

# Mixed Dark Matter in a Left-Right Extension of SM

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based on arXiv:1604.06100v1  
with A. Berlin, P. Fox and D. Hooper.



# Outline

- Introduction
- Left-Right extension of SM
  - \* Higgs Sector
  - \* Gauge Boson Sector
  - \* Fermion Sector
- Dark Matter Phenomenology
  - \* Singlet Triplet
  - \* Singlet Bidoublet
  - \* Triplet Bidoublet
- Conclusions

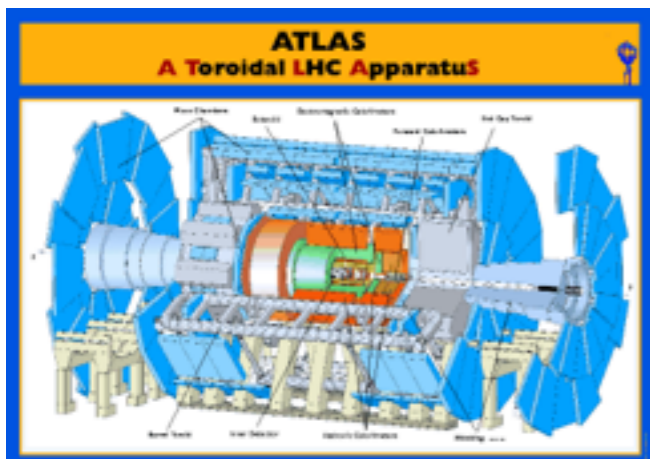
# Motivation

MAIN GOAL : Discuss new DM scenario in Left-right model.

Motivation 1: 8 TeV ATLAS & CMS diboson excess.



Motivation 2: Interesting Dark Matter phenomenology with detection prospects at future experiments.



Recap

# Diboson excesses at LHC Run 1

1.8 - 2.0 TeV  
Range

ATLAS

$X \rightarrow V(J)V(J)$  Local:  $3.4\sigma$  (WZ),  $2.6\sigma$  (WW),  $2.9\sigma$  (ZZ)

Global:  $2.5\sigma$  (WZ)

arXiv:1506.00962

$X \rightarrow jj$   $1\sigma$  (jj)

