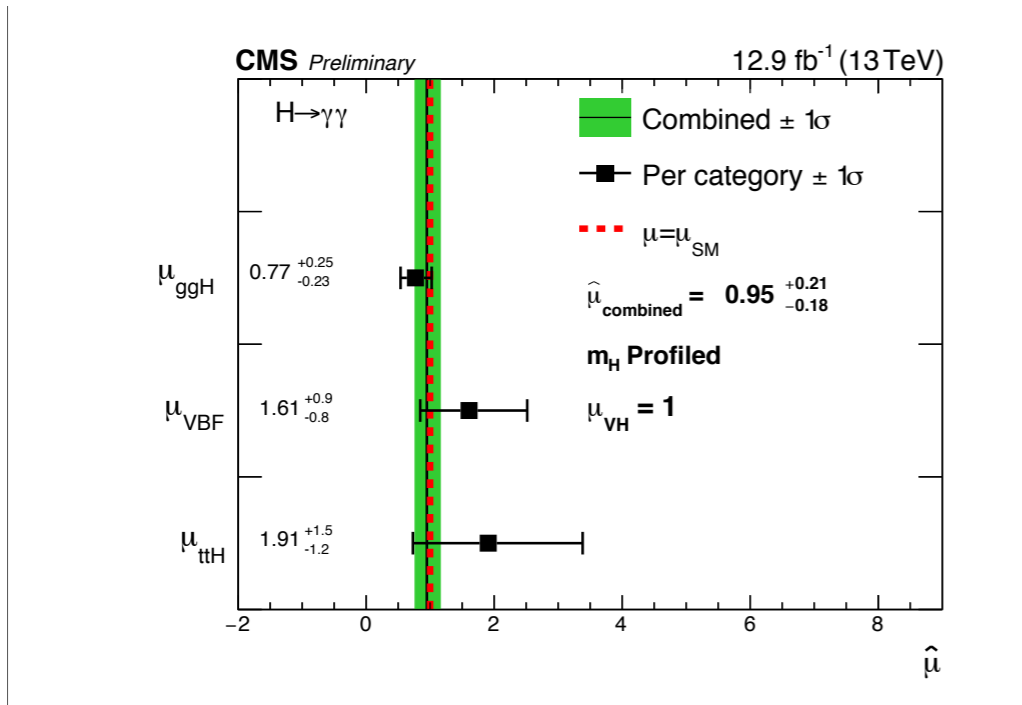


Physics opportunities at Future Hadron collider

LianTao Wang
University of Chicago

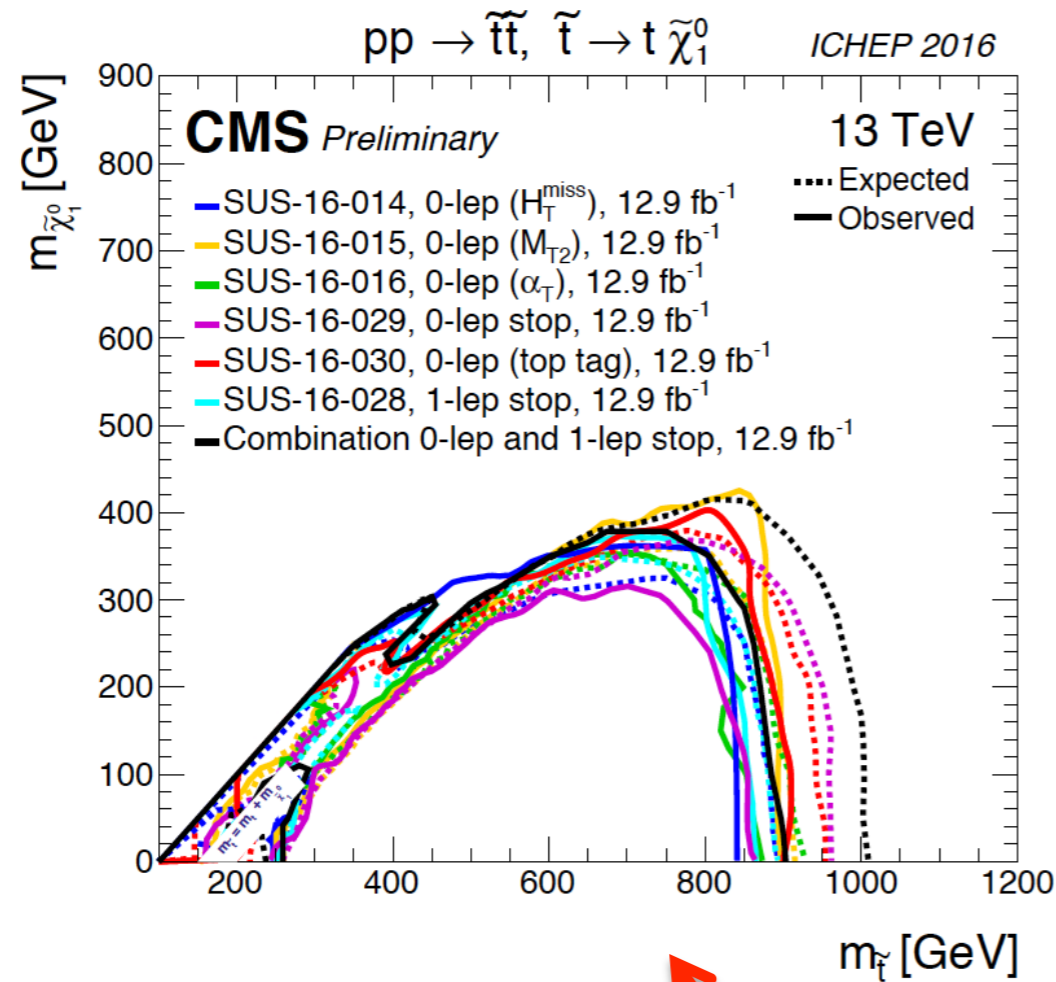
44th SSI. August 19, 2016

Current status



Found Higgs

Current status



ATLAS Exotics Searches* - 95% CL Exclusion
Status: August 2016

ATLAS Preliminary
 $\int \mathcal{L} dt = (3.2 - 20.3) \text{ fb}^{-1}$ $\sqrt{s} = 8, 13 \text{ TeV}$

Model	ℓ, γ	Jets†	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference
Extra dimensions	ADD $G_{KK} + g/q$	$\geq 1j$	Yes	3.2	M_0 6.58 TeV	$n=2$ 1604.07773
	ADD non-resonant $\ell\ell$	$2e, \mu$	-	20.3	M_0 4.7 TeV	$n=3$ HLZ 1407.2410
	ADD QBH $\rightarrow \ell q$	$1e, \mu$ $1j$	-	20.3	M_0 5.2 TeV	$n=6$ 1311.2006
	ADD QBH	-	$\geq 2j$	15.7	M_0 8.7 TeV	$n=6$ ATLAS-CONF-2016-069
	ADD BH high Σp_T	$\geq 1e, \mu$	$\geq 2j$	3.2	M_0 8.2 TeV	$n=6, M_D = 3 \text{ TeV, rot BH}$ 1606.02265
	ADD BH multijet	-	$\geq 3j$	3.6	M_0 9.55 TeV	$n=6, M_D = 3 \text{ TeV, rot BH}$ 1512.02586
	RS1 $G_{KK} \rightarrow \ell\ell$	$2e, \mu$	-	20.3	G_{KK} mass 2.68 TeV	$k/M_{Pl} = 0.1$ 1405.4123
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2γ	-	3.2	G_{KK} mass 3.2 TeV	$k/M_{Pl} = 0.1$ 1606.03833
	Bulk RS $G_{KK} \rightarrow WW \rightarrow qq\ell\nu$	$1e, \mu$ $1j$	Yes	13.2	G_{KK} mass 1.24 TeV	$k/M_{Pl} = 1.0$ ATLAS-CONF-2016-062
	Bulk RS $G_{KK} \rightarrow HH \rightarrow bbbb$	-	$\geq 4b$	13.3	G_{KK} mass 360-860 GeV	$k/M_{Pl} = 1.0$ ATLAS-CONF-2016-049
Bulk RS $G_{KK} \rightarrow tt$	$1e, \mu$ $\geq 1b, \geq 1J/2j$	Yes	20.3	G_{KK} mass 2.2 TeV	$BR = 0.925$ 1505.07018	
2UED / RPP	$1e, \mu$ $\geq 2b, \geq 4j$	Yes	3.2	KK mass 1.46 TeV	Tier (1,1), $BR(A^{(1,1)} \rightarrow \tau\tau) = 1$ ATLAS-CONF-2016-013	
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2e, \mu$	-	13.3	Z' mass 4.05 TeV	ATLAS-CONF-2016-045
	SSM $Z' \rightarrow \tau\tau$	2τ	-	19.5	Z' mass 2.02 TeV	1502.07177
	Leptophobic $Z' \rightarrow bb$	-	$2b$	3.2	Z' mass 1.5 TeV	1603.08791
	SSM $W' \rightarrow \ell\nu$	$1e, \mu$	Yes	13.3	W' mass 4.74 TeV	ATLAS-CONF-2016-061
	HVT $W' \rightarrow WZ \rightarrow qq\nu\nu$ model A	$0e, \mu$ $1j$	Yes	13.2	W' mass 2.4 TeV	ATLAS-CONF-2016-082
	HVT $W' \rightarrow WZ \rightarrow qq\nu\nu$ model B	-	$2j$	15.5	W' mass 3.0 TeV	ATLAS-CONF-2016-055
CI	HVT $V' \rightarrow WH/ZH$ model B	multi-channel	-	3.2	V' mass 2.31 TeV	1607.05621
	LRSM $W'_2 \rightarrow tb$	$1e, \mu$ $2b, 0-1j$	Yes	20.3	W' mass 1.92 TeV	1410.4103
	LRSM $W'_2 \rightarrow tb$	$0e, \mu$ $\geq 1b, 1j$	-	20.3	W' mass 1.76 TeV	1408.0886
DM	CI $qqqq$	-	$2j$	15.7	A 19.9 TeV $\eta_{LL} = -1$	ATLAS-CONF-2016-069
	CI $\ell\ell qq$	$2e, \mu$	-	3.2	A 25.2 TeV $\eta_{LL} = -1$	1607.03669
	CI $uutt$	$2(SS) \geq 3e, \mu \geq 1b, \geq 1j$	Yes	20.3	A 4.9 TeV $ C_{\text{coil}} = 1$	1504.04605
LO	Scalar LO 1 st gen	$2e$	$\geq 2j$	3.2	LO mass 1.1 TeV	$g_s = 0.25, g_t = 1.0, m(\chi) < 250 \text{ GeV}$ 1604.07773
	Scalar LO 2 nd gen	2μ	$\geq 2j$	3.2	LO mass 1.05 TeV	$g_s = 0.25, g_t = 1.0, m(\chi) < 150 \text{ GeV}$ 1604.01306
	Scalar LO 3 rd gen	$1e, \mu$	$\geq 1b, \geq 3j$	Yes	20.3	LO mass 640 GeV
Heavy quarks	VLO $TT \rightarrow Ht + X$	$1e, \mu$ $\geq 2b, \geq 3j$	Yes	20.3	T mass 855 GeV	T in (TB) doublet 1505.04306
	VLO $YY \rightarrow Wb + X$	$1e, \mu$ $\geq 1b, \geq 3j$	Yes	20.3	Y mass 770 GeV	Y in (BY) doublet 1505.04306
	VLO $BB \rightarrow Hb + X$	$1e, \mu$ $\geq 2b, \geq 3j$	Yes	20.3	B mass 735 GeV	isospin singlet 1505.04306
	VLO $BB \rightarrow Zb + X$	$2l \geq 3e, \mu$ $\geq 2l \geq 1b$	-	20.3	B mass 755 GeV	B in (BY) doublet 1409.5500
	VLO $QQ \rightarrow WqWq$	$1e, \mu$ $\geq 4j$	Yes	20.3	Q mass 690 GeV	1509.04261
	VLO $T_{5/3} T_{5/3} \rightarrow WtWt$	$2(SS) \geq 3e, \mu \geq 1b, \geq 1j$	Yes	3.2	$T_{5/3}$ mass 990 GeV	ATLAS-CONF-2016-032
Excited fermions	Excited quark $q^* \rightarrow q\gamma$	1γ	$1j$	3.2	q^* mass 4.4 TeV	only u^* and d^* , $\Lambda = m(q^*)$ 1512.05910
	Excited quark $q^* \rightarrow qg$	-	$2j$	15.7	q^* mass 5.6 TeV	only u^* and d^* , $\Lambda = m(q^*)$ ATLAS-CONF-2016-069
	Excited quark $b^* \rightarrow b\gamma$	-	$1b, 1j$	8.8	b^* mass 2.3 TeV	ATLAS-CONF-2016-060
	Excited quark $b^* \rightarrow Wt$	$1 \text{ or } 2e, \mu$ $1b, 2-0j$	Yes	20.3	b^* mass 1.5 TeV	$f_b = f_t = f_c = 1$ 1510.02664
	Excited lepton ℓ^*	$3e, \mu$	-	20.3	ℓ^* mass 3.0 TeV	$\Lambda = 3.0 \text{ TeV}$ 1411.2921
	Excited lepton ν^*	$3e, \mu, \tau$	-	20.3	ν^* mass 1.6 TeV	$\Lambda = 1.6 \text{ TeV}$ 1411.2921
Other	LSTC $a_T \rightarrow W\gamma$	$1e, \mu, 1\gamma$	Yes	20.3	a_T mass 960 GeV	1407.8150
	LRSM Majorana ν	$2e, \mu$ $2j$	-	20.3	N^* mass 2.0 TeV	1506.06020
	Higgs triplet $H^{\pm\pm} \rightarrow ee$	$2e$ (SS)	-	13.9	$H^{\pm\pm}$ mass 570 GeV	DY production, $BR(H^{\pm\pm} \rightarrow ee)=1$ ATLAS-CONF-2016-051
	Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$	$3e, \mu, \tau$	-	20.3	$H^{\pm\pm}$ mass 400 GeV	DY production, $BR(H^{\pm\pm} \rightarrow \ell\tau)=1$ 1411.2921
	Monotop (non-res prod)	$1e, \mu$ $1b$	Yes	20.3	spin-1 invisible particle mass 657 GeV	$a_{\text{non-res}} = 0.2$ 1410.5404
	Multi-charged particles	-	-	20.3	multi-charged particle mass 785 GeV	DY production, $ g = 5e$ 1504.04188
Magnetic monopoles	-	-	7.0	monopole mass 1.34 TeV	DY production, $ g = 1g_D, \text{spin } 1/2$ 1509.08059	

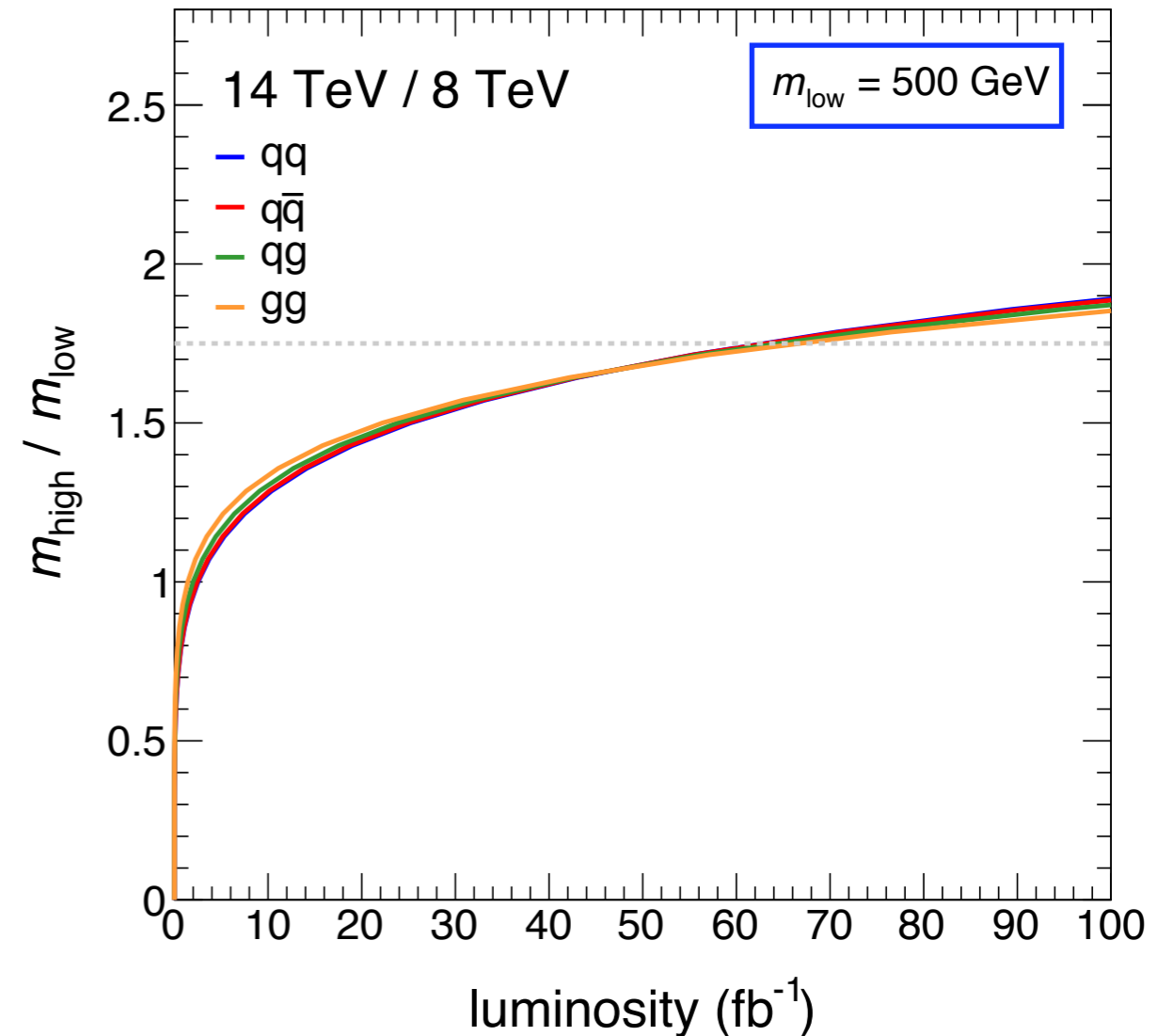
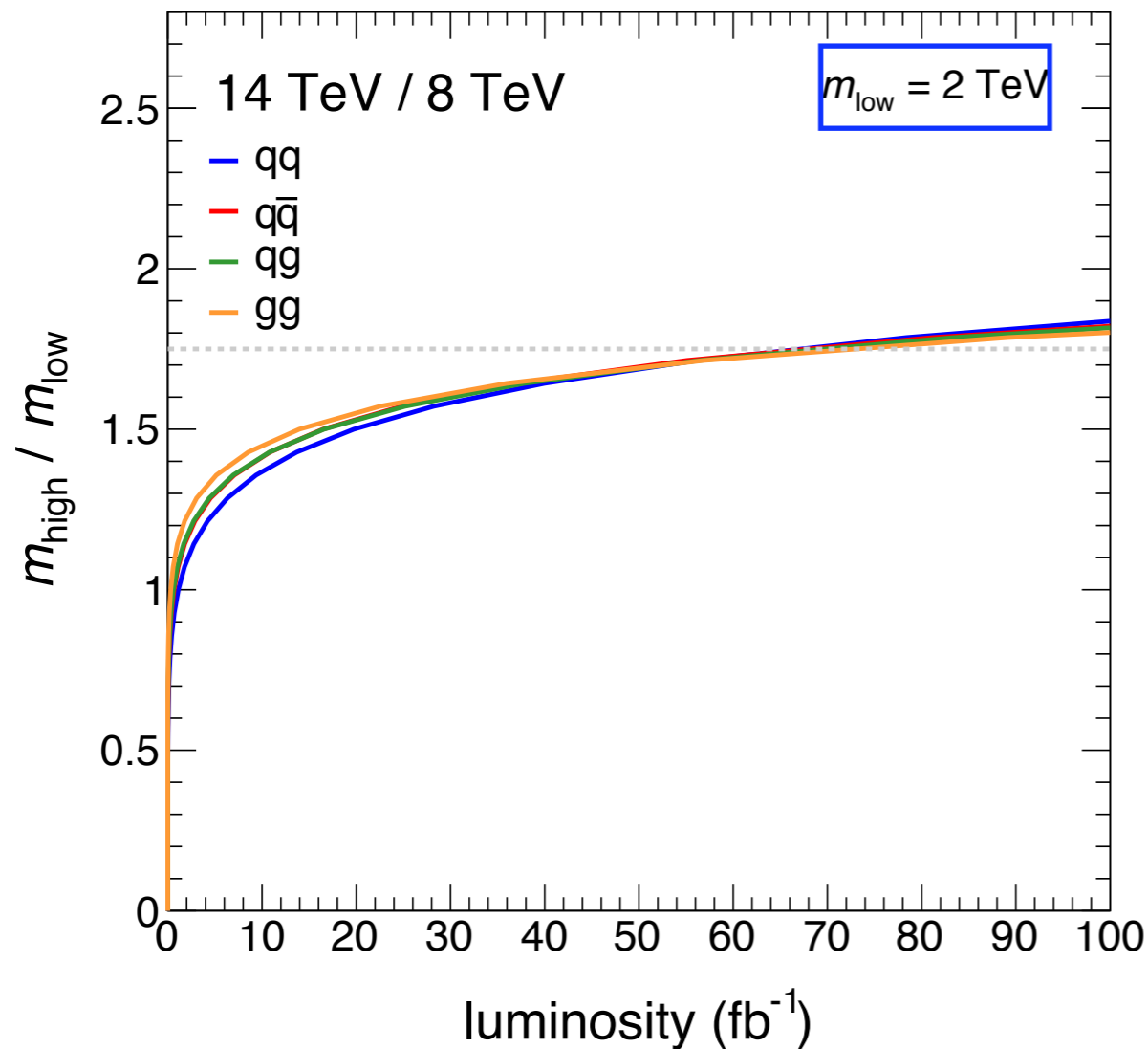
*Only a selection of the available mass limits on new states or phenomena is shown. Lower bounds are specified only when explicitly not excluded.

- No "early" discovery.
- Disappointed? Yes.
- Surprised? Not much.

As data accumulates

Run I limit 2 TeV, e.g. pair of 1 TeV gluino.

500 GeV, e.g. pair of 250 GeV electroweak-ino



Rapid gain initial 10s fb^{-1} , slow improvements afterwards.

Reaching the “slow” phase after Moriond 2017

LHC Run 2 will continue to pursue a broad physics program.

Of course, there are gaps in to be filled, new signals to be looked at.

Still room for discovery.

This Lecture

This Lecture

- Focus on longer term future.

This Lecture

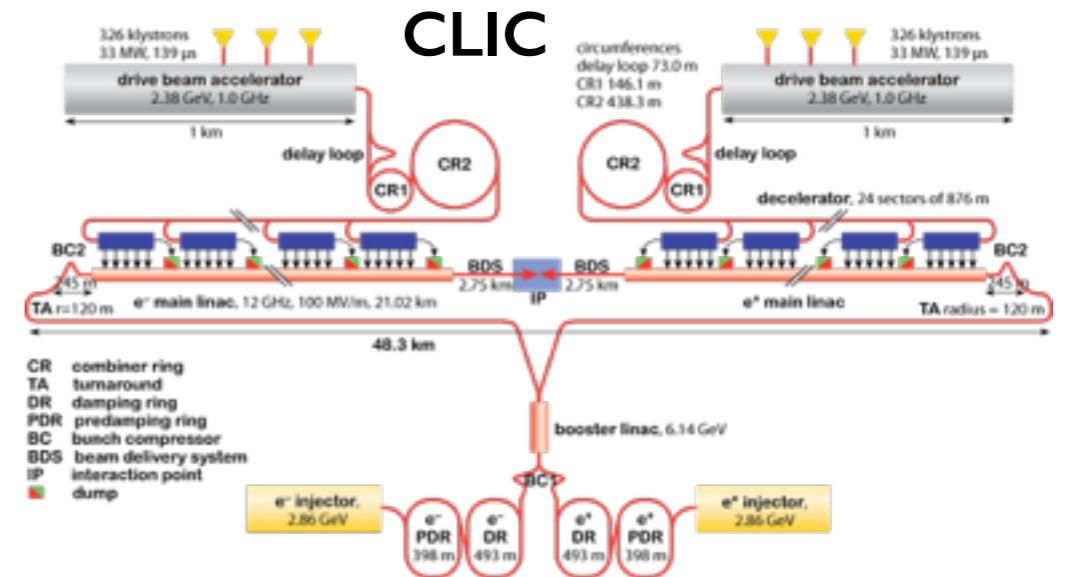
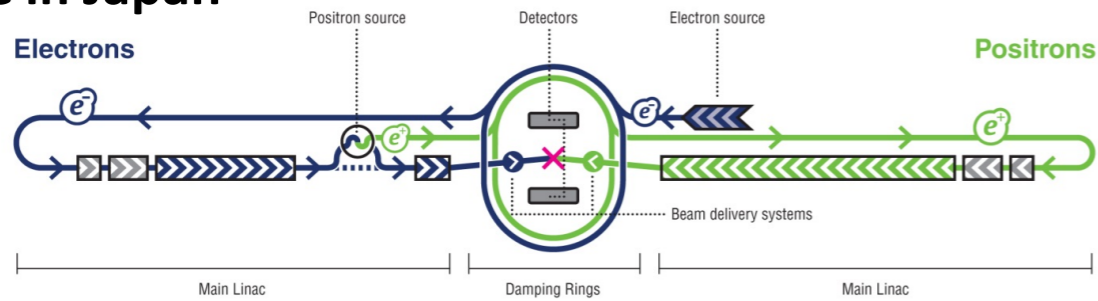
- Focus on longer term future.
- Assuming no discovery of new particle at the LHC.

This Lecture

- Focus on longer term future.
- Assuming no discovery of new particle at the LHC.
- Physics case for future hadron collider
 - ▶ Cover significant ground beyond the LHC.
 - ▶ Answering important questions beyond the reach of the LHC

Beyond the LHC, future facilities

ILC in Japan

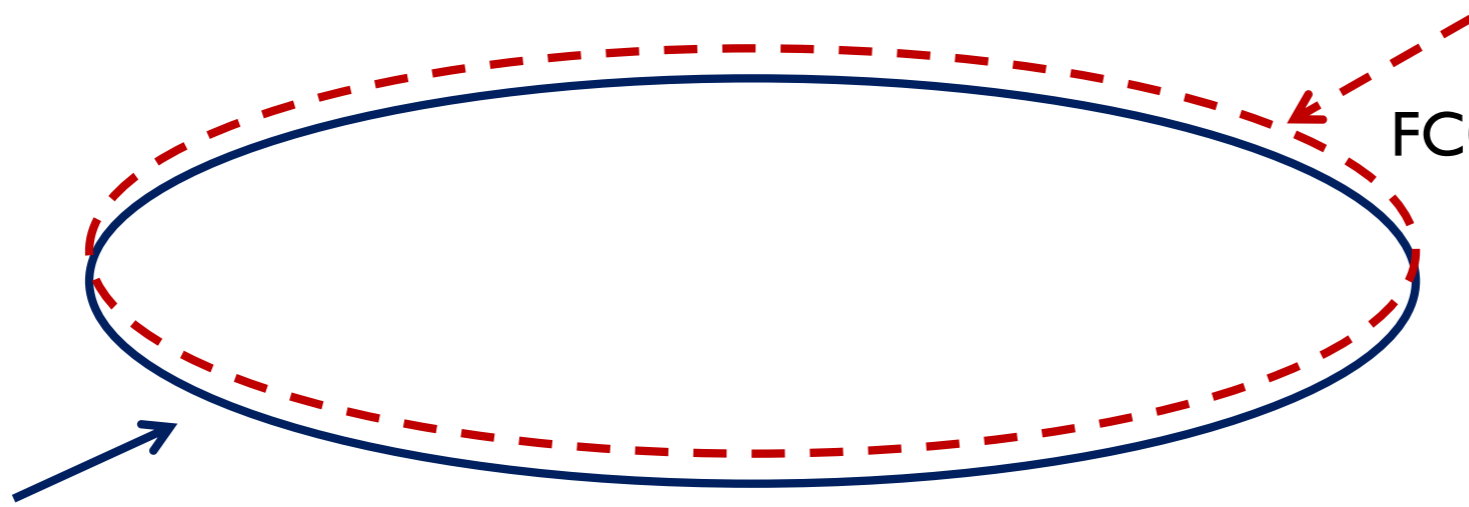


Circular. “Scale up” LEP+LHC

~100 TeV

pp collider

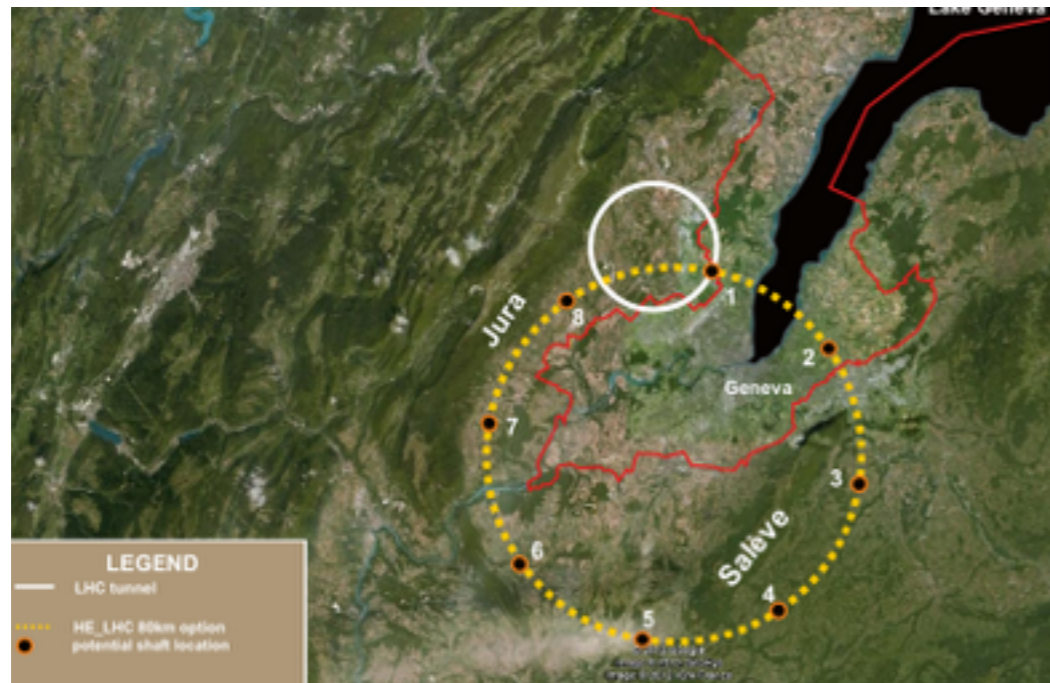
FCC-hh (CERN), SppC(China)



250 GeV **e⁻e⁺ Higgs Factory**

FCC-ee (CERN), CEPC(China)

Future circular colliders



CERN

Higgs factory: FCC-ee
pp Collider: FCC-hh



China.

Higgs factory: CEPC
pp Collider: SppC

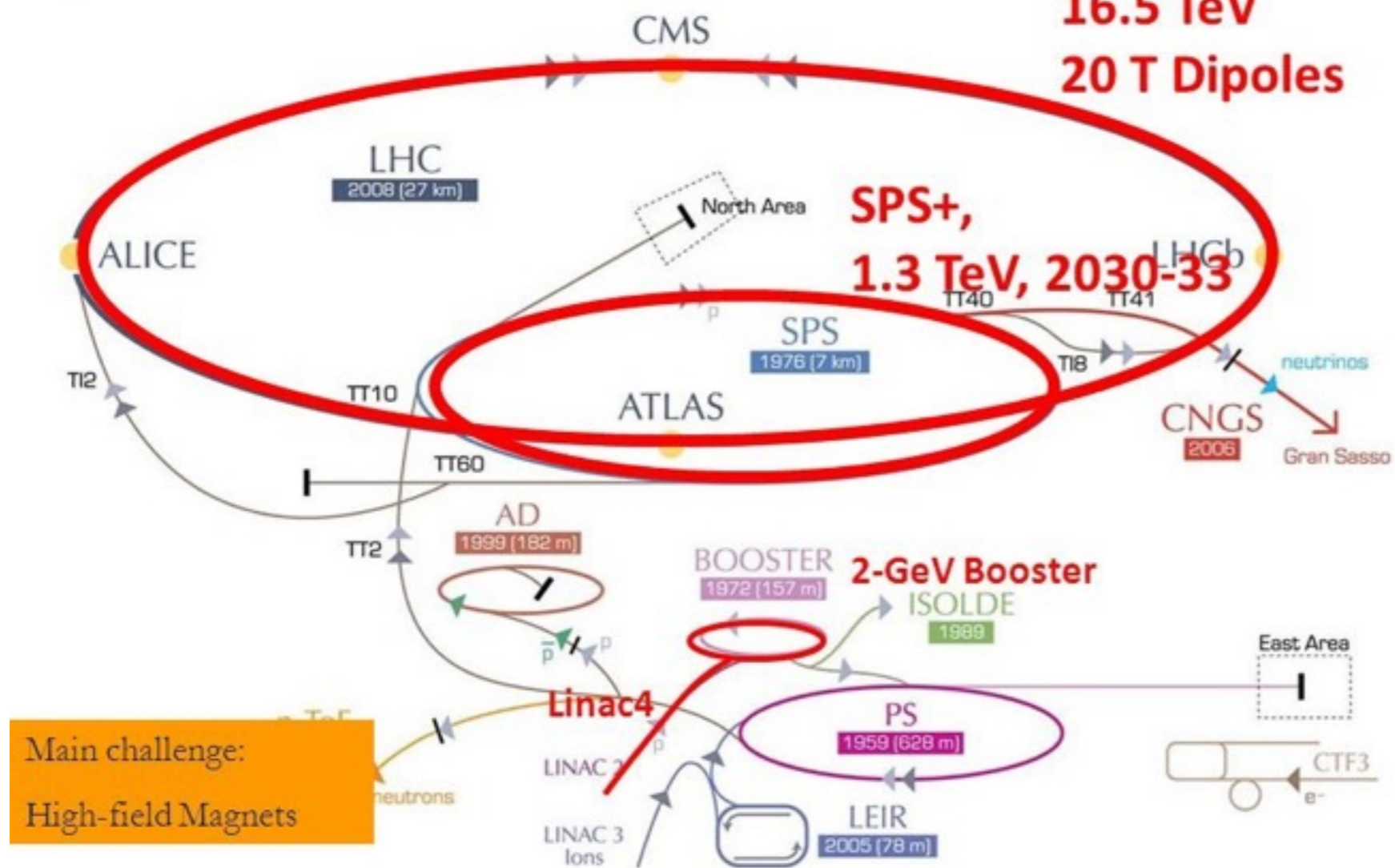
HE-LHC

High-Energy LHC (HE-LHC)?

