

# Sphalerons far from equilibrium and associated phenomena

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Based on MM, S. Schlichting, R. Venugopalan  
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Stony Brook University



# Outline

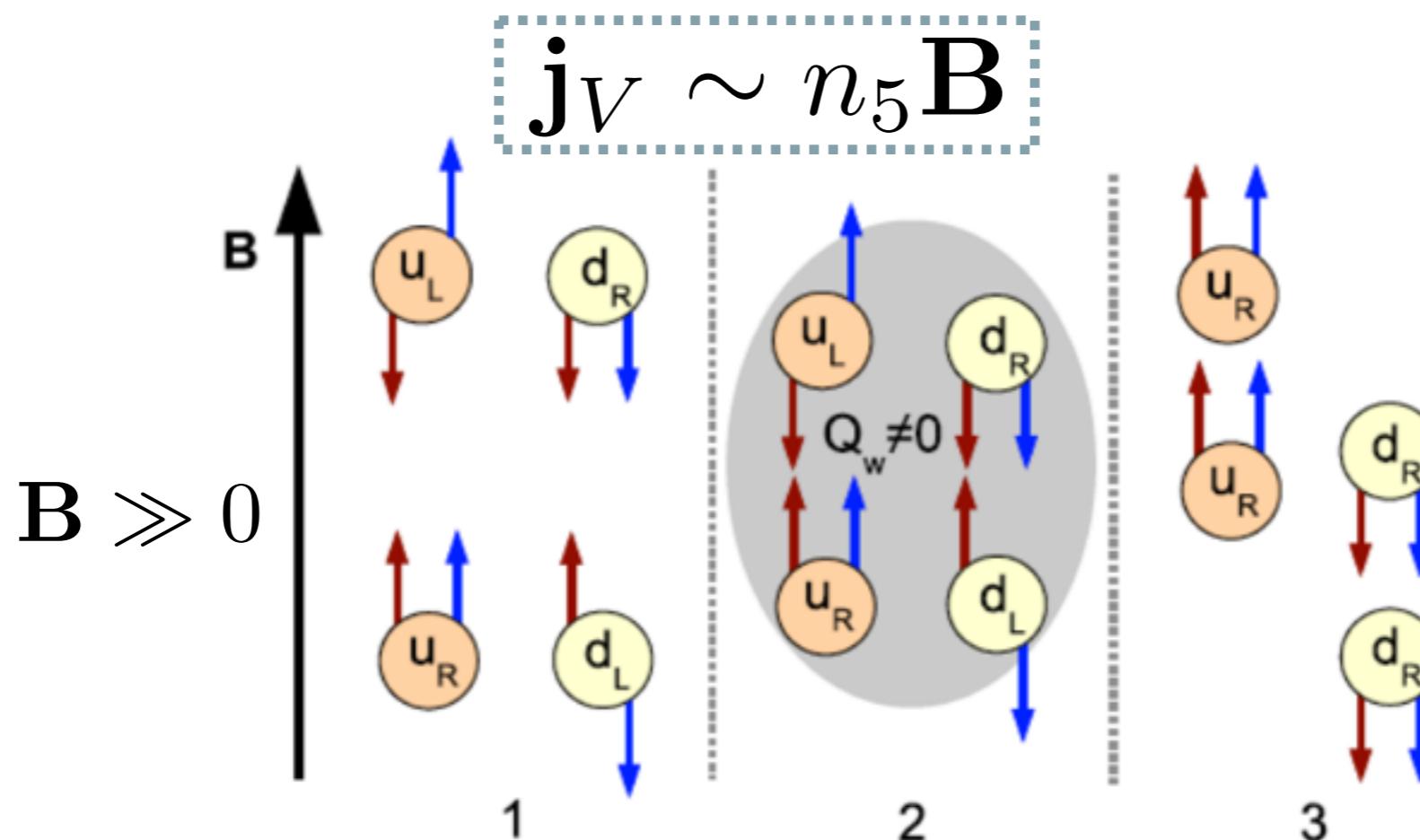
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- Motivation- CME
- Topology in real-time lattice gauge theory
- Far from equilibrium sphalerons
- Outlook to future work



# Motivation-CME

- Chiral Magnetic Effect (Kharzeev, McLerran, Warringa NPA803, 227 (2008); Fukushima, Kharzeev, Warringa PRD78 (2008) 074033)
- Vector Current= Axial charge imbalance + U(1) B-field



(Kharzeev Prog. Part. Nucl. Phys. 75 (2014) 133-151)

Red=momentum; Blue=spin;

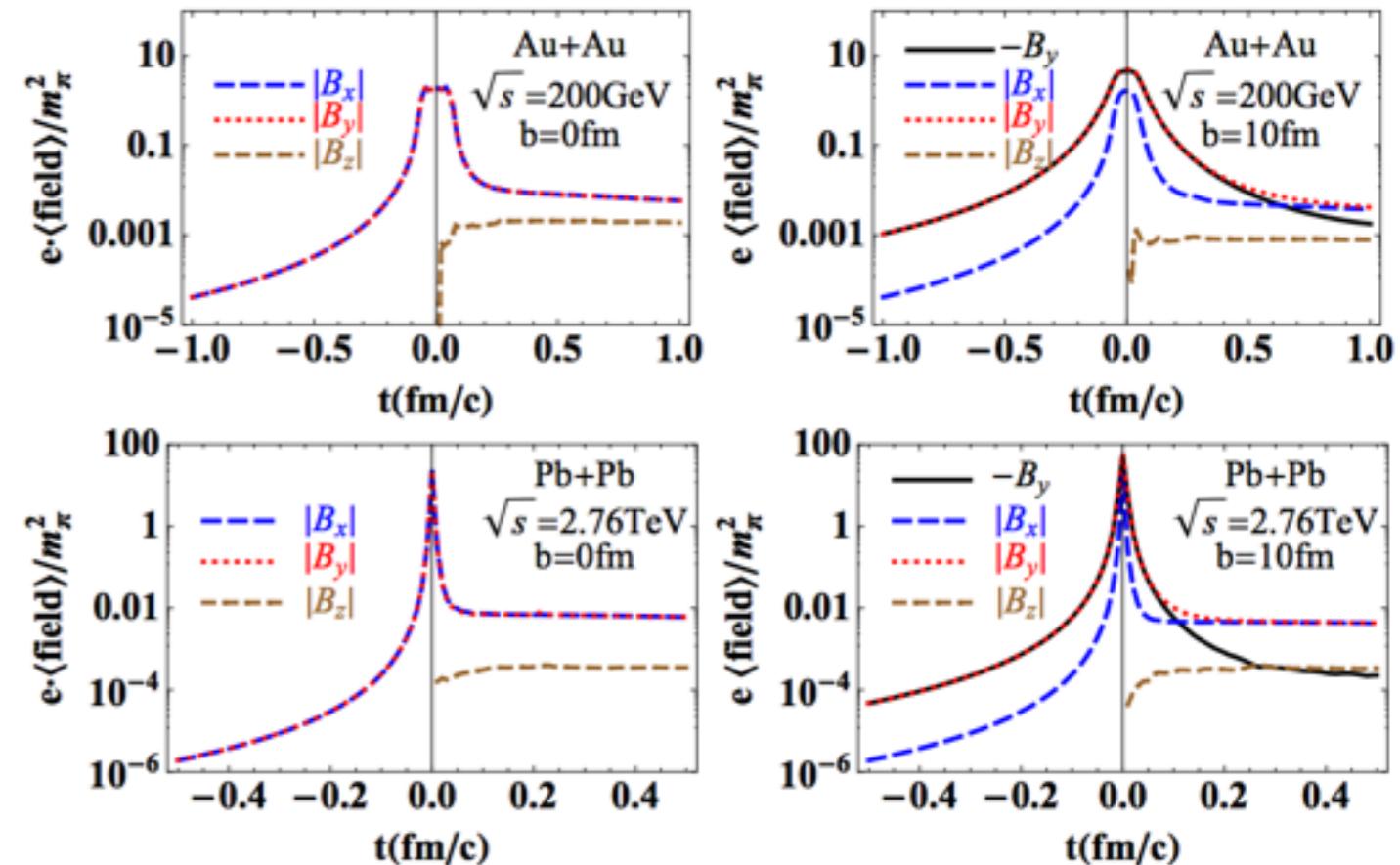
# Motivation-CME

$$\mathbf{j}_V \sim n_5 \boxed{\mathbf{B}}$$

- Magnetic field strongest at early times
  - $e\mathbf{B} \equiv (m_\pi)^2 \sim 10^{18} \text{G}$
- Dependent on impact parameter and beam energy
- Very short lifetime
  - $t \sim 0.1 - 0.2 \frac{\text{fm}}{\text{c}}$
  - may be increase by conductivity
- Much recent progress

(Deng and Huang, PRC85, 044907; McLerran and Skokov NPA929 (2014) 184-190;  
Tuchin et al PRC91 (2015) 064902, arXiv:1604.04572)

## Theoretical model calculations



(Deng and Huang, PRC85, 044907)

# Motivation-CME

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- In order to fully understand CME, need to understand initial axial charge
- Axial charge production governed by anomaly (t'Hooft; Adler, Bell, Jackiw)

$$\partial_\mu j_5^\mu = 2m_f \bar{q} i \gamma_5 q - \frac{g^2}{16\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$$