arXiv:1407.1376

CMS

$X \rightarrow V(J)V(J)$   $1.4\sigma$

arXiv:1405.1994

$X \rightarrow jj$   $2\sigma$

$W' \rightarrow N_R e$   $2.2\sigma$

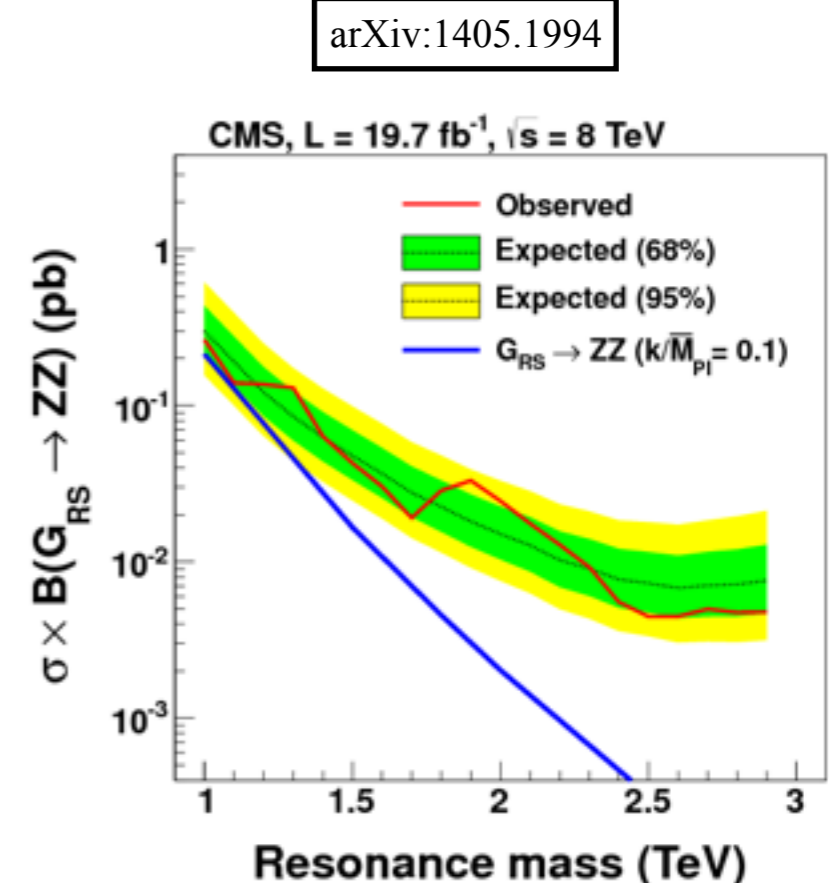
