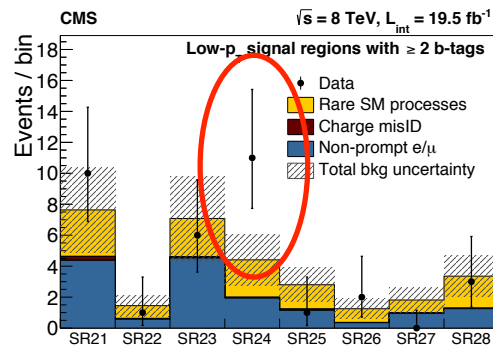
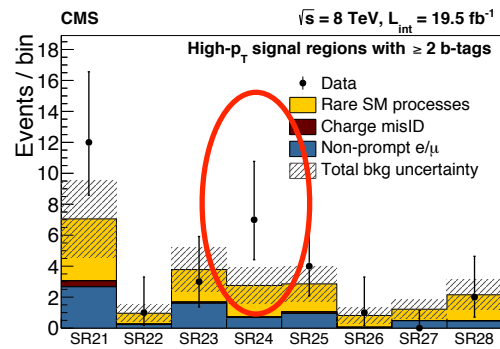


# Same Sign Lepton Excesses

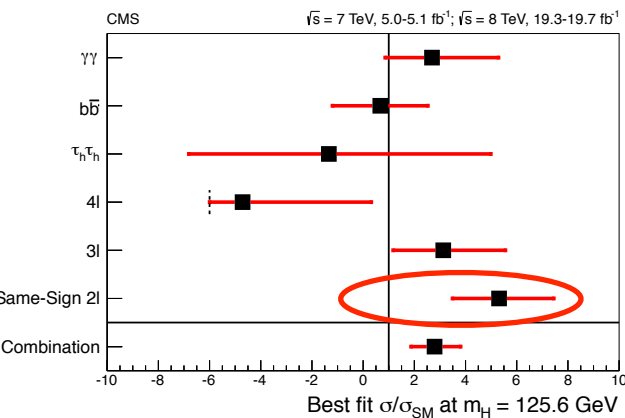


SR1b $1_{bin}$	Total	$ee$	$e\mu$	$\mu\mu$
Observed events	10	6	4	0
Total expected background events	$4.7 \pm 2.1$	$1.4 \pm 0.8$	$2.1 \pm 1.1$	$1.2 \pm 0.4$
<b>Components of the background</b>				
$t\bar{t}V$ , $t\bar{t}H$ , $tZ$ and $t\bar{t}t$	$2.5 \pm 1.7$	$0.6 \pm 0.3$	$1.2 \pm 1.0$	$0.7 \pm 0.3$
Dibosons and tribosons	$0.9 \pm 0.4$	$0.10 \pm 0.04$	$0.3 \pm 0.1$	$0.5 \pm 0.3$
Fake leptons	$0.8^{+1.2}_{-0.8}$	$0.4^{+0.7}_{-0.4}$	$0.4^{+0.5}_{-0.4}$	$< 0.1$
Charge-flip electrons	$0.5 \pm 0.1$	$0.3 \pm 0.1$	$0.3 \pm 0.1$	–
$p(s=0)$	0.07	0.01	0.18	0.50

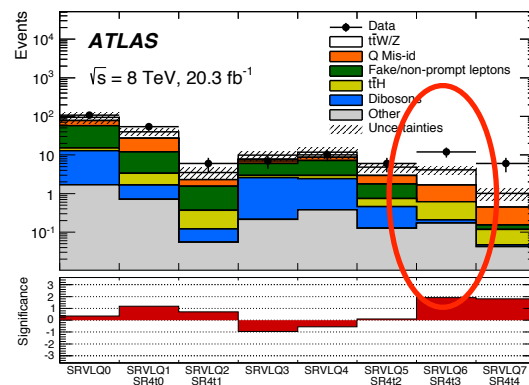
CMS (SUSY), <http://arxiv.org/abs/1311.6736>

(24 signal regions in paper)

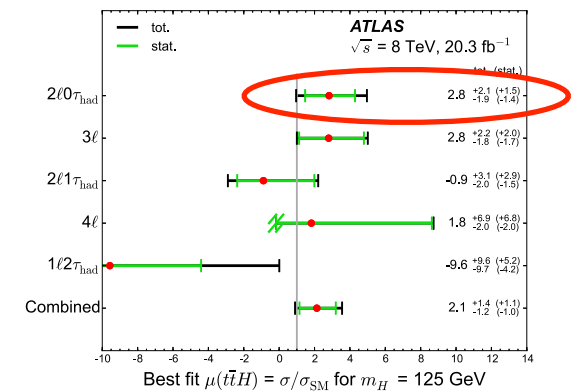
ATLAS (SUSY), <http://arxiv.org/abs/1404.2500> (5 signal regions in paper)



CMS (ttH), <http://arxiv.org/abs/1408.1682>



ATLAS (TT), <http://arxiv.org/abs/1504.04605>



ATLAS (ttH), <http://arxiv.org/abs/1506.05988>

The ATLAS analyses are correlated, and same for CMS  
 So, ~2 analyses and excesses are  $< 3 \sigma$   
 Worth keeping an eye on? Sure.

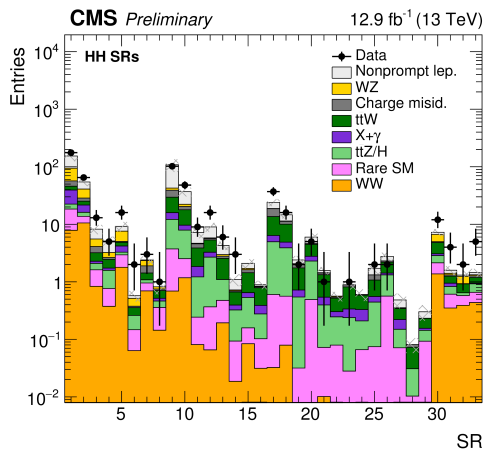
# 13 TeV

ATLAS-CONF-2016-037

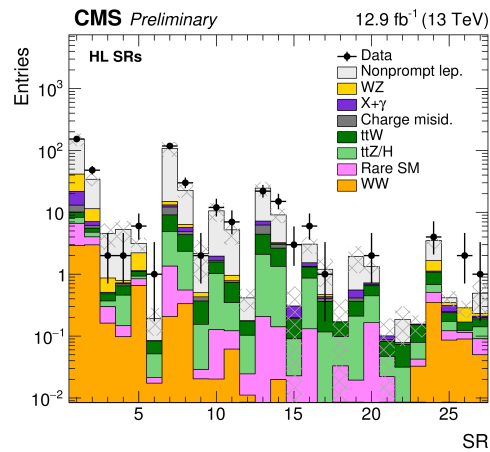
	SR3L1	SR3L2	SR0b1	SR0b2	SR1b
Observed	6	2	5	0	12
Total SM background	$6.1 \pm 2.2$	$1.2 \pm 0.5$	$8.8 \pm 2.9$	$1.6 \pm 0.8$	$11.4 \pm 2.8$
$t\bar{t}Z$	$0.69 \pm 0.25$	$0.10 \pm 0.04$	$0.45 \pm 0.18$	$0.10 \pm 0.04$	$1.6 \pm 0.6$
$t\bar{t}W$	$0.09 \pm 0.04$	$0.02 \pm 0.01$	$0.45 \pm 0.17$	$0.13 \pm 0.06$	$2.0 \pm 0.7$
Diboson	$4.2 \pm 2.0$	$0.7 \pm 0.4$	$3.7 \pm 1.9$	$0.7 \pm 0.5$	$0.5 \pm 0.4$
Rare	$0.8 \pm 0.4$	$0.21 \pm 0.13$	$0.8 \pm 0.4$	$0.18 \pm 0.12$	$2.7 \pm 0.9$
Fake/non-prompt leptons	$0.29 \pm 0.29$	$0.15 \pm 0.15$	$2.9 \pm 2.0$	$0.4 \pm 0.5$	$3.3 \pm 2.1$
Charge-flip	–	–	$0.50 \pm 0.09$	$0.08 \pm 0.03$	$1.43 \pm 0.19$

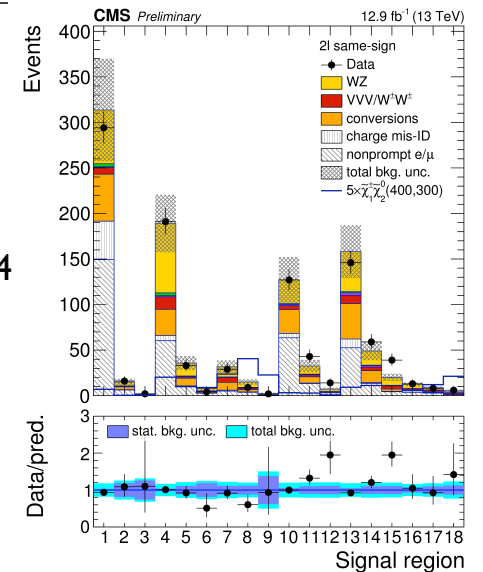
	SR3b	SR1b-GG	SR1b-DD	SR3b-DD
Observed	2	2	12	4
Total SM background	$1.6 \pm 0.6$	$1.7 \pm 0.5$	$12.0 \pm 2.7$	$1.9 \pm 0.8$
$t\bar{t}Z$	$0.19 \pm 0.07$	$0.26 \pm 0.08$	$2.8 \pm 0.9$	$0.30 \pm 0.10$
$t\bar{t}W$	$0.17 \pm 0.06$	$0.33 \pm 0.11$	$1.8 \pm 0.6$	$0.18 \pm 0.07$
Diboson	$< 0.1$	$0.08 \pm 0.19$	$0.6 \pm 0.4$	$< 0.1$
Rare	$0.89 \pm 0.31$	$0.64 \pm 0.34$	$2.6 \pm 1.3$	$0.8 \pm 0.4$
Fake/non-prompt leptons	$0.2 \pm 0.5$	$0.21 \pm 0.33$	$2.5 \pm 1.7$	$0.5 \pm 0.6$
Charge-flip	$0.14 \pm 0.03$	$0.18 \pm 0.07$	$1.74 \pm 0.22$	$0.14 \pm 0.03$



CMS-PAS-SUS-16-020



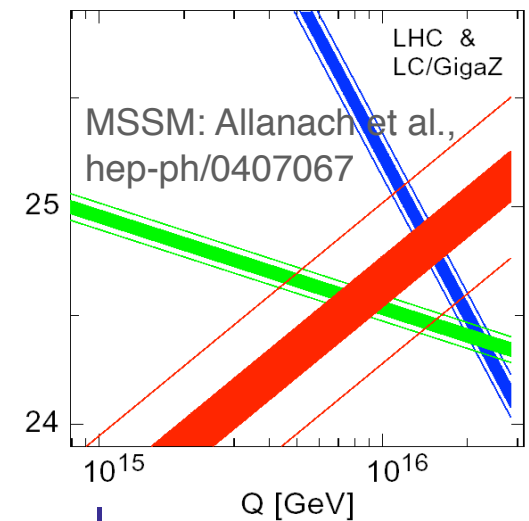
CMS-PAS-SUS-16-024



Not much there....

# Not SUSY?

- ❖ SUSY theories (and others with full or partial set of SM-partners) have a number of attractive features
  - “Explanation” for low Higgs mass (and sometimes EWSB)
  - Gauge coupling unification (often)
  - Dark matter candidate (if introduce a new parity, natural in UED, ~ad-hoc in SUSY)
  - No new interactions (often)
- ❖ But answering those questions comes at a large cost
  - Many new particles, with masses and mixing angles
  - Need to explain why mass scale is so low (or high), spin?



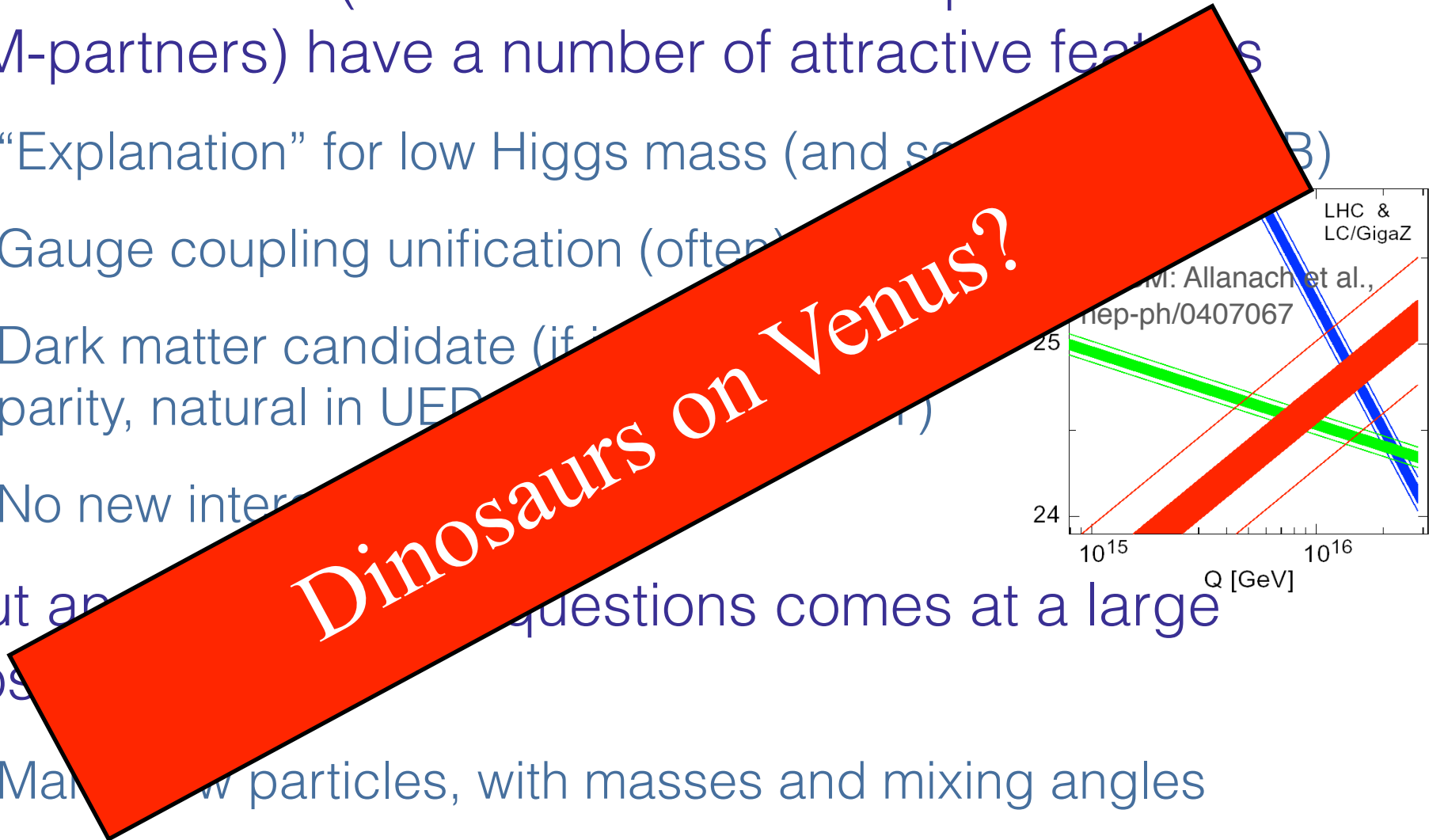
# Not SUSY?

❖ SUSY theories (and others with full or partial set of SM-partners) have a number of attractive features

- “Explanation” for low Higgs mass (and so for the hierarchy problem)
- Gauge coupling unification (often)
- Dark matter candidate (if  $R$ -parity, natural in UED)
- No new interactions

❖ But an important question comes at a large scale

- Many new particles, with masses and mixing angles
- Need to explain why mass scale is so low (or high), spin?

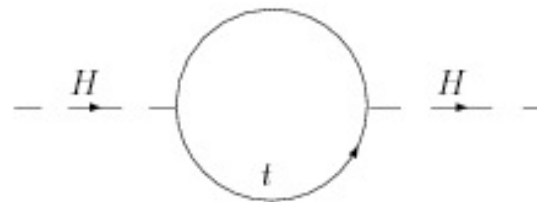




# Less Ambitious

# Giving up on Dark Matter

- ❖ Electroweak-scale WIMPs fit the data well
  - But maybe hard/impossible to produce at colliders
- ❖ Or dark matter not WIMPs at all
- ❖ Back to problem #1:



A Feynman diagram showing a top quark loop. Two external Higgs boson lines, labeled 'H', enter from the left and exit to the right. The loop is a circle with an arrow indicating a clockwise path, labeled 't' at the bottom. A blue arrow points from the diagram to the right, leading to the mathematical expression  $\frac{3}{16\pi^2} y_t^2 E^2$ .

➔ Top partner!

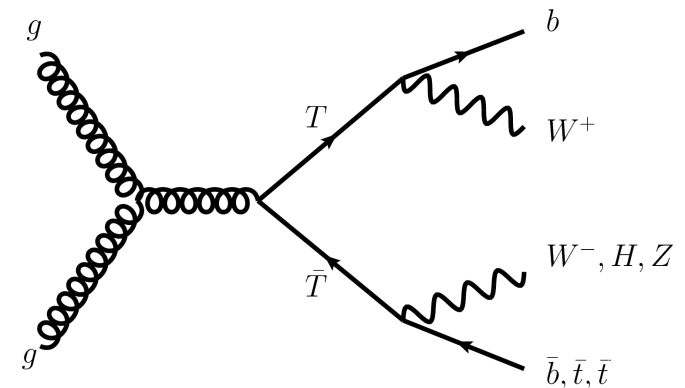
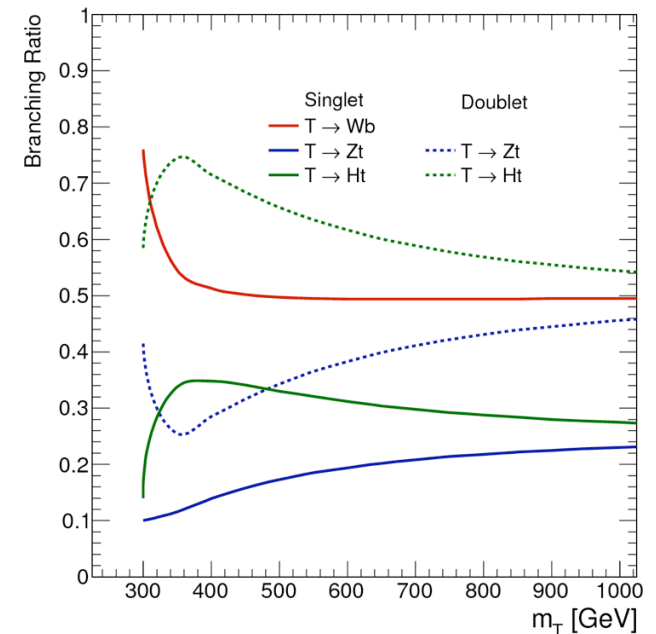
# Singlets, Doublets, ...

❖ Vector-like top partners (still fermions) less constrained by flavor....

- Opens up decay modes
- Top partner partners:
  - $\chi^{5/3}$
  - ...

❖ Rich set of signatures

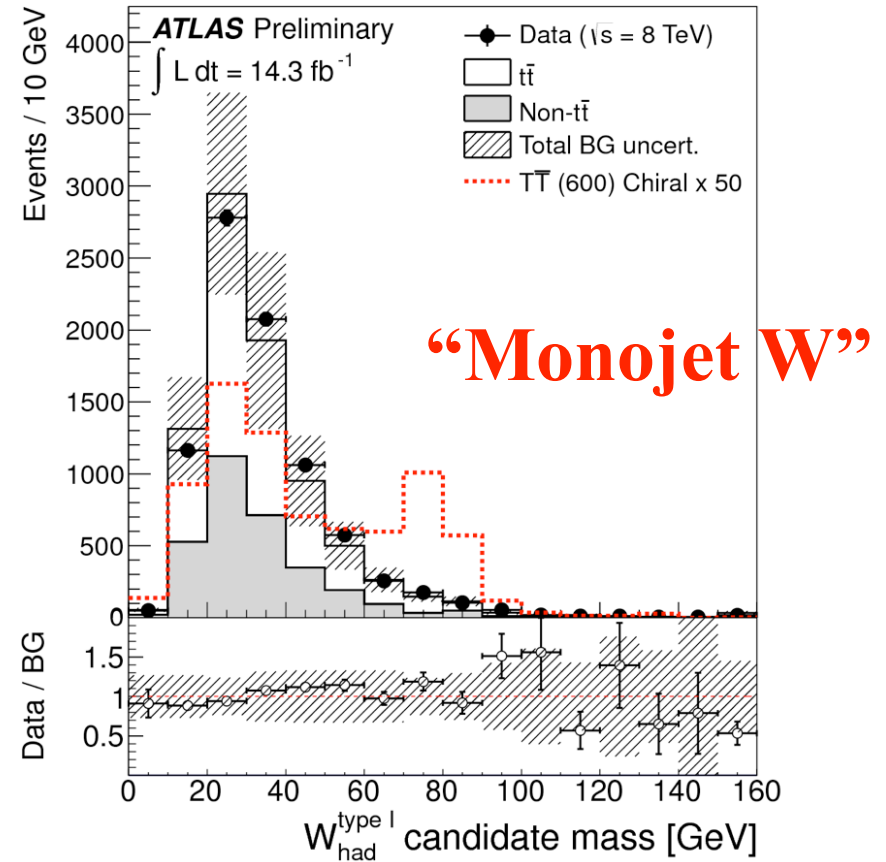
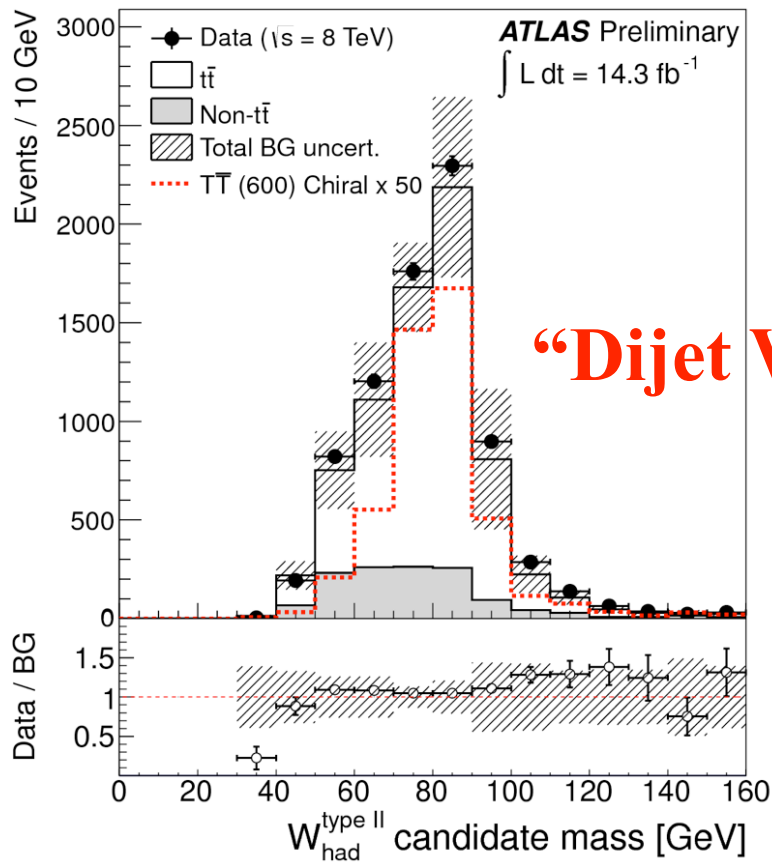
- Just no huge MET
- At least not systematically...



# W's Can Be Light

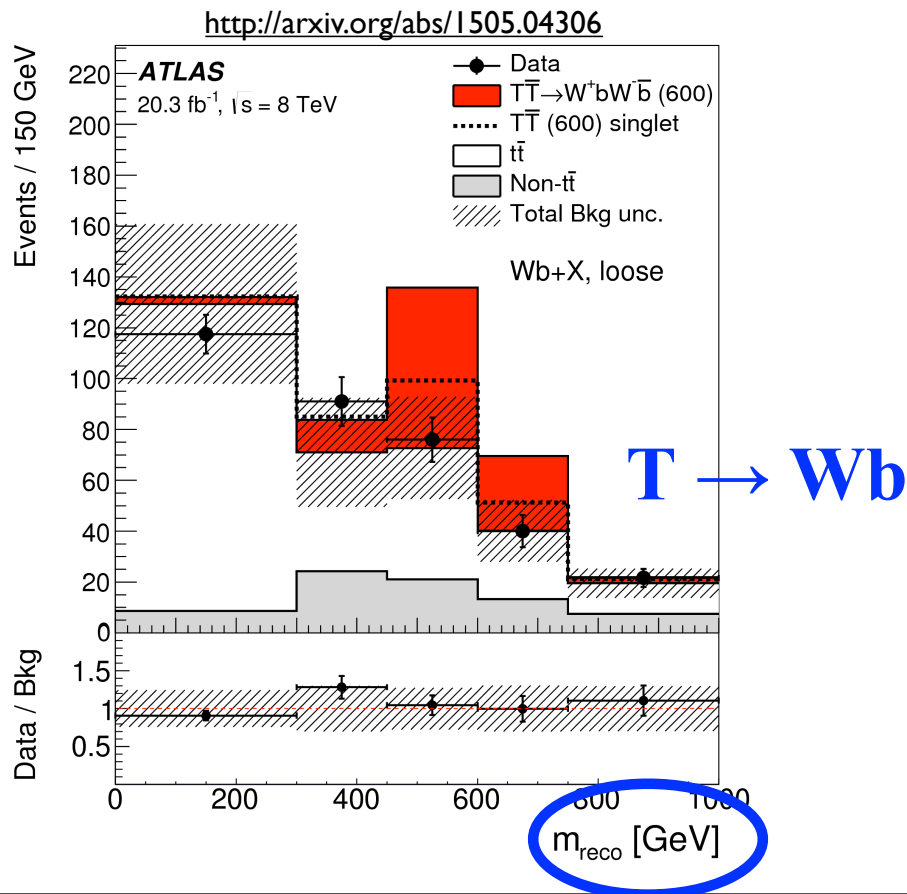
❖  $T \rightarrow Wb$  with  $m_T \sim 600$  GeV

➔  $W$  will be boosted, and if decays hadronically  $\rightarrow$  single jet



# Wb versus Ht

- ❖  $T \rightarrow Wb$  yields the same final state as  $t \rightarrow Wb$ 
  - Need to discriminate, e.g. reconstruct  $m_T$

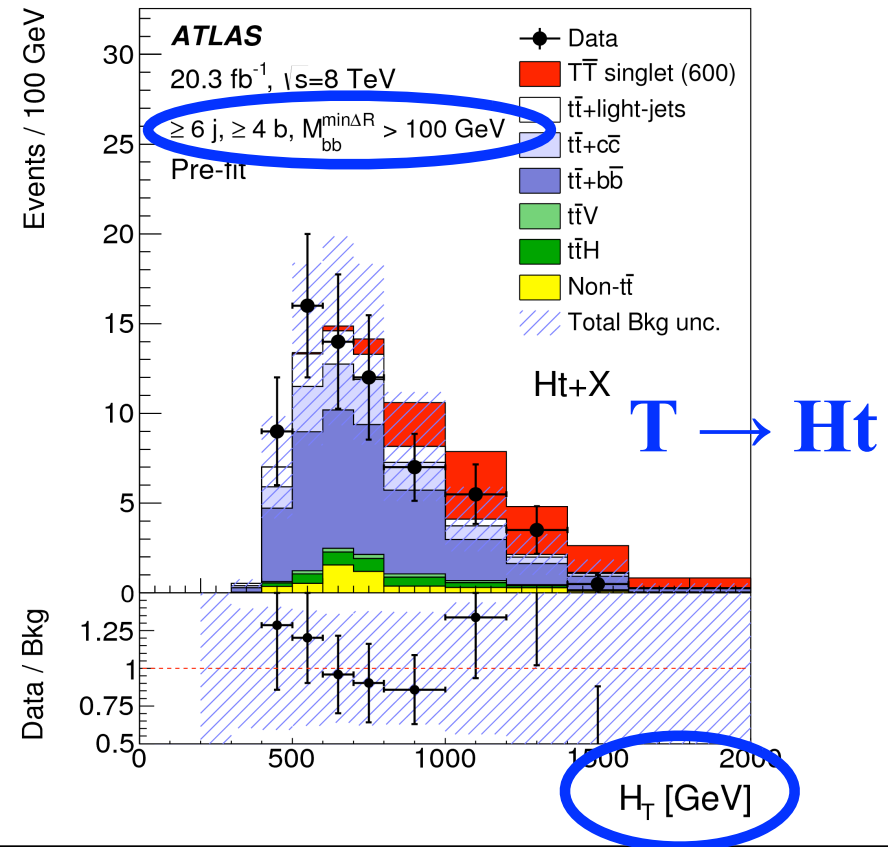
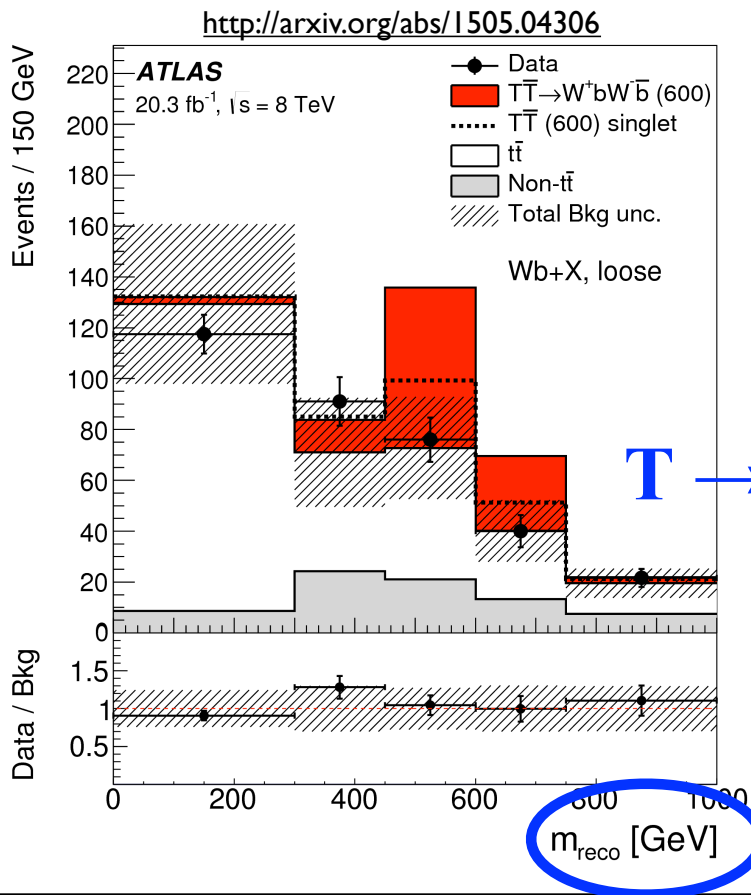