HE-LHC >2035

16.5 TeV

20 T Dipoles

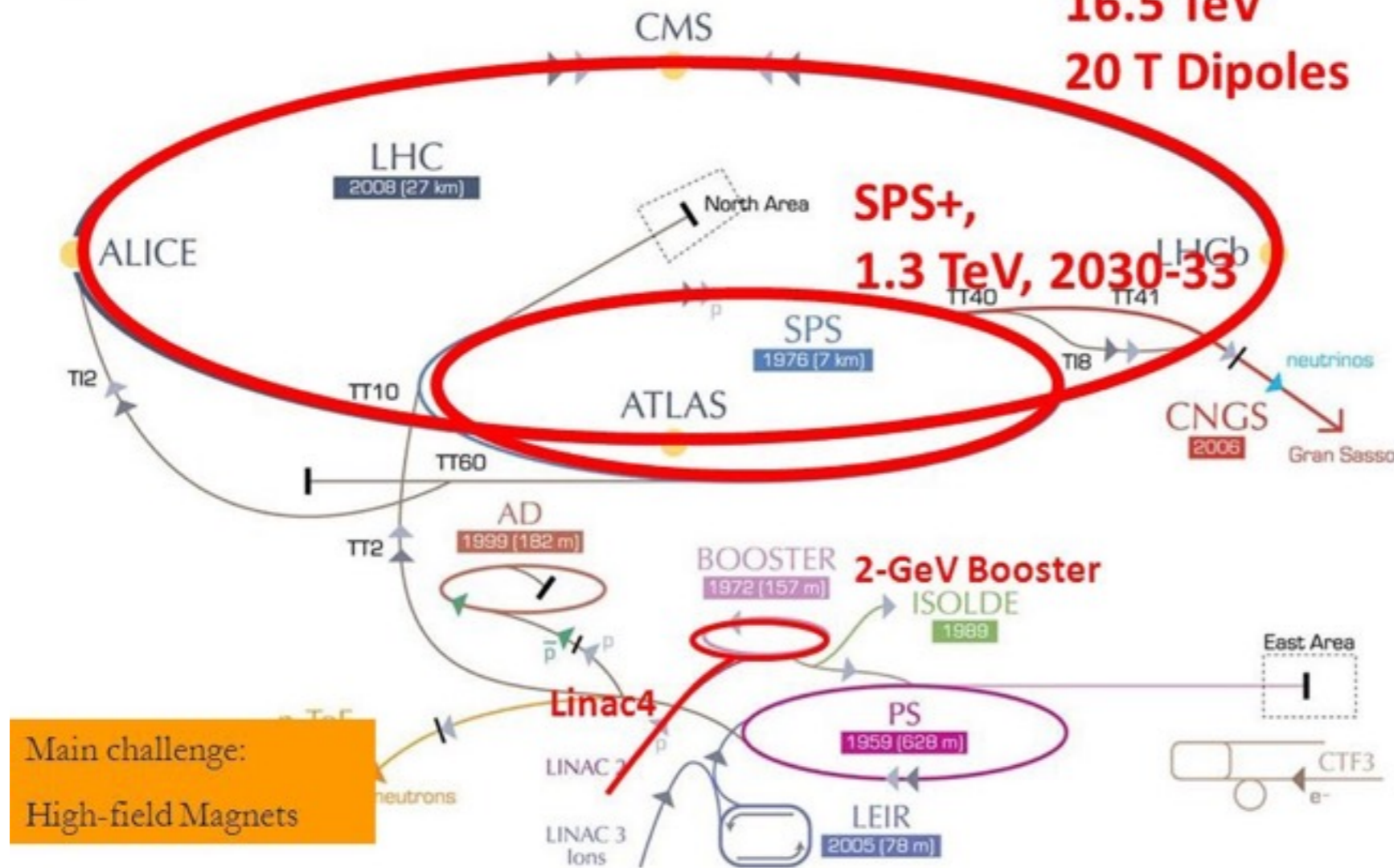


– 28 TeV more realistically?

HE-LHC

High-Energy LHC (HE-LHC)?

HE-LHC >2035
16.5 TeV
20 T Dipoles

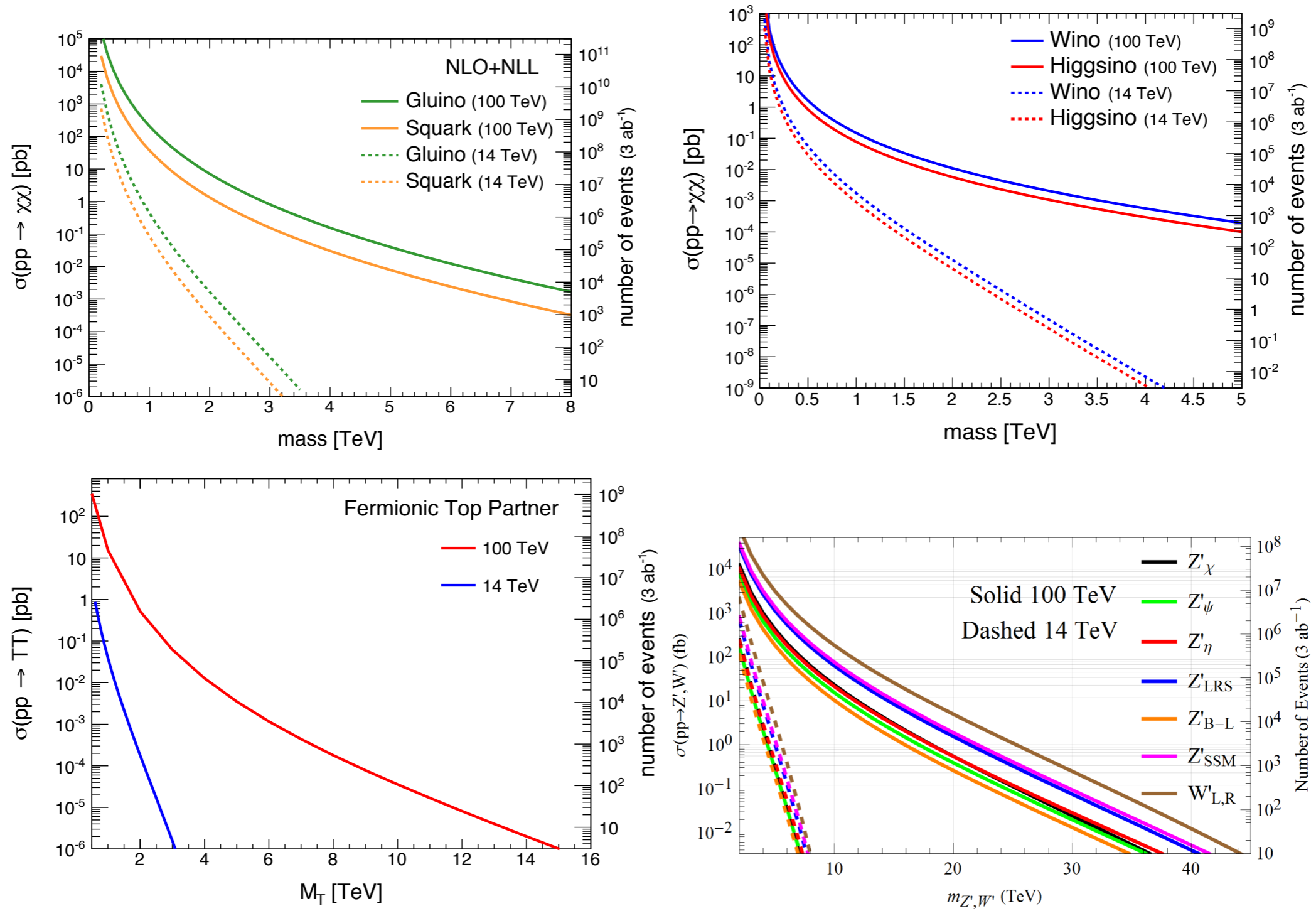


– 28 TeV more realistically?

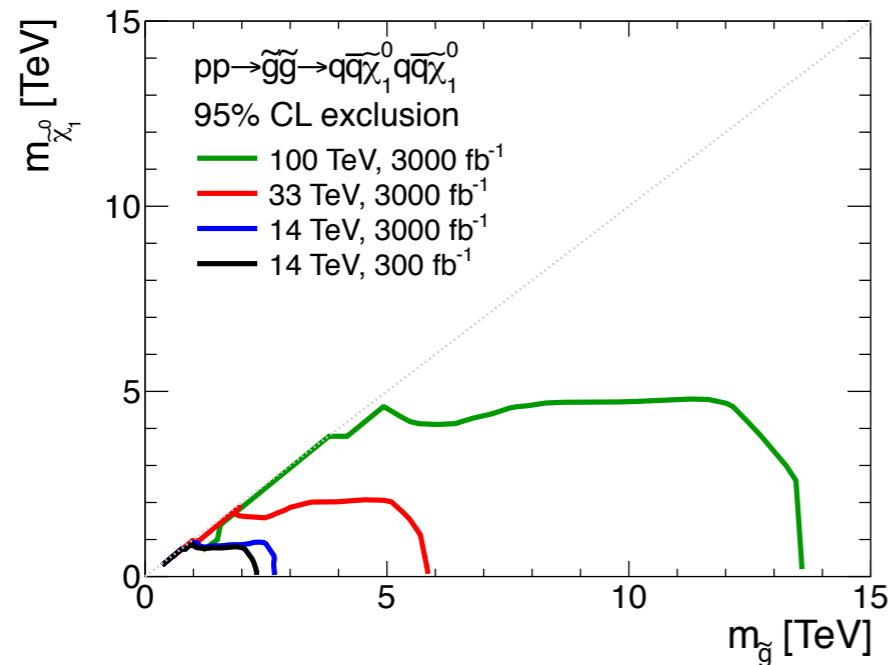
Will focus on 100 TeV collider here.

Basic physics capability

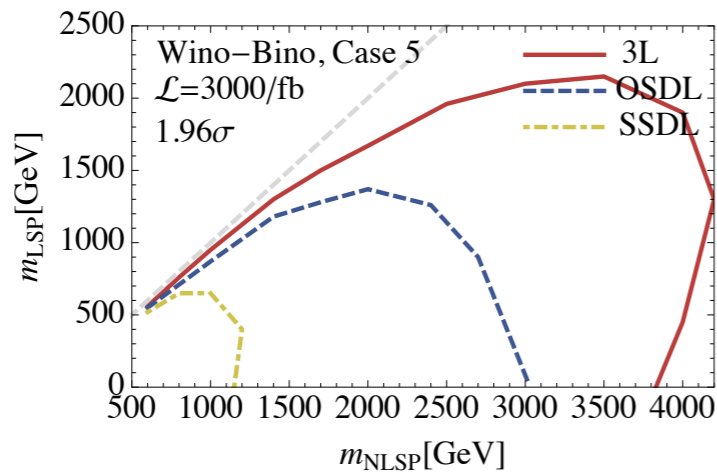
100 TeV pp collider, a big step in energy



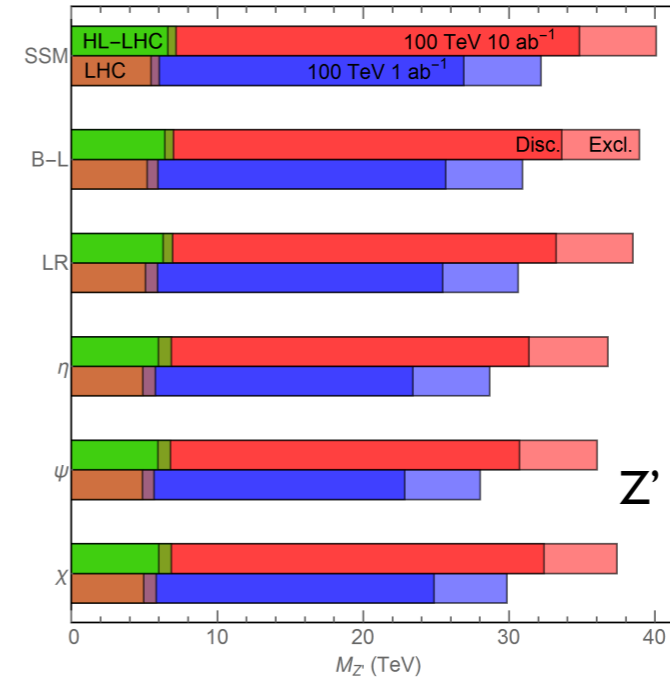
A big step forward in the energy frontier



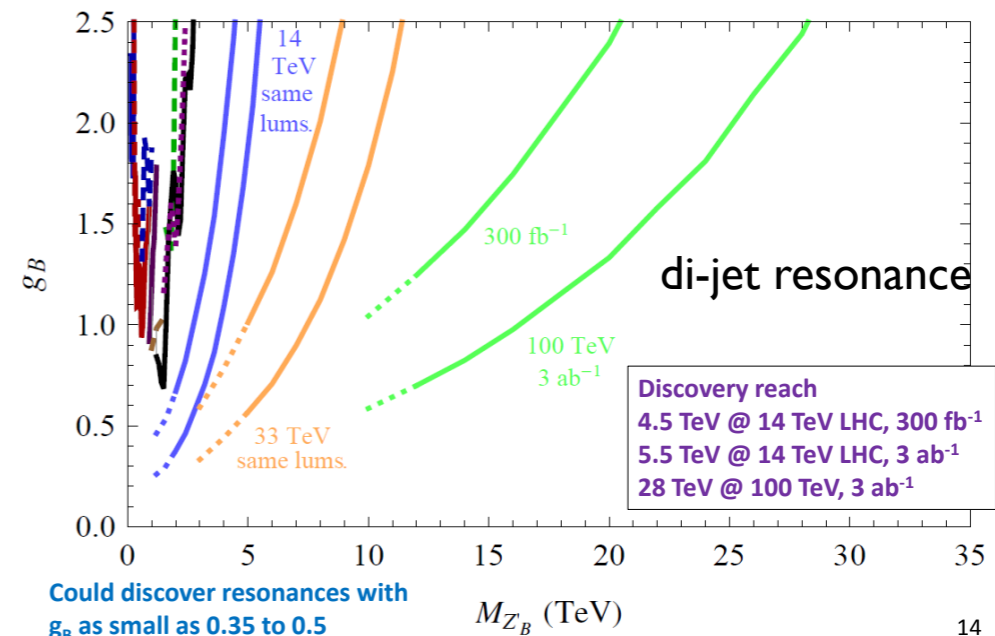
Cohen et al, 2013



Gori, Jung, LTW, Wells, 2014



Han, Langacker, Liu, LTW, to appear

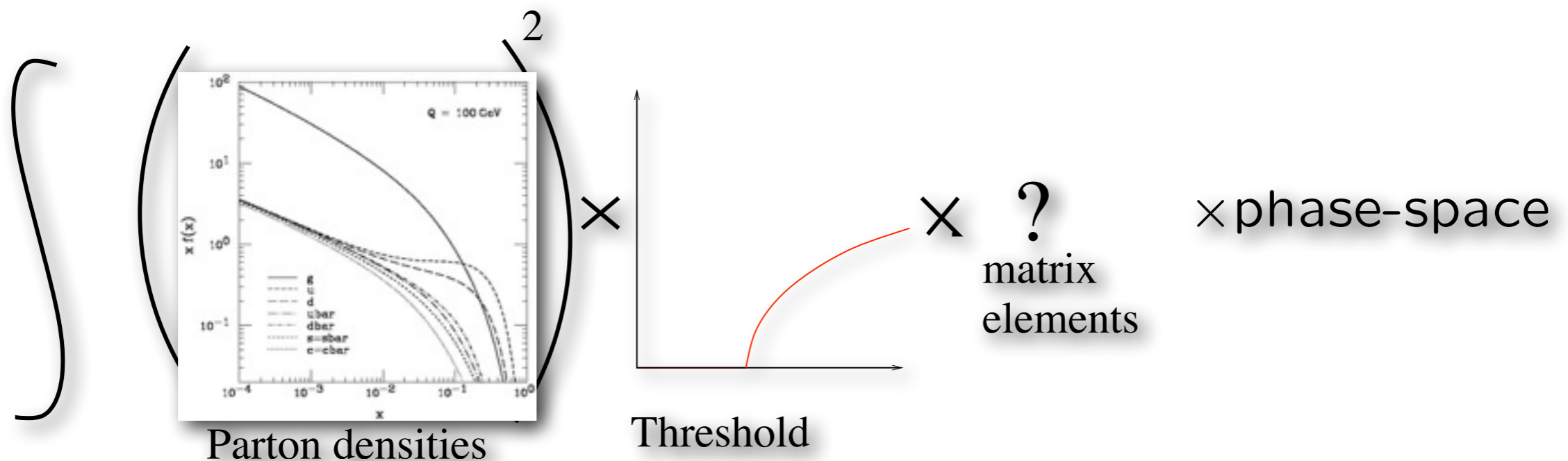


Felix Yu, 2013

cross the board: x 5(more) improvement, into (10)TeV regime

Production of new physics particles

- Schematics of production at hadron colliders.
- Dominated by parton densities and thresholds (mass and cut).



$$\frac{d^2\sigma(a, b \rightarrow \dots)}{d\hat{s} dY} = \frac{1}{\hat{s}} \sum_{a,b} x_1 f_a(x_1) x_2 f_b(x_2) \hat{\sigma}(a, b \rightarrow \dots)$$

Partonic cross section

Parton luminosity

- The cross section can be written as

$$\sigma = \sum_{a,b} \int d\tau \frac{dL_{ab}}{d\tau} \hat{\sigma}$$

$\frac{dL_{ab}}{d\tau}$ ← parton luminosity

$\tau = \frac{\hat{s}}{S} = x_1 x_2$

S : center of mass energy
 \hat{s} : parton center of mass energy

$$L_{ab}(\tau) = \frac{1}{1 + \delta_{ab}} \int_{\tau}^1 \frac{dx}{x} \left[f_a(x) f_b\left(\frac{\tau}{x}\right) + f_a\left(\frac{\tau}{x}\right) f_b(x) \right]$$

Parton luminosity

- The cross section can be written as

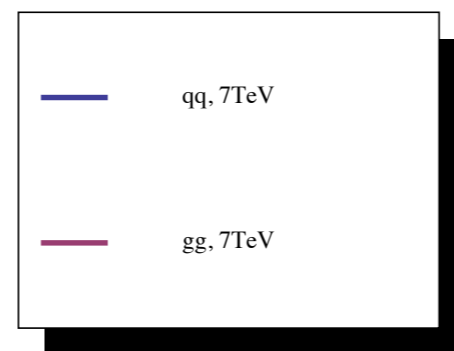
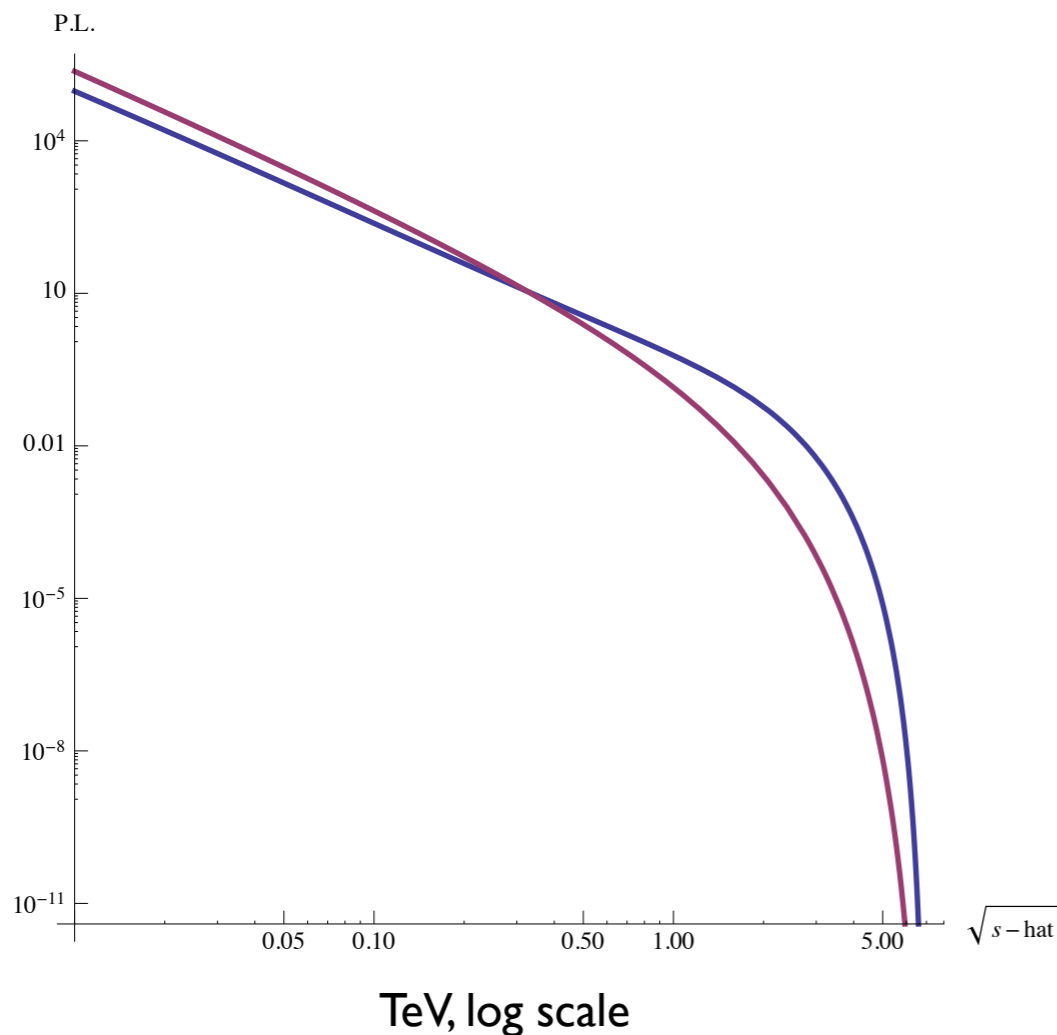
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parton luminosity

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S : center of mass energy
 \hat{s} : parton center of mass energy

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Very sharp falling

$$\propto \frac{1}{\tau^a}, \quad a \sim 3 - 7$$

Falls by a factor of 10 for every 600 GeV

⇒ Production dominantly on threshold

Rough estimates of discovery reach

$$\sigma \sim L_p \cdot \hat{\sigma} \sim \frac{1}{\tau^a} \hat{\sigma}$$

L_p : parton luminosity, $\hat{\sigma}$: parton cross section

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Production of new physics particle of mass M

Fast falling parton luminosity \Rightarrow
dominant contribution from
parton cross section near threshold

$$\hat{s} \sim M^2 \rightarrow \tau \sim \frac{M^2}{S}$$

$$\hat{\sigma} \sim \frac{1}{M^2}$$

Rough estimates of discovery reach

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$$\hat{\sigma} \sim \frac{1}{M^2}$$

Number of new physics particle produced:

$$N = \sigma \cdot \mathcal{L}$$

\mathcal{L} : luminosity

Discovery reach

Consider 2 colliders.

Collider 1: $E_{\text{cm}} = E_1$, or $S_1 = E_1^2$. Collider 2: $E_{\text{cm}} = E_2$, or $S_2 = E_2^2$.

$$E_2 > E_1$$

Reach for new physics at these 2 colliders

Collider 1: M_1 . Collider 2: M_2 .

Discovery reach

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Reach for new physics at these 2 colliders

Collider 1: M_1 . Collider 2: M_2 .

Assume the reach is obtained from the same number of signal events

$$\frac{1}{\tau_1^a} \frac{1}{M_1^2} \mathcal{L}_1 = \frac{1}{\tau_2^a} \frac{1}{M_2^2} \mathcal{L}_2 \quad \text{used} \quad \hat{\sigma} \sim \frac{1}{M^2}$$

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We have

$$\frac{M_2}{M_1} = \left(\frac{S_2}{S_1} \right)^{1/2} \left(\frac{S_1 \mathcal{L}_2}{S_2 \mathcal{L}_1} \right)^{\frac{1}{2a+2}} \quad \text{used} \quad \hat{s} \sim M^2 \rightarrow \tau \sim \frac{M^2}{S}$$

Discovery reach

$$\frac{M_2}{M_1} = \left(\frac{S_2}{S_1}\right)^{1/2} \left(\frac{S_1 \mathcal{L}_2}{S_2 \mathcal{L}_1}\right)^{\frac{1}{2a+2}}$$

$M_2 > M_1$ if $S_2 > S_1$

Large gain with higher energy

Discovery reach

$$\frac{M_2}{M_1} = \left(\frac{S_2}{S_1}\right)^{1/2} \left(\frac{S_1 \mathcal{L}_2}{S_2 \mathcal{L}_1}\right)^{\frac{1}{2a+2}}$$

$$M_2 > M_1 \text{ if } S_2 > S_1$$

Large gain with higher energy

If we want $\frac{M_2}{M_1} \sim \frac{E_2}{E_1} = \left(\frac{S_2}{S_1}\right)^{1/2}$ We need $\frac{S_2}{S_1} = \frac{\mathcal{L}_2}{\mathcal{L}_1}$

That is, a factor of 50 more luminosity going from 14 TeV to 100 TeV.

From HL-LHC, we will have 3 ab⁻¹. For 100 TeV, we need 150 ab⁻¹.

A lot!

However, situation is actually better.

Discovery reach

$$\frac{M_2}{M_1} = \left(\frac{S_2}{S_1}\right)^{1/2} \left(\frac{S_1 \mathcal{L}_2}{S_2 \mathcal{L}_1}\right)^{\frac{1}{2a+2}}$$

a is large (3-7).

The second factor on r.h.s is increasing slowly with large luminosity
i.e., not losing that much without very large luminosity.

Discovery reach

$$\frac{M_2}{M_1} = \left(\frac{S_2}{S_1}\right)^{1/2} \left(\frac{S_1 \mathcal{L}_2}{S_2 \mathcal{L}_1}\right)^{\frac{1}{2a+2}}$$

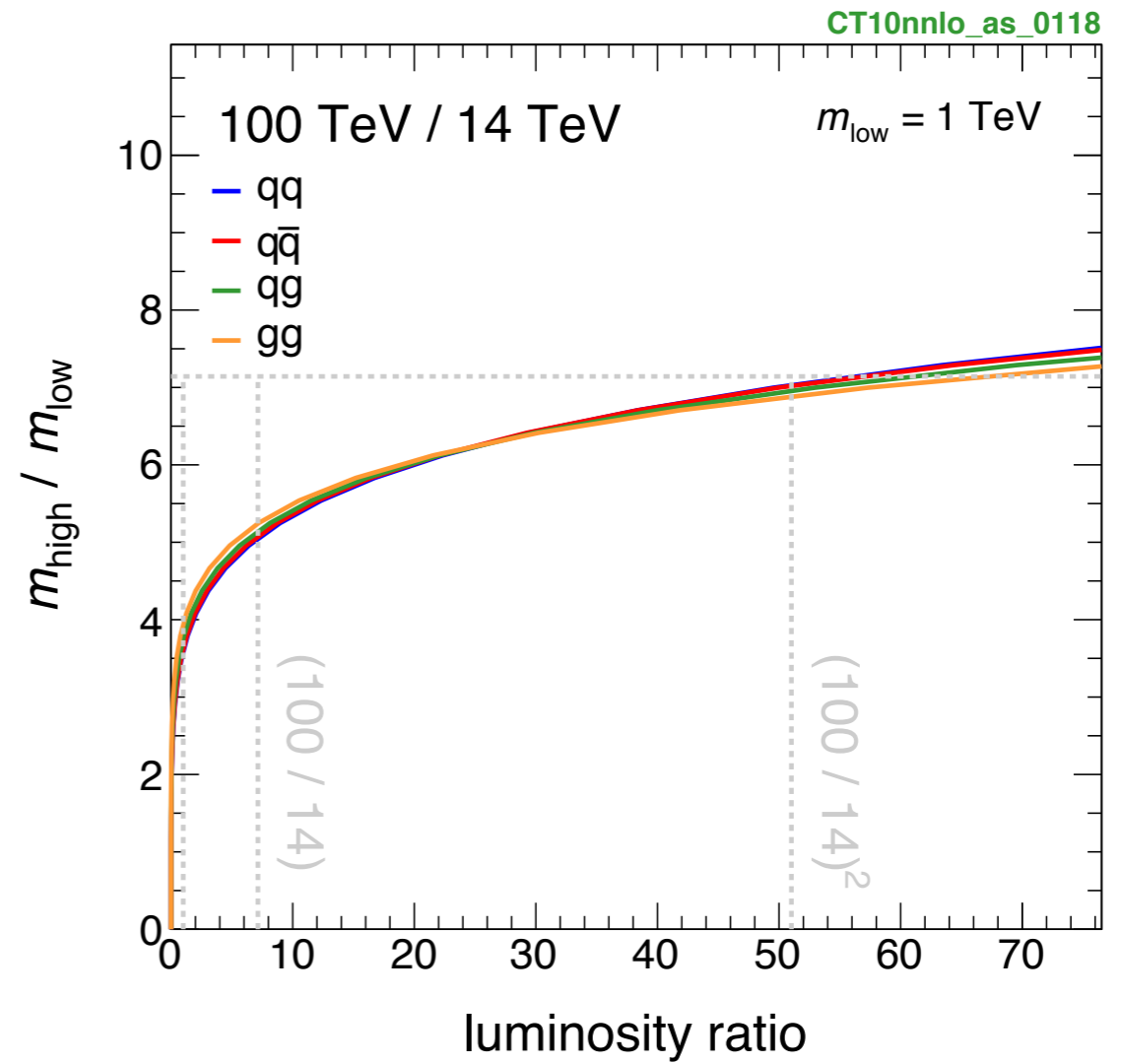
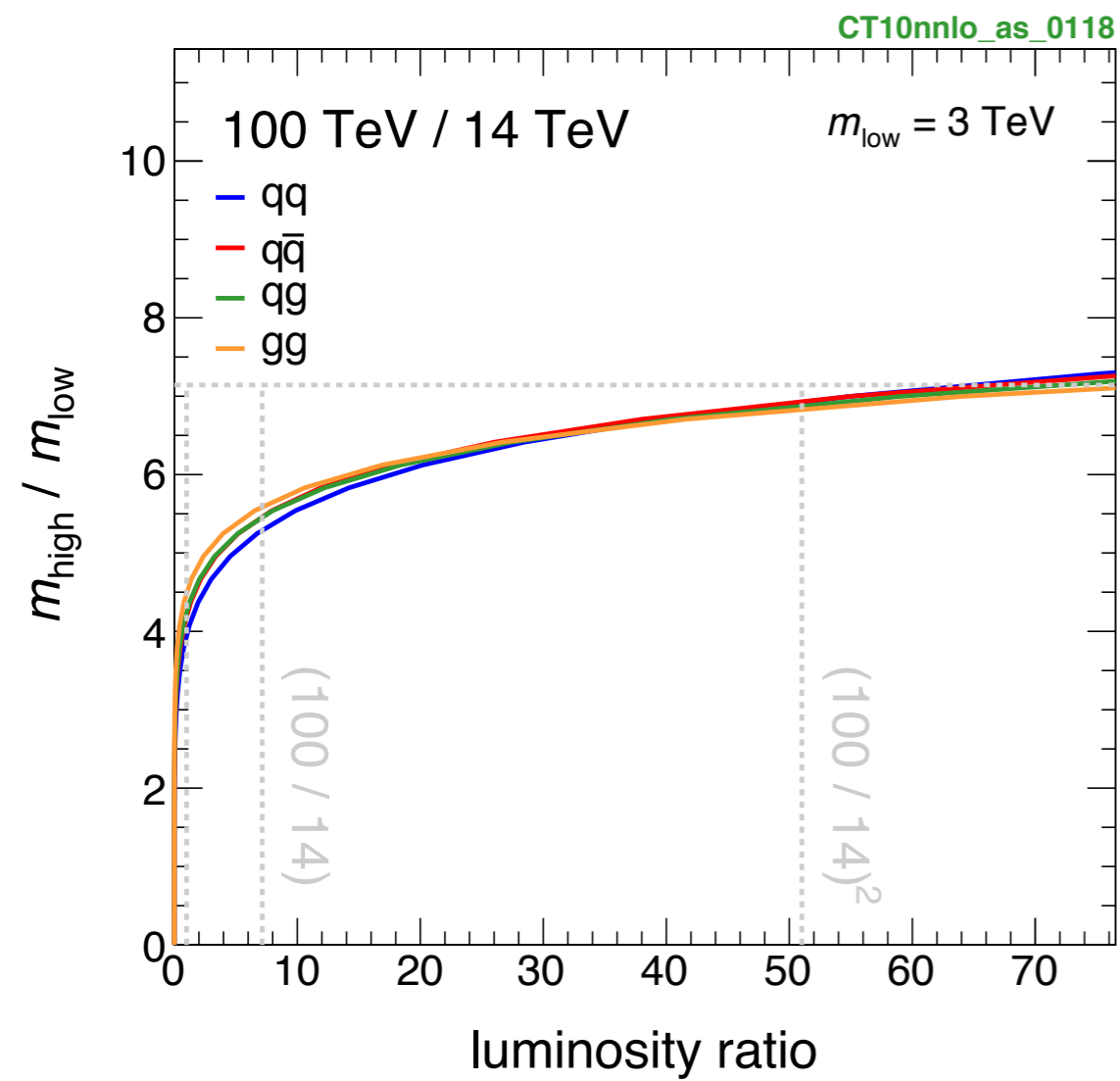
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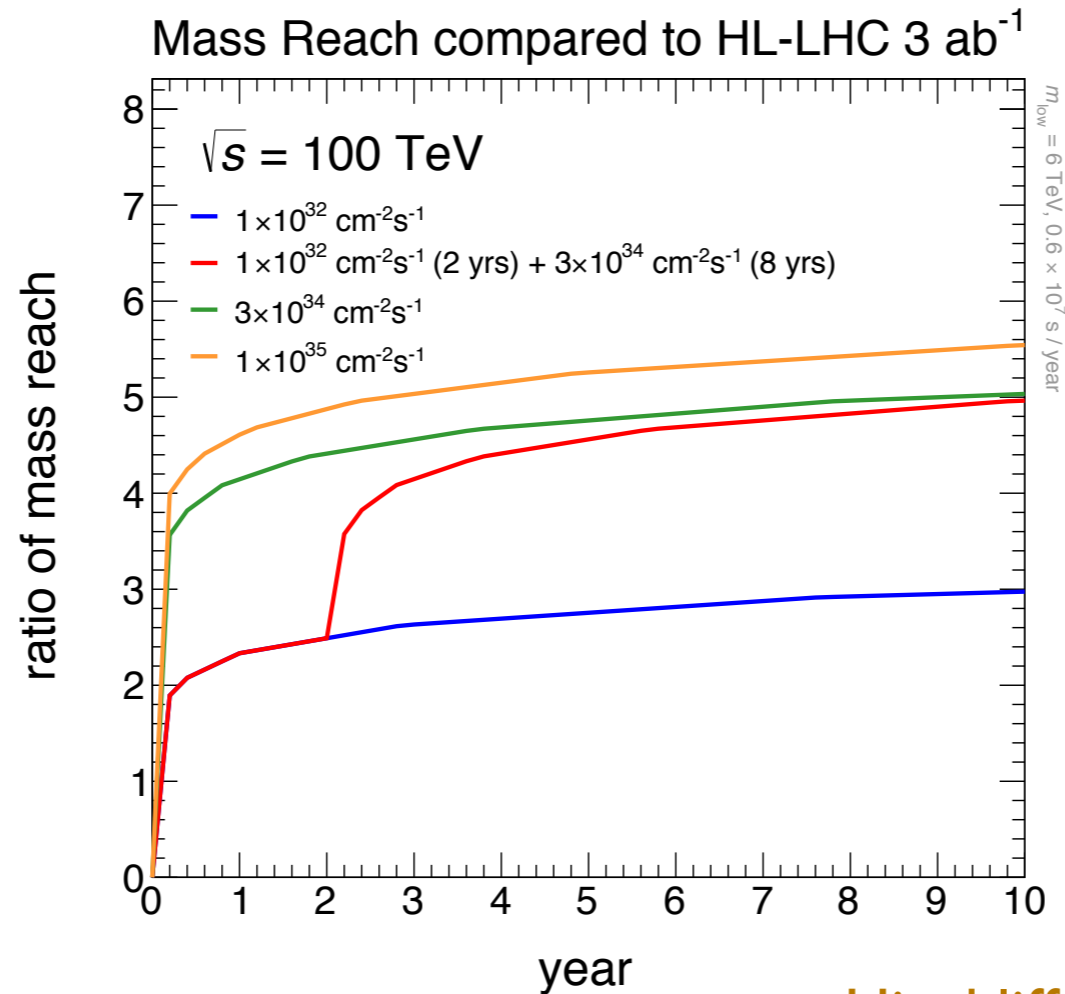
In particular, for the same collider, as luminosity increases

$$\frac{M_2}{M_1} = \exp\left(\frac{1}{2a+2} \log(\mathcal{L}_2/\mathcal{L}_1)\right) \simeq 1 + \frac{1}{2a+2} \log(\mathcal{L}_2/\mathcal{L}_1)$$

Discovery reach



100-ish TeV pp collider




Hinchliffe, Kotwal, Mangano, Quigg, LTW

A factor of about 5 increase in reach
with modest luminosity

Status of circular collider studies

- In the past 2 years, many studies of the physics reaches of the circular colliders have been carried out.
 - ▶ On both FCC and CEPC/SppC.
- Preliminary physics case has been made.
- Active efforts in trying to make it happen. Prospect will be clearer in the coming several years.