Axial current  $j_5^\mu = (n_5, \mathbf{j}_5)$

Fermion field contribution:  
~fermion mass

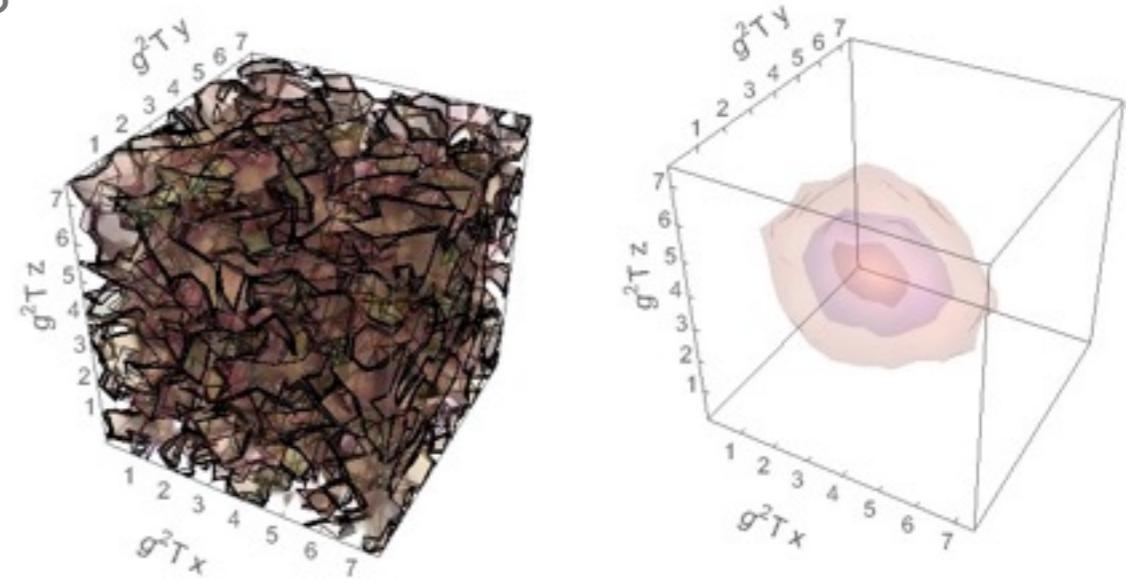
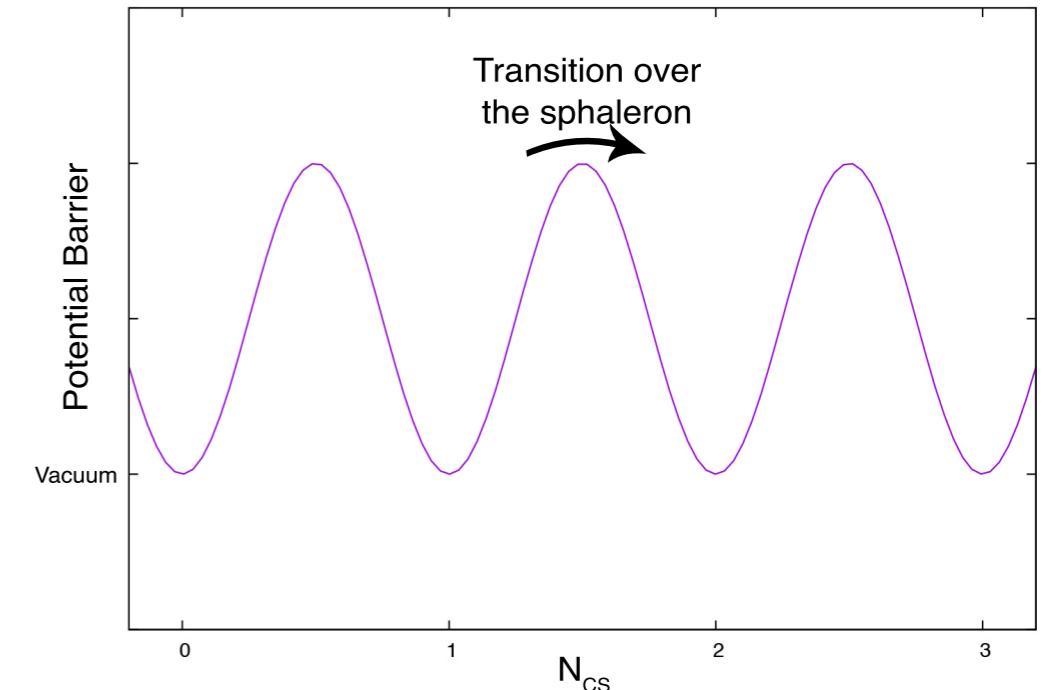
Gauge field contribution:  
receives contributions from  
topological transitions, field  
strength fluctuations, ...

The diagram illustrates the decomposition of the axial current. A horizontal line at the top represents the total axial current  $\partial_\mu j_5^\mu$ . It branches into two parts below: a vertical line representing the fermion field contribution  $2m_f \bar{q} i \gamma_5 q$ , and a diagonal line representing the gauge field contribution  $\frac{g^2}{16\pi^2} F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$ .

# Axial charge

$$\partial_\mu j_5^\mu{}_f \sim F_{\mu\nu}^a \tilde{F}_a^{\mu\nu} \propto \mathbf{E}^a \mathbf{B}^a$$

- Axial charge receives contributions
  - Topological transitions - present focus
  - Real time: sphaleron transitions  
*(Klinkhamer, Manton PRD30 (1984) 2212)*
    - Sphaleron= Greek for ‘ready to fall’
  - Field strength fluctuations



# Axial charge

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- Sphaleron transitions correspond to integer change in Chern-Simons number

$$\Delta N_{CS} = \frac{g^2}{8\pi} \int d^4x \mathbf{E}^a \cdot \mathbf{B}^a \quad N_{CS}(t) = -2N_f \int d^3x j_5^0$$

- Determine global to axial charge production via the anomaly by tracking change in Chern-Simons number

$$\Delta N_5 = \int d^4x j_5^0 = -2N_f \Delta N_{CS}$$

# Sphalerons in equilibrium

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- Sphaleron transitions are given by modes on the order of the magnetic screening length  $\Lambda_s \sim \frac{1}{g^2 T}$   
(Arnold and McLerran, PRD37, 1020 (1988))
- In equilibrium, sphaleron transitions dominate long time Chern-Simons number diffusion- probabilistic

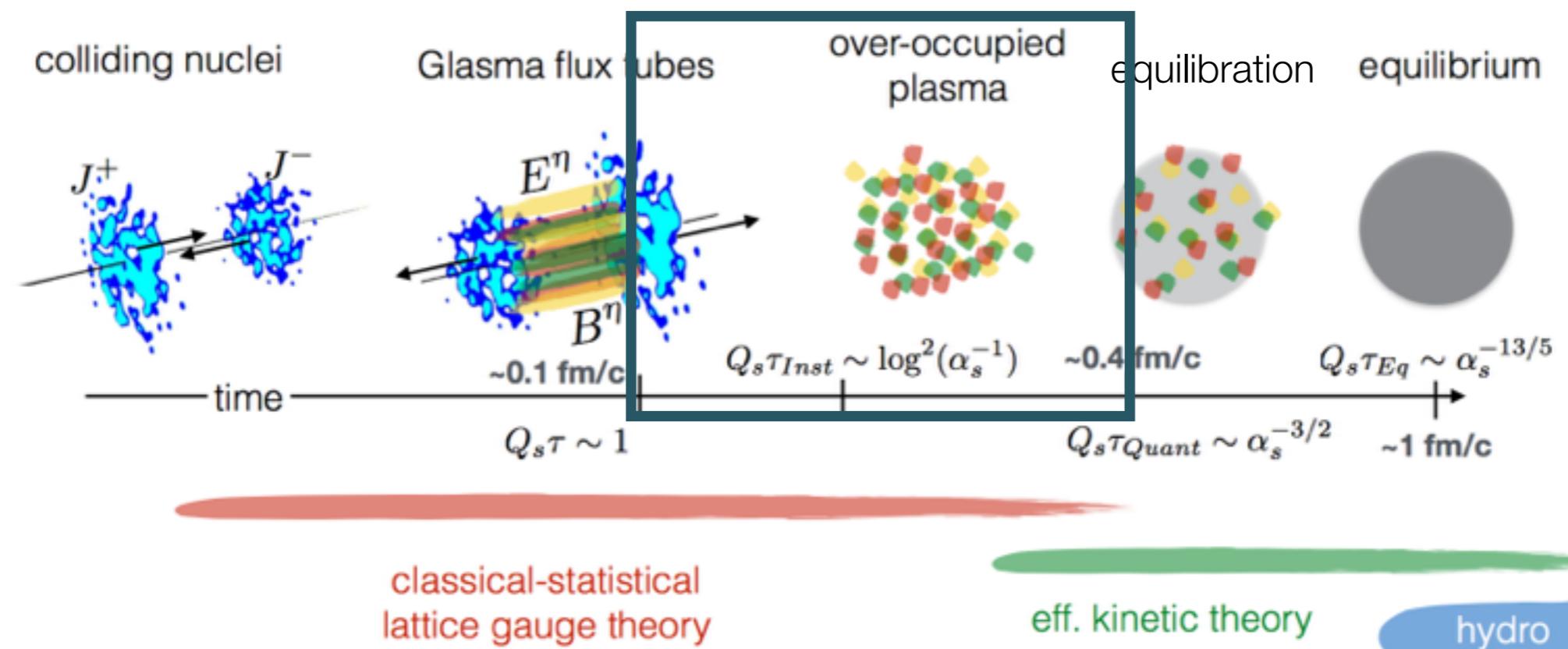
- Phase space:  $(\alpha_S T)^3$       Typical freq:  $(\alpha_S^2 T)$   
(Arnold, Son, and Yaffe, PRD55, 6264 (1997);  
Bodeker PLB426 (1998) 351-360)
- Integer random walk with known diffusion constant

$$\Gamma_{sph}^{eq} = \lim_{\delta t \rightarrow \infty} \frac{\langle (N_{CS}(t + \delta t) - N_{CS}(t))^2 \rangle}{V \delta t} = \kappa \alpha_S^5 T^4$$

- Most recent calculations  $\kappa \sim 25$   
(Moore and M. Tassler, JHEP 02, 105 (2011))

# The glasma

- Magnetic field is understood to be strongest at  $\leq 1\text{fm}/c$
- This is when system is far from equilibrium, gluon dominated
  - Described by the ‘glasma’  
(McLerran and Lappi, NPA772 (2006), Krasnitz and Venugopalan, NPB557 (1999) 237)
  - Necessary to study axial charge generation at this time



(Adapted from S.Schichtling Bielefeld 2016)