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❖  $T \rightarrow Wb$  yields the same final state as  $t \rightarrow Wb$

- Need to discriminate, e.g. reconstruct  $m_T$

❖  $T \rightarrow Ht$ :  $t\bar{t}HH$ , so  $WWbbbb$

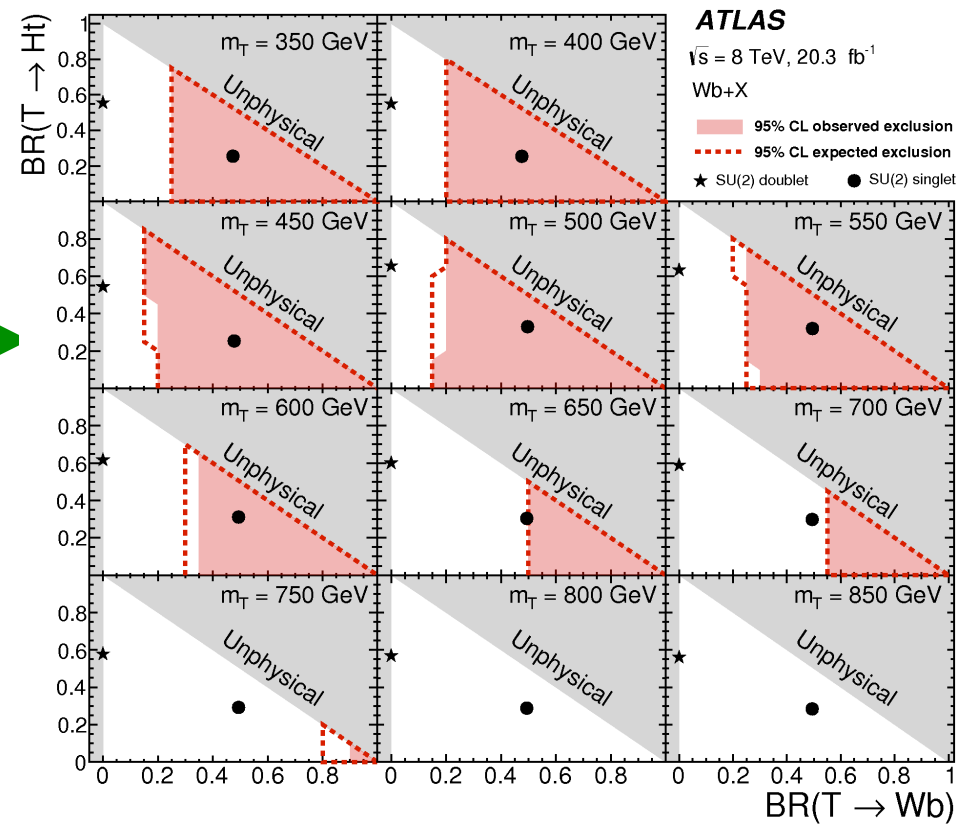
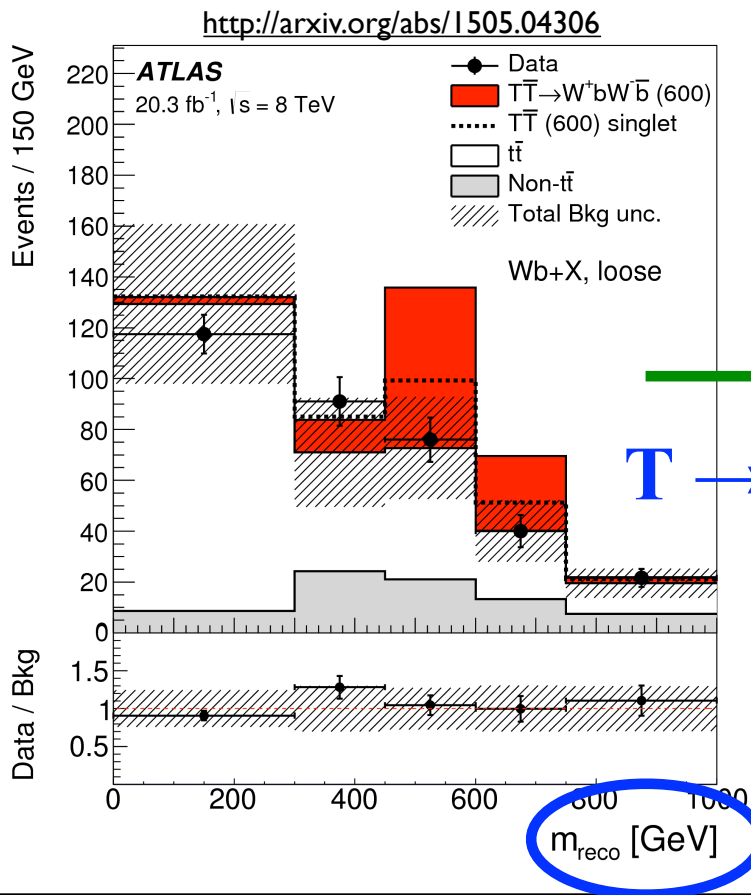


# Wb versus Ht

❖  $T \rightarrow Wb$  yields the same final state as  $t \rightarrow Wb$

- Need to discriminate, e.g. reconstruct  $m_T$

❖  $T \rightarrow Ht$ :  $t\bar{t}HH$ , so  $WWb\bar{b}\bar{b}\bar{b}$

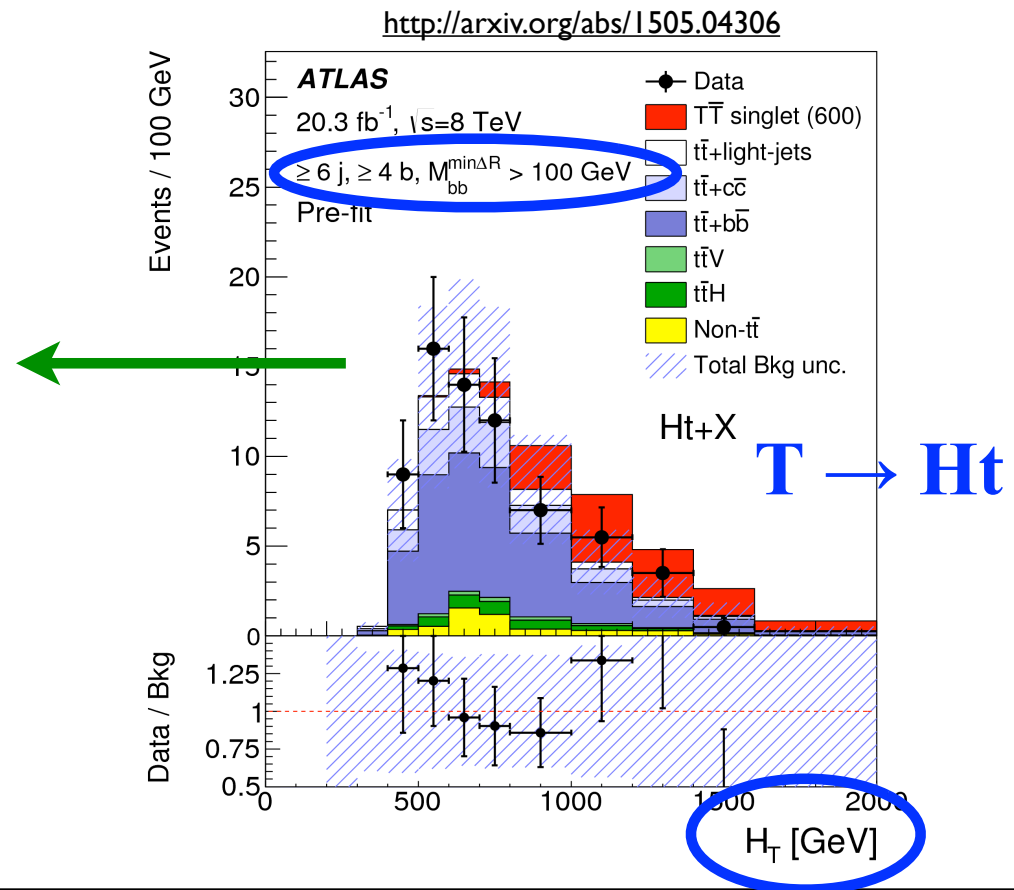
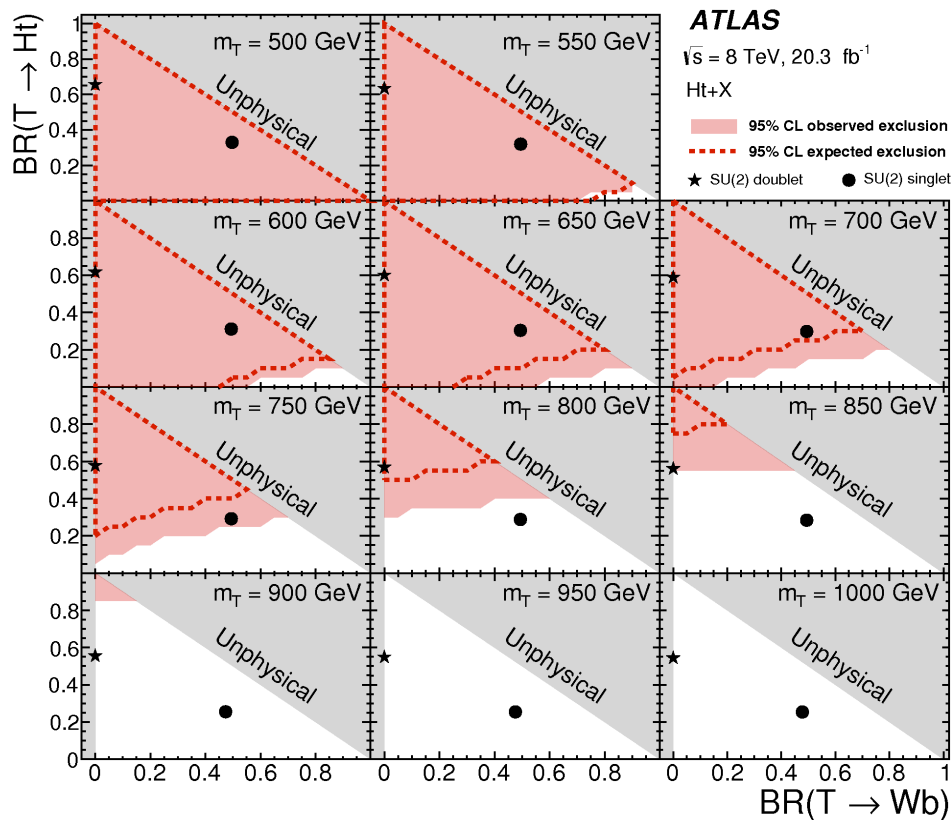


# Wb versus Ht

❖  $T \rightarrow Wb$  yields the same final state as  $t \rightarrow Wb$

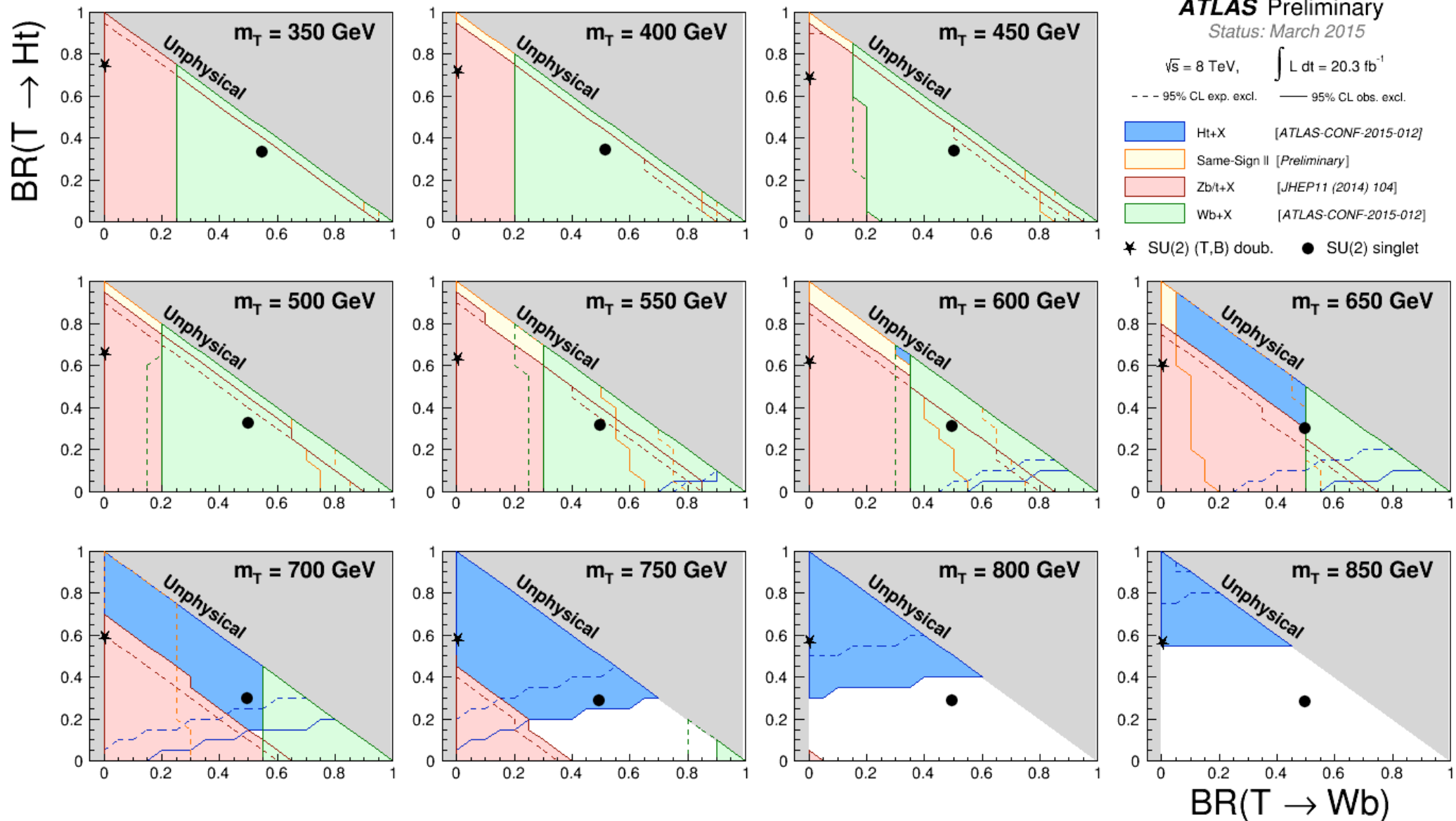
- Need to discriminate, e.g. reconstruct  $m_T$

❖  $T \rightarrow Ht$ :  $ttHH$ , so  $WWbbbb$

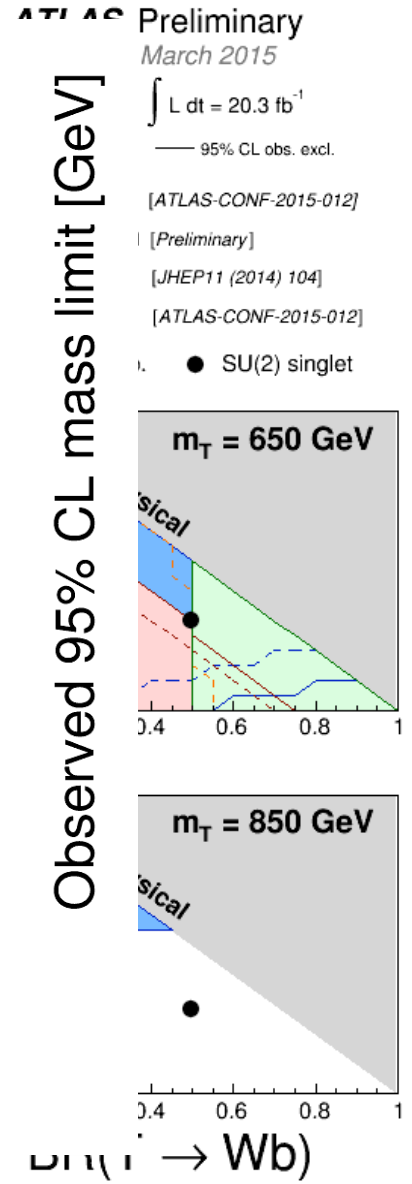
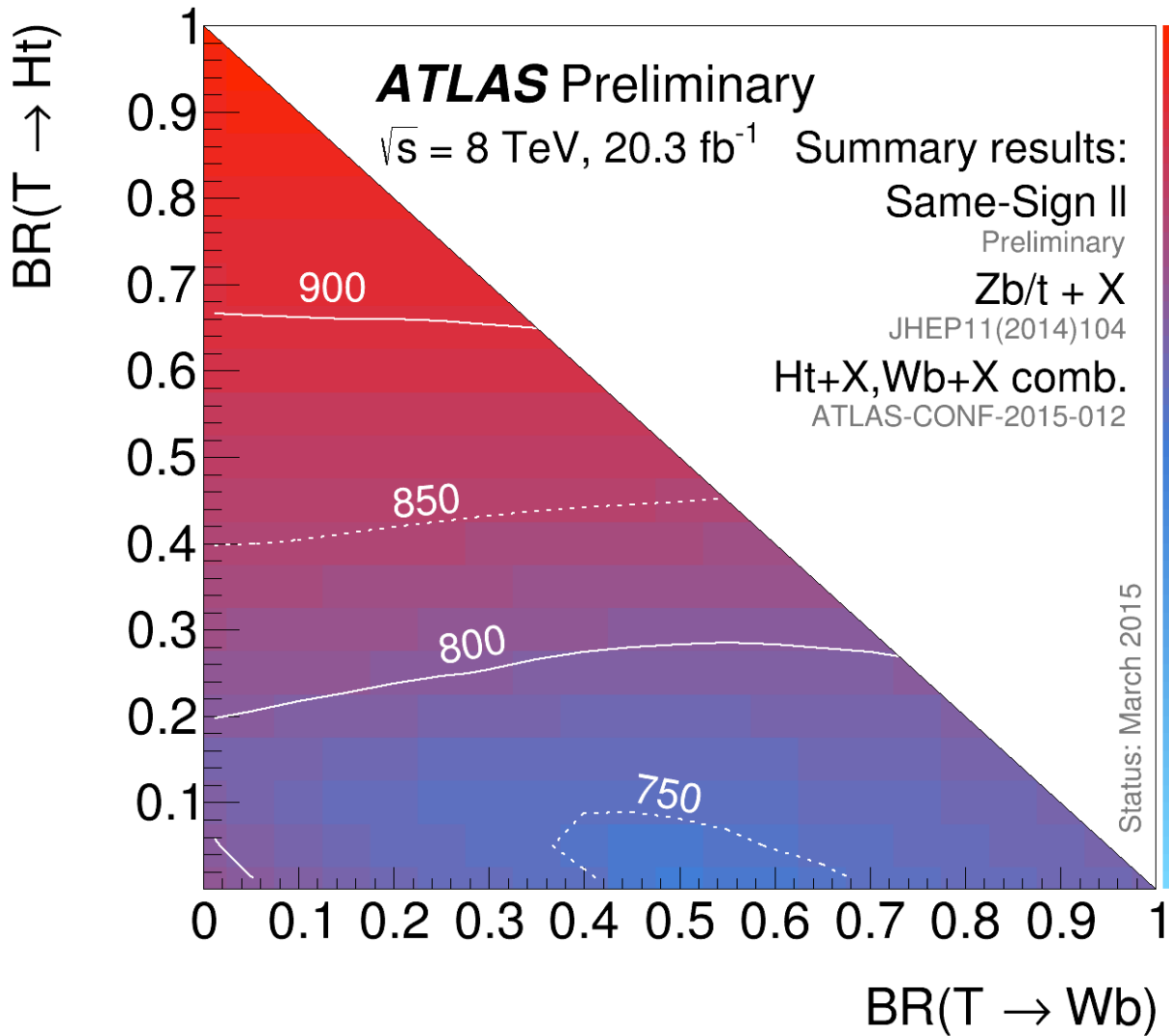
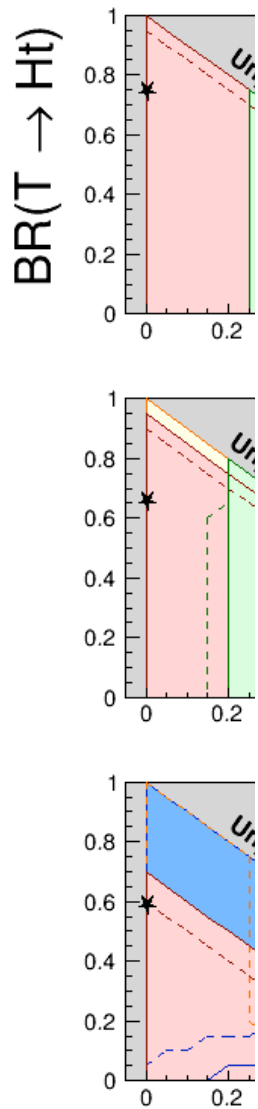




# All Together Now



# Presented Differently



# Anecdotes From the Field (II)

- ❖ It turns out these targeted searches maybe weren't useful (or all that well optimized)
  - “In particular, we find that missing energy searches designed for superparticle production, supply superior sensitivity for vector-like quarks than the dedicated new quark searches themselves”
    - Biekötter, Hewett et al. [arXiv:1608.01312](https://arxiv.org/abs/1608.01312)
  - “Owing to the much higher production cross-sections of heavy top quarks as compared to stops, masses up to  $m_T \approx 850$  GeV can be excluded from the Run 1 stop searches”
    - Kraml et al. [arXiv:1607.02050](https://arxiv.org/abs/1607.02050)
- ❖ The moral of the story...

# Also

## ❖ Somewhat simplistic: other new particles?

$T, B$	$qH$	$ql^{+l^{-}}$ $Z, LQ$	$q E_T^{\text{miss}}$ $Z, LQ$	$ql^{+\nu}$ $W, LQ$	$qqq$ $Z/W$	$qW^{+}W^{-}$ $H, Q$	$qZH$ $Q$	$qHH$ $Q$	$qW^{+}Z$ $Q$	$qW^{+}H$ $Q$
$qH$	D	-								
$ql^{+l^{-}}$	A	A	-							
$q E_T^{\text{miss}}$	C	A	C	-						
$ql^{-\nu}$	D	A	C	A	-					
$qqq$	E	A	B/C	A/B	A/E	-				
$qW^{+}W^{-}$	B	A	B	A	A/B	A	-			
$qZH$	A	A	A	A	A	A	A	-		
$qHH$	D	A	C	B	B/C	B	A	C	-	
$qW^{-}Z$	A	A	A	A	A	A	A	A	A	-
$qW^{-}H$	D	A	C	A	A/C	A	A	C	A	B

Table 1: A:  $\geq 2$  leptons ( $Z, \geq 3 W, l + W$ ). B: 1 lepton ( $2W$ ). C:  $E_T^{\text{miss}}$ . D:  $VH, HH, tH$ . E:  $W/H/t$ -jets, all jets

$X, Y$	$ql^{+\nu}$ $W, LQ$	$qqq$ $W$	$qW^{+}Z$ $Q$	$qW^{+}H$ $Q$
$ql^{-\nu}$	A	-		
$qqq$	A/B	A/E	-	
$qW^{-}Z$	A	A	A	-
$qW^{-}H$	A	D	A	B

Table 2: A:  $\geq 2$  leptons ( $Z, \geq 3 W, l + W$ ). B: 1 lepton ( $2W$ ). C:  $E_T^{\text{miss}}$ . D:  $VH, HH, tH$ . E:  $W/H/t$ -jets, all jets

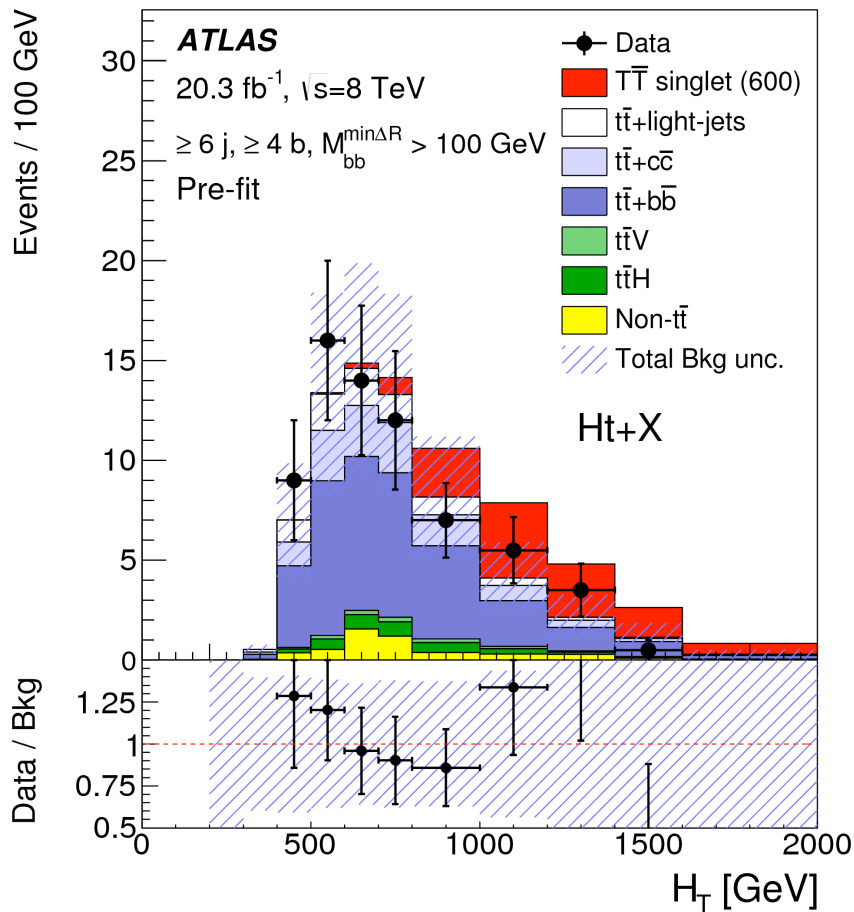
G Cacciapaglia, GB  
in Les Houches 2015

# Uncertainties

# Systematic Uncertainties

- ❖ Statistical uncertainties are easy: with limited number of events (and experiments), precision on a measurement is limited
- ❖ Systematic uncertainties vastly more complex
  - Example: measure a cross-section:  $\sigma = \frac{N_{\text{events}}}{LA\epsilon}$
  - L is the integrated luminosity, A the acceptance,  $\epsilon$  the efficiency
    - Statistical uncertainty comes from  $N_{\text{events}}$
    - Systematic uncertainties arise from limited knowledge of L, A and  $\epsilon$ 
      - ▶ L is estimated from Van der Meer scans
      - ▶ A typically depends on parton distribution functions
      - ▶ efficiency is a convolution of many experimental uncertainties

# Example



❖  $H_T$  is the sum of scalar energies of jets, leptons,...

- If the jet energy scale is different between data and MC, comparison is wrong
- If the jet energy scale dependence on jet energy is wrong, distort shape
- etc.

❖ But how do I determine the jet energy scale uncertainty?

- testbeams (single pions)
- dijet balance
- $\gamma/Z$ +jet balance
- ...

❖ What about the MC generation itself?

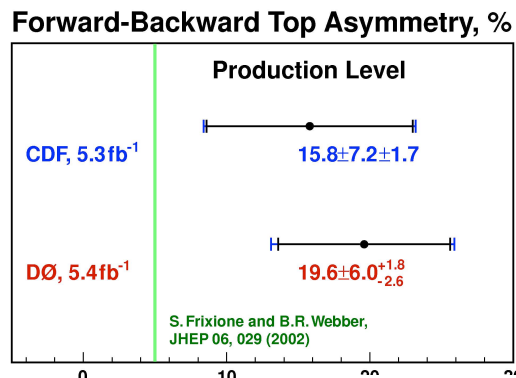
# Simulation



# Anecdotes From the Field (III)

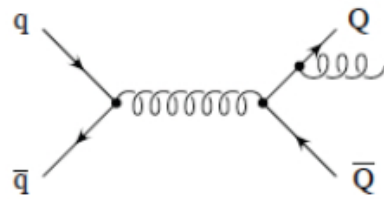
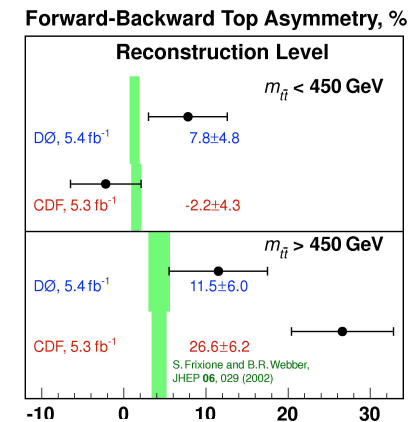
## ❖ $t\bar{t}$ charge asymmetry at the Tevatron

- At Feynman diagram level, NLO effect (Tevatron is proton-anti-proton collider)

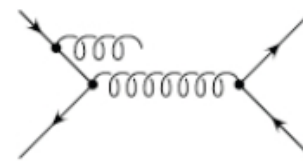


<http://arxiv.org/abs/1107.4995>

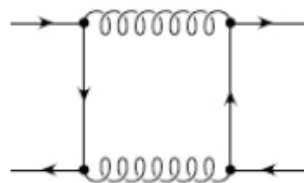
Ca. 2010, big fuss:  
much larger than SM!



