

Higher order QCD Corrections for the $t\bar{t}$ cross section

M. Czakon



TOP2016, Olomouc, Czech Republic, 19-23 September 2016

What happened since the last episode

Off-shell effects at NLO in QCD

- $t\bar{T}+\text{jet} \rightarrow$ *talk by M. Worek*
- $t\bar{T}+\text{Higgs}$

Off-shell effects at NLO in EW \rightarrow *talk by S. Pozzorini*

$t\bar{T} + 3 \text{ jets}$ at NLO in QCD \rightarrow *talk by S. Pozzorini*

Parton shower matching with improved resonance treatment

Single-top at NNLO in QCD in the NWA \rightarrow *talk by F. Tramontano*

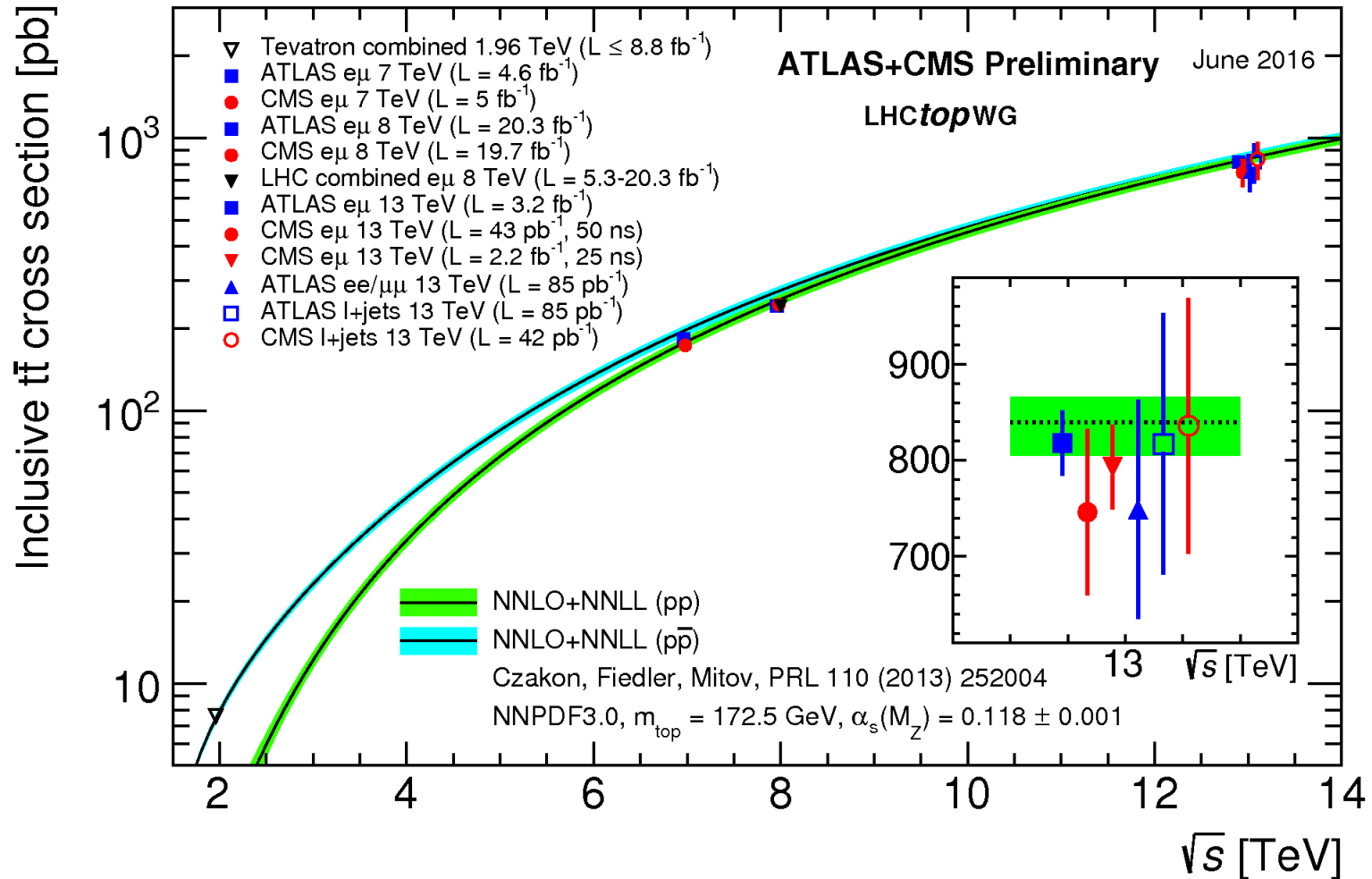
Four-loop relation between the $\overline{\text{MS}}$ and on-shell mass definitions

EW corrections at NLO with photon PDF contributions \rightarrow talk by I. Tsinikos

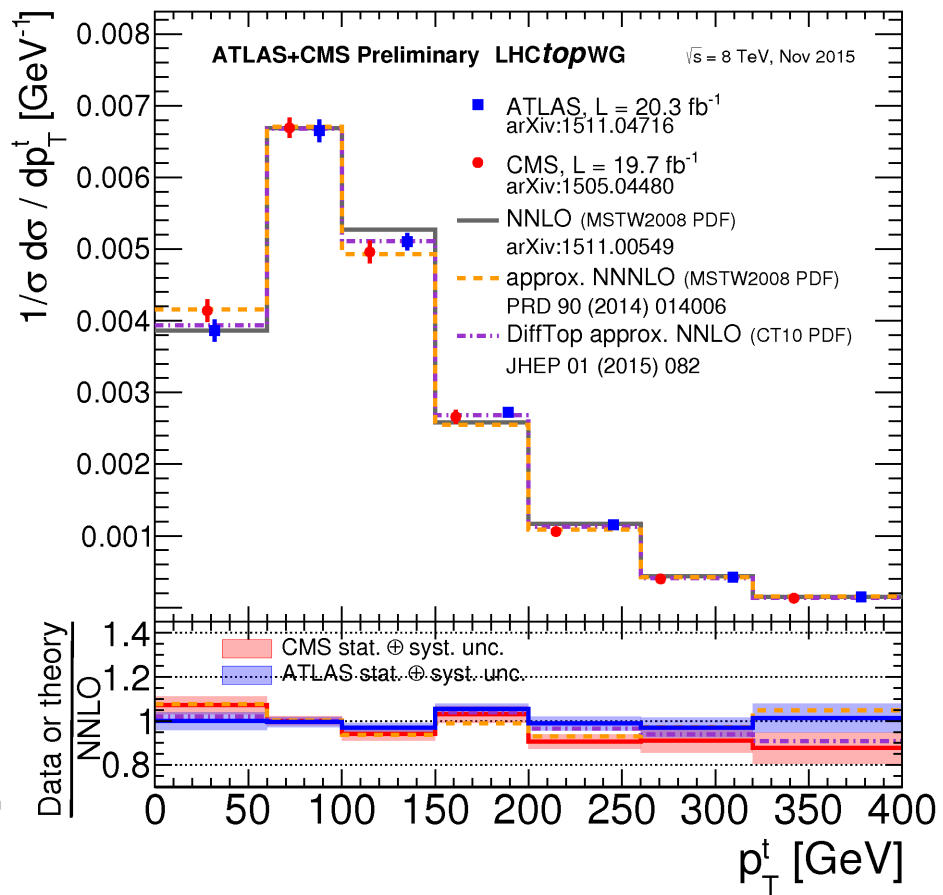
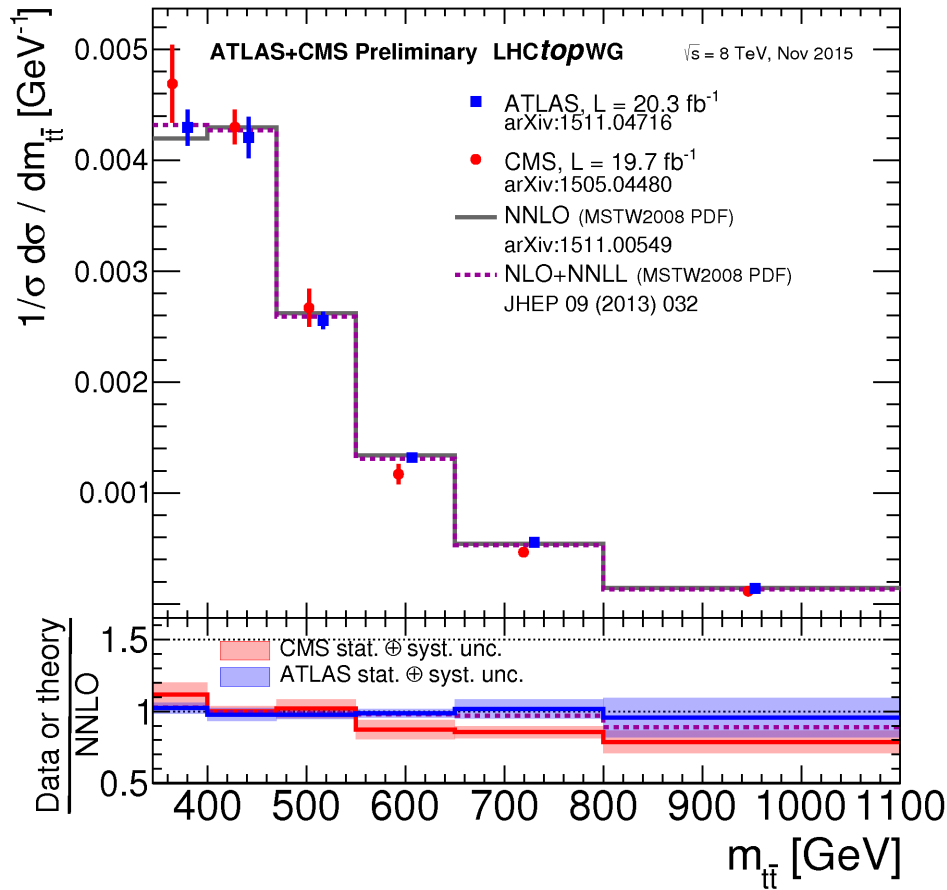
Boosted-top resummation