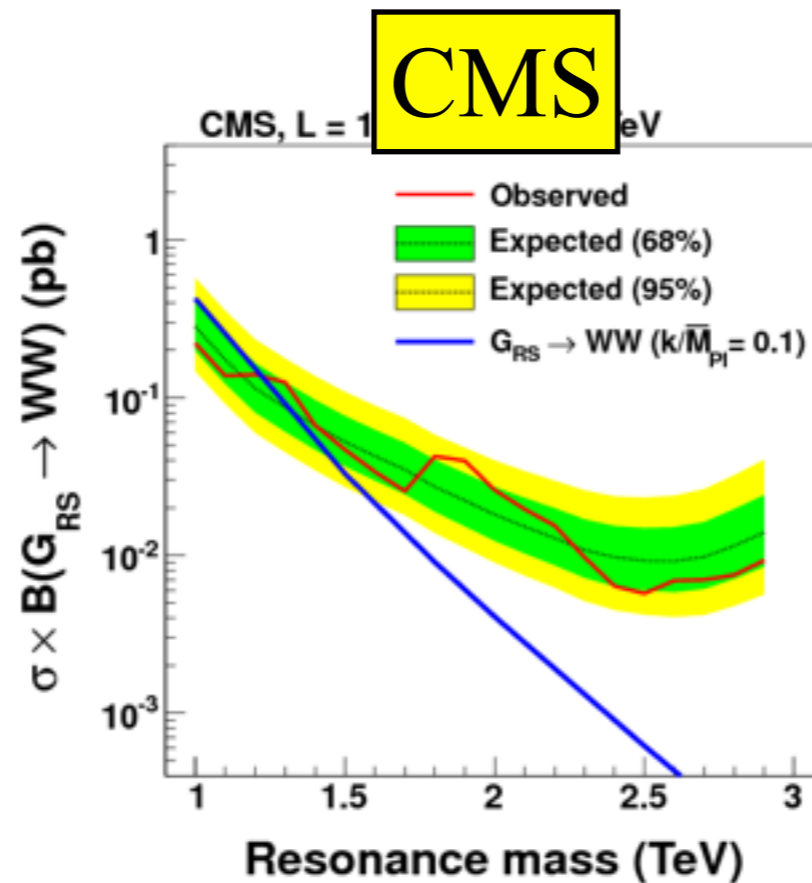
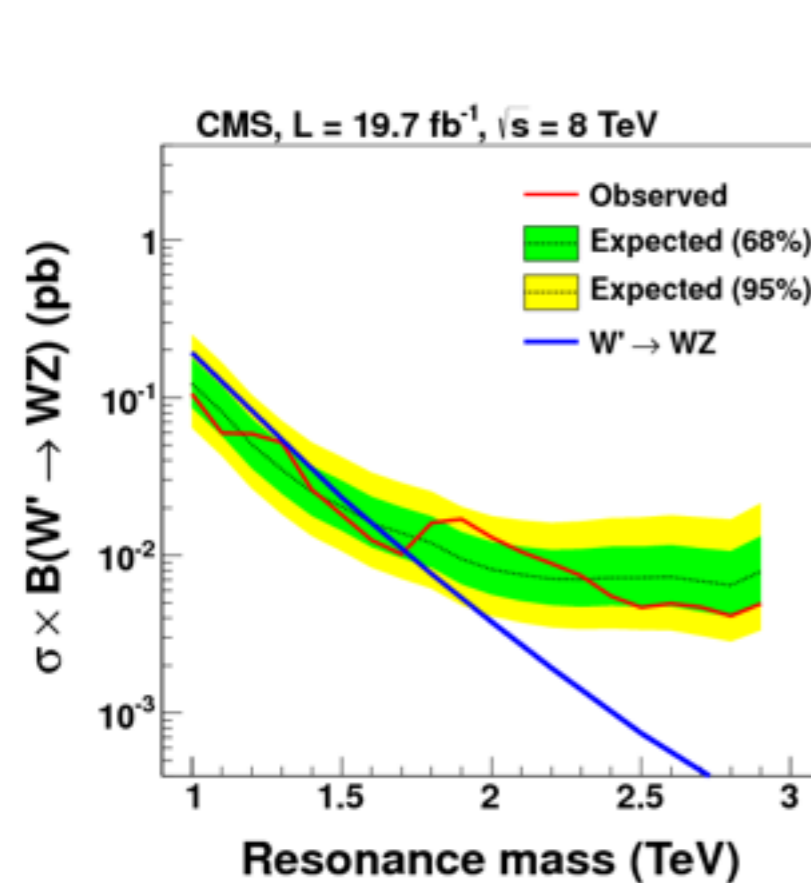
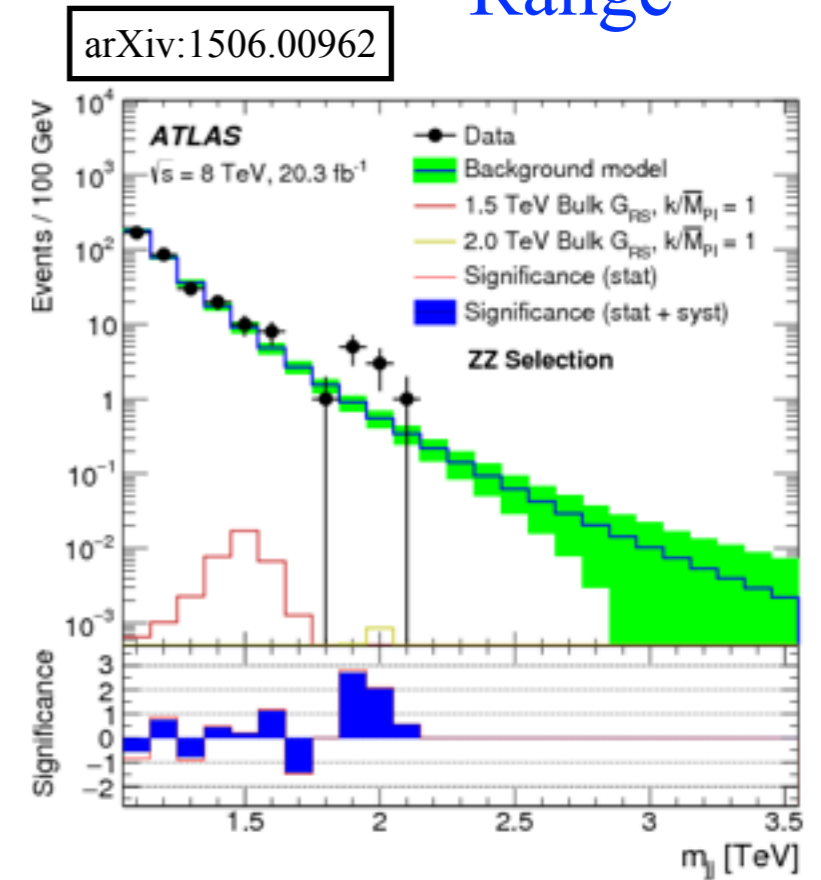
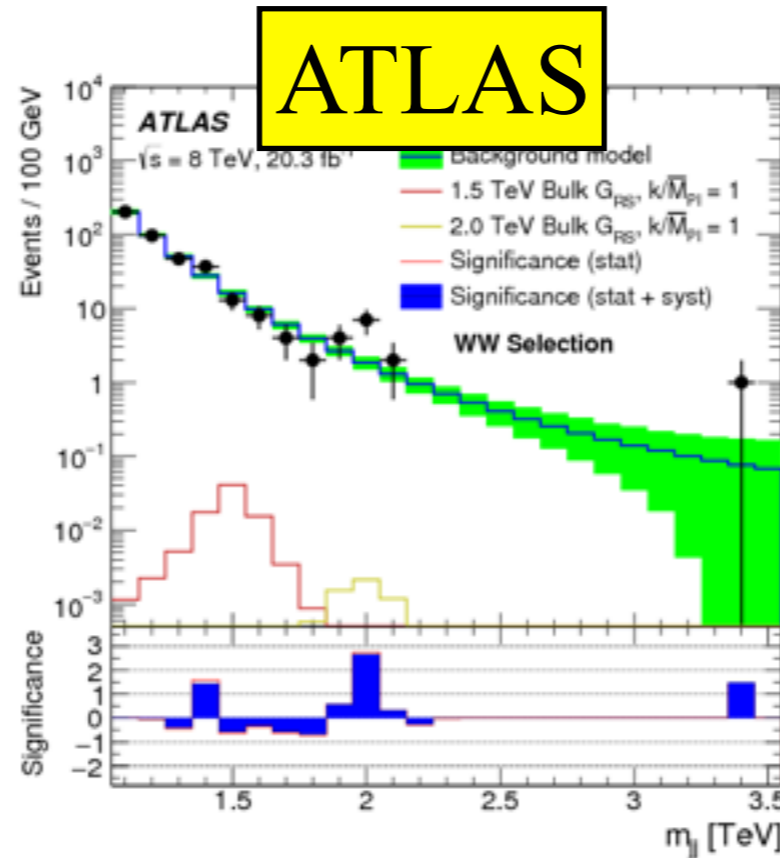
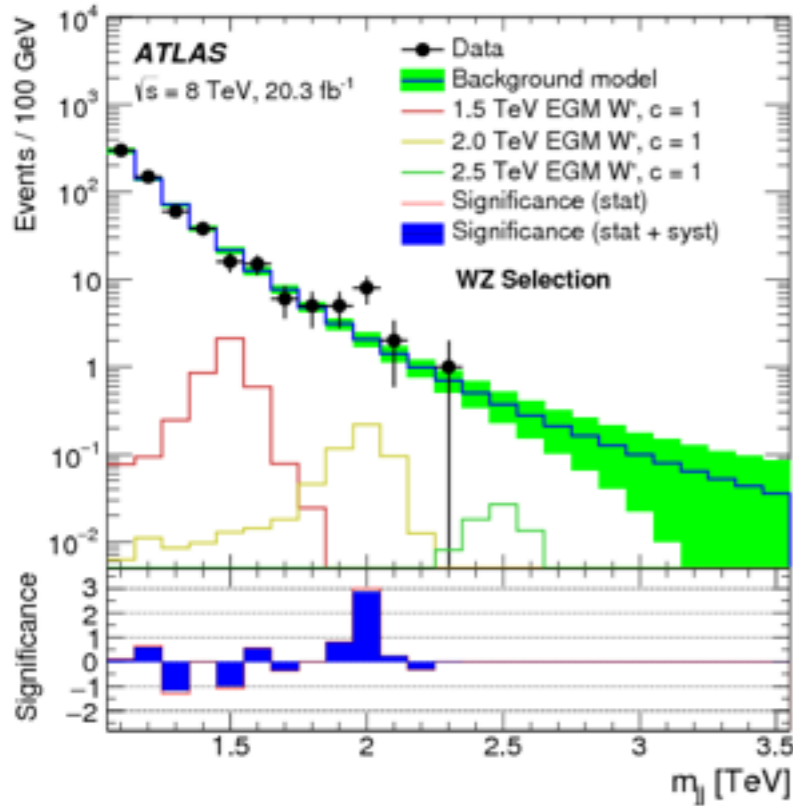
arXiv:1407.3683