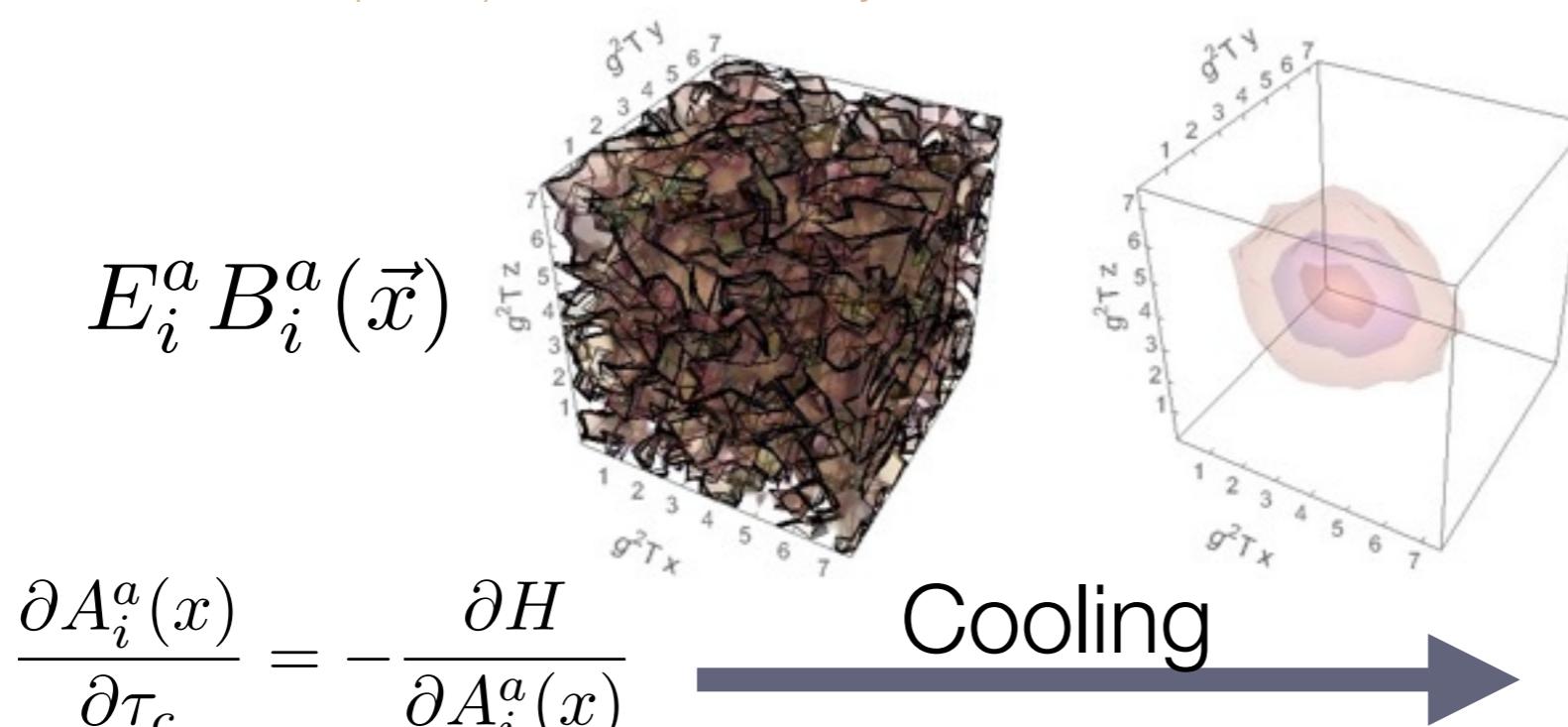
# Early times after HIC

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- Initially gluon occupation much larger than quark
  - Study classical Yang-Mills  
(McLerran, Venugopalan PRD49 2233 (1994))
- Non-perturbative large phase-space density of gluon
  - $f(p \sim Q_s) \sim \frac{1}{\alpha_S}$
  - Amenable to classical-statistical lattice description  
(Berges, et al PLB 681 2009; Berges et al PRD89, 074011 (2014))
- 2D boost-invariant case previously studied  
(Kharzeev, Krasnitz, Venugopalan PLB545, 298 (2002))
- Neglect longitudinal expansion and consider SU(2) for simplicity

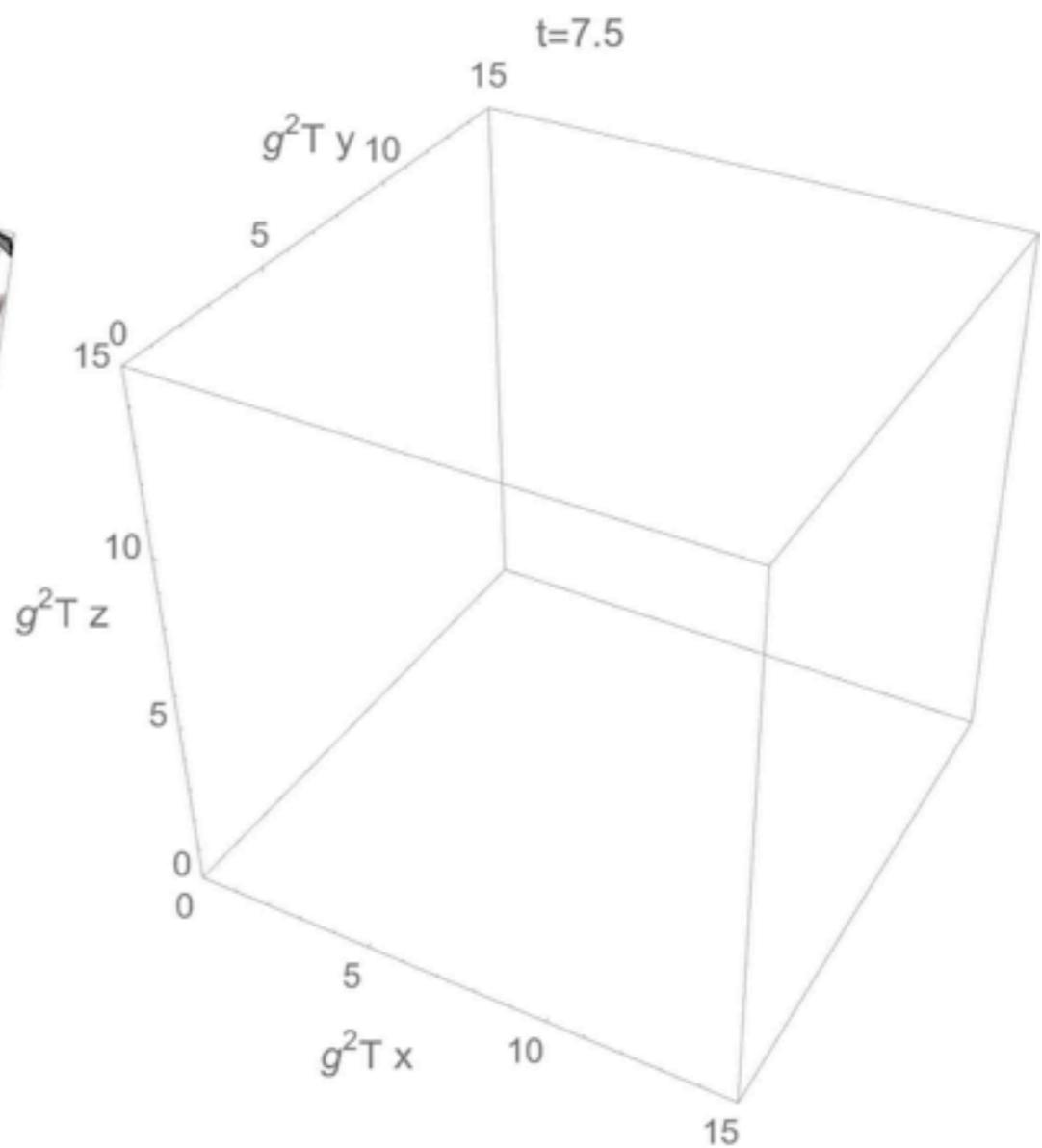
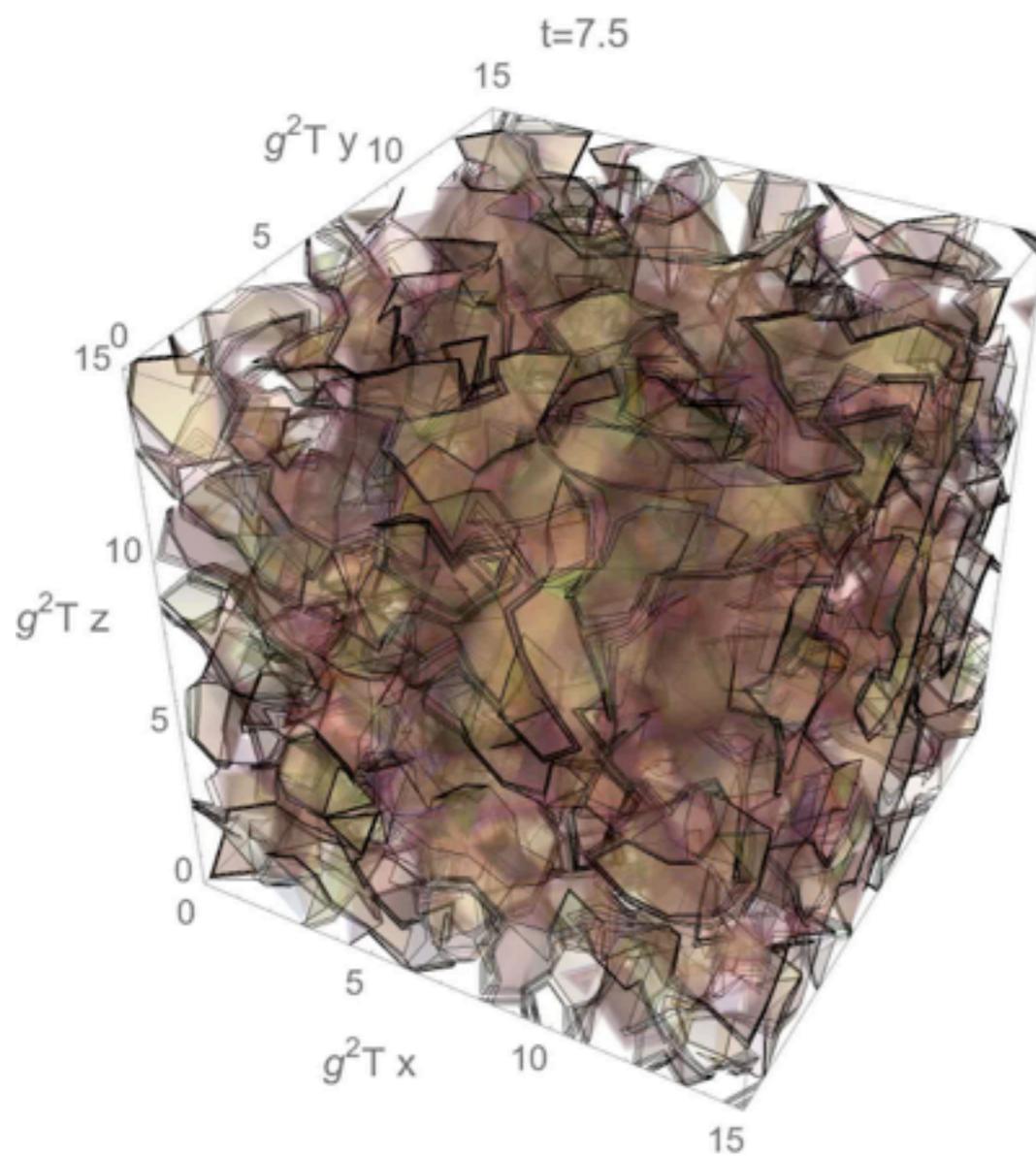
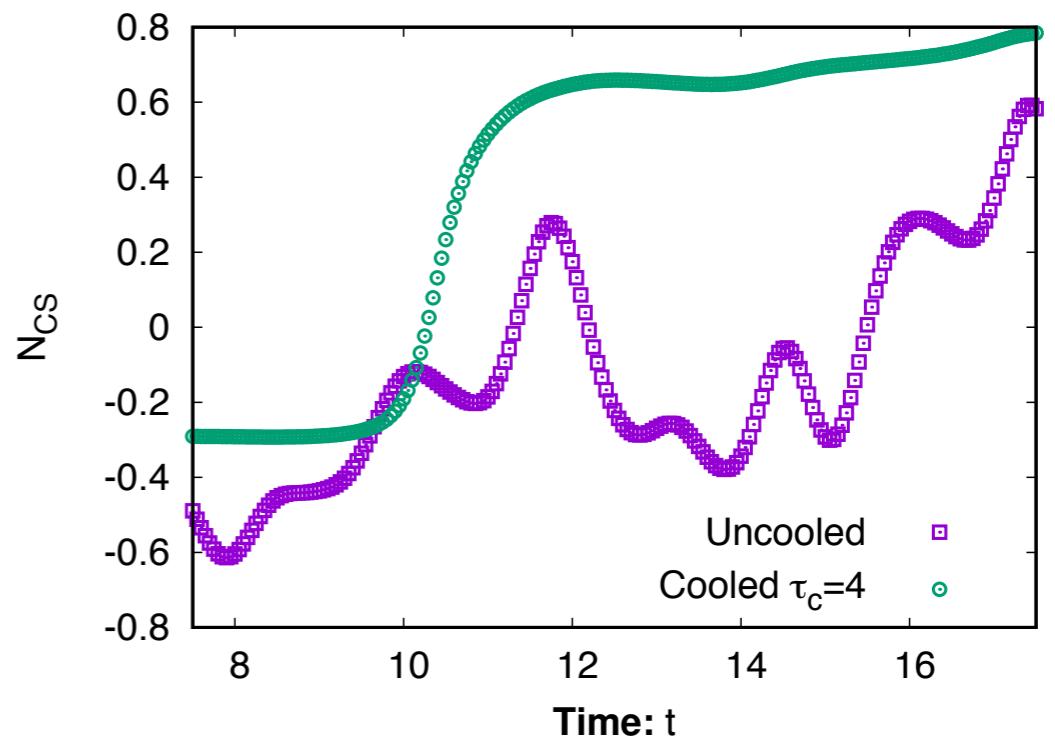
# Topology on the lattice