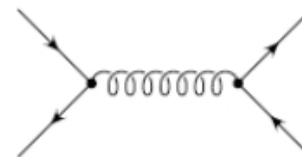
(a) Final state gluon bremsstrahlung FSR



(b) Initial state gluon bremsstrahlung ISR



(c) double virtual gluon exchange

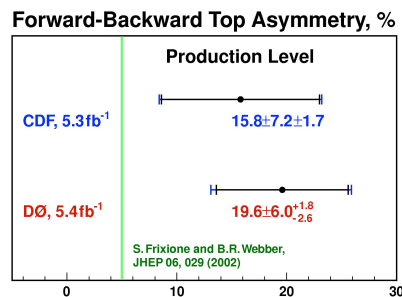


(d) Born diagram

# Anecdotes From the Field (III)

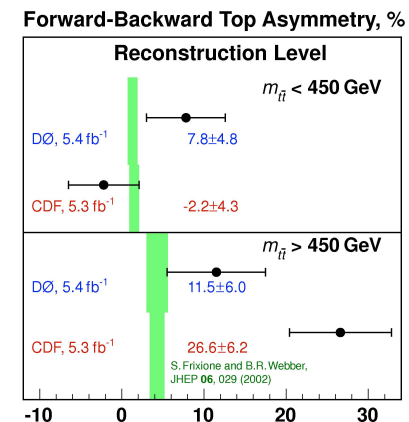
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- At Feynman diagram level, NLO effect (Tevatron is proton-anti-proton collider)



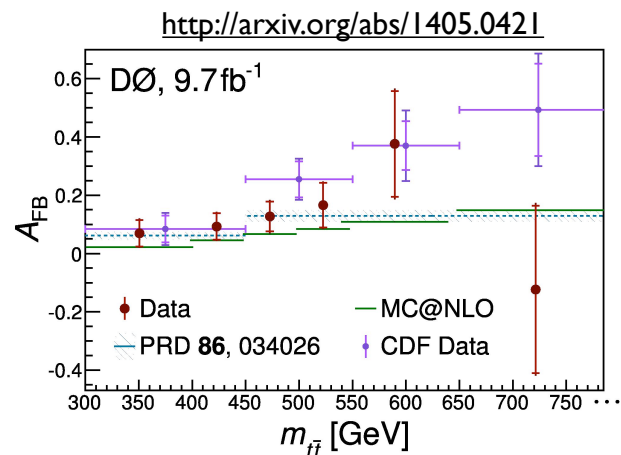
<http://arxiv.org/abs/1107.4995>

Ca. 2010, big fuss:  
much larger than SM!



- In real life, already exists at ~LO!

- Shown it is there in Pythia: parton shower, recoils! Skands et al: <http://arxiv.org/abs/1205.1466>



- Many scales in events Brodsky, Wu <https://arxiv.org/abs/1205.1232>

(My) current conclusion:  
no BSM physics here: just reality  
vs Feynman diagrams

# About Generators...

- ❖ We use four kinds of Monte Carlo generators
  - “Calculators” (often NNLO) do not actually generate events, they just calculate some (limited) distributions, like  $W p_T$
  - Traditional  $2 \rightarrow 2$  generators: LO, e.g.  $q\bar{q} \rightarrow WZ$ 
    - Include parton shower, i.e. QCD radiation, and hadronization to jets
    - pythia and herwig
  - “Matrix Element”  $2 \rightarrow n$  ( $n < 9$ ): LO, e.g.  $q\bar{q} \rightarrow evjjjj$ 
    - Necessary to generate events with multiple hard jets
    - Require matching to parton shower to avoid double counting
  - NLOwPS  $2 \rightarrow n$  generators: include NLO corrections
    - I.e. in a sense they are  $2 \rightarrow n$  with virtual corrections

# Correction Factors

- ❖ Of course, always limitations, so “K-factors” needed
  - Different ones for heavy flavor etc..... (DØ) convention to avoid confusion....
  - **K-factor is purely theoretical, and denotes a (N)NLO/LO ratio of cross sections;**
  - **K'-factor is also theoretical, and denotes a (N)NLO/LL ratio of cross sections.**  
According to Steve, ALPGEN cross sections are **Leading Log**;
  - **S-factor is empirical, and comes on top of K or K'** to bring MC in agreement with data. MC should be initially normalized to luminosity, and all correction (a.k.a. scale) factors should be applied (trigger, ID...);
  - **HF-factor is, in principle, theoretical, but in practice only theory inspired.**  
It tells you by how much heavy flavor production should be increased, on top of K or K', and possibly S;
  - **S\_HF-factor is empirical, and comes on top of K or K', S, and HF, to bring MC in agreement with data, after b-tagging.**

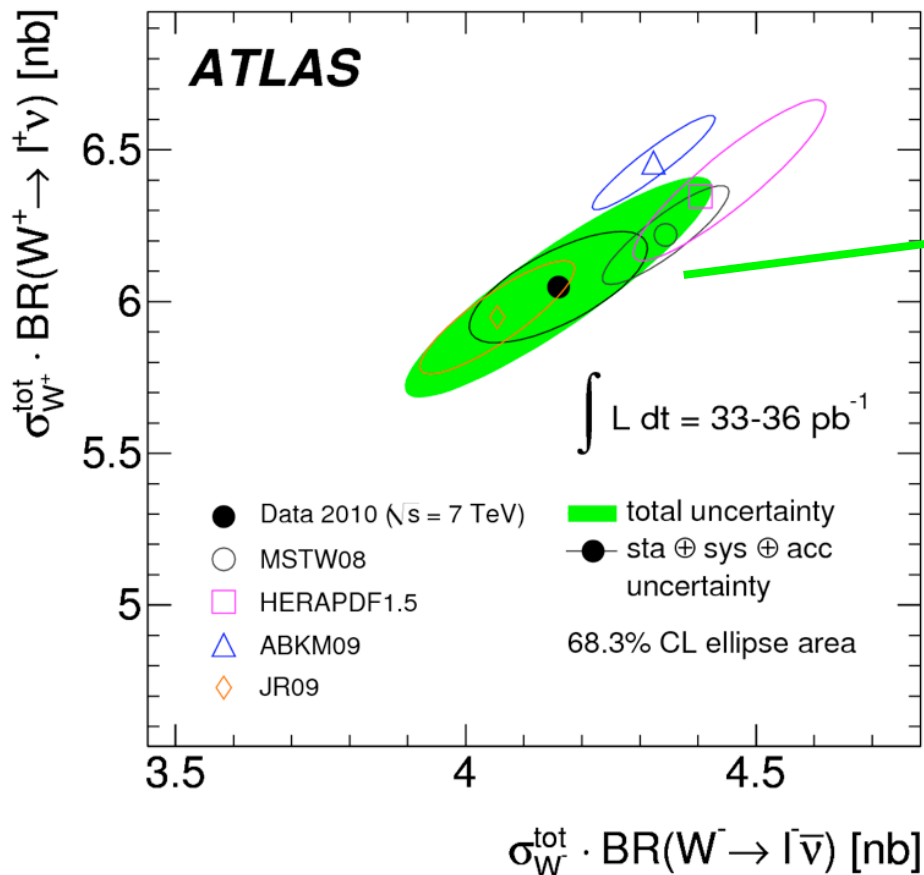
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In addition to WIZ-PT reweighting

# Sometimes Physics Helps

- ❖ At the LHC, produce more  $W^+$  than  $W^-$ 
  - Can exploit that to normalize  $W^+$  jets

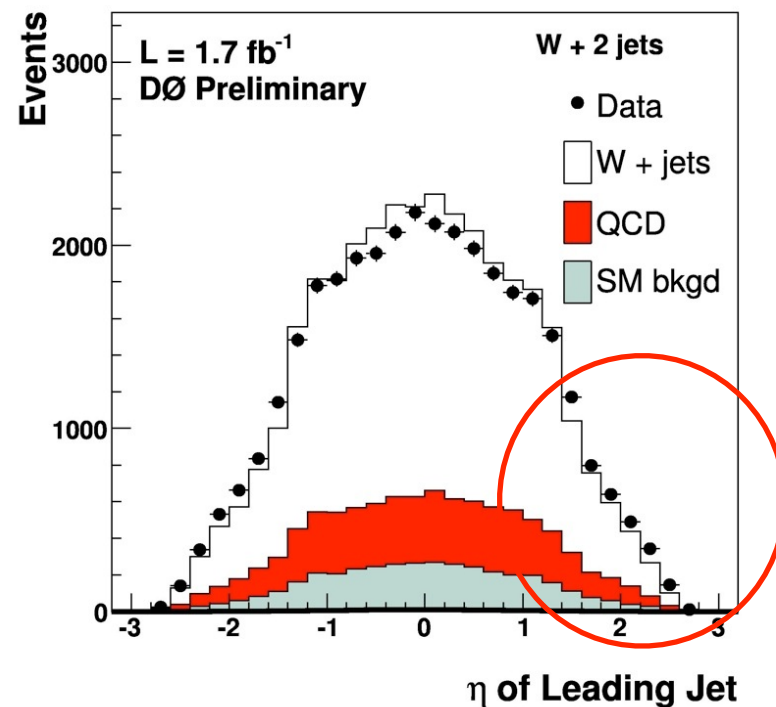


$$(N_{W^+} + N_{W^-})^{\text{exp}} = \left( \frac{r_{\text{MC}} + 1}{r_{\text{MC}} - 1} \right) (N_{W^+} - N_{W^-})^{\text{data}}$$

But what about shape??

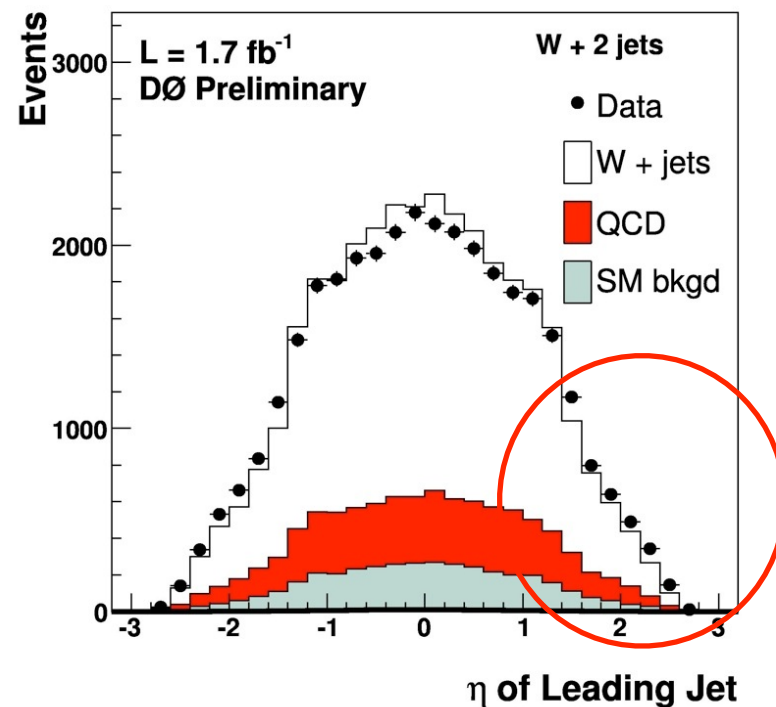
# Anecdotes From the Field (IV)

- ❖ Pile-up events (“minimum bias”) do produce jets



# Anecdotes From the Field (IV)

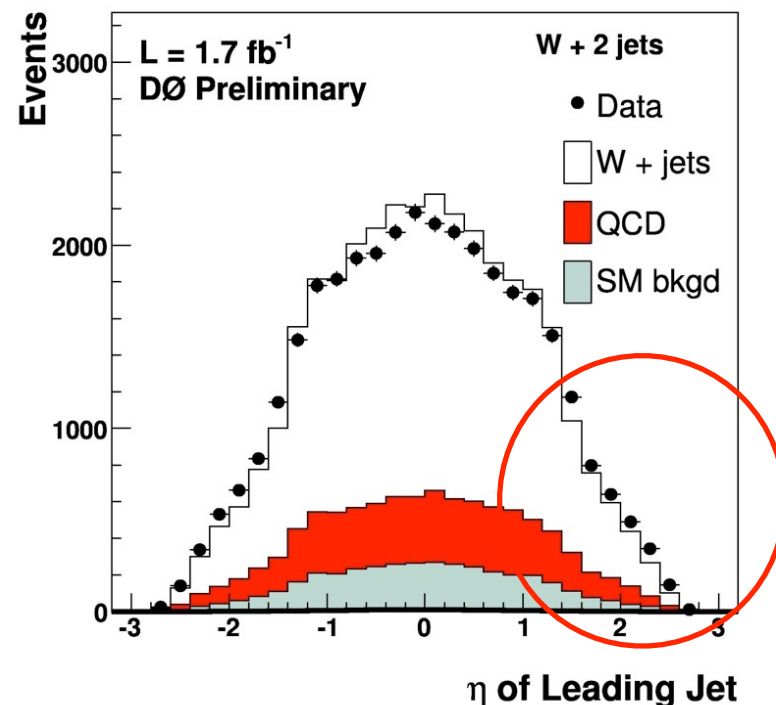
- ❖ Pile-up events (“minimum bias”) do produce jets
  - At high  $\mathcal{L}$ , require that tracks pointing to jets originate from same vertex as lepton
  - High  $\eta$  excess disappeared!





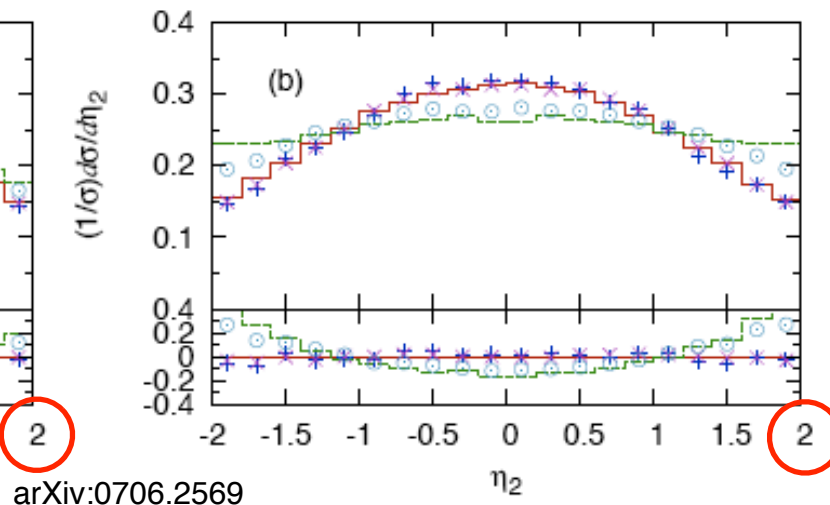
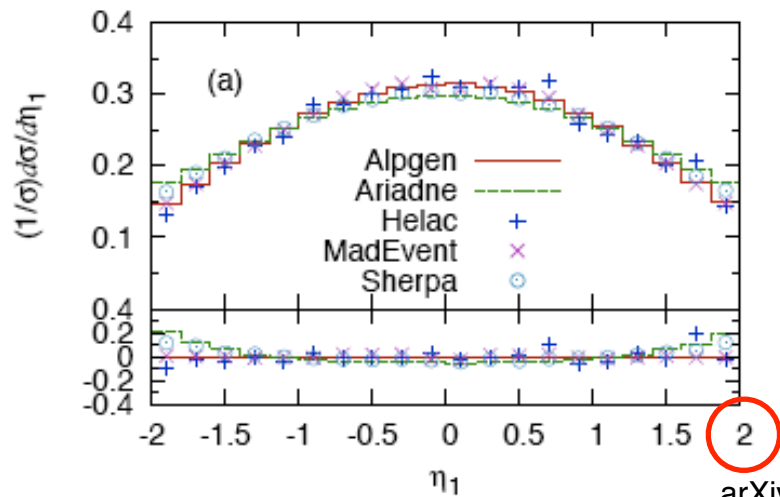
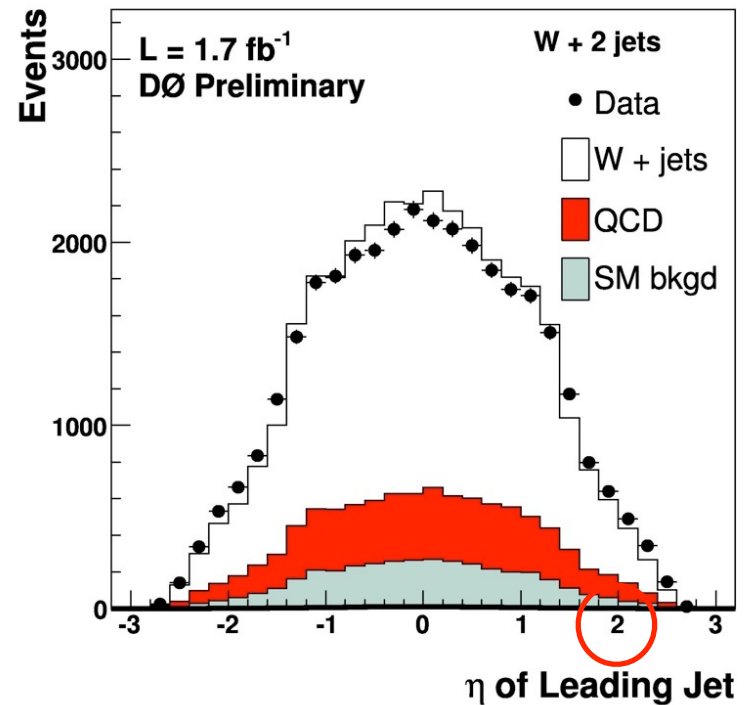
# Anecdotes From the Field (IV)

- ❖ Pile-up events (“minimum bias”) do produce jets
  - At high  $\mathcal{L}$ , require that tracks pointing to jets originate from same vertex as lepton
  - High  $\eta$  excess disappeared!
  - Eta-dependence of jet-vertex match turns out to have shape very similar to excess
  - After correcting for this, excess was back....



# So...

- ❖ After all K/K'/S/HF-factors and boson  $p_T$  reweighting:
- ❖ Similar angular differences between generators: reweigh alpgen to sherpa



Alpgen, MadEvent,  
Helac with MLM,  
Sherpa and Ariadne  
with CKKW  
matching

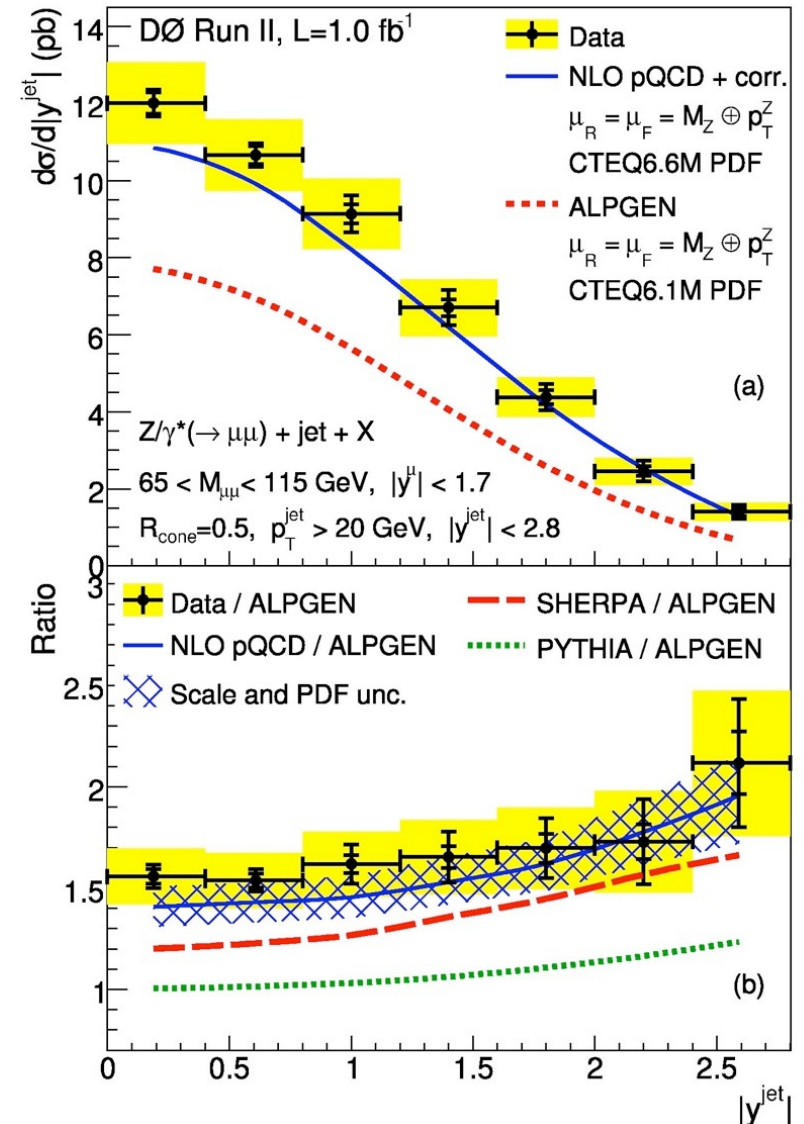
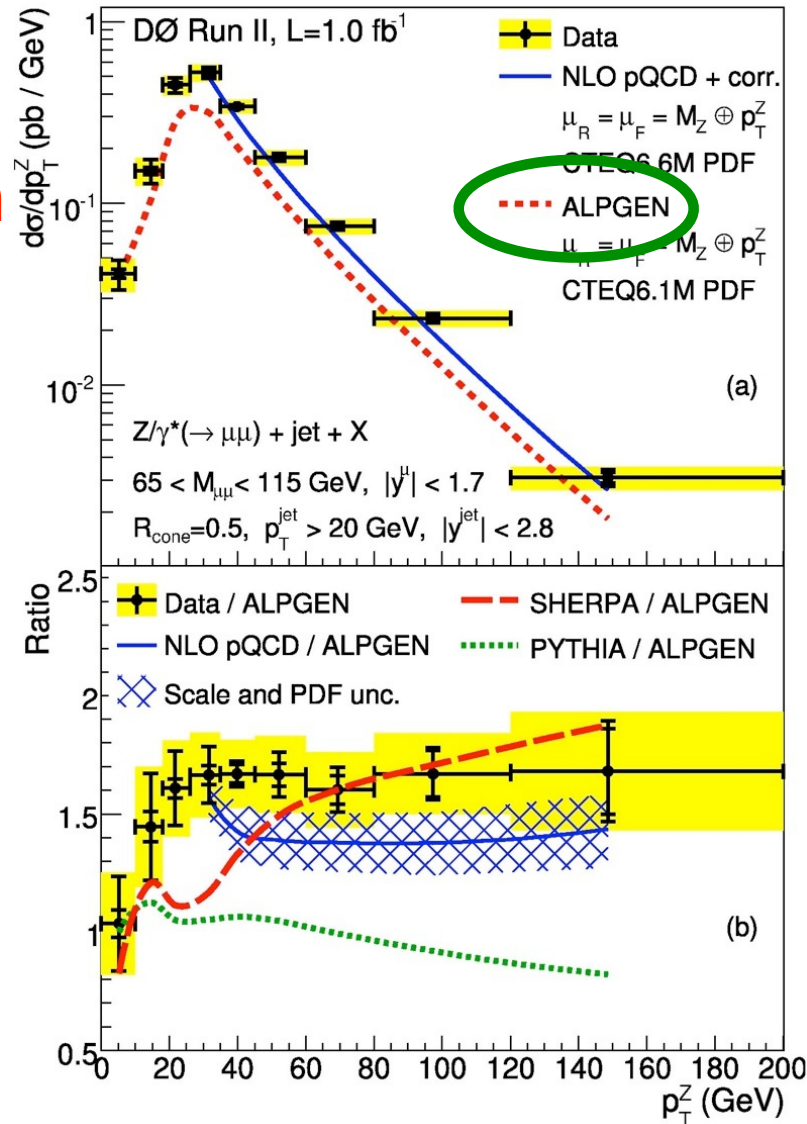
arXiv:0706.2569

# Z ( $\rightarrow ll$ ) + jets

❖ Can get a clean sample, check if our simulation reproduces the data

⇒ yes, with  
~expected deviations

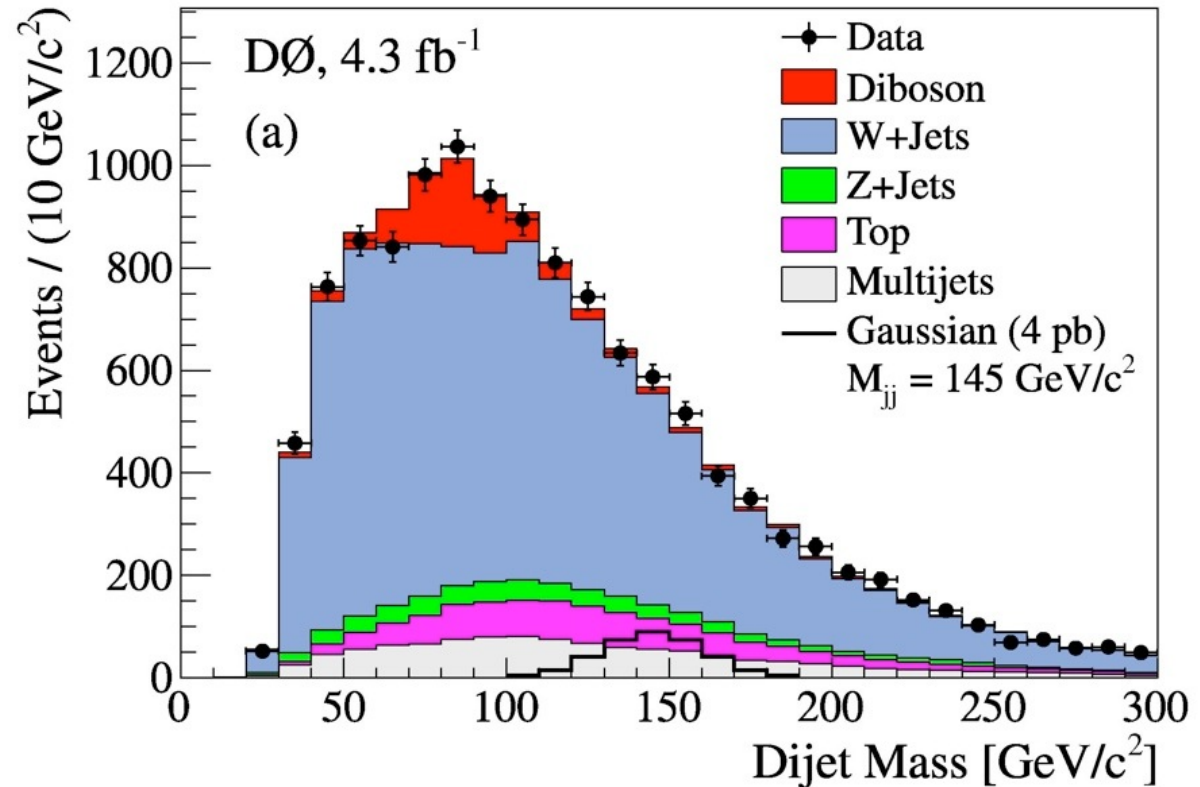
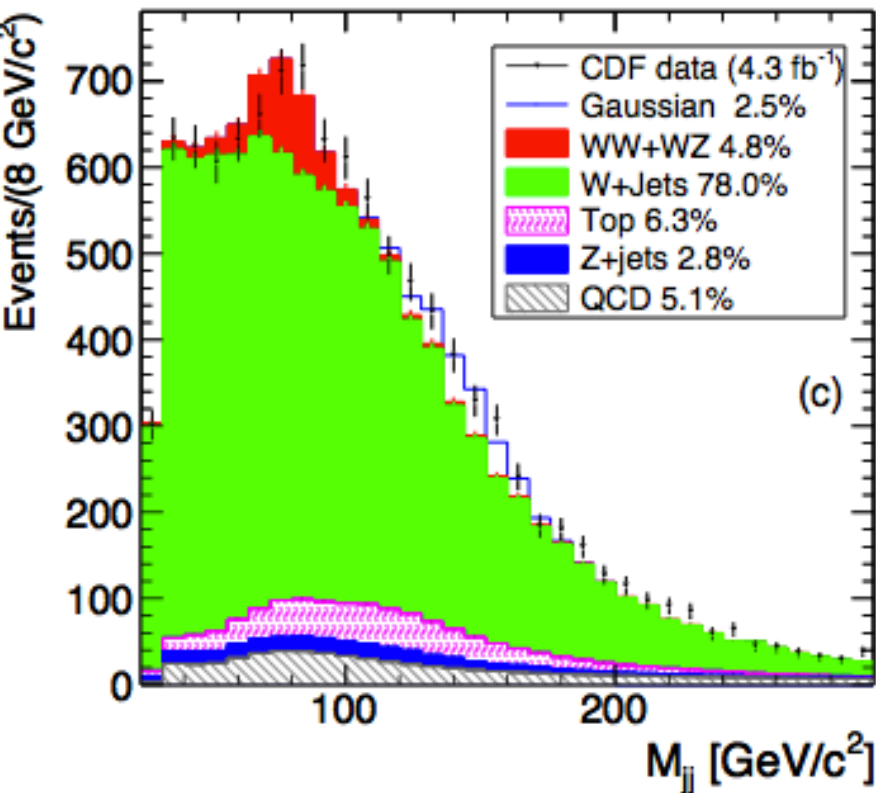
Need  
reweighting  
of MC



# Anecdotes From the Field (V)

## ❖ Searched for WW/WZ in $\ell\nu jj$

[Phys.Rev.Lett.106:171801](#)



## ❖ The background here is not SM, it is uncorrected alpgen!!

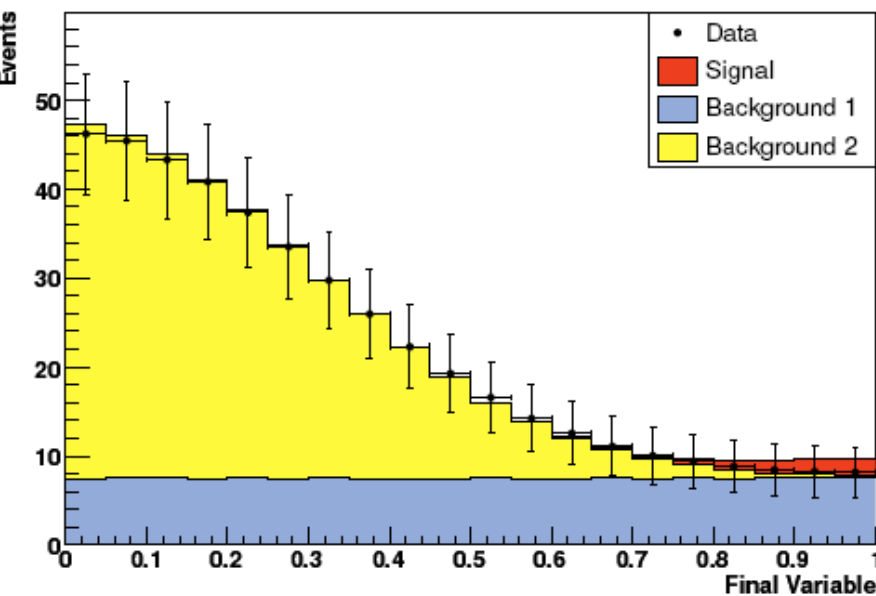
- But this is not the issue.....