Status of circular collider studies

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 - ▶ On both FCC and CEPC/SppC.
- Preliminary physics case has been made.  rest of this
- Active efforts in trying to make it happen. this lecture
Prospect will be clearer in the coming several years.

Open questions beyond LHC

- Nature of electroweak symmetry breaking.
- Naturalness.
- Dark matter.
-

Need to go beyond

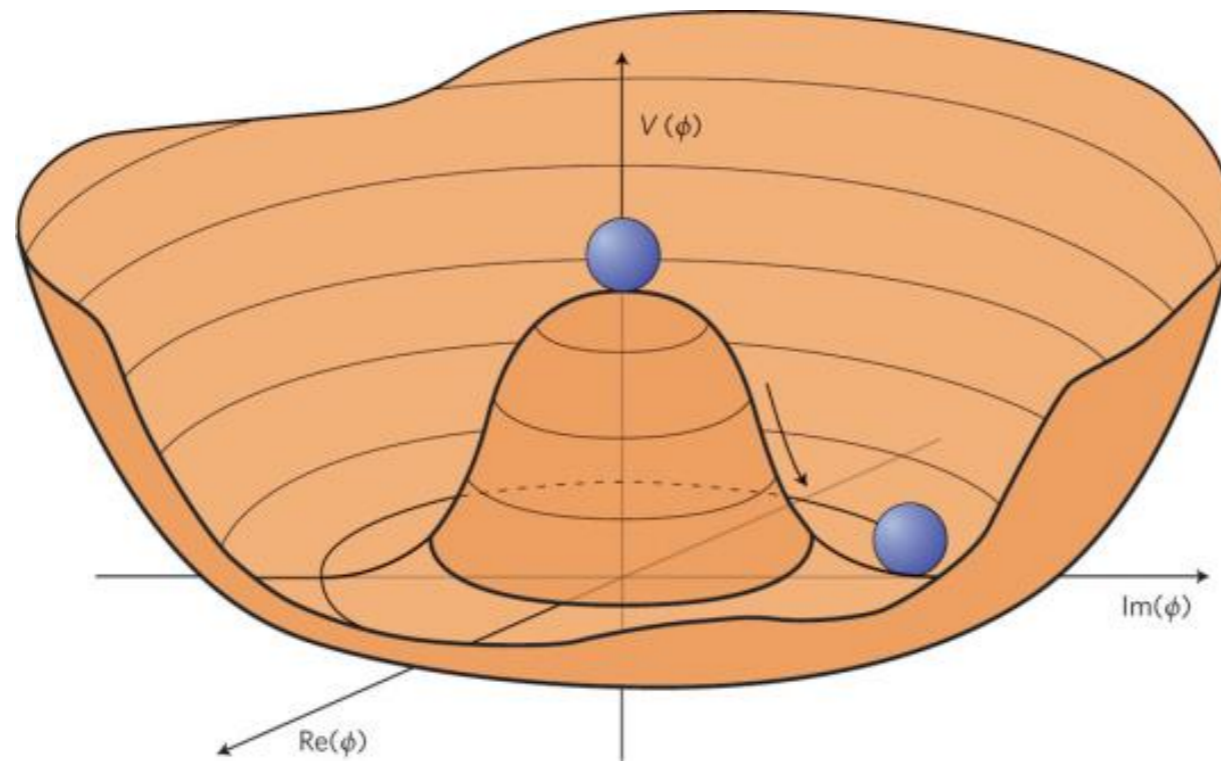
Nature of electroweak symmetry breaking

Higgs is special

particle	spin
quark: u, d,...	1/2
lepton: e...	1/2
photon	1
W,Z	1
gluon	1
Higgs	0

**h: a new kind of
elementary particle**

“Simple” picture: Mexican hat

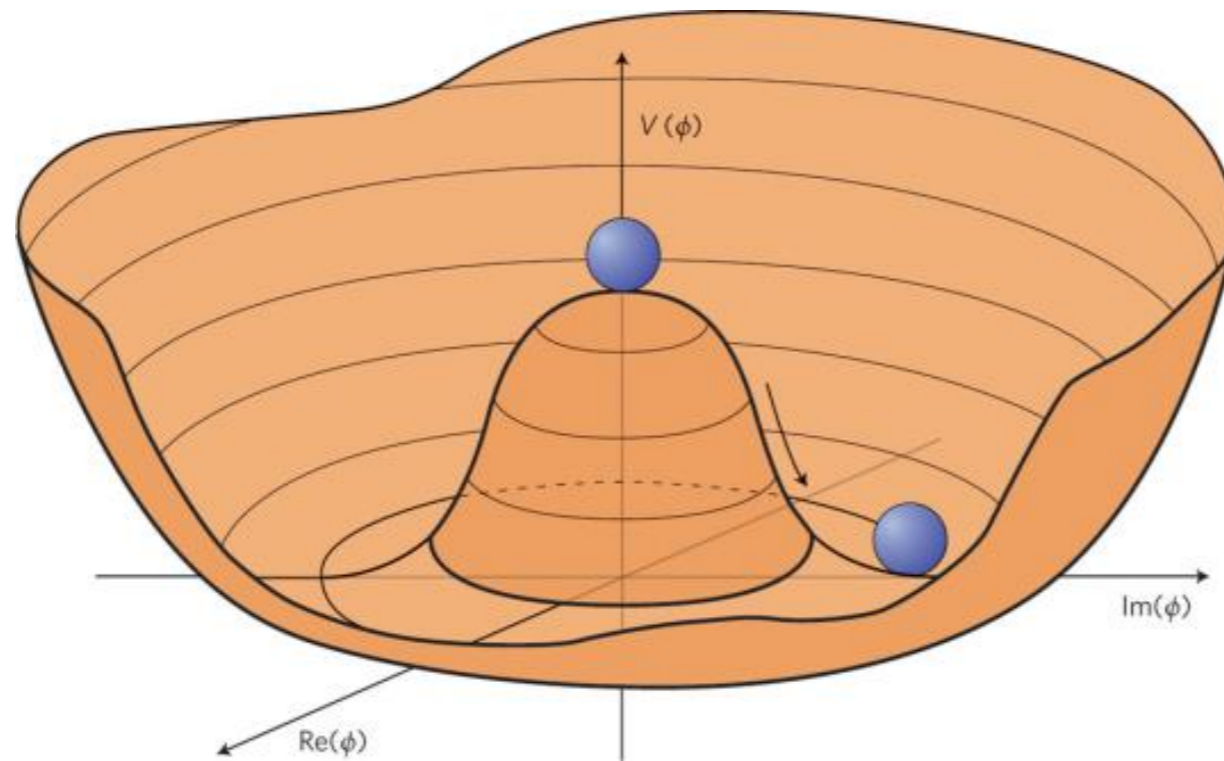


$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4}h^4$$

$$\langle h \rangle \equiv v \neq 0 \rightarrow m_W = g_W \frac{v}{2}$$

Similar to, and motivated by Landau-Ginzburg theory of superconductivity.

"Simple" picture: Mexican hat

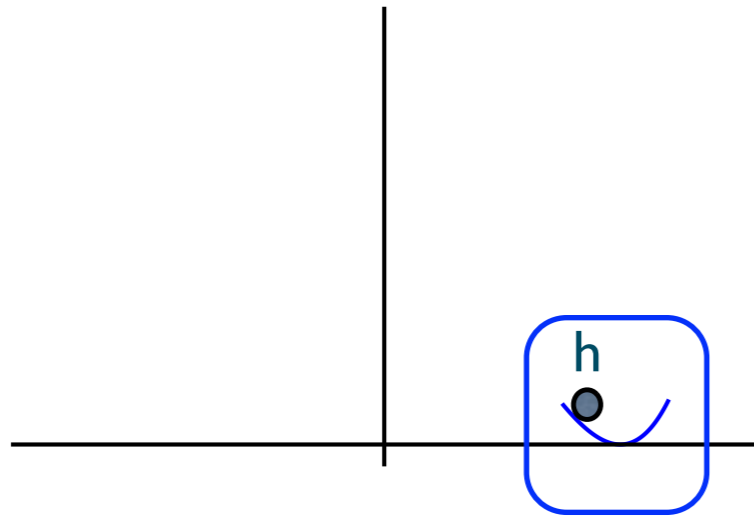


$$V(h) = \frac{1}{2}\mu^2 h^2 + \frac{\lambda}{4} h^4$$
$$\langle h \rangle \equiv v \neq 0 \rightarrow m_W = g_W \frac{v}{2}$$

Similar to, and motivated by
Landau-Ginzburg theory
of superconductivity.

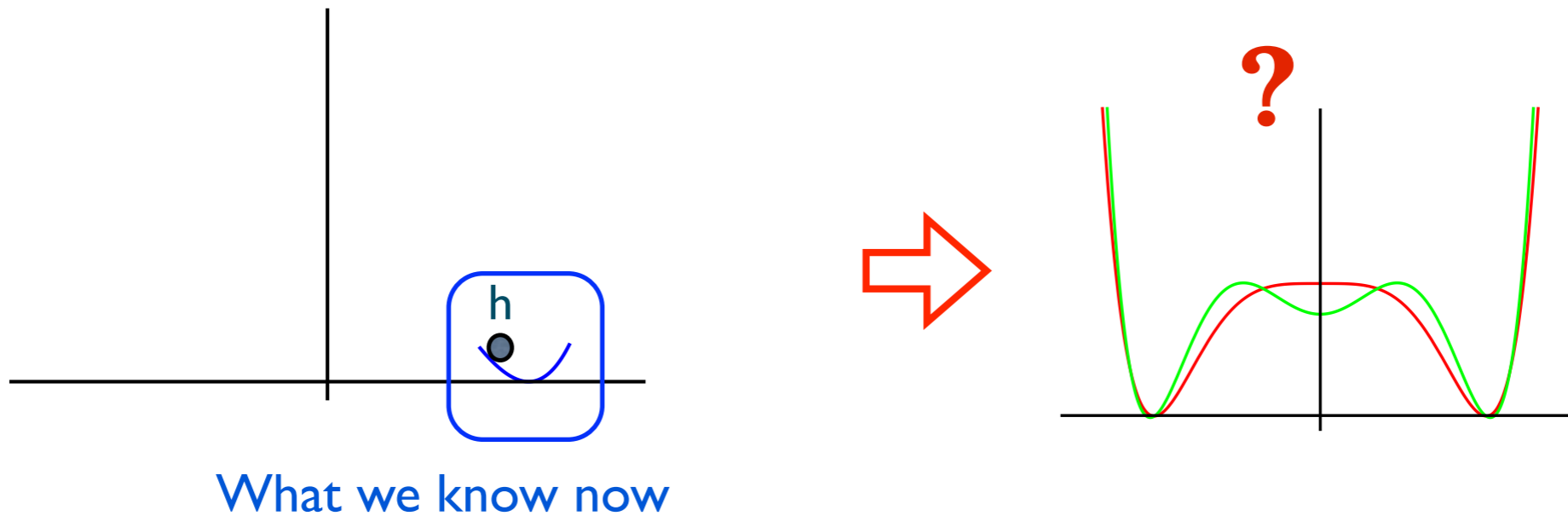
However, this simplicity is deceiving.
Parameters not predicted by theory. Need new physics

Not even sure about "Mexican hat".



What we know now

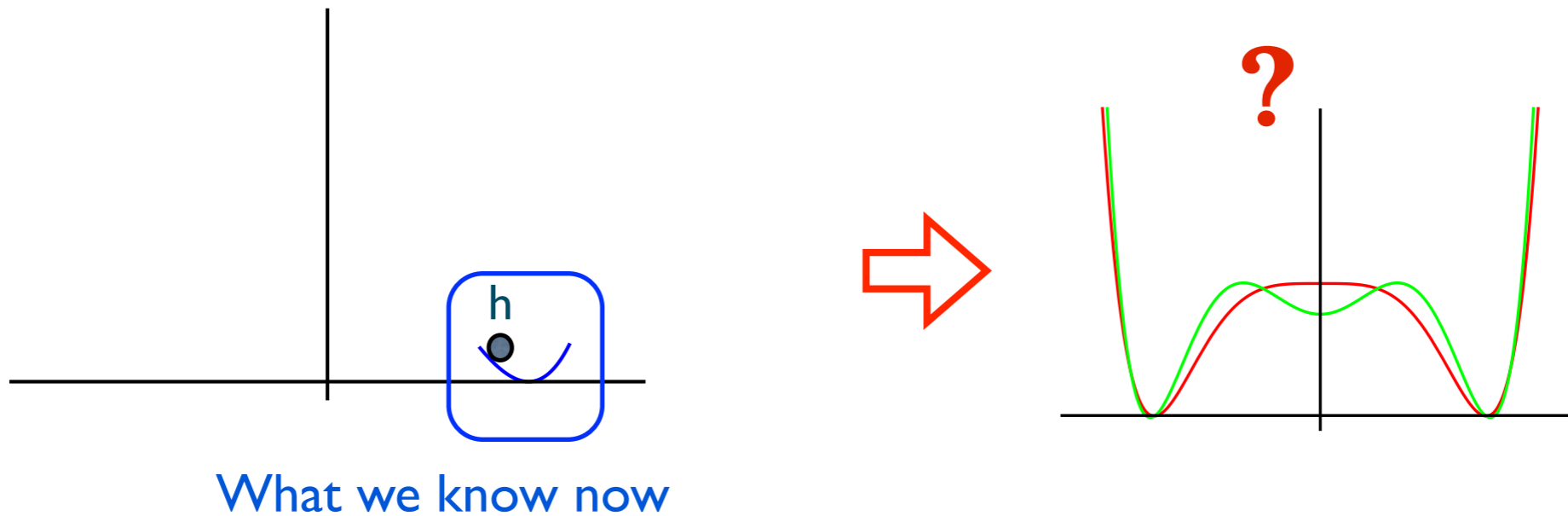
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Is the EW phase transition first order?

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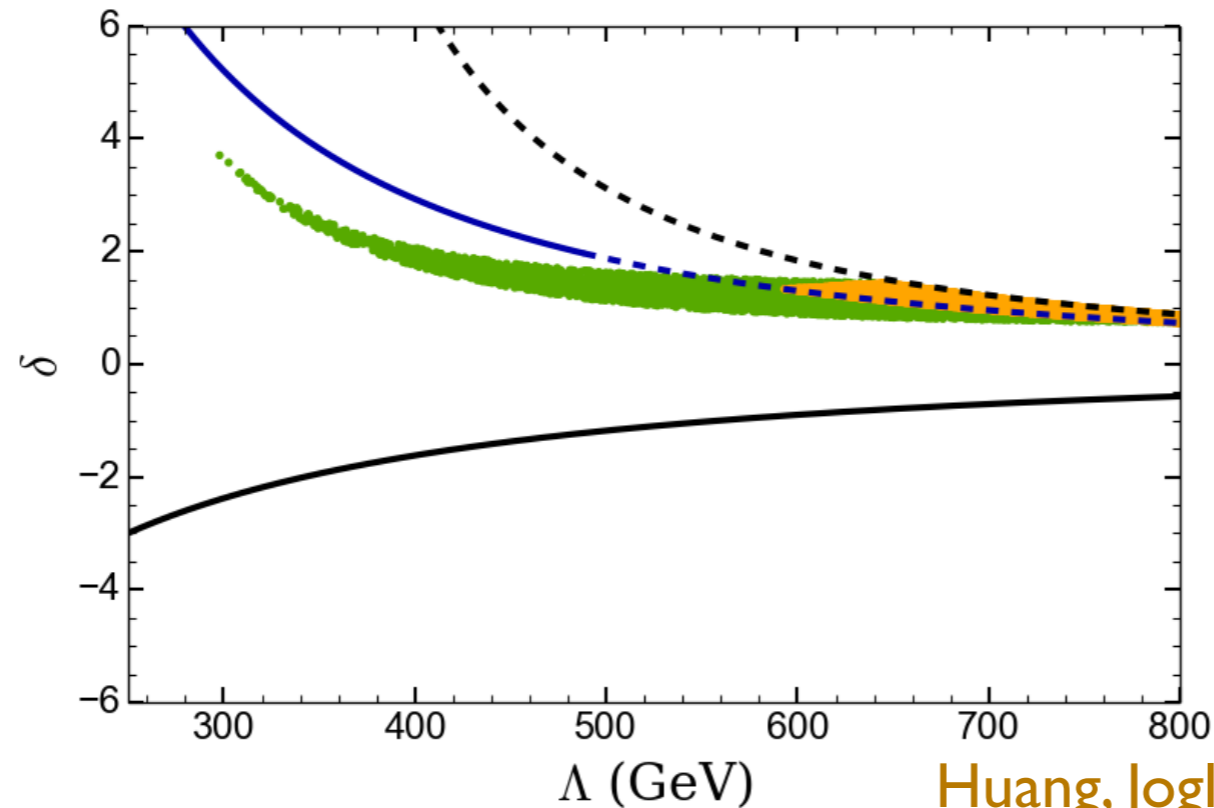
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Is the EW phase transition first order?

LHC can not distinguish these definitively.

1st order phase transition

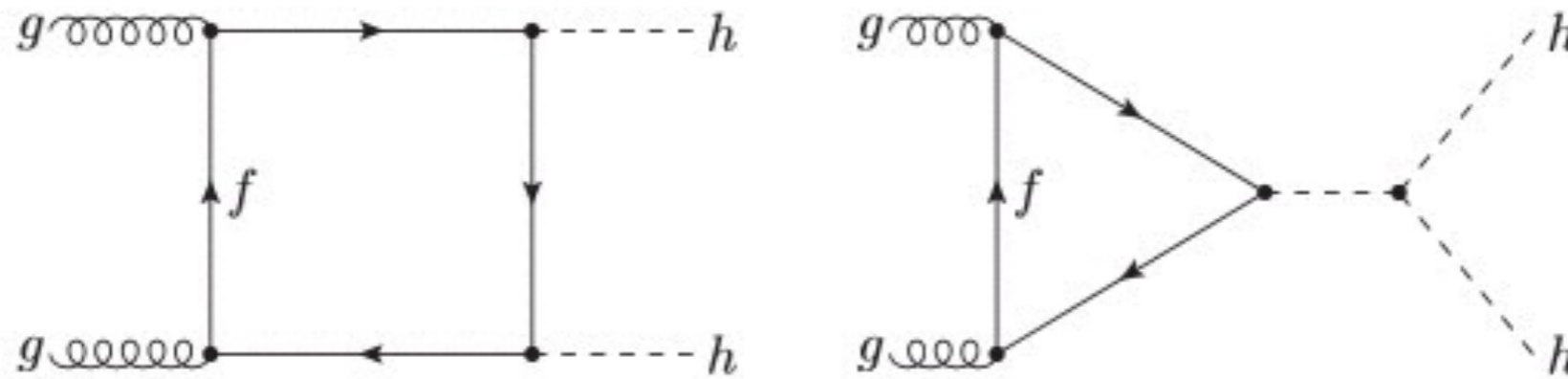
⇒ large modification of trilinear coupling



Huang, Joglekar, Li, Wagner, 1512.00068

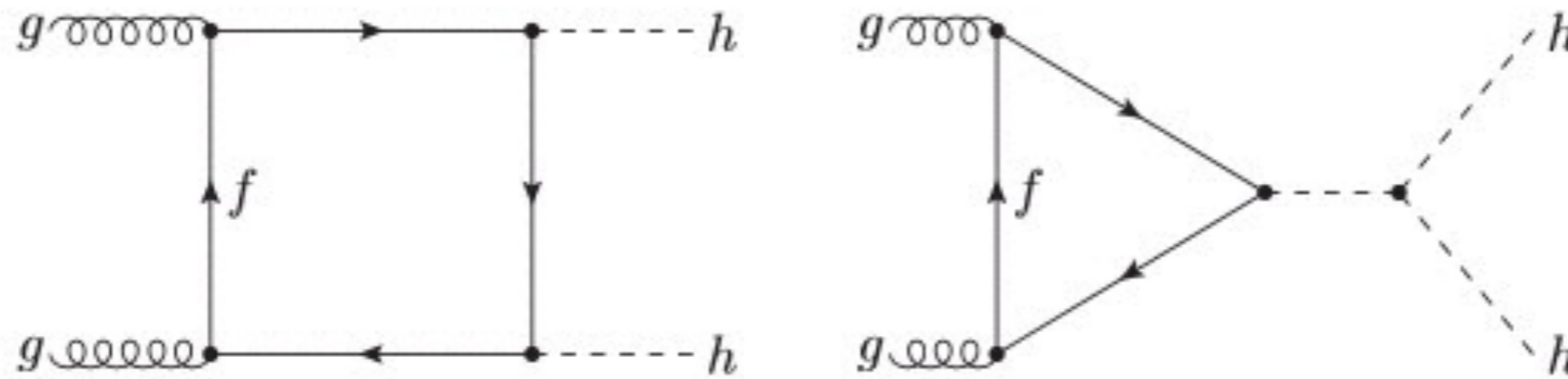
$$V(h) = \frac{m^2}{2} h^2 + \lambda h^4 + \frac{1}{\Lambda^2} h^6 + \dots$$

Measuring triple Higgs



$f = \text{top}, \dots$ Many possible final state. Very difficult channel.

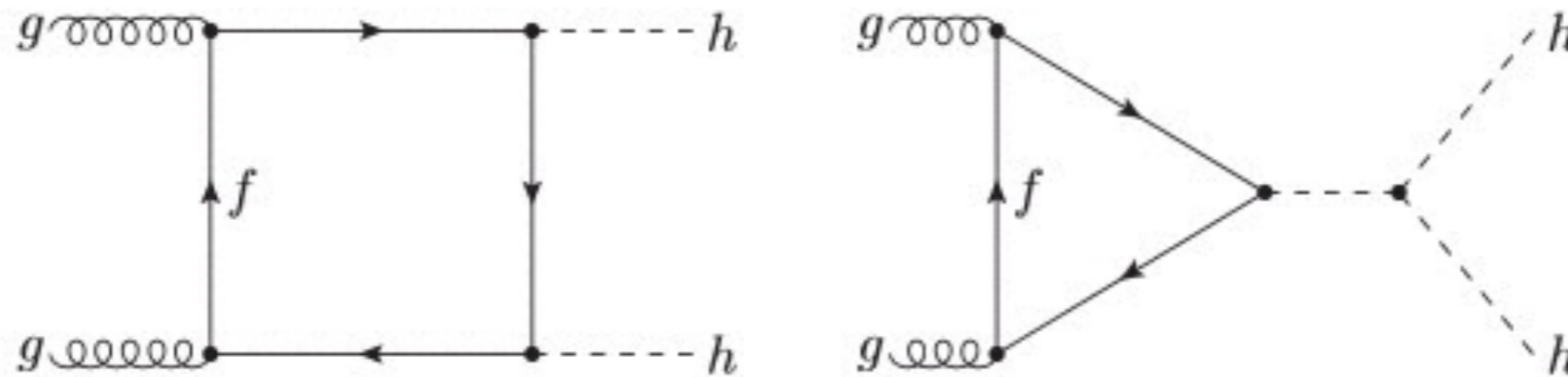
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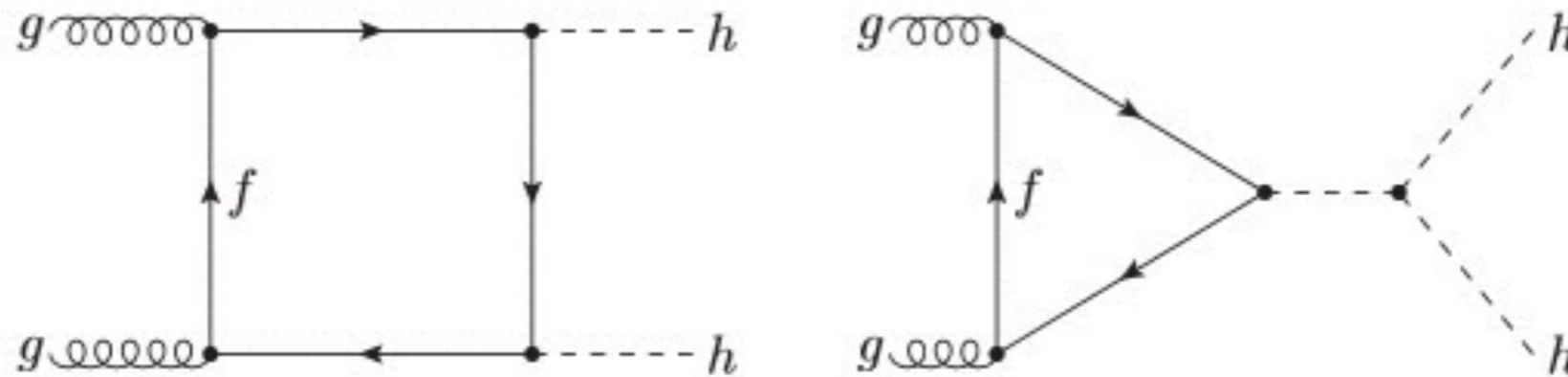
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Triple Higgs coupling at 100 TeV pp collider 30 ab^{-1}
Some preliminary studies, incomplete not fully realistic.

$$\frac{\lambda}{\lambda_{\text{SM}}} \in \begin{cases} [0.891, 1.115] & \text{no background syst.} \\ [0.882, 1.126] & 25\% hh, 25\% hh + \text{jet} \\ [0.881, 1.128] & 25\% hh, 50\% hh + \text{jet} \end{cases}$$

Barr, Dolan, Englert, de Lima, Spannowsky

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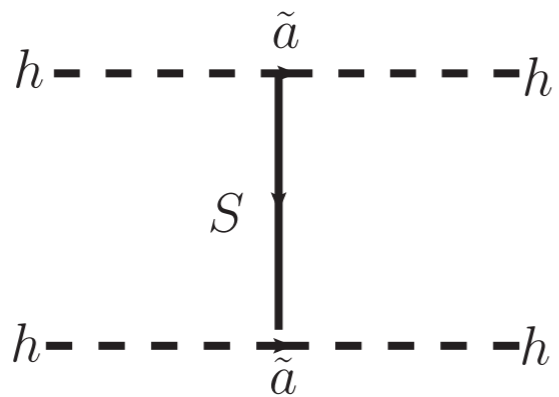
Barr, Dolan, Englert, de Lima, Spannowsky

ILC 500: 27%

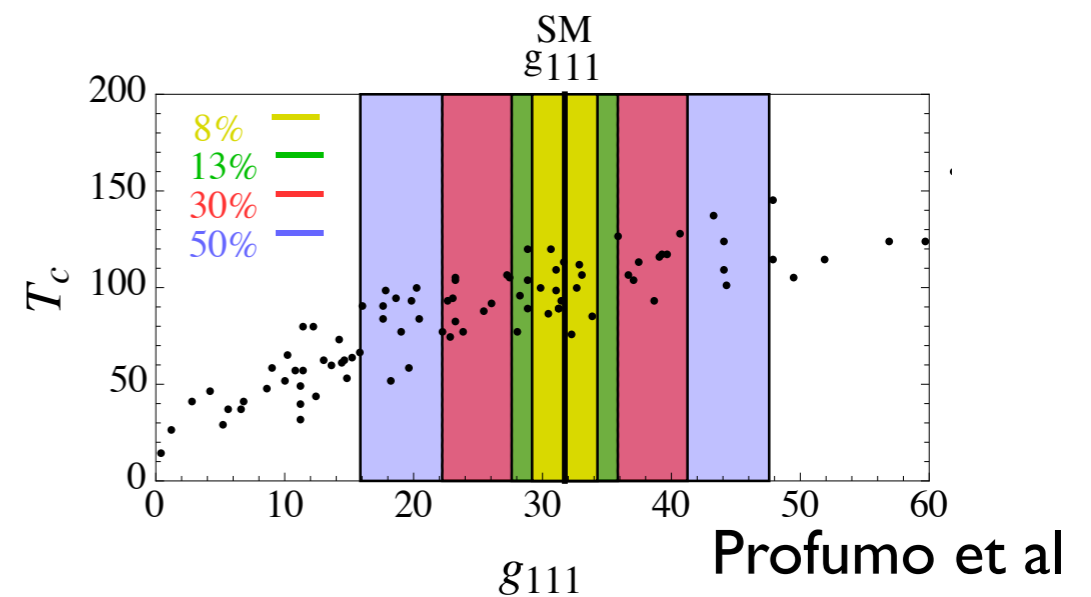
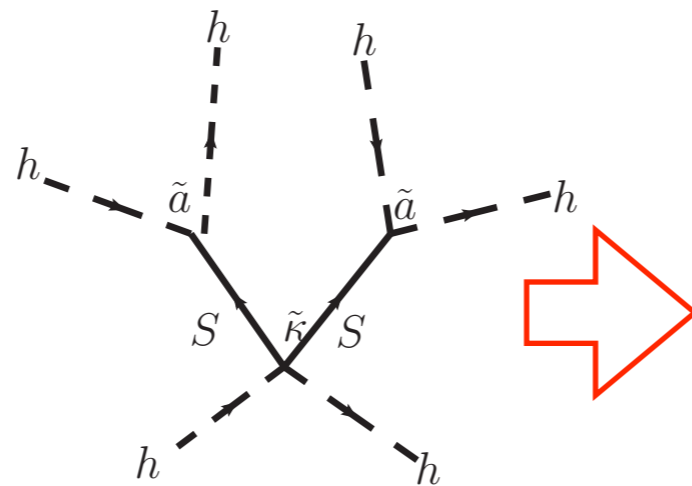
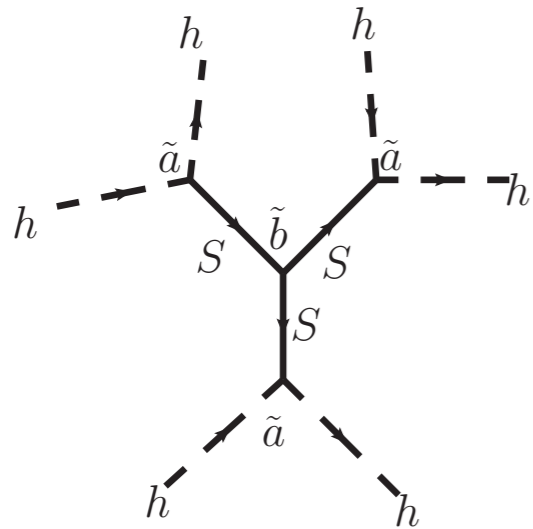
ILC ultimate, 1 TeV 5 ab^{-1} : 10%

Simple example: Generic singlet model

$$m^2 h^\dagger h + \tilde{\lambda} (h^\dagger h)^2 + m_S^2 S^2 + \tilde{a} S h^\dagger h + \tilde{b} S^3 + \tilde{\kappa} S^2 h^\dagger h + \tilde{h} S^4$$