$W' \rightarrow Wh$   $2.8\sigma$

PAS-EXO-14-010

# Diboson excesses at LHC Run 1

1.8 - 2.0 TeV  
Range





Has there been a 2 TeV resonance sighting?



13 TeV not clear enough to confirm or rule out

I want to believe in 14 TeV

All 5/6 excesses consistent with massive  $W'$  boson connected to BSM theories

## Left-Right Extension of SM

Introduced in 1974-75 to primarily Understand P-violation in SM.

At Lower Energies L-R models described by gauge group:

$$SU(3)_c \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$



# Higgs Sector

Extended Higgs sector:

$$\Delta_L \in (3, 1, 2) \quad \Delta_R \in (1, 3, 2) \quad \phi \in (2, 2, 0)$$

Spontaneous Symmetry Breaking at TeV scale:

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$



$$\langle \Delta_R \rangle = v_R$$

$$SU(2)_L \otimes U(1)_Y$$

$$SU(2)_L \otimes U(1)_Y$$



$$\langle \phi \rangle \sim v$$

$$U(1)_{EM}$$

Right Higgs Triplet is parametrized:

$$\Delta_R = \begin{pmatrix} \frac{g_R}{\sqrt{2}g_L} \frac{M_W}{M_{W'}} c_{2\beta} H^+ & \Delta^{++} \\ v_R + \frac{1}{\sqrt{2}} \Delta^0 & -\frac{g_R}{\sqrt{2}g_L} \frac{M_W}{M_{W'}} c_{2\beta} H^+ \end{pmatrix}$$

Higgs Bidoublet is parametrized

$$\phi = \begin{pmatrix} c_\beta v + \frac{1}{\sqrt{2}} (c_\beta h + s_\beta H + i s_\beta A) & c_\beta H^+ \\ s_\beta H^- & s_\beta v + \frac{1}{\sqrt{2}} (s_\beta h - c_\beta H + i c_\beta A) \end{pmatrix}$$

# Gauge Boson Sector

After SSB Gauge Bosons get masses.