NNLO differential distributions with dynamical scales

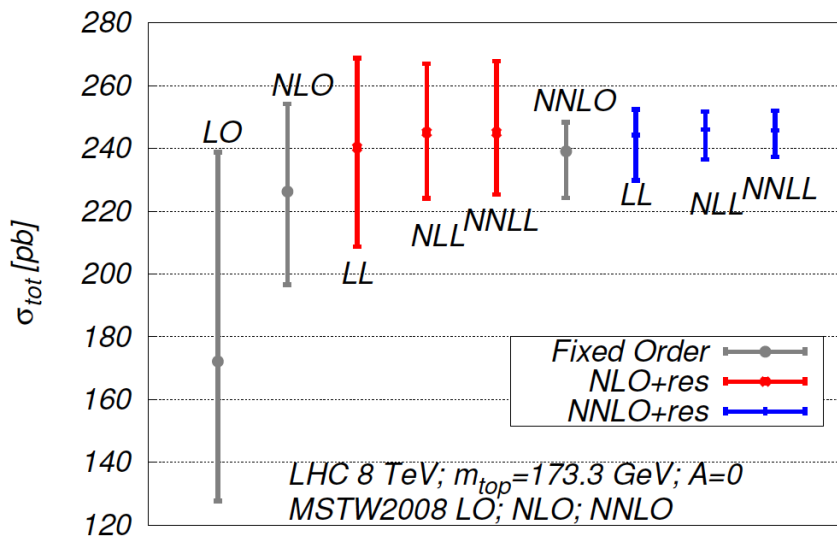
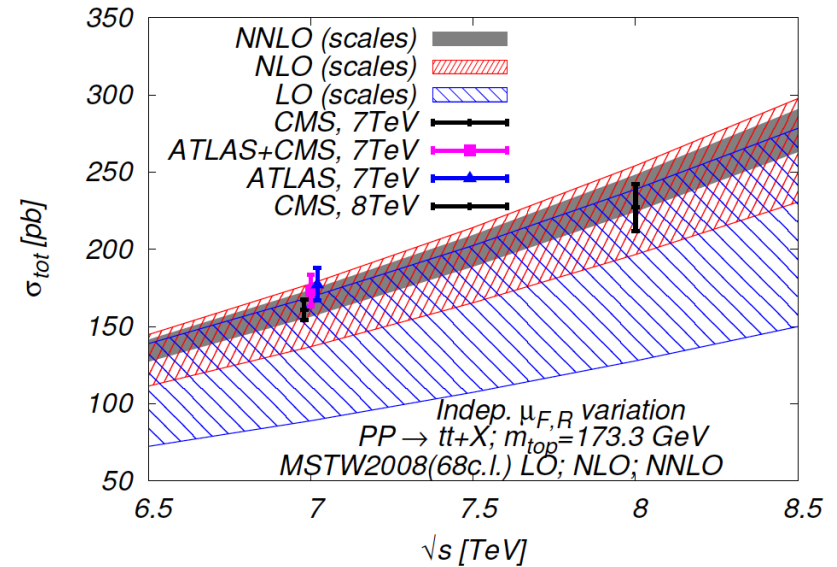
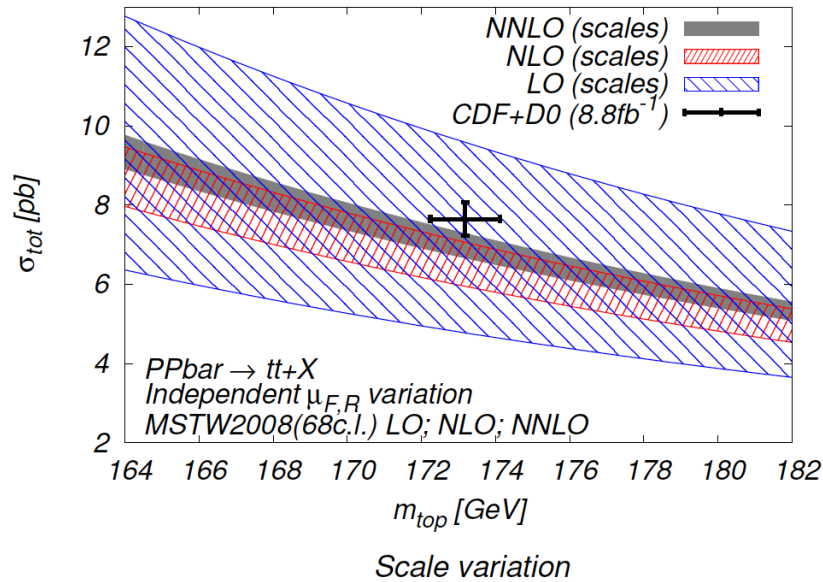
Total Cross Sections