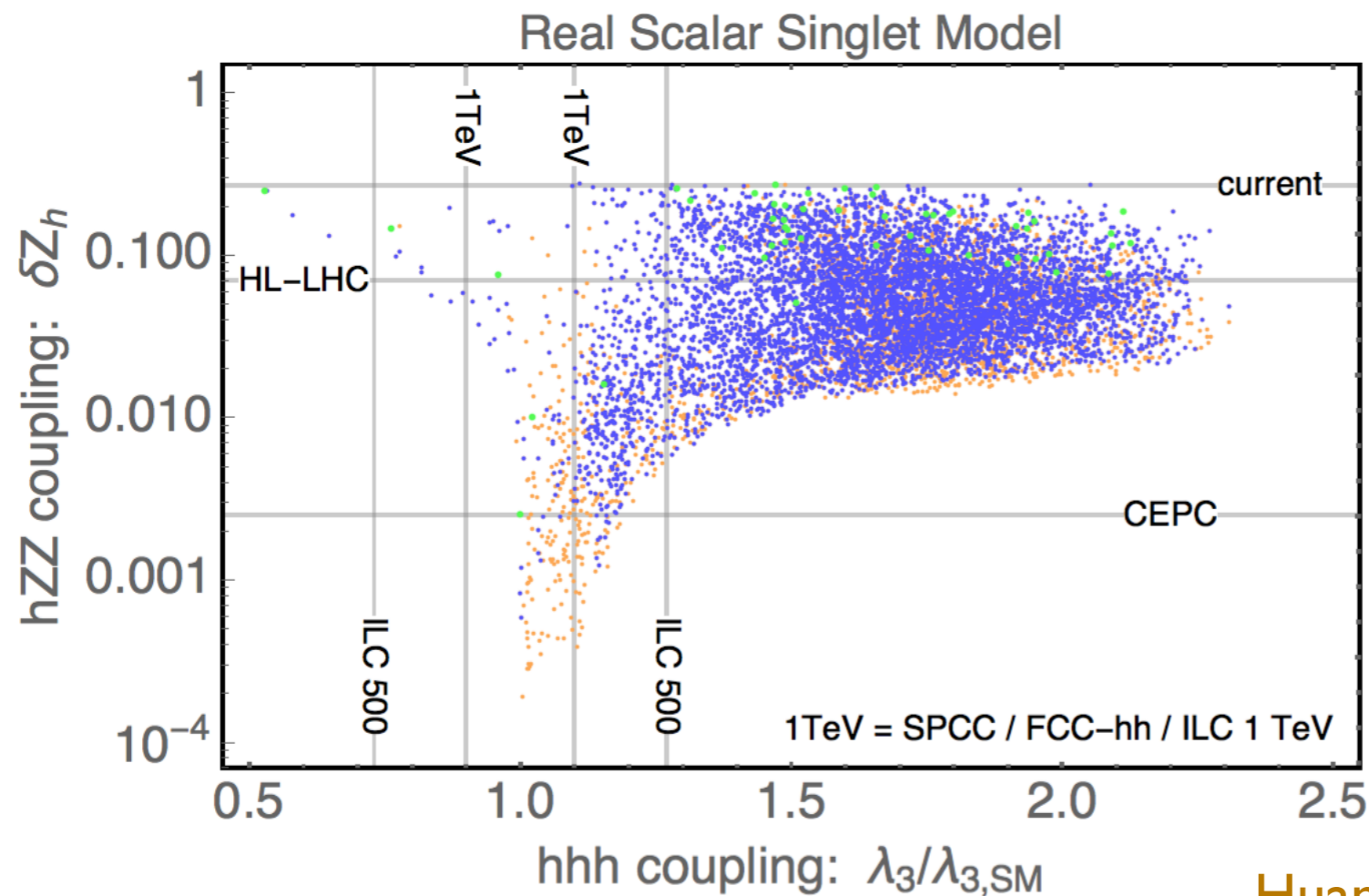


shift in h-Z coupling > %
Higgs factory important



O(1) deviation in triple Higgs coupling

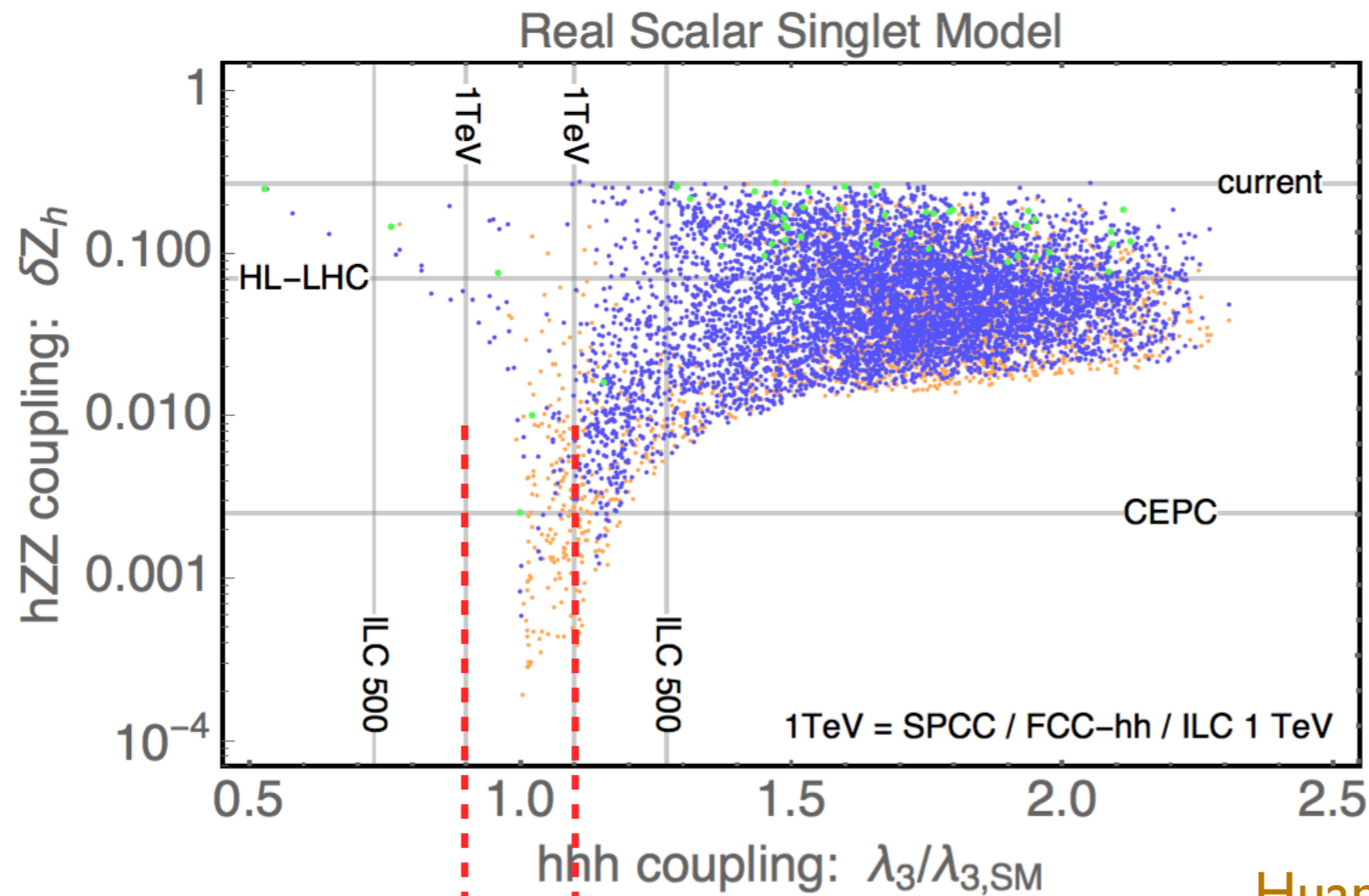
Also considering Higgs factories



Huang, Long, LTW, in progress

Orange = first order phase transition, $v(T_c)/T_c > 0$
Blue = “strongly” first order phase transition, $v(T_c)/T_c > 1.3$
Green = very strongly 1PT, could detect GWs at eLISA

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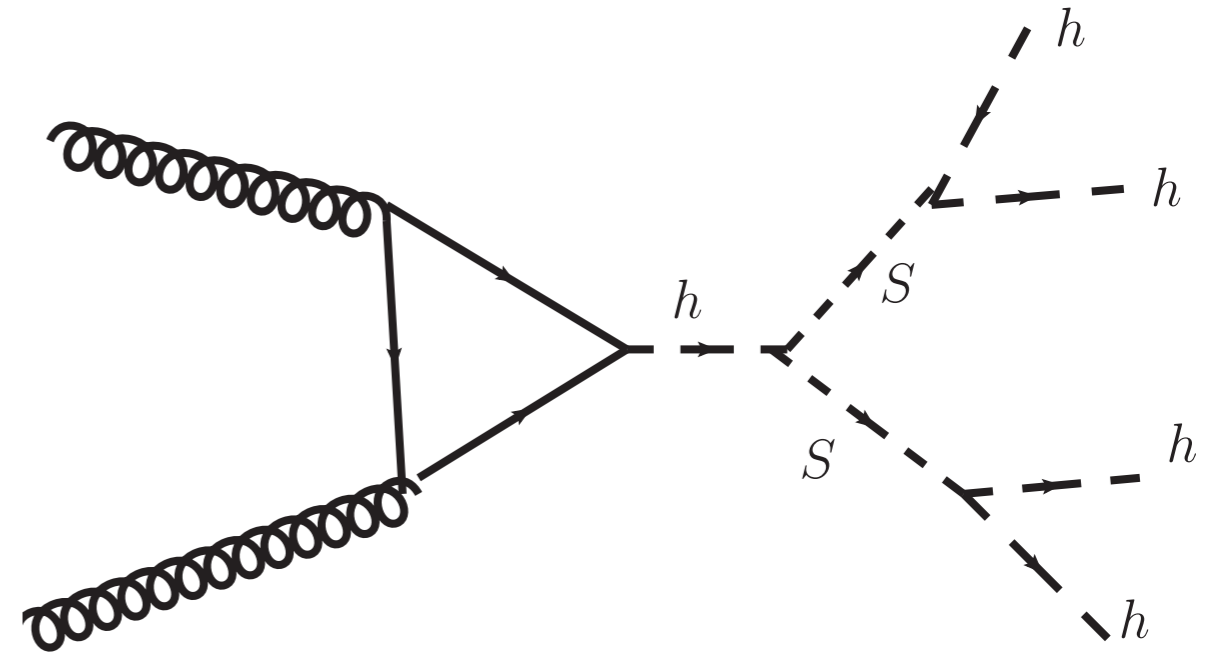
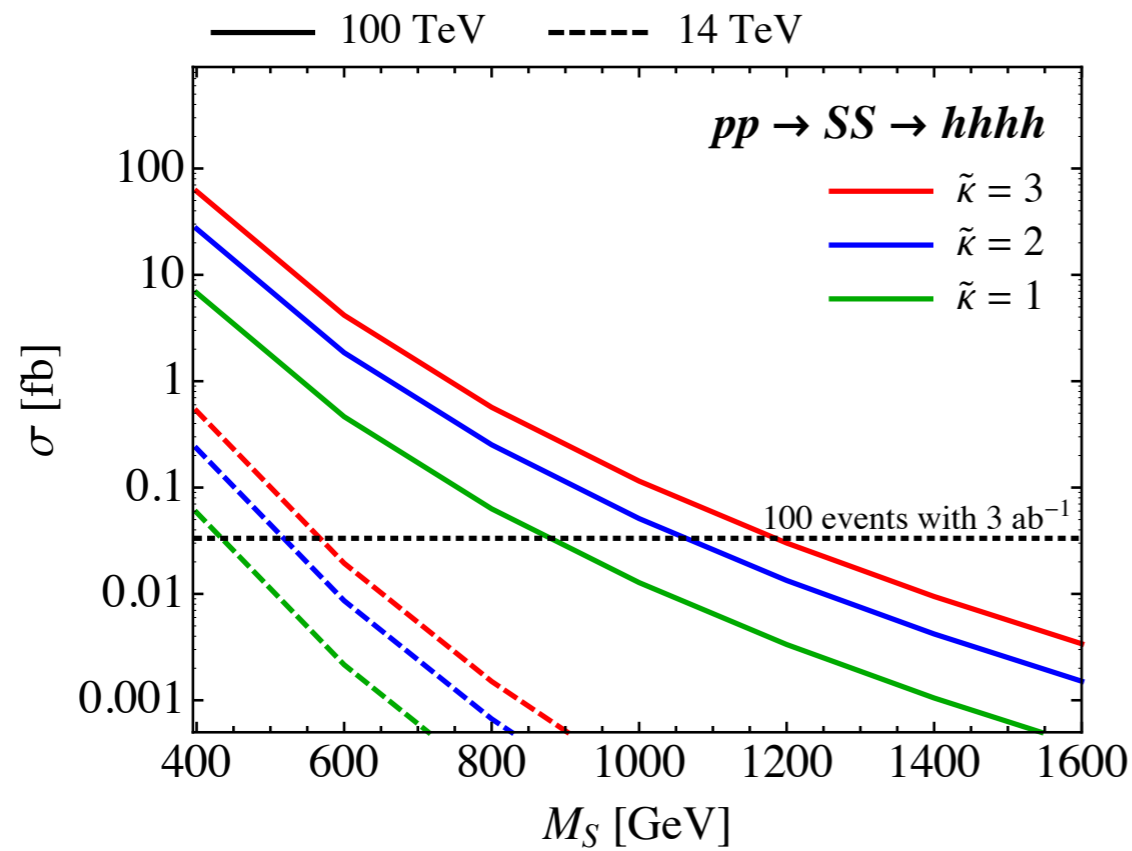


Huang, Long, LTW, in progress

100 TeV pp

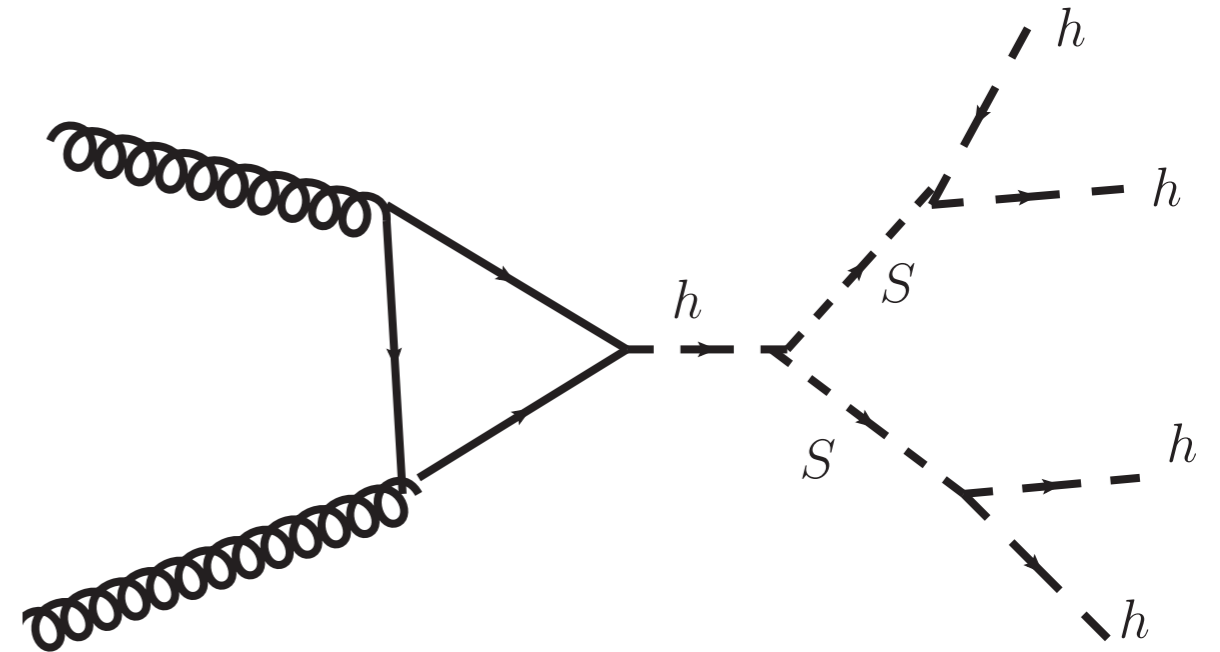
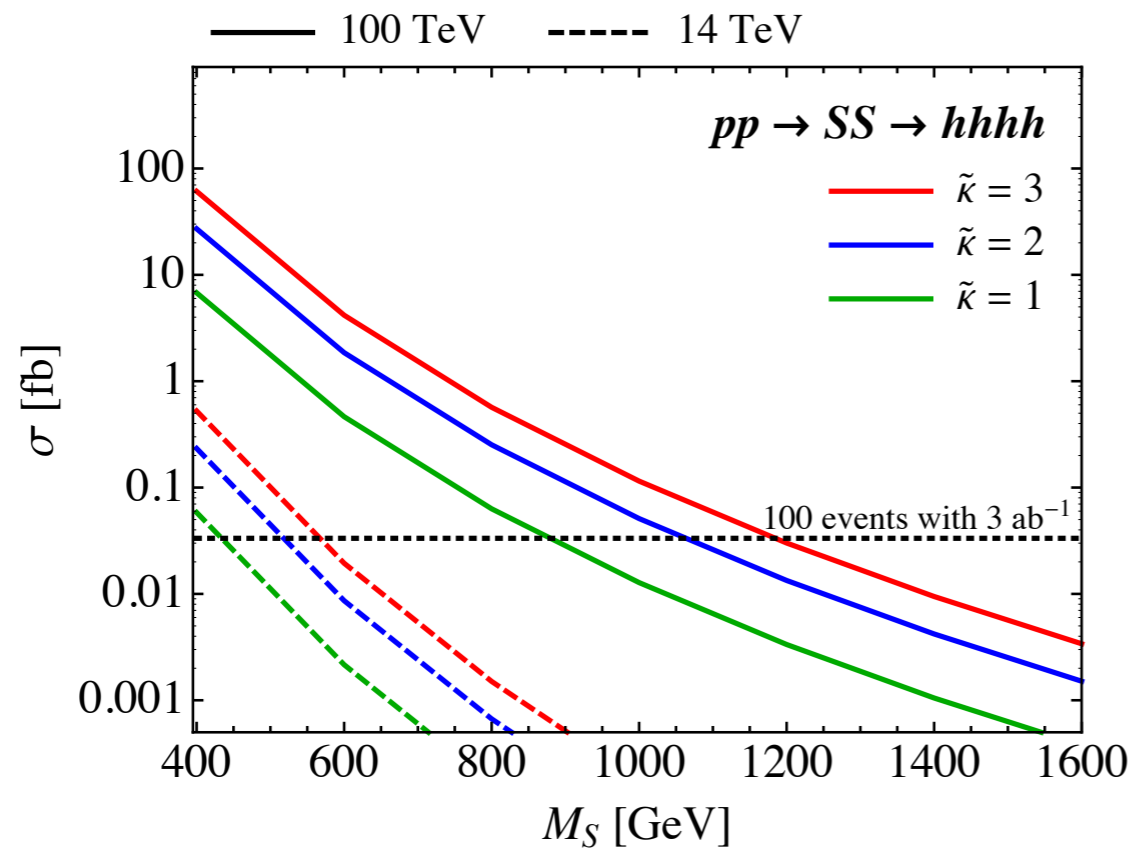
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Singlet search at 100 TeV



- 4 Higgs final state with decent rate.
- Good discovery potential.

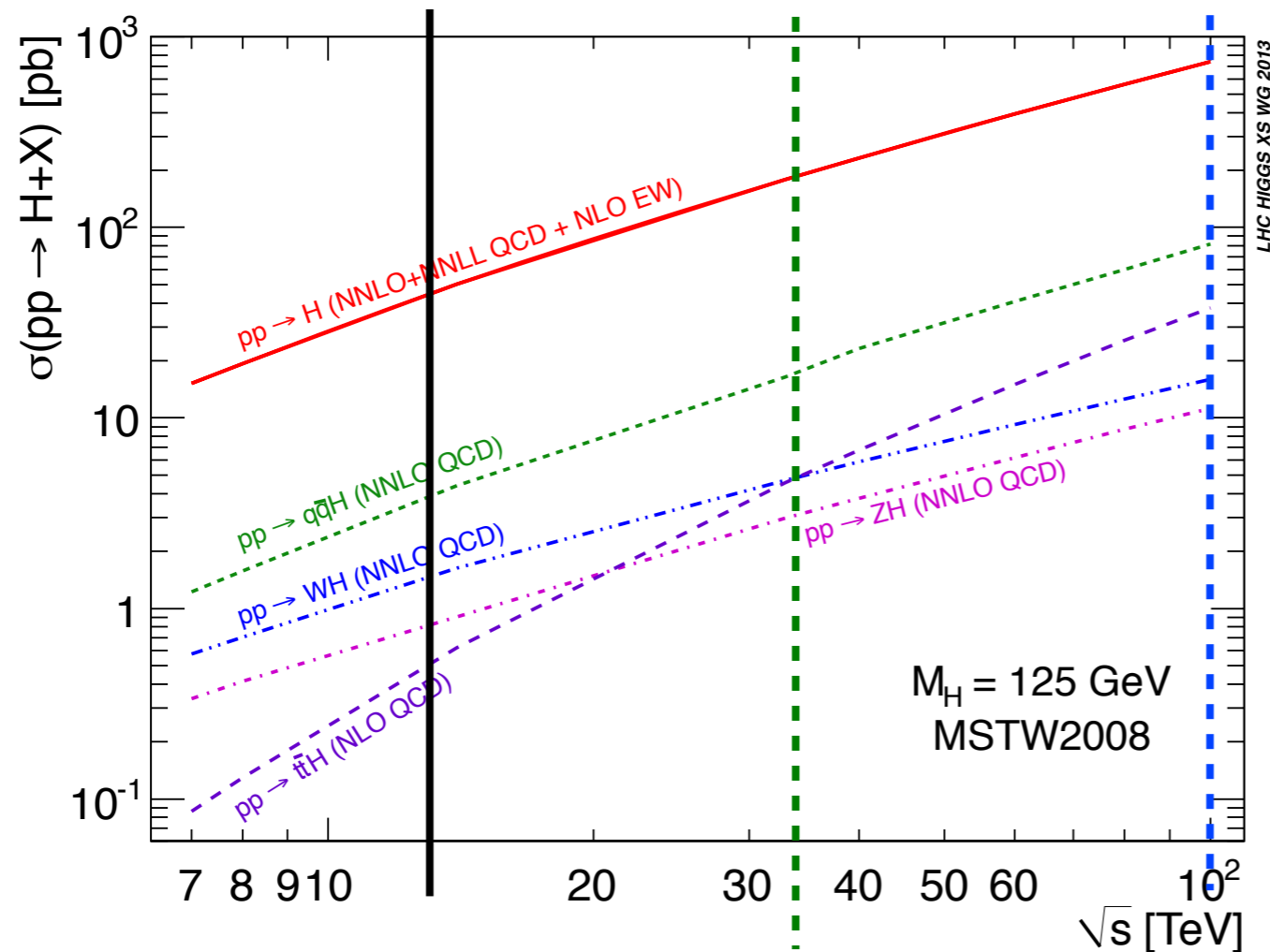
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Combination of Higgs factory and 100 TeV pp collider can go very long way in understanding EWSB

More Higgs physics at hadron collider



of Higgses in 3 ab^{-1}

100 TeV > 2 billion

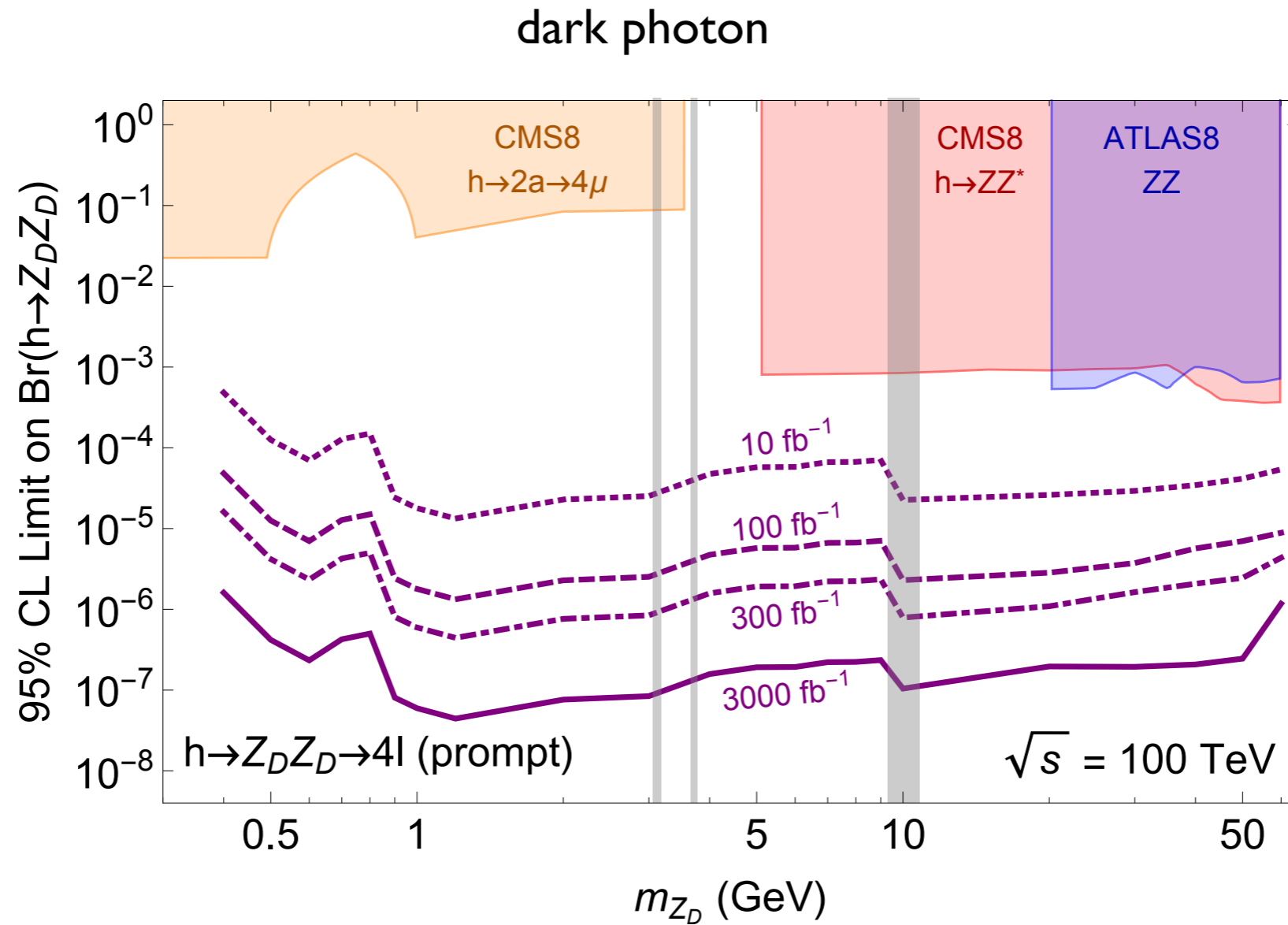
33 TeV > 500 million

14 TeV > 150 million

In comparison, $O(\text{million})$
Higgs at Higgs factories

Can look for very rare and distinct Higgs signal.

New physics Higgs rare decays

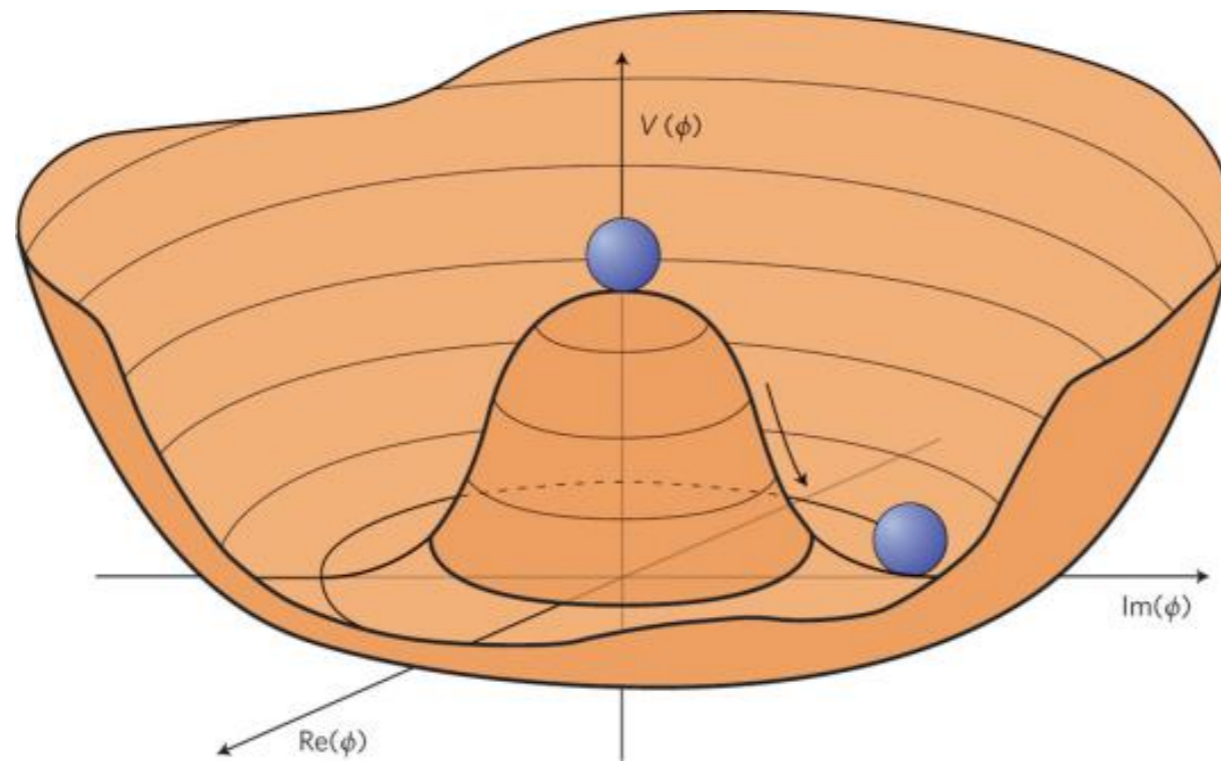


Curtin, Gori, Shelton

There are certainly more examples.

Naturalness

Explaining the Higgs potential.



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Explaining electroweak scale $O(100)$ GeV

Explaining EWSB: naturalness

..... M: The energy scale of new physics responsible for EWSB



Electroweak scale, 100 GeV.

$m_h, m_W \dots$

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What is M? Can it be very high, such as $M_{\text{Planck}} = 10^{19}$ GeV, ...?

If so, why is so different from 100 GeV?



Electroweak scale, 100 GeV.

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Naturalness of electroweak symmetry breaking

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TeV new physics.
Naturalness motivated



Electroweak scale, 100 GeV.
 $m_h, m_W \dots$

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- No large cancellation $\Rightarrow m_h^2$ (physical) $\approx (M_{1,2})^2$
 - ▶ **$M \approx 100$ GeV – TeV, new physics at TeV scale!**

Is fine-tuning ok?