- Solve classical Yang-Mills in real time 3D spatial lattice
- Calculate  $\frac{dN_{CS}}{dt} = \frac{1}{8\pi^2} \int d^3x E_i^a B_i^a$
- Lattice  $E_i^a B_i^a$  not a total derivative, plagued by UV noise
- Need gradient flow (cooling) to see topological contribution  
(Moore NPB480 (1996) 657-688; Ambjorn and Krasnitz, NPB506, 387 (1997))



(MM, Schlichting,  
Venugopalan PRD 074036)

# Sphaleron in action



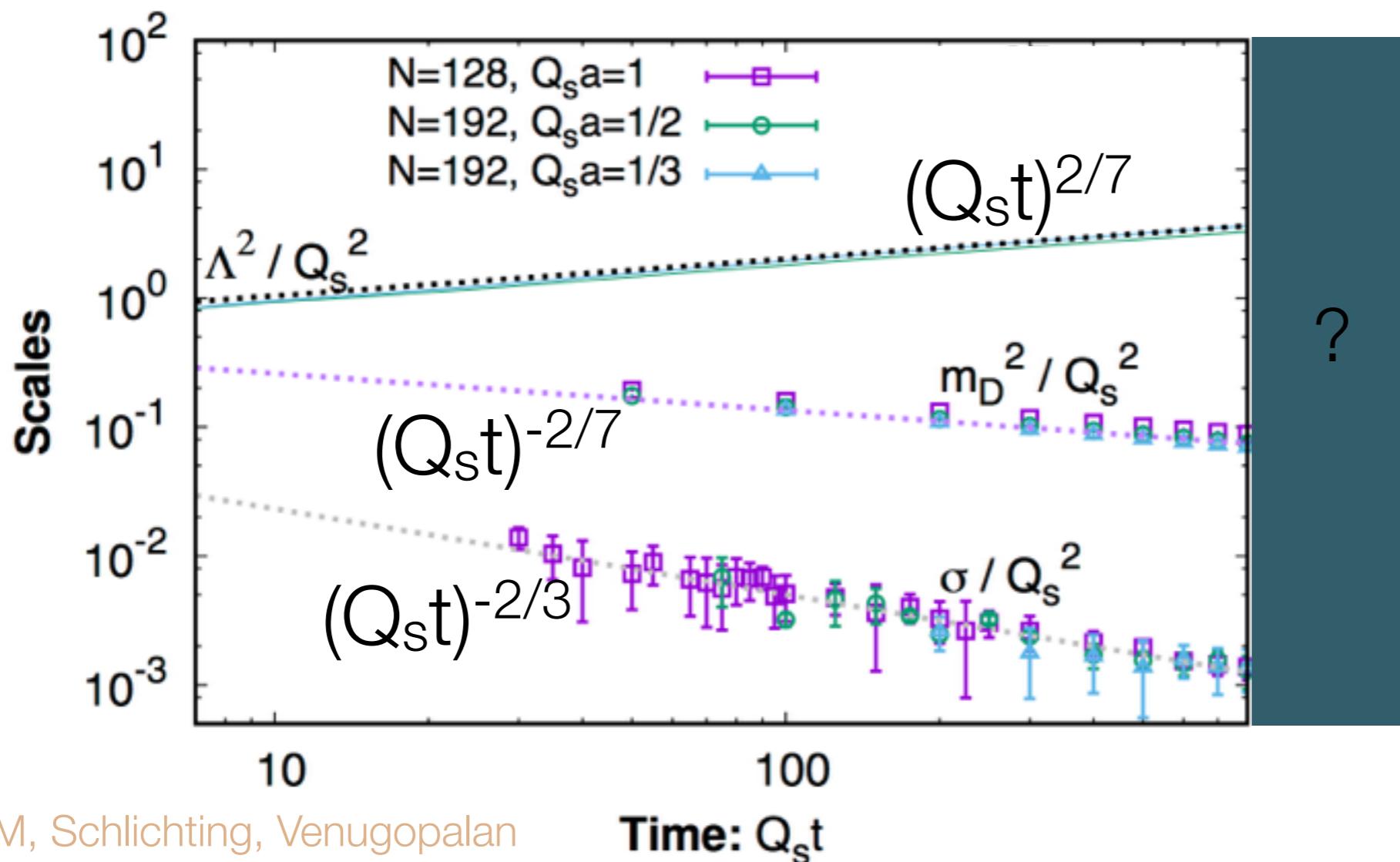
# Dynamical scales at a glance

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- Hard scale  $Q_s \longrightarrow \Lambda^{eq} \sim T$ 
  - Average typical momentum of hard modes  $\Lambda(t)^2 = \frac{2}{3} \frac{\int d^3p p^3 f(p)}{\int d^3p p f(p)}$   
(Kurkela and Moore, JHEP 1112, 044 (2011); Schlichting PRD86, 065008 (2012))
- Electric screening scale  $Q_s \longrightarrow m_D^{eq} \sim gT$ 
  - Appeal to perturbative formula  $m_D^2 = 4g^2 N_c \int \frac{d^3p}{(2\pi)^3} \frac{f(p)}{p}$   
(Kurkela and Moore, PRD86, 056008 (2012))
- Magnetic screening scale  $Q_s \longrightarrow \Lambda_s^{eq} \sim \sqrt{\sigma} \sim g^2 T$ 
  - Measure spatial string tension  
(Dumitru, Lappi, and Nara, PLB734, 7 (2014))
  - Large area fall limit of spatial Wilson loop  
 $\langle W(A, t) \rangle \approx \exp(-\sigma A)$

# Dynamical separation of scales

- Kinetic theory predictions exist (Blaizot et al, NPA873, 68 (2012))
  - 2/3 in agreement