Differential Cross Sections



Perturbation Theory Convergence



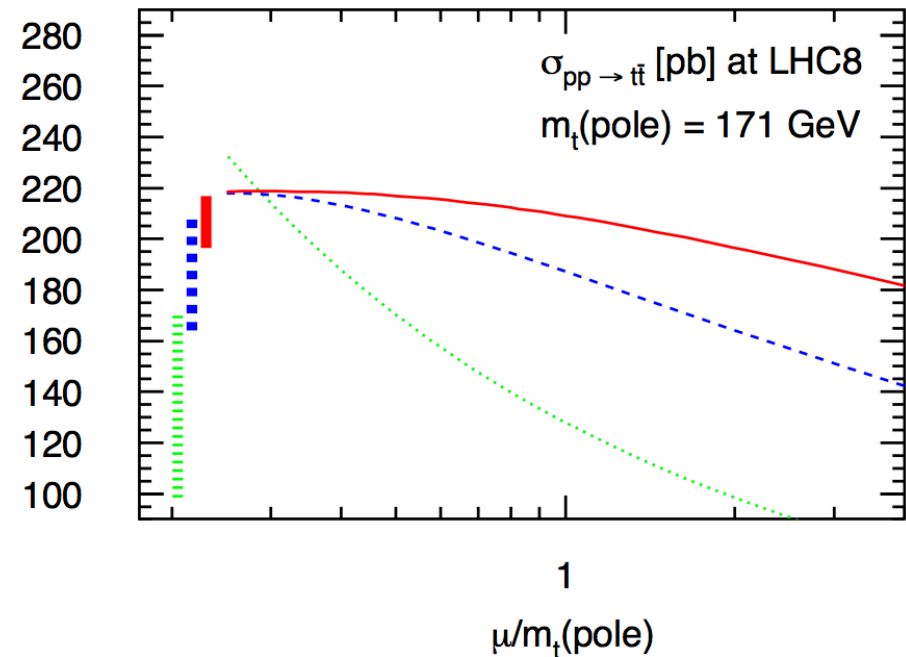
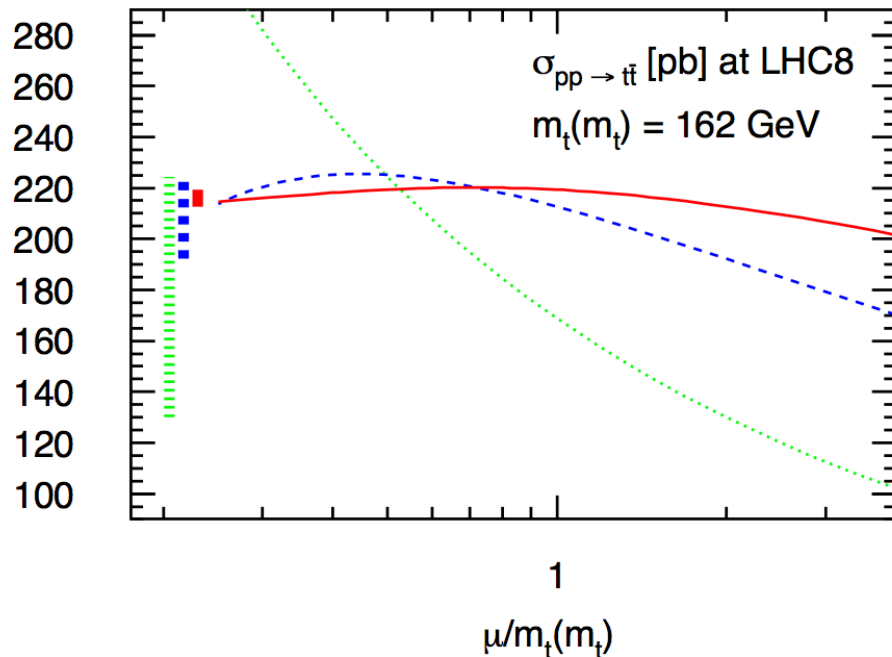
Concurrent uncertainties:

- Scales ~ 3%
- pdf (at 68%cl) ~ 2-3%
- α_S (parametric) ~ 1.5%
- m_{top} (parametric) ~ 3%

Soft gluon resummation makes a difference: **5% \rightarrow 3%**

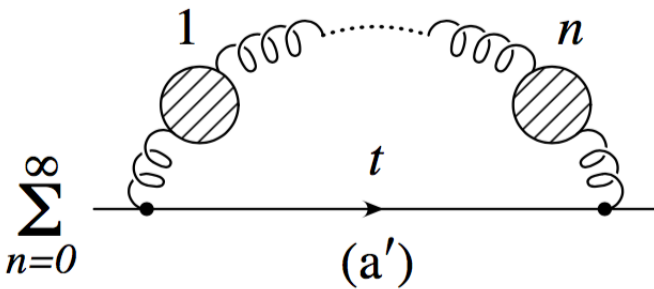
Perturbation Theory Convergence

- It has been argued that it is better to use the $\overline{\text{MS}}$ mass to improve convergence
- Is there a better scale in the on-shell scheme?
- Relevant for differential Monte Carlo description



Ambiguity of the Pole Mass

- Pole mass defined by an asymptotic series



$$\sum_{n=0}^{\infty} \Sigma^{(1)}(m_R, a) = \frac{16m_R}{3\beta_0} \sum_{n=0}^{\infty} c_n a^{n+1}$$

$$a \equiv \frac{\beta_0 \alpha_s(m_R)}{4\pi}$$

$$c_n \xrightarrow{n \rightarrow \infty} e^{-C/2} 2^n n!$$

- **Renormalon ambiguity: the series is not Borel summable**
- Ambiguity proportional to Λ_{QCD} , but with what coefficient?
- Relation to $\overline{\text{MS}}$ mass up to 4-loops

$$m_P = 163.643 + 7.557 + 1.617 + 0.501 + (0.195 \pm 0.005) \text{ GeV}$$

Marquard, Smirnov, Smirnov, Steinhauser '15

- Most recent estimate of the ambiguity

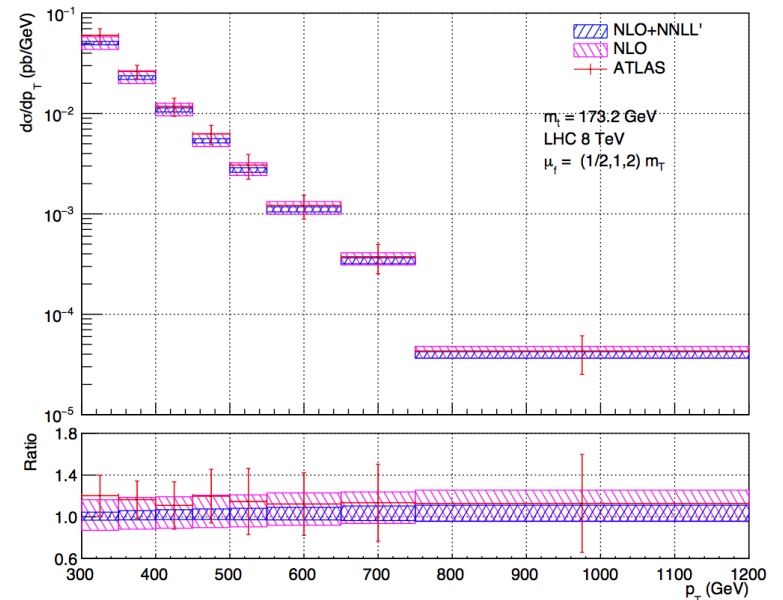
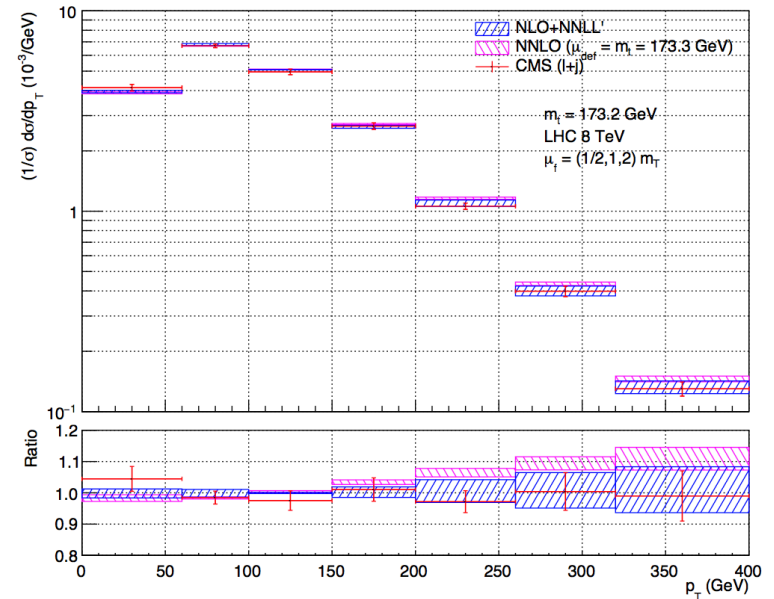
$$\delta^{(5+)} m_P = 0.250_{-0.038}^{+0.015} (N) \pm 0.001 (c_4) \pm 0.010 (\alpha_s) \pm 0.071 (\text{ambiguity}) \text{ GeV}$$

Beneke, Marquard, Nason, Steinhauser arXiv:1605.03609

Boosted Top Resummation

- Soft-gluon resummation on top of top-quark fragmentation
- Transverse momentum distribution modified by dynamical scales and resummation
- At low p_T better description of CMS data, slightly worse for ATLAS (not shown)
- Larger scale dependence?