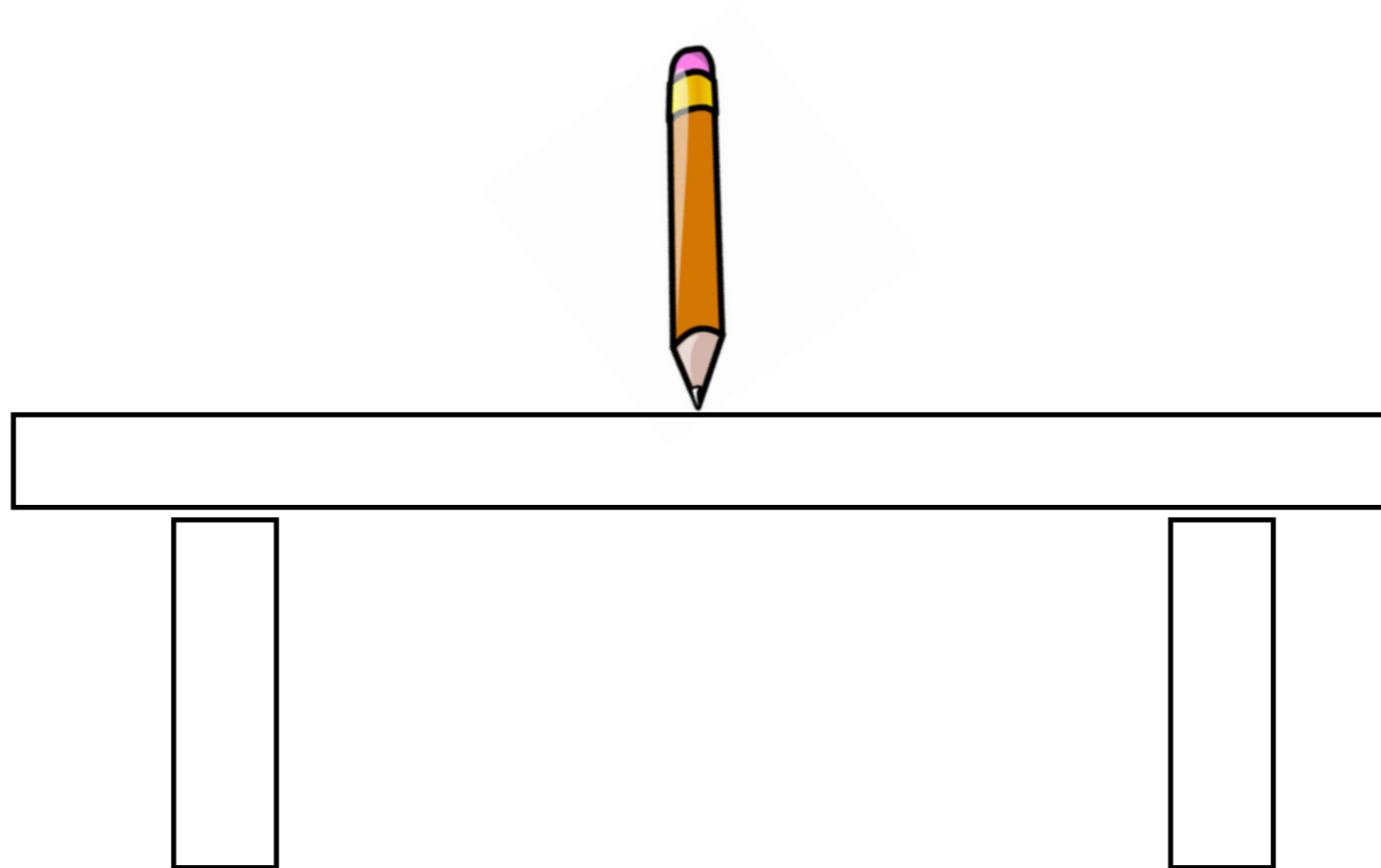
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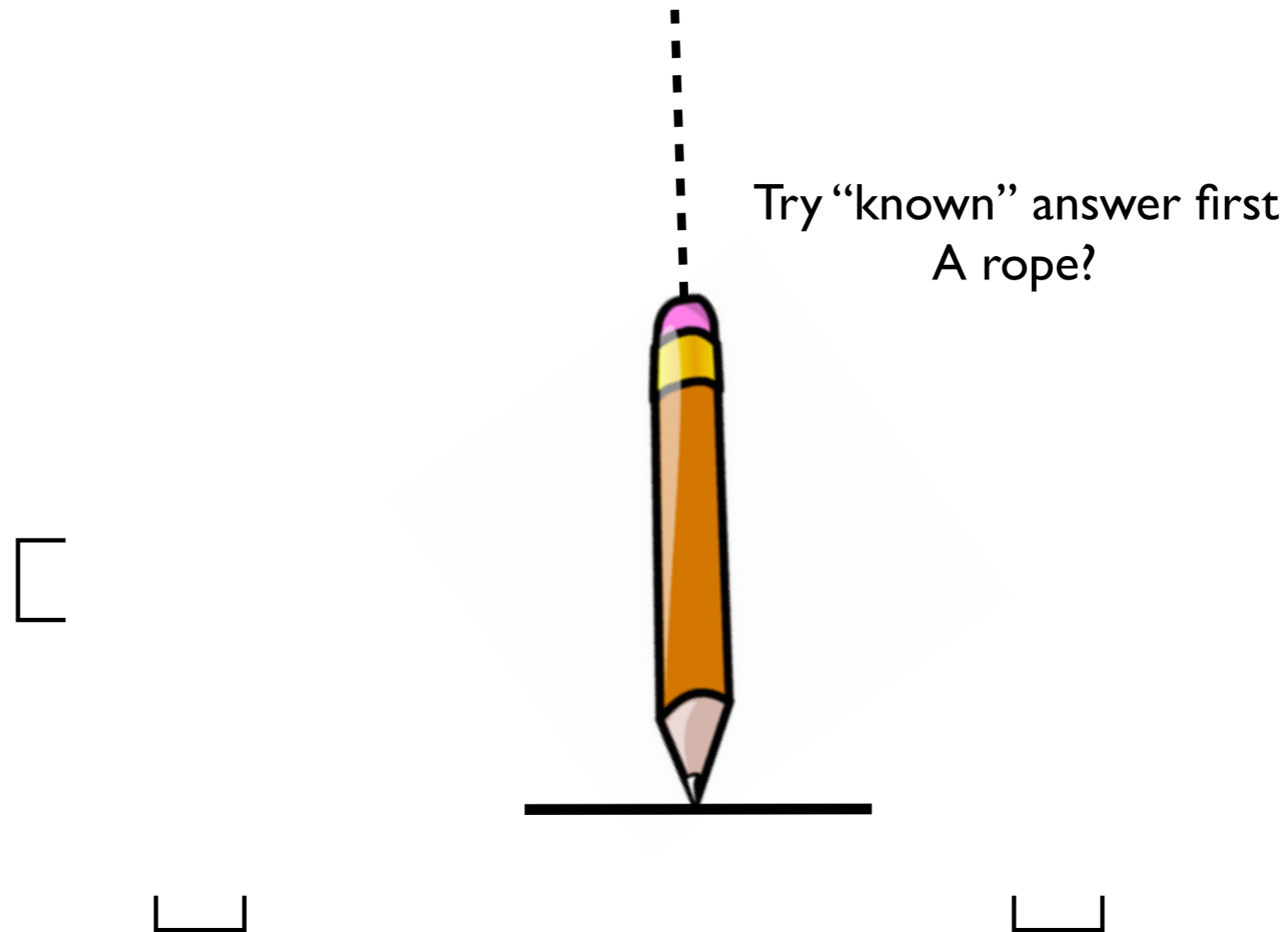


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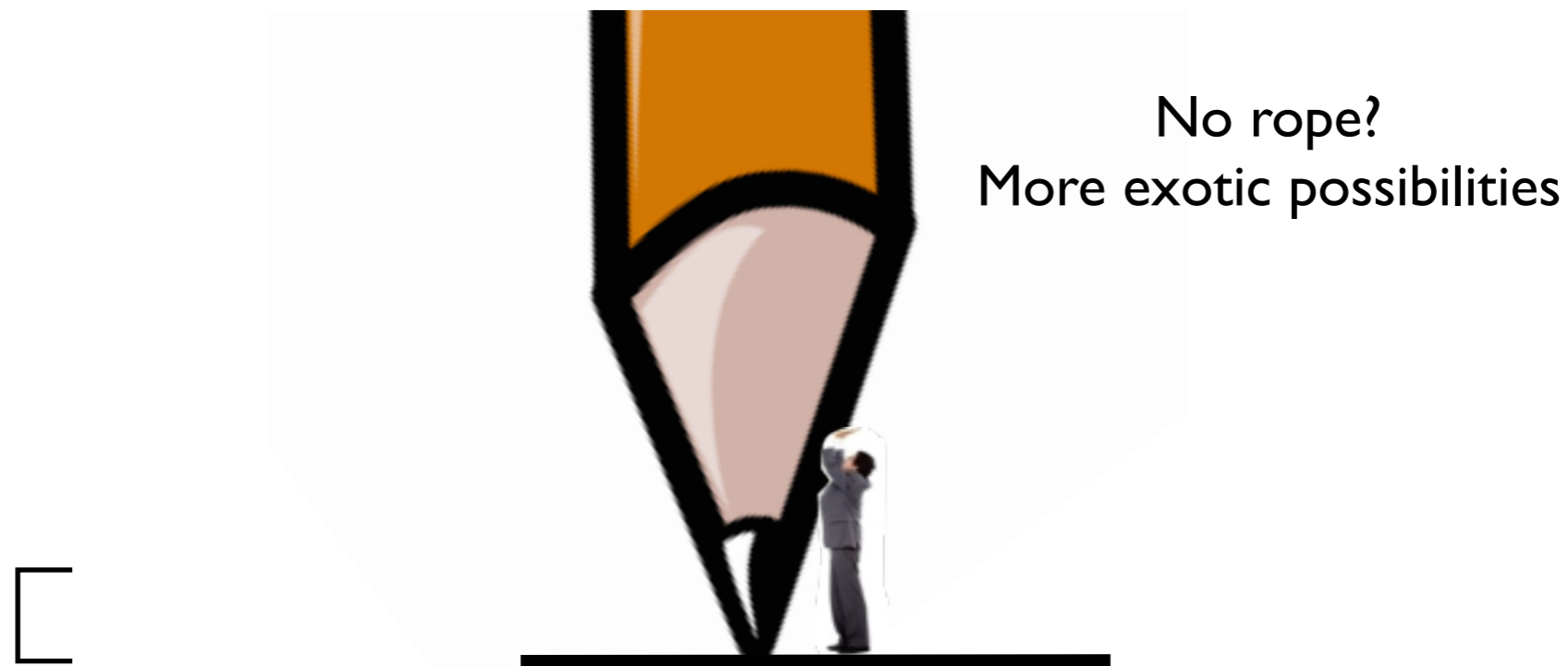


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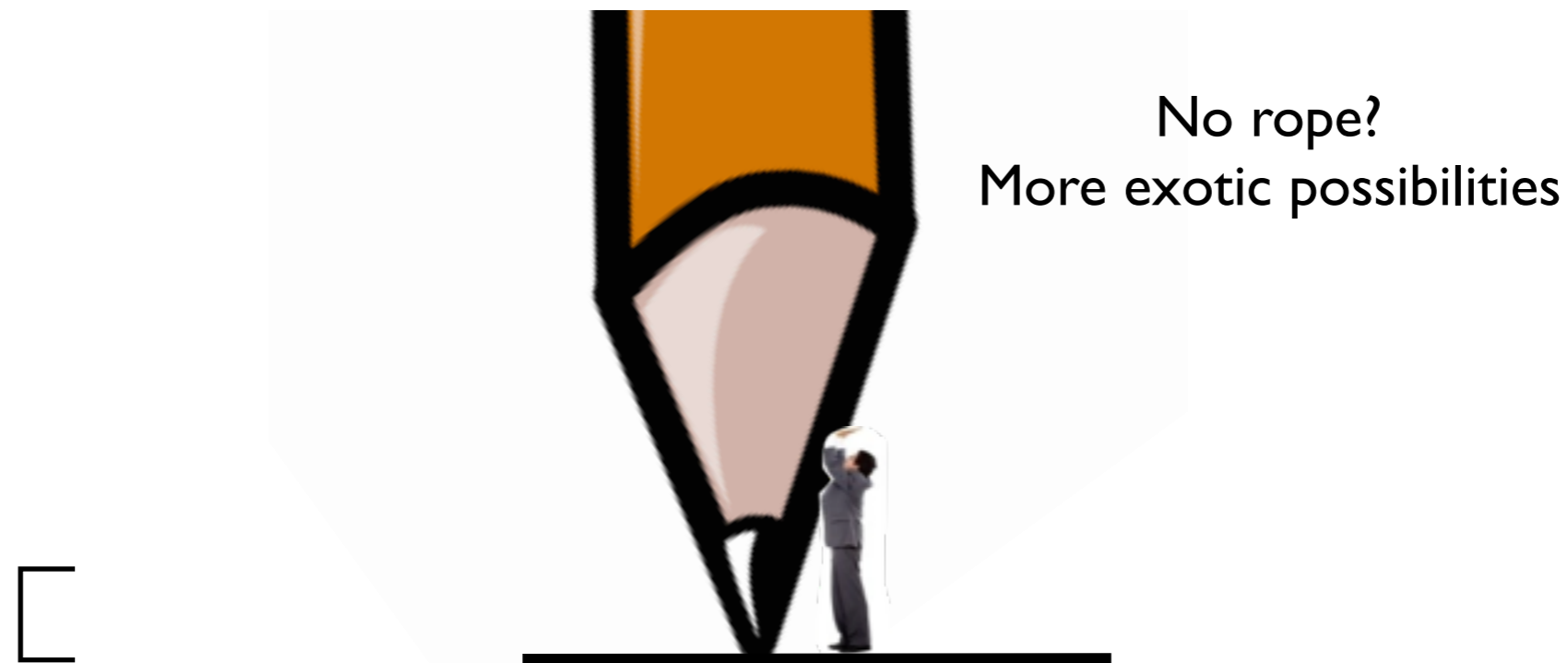


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Similarly, we have been searching for an explanation for the fine-tuning of Higgs mass $O(10^{-32})$

Another fine-tuning problem

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- Time to completely give up on this “conventional” naturalness? No!

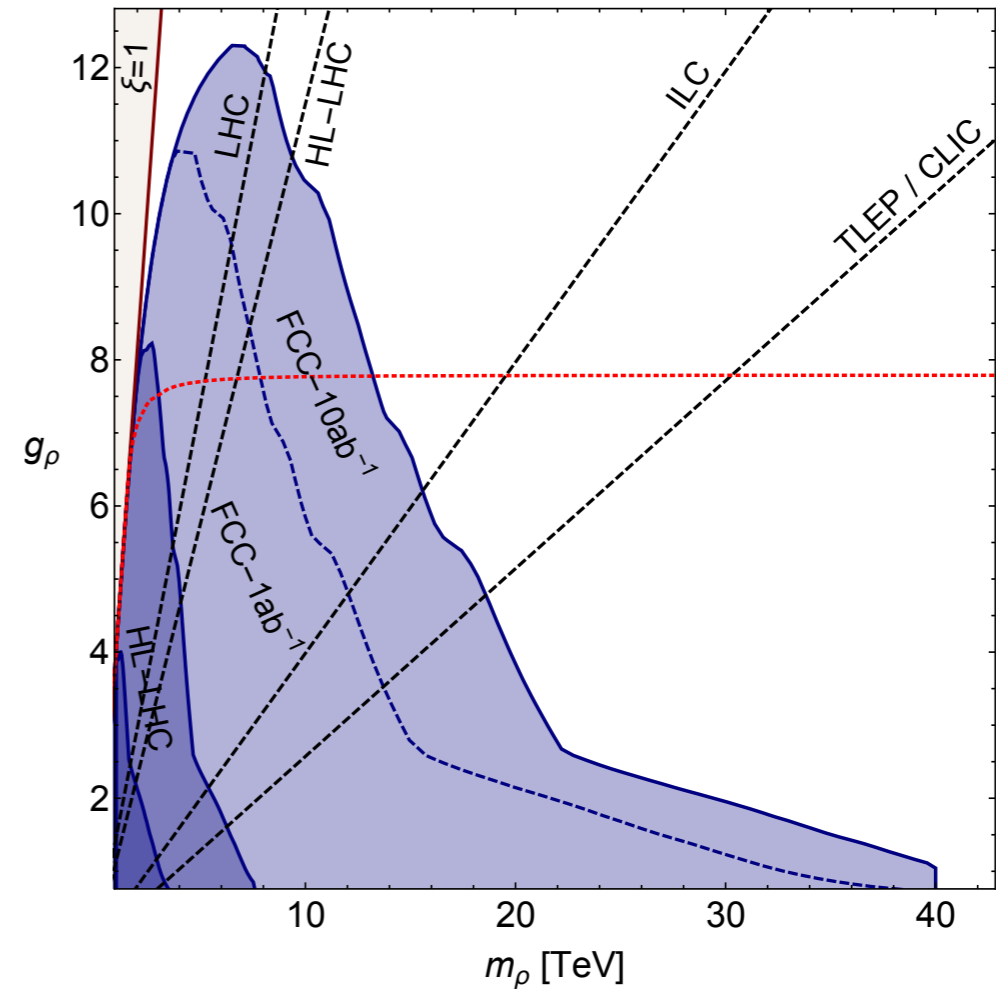
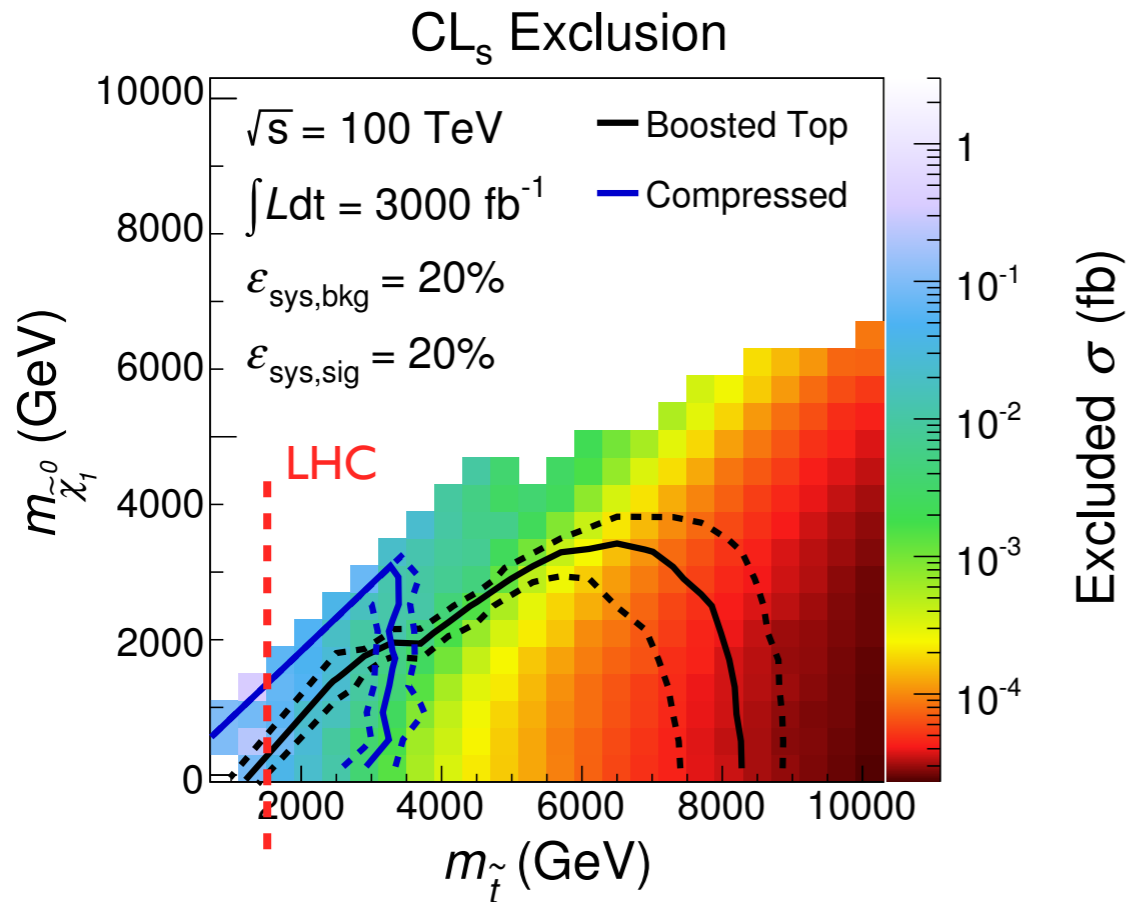
“Alternatives”

- Connection with cosmological evolution?
- Unique vacuum vs landscape
 - ▶ Dynamics vs selection.
- Dramatic new phenomena in quantum field theory
 - ▶ UV-IR connection. etc.
- Dramatic paradigm shifts. Very interesting.
- Too important to completely give up on the conventional notion of naturalness after the LHC.

Test naturalness at 100 TeV collider

Cohen et. al., 2014

Pappadopulo, Thamm, Torre, Wulzer, 2014

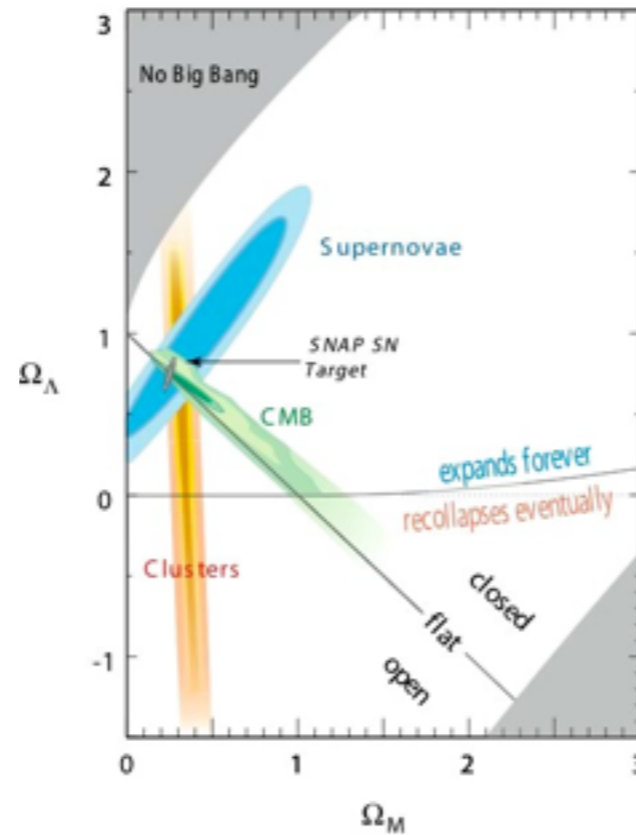
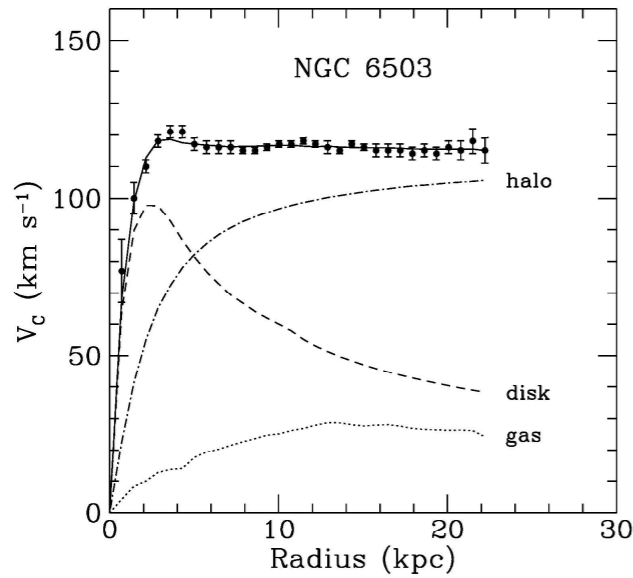


- tune proportional to $(M_{\text{new physics}})^2$.
 - ▶ Much better test than LHC, by orders of magnitude!
 - ▶ Potential for discovery (would be a victory for naturalness).

Testing WIMP Dark Matter

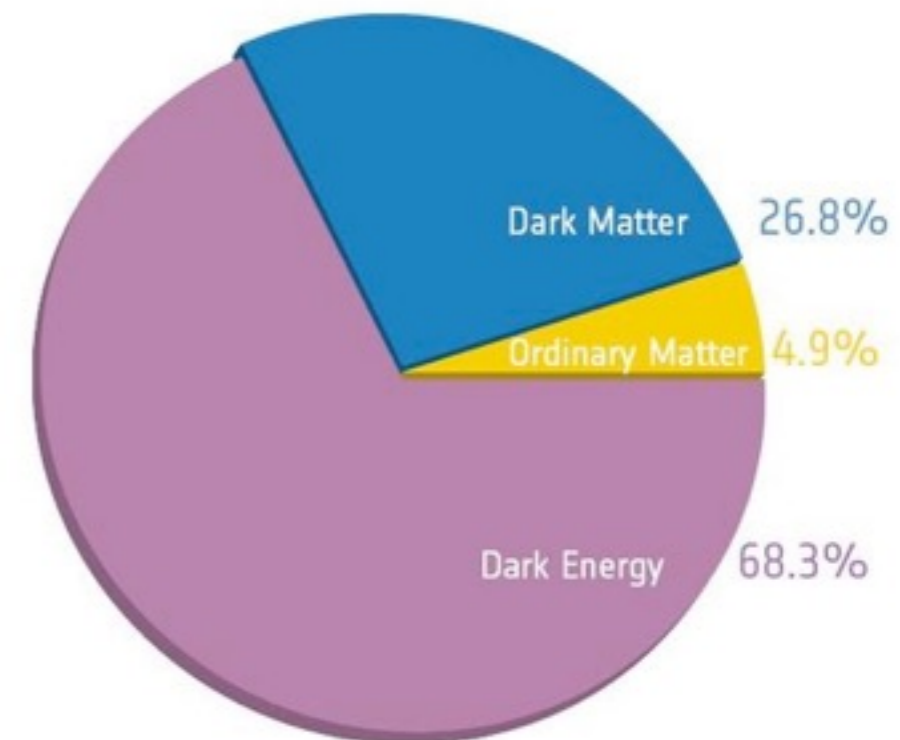
$$M_{\text{WIMP}} \leq 1.8 \text{ TeV} \left(\frac{g^2}{0.3} \right)$$

Dark matter

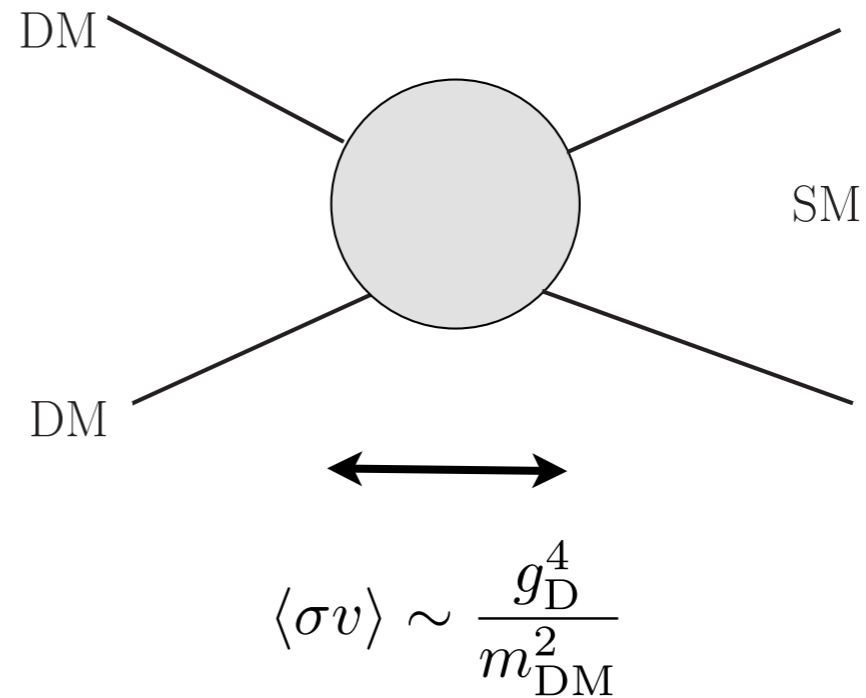
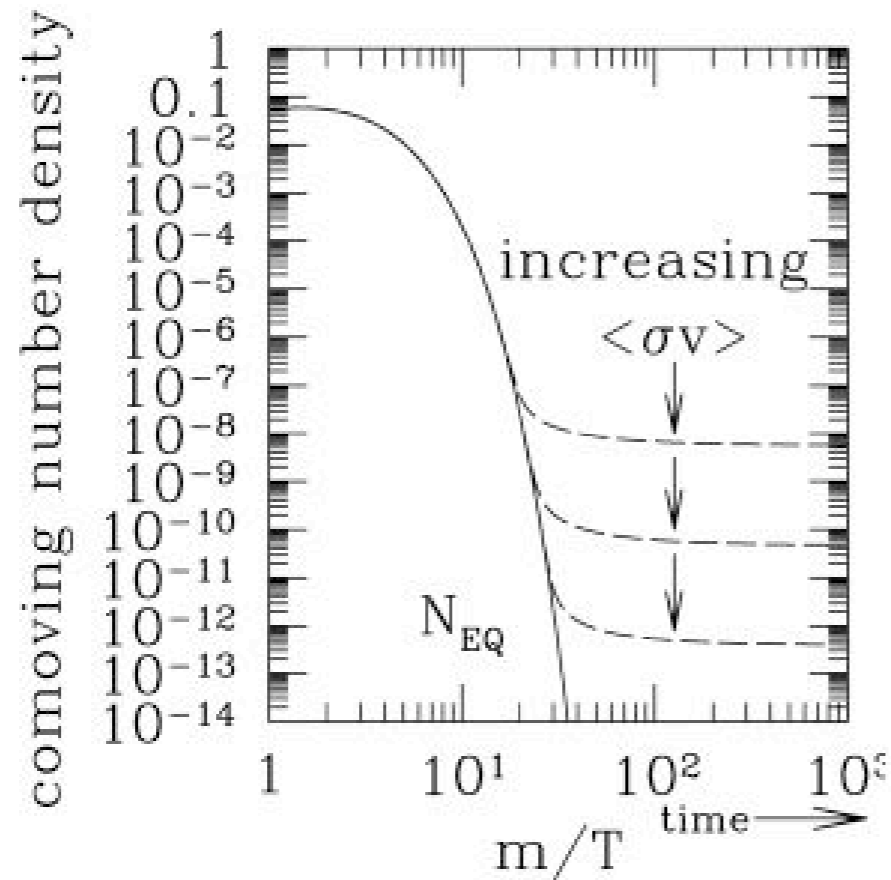


Not required by theory. It is there.
Only seen its gravitational interaction.

We have to understand them better.
Collider search is a key approach.

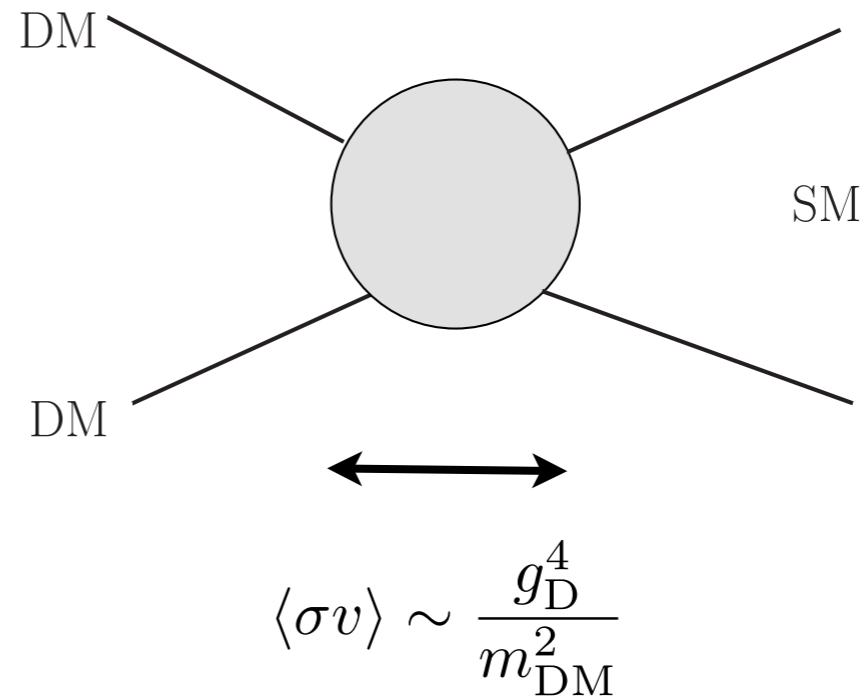
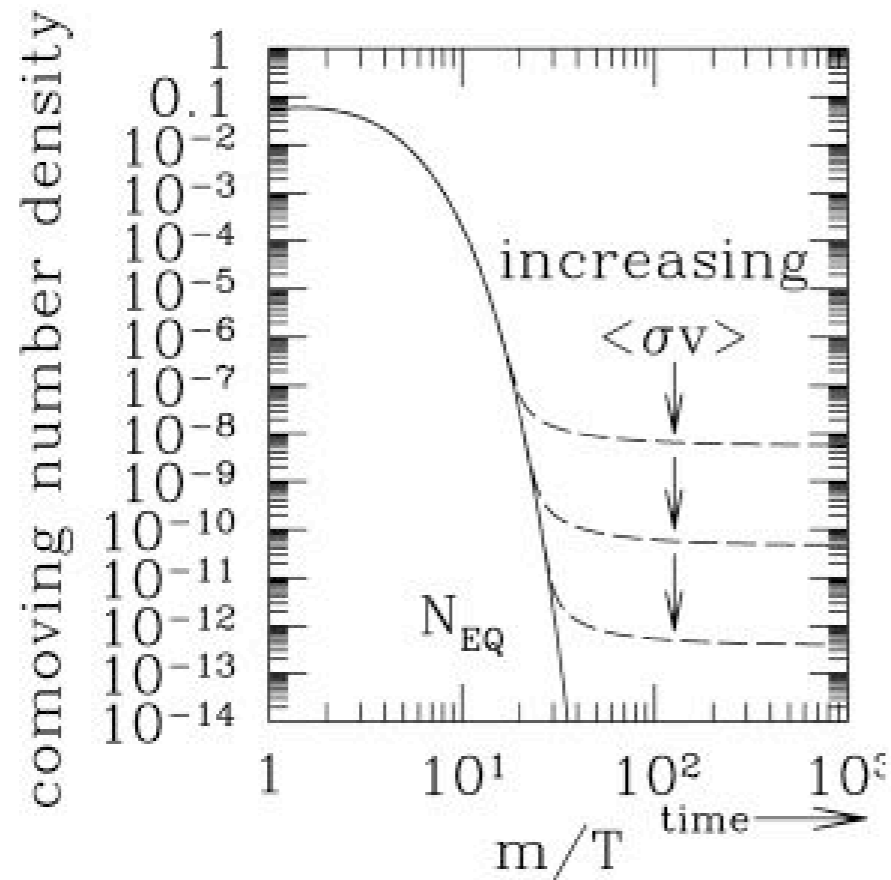


WIMP scenario.



- Thermal equilibrium in the early universe.
- If $g_D \sim 0.1$ $M_D \sim 10$ s GeV - TeV
 - ▶ We get the right relic abundance of dark matter.
- Major hint for weak scale new physics!