# Systematics Profiling

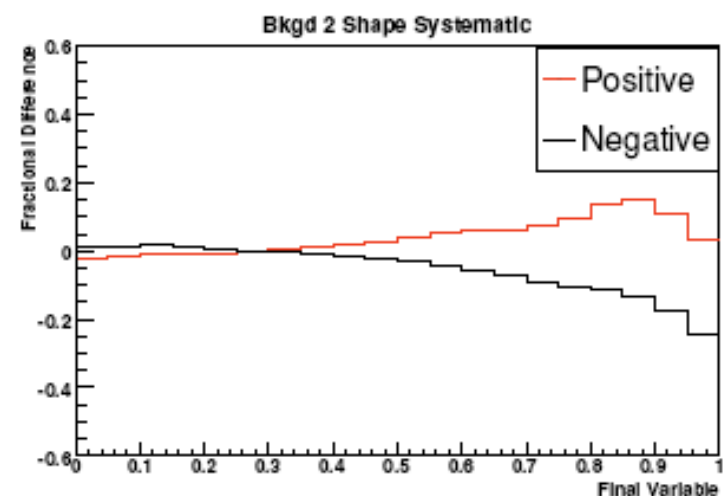
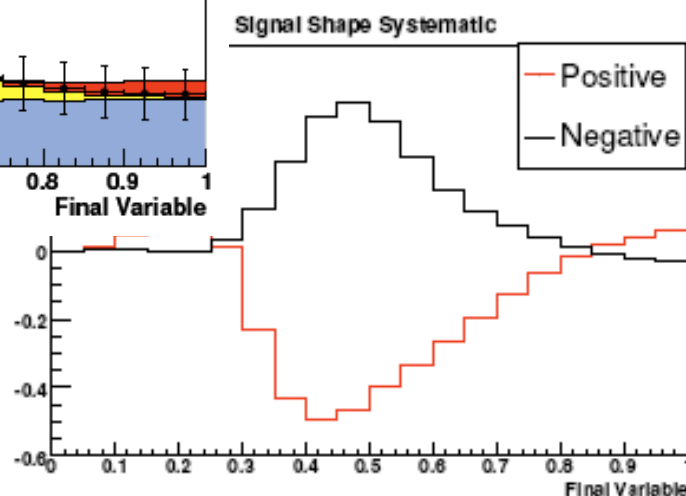
- ❖ Systematic uncertainties are propagated through the full analysis chain to the discriminating distribution
  - E.g. we repeat the analysis with jet energy scale shifted up & down by  $1\sigma$
  - Some systematic uncertainties affect shape (jet/lepton/photon reconstruction efficiency, energy scale and resolution,  $p_T$  distributions, background models), others only normalization (lepton reconstruction efficiencies and momentum calibration, background normalizations, theoretical cross-sections and luminosity)
  - Systematic uncertainties are treated as nuisance parameters when fitting signal+background to the data
    - I.e. modify signal and background shape
    - Can be fixed, or allowed to change

# Systematics Profiling

- ❖ Nuisance parameters tend to be correlated, but not 100%, among backgrounds
  - Can affect rates, shapes, or both (in any distribution), and often asymmetric and non-gaussian



## Toy Example (W. Fisher)

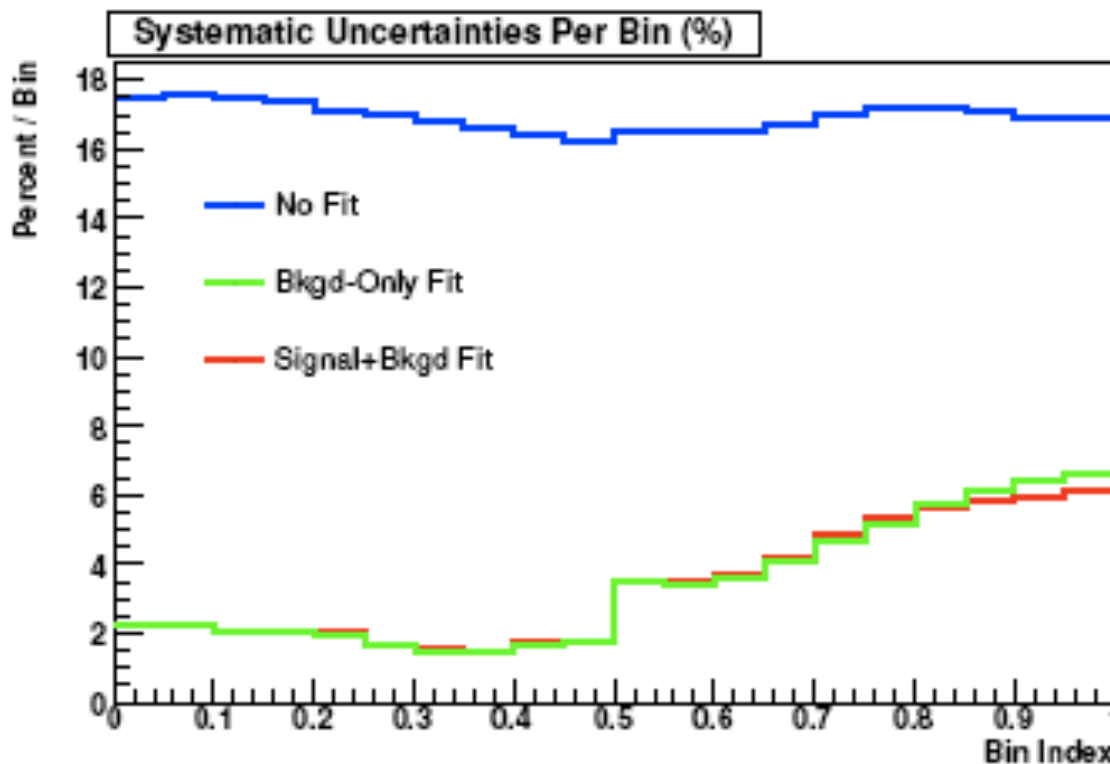




- ❖ Generate pseudo-experiments (events in bins according to poisson), then for each experiment vary nuisance parameters
  - Variations in background (& S+B) prediction
    - Compare results to data using log-likelihood ratio
- ❖ We can maximize likelihood ratio as a function of nuisance parameters → constrain them
  - I.e. use full shape of distribution(s) to see which background uncertainties are over/underestimated
    - Of course limited to size of statistical fluctuations
  - Can remove bins with large S/B if needed
    - Mostly important if uncertainties lead to similar shape distortions
  - Want enough background-rich phase space in fit!
    - Even include control regions

## ❖ Test example:

- Data constructed to disagree with background-only hypothesis (wrong estimates for background uncertainties)
- But to agree with background-only better than signal+background
- Improvement quite spectacular (by construction in example)

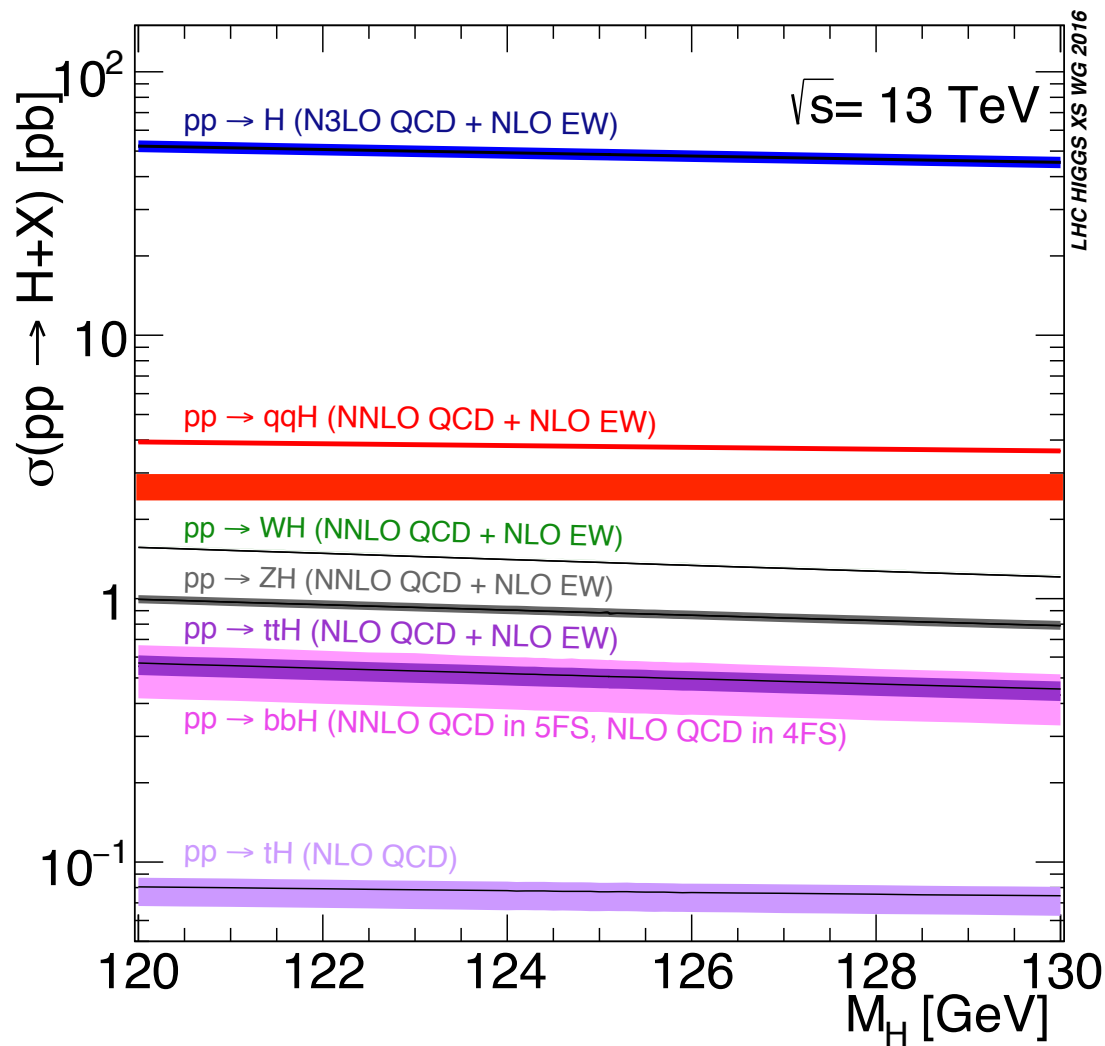




# ttH

- ❖ ttH important measurement and discovery channel
  - If  $m_H$  had been smaller, no  $H \rightarrow \gamma\gamma$ , at 14 TeV S/B better for ttH than VH
  - Direct probe of t-H coupling
- ❖ Studied pre-data and looked good, but ttbb only known at LO at that time...
- ❖ In 2009, Bredenstein, Denner, Dittmaier and Pozzorini published ttbb@NLO ([JHEP 1003 \(2010\) 021](#) and earlier)
  - ttbb twice as high as expected
  - At higher Higgs  $p_T$ , significant shape distortions
  - Would ttH be observable at all?
  - How can we understand ttbb with sufficient accuracy?

# Production

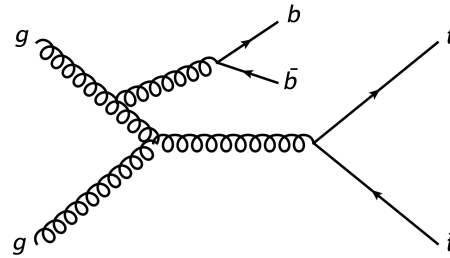


ttbb

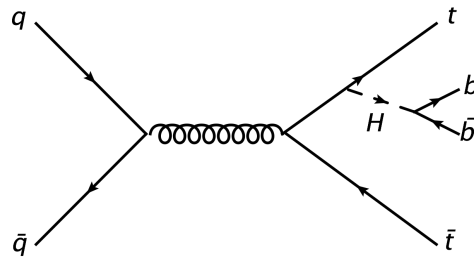
13 fb<sup>-1</sup>: ~7000 ttH events, 4200 ttH, H → bb

# Decay

Control Regions

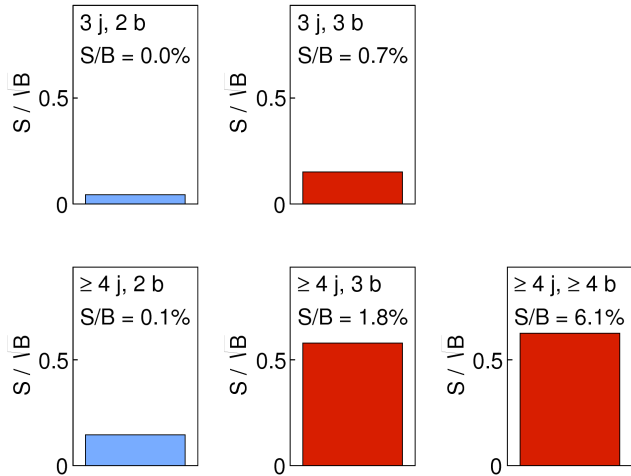


Signal Regions

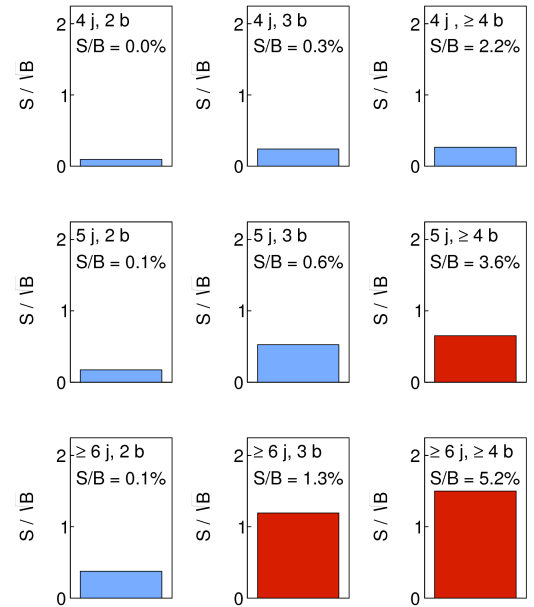


ATLAS-CONF-2016-080

ATLAS Simulation Preliminary  
 $\sqrt{s} = 13 \text{ TeV}, 13.2 \text{ fb}^{-1}$   
 Dilepton

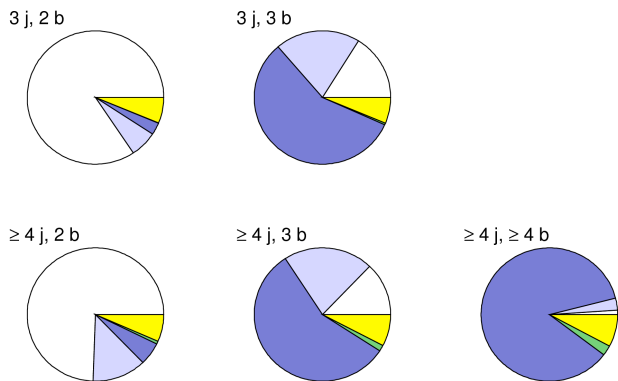


ATLAS Simulation Preliminary  
 $\sqrt{s} = 13 \text{ TeV}, 13.2 \text{ fb}^{-1}$   
 Single Lepton



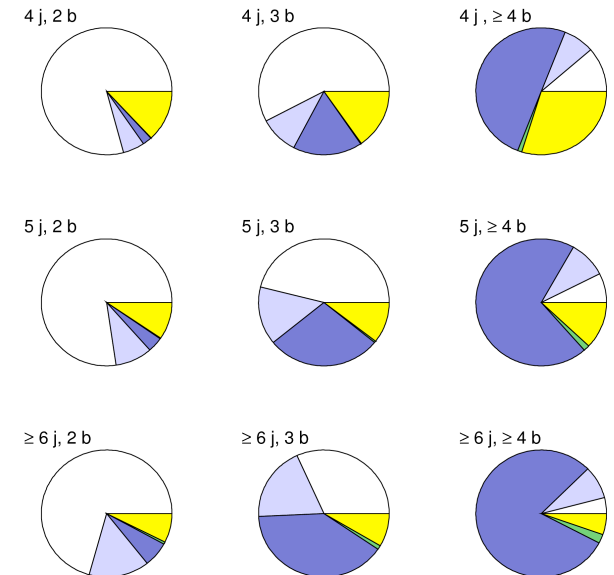
ATLAS Simulation Preliminary  
 $\sqrt{s} = 13 \text{ TeV}$   
 Dilepton

$\square$   $t\bar{t}$  + light  
 $\square$   $t\bar{t}$  +  $\geq 1c$   
 $\square$   $t\bar{t}$  +  $\geq 1b$   
 $\square$   $t\bar{t}$  + V  
 $\square$  Non- $t\bar{t}$



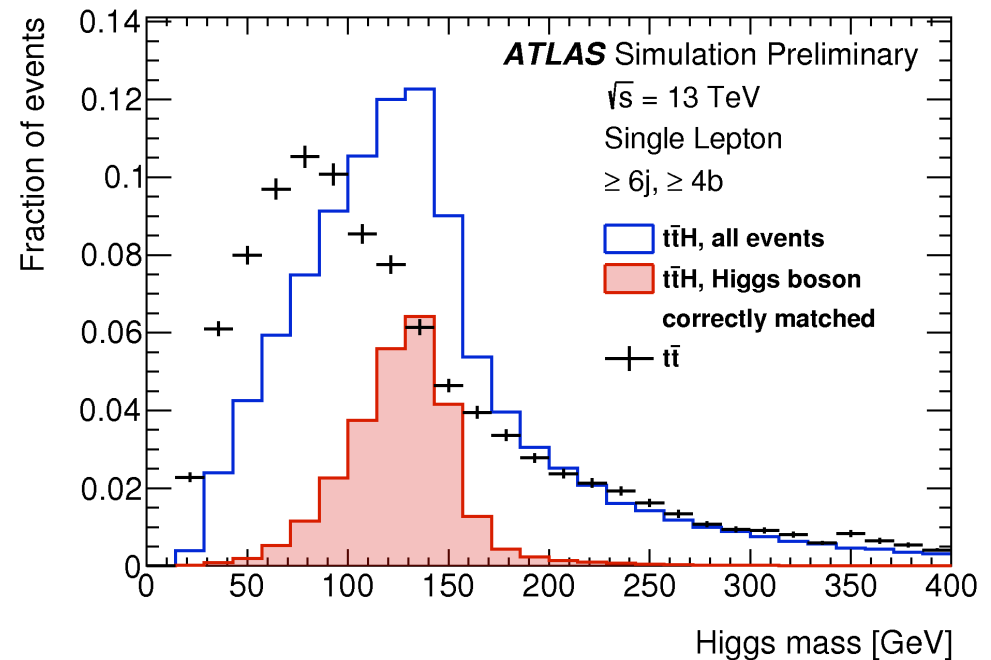
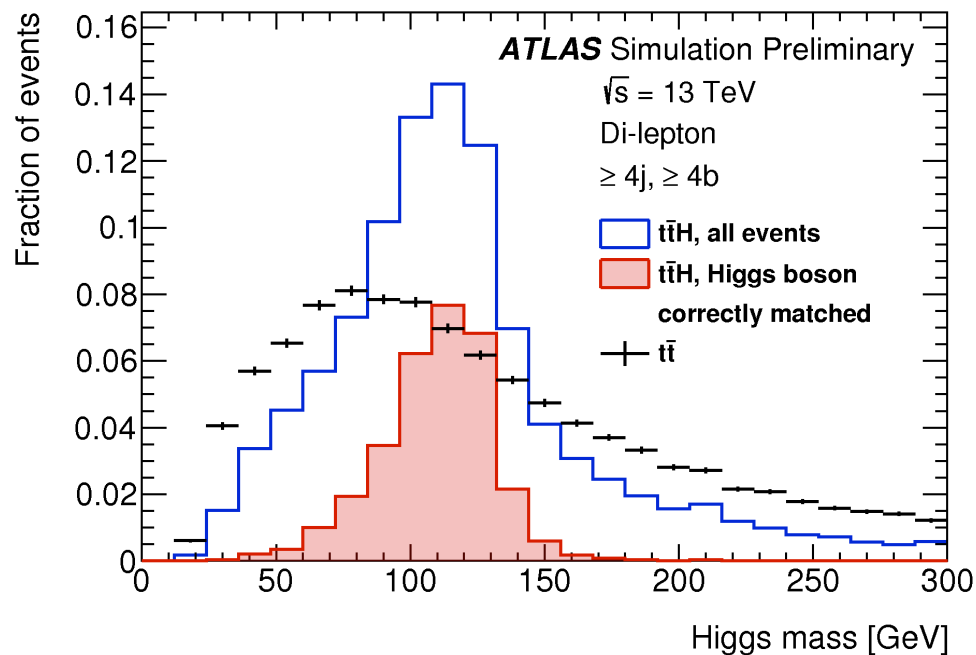
ATLAS Simulation Preliminary  
 $\sqrt{s} = 13 \text{ TeV}$   
 Single Lepton

$\square$   $t\bar{t}$  + light  
 $\square$   $t\bar{t}$  +  $\geq 1c$   
 $\square$   $t\bar{t}$  +  $\geq 1b$   
 $\square$   $t\bar{t}$  + V  
 $\square$  Non- $t\bar{t}$



# $m_{bb}$

- ❖ Clearly,  $m_{bb}$  is a good discriminating variable
  - But resolution is poor
  - Include other variables, combine using multivariate tools (NN, BDT)

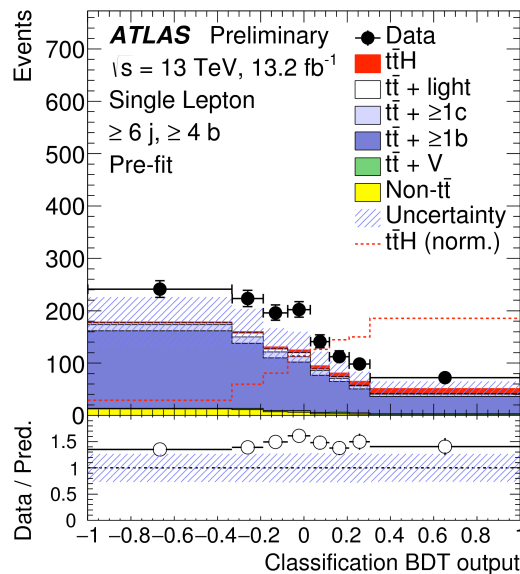
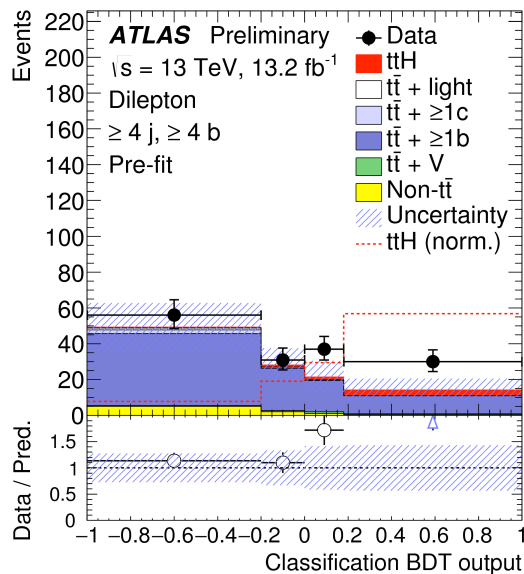
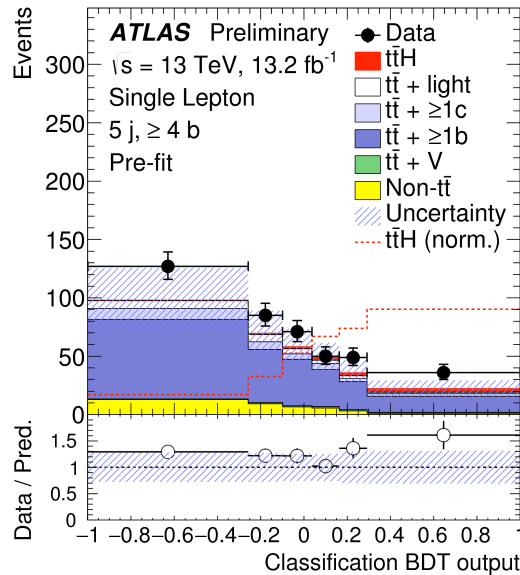
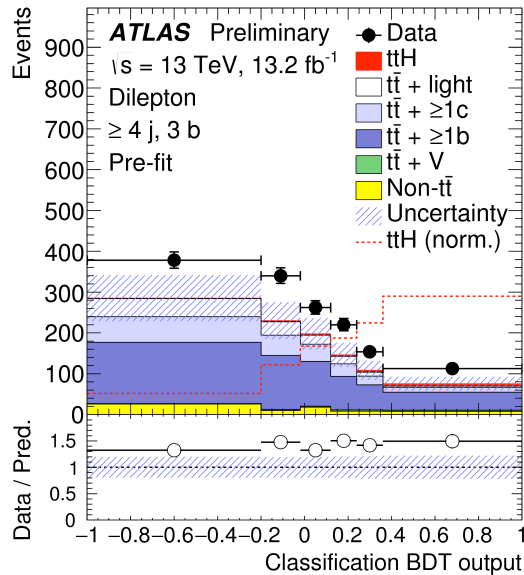


# MVA

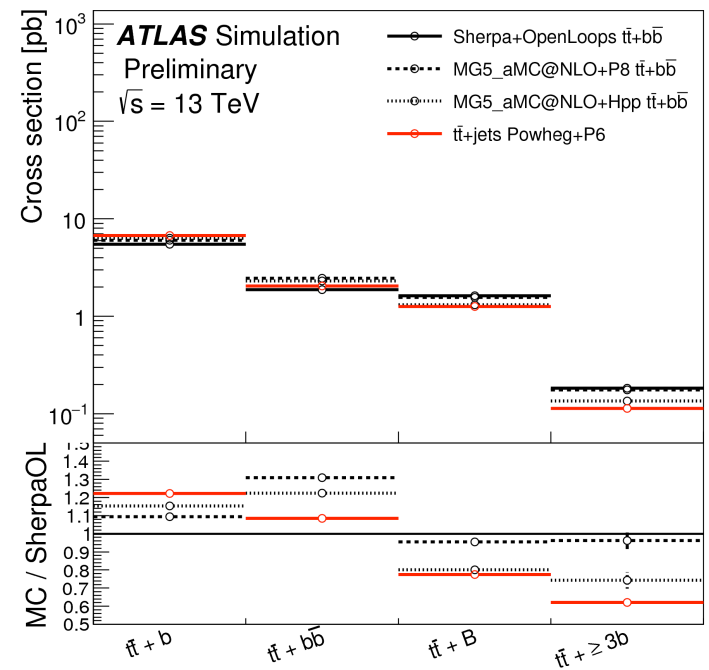
Variable	Definition	Region		
		$\geq 4j, \geq 4b$	$\geq 4j, 3b$	$3j, 3b$
General kinematic variables				
$\Delta\eta_{bb}^{avg}$	Average $ \Delta\eta $ among pairs of $b$ -jets	✓	–	–
$\Delta\eta_{bb}^{max}$	Maximum $\Delta\eta$ between any two $b$ -jets	–	✓	✓
$\Delta\eta_{jj}^{avg}$	Average $\Delta\eta$ among jet pairs	–	✓	–
$\Delta R_{bb}^{max p_T}$	$\Delta R$ between the two $b$ -tagged jets with the largest vector sum $p_T$	✓	✓	✓
$\Delta R_{bb}^{Higgs}$	$\Delta R$ between the two $b$ -tagged jets with mass closest to the Higgs boson mass	✓	–	–
$\Delta R_{bb}^{max m}$	$\Delta R$ between the two $b$ -jets with the largest invariant mass	✓	✓	✓
$m_{bb}^{max p_T}$	Mass of the two $b$ -tagged jets with the largest vector sum $p_T$	–	–	✓
$m_{bb}^{Higgs}$	Mass of the two $b$ -tagged jets closest to the Higgs boson mass	✓	✓	✓
$m_{bb}^{min}$	Minimum mass of two $b$ -tagged jets	–	–	✓
$m_{bb}^{min \Delta R}$	Mass of the combination of the two $b$ -tagged jets with the smallest $\Delta R$	✓	✓	✓
$p_{T,b}^{min}$	Minimum $b$ -tagged jet $p_T$	–	–	✓
$H_T^{all}$	Scalar $p_T$ sum of all leptons and jets	–	✓	✓
$N_{bb}^{Higgs 30}$	Number of $b$ -jet pairs with invariant mass within 30 GeV of the Higgs boson mass	✓	–	✓
$N_{jj}^{Higgs 30}$	Number of jet pairs with invariant mass within 30 GeV of the Higgs boson mass	–	✓	–
Aplanarity	$1.5\lambda_2$ , where $\lambda_2$ is the second eigenvalue of the momentum tensor [42] built with all jets	✓	✓	✓
Centrality	Sum of the $p_T$ divided by sum of the $E$ for all jets and both leptons	✓	–	✓
$H_{2jets}$	Third Fox–Wolfram moment computed using all jets	–	✓	–
$H_{4all}$	Fifth Fox–Wolfram moment computed using all jets and leptons	–	–	✓
Variables from reconstruction BDT output				
BDT output		✓*	✓*	–
$m_H$	Higgs boson mass	✓(*)	✓(*)	–
$\Delta\eta_{H,l}^{min}$	Minimum $\Delta\eta$ between the Higgs boson and a lepton	✓*	✓	–
$\Delta\eta_{H,l}^{max}$	Maximum $\Delta\eta$ between the Higgs boson and a lepton	✓*	✓	–
$\Delta\eta_{H,b}^{min}$	Minimum $\Delta\eta$ between the Higgs boson and a $b$ -jet	✓*	–	–