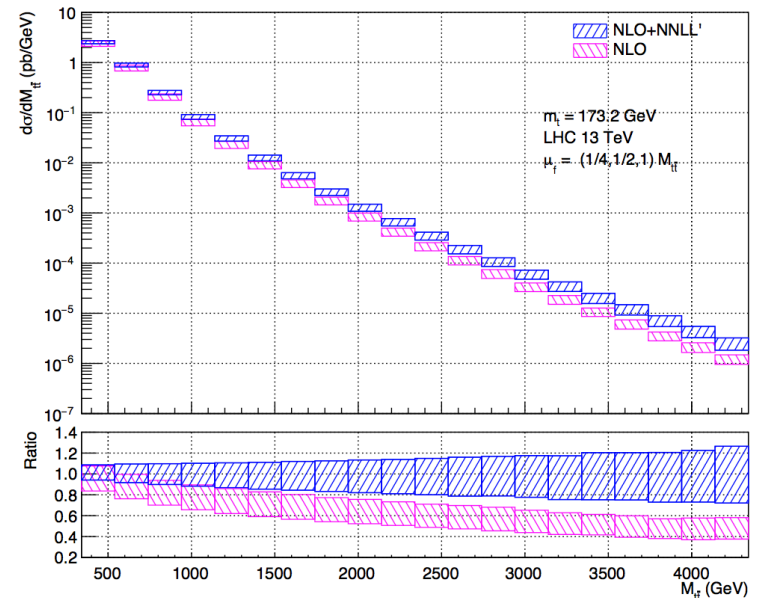
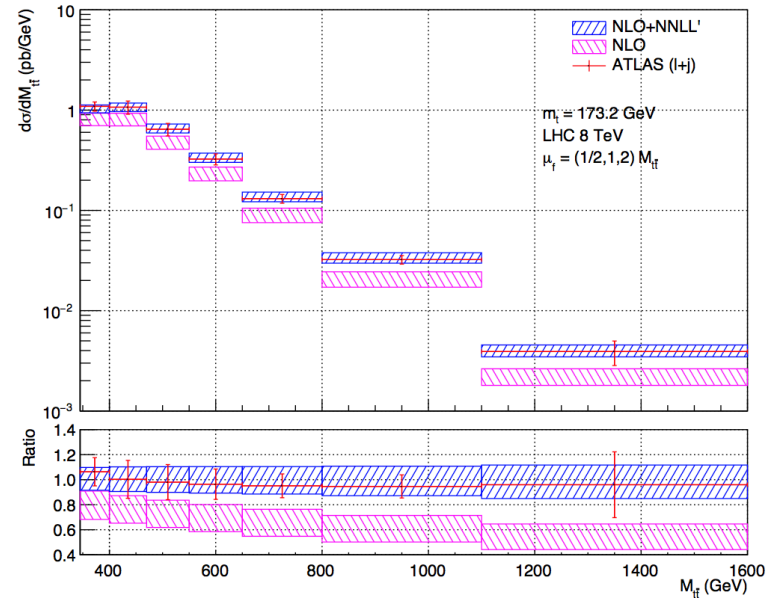
Pecjak, Scott, Wang, Yang '15



Boosted Top Resummation

- Observable dependent scale
- Results presented for 13 TeV as well
- At some point consistent matching to NNLO will become necessary
- When is true resummation needed?

Pecjak, Scott, Wang, Yang '15

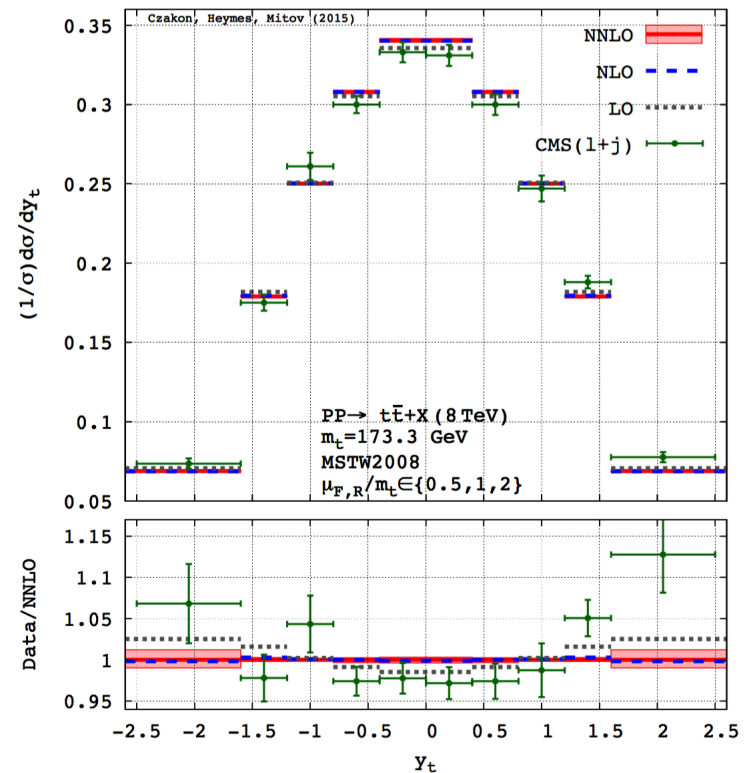
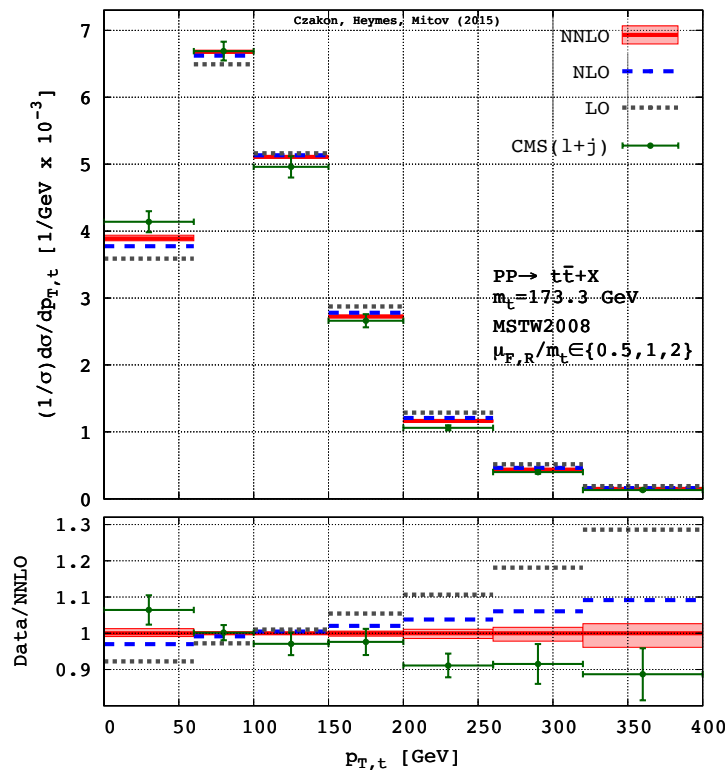


Differential Distributions

Typical differential distributions are:

1. transverse momentum of the top-quark and the top-quark pair
2. rapidity of the top-quark and the top-quark pair
3. invariant mass of the top-quark pair