WIMP mass

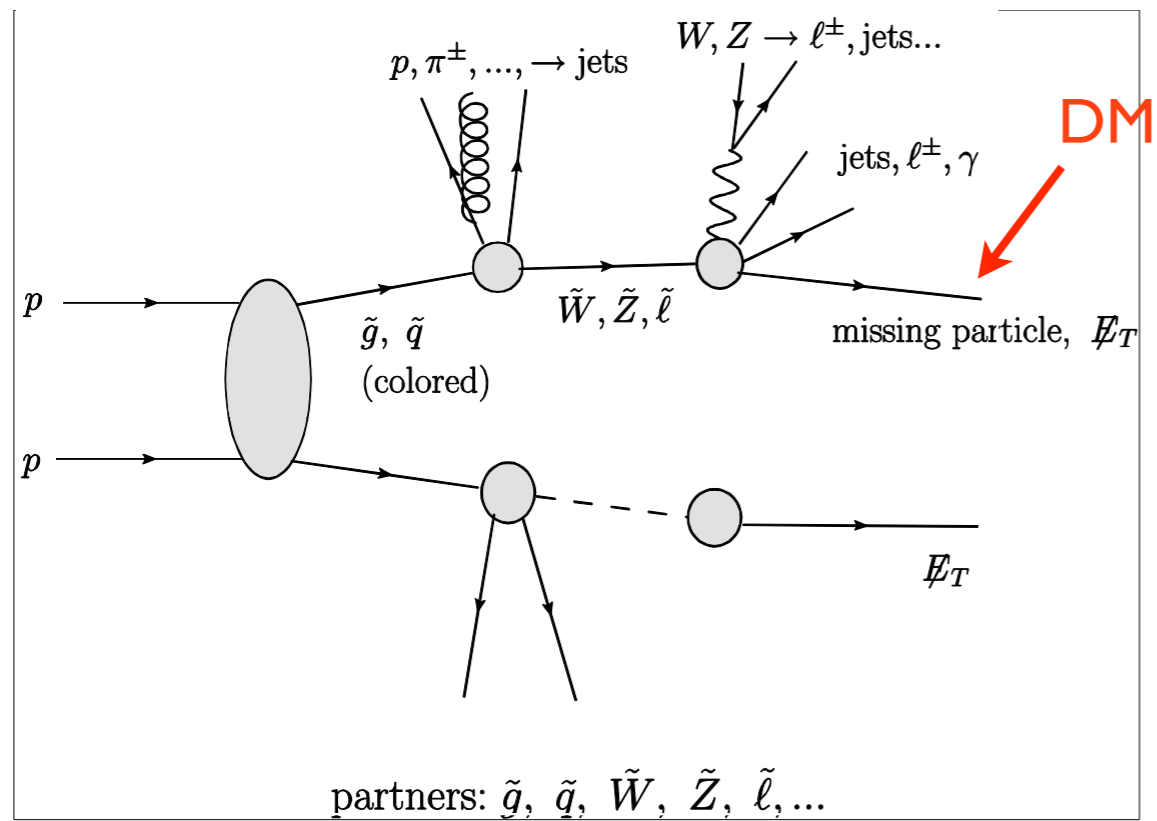


- More precisely, to get the correct relic abundance

$$M_{WIMP} \leq 1.8 \text{ TeV} \left(\frac{g^2}{0.3} \right)$$

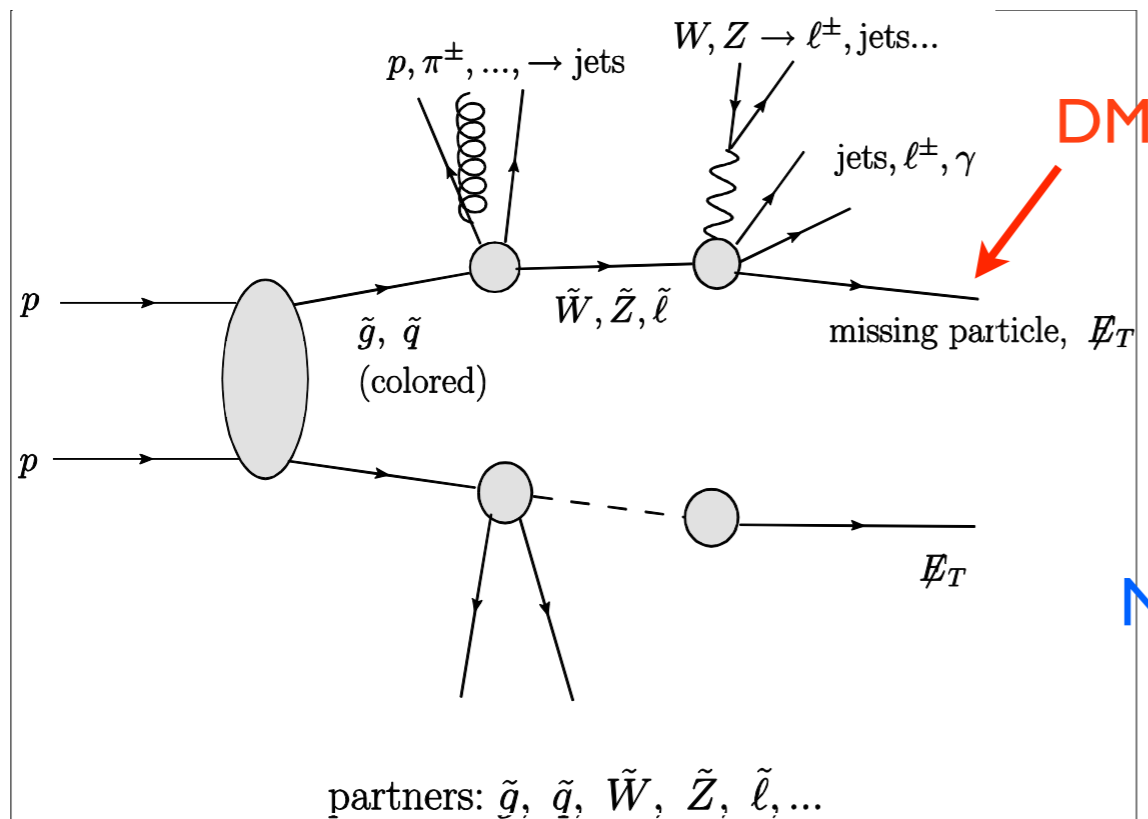
TeV-ish in simplest models

The story I grew up with

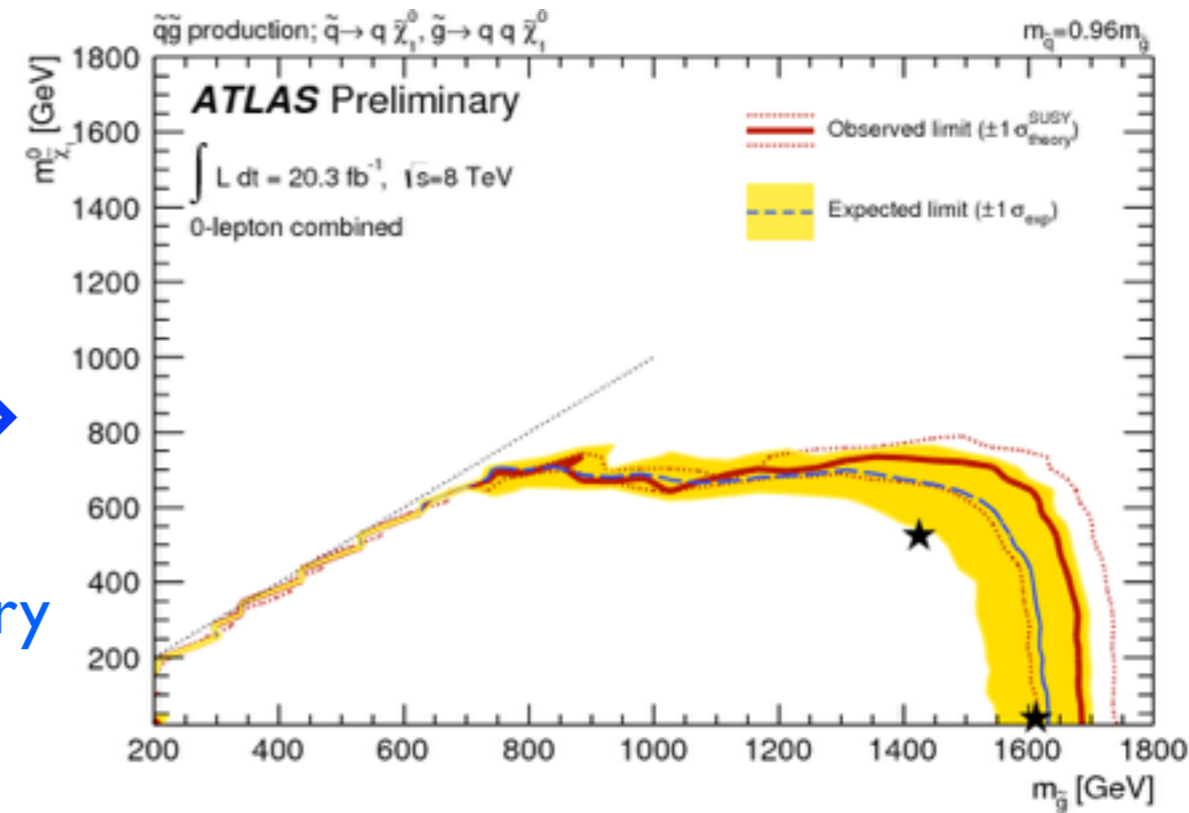


- WIMP is part of a complete model at weak scale.
- It's produced as part of the NP signal, shows up as missing energy.
 - ▶ Dominated by colored NP particle production: eg. gluino.
- The reach is correlated with the rest of the particle spectrum.

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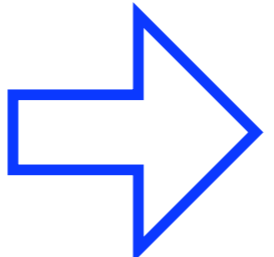
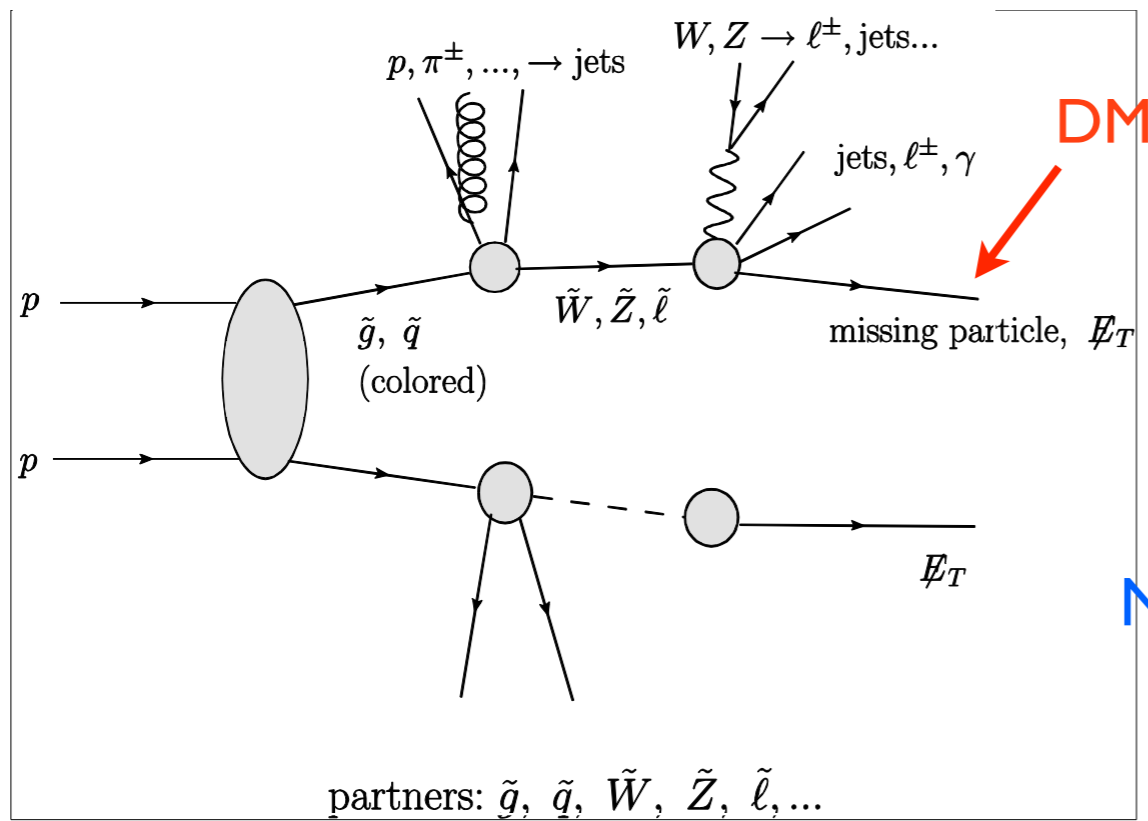


➡
No discovery yet

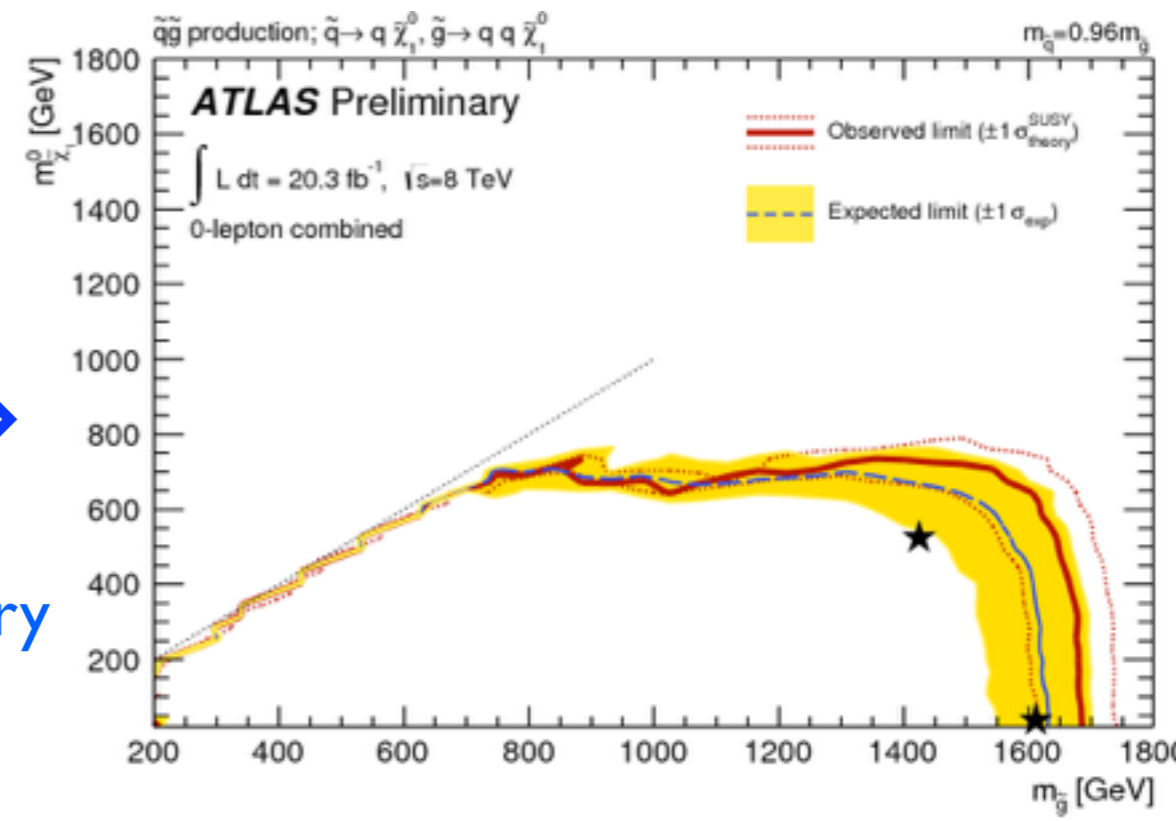


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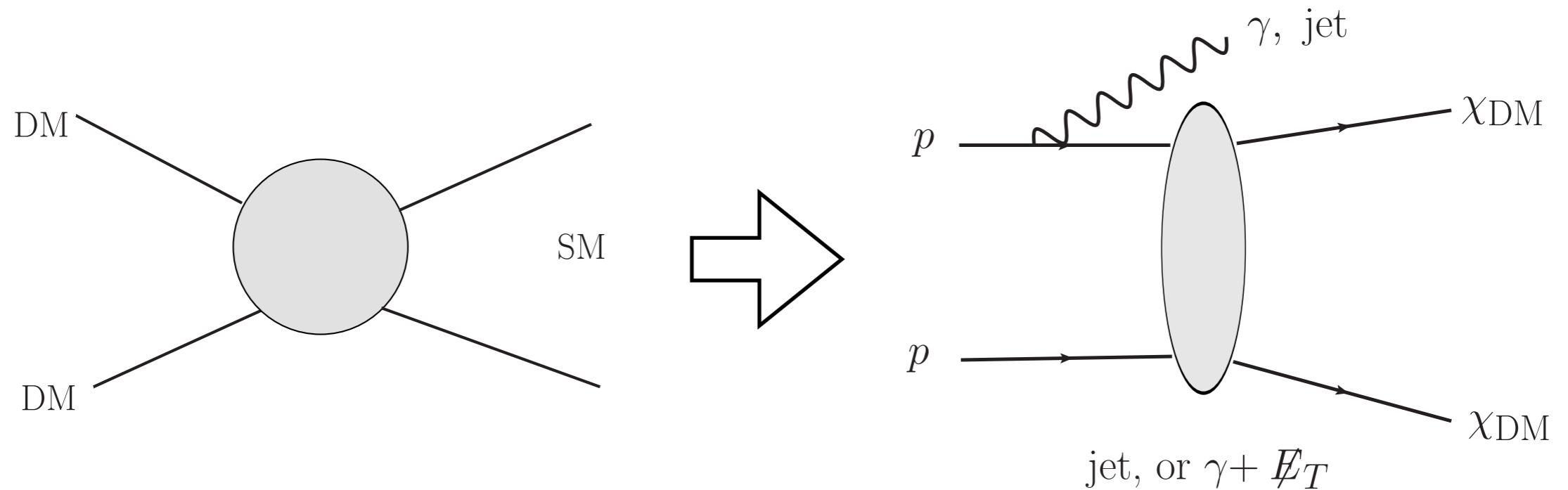
No discovery yet



Of course, still plausible at the LHC, will keep looking.
Higher energy \Rightarrow higher reach

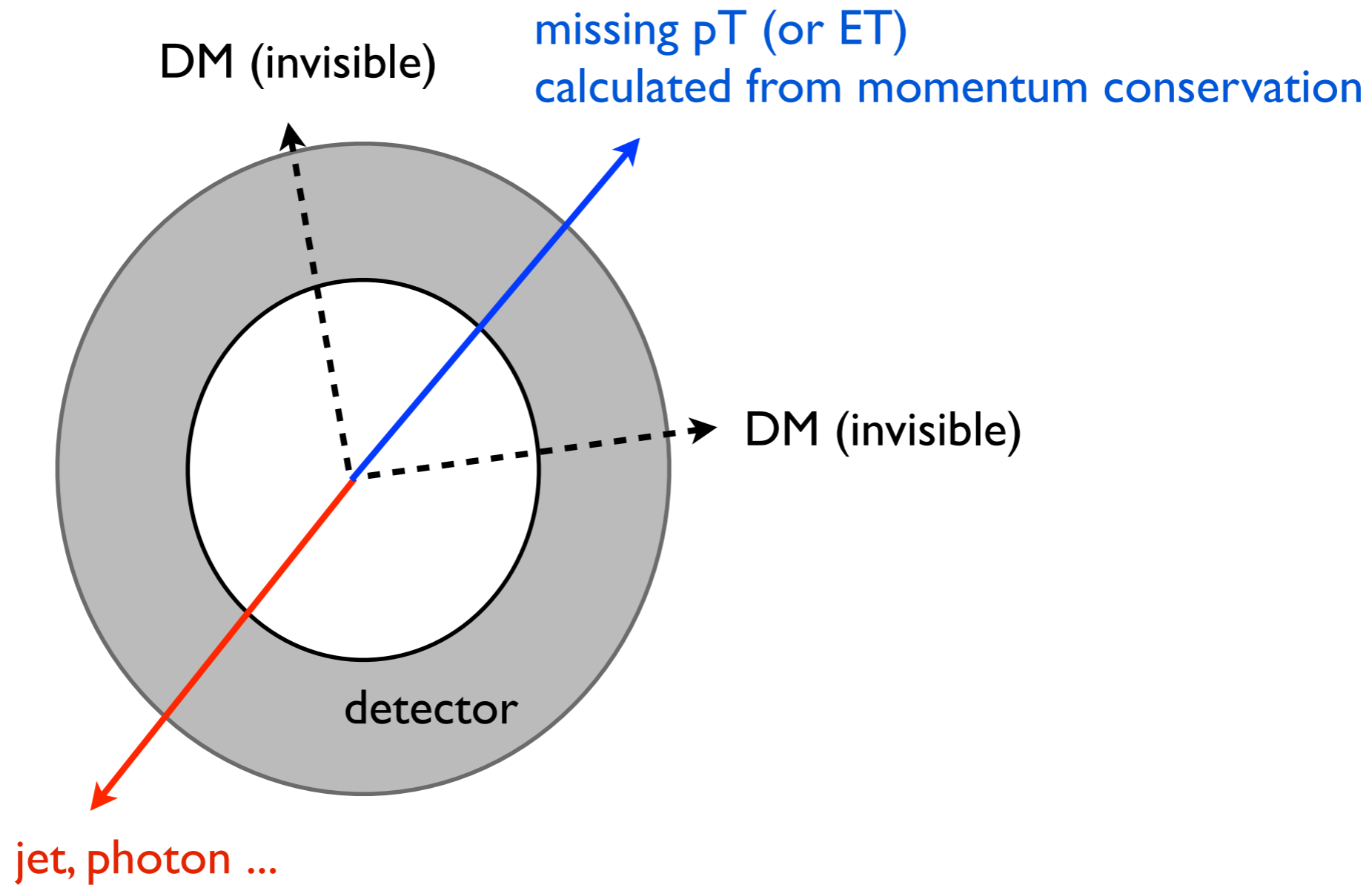
Basic channel

- pair production + additional radiation.

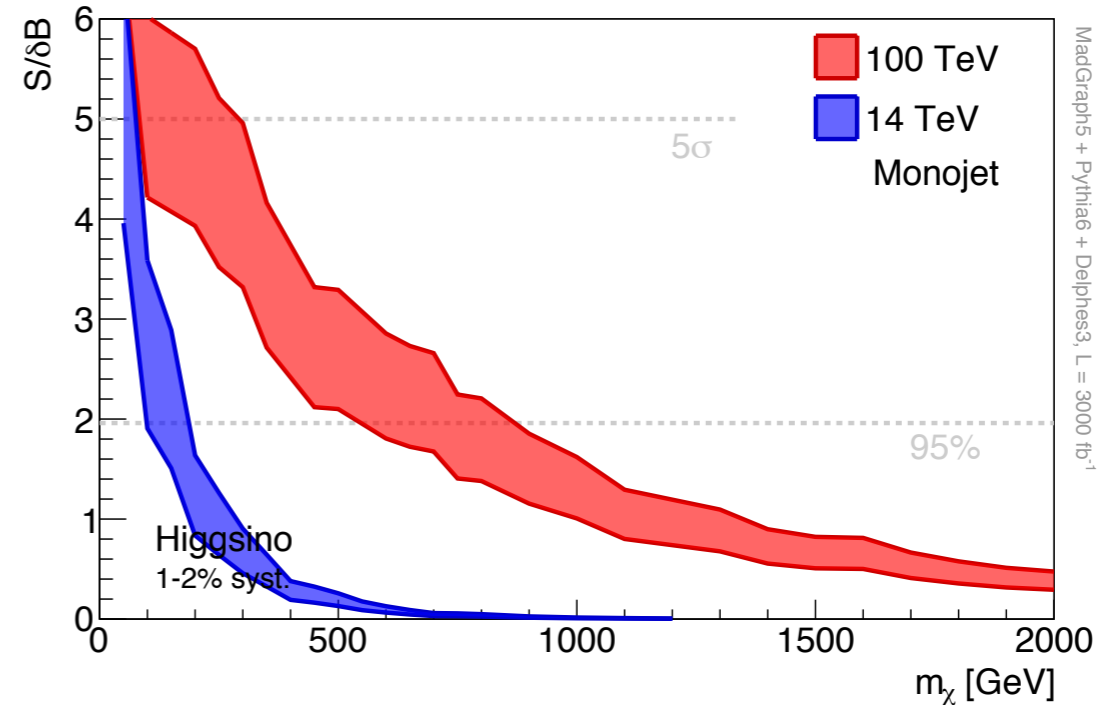
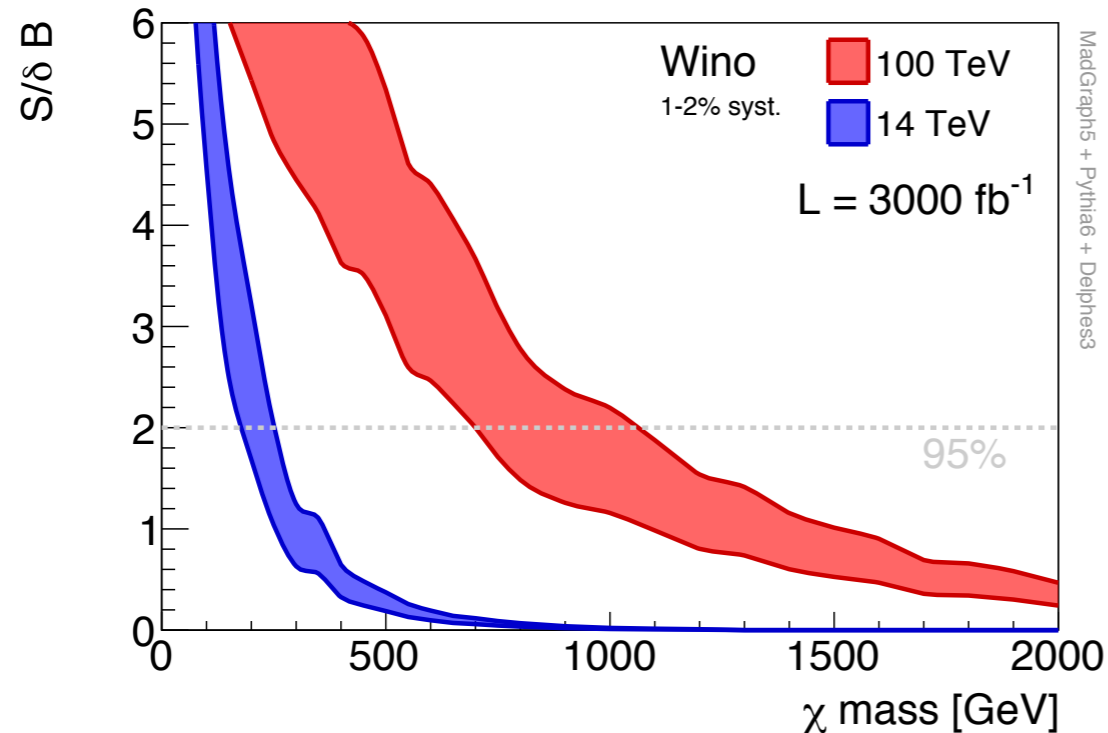


- Mono-jet, mono-photon, mono-...
- Have become "Standard" LHC searches.

Mono-X signature



Dark matter (mono-jet)



$$M_{\text{WIMP}} \leq 1.8 \text{ TeV} \left(\frac{g^2}{0.3} \right)$$

- LHC only coverage very limited. Rate, systematics...
- 100 TeV pp collider can probe the “bulk” of WIMP parameter space.

Very degenerate, disappearing track.

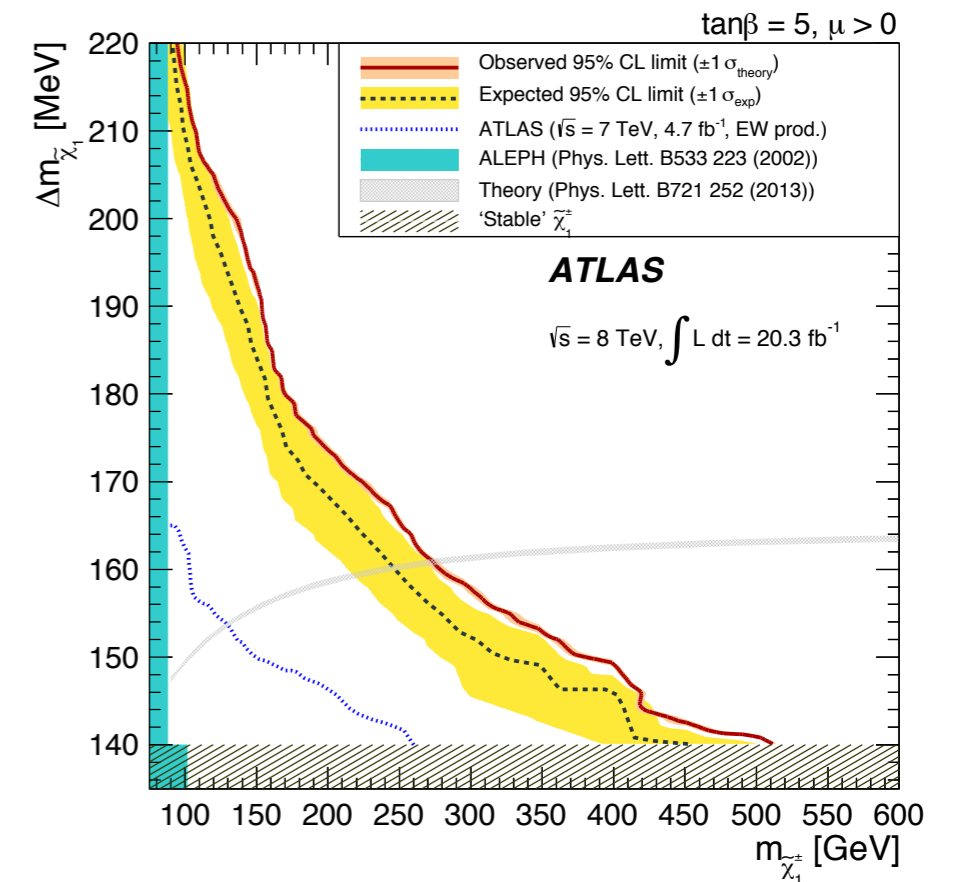
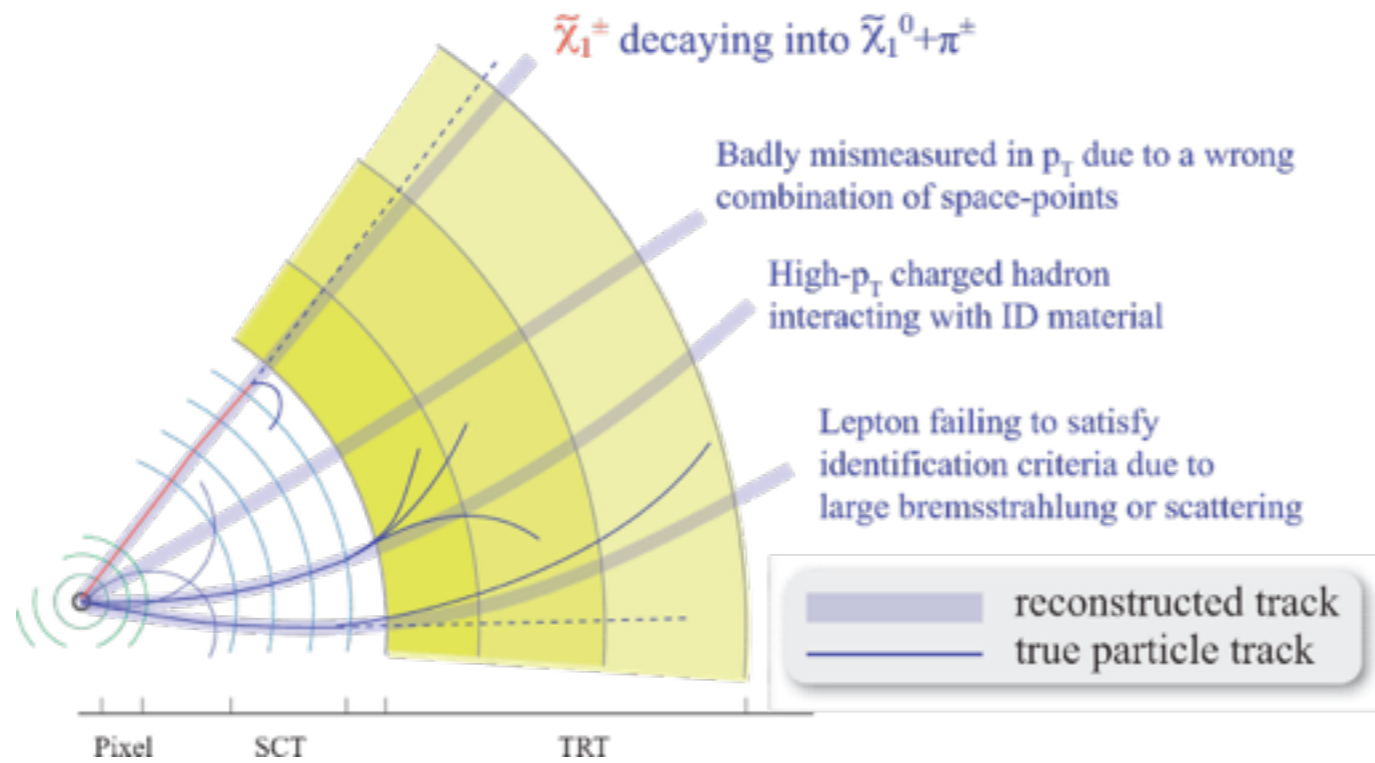
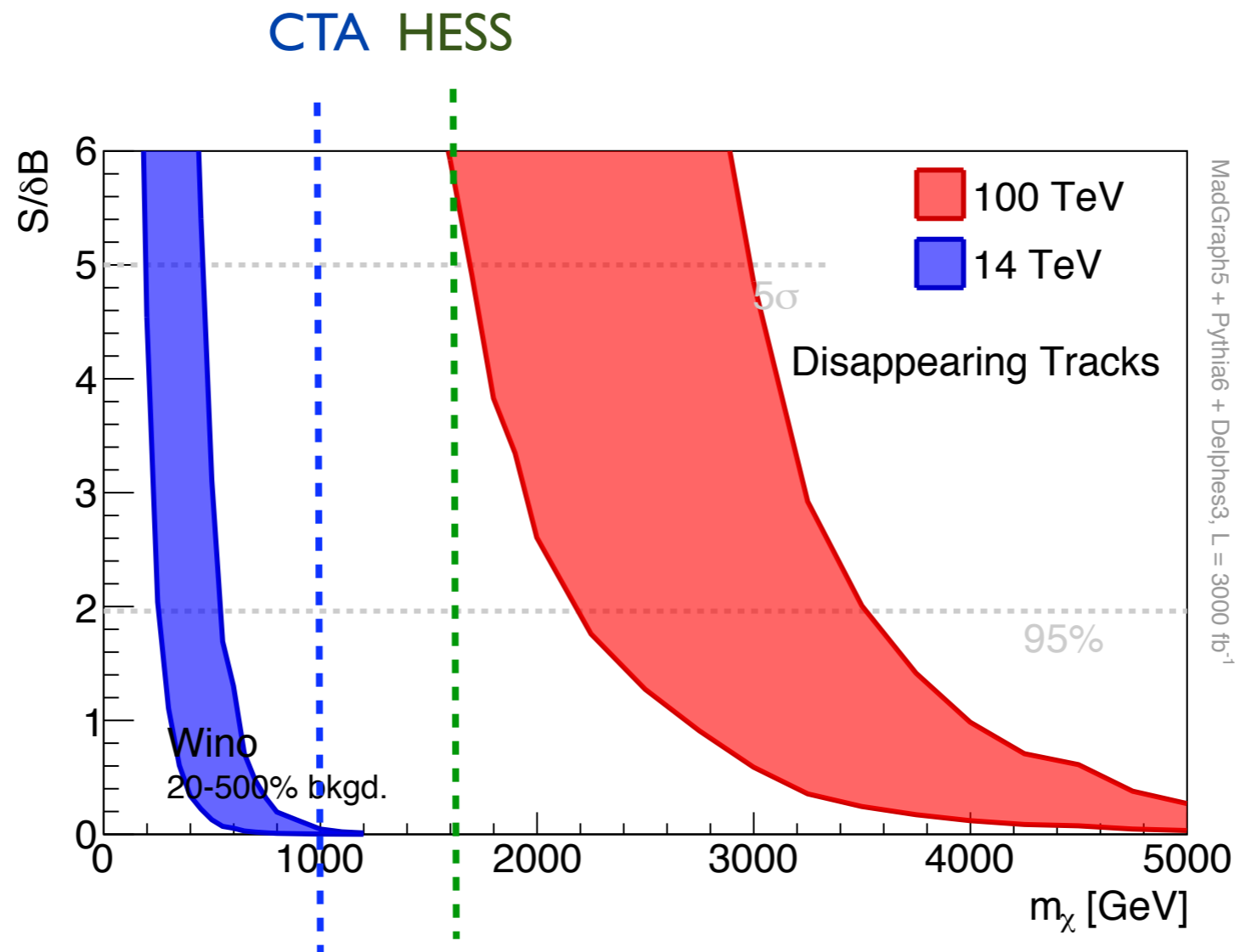


Figure from ATLAS disappearing track search twiki

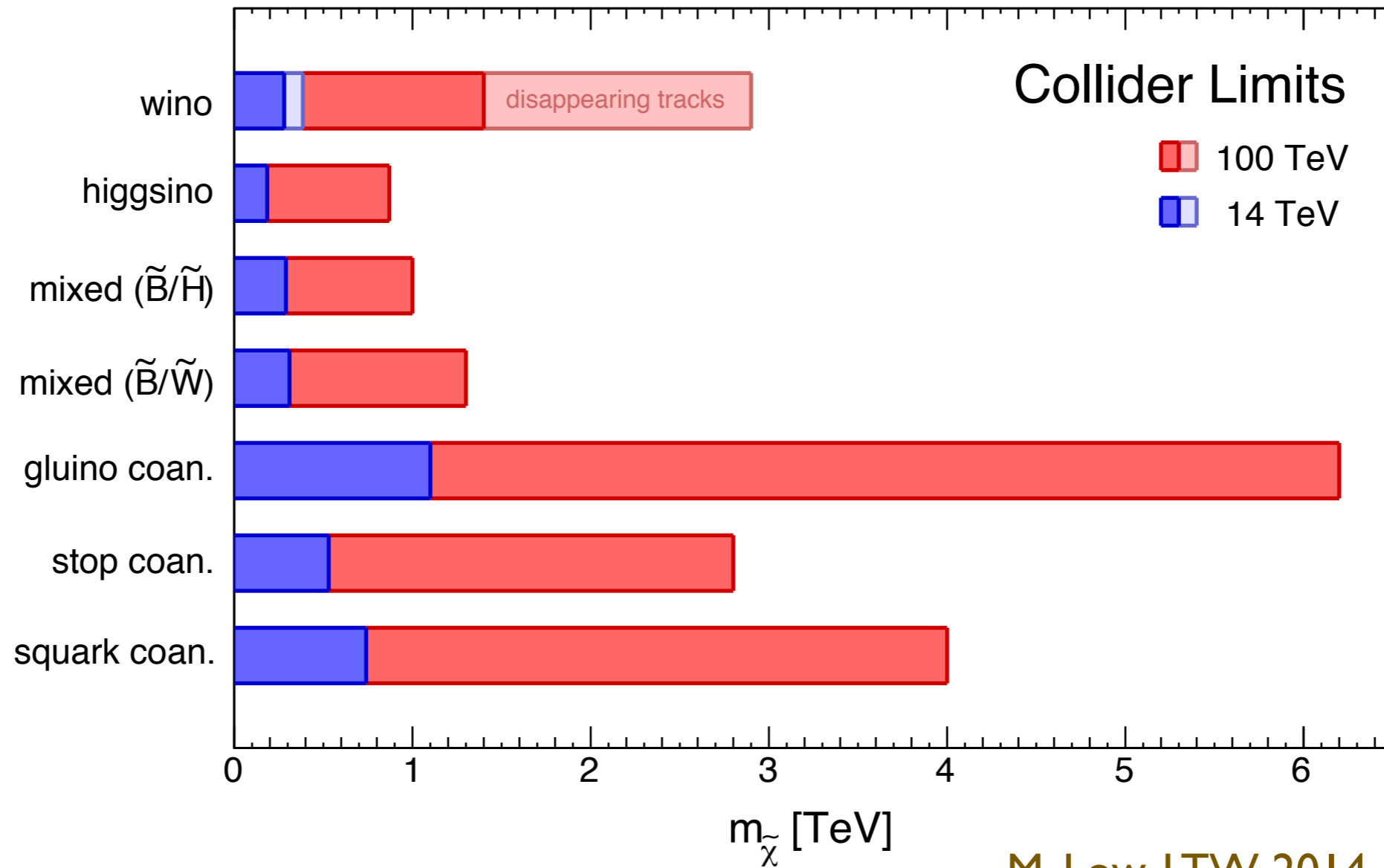
- Main decay mode $\chi^\pm \rightarrow \pi^\pm + \chi^0$.
- Charge track $\approx 10(s)$ cm
- Impressive limit at the LHC already.