Variable	Definition	Region		
		$\geq 6j, \geq 4b$	$\geq 6j, 3b$	$5j, \geq 4b$
General kinematic variables				
$\Delta R_{bb}^{avg}$	Average $\Delta R$ for all $b$ -tagged jet pairs	✓	✓	✓
$\Delta R_{bb}^{max p_T}$	$\Delta R$ between the two $b$ -tagged jets with the largest vector sum $p_T$	✓	–	–
$\Delta\eta_{jj}^{max}$	Maximum $\Delta\eta$ between any two jets	✓	✓	✓
$m_{bb}^{min \Delta R}$	Mass of the combination of the two $b$ -tagged jets with the smallest $\Delta R$	✓	✓	–
$m_{jj}^{min \Delta R}$	Mass of the combination of any two jets with the smallest $\Delta R$	–	–	✓
$m_{bj}^{max p_T}$	Mass of the combination of a $b$ -tagged jet and any jet with the largest vector sum $p_T$	–	✓	–
$p_T^{jet5}$	$p_T$ of the fifth leading jet	✓	✓	✓
$N_{bb}^{Higgs 30}$	Number of $b$ -jet pairs with invariant mass within 30 GeV of the Higgs boson mass	✓	–	✓
$N_{40}^{jet}$	Number of jets with $p_T \geq 40 GeV$	–	✓	–
$H_T^{had}$	Scalar sum of jet $p_T$	–	✓	✓
$\Delta R_{lep-bb}^{min \Delta R}$	$\Delta R$ between the lepton and the combination of the two $b$ -tagged jets with the smallest $\Delta R$	–	–	✓
Aplanarity	$1.5\lambda_2$ , where $\lambda_2$ is the second eigenvalue of the momentum tensor [42] built with all jets	✓	✓	✓
Centrality	Scalar sum of the $p_T$ divided by sum of the $E$ for all jets and the lepton	✓	✓	✓
$H1$	Second Fox–Wolfram moment computed using all jets and the lepton	✓	✓	✓
Variables from reconstruction BDT output				
BDT output		✓*	✓*	✓*
$m_H$	Higgs boson mass	✓	✓	✓
$m_{H,blep top}$	Mass of Higgs boson and $b$ -jet from leptonic top	✓	–	–
$\Delta R_{Higgs bb}$	$\Delta R$ between $b$ -jets from the Higgs boson	✓	✓	✓
$\Delta R_{H,t\bar{t}}$	$\Delta R$ between Higgs boson and $t\bar{t}$ system	✓*	✓*	✓*
$\Delta R_{H,lep top}$	$\Delta R$ between Higgs boson and leptonic top	✓	–	–
$\Delta R_{H,bhad top}$	$\Delta R$ between Higgs boson and $b$ -jet from hadronic top	–	✓*	✓*

# BDT Outputs

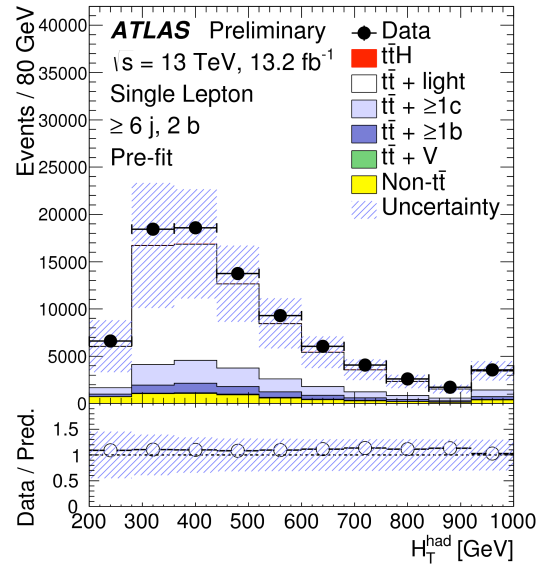
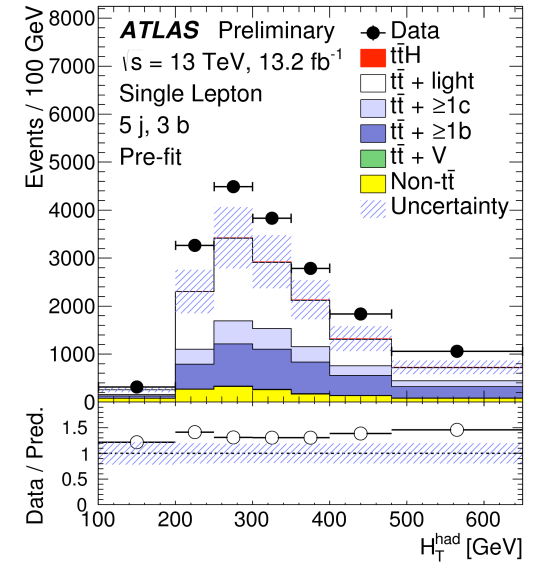
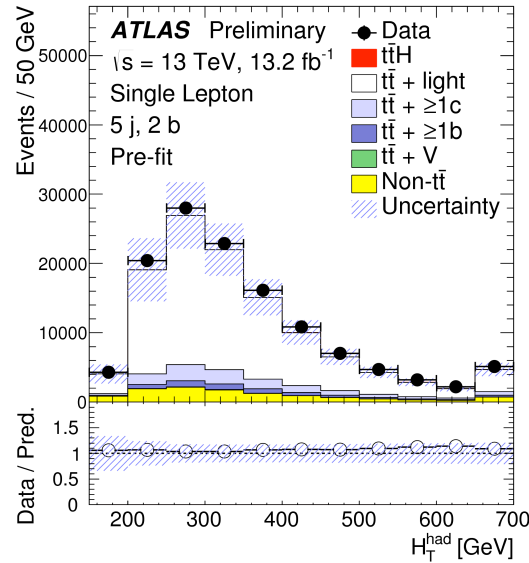
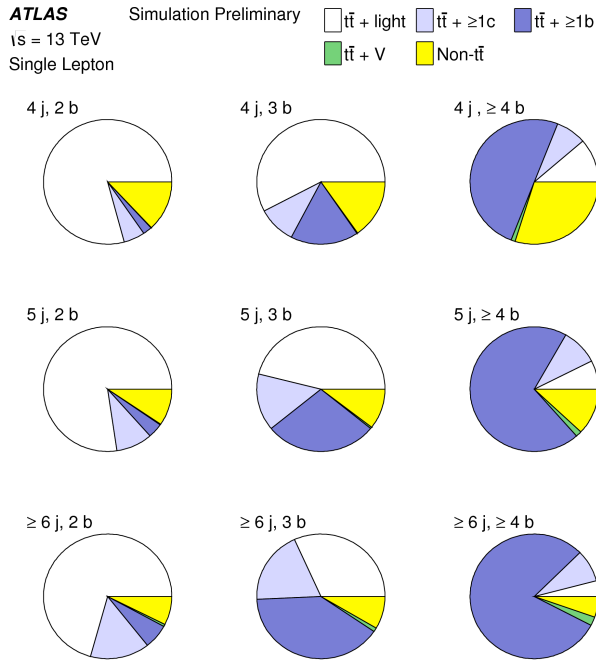


Uncertainties much larger than signal!



And then there are shape uncertainties

# Closer Look



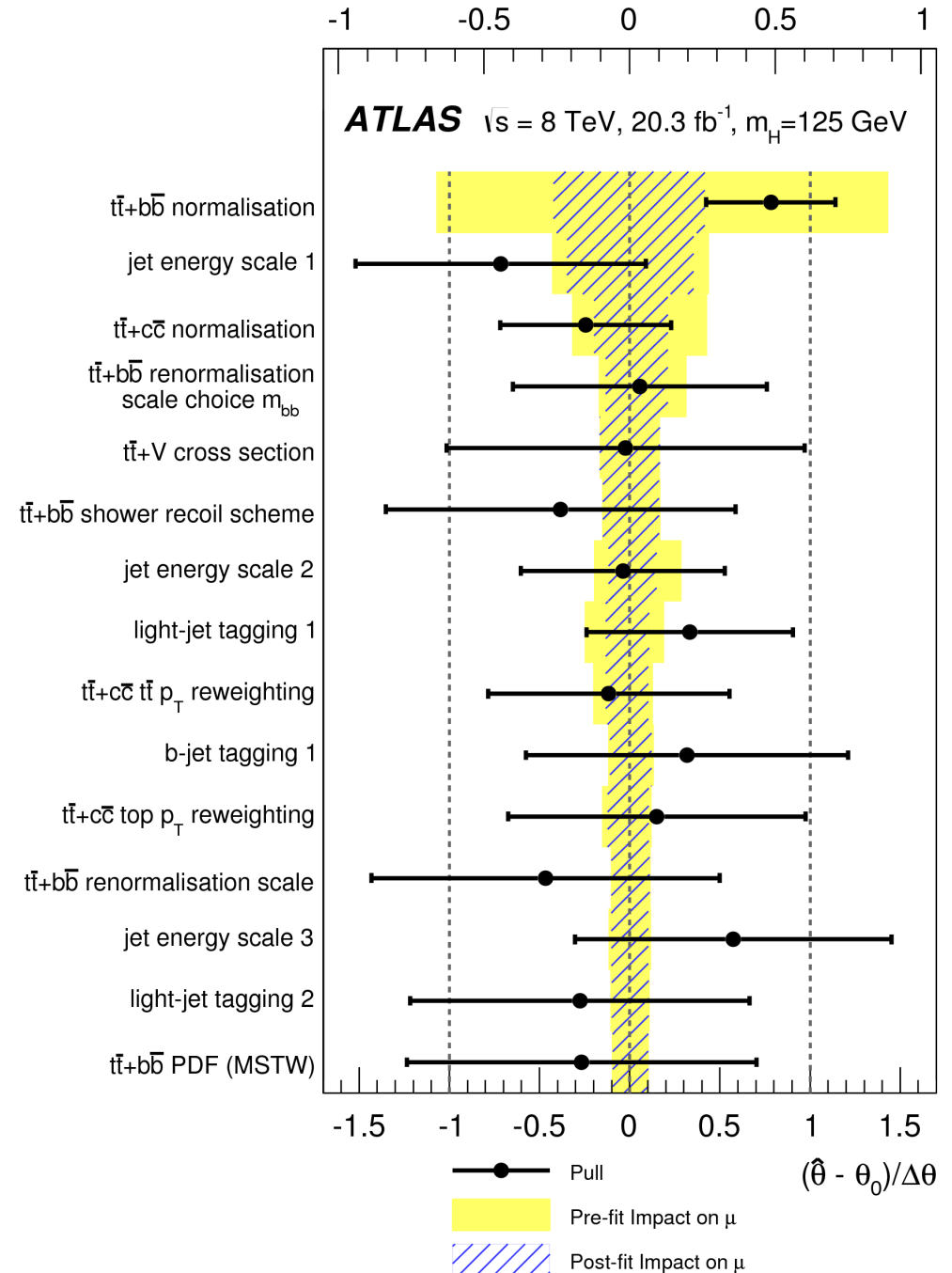
Strong handle  
 on  $t\bar{t}$ +light,  
 $t\bar{t}c$  and  $t\bar{t}b$   
 contributions



Include in fit!

# Fit Results

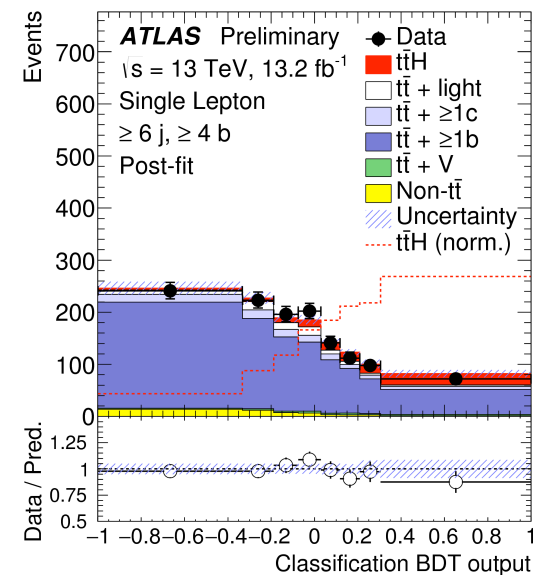
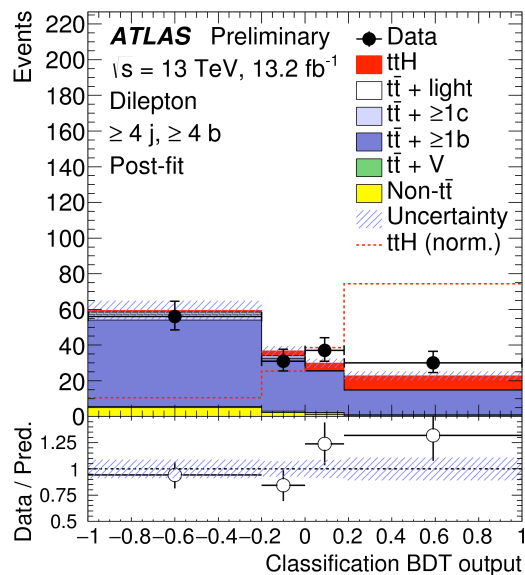
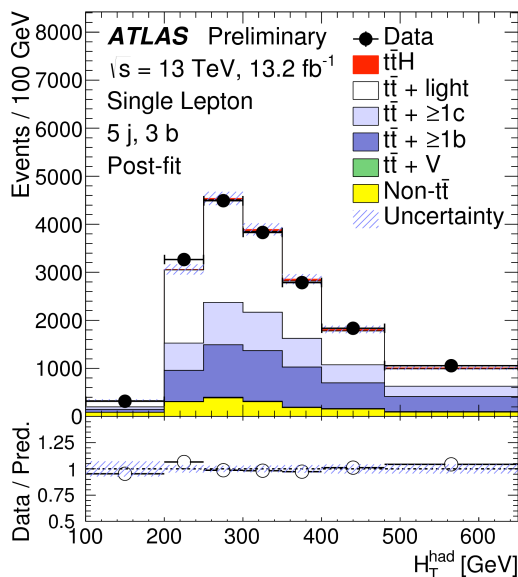
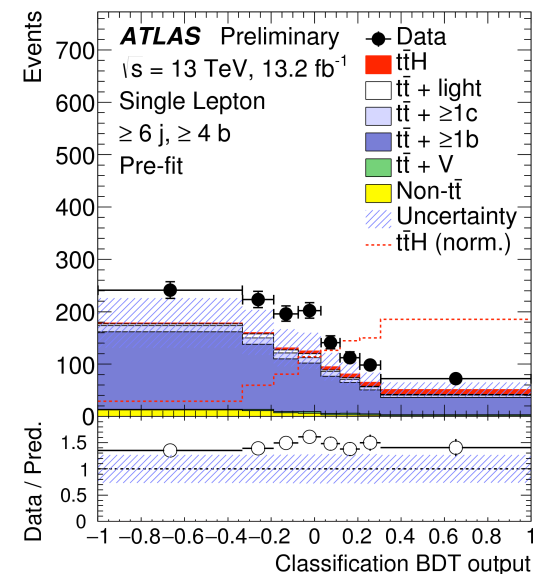
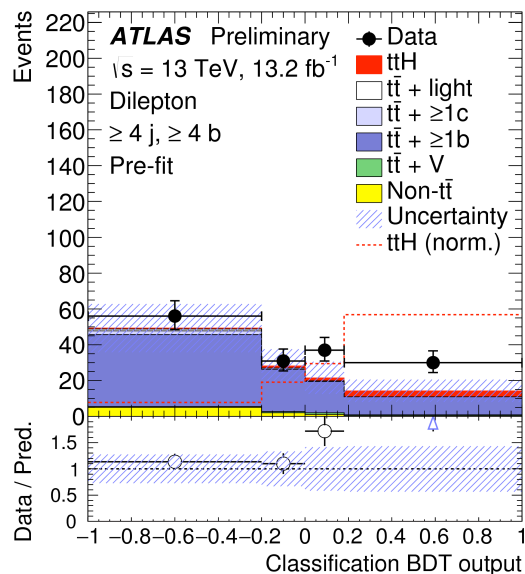
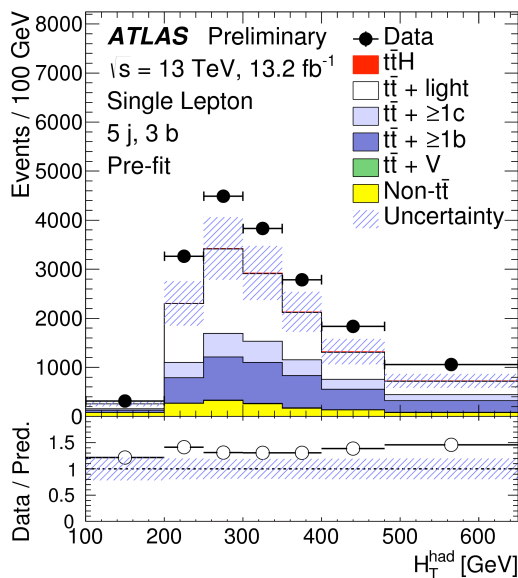
- ❖ Need to compare starting point and results
  - Pathologies due to lack of MC stats in some areas, strong correlations, ...
- ❖ Crucial to design analysis with good control regions: fit helps address least understood systematics



ATLAS  $ttH$  search: [arXiv:1503.05066](https://arxiv.org/abs/1503.05066)

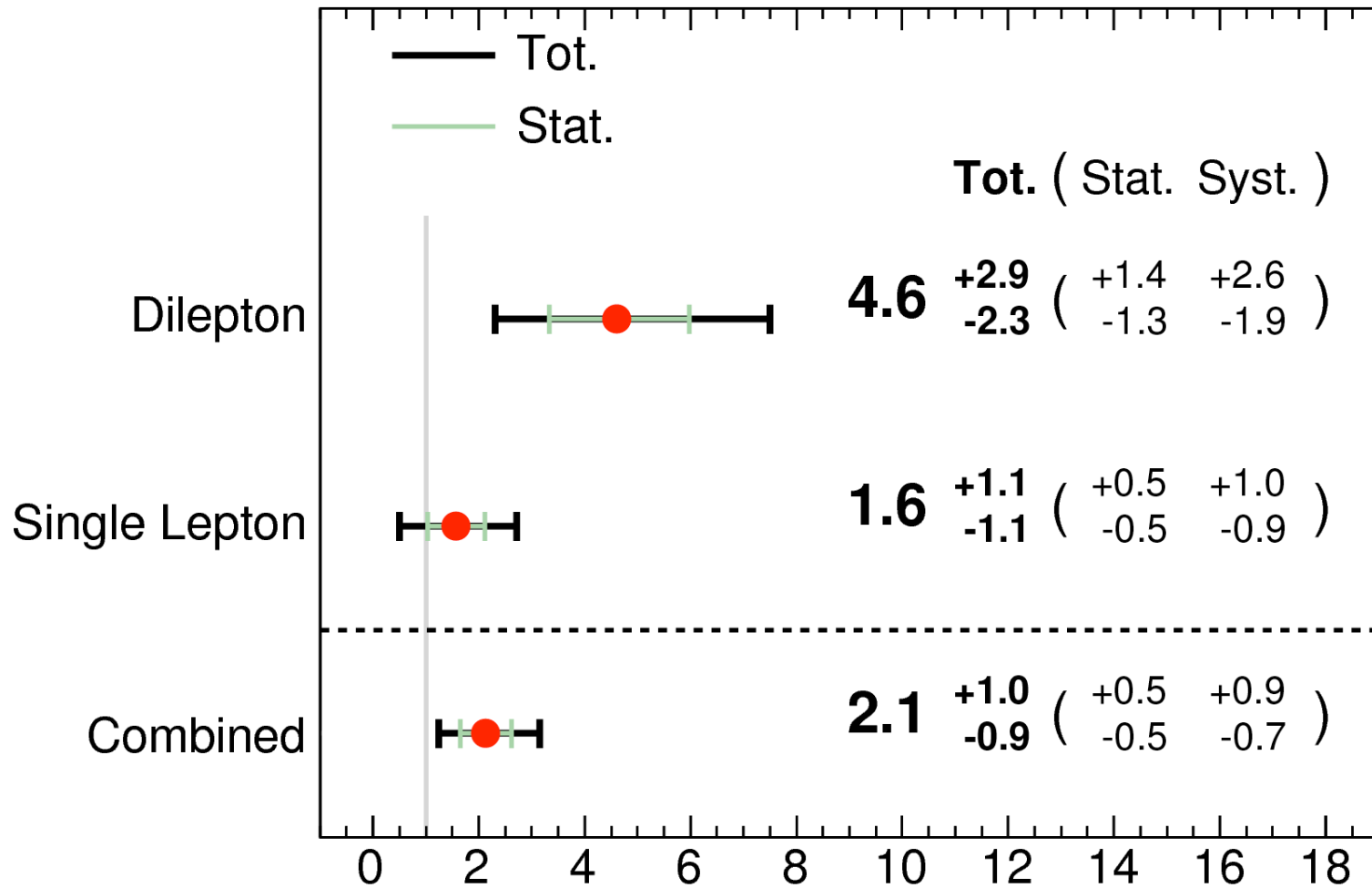


# Pre- vs Post-Fit



# Oh Yeah

**ATLAS** Preliminary  $t\bar{t}H$  ( $b\bar{b}$ ),  $\sqrt{s} = 13$  TeV,  $13.2 \text{ fb}^{-1}$



Best fit  $\mu = \sigma^{t\bar{t}H} / \sigma_{SM}^{t\bar{t}H}$  for  $m_H = 125$  GeV

# So, Physics Analysis

## ❖ Start from:

- “*How well* do we understand data *and* the SM?”
- How confident are we in corrections we apply?

## ❖ Given that:

- Which measurements can we make? What do we need to do to improve our understanding?

## ❖ Balance the work!

- Early, low background searches
  - Detailed understanding/verification of SM predictions
- Complementary measurements!**

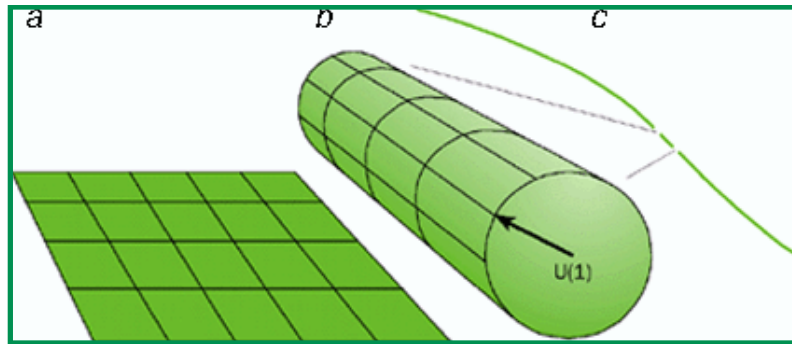
## ❖ Increasingly complex searches

- Tough backgrounds, hard work
- Don't scorn multivariate and statistical tools

# Excesses

# Extra Dimensions

- ❖ A promising approach to quantum gravity consists in adding extra space dimensions: string theory
  - Additional space dimensions are hidden, presumably because they are compactified



Source: PhysicsWorld

- ❖ Radius of compactification usually assumed to be at the scale of gravity, i.e.  $10^{18}$  GeV
  - In '90 Antoniadis realized they may be much larger...

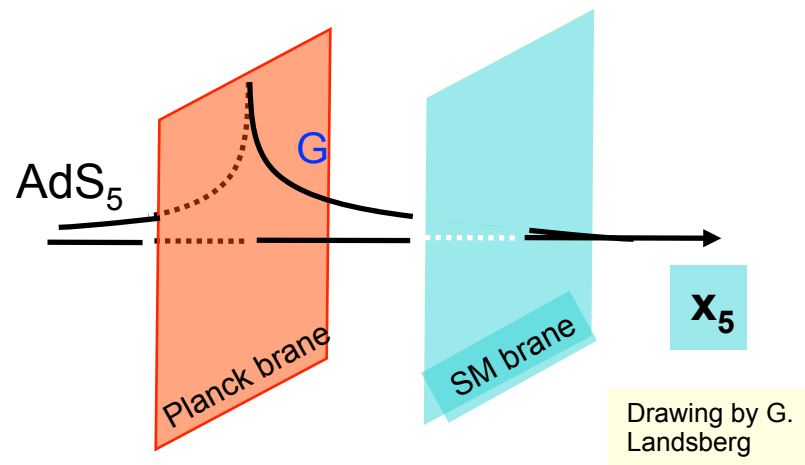
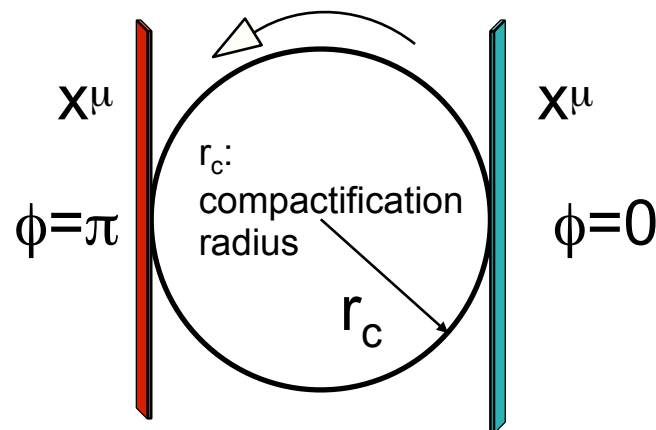
Phys.Lett.B246:377-384,1990

# Warped Extra Dimensions

## ❖ “Simple” Randall-Sundrum model:

Randall & Sundrum, Phys.Rev.Lett. 83 (1999) 3370

- SM confined to a brane, and gravity propagating in an extra dimension
- The metric in the extra dimension is “warped” by a factor  $\exp(-2kr_c\phi)$
- (Requires 2 branes)



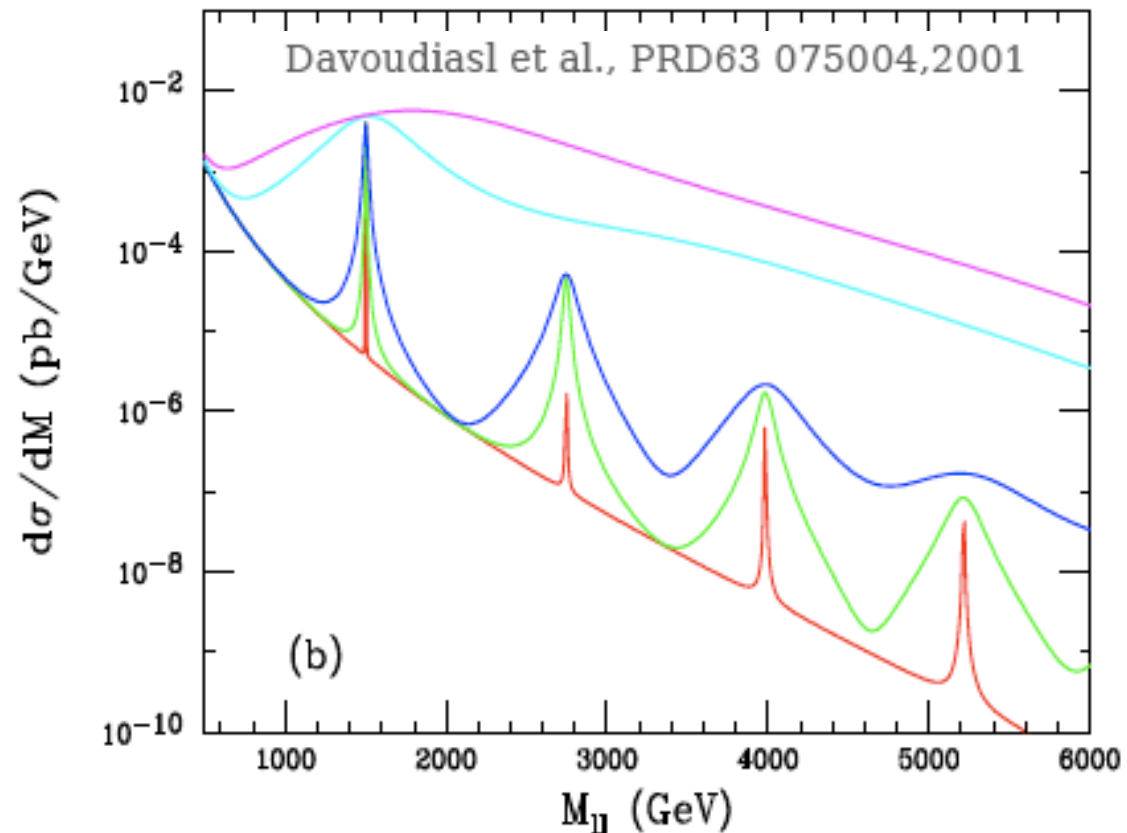
# Graviton Excitations

❖ In RS, get a few massive graviton excitations

- Widths depend on warp factor  $k$
- Mass separation = zeros of Bessel function

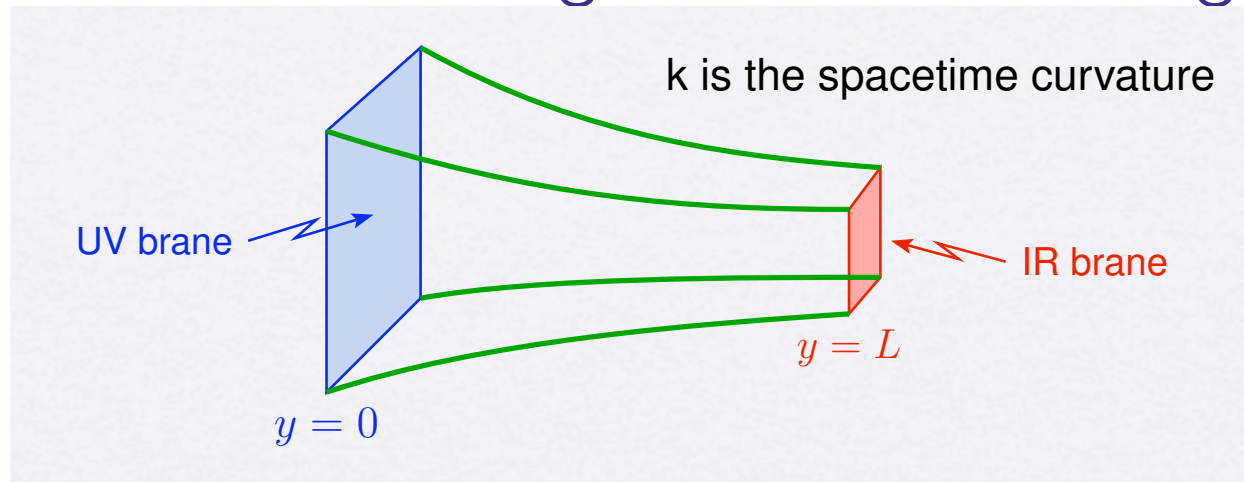
➔ Smoking gun!