Compatible with 2 TeV excess

$$M_W = \frac{g_L v}{\sqrt{2}}$$

$$M_{W'} = g_R v_R$$

$$M_A^2 = 0$$

$$M_Z^2 = (g_L^2 + g_Y^2) \frac{v^2}{2}$$

$$M_{Z'}^2 = 2(g_R^2 + g_{B-L}^2) v_R^2$$

# Fermion Sector

Parity is restored in LR model in Fermion sector:

$$Q_L = \begin{pmatrix} u \\ d \end{pmatrix}_L$$

$$Q_R = \begin{pmatrix} u \\ d \end{pmatrix}_R$$

$$Q = T_{3L} + T_{3R} + \frac{Y_{B-L}}{2}$$

$$L_l = \begin{pmatrix} \nu \\ l \end{pmatrix}_L$$

$$L_R = \begin{pmatrix} N \\ l \end{pmatrix}_R$$

# Dark Matter Phenomenology

$SU(2)_R$ multiplets	Mixing?
Singlet-Doublet	No ✘
Singlet-Triplet	Yes ✔
Singlet-Bidoublet	Yes ✔
Doublet-Triplet	No ✘
Doublet-Bidoublet	No ✘
Triplet-Bidoublet	Yes ✔

mixing through renormalizable Yukawa Interactions with Triplet & Bidoublet Higgs

Similar studies in non LR models

C. Cheung & D. Sanford. (arXiv:1311.5896)

T. Tait & Z. Yu. (arXiv:1601.01354)



# Singlet-Triplet

Fields	Charges	Spin
S	(1,1,0)	1/2
T <sub>1</sub>	(1,3,2)	1/2
T <sub>2</sub>	(1,3,-2)	1/2

Fermion Triplets parametrized as:

$$T_1 = \begin{pmatrix} t_1^+ / \sqrt{2} & t_1^{++} \\ t_1^0 & -t_1^+ / \sqrt{2} \\ t_1^- & -t_1^- / \sqrt{2} \end{pmatrix} \quad T_2 = \begin{pmatrix} t_2^+ / \sqrt{2} & t_2^{++} \\ t_2^0 & -t_2^+ / \sqrt{2} \\ t_2^- & -t_2^- / \sqrt{2} \end{pmatrix}$$