Wino



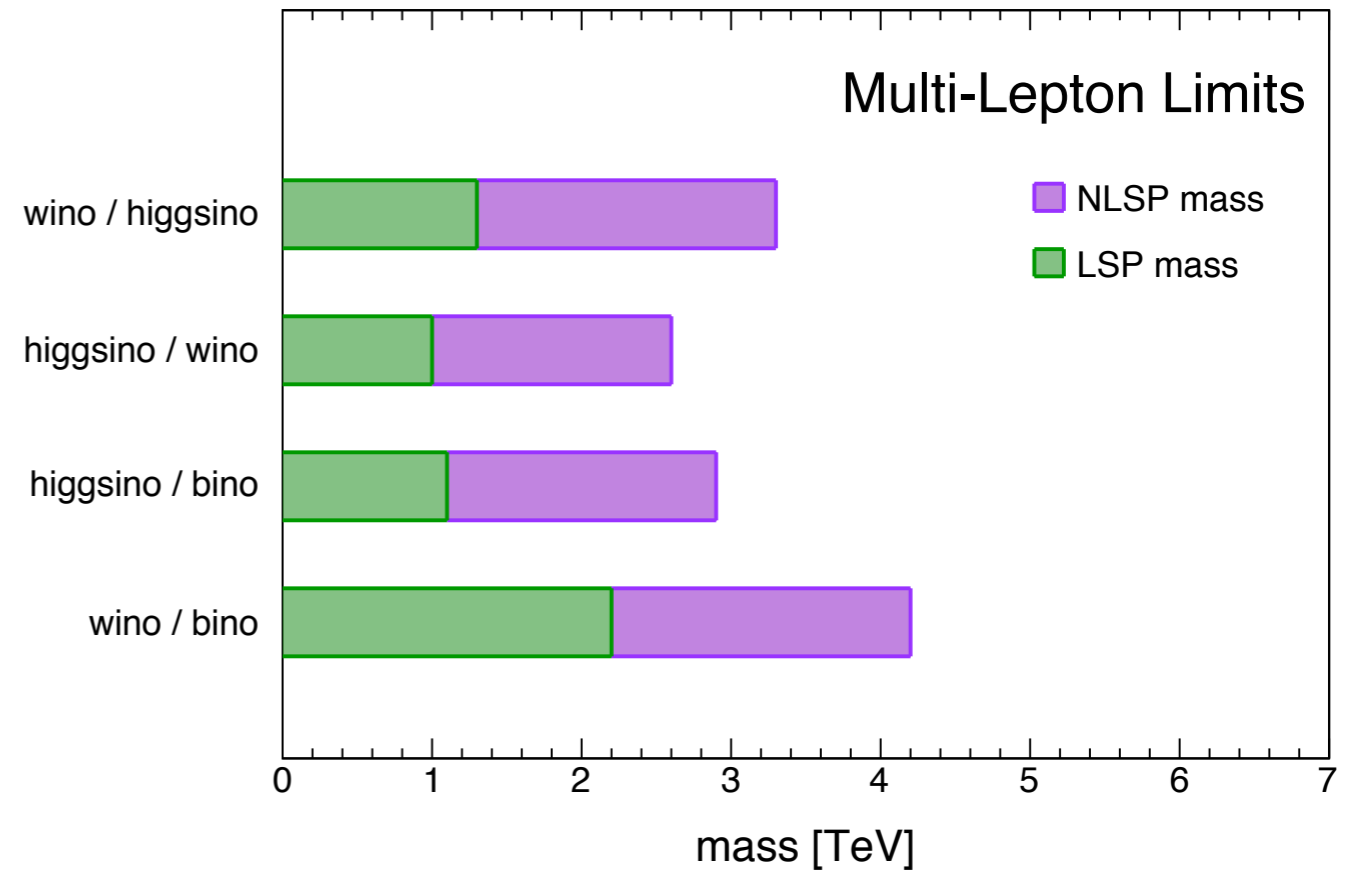
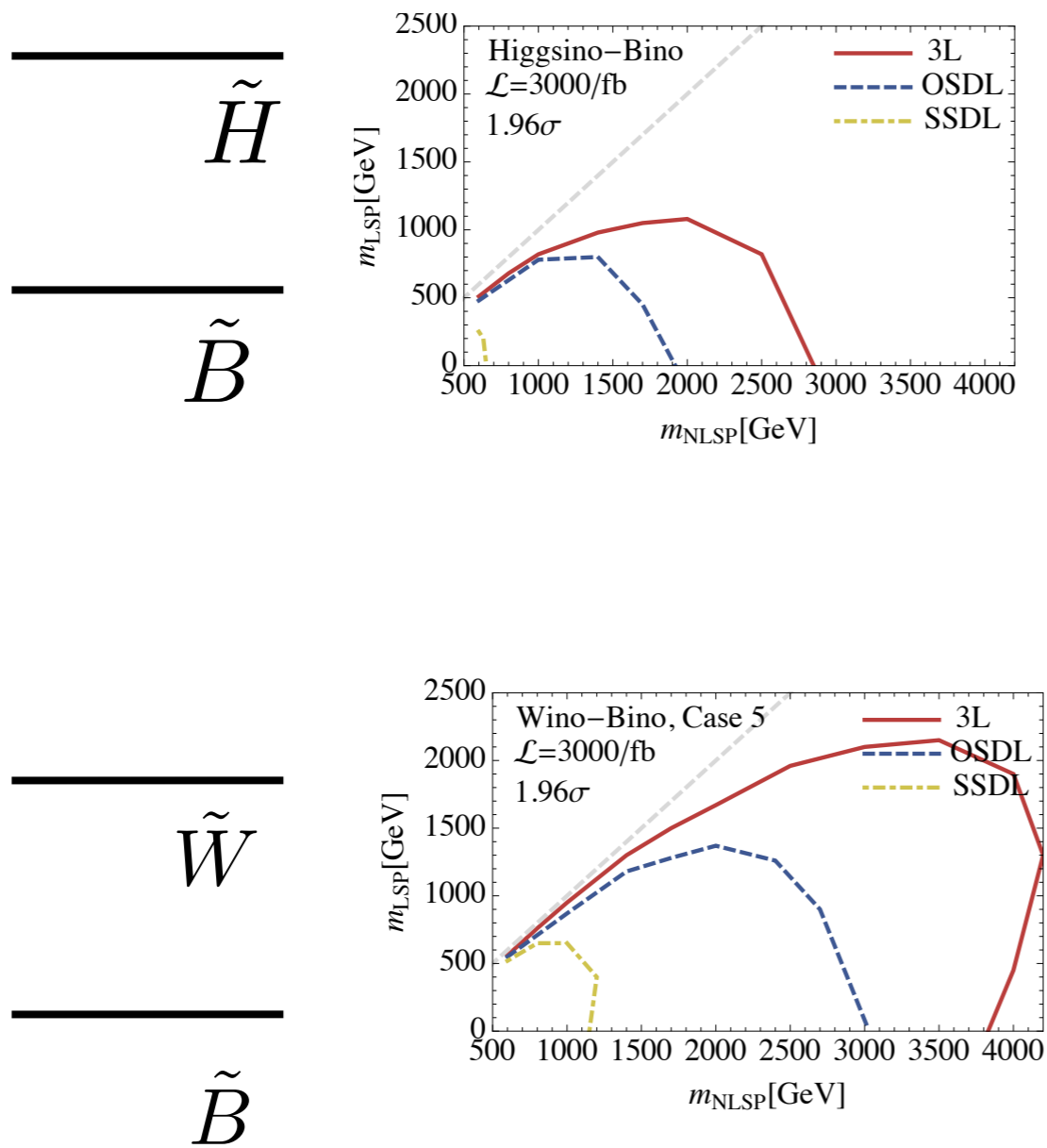
- “Completely cover” the wino parameter space.

Mono-jet



$$M_{\text{WIMP}} \leq 1.8 \text{ TeV} \left(\frac{g^2}{0.3} \right)$$

With cascade decays

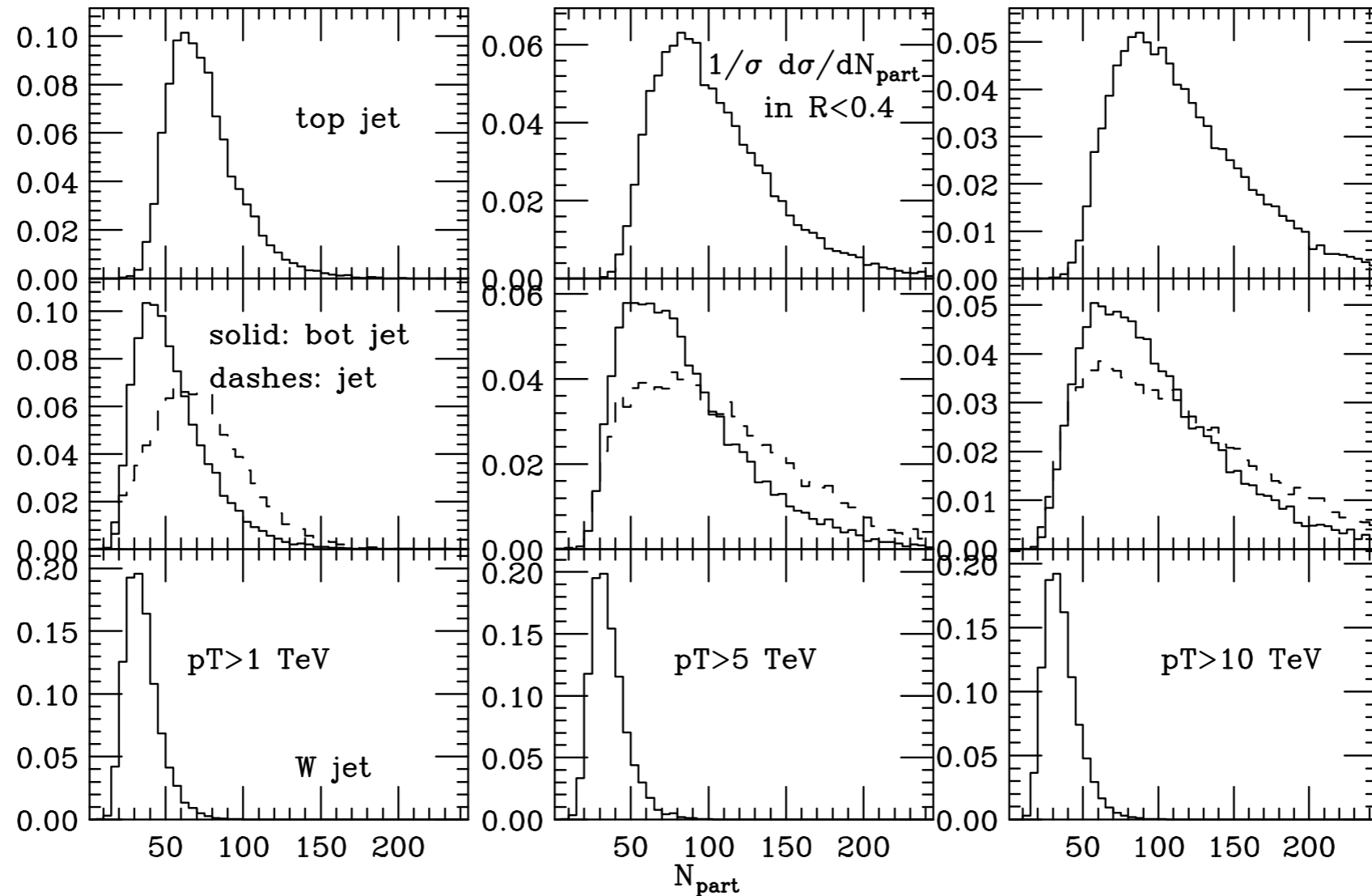


Gori, Jung, Wang, Wells, 2014

Decay \Rightarrow leptons \Rightarrow stronger limits

More novelties at a 100 TeV collider

- Bigger, messier jets.

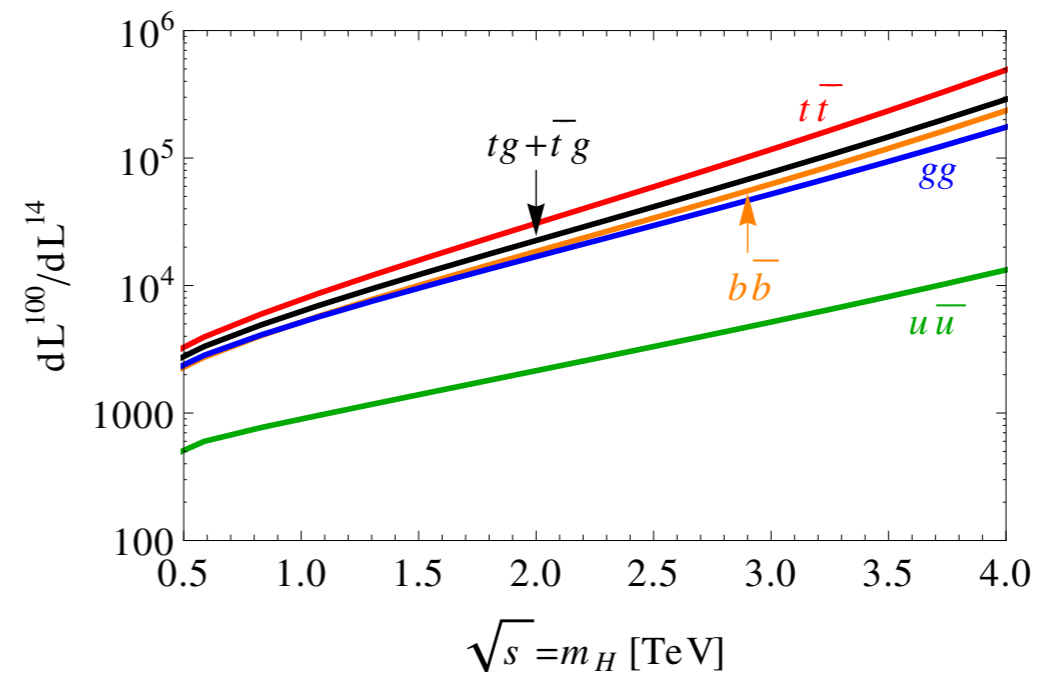
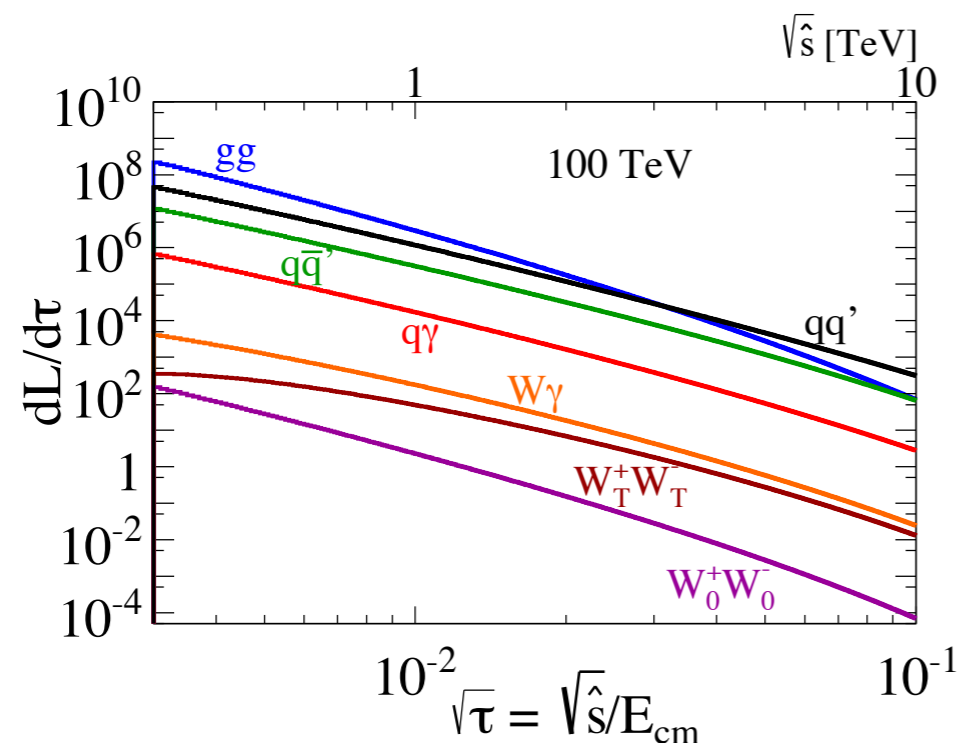


LHC triggered a revolution in jet technology.

100 TeV pp collider demands more!

More novelties at a 100 TeV collider

- SM EW scale particles become very light.
- W/Z/t/h
 - ▶ Treating them as part of the "PDF".



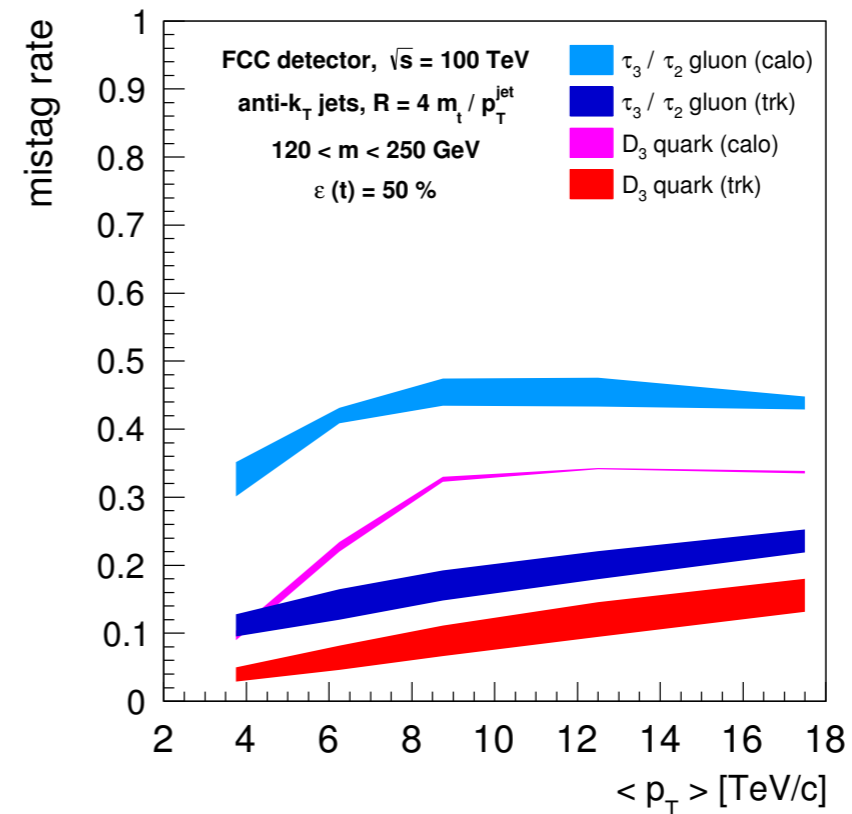
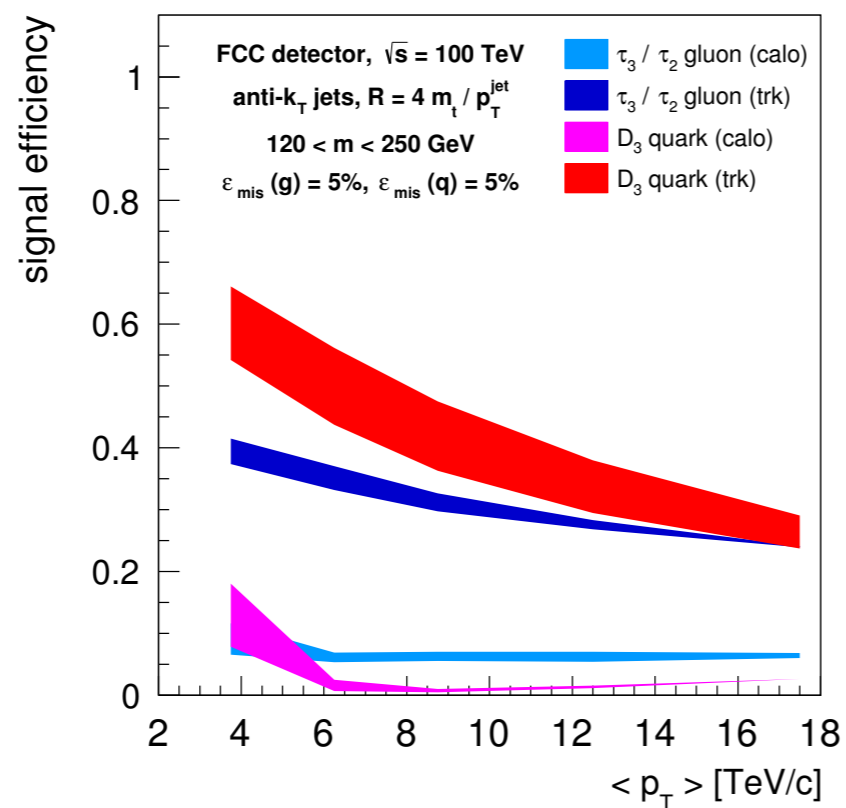
We learned a lot about going from 4 \rightarrow 5 flavors (doing bottom quark properly).

Similar strategy here (?)

More novelties at a 100 TeV collider

- SM EW scale particles become very like.
- Tagging W/Z/t/h as "fat" jets
 - ▶ Not so fat any more, using tracks.

Larkoski, Maltoni, Selvaggi, 2015



New strategies?

Why 100 TeV?

- A benchmark used in the studies.
- Of course, higher is better!
- However, technological + cost constraints
- 100-ish seems to be the best we can do at the moment.
- With further design and physics studies, the number can change.
- A discovery at the LHC can dramatic change the plan.

Comments

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- Physics case of next generation high energy pp collider “obvious”.

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- Physics case of next generation high energy pp collider “obvious”.
- Without LHC discovery.
 - ▶ Physics case for a 100 TeV pp collider stronger than HE-LHC at 28 TeV. Need a big step.
 - ▶ Cost+technological challenge. Perhaps easier to “sell” only as a second step of a circular Higgs factory in longer term.

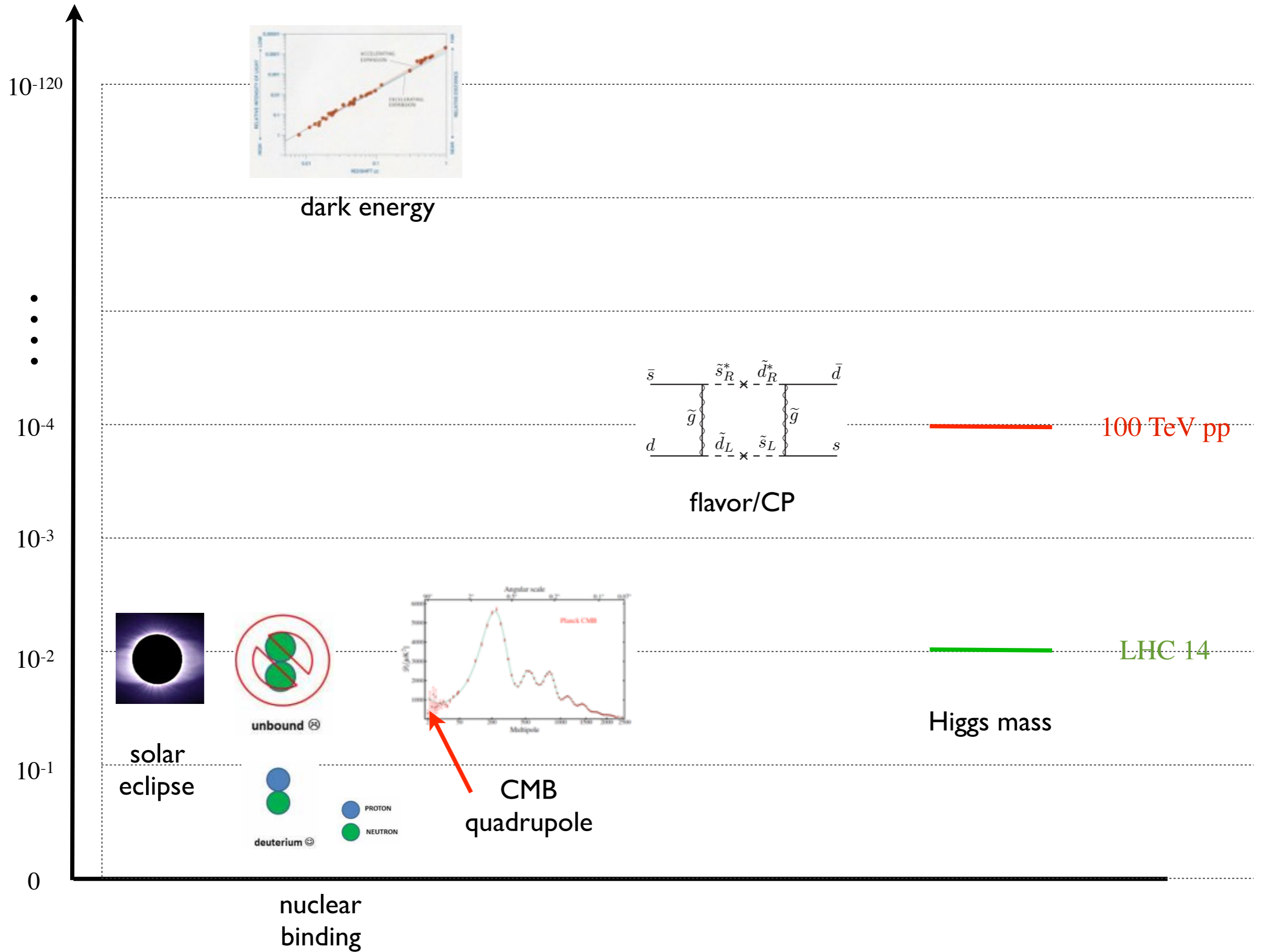
More opportunities and challenges

- Better SM theory calculation needed for taking full advantage of energy and luminosity.
- Many more NP channels, e.g. flavor (violating) physics at 10s TeV?
- Full set of Higgs measurements at 100 TeV, more careful study.
- Physics driven (such as dark matter search) novel detector designs.
- We will do much better than we know now in a couple of decades. cf. LHC vs SppS.



A lot to look forward to!

extras

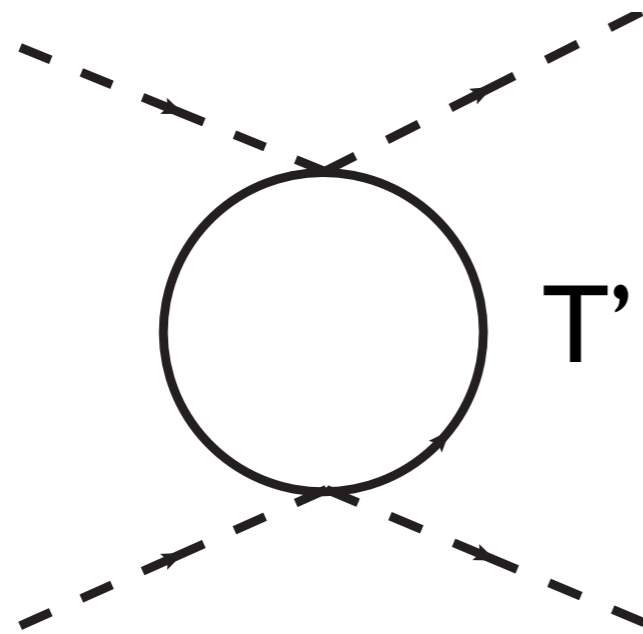


If we made a discovery at run 2

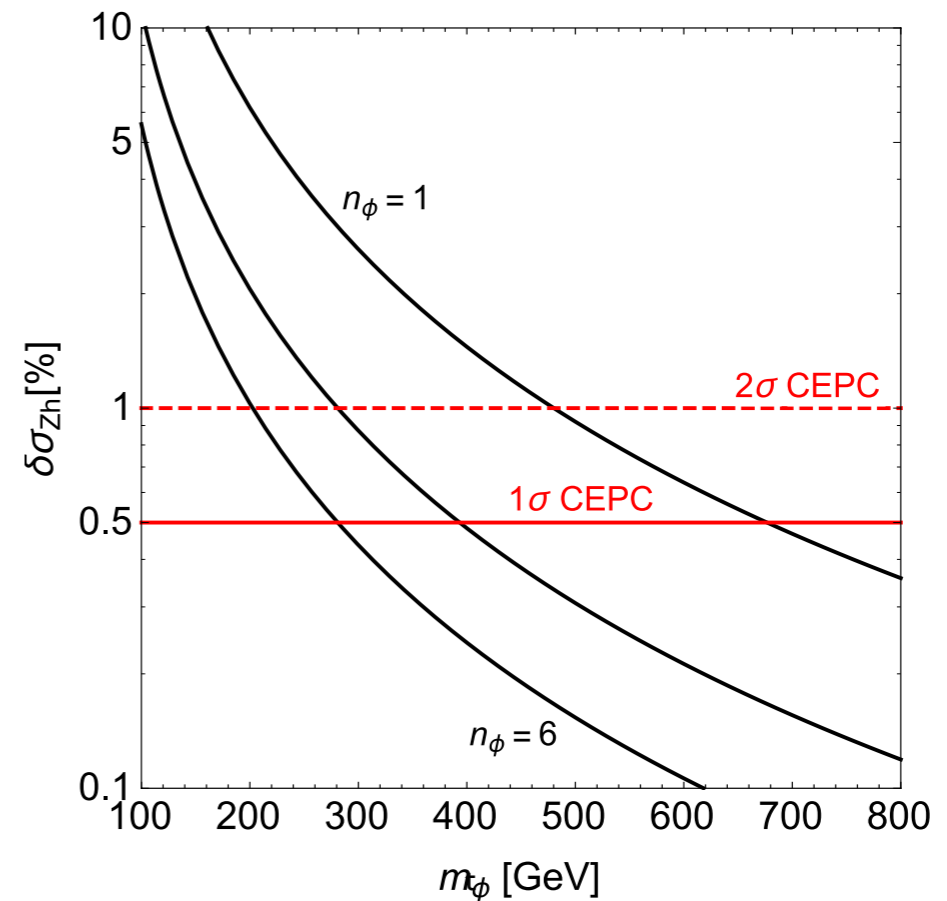
- Beginning of a new era. Seeing the first sign of a new layer of new physics.
- However, it is unlikely to discover the full set of the particles, since we have not see anything yet.
- Typically, going from 8 TeV to 14 TeV increase the reach at most by a factor of 2.
- However, many models feature particles with masses spread at least factor of several apart.
- Won't be able to see everything.
- LHC discovery will set the stage for our next exploration, in particular at a 100 TeV pp collider.