(BRs also different than  $Z'$ :  
e.g.  $\gamma\gamma$  allowed)



# Hierarchies

- ❖ Physics on a curved gravitational background:



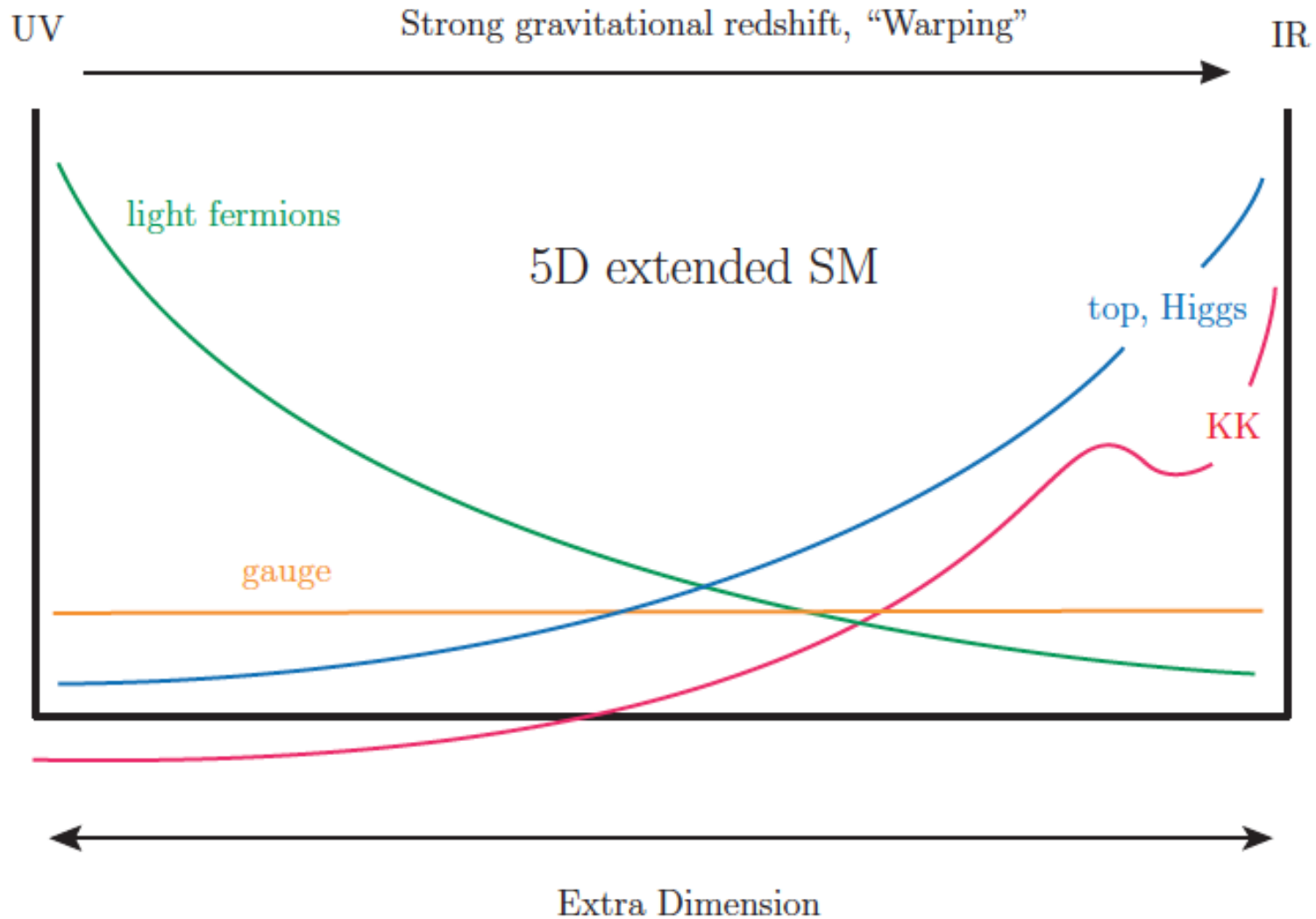
- ❖ Scales depend on position along extra dimensions
  - UV brane scale is  $M_{\text{Pl}} = 2 \times 10^{18}$  GeV
  - IR brane scale is  $M_{\text{Pl}} e^{-kL} \sim 1$  TeV if  $kL \sim 30$
- ❖ If were to localize Higgs on IR brane, naturally get EW scale  $\sim 1$  TeV (from geometry!)



# Flavor

- ❖ Interesting variation has fermions located along the extra dimension
  - Fermion masses generated by geometry
  - Heavier fermions are closer to IR brane, and gauge boson excitations as well
    - Gauge boson excitations expected to have masses in the 3-4 TeV range (bounds from precision measurements)
    - Couple mainly to top/W/Z (!)
  - Flavor changing determined by overlap of fermion “wave function” in the ED
    - Nice suppression of FCNC etc.

# Graphically

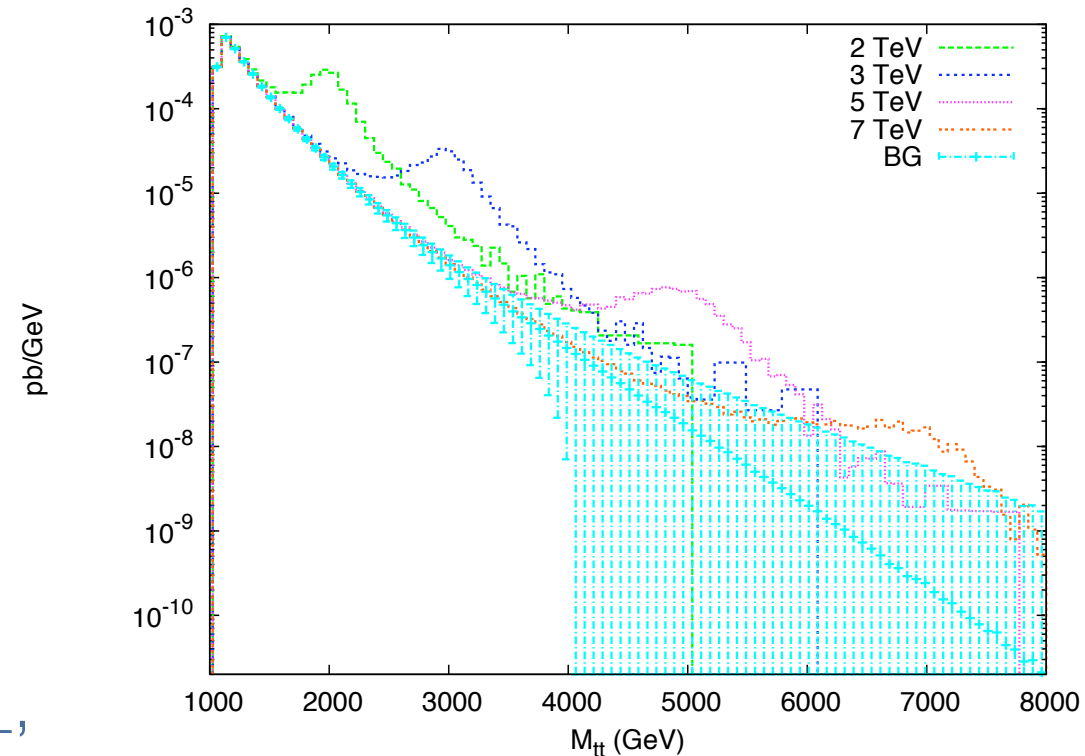


(From K Agashe et al, [arXiv:1608.00526](https://arxiv.org/abs/1608.00526))

# Gauge Boson Excitations

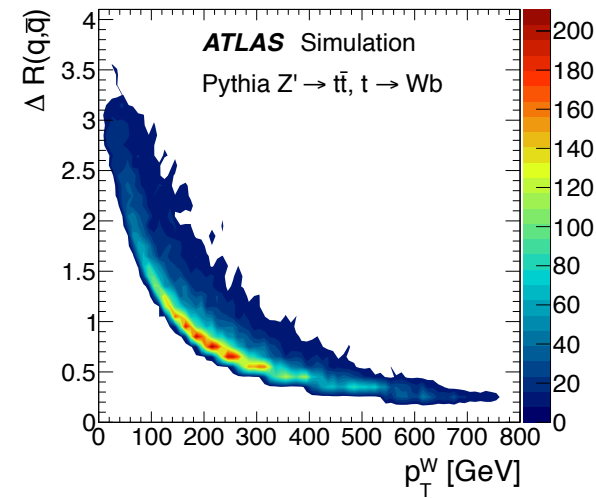
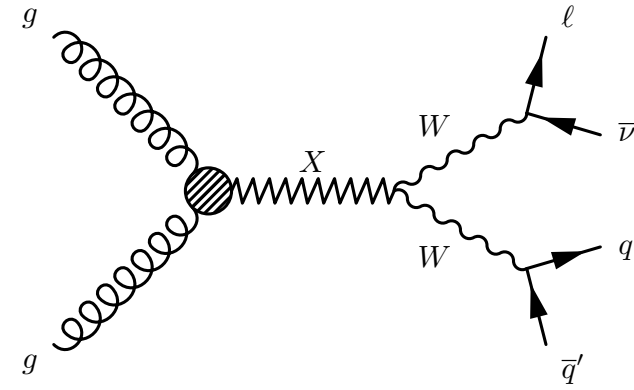
- ❖ Excitations of the gauge bosons are very promising channels for discovery
  - Couplings to light fermions are small
  - Small production cross-sections
  - Large coupling to top,  $W_L$ ,  $Z_L$
  - Look for  $t\bar{t}$ ,  $WW$ ,  $ZZ$  resonances (that can be wide)

B. Lillie et al., JHEP 0709:074,2007



# Dibosons

- ❖ Signature of Randall-Sundrum excitations as well as  $W'$ ,  $Z'$ 
  - Can look for e.g.  $W' \rightarrow WZ$ ,  $WH$ 
    - Many final state options:  $lvll$ ,  $lvqq$ ,  $qqll$ ,  $qqqq$ , ...
    - Three leptons  $\rightarrow$  low background but low branching ratio: good at low mass where backgrounds are large
    - Fully hadronic  $\rightarrow$  high branching ratio but substantial multi-jet background: good at high mass where cross-section is lower
  
- ❖ For high mass  $W'$ ,  $Z'$  decay products are boosted...ok for leptonic decays,
  - ... but hadronic decay products merge:
    - $\Delta R \sim 2m/p_T \Rightarrow$  for  $p_T \sim 500$  GeV,  $\Delta R \sim 0.4$ , typical jet size



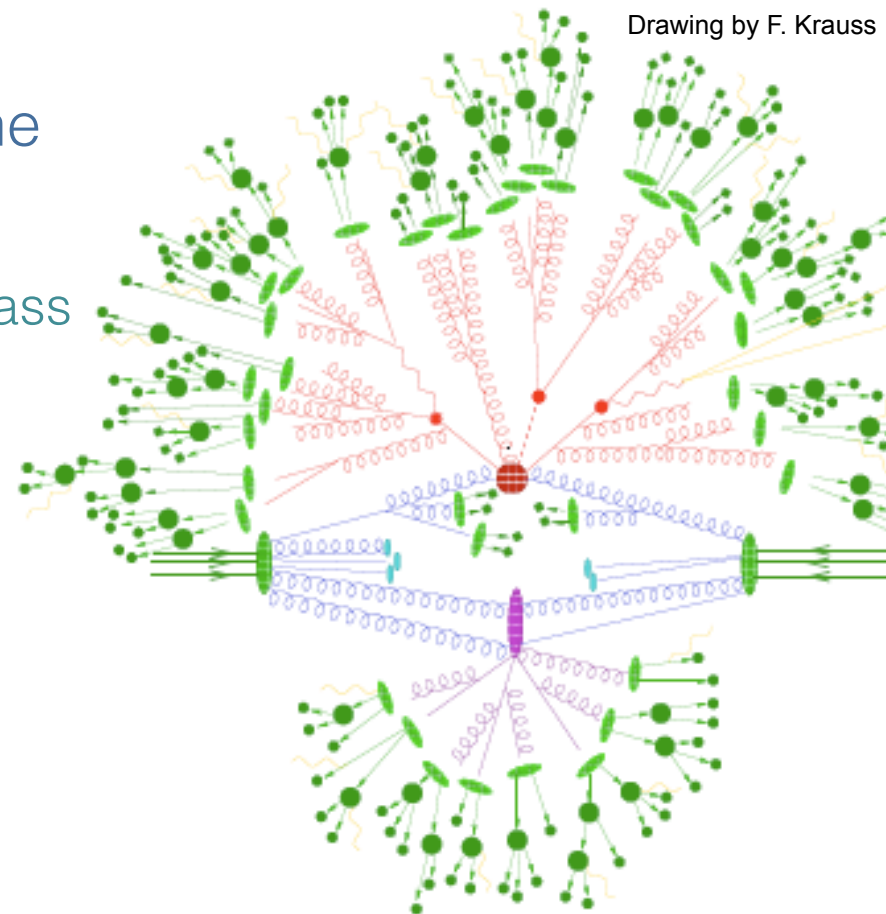
# Fully Hadronic Decays

- ❖ Decay hadrons reconstructed as a single jet
  - But even if it looks like a single jet, it originates from a massive particle decaying to two (W, Z, H) or three (top) hard partons, not one

- If I measured each of the partons in the jet perfectly, I would be able to:
  - Reconstruct the “originator’s” invariant mass
  - Reconstruct the direct daughter partons

But

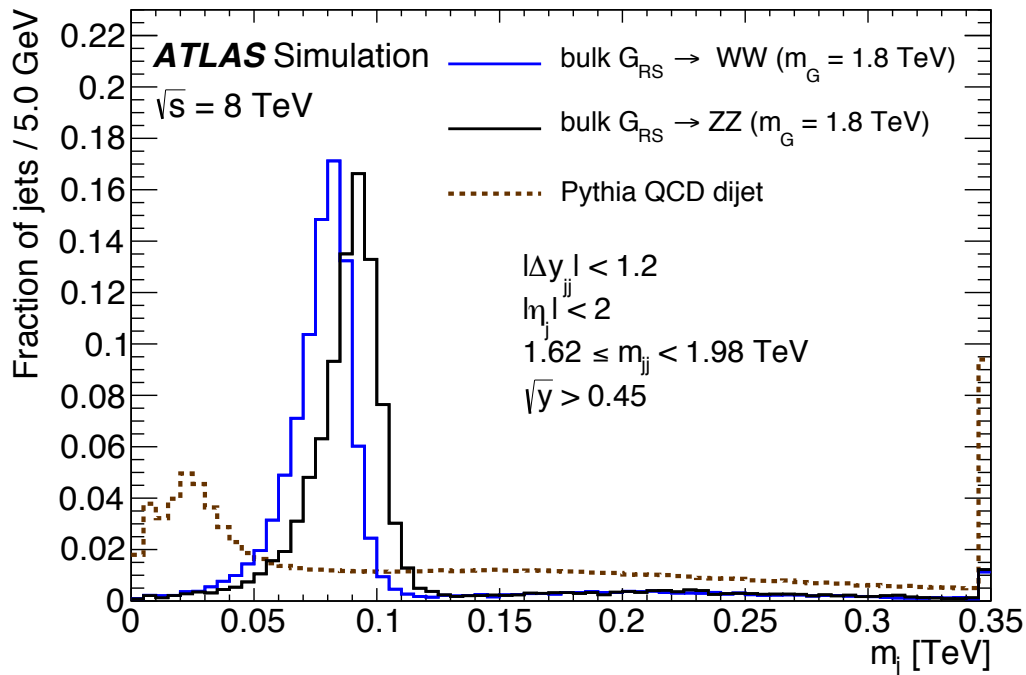
- Quarks hadronize → cross-talk
- My detector can't resolve all individual hadrons



# Jet Mass

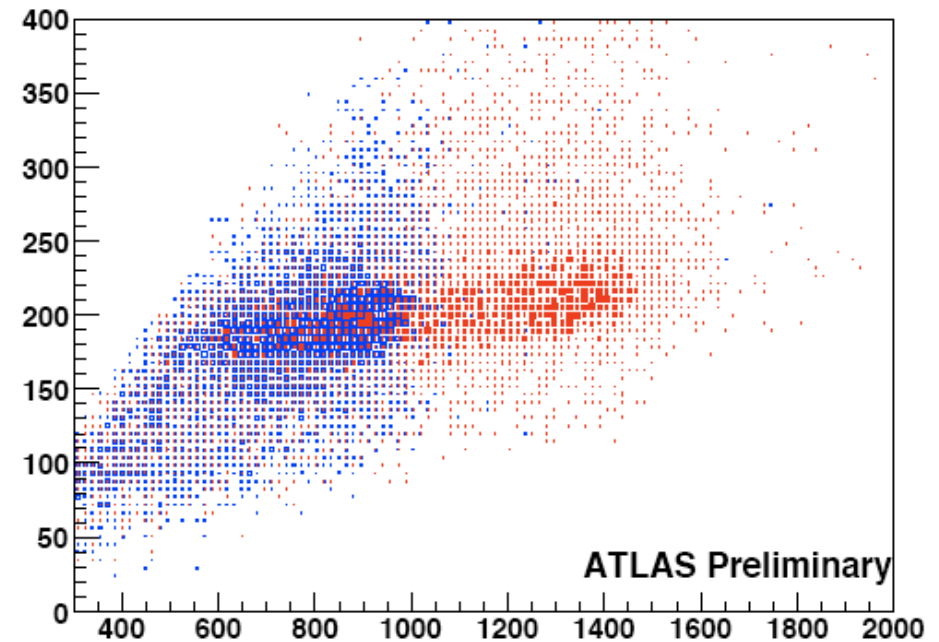
- ❖ Jet mass: invariant mass of all jet constituents
  - In principle, close to object mass
  - and invariant!

ATLAS, arXiv:1506.00962



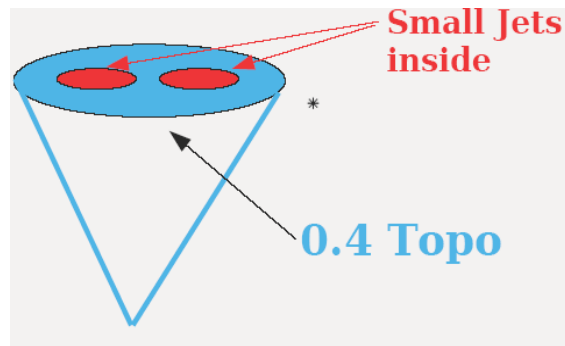
Jet Mass vs pT

For top quarks...



# Subjects

- ❖ Jet mass is not sensitive to structure
  - Can't tell whether a jet is isotropic or not
- ❖ Expect “blobs” with higher concentration of energy for jets from top/W/Z decays



- ❖ Multiple ways of exploiting this....
  - $k_T$  splitting scales, “mass drop”, ...

J. M. Butterworth, B. E. Cox, and J. R. Forshaw, *Phys. Rev.* **D65** (2002) 096014

# $k_T$ Splitting Scales

❖  $k_T$  jet algorithm is much better suited to understand jet substructure than cone:

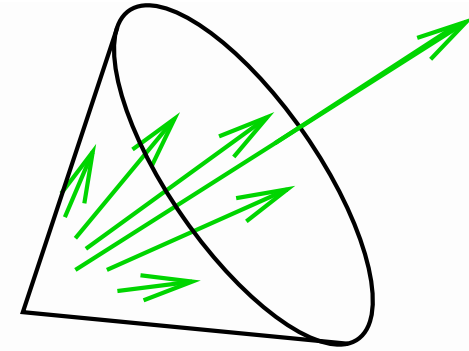
- Cone maximizes energy in an  $\eta \times \phi$  cone
- $k_T$  is a “nearest neighbor” clusterer

$$y_2 = \min(E_a^2, E_b^2) \cdot \theta_{ab}^2 / p_{T(jet)}^2$$

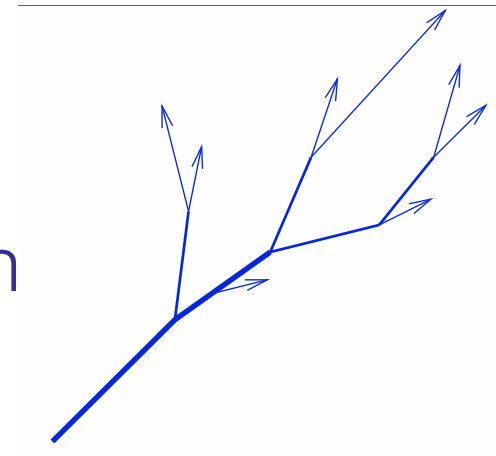
$$Y \text{ scale} = \sqrt{p_{T(jet)}^2 \cdot y_2}$$

❖ Can use the  $k_T$  algorithm on jet constituents and get the (y-)scale at which one switches from 1  $\rightarrow$  2 ( $\rightarrow$  3 etc.) jets

- Scale is related to mass of the decaying particle



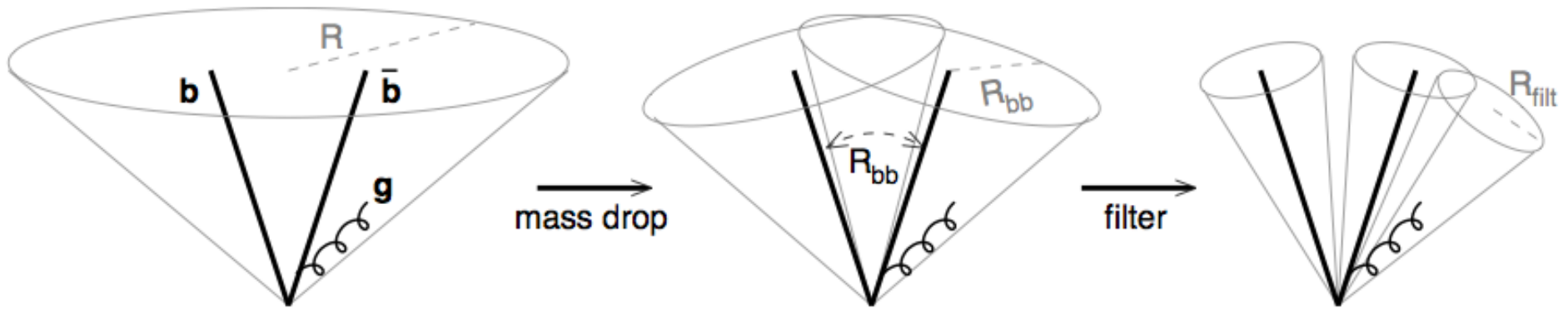
Cone



$k_T$



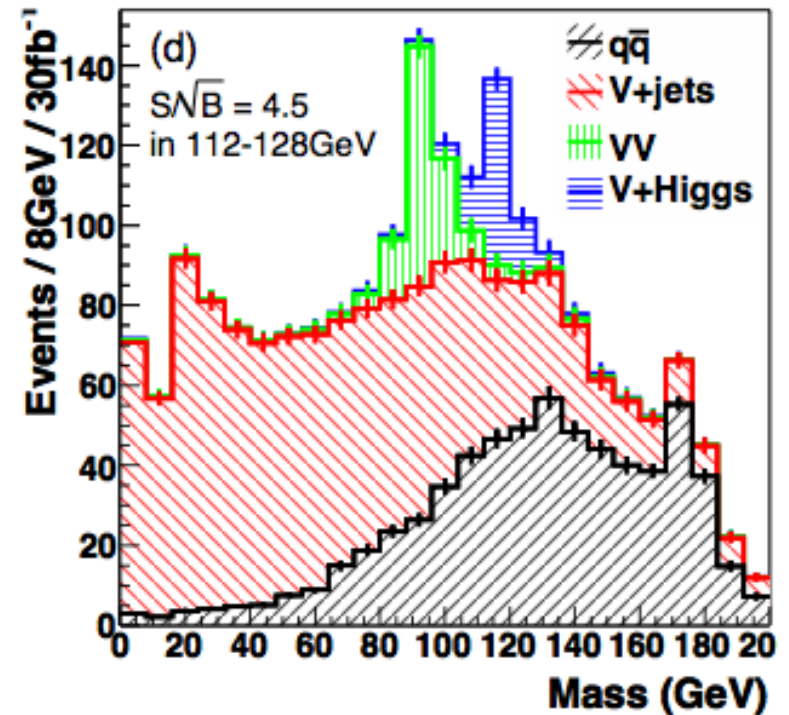
# Mass Drop



**BDRS, Phys.Rev.Lett. 100 (2008) 242001**

❖ Introduced to recover  $(W/Z)H \rightarrow (W/Z)bb$  at the LHC

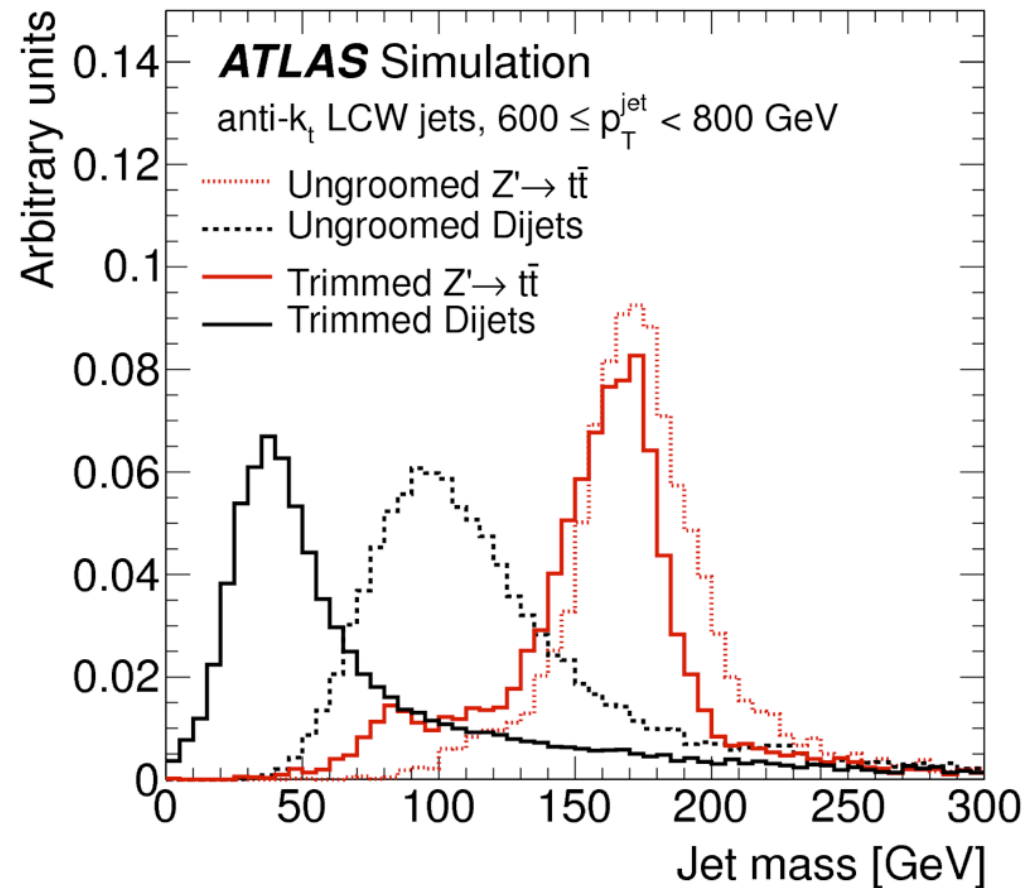
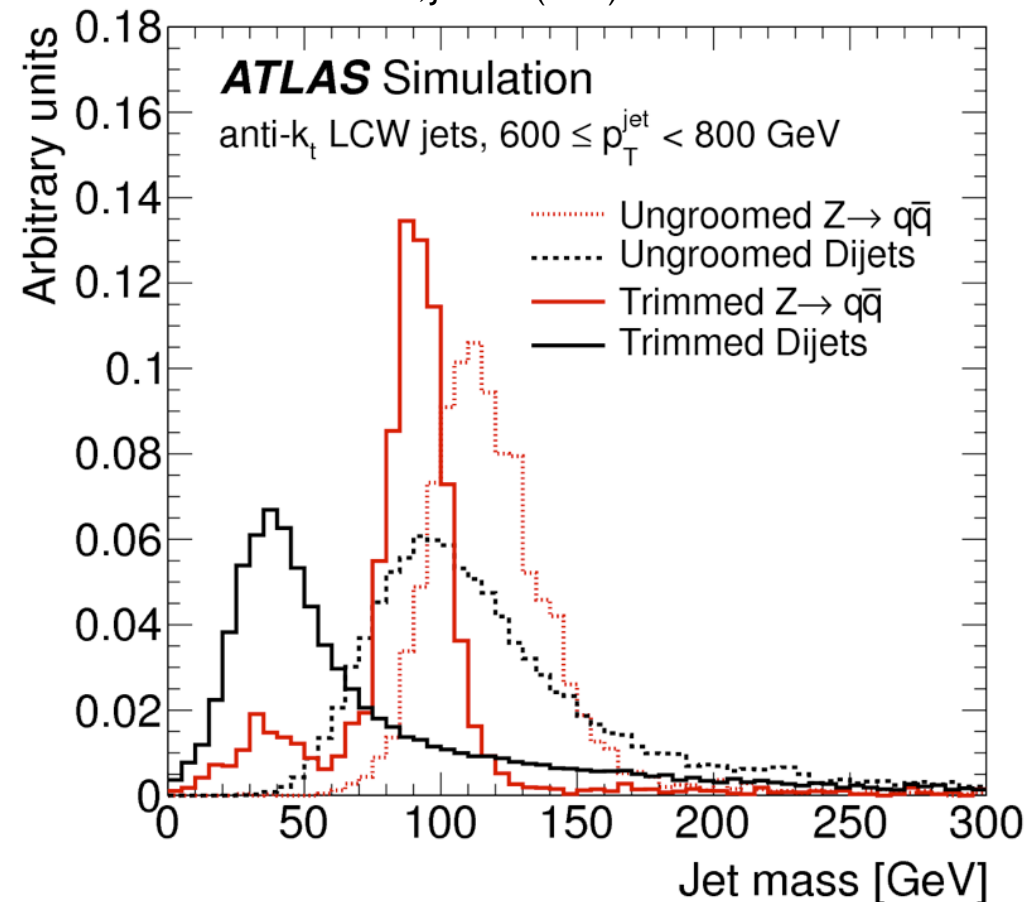
- In boosted regime,  $Wbb$  and  $Zbb$  backgrounds easier to manage
- But strategy can be used more generically