Difference between normalized and absolute distributions

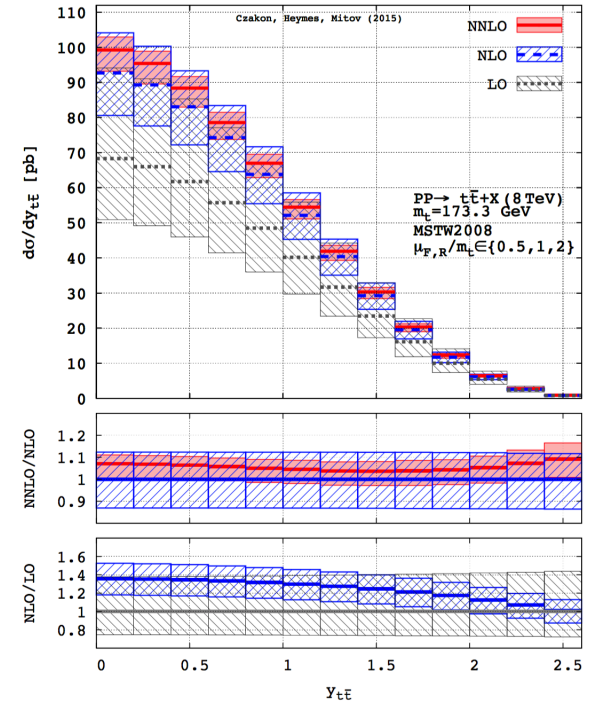
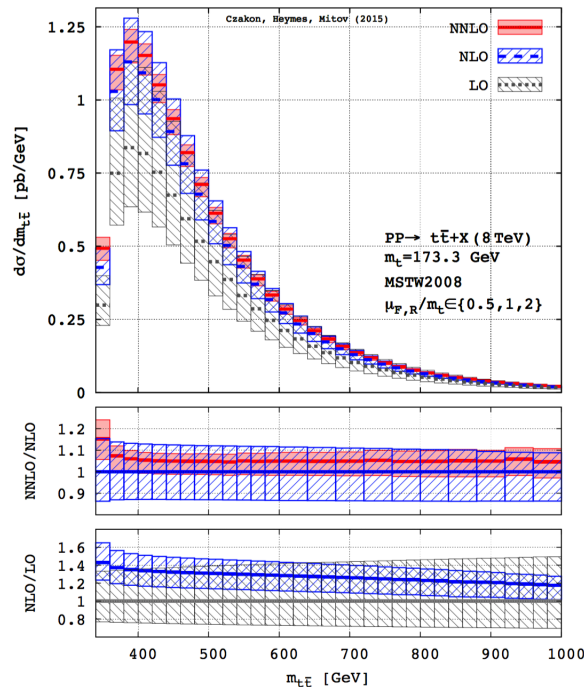
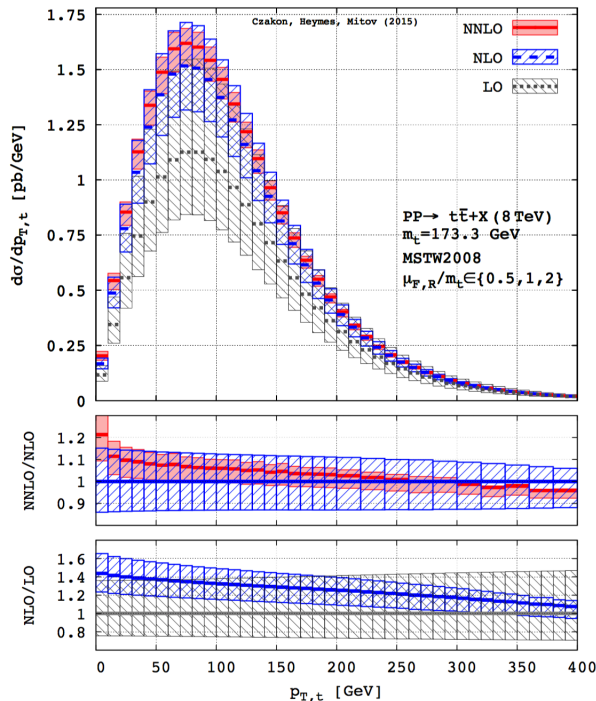


Differential Distributions

Typical differential distributions are:

1. transverse momentum of the top-quark and the top-quark pair
2. rapidity of the top-quark and the top-quark pair
3. invariant mass of the top-quark pair

Difference between normalized and absolute distributions



Dynamical Scales

- Over extended kinematical ranges it is necessary to use dynamical scales
- Examples in the case of top-quark pair production:

$$\mu_0 \sim m_t ,$$

$$\mu_0 \sim m_T = \sqrt{m_t^2 + p_T^2} ,$$

$$\mu_0 \sim H_T = \sqrt{m_t^2 + p_{T,t}^2} + \sqrt{m_t^2 + p_{T,\bar{t}}^2} ,$$

$$\mu_0 \sim H'_T = \sqrt{m_t^2 + p_{T,t}^2} + \sqrt{m_t^2 + p_{T,\bar{t}}^2} + \sum_i p_{T,i} ,$$

$$\mu_0 \sim E_T = \sqrt{\sqrt{m_t^2 + p_{T,t}^2} \sqrt{m_t^2 + p_{T,\bar{t}}^2}} ,$$

$$\mu_0 \sim H_{T,\text{int}} = \sqrt{(m_t/2)^2 + p_{T,t}^2} + \sqrt{(m_t/2)^2 + p_{T,\bar{t}}^2} ,$$

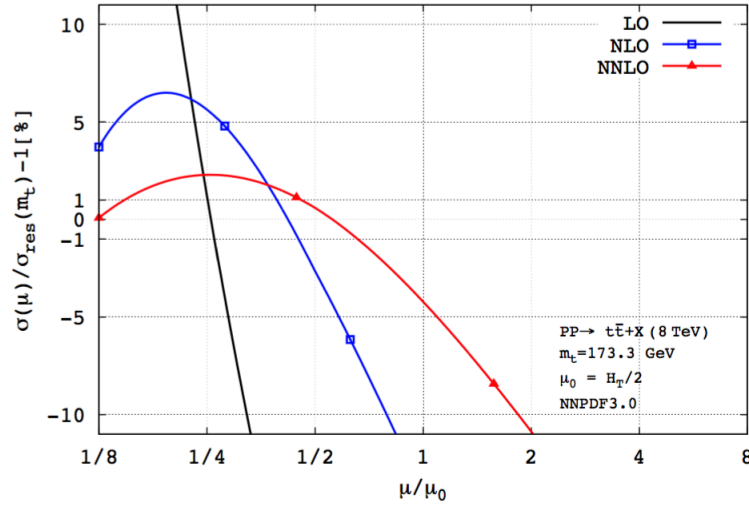
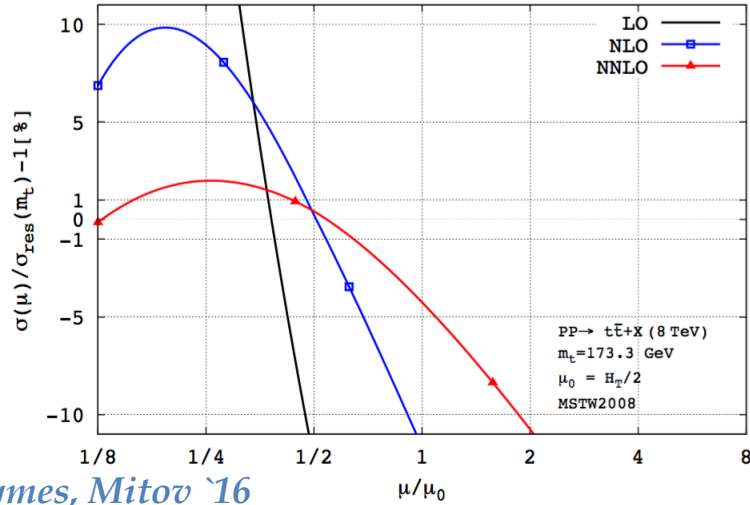
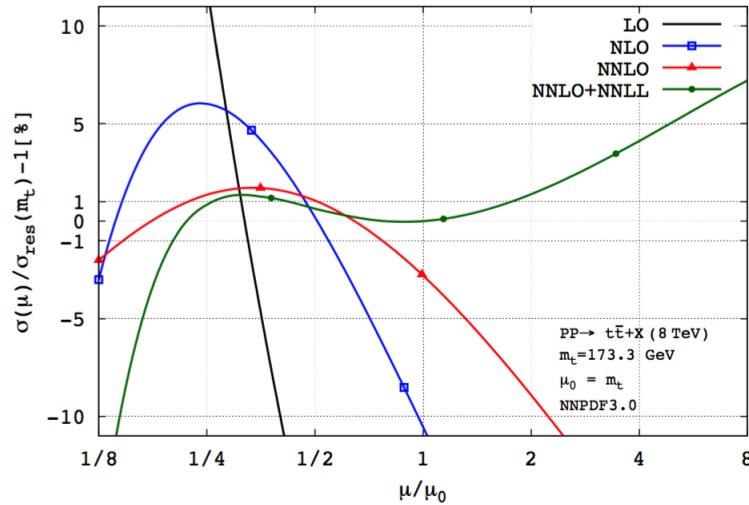
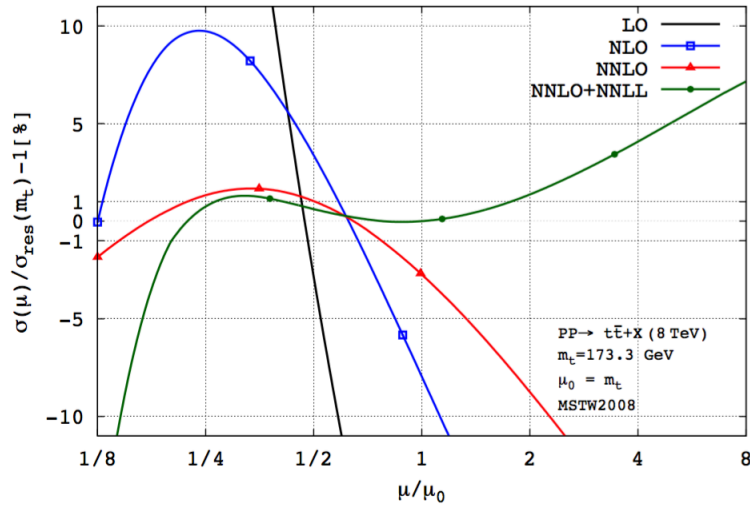
$$\mu_0 \sim m_{t\bar{t}} ,$$