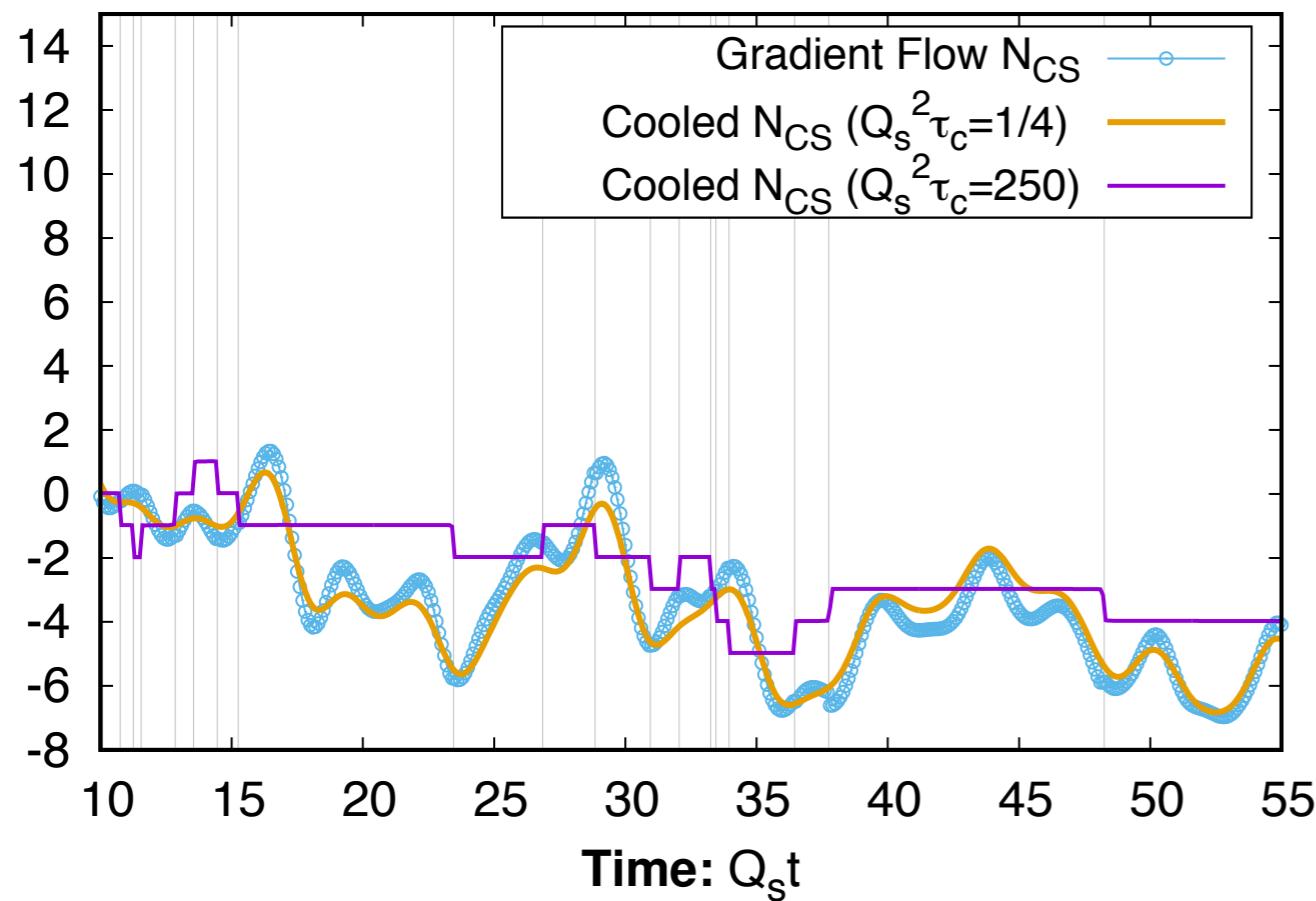
At equilibrium:  
Hard Scale  
 $\Lambda^{eq} \sim T$

Electric Screening  
 $m_D^{eq} \sim gT$

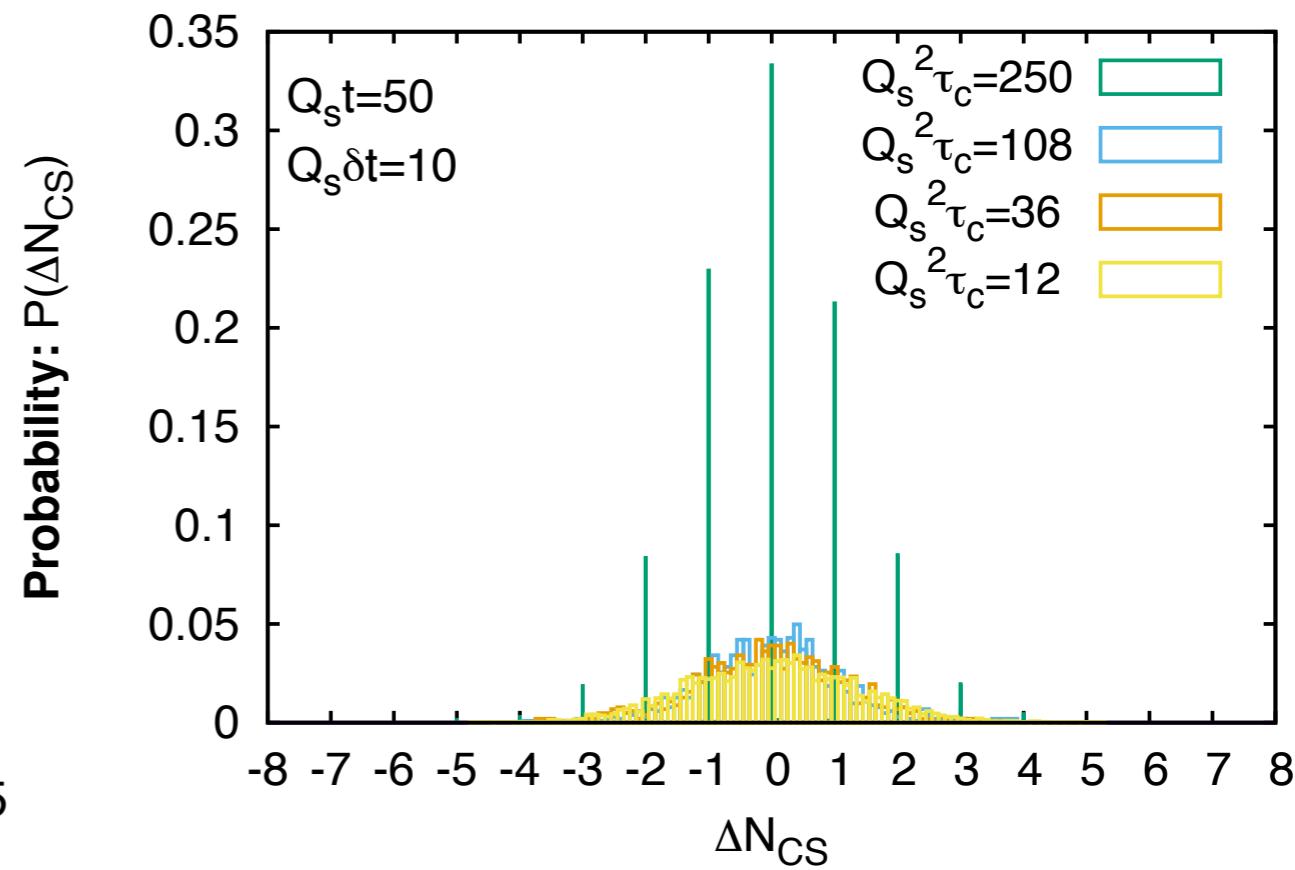
Magnetic Screening  
 $\Lambda_s^{eq} \sim \sqrt{\sigma} \sim g^2 T$

# Sphalerons in the glasma

- Detect integer changes in Chern-Simons number for single configuration
- Histogram Chern-Simons diffusion of many configurations

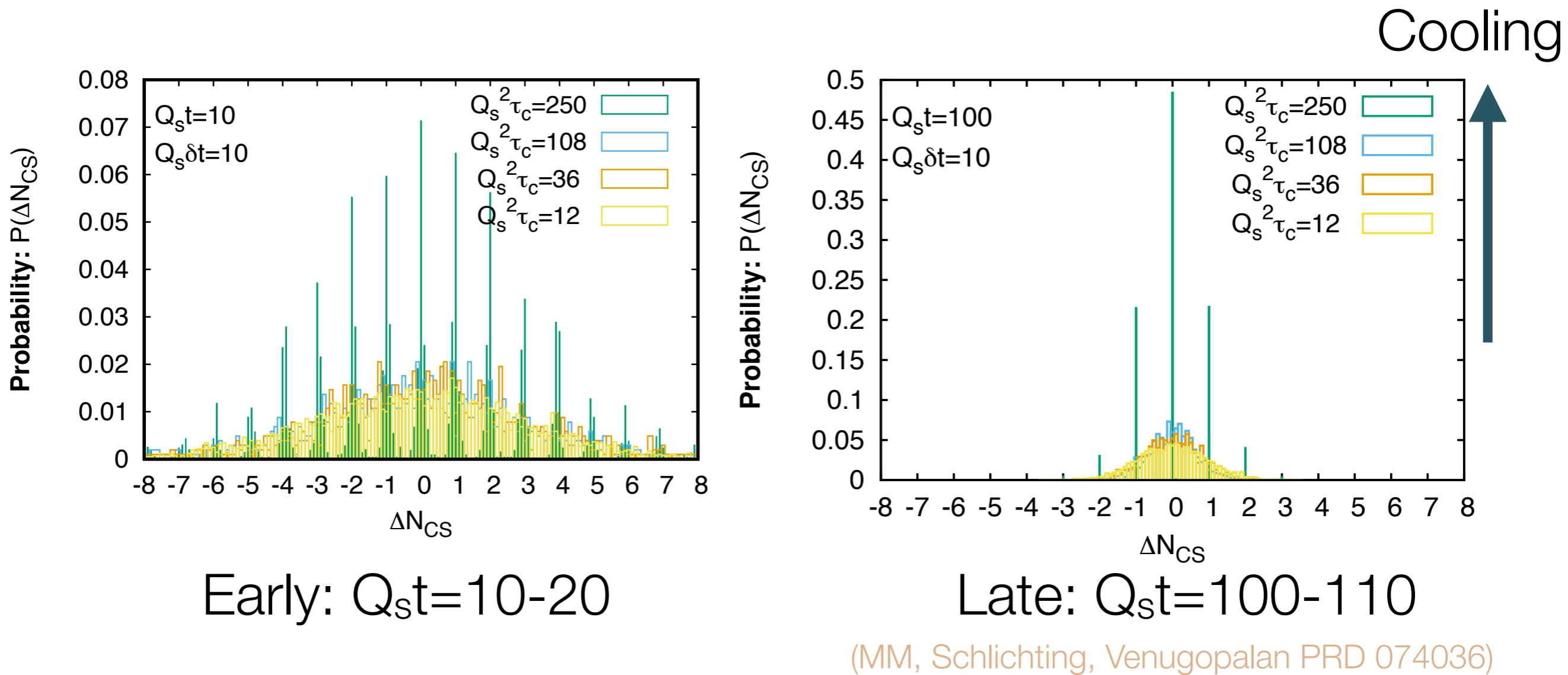


(MM, Schlichting, Venugopalan PRD 074036)