# Grooming

- ❖ Decluster (or recluster with small R), and remove soft stuff
  - Clean up soft QCD radiation/connection to underlying event

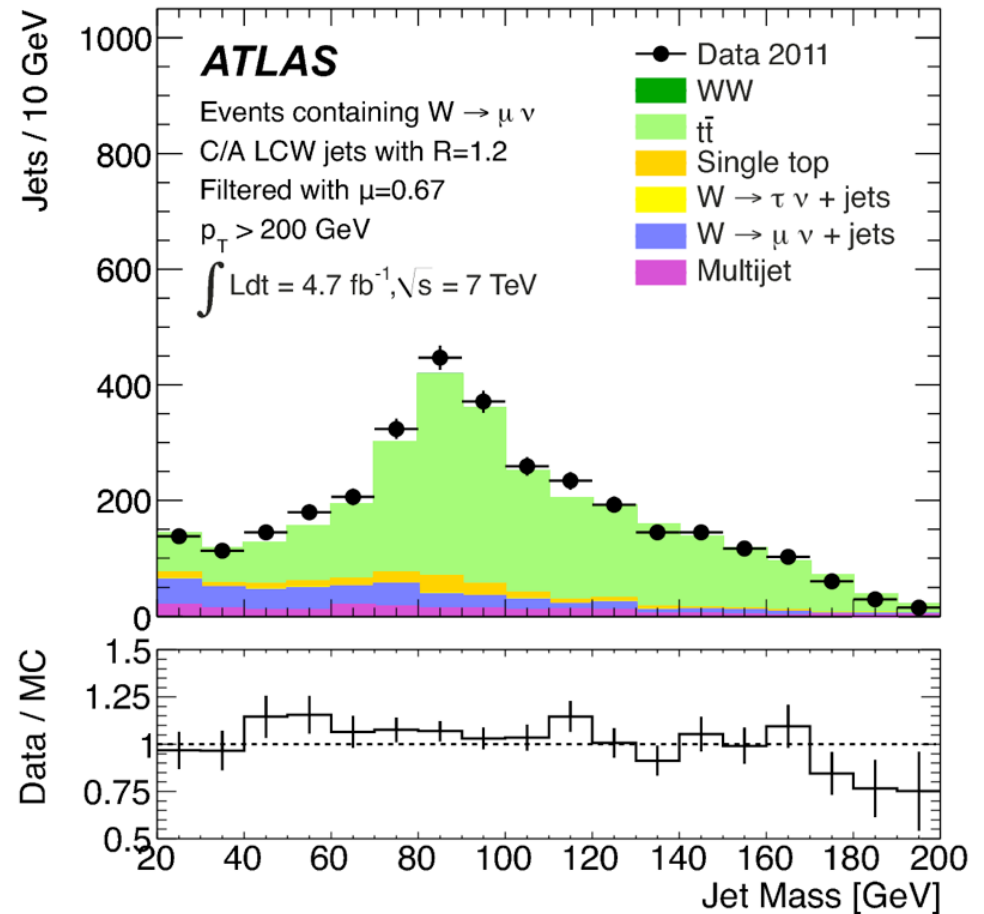
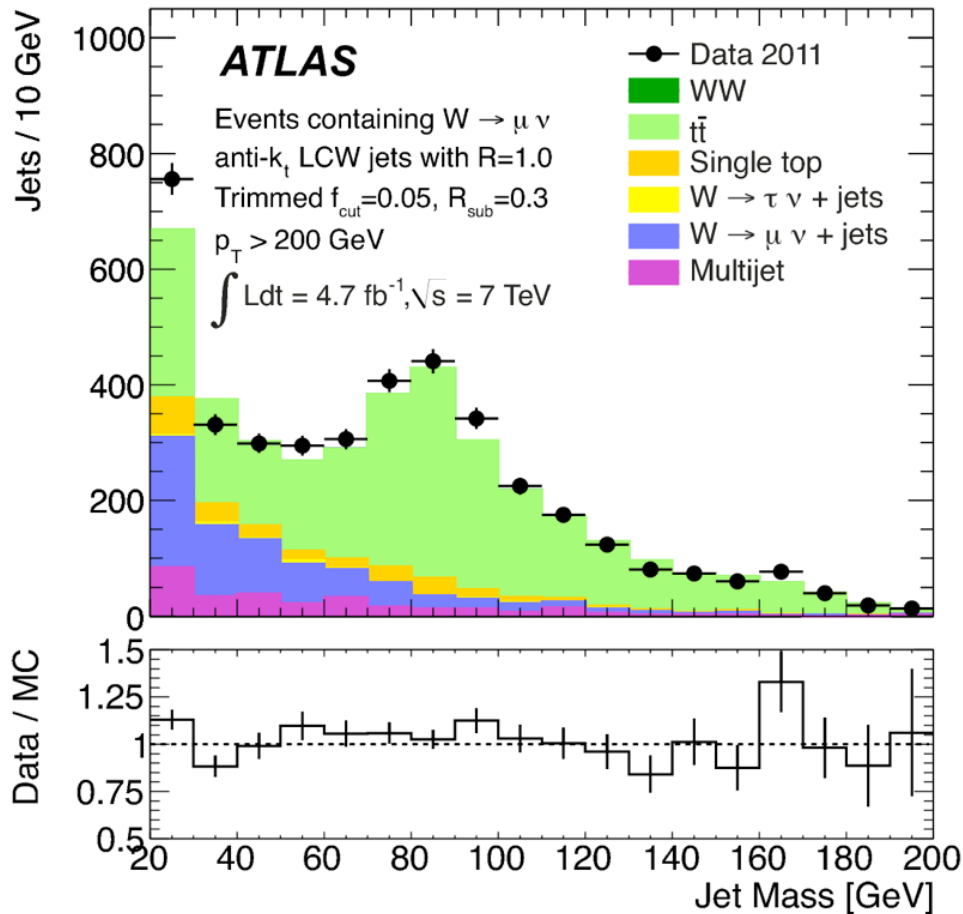
ATLAS, JHEP09 (2013) 076



# Proof of Principle

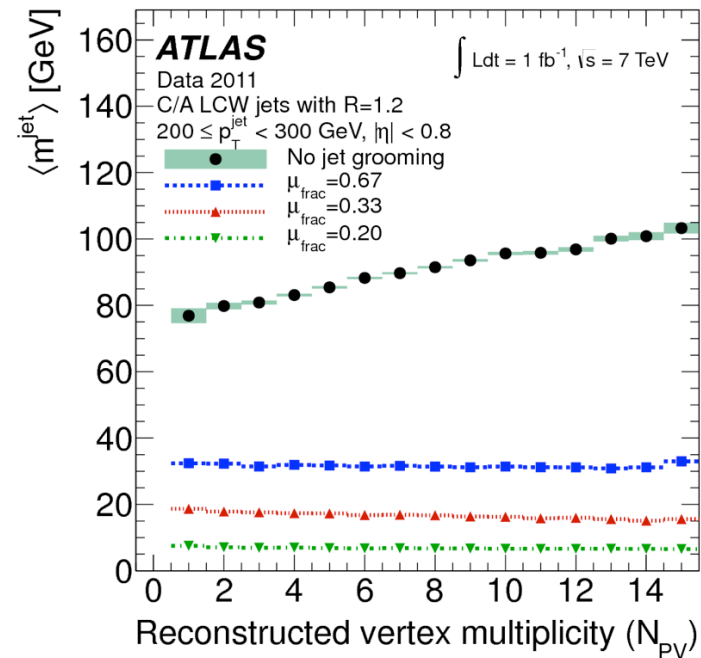
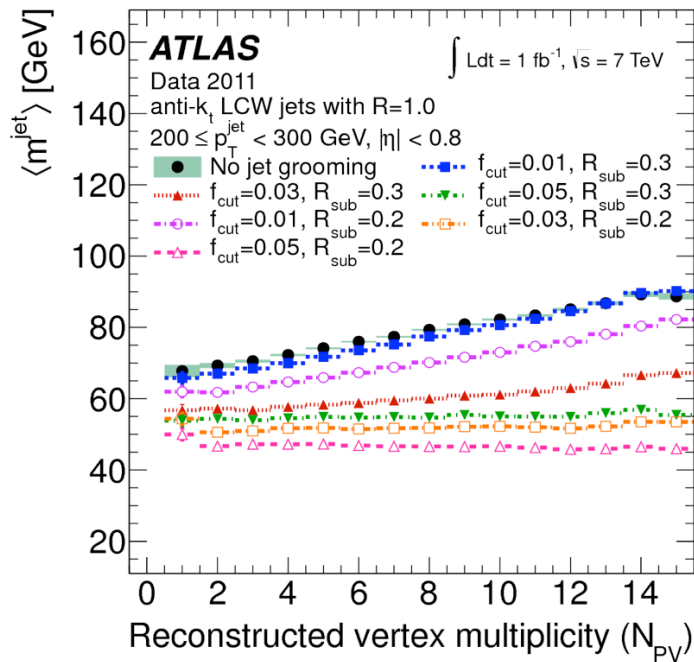
## ❖ W-jets from top quarks

ATLAS, JHEP09 (2013) 076



# Added Benefits

- ❖ Pile-up is a big deal at hadron colliders
  - Low- $p_T$ , “uninteresting” QCD will always have a much larger cross-section than rare processes we’re hunting

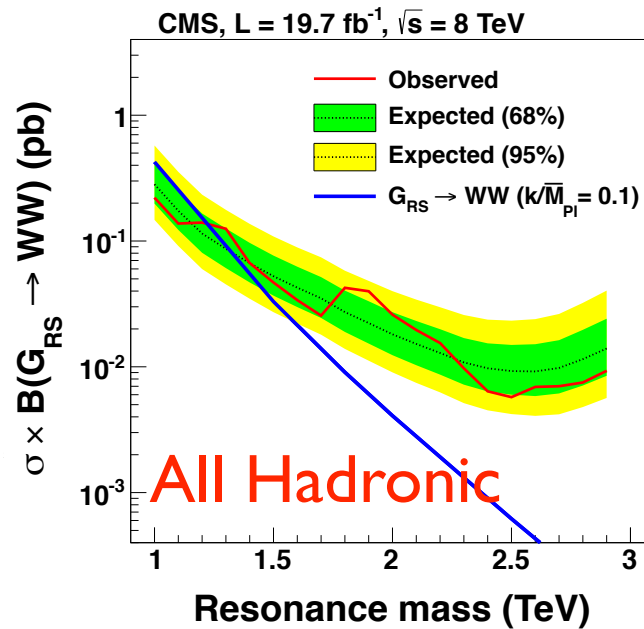
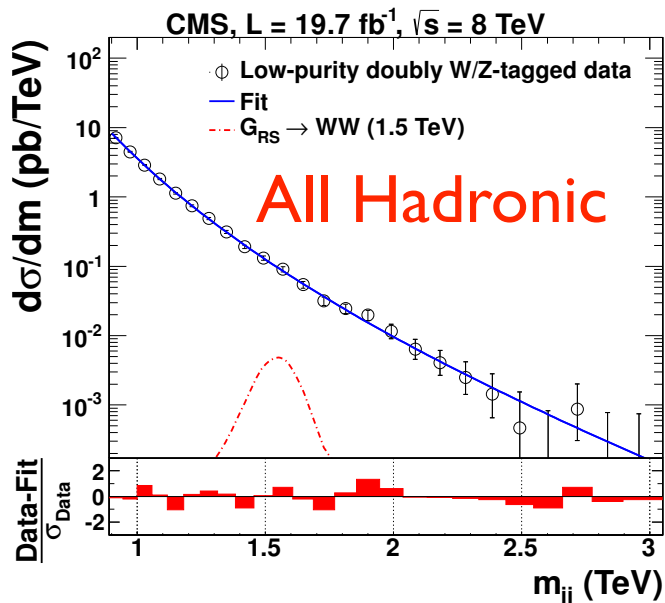
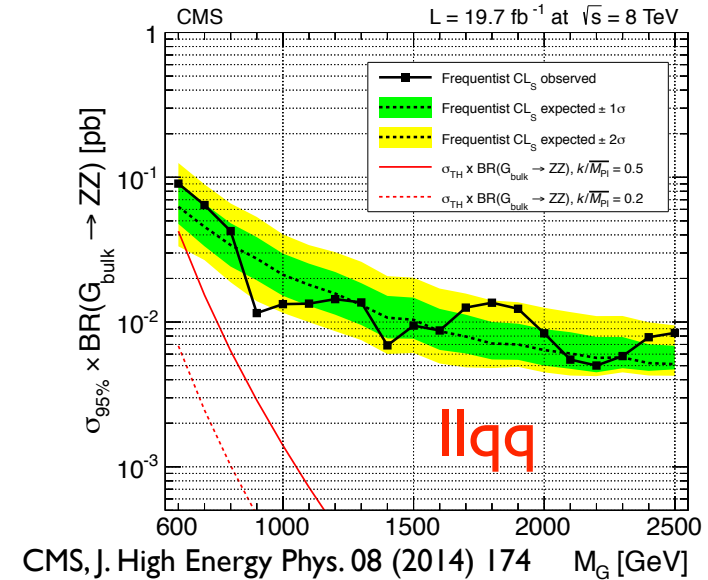
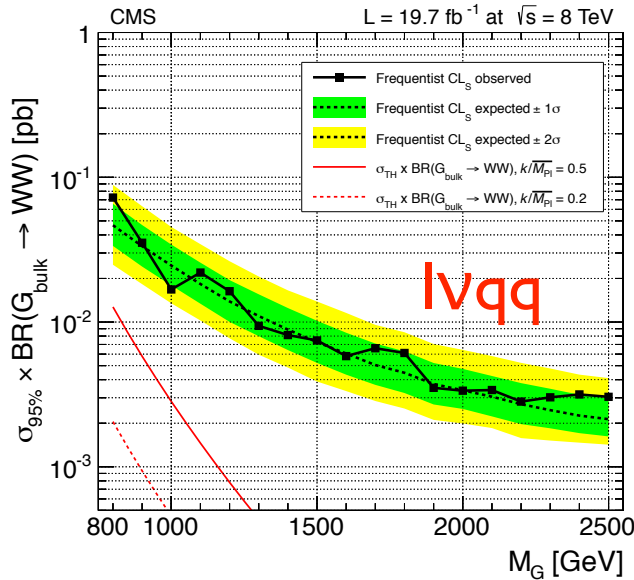
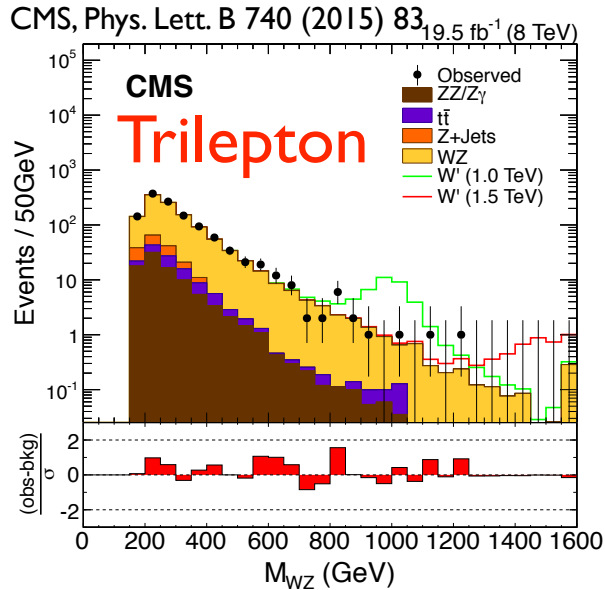


**!Optimal parameter set/strategy is detector-dependent!**

# Many More Techniques

- ❖ Whole “jet structure” community exists
  - Reports of BOOST workshops a very useful resource:
    - **Boosted objects: A Probe of beyond the Standard Model physics, A. Abdesselam et al, Eur.Phys.J. C71 (2011) 1661; Jet Substructure at the Tevatron and LHC: New results, new tools, new benchmarks, A. Abdesselam et al, J.Phys. G39 (2012) 063001**
  - Direct comparison of multiple taggers, and “groomers”
  - More tools have been developed, and also more extensive non-perturbative calculations of the jet structure
  - Many of the tools available in the fastjet library (Cacciari, Salam, Soyez)
    - <http://www.lpthe.jussieu.fr/~salam/fastjet/>

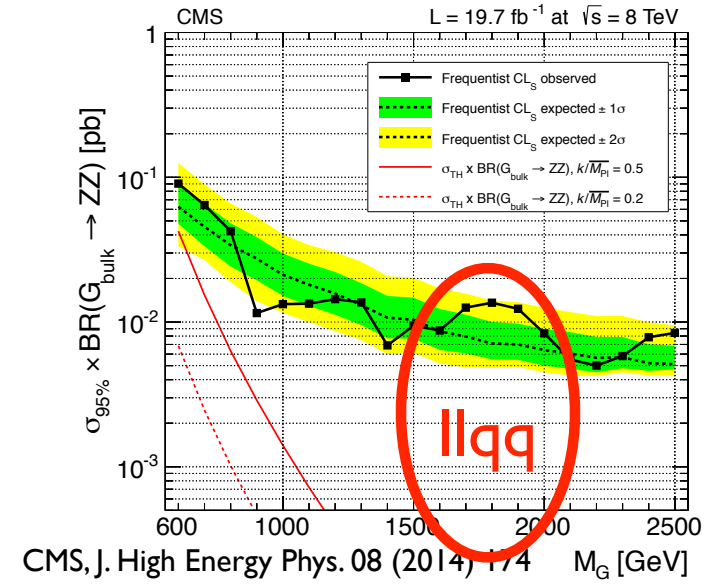
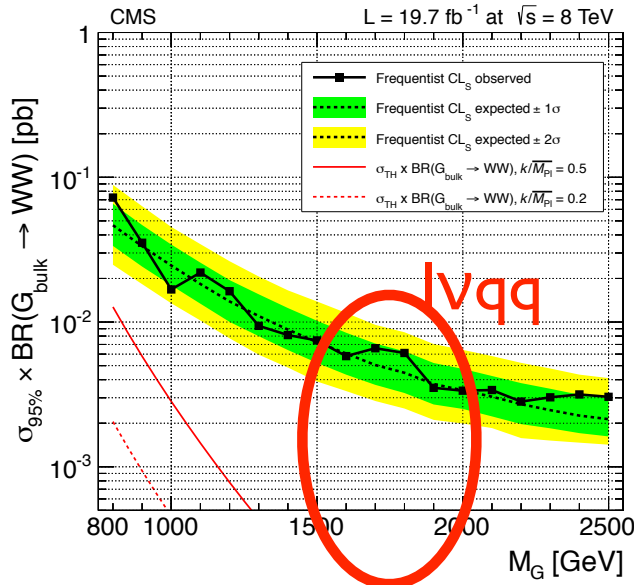
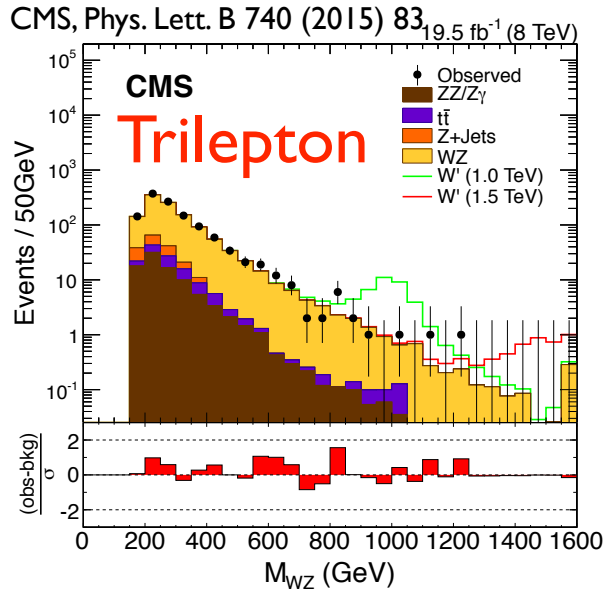
# Diboson Resonance Results: CMS



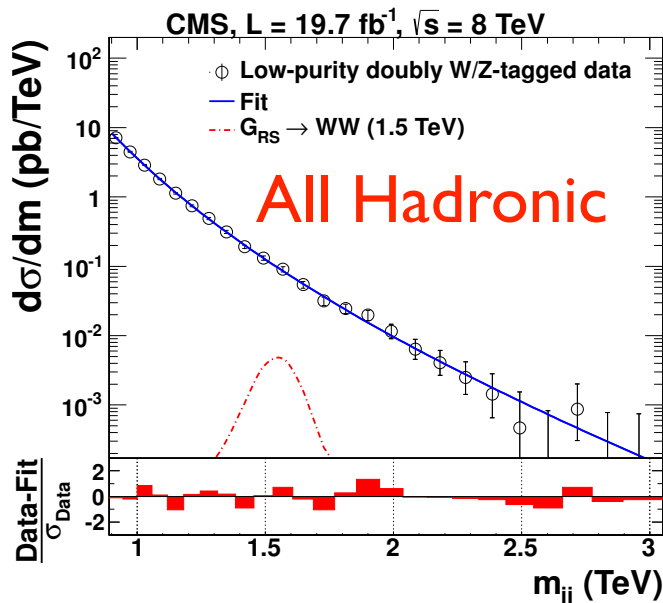
No evidence...

CMS, J. High Energy Phys. 08 (2014) 173

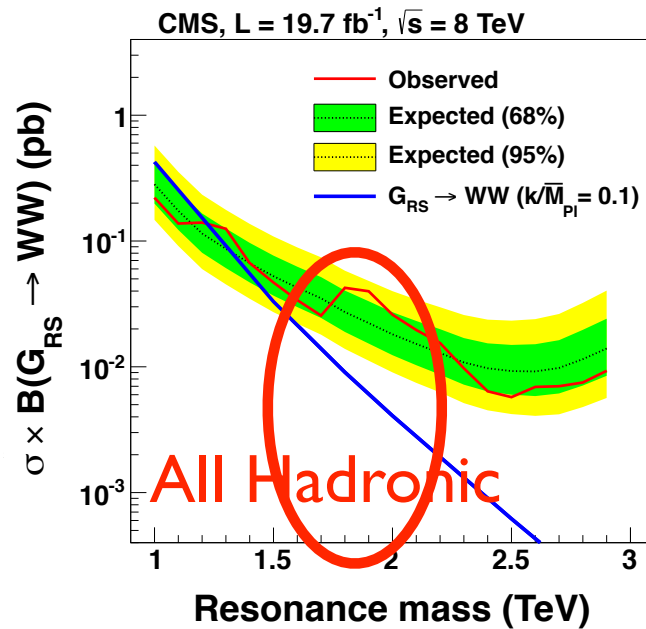
# Diboson Resonance Results: CMS



CMS, J. High Energy Phys. 08 (2014) 174



CMS, J. High Energy Phys. 08 (2014) 173

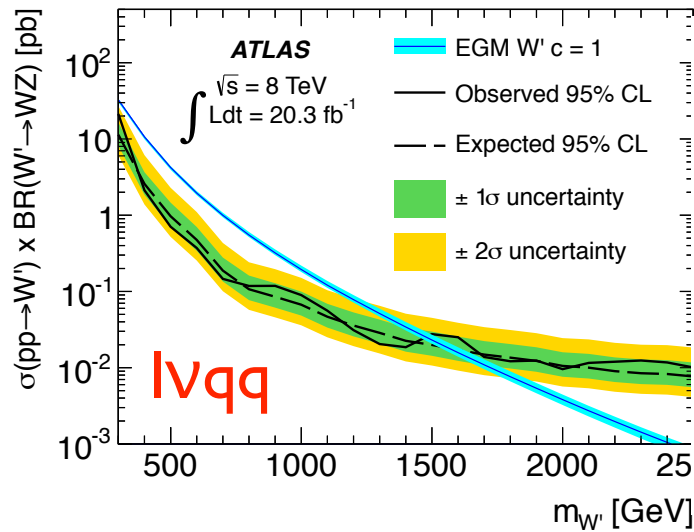
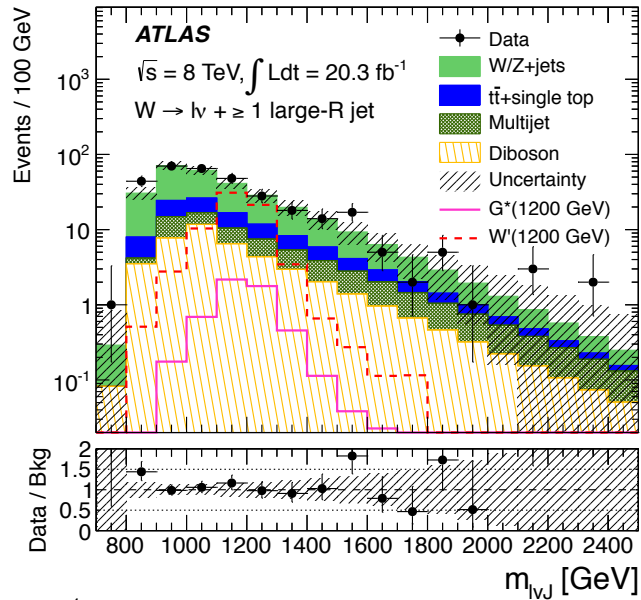


No evidence...  
but a coincidence?

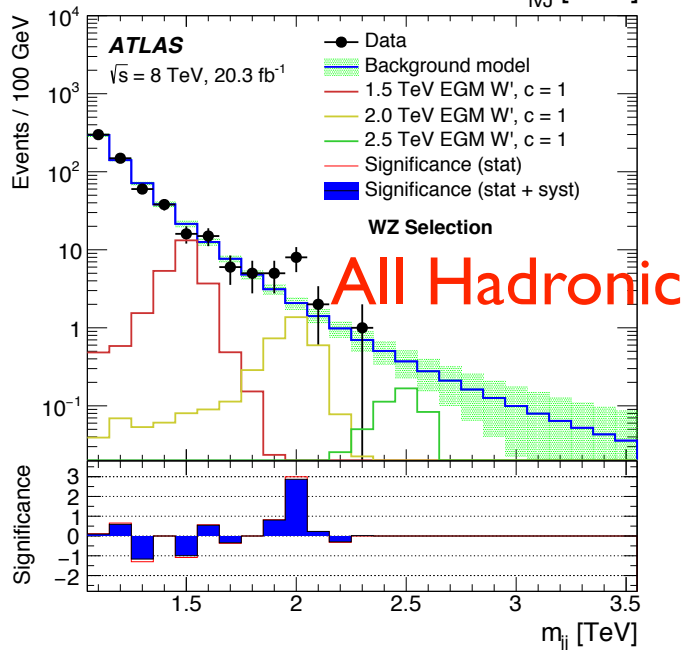
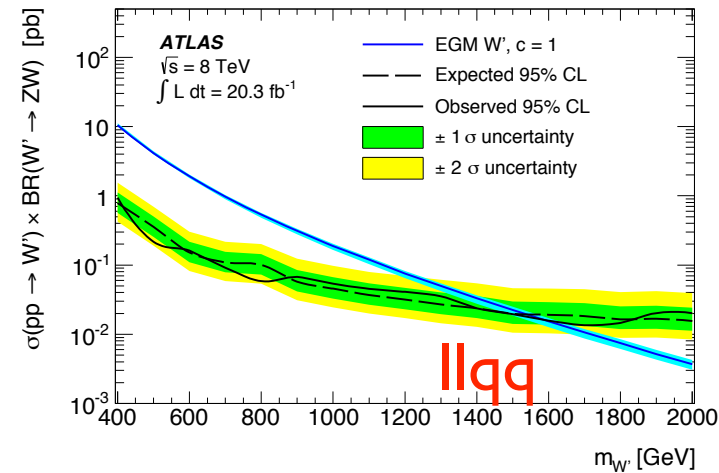


# Diboson Resonance Results: ATLAS

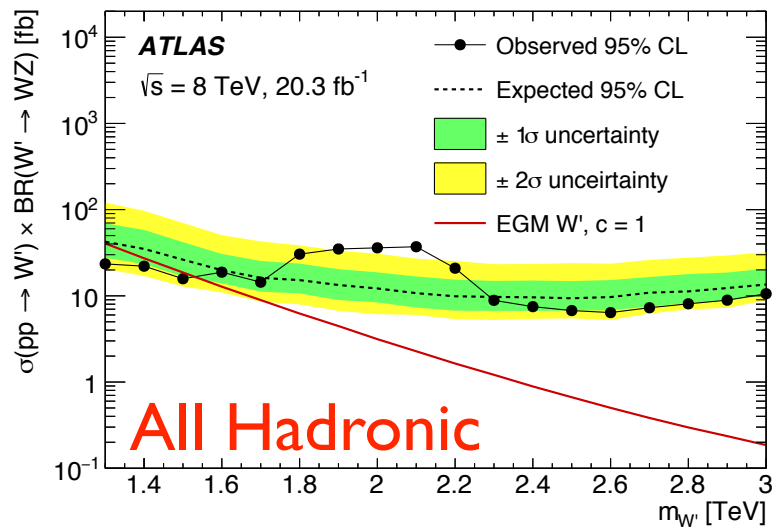
ATLAS, Eur. Phys. J. C (2015) 75:209



ATLAS, Eur. Phys. J. C (2015) 75:69



ATLAS, arXiv:1506.00962

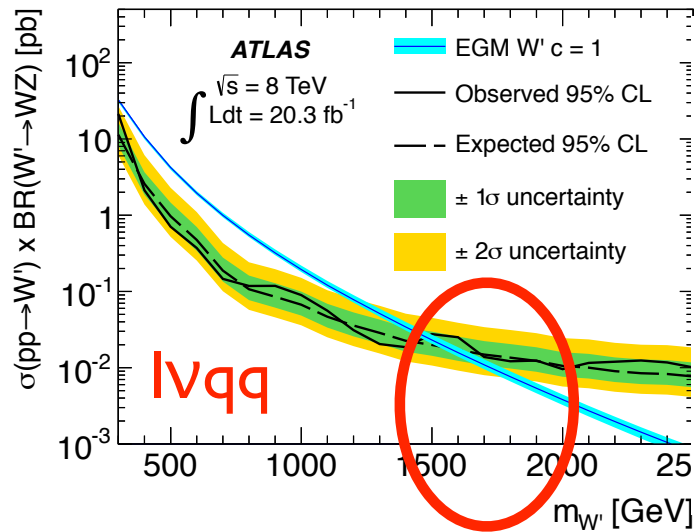
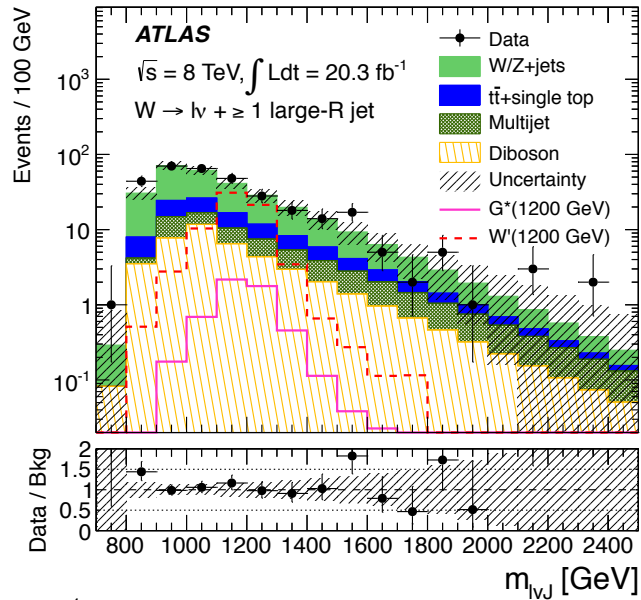


No evidence...