Our recommendation for p_T (but $1/2$)

Our recommendation for the others (but $1/4$)

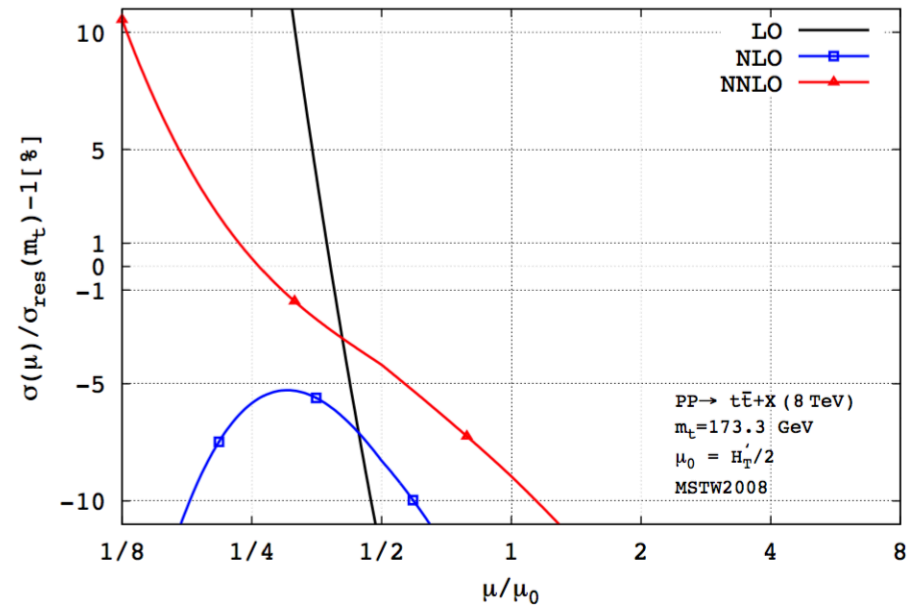
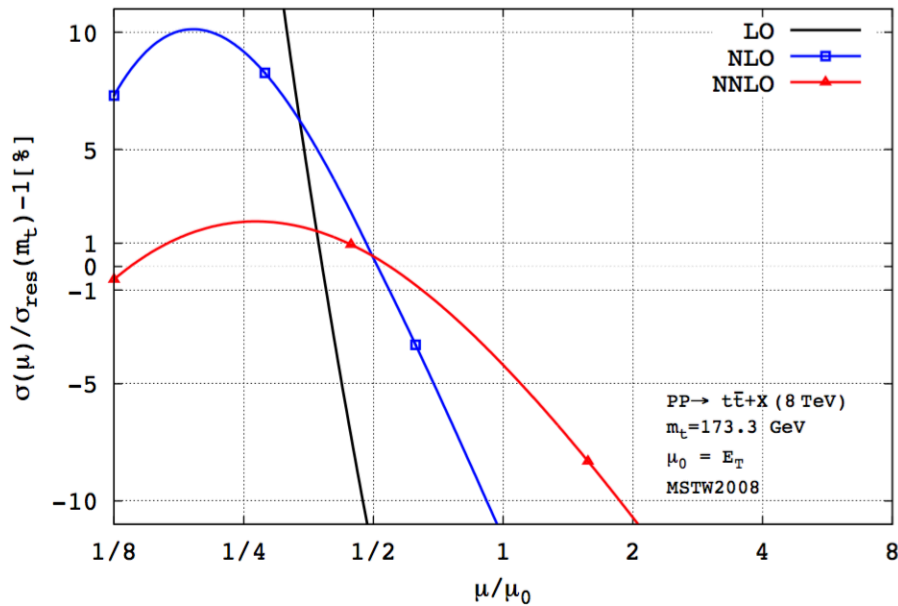
Dynamical Scales

- Dynamical scales modify the total cross section
- Because of threshold enhancement close results from an “average” fixed scale



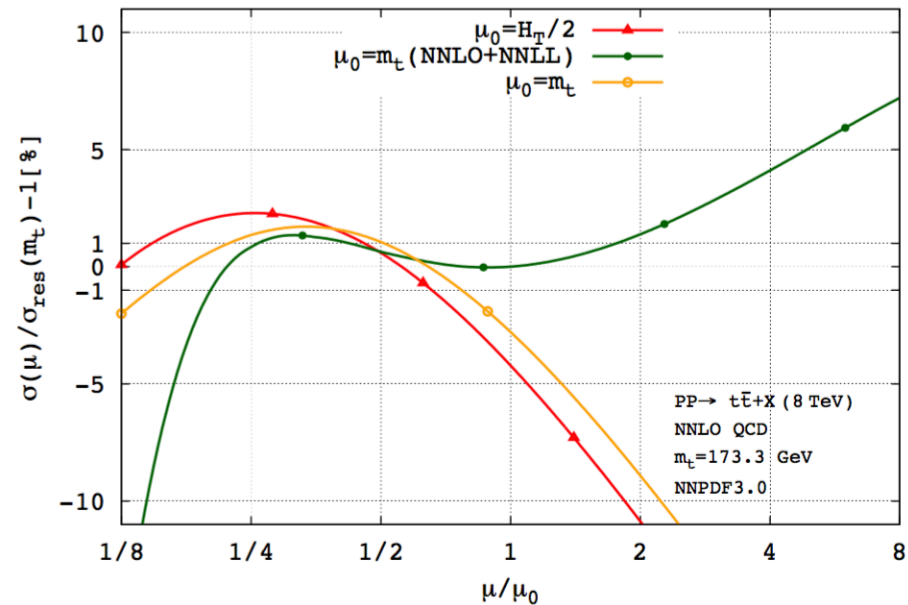
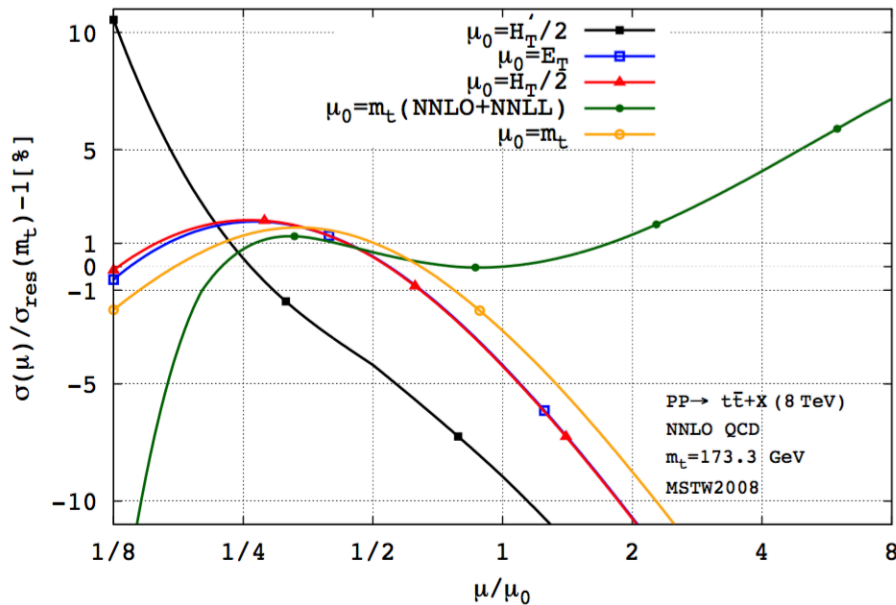
Dynamical Scales