Can isolate transitions from background

# Sphalerons in the glasma



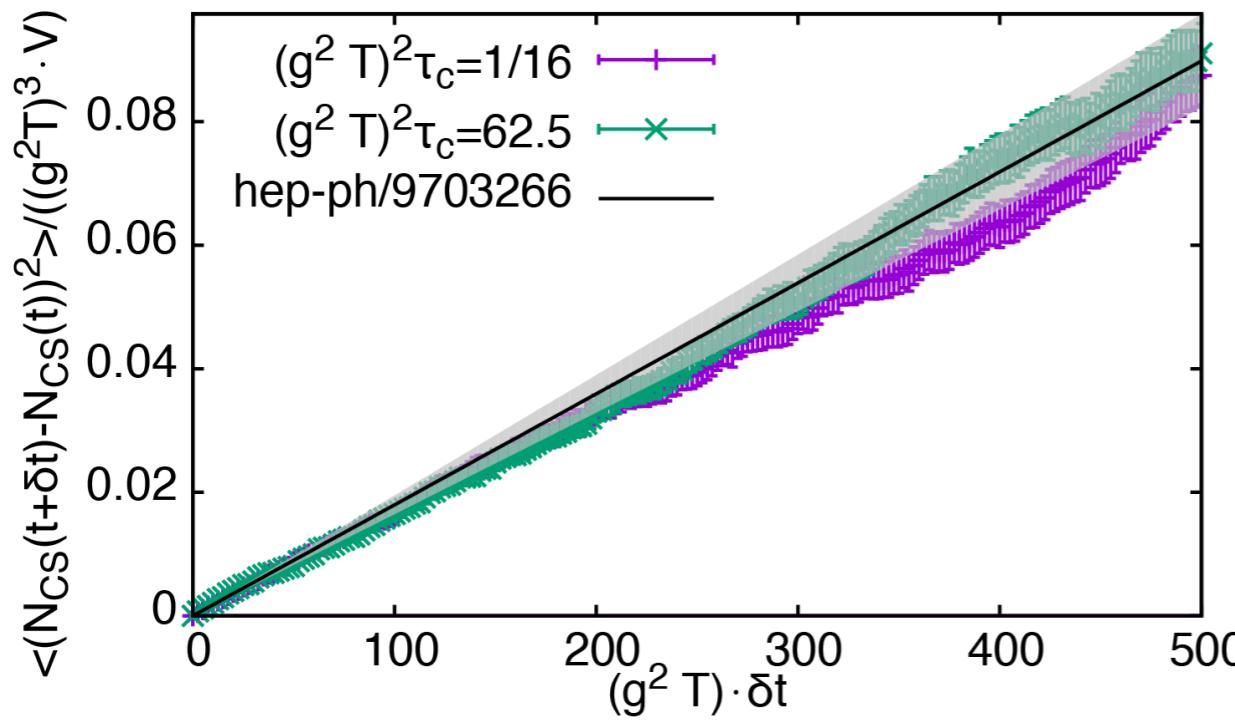
(MM, Schlichting, Venugopalan PRD 074036)

Most transitions happen at early times

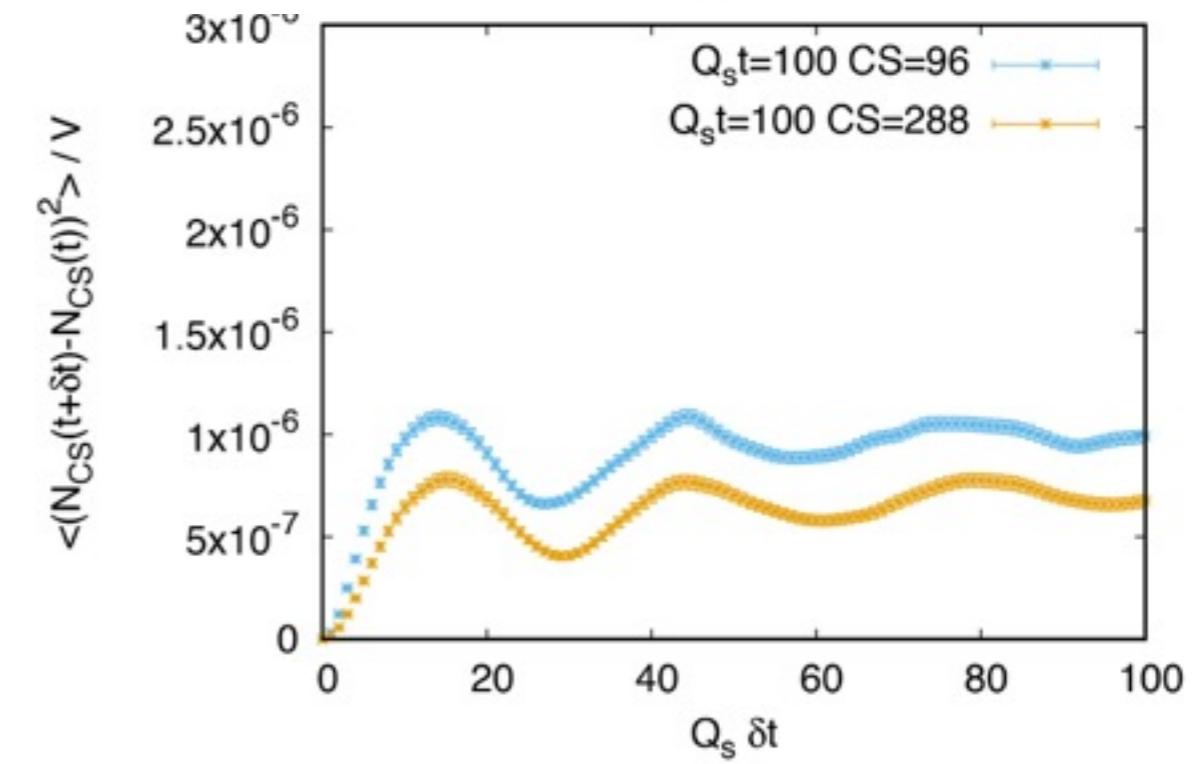
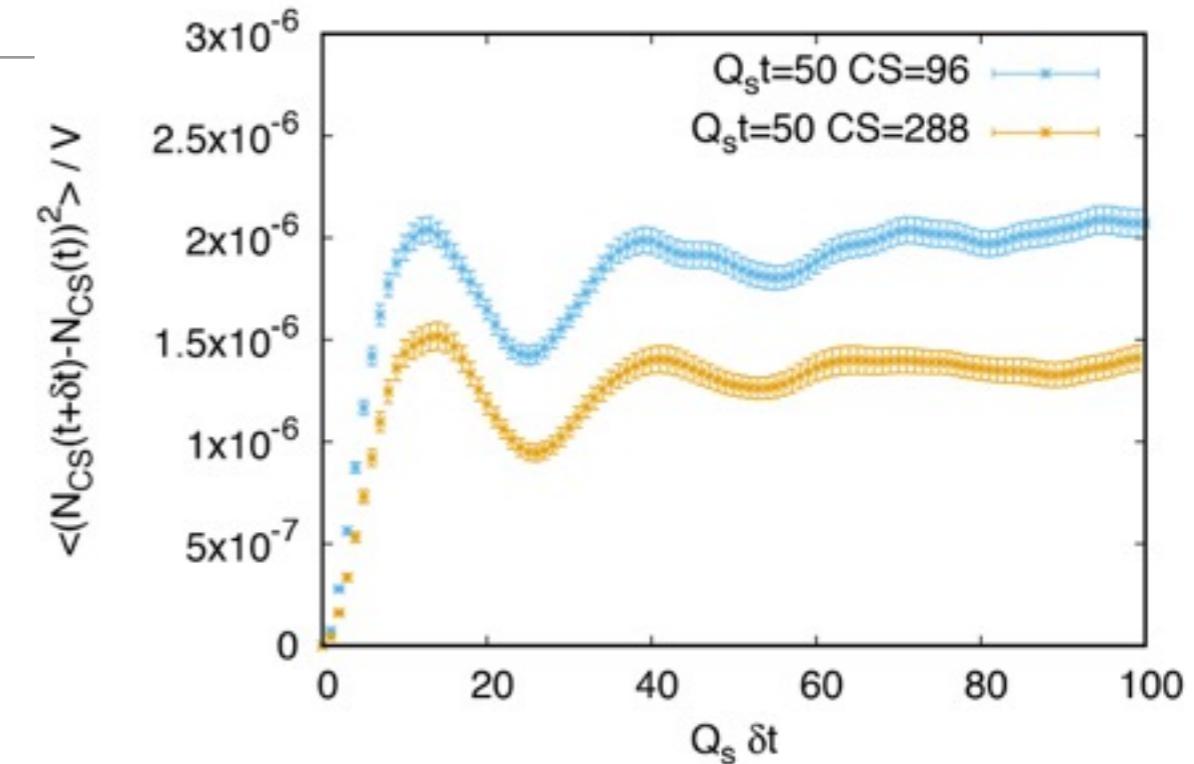
# Chern-Simons number correlations