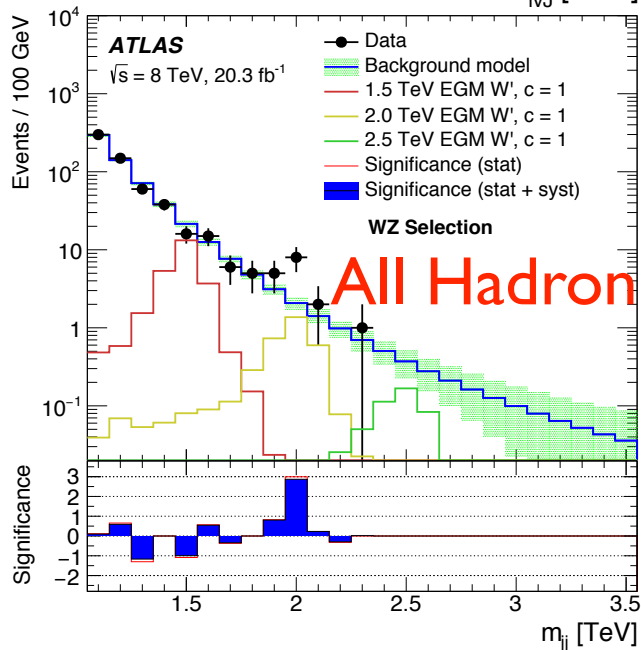
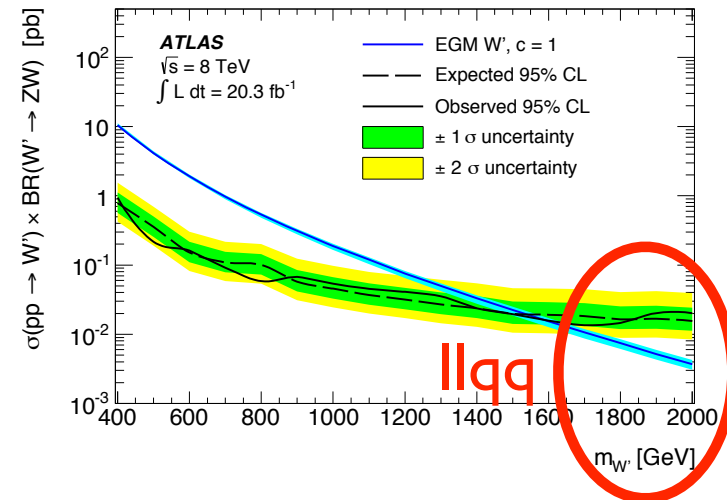


# Diboson Resonance Results: ATLAS

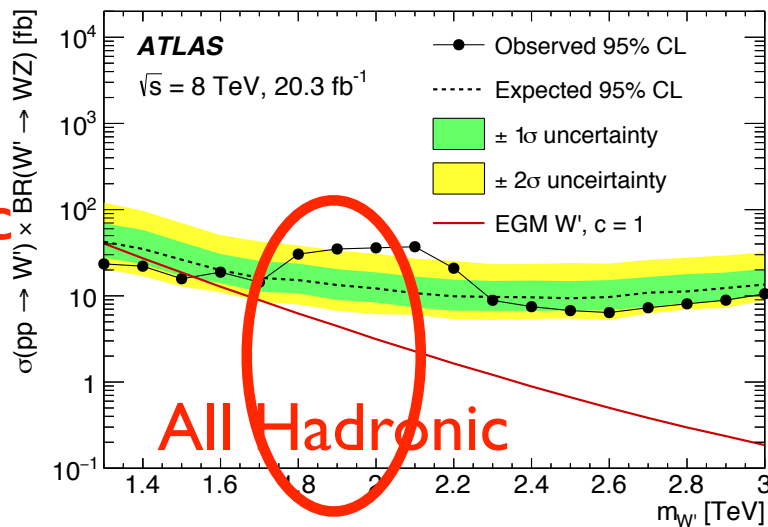
ATLAS, Eur. Phys. J. C (2015) 75:209



ATLAS, Eur. Phys. J. C (2015) 75:69



ATLAS, arXiv:1506.00962

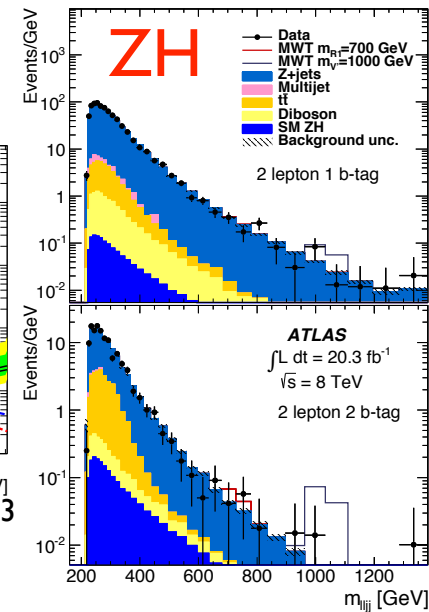
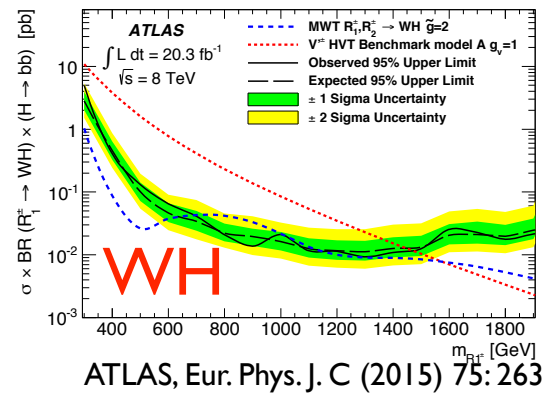
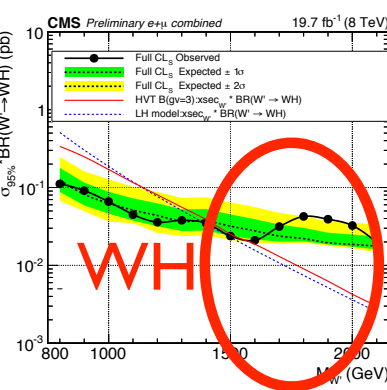
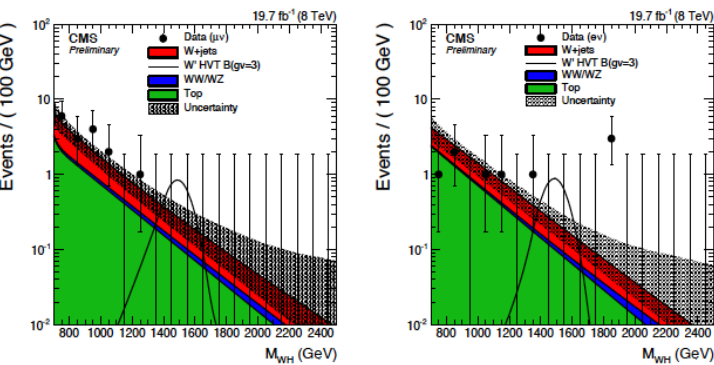


No evidence...  
 lesser coincidence?

# Decays to VH?

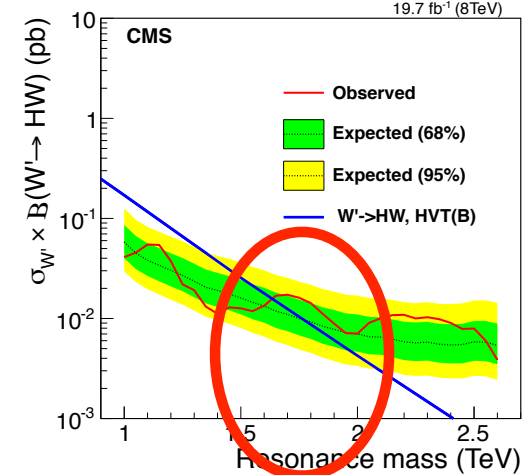
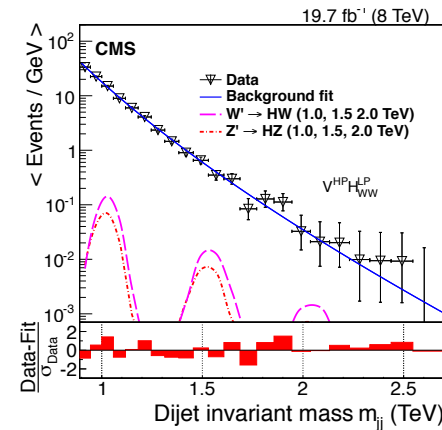
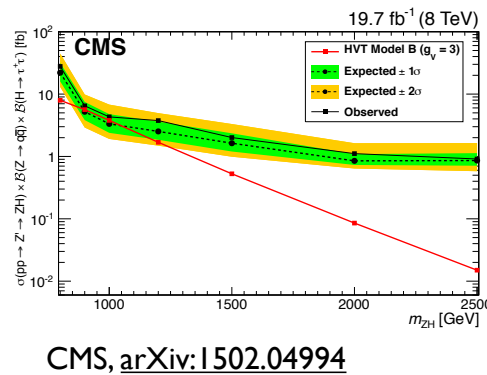
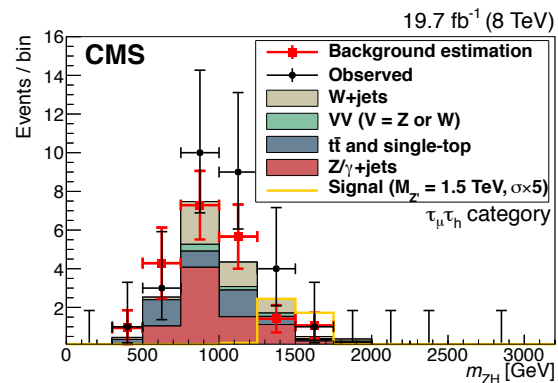
- Both ATLAS & CMS have analyses with b-tags  
 $\Rightarrow$  reduced acceptance

CMS-PAS-EXO-14-010



- CMS ZH  $\Rightarrow$   $\tau\tau$ J and all-hadronic VH

CMS, arXiv:1506.01443

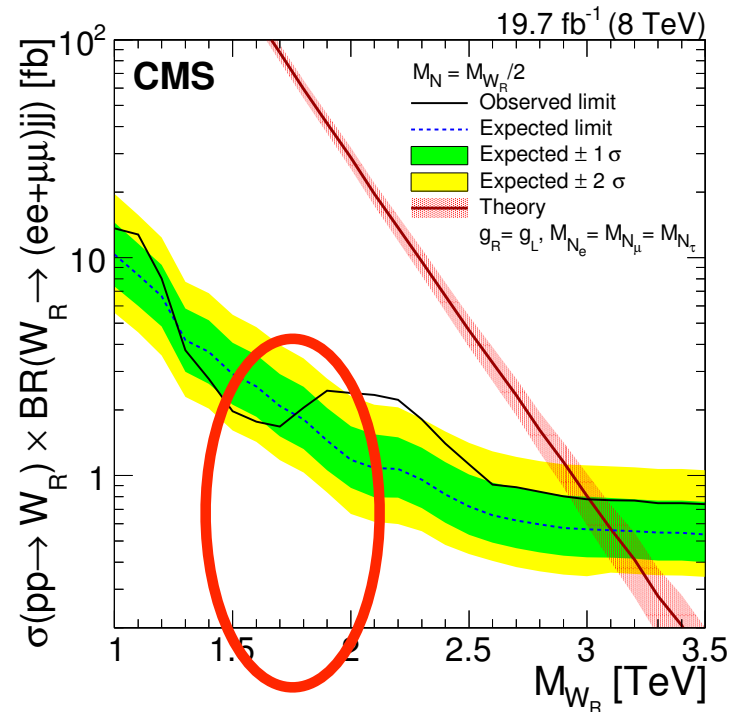
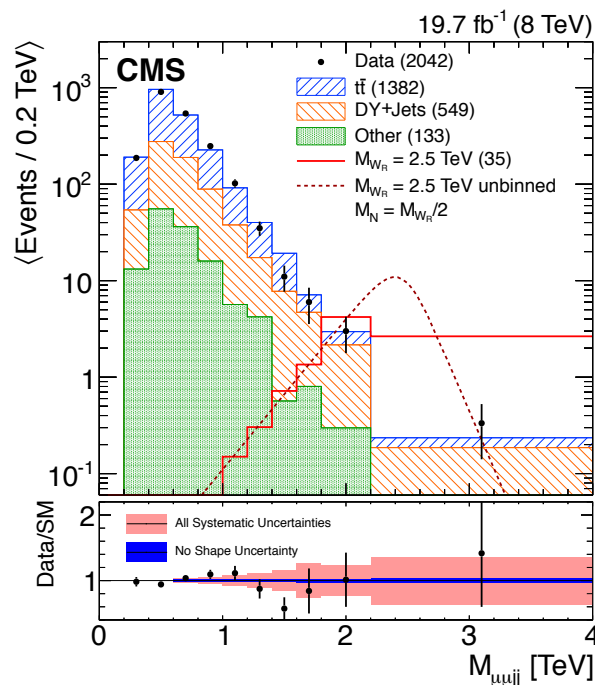
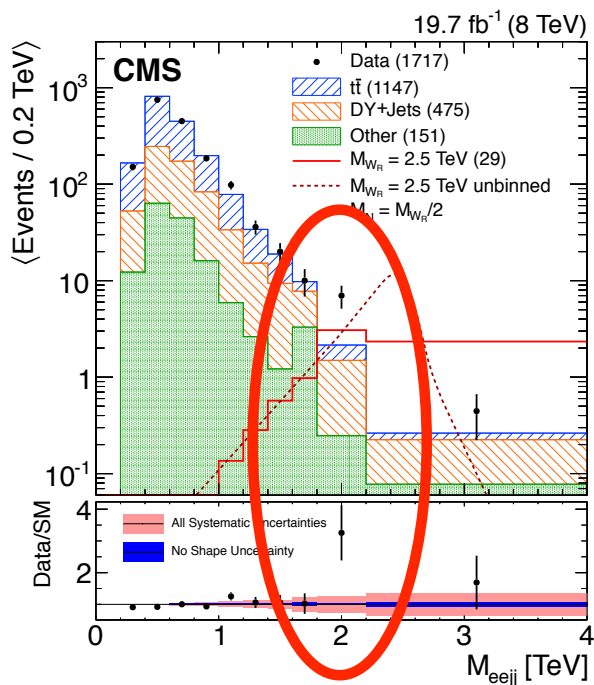


# Speculating...

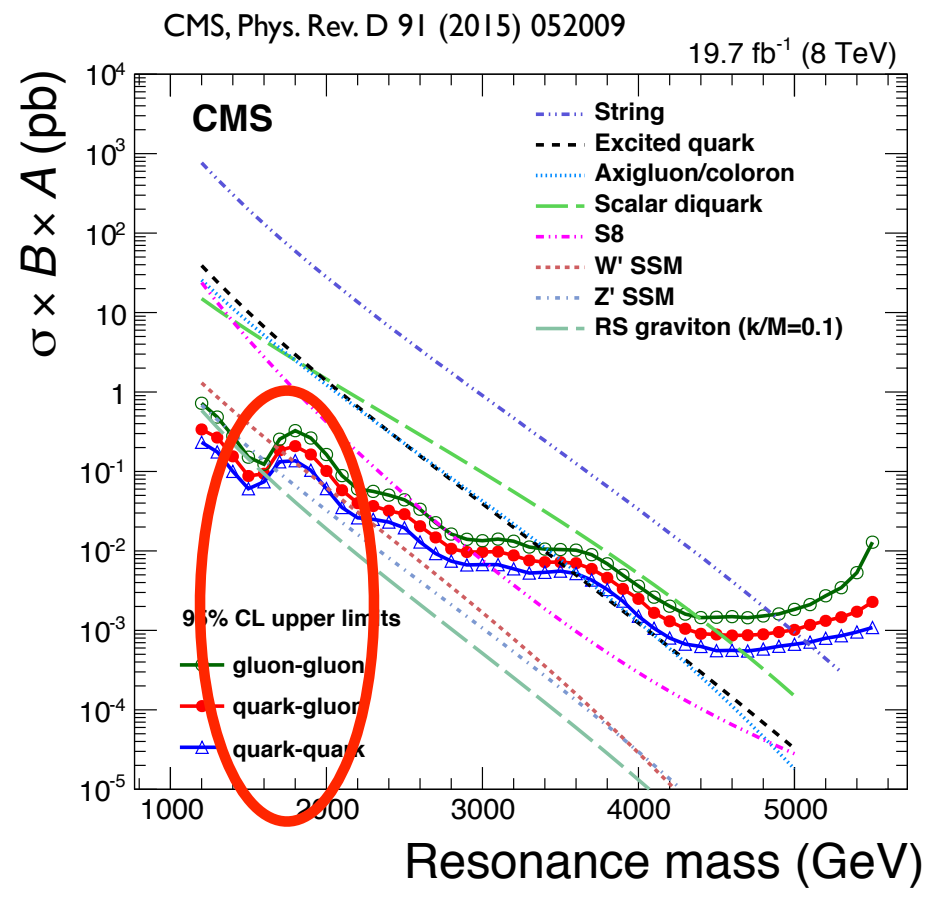
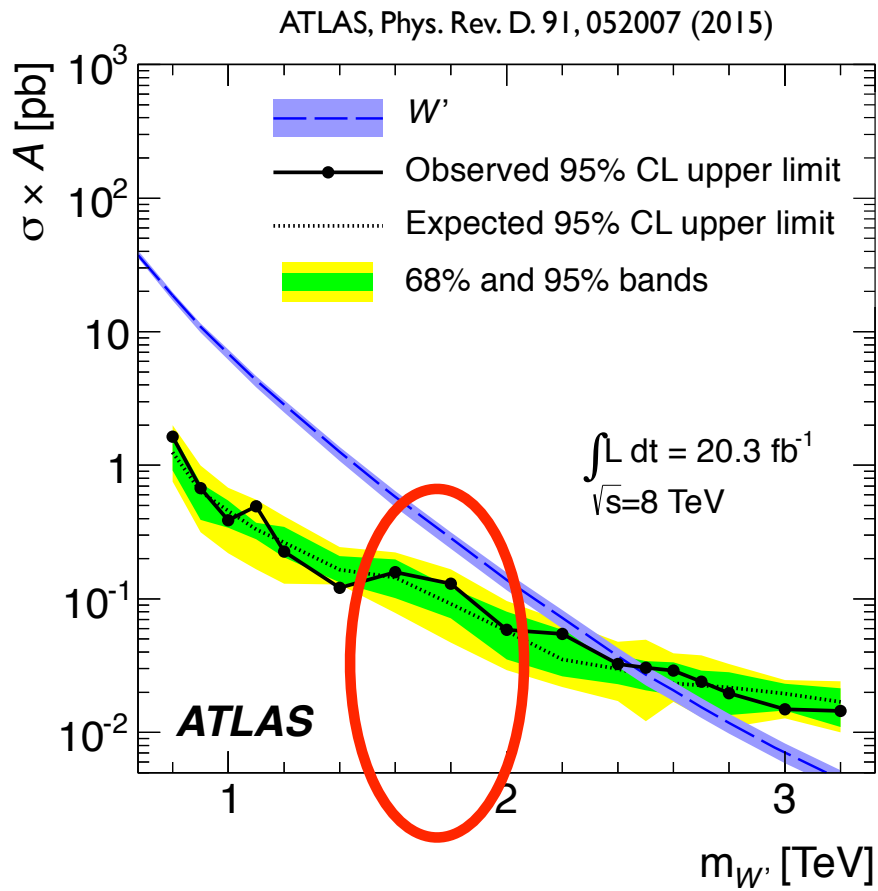
## ❖ Hints for a $W'$ ?

- If it's  $W_R$  it might decay via  $\nu_R$ :  $W_R \rightarrow \ell \nu_R^{(*)} \rightarrow \ell(\ell W_R^*) \rightarrow \ell(\ell(q\bar{q}'))$
- Dilepton + two jets final state

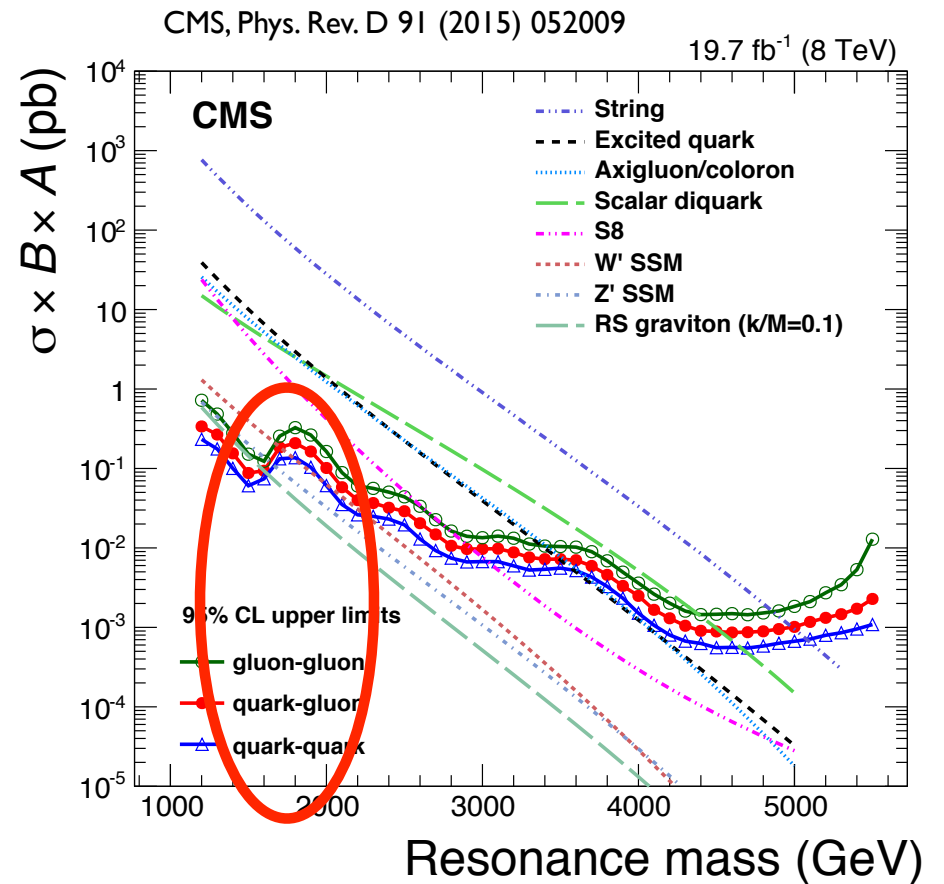
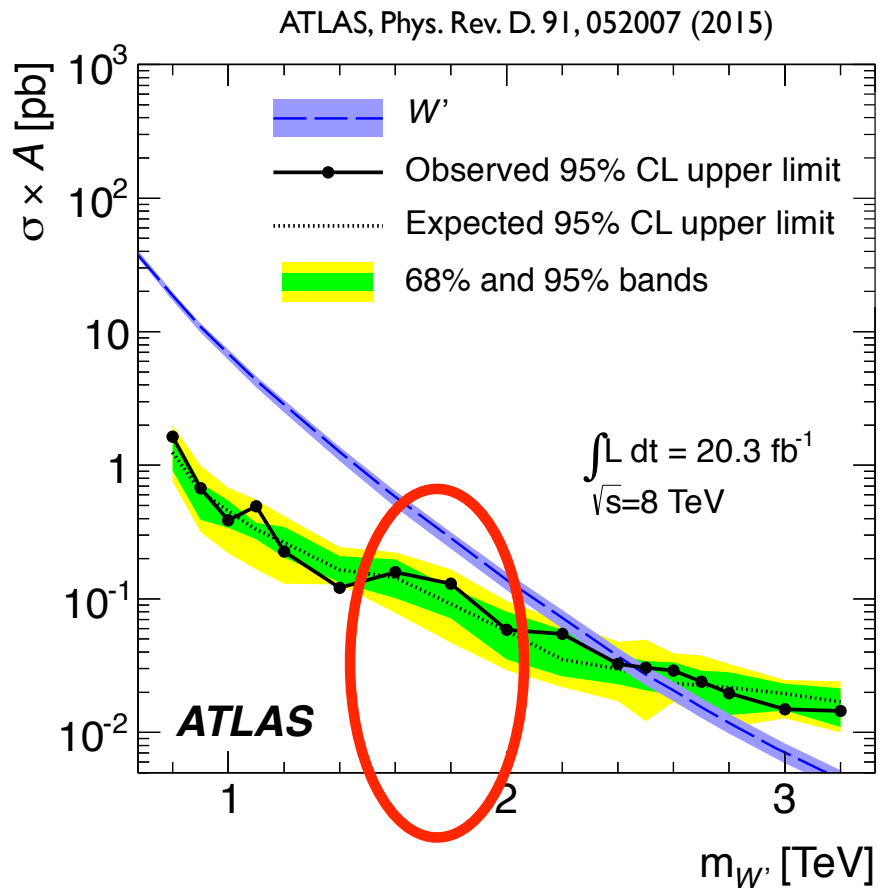
CMS, Eur. Phys. J. C 74 (2014) 3149



# And Dijets



# And Dijets

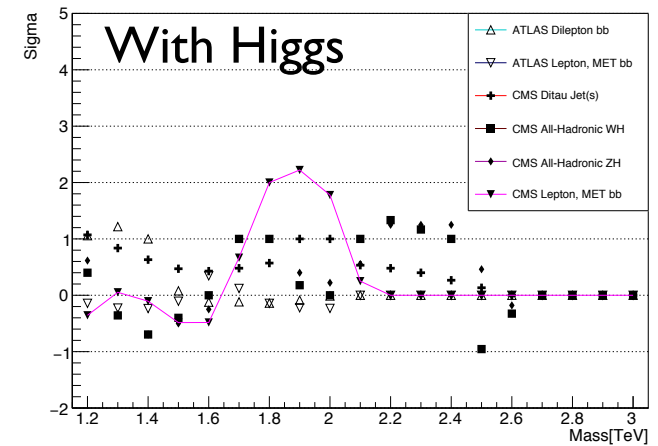
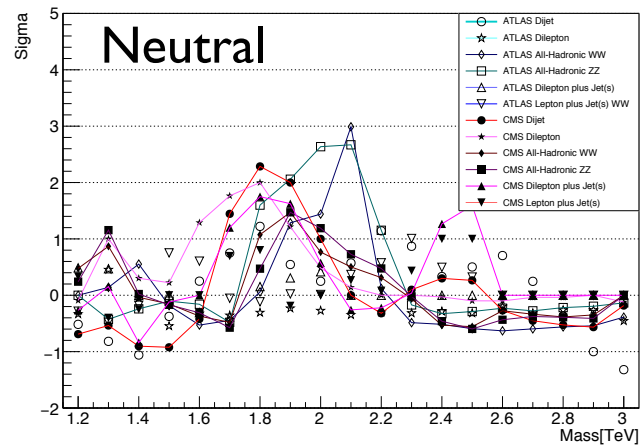
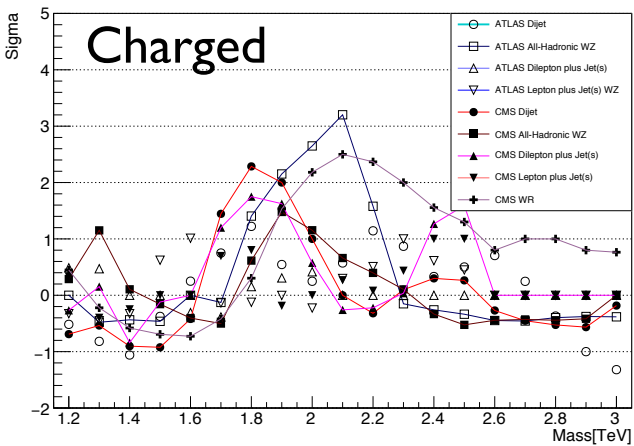


➔ Intriguing number of 1-2 sigma excesses at  $\sim 1.9$  TeV (mostly at edge of kinematic range):

- 2 in ATLAS: all-hadronic diboson, dijet
- 5 in CMS: all-hadronic,  $llj$ ,  $WH$  diboson,  $W_R$ , dijet, dilepton

# Comparing Excesses

❖ Are they all at the same mass?



❖ Somewhat... needs 2016 data

# We Should Have Excesses!

## ❖ Run 1:

- ATLAS: ~50 exotics + ~40 susy papers @ 8 TeV
- CMS: ~50 exotics + ~25 susy papers @ 8 TeV
- Expect ~eight false  $2\sigma$  (excesses+deficits) and maybe one  $3\sigma$

## ❖ Expect ~four $>2\sigma$ and $0.25 >3\sigma$ accidental excesses

## ❖ $3\sigma$ :

- ATLAS Z+MET

## ❖ $2+\sigma$ :

- $H \rightarrow \mu\tau$
- ~2 in same-sign dileptons
- ~4 in  $W'/Z'$ -type searches

## ❖ Count is ~double...

## ❖ ... and false positives not expected to cluster...

# Thanks

(and mainly: question everything!)