- Some scales behave suspiciously, while seeming perfectly reasonable



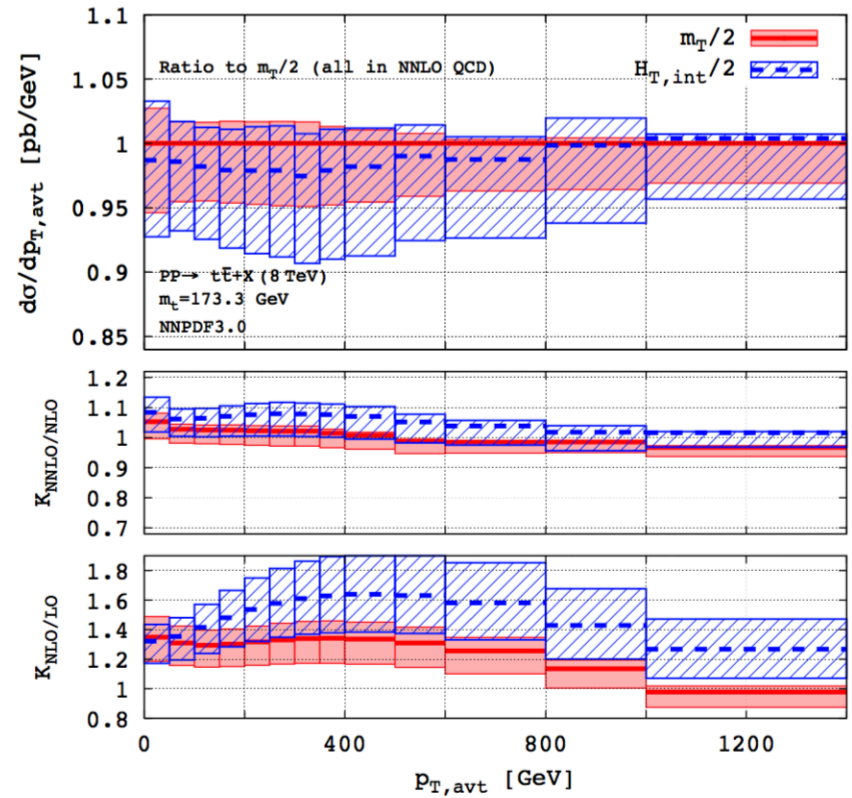
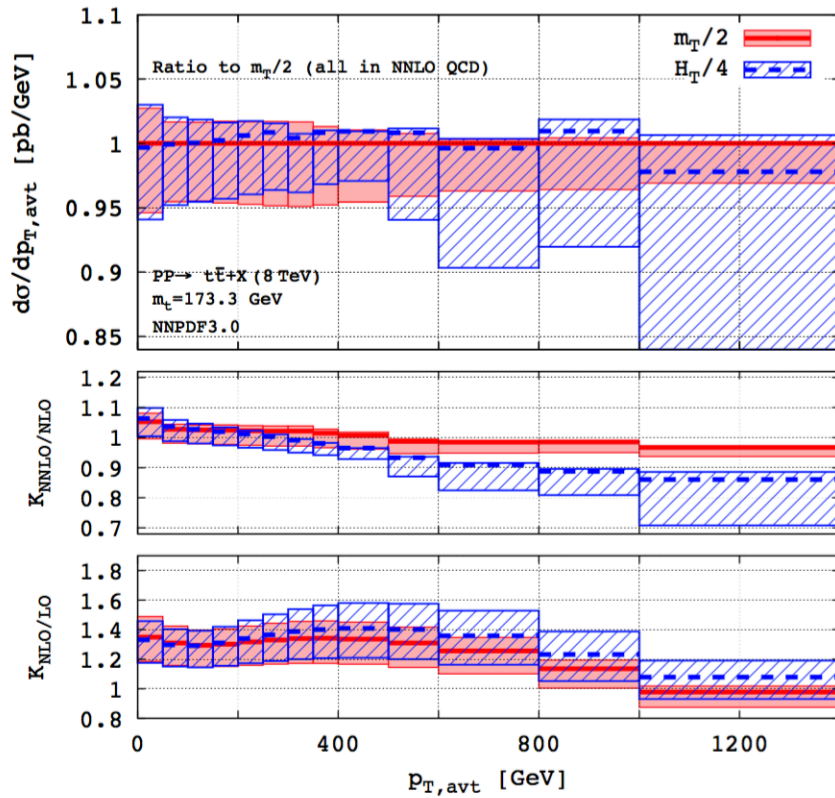
Dynamical Scales

- A comparison of different scales at highest precision
- Different PDF sets



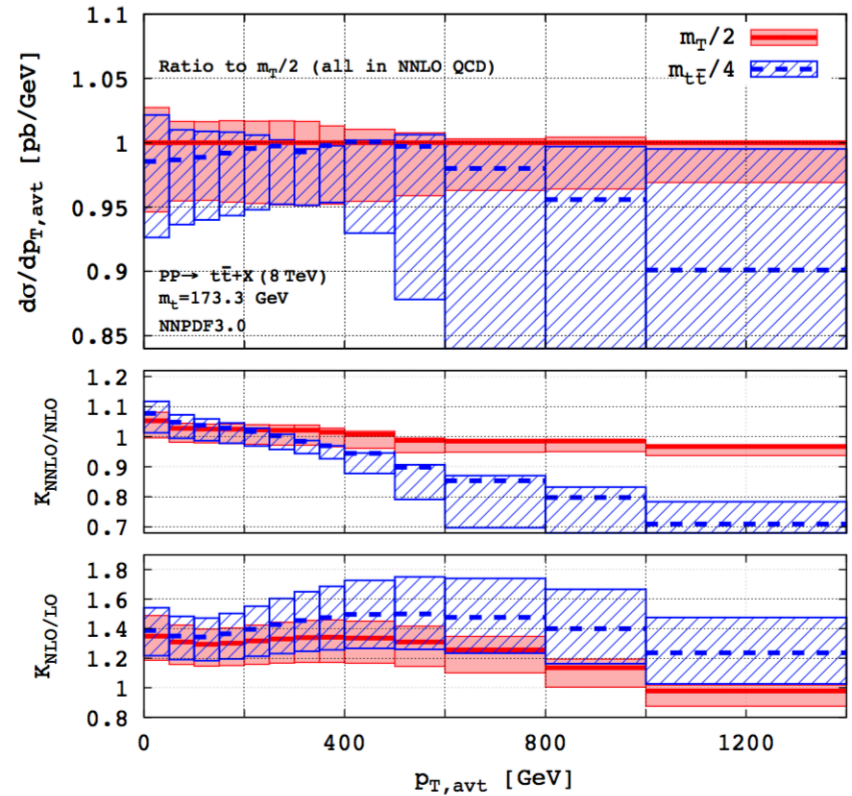
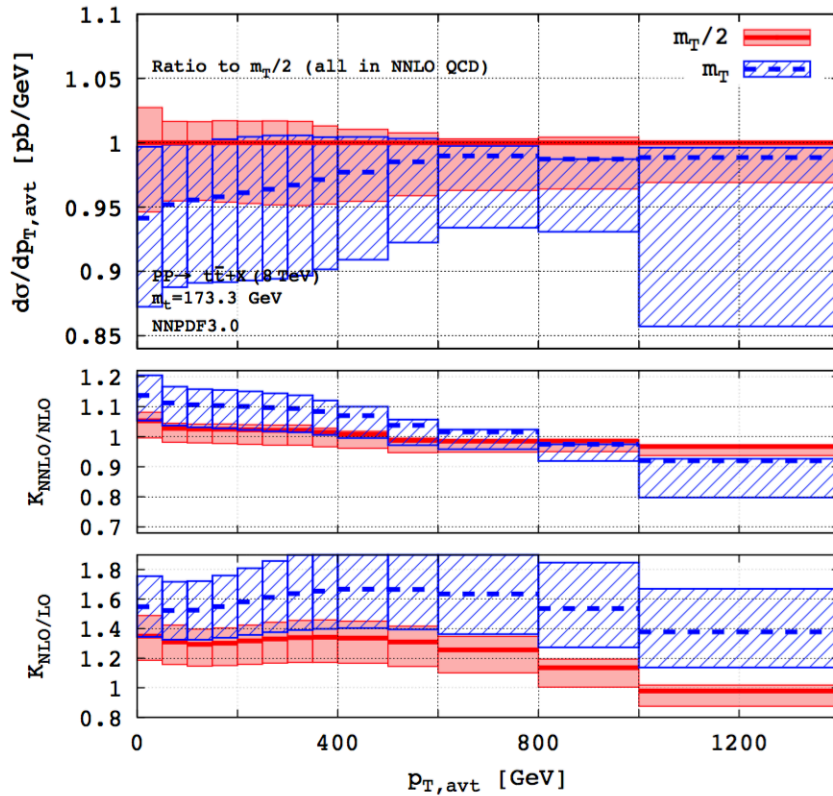
Dynamical Scales