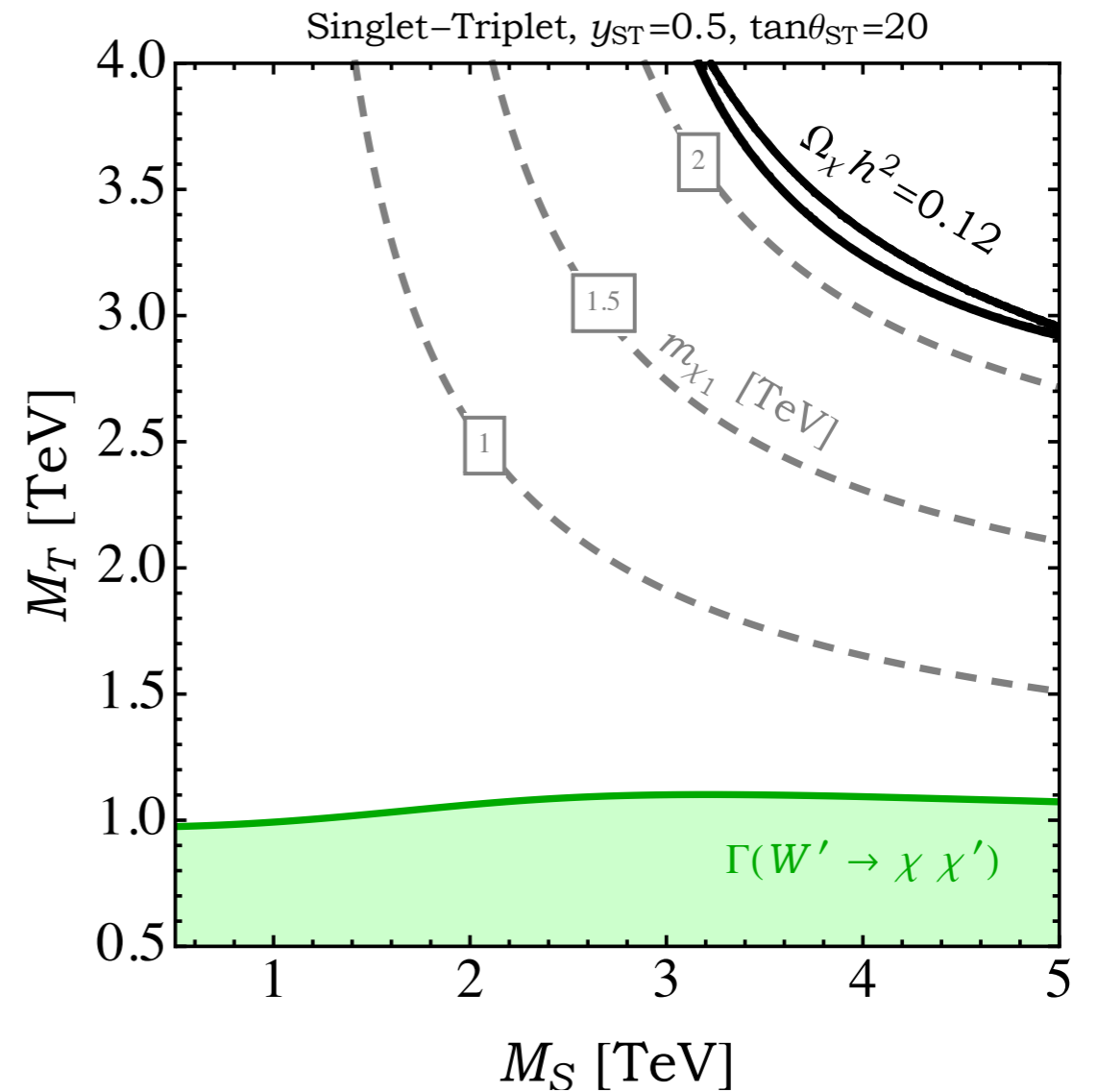
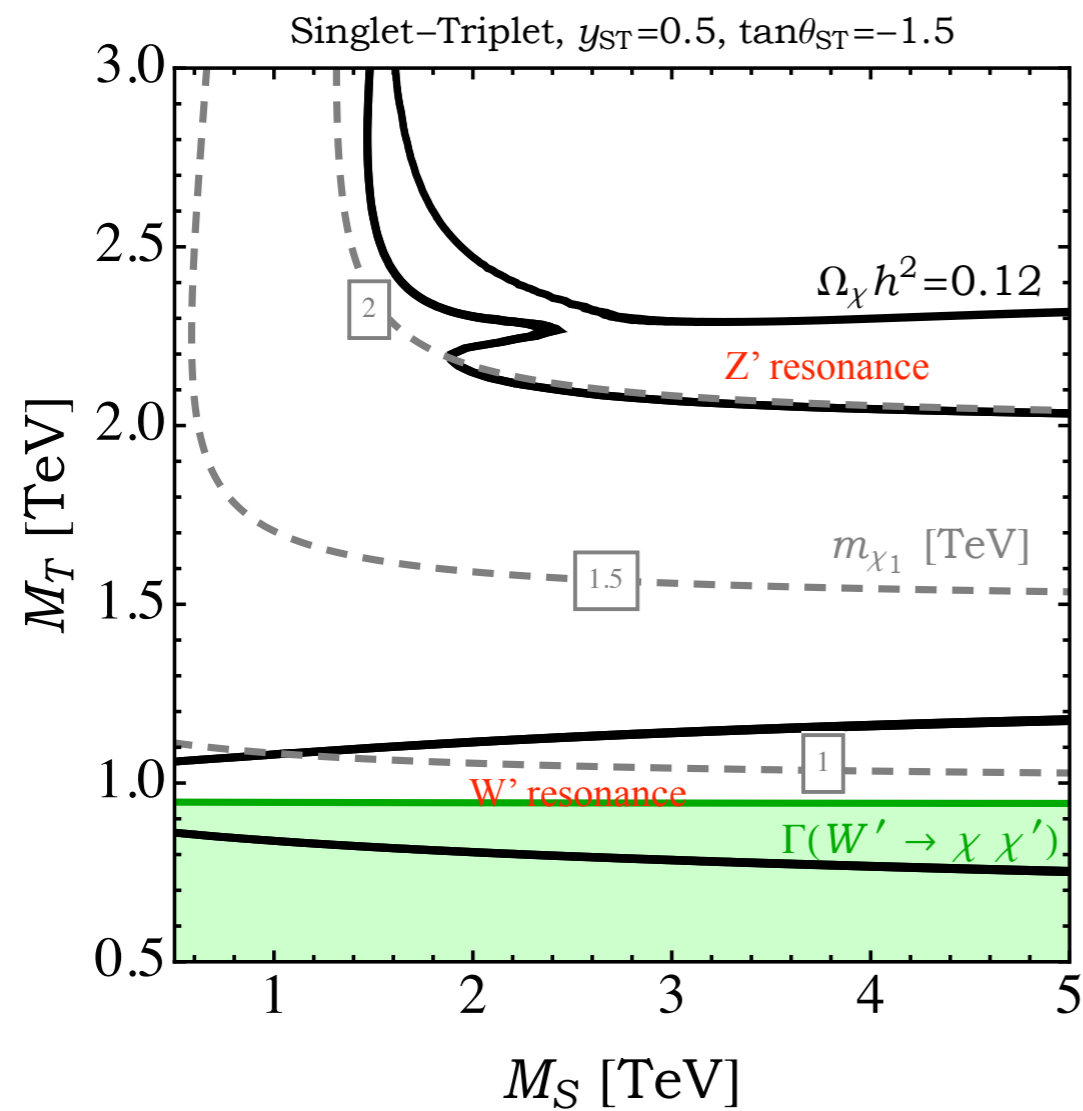
Singlet-Triplet Lagrangian:

$$\mathcal{L} \supset S^\dagger i \bar{\sigma}^\mu \partial_\mu S + \text{tr}(T_1^\dagger i \bar{\sigma}^\mu D_\mu T_1) + \text{tr}(T_2^\dagger i \bar{\sigma}^\mu D_\mu T_2) - \left[ \frac{1}{2} M_S S^2 + M_T \text{tr}(T_1 T_2) \right. \\ \left. + \lambda_1 S \text{tr}(T_1 \Delta_R^\dagger) + \lambda_2 S \text{tr}(T_2 \Delta_R) + \text{h.c.} \right]$$

We use MicrOMEGAs and MadDM to compute relic density and direct detection.

$$y_{ST} = \sqrt{\lambda_1^2 + \lambda_2^2} \quad \tan \theta_{ST} = \lambda_1 / \lambda_2$$