Neutral naturalness

Twin Higgs. Chacko et al. Talk by Craig



Top partner only couple to Higgs.
Wavefunction renormalization
Induce shift in Higgs coupling.



Craig, Englert, McCullough, 2013

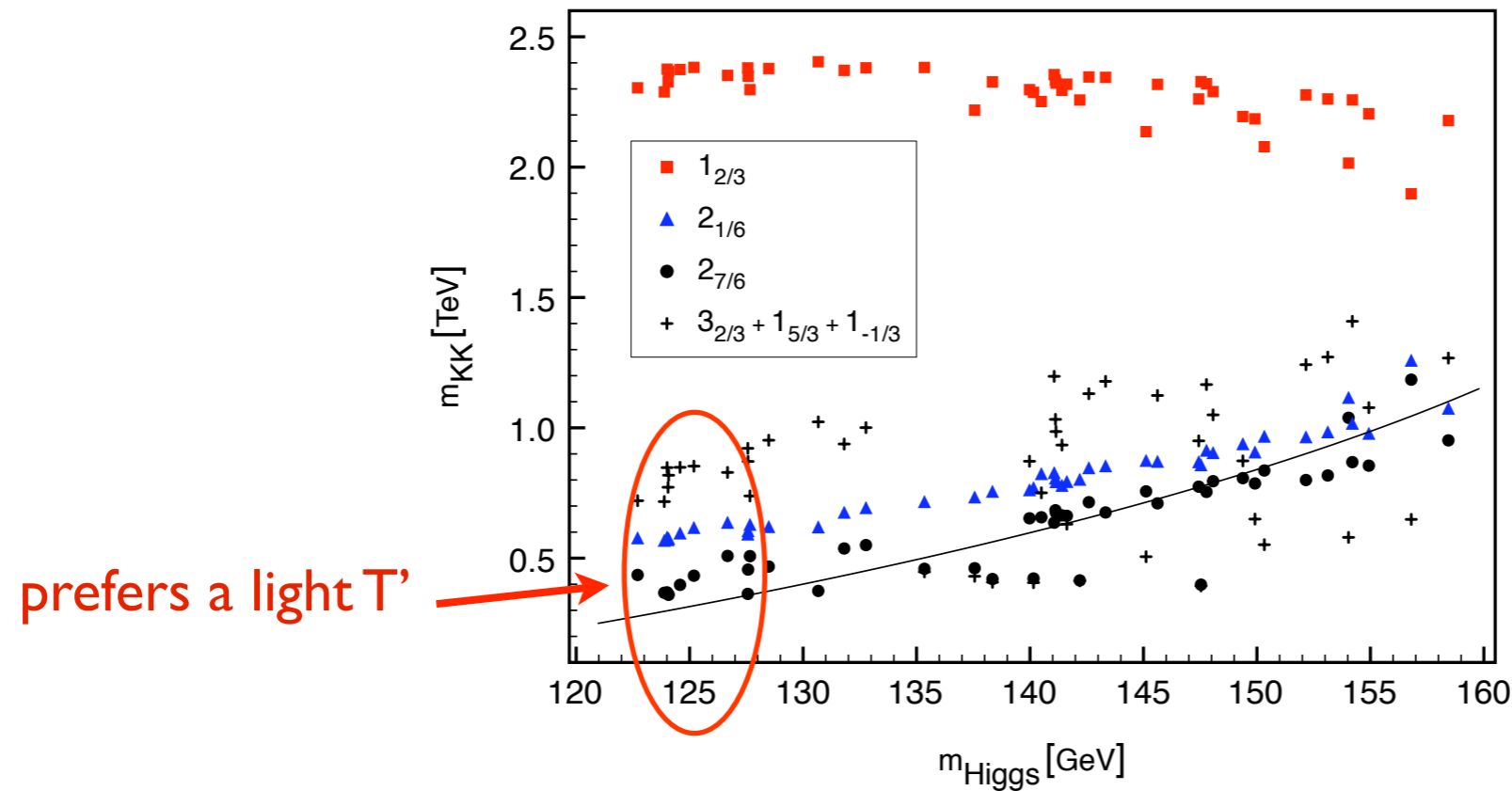
- LHC reach poor. Theory can be completely natural.
- Higgs factory can test this.

Need to consider UV completions for neutral top partners

- Induce measurable shifts in Higgs couplings, precision observables.
- UV completions can be directly probed at 100 TeV.
- Combination of precision measurement and direct search at 100 TeV pp collider can test naturalness.

Compositeness and top partner

Wulzer's talk

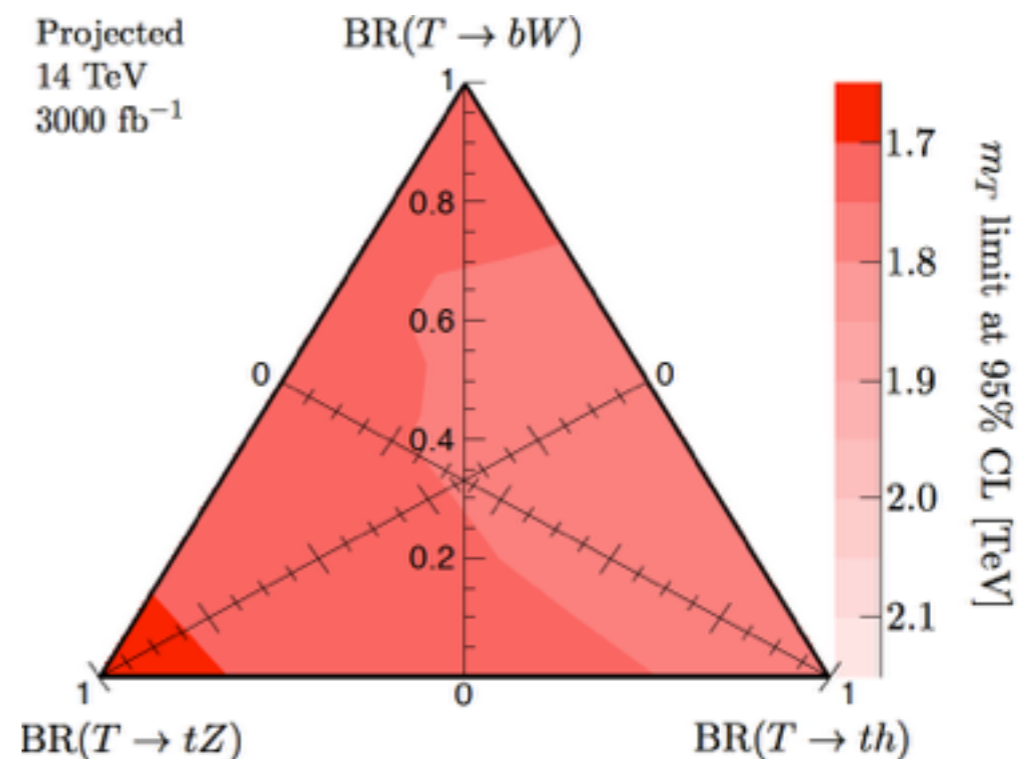
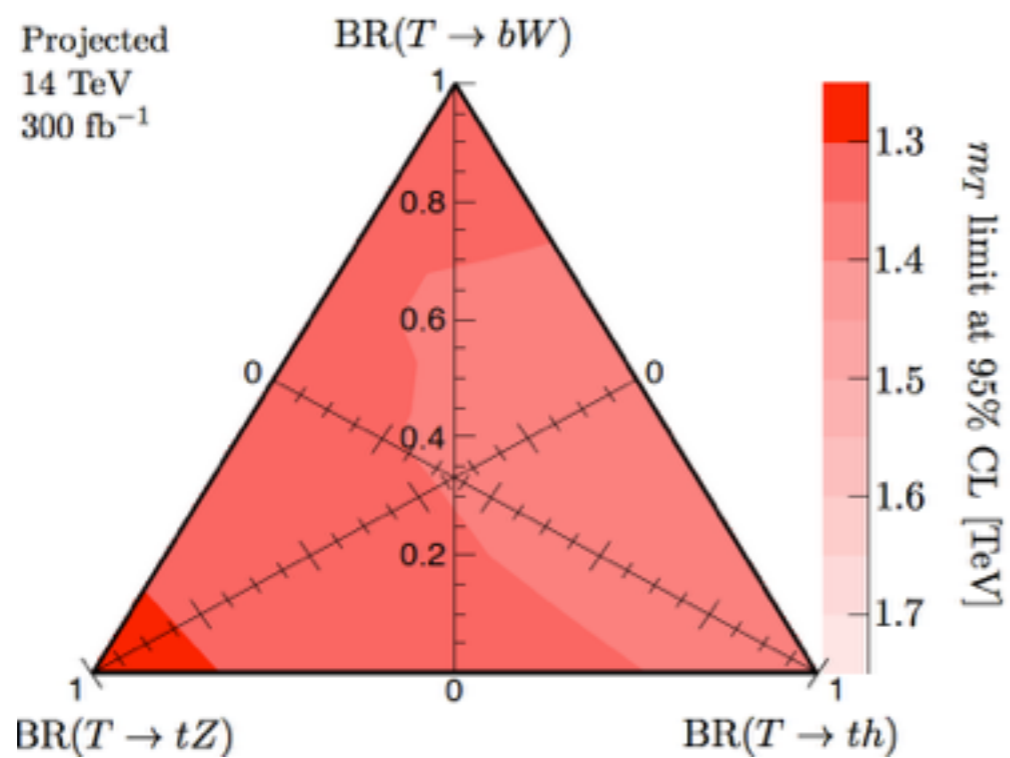
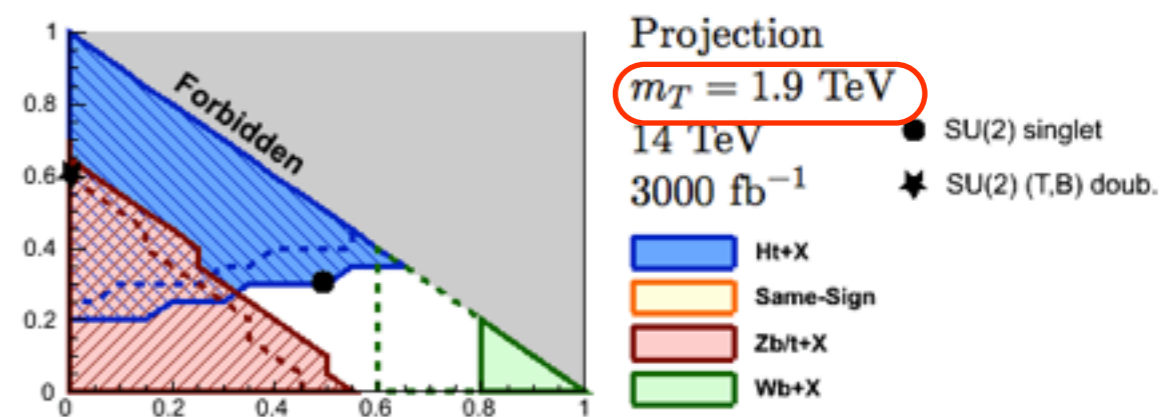
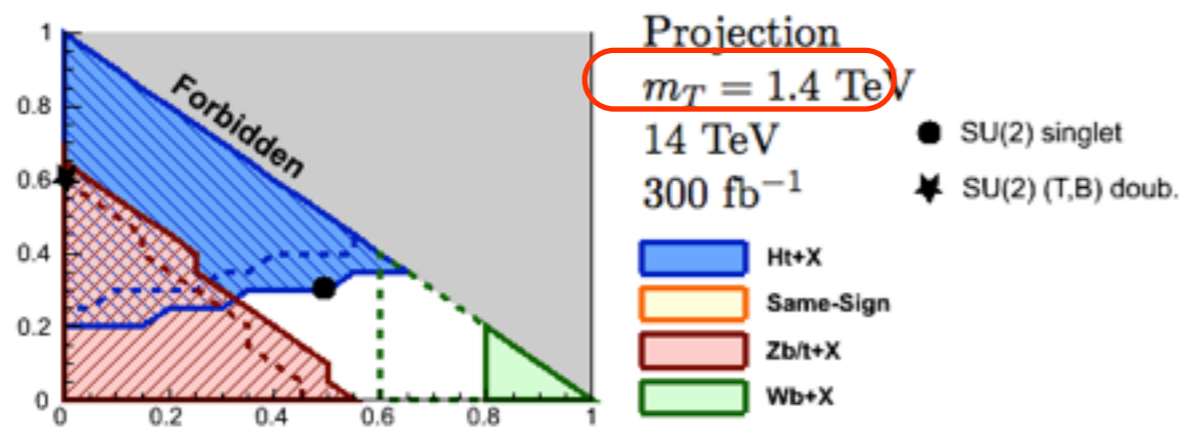


Contino, Da Rold, Pomarol, 2006

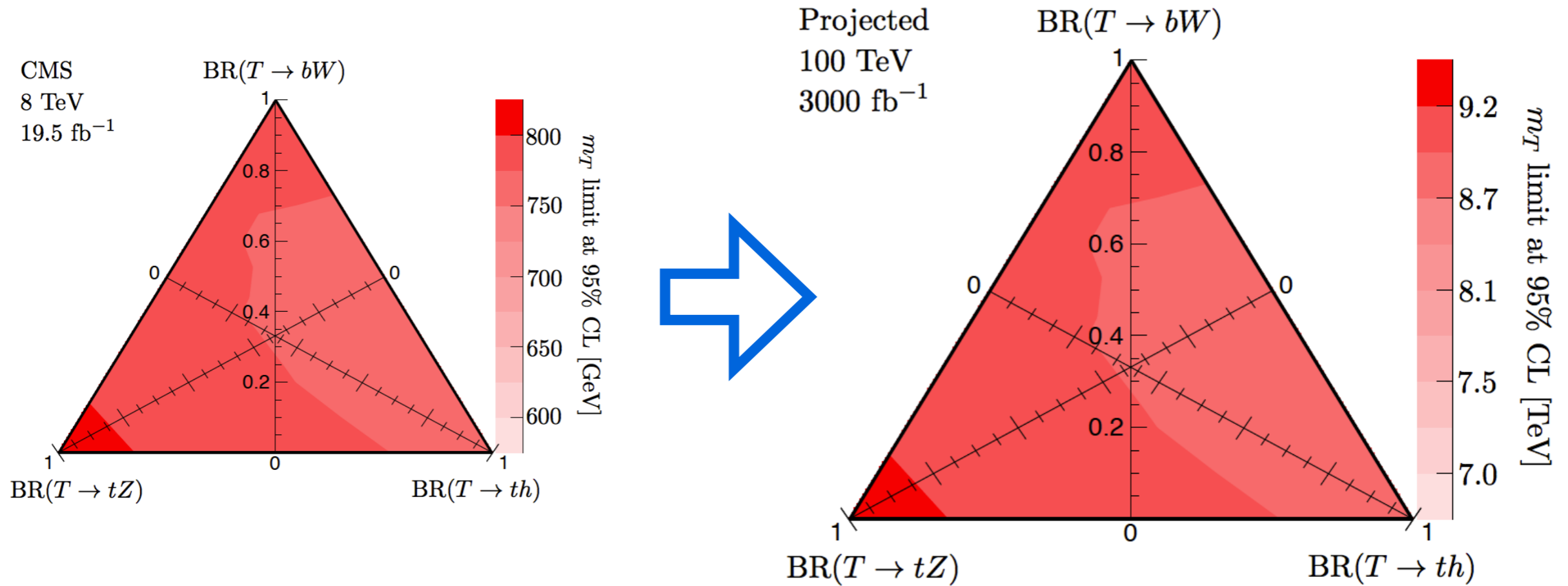
- Plays a crucial role in EWSB.

For a comprehensive discussion, see
De Simone, Matsedonskyi, Rattazzi, Wulzer, 1211.5663

LHC 14 should cover (most of) it.

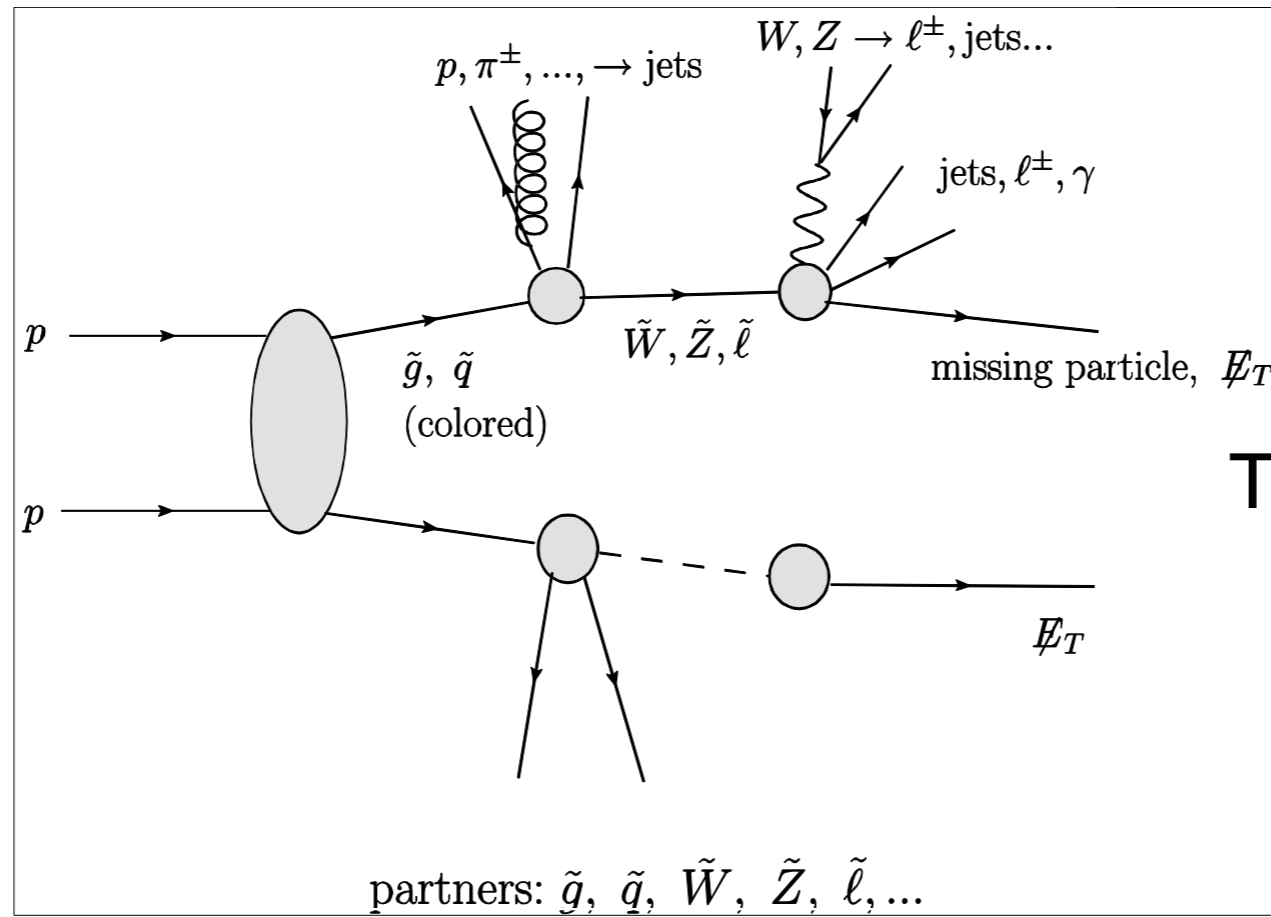


Going up to 100 TeV



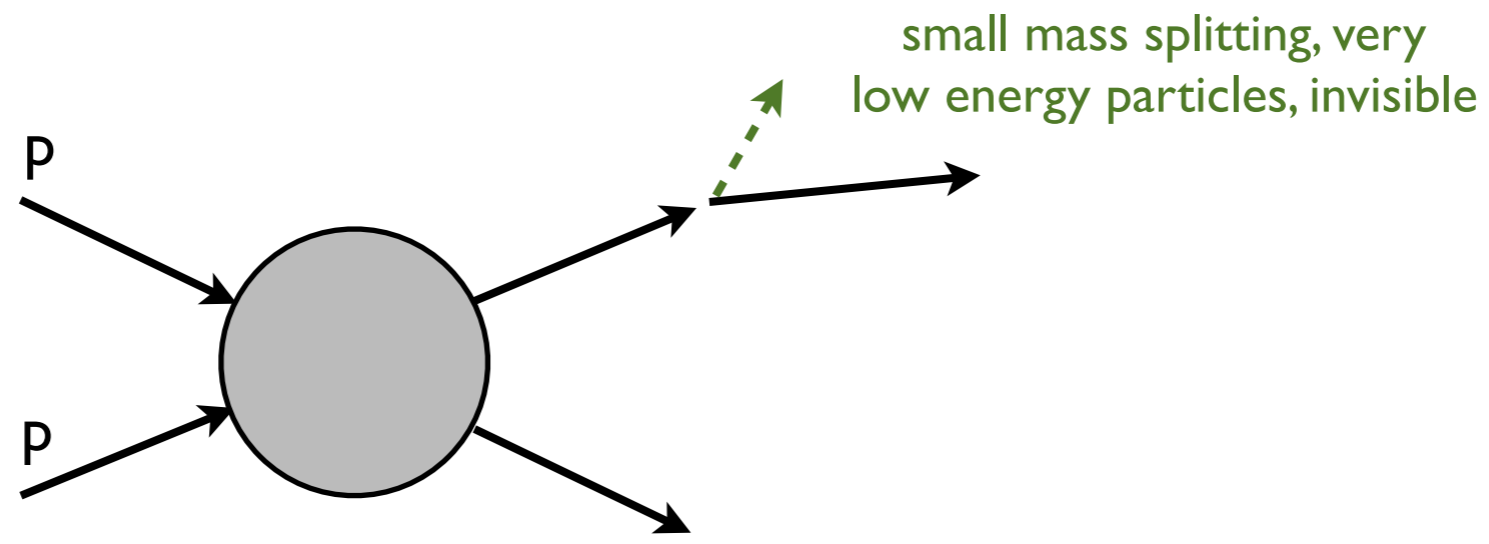
- Again, room for improvement by using single production, boosted technique, etc.

SUSY DM signal in the compressed case

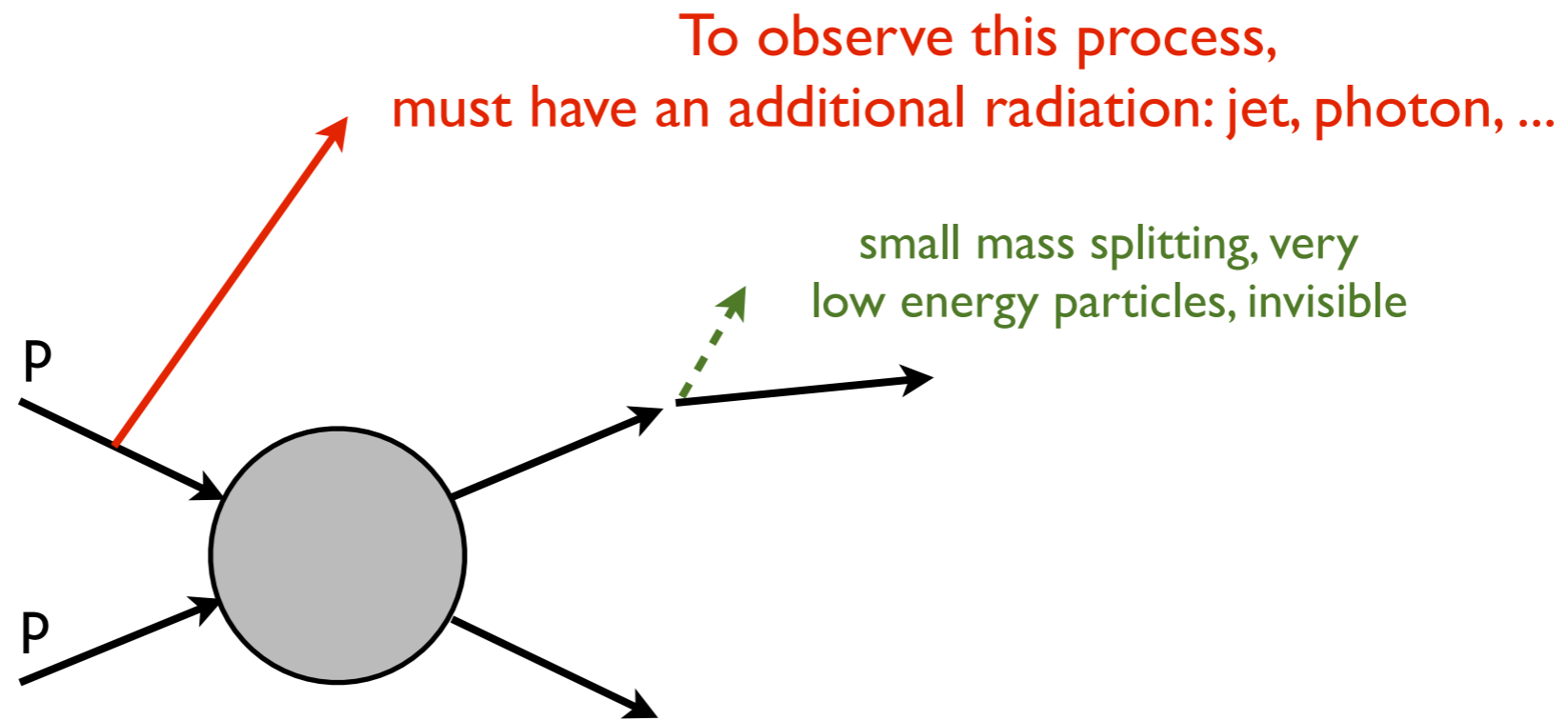


The "usual" story

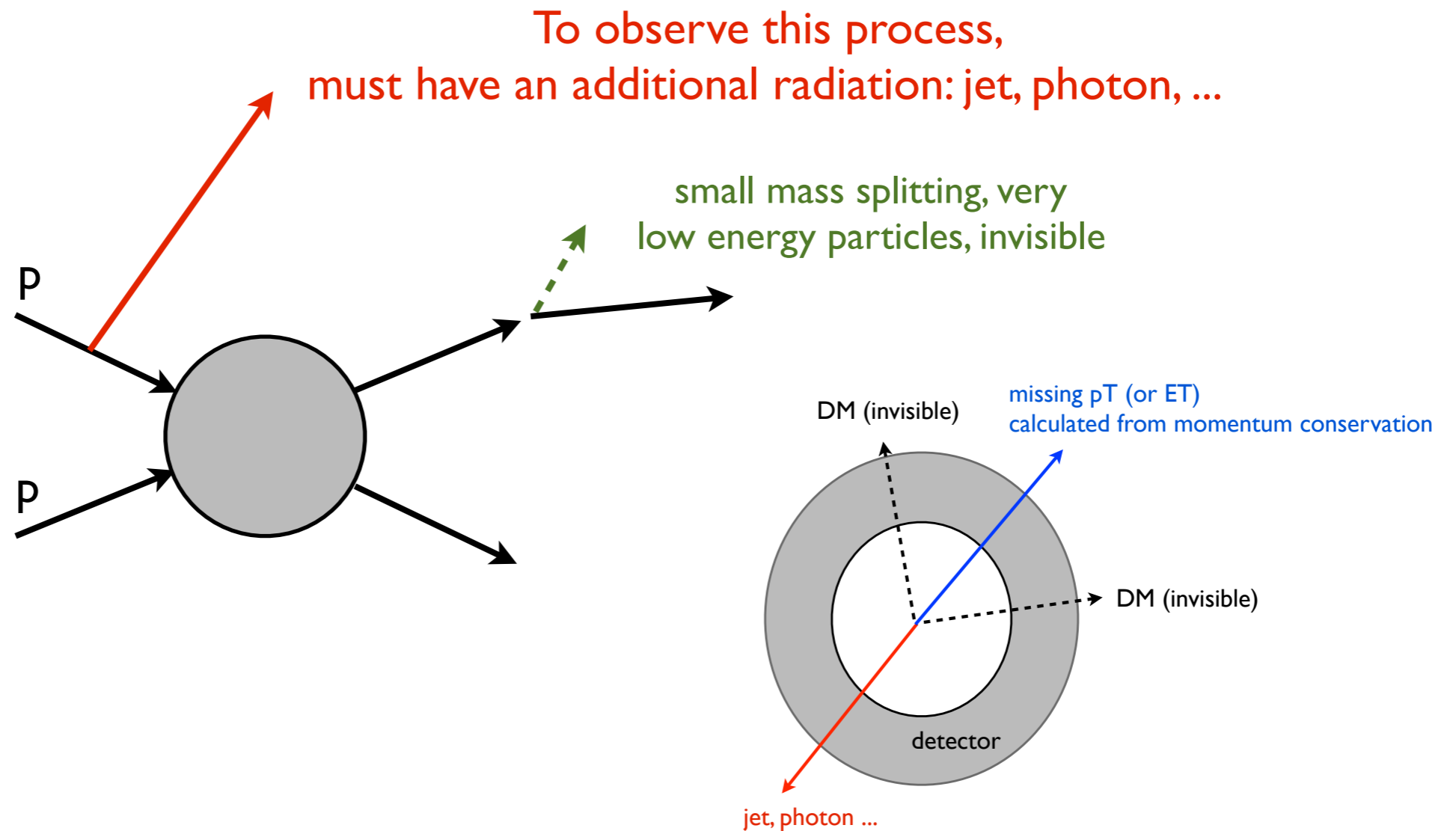
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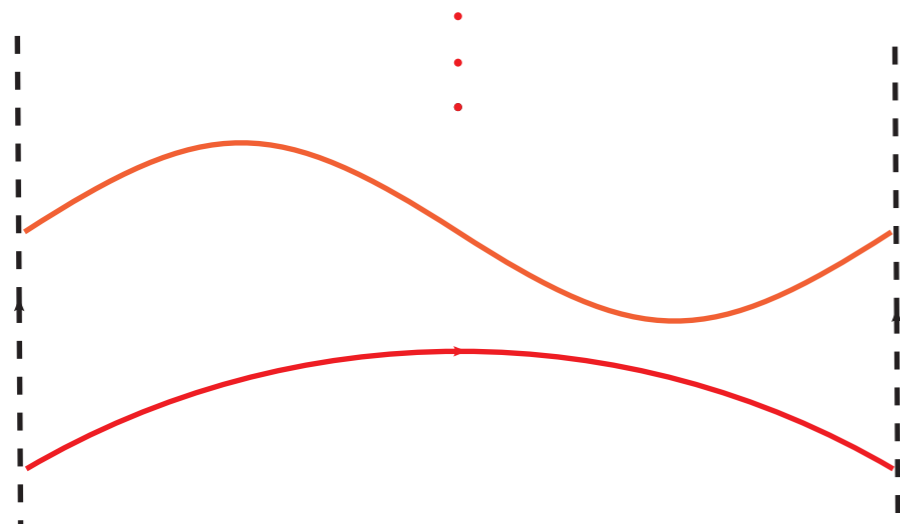
SUSY DM signal in the compressed case



- Back to the basic mono-jet, mono-photon...

Higgs mass in quantum theory.

Quantum fluctuation: Zero point energy



$$\mathcal{H}_{\text{quant}} = \sum_{\vec{p}} \frac{1}{2} \hbar \omega_{\vec{p}} \simeq \int^{|\vec{p}| < \Lambda} \frac{d^3 \vec{p}}{(2\pi)^3} \hbar \omega_{\vec{p}}$$
$$\omega_{\vec{p}} = \sqrt{\vec{p}^2 + m^2} \quad (\hbar = 1)$$

Λ : a cut-off.

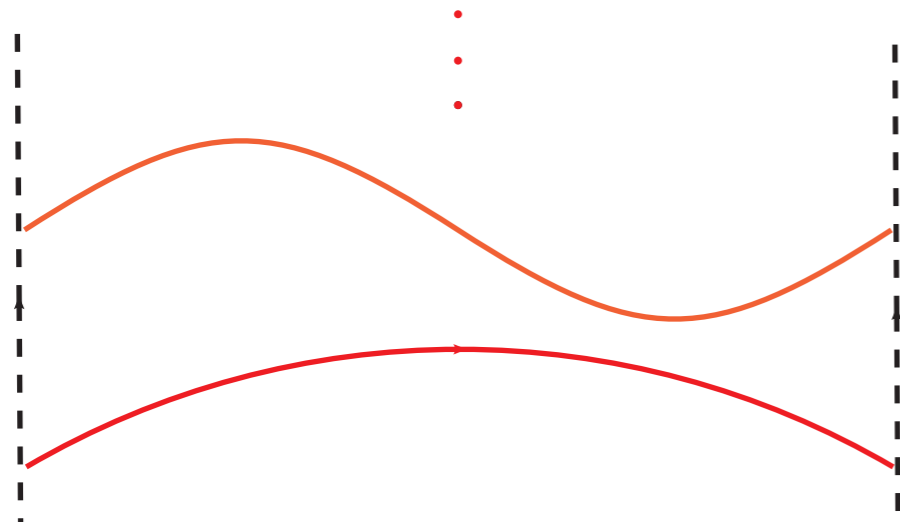
The energy scale of new physics.

Standard Model: include fluctuations of W boson, top quark, ...

$$m_W = g_2 h, \quad m_{\text{top}} = y_t h \quad \mathcal{H}_{\text{quant}} \simeq \frac{9}{64\pi^2} g_2^2 \Lambda^2 h^2 - \frac{3}{8\pi^2} y_t^2 \Lambda^2 h^2 + \dots$$

Higgs mass in quantum theory.

Quantum fluctuation: Zero point energy



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– Renormalization: $m_h^2(\text{physical}) = m_0^2 + c \Lambda^2$

► m_0^2 can always be adjusted to give correct $m_h^2(\text{physical})$.

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- No large cancellation $\Rightarrow m_h^2$ (physical) $\approx c \Lambda^2$
Naturalness criterion leads to a prediction of the mass scale of new physics!!

Rate for double Higgs production.