- Improvements of convergence with “reasonable” scales



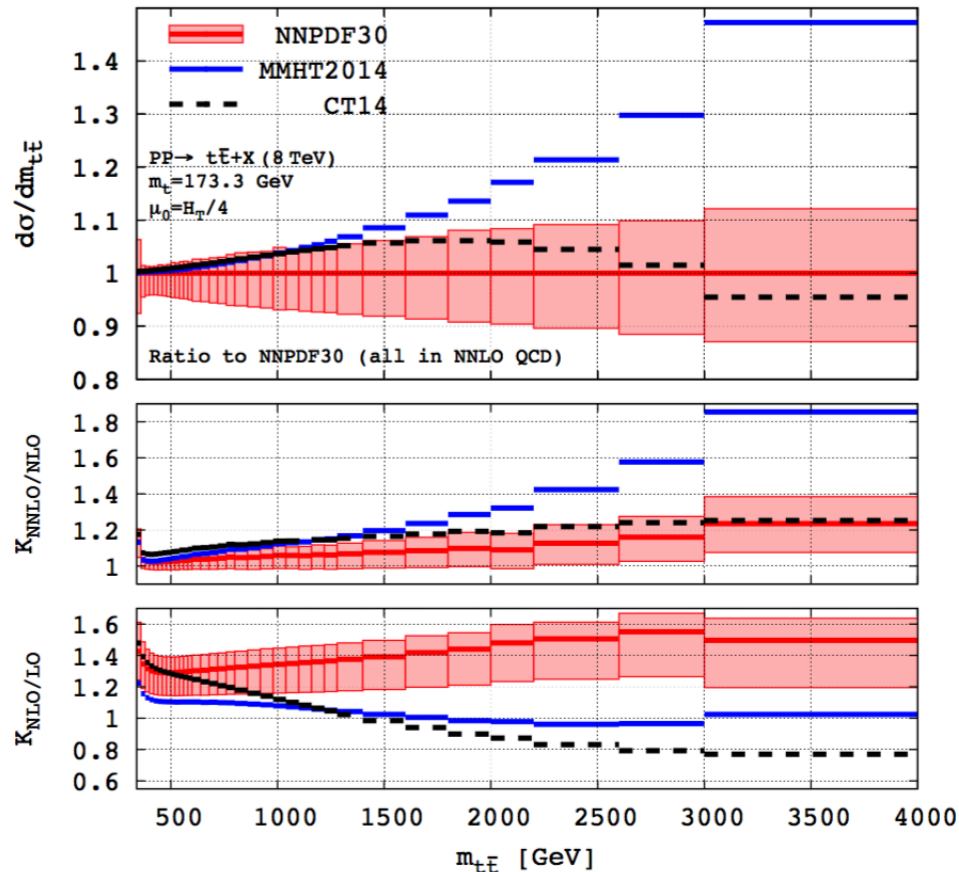
Dynamical Scales

- Improvements of convergence with “reasonable” scales
- Problems in the case of “less reasonable” scales



Reliability of PDF Sets

- Above a certain invariant mass no more precise predictions
- Use the distributions to improve PDFs?



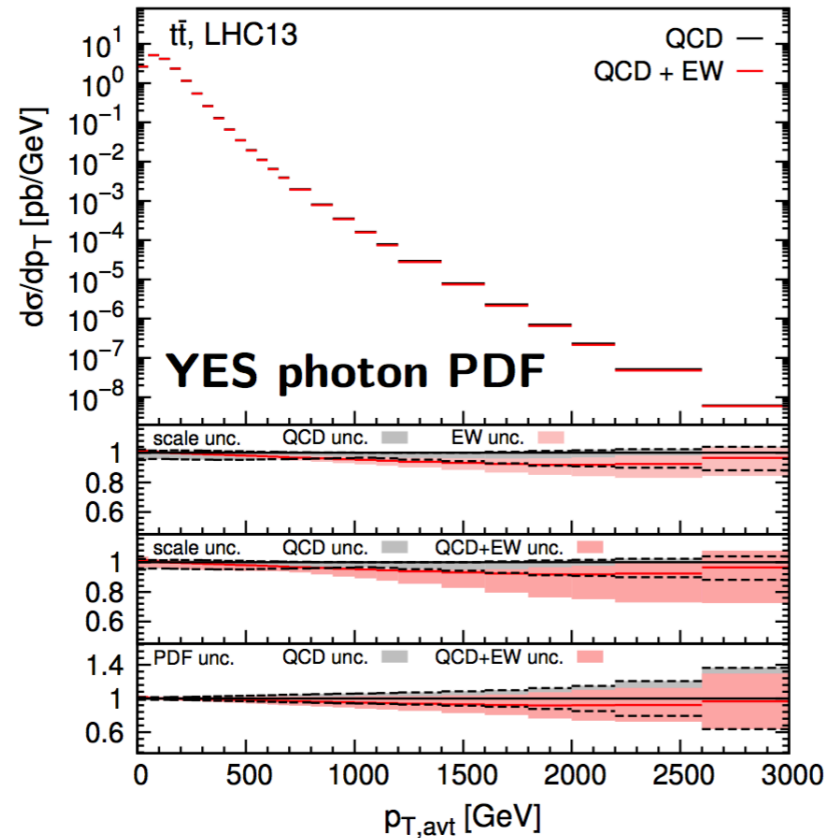
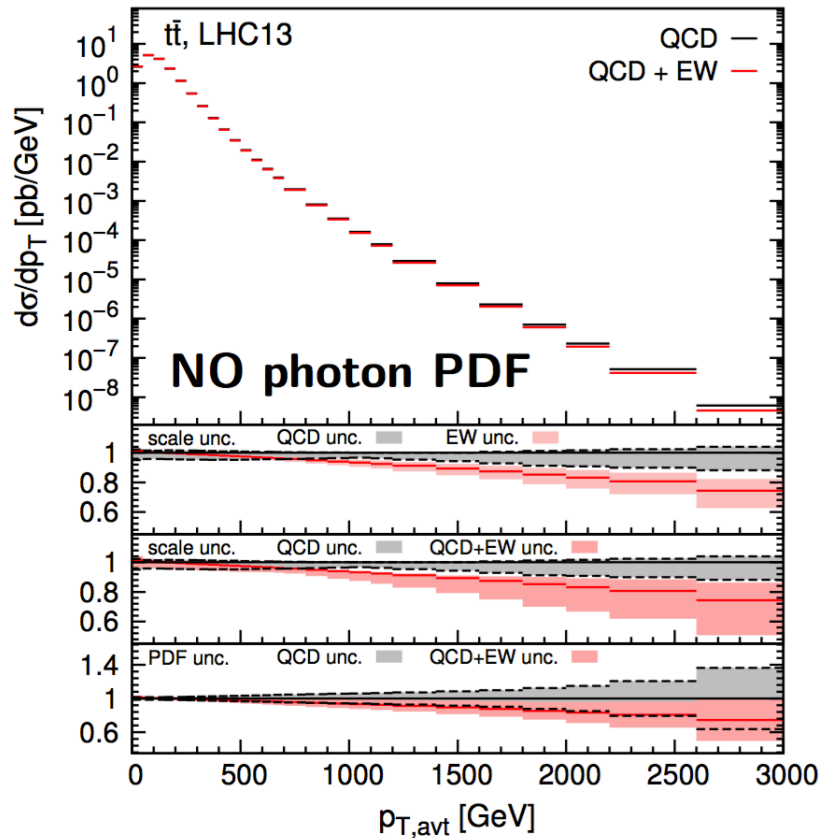
Concluding Remarks

- High precision should be associated with fixed order perturbation theory:
 - Clear advantage: not many ambiguities
 - But: beware of range of applicability
 - Currently at **next-to-next-to-leading order** for on-shell production
MC, Bärrnreuther, Fiedler, Heymes, Mitov '12 - '16
 - Partial independent results by:
Abelof, Gehrmann-De Ridder, Maierhofer, Pozzorini '14
Catani, Grazzini, Torre '14 - '15
 - Currently substantial effort to include Narrow Width Approximation

- **Advertisement**
 - Combination with electroweak corrections
see talk by I. Tsinikos on Tuesday evening

Concluding Remarks

- High precision should be associated with fixed order perturbation theory:



Preliminary: MC, D. Heymes, A. Mitov, D. Pagani, I. Tsinikos, M. Zaro

- Advertisement

- Combination with electroweak corrections
see talk by I. Tsinikos on Tuesday evening