(MM, Schlichting,  
Venugopalan PRD 074036)

- In equilibrium: random walk
- Out of equilibrium: Non-Markovian
  - Simple probabilistic picture breaks down



Thermal Equilibrium

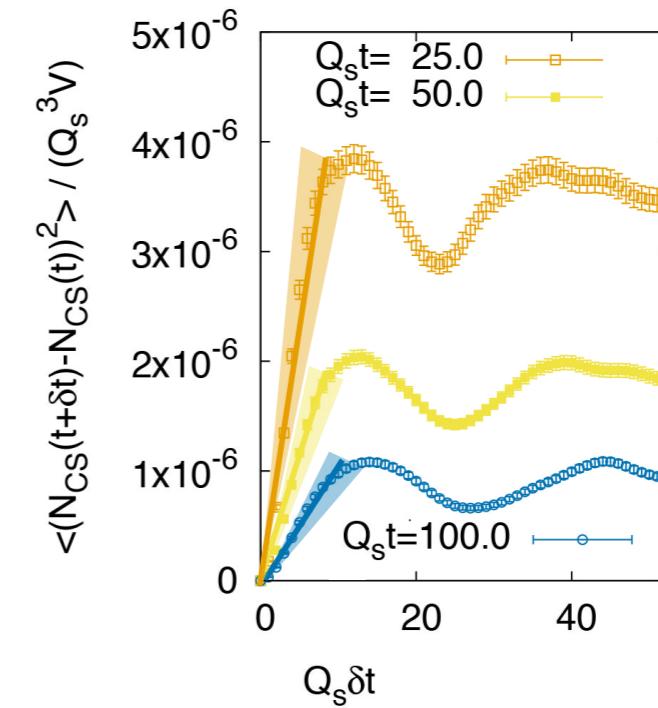
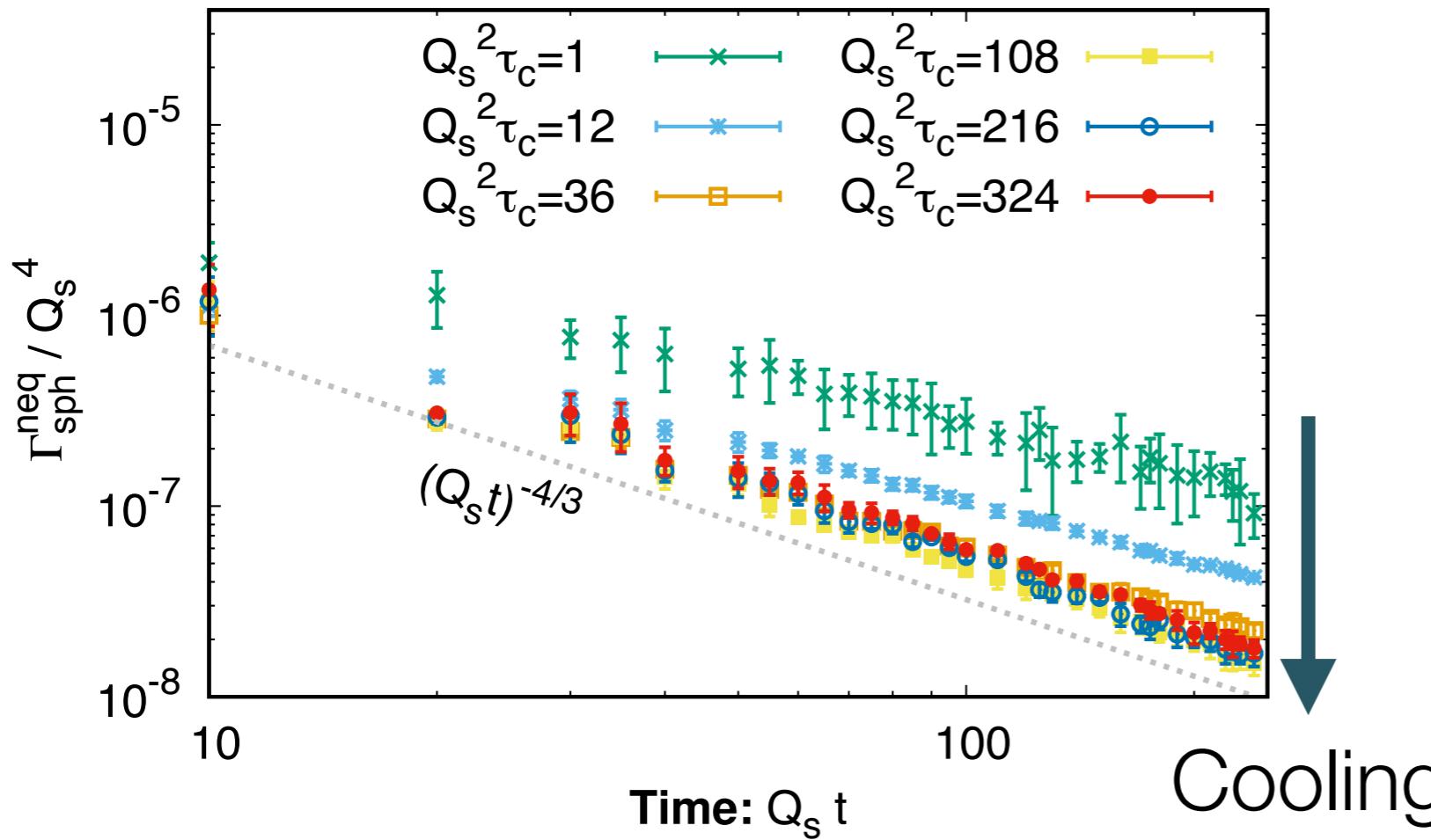


Glasma

# Non-equilibrium sphaleron transition rate

- Define non-equilibrium transition rate as initial rise in autocorrelation function

$$\Gamma_{sph}^{neq}(t) = \left\langle \frac{(N_{CS}(t + \delta t) - N_{CS}(t))^2}{V \delta t} \right\rangle_{Q_s \delta t < 10}$$



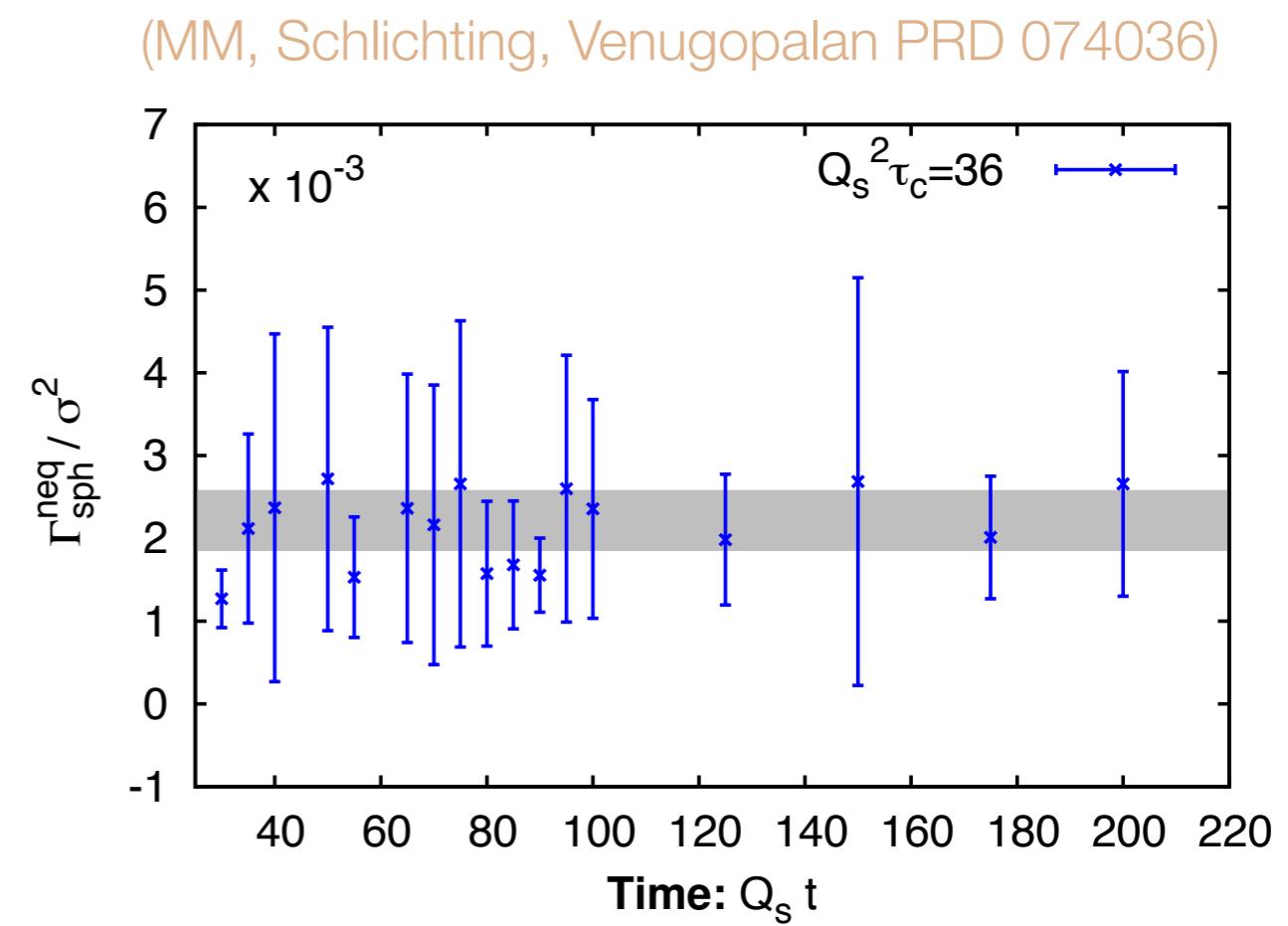
- Strongly time dependent
- Contributions from field strength fluctuations

# Non-equilibrium sphaleron transition rate

- In equilibrium, sphaleron rate controlled by magnetic modes
- From dynamical separation of scales, non-equilibrium sphaleron rate controlled by modes of order magnetic screening

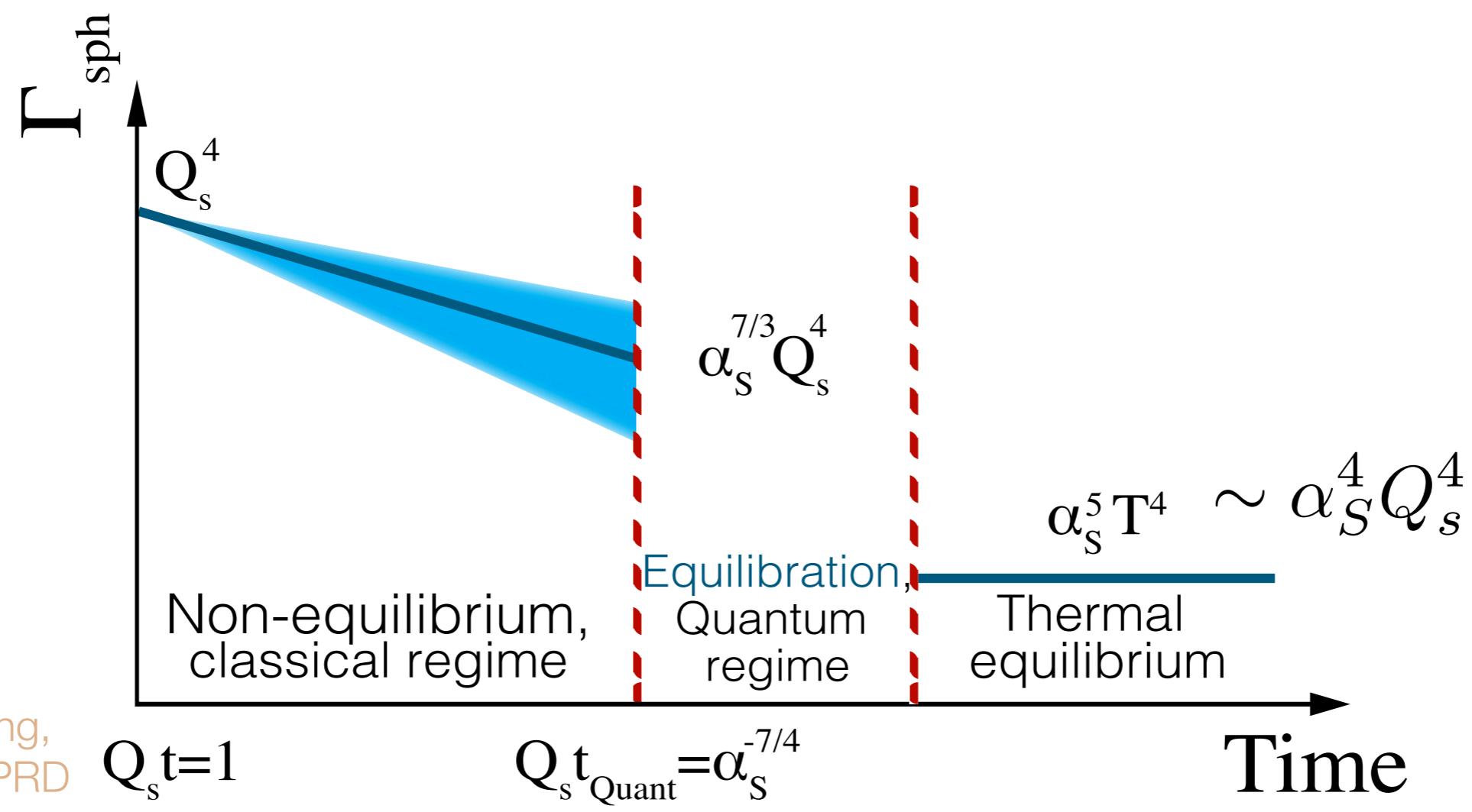
$$\Gamma_{sph}^{neq}(t) = 2 \times 10^{-2} \sigma^2(t)$$

$$\sigma^2(t) \approx Q_s^2 (Q_s t)^{-2/3}$$



# Sphalerons in non-Abelian plasmas

- Glasma (early times):  $\Gamma_{sph}^{neq} \approx Q_s^4 \rightarrow \Gamma_{sph}^{neq}(t) \approx Q_s^4(Q_s t)^{-4/3}$
- Equilibrium (late times):  $\Gamma_{sph}^{eq} \approx \alpha_S^5 T^4$



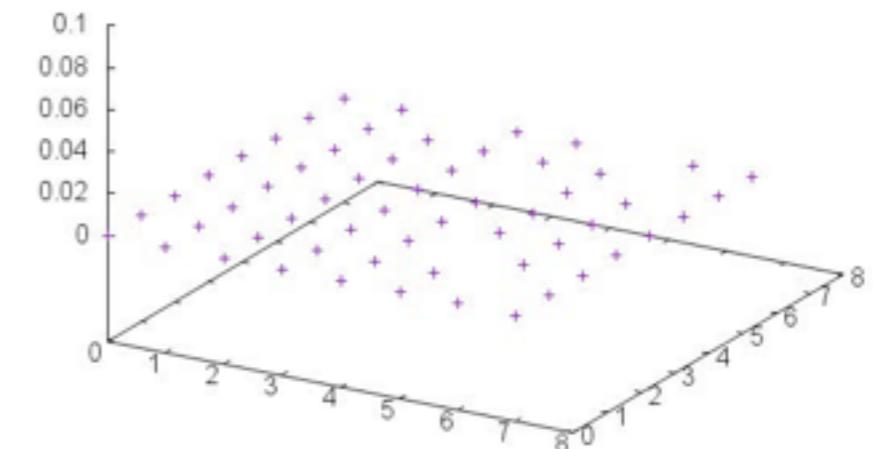
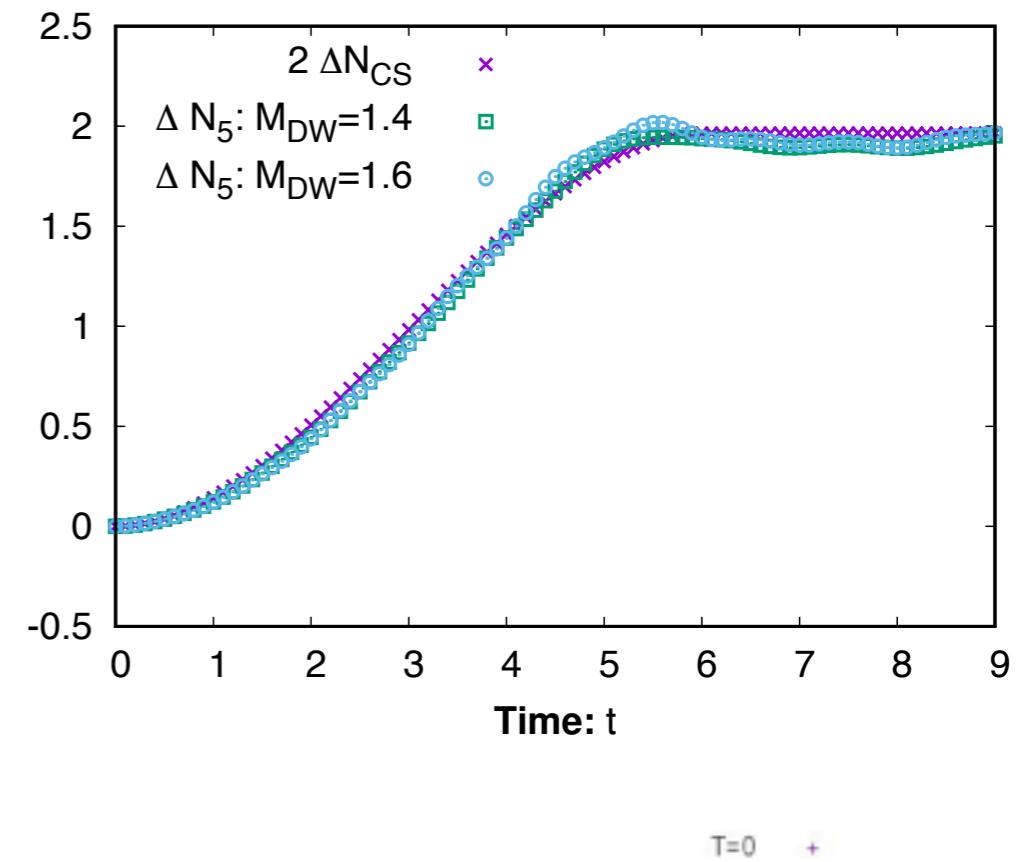
# Implications for the CME

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- Generalizing for longitudinal expansion in the glasma
  - Initial times appear to be dominant in all respects
    - Greatest magnetic field
    - Dominant amount of axial charge (should be) generated
      - Glasma:  $\Gamma_{sph}^{neq} \sim Q_s^4$     Equilibrium:  $\Gamma_{sph}^{eq} \sim \alpha_S^5 T^4 \sim \alpha_S^4 Q_s^4$
  - However, need longitudinal expansion to address definitively and be more quantitative
    - Need to understand topological transitions on expanding lattice
    - Need to understand magnetic screening in anisotropic plasmas

# Future outlook

- Add lattice fermions, study real-time generation of axial charge and formation of CME/CMW with same non-equilibrium gauge configurations
  - Overlap fermions with background magnetic field
    - Work started with S. Schlichting, S. Sharma
- Sphaleron inspired initial condition for anomalous hydrodynamics
  - Discussions begun with B. Schenke, D. Kharzeev, Y. Hirono,



Thanks!