---  $M_{\chi_1^0}$  [TeV]  
 —  $\Omega_\chi h^2 = 0.12$



**MicrOMEGAs.**



# Singlet-Bidoublet

Fields	Charges	Spin
S	(1,1,0)	1/2
B	(2,2,0)	1/2

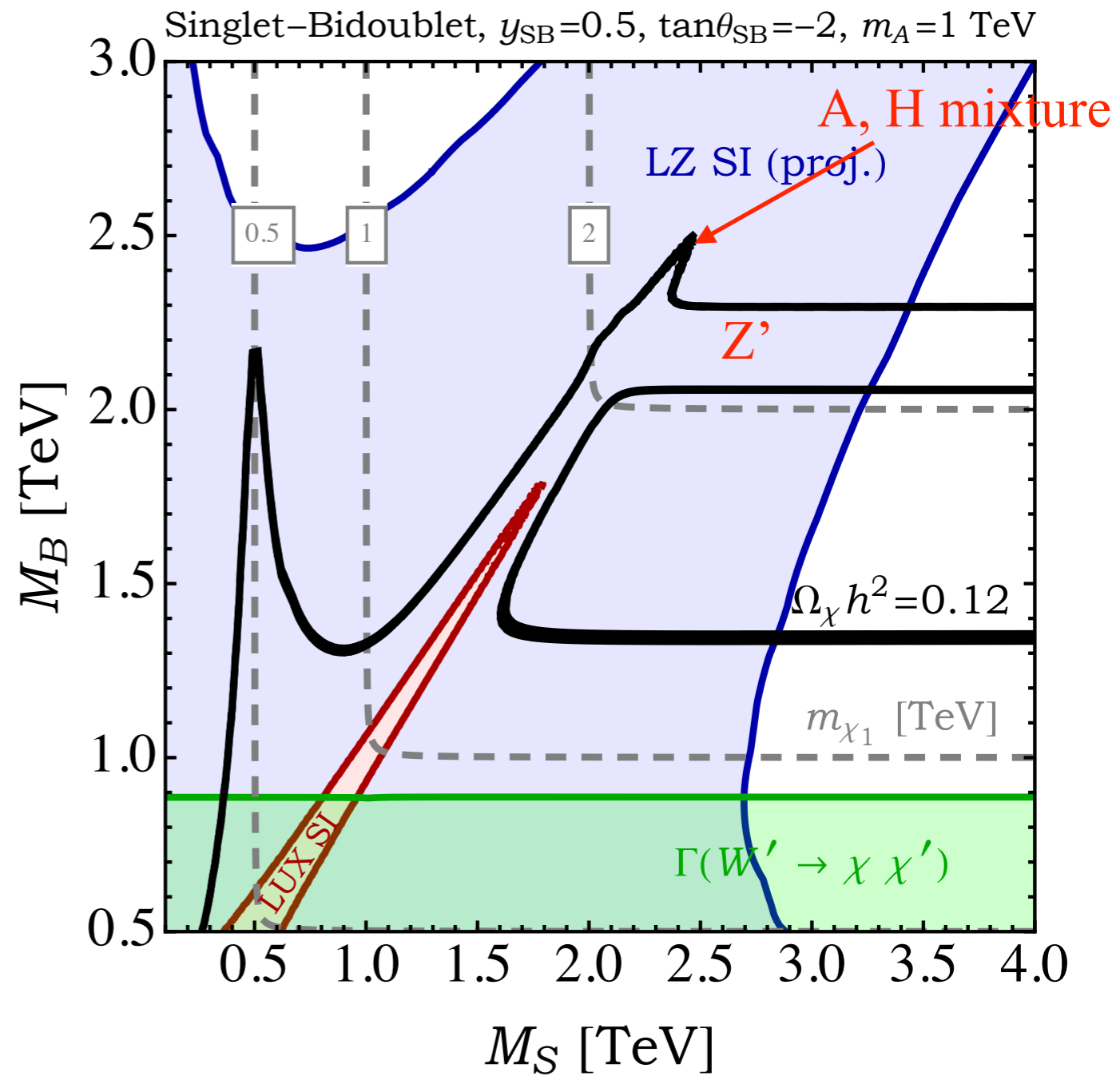
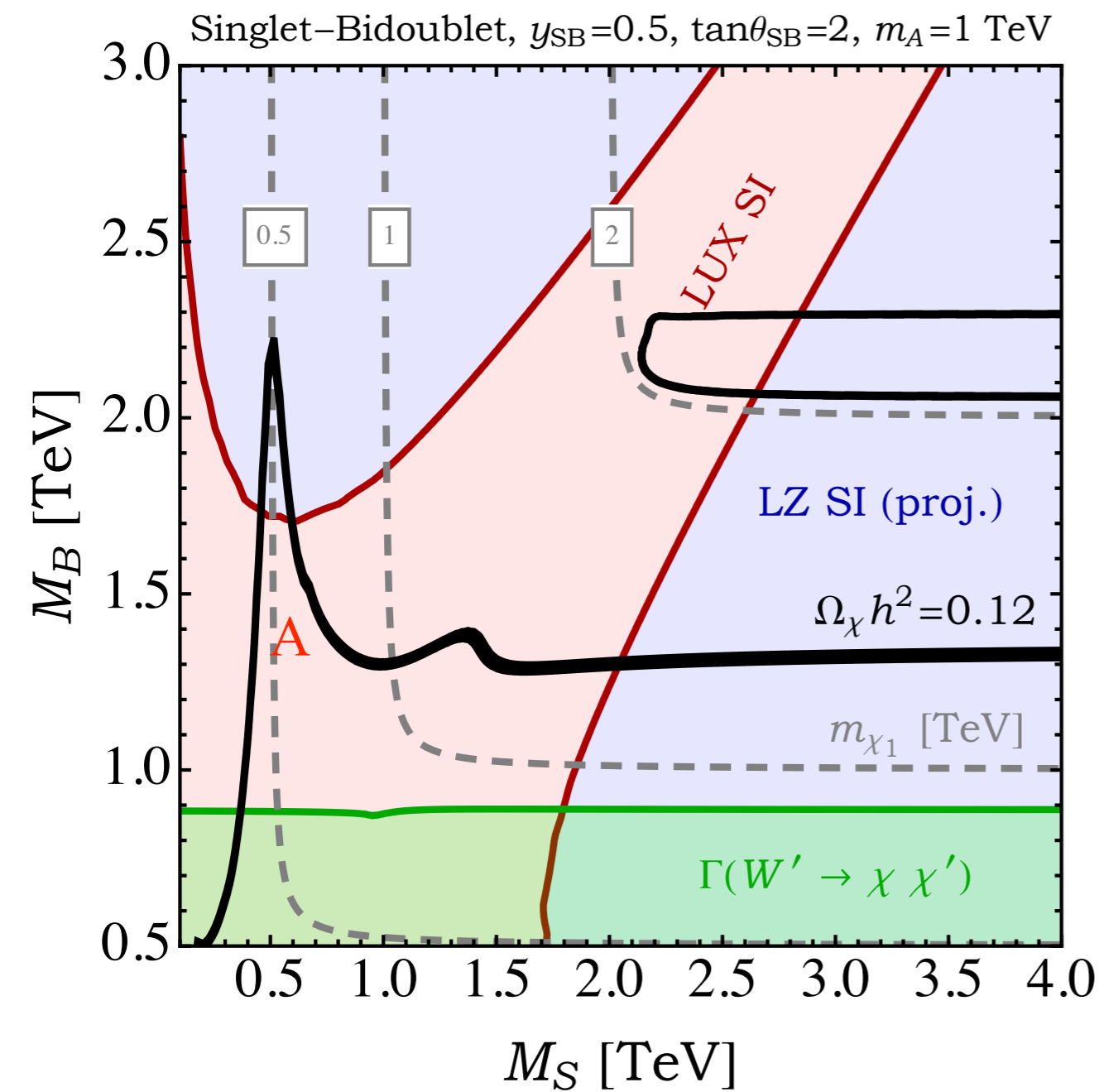
Fermion Bidoublet parametrized as:

$$B = \begin{pmatrix} b_1^0 & -b_2^+ \\ b_1^- & b_2^0 \end{pmatrix}$$

Singlet-Bidoublet Lagrangian:

$$\mathcal{L} \supset S^\dagger i\bar{\sigma}^\mu \partial_\mu S + tr(B^\dagger i\bar{\sigma}^\mu D_\mu B) - \left[ \frac{1}{2} M_s S^2 + \frac{1}{2} M_B tr(B\tilde{B}^\dagger) \right. \\ \left. + \lambda S tr(B\phi^\dagger) + \tilde{\lambda} S tr(B\tilde{\phi}^\dagger) + \text{h.c.} \right]$$

Parametrization  $y_{SB} = \sqrt{\lambda^2 + \tilde{\lambda}^2}$   $\tan \theta_{SB} = \lambda/\tilde{\lambda}$  ← Yukawas



# Triplet-Bidoublet

Fields	Charges	Spin
T	(1,3,0)	1/2
B	(2,2,0)	1/2

Fermion Bidoublet parametrized:

$$B = \begin{pmatrix} b_1^0 & -b_2^+ \\ b_1^- & b_2^0 \end{pmatrix}$$

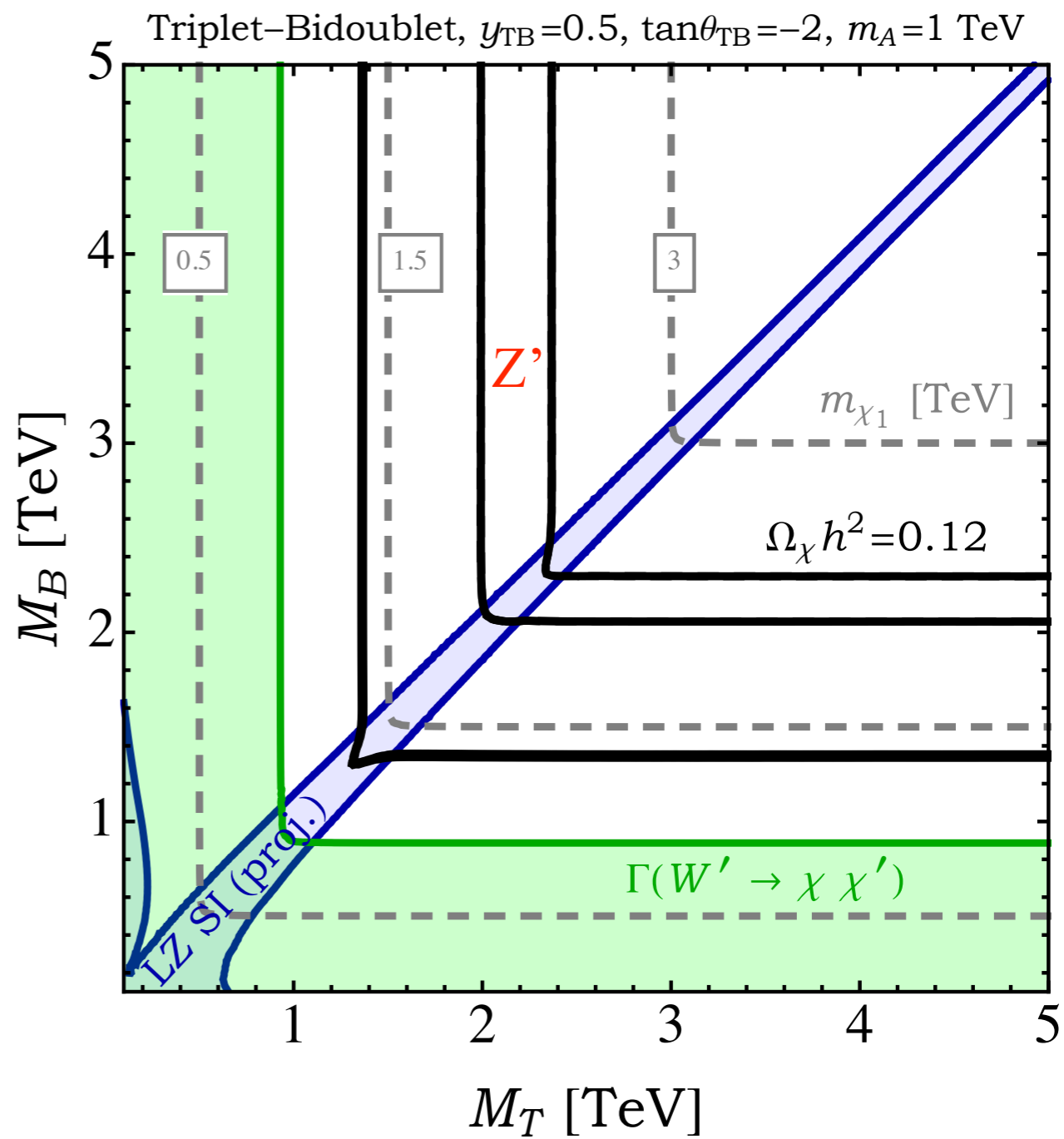
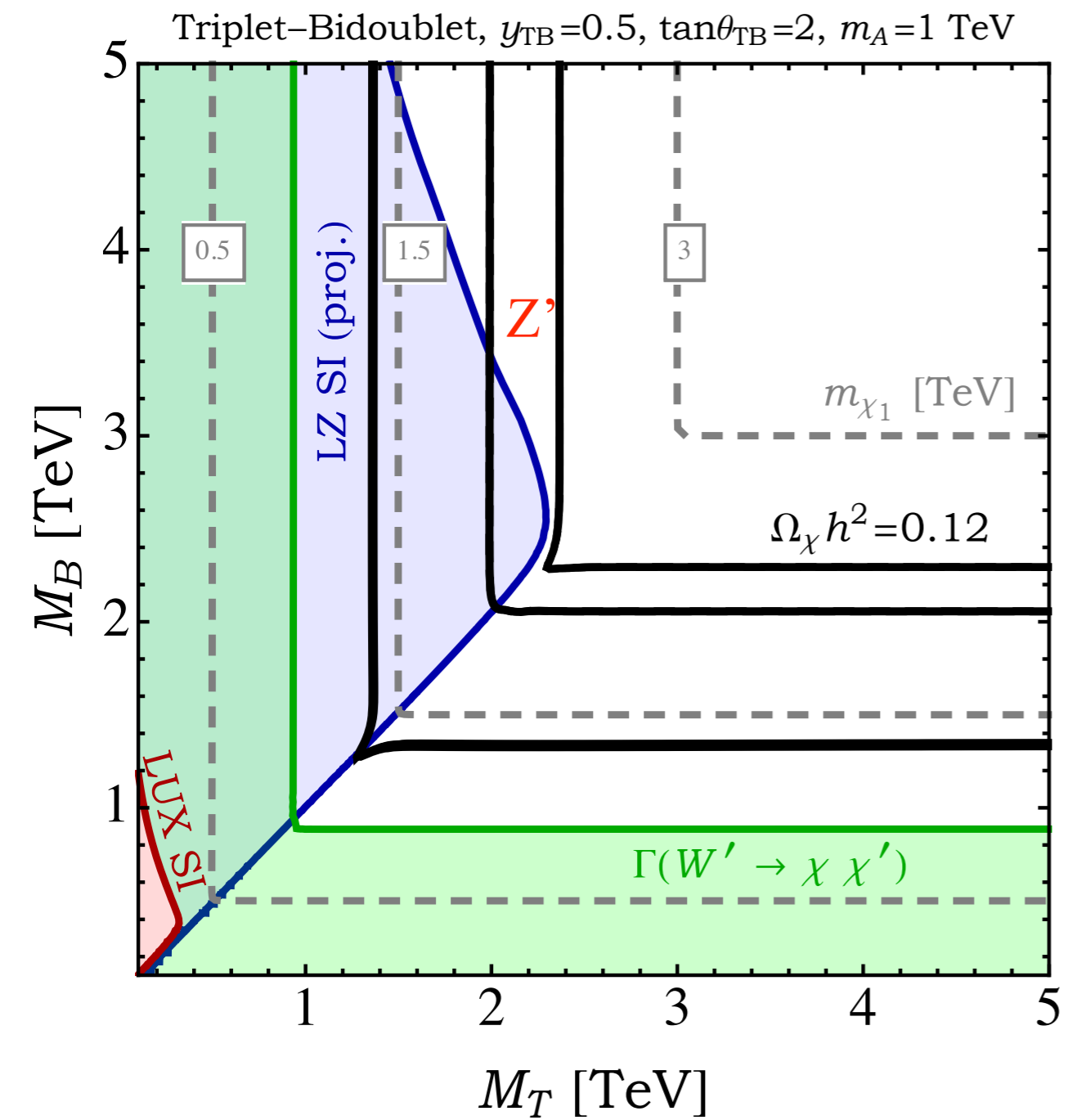
Fermion Triplet parametrized:

$$T = \begin{pmatrix} t^0/\sqrt{2} & t_2^+ \\ t_1^- & -t^0/\sqrt{2} \end{pmatrix}$$

Triplet-Bidoublet Lagrangian:

$$\mathcal{L} \supset tr(T^\dagger i\bar{\sigma}^\mu D_\mu T) + tr(B^\dagger i\bar{\sigma}^\mu D_\mu B) - \left[ \frac{1}{2} M_T tr(T^2) + \frac{1}{2} M_B tr(B\tilde{B}^\dagger) \right. \\ \left. + \lambda tr(BT\phi^\dagger) + \tilde{\lambda} tr(BT\tilde{\phi}^\dagger) + \text{h.c.} \right]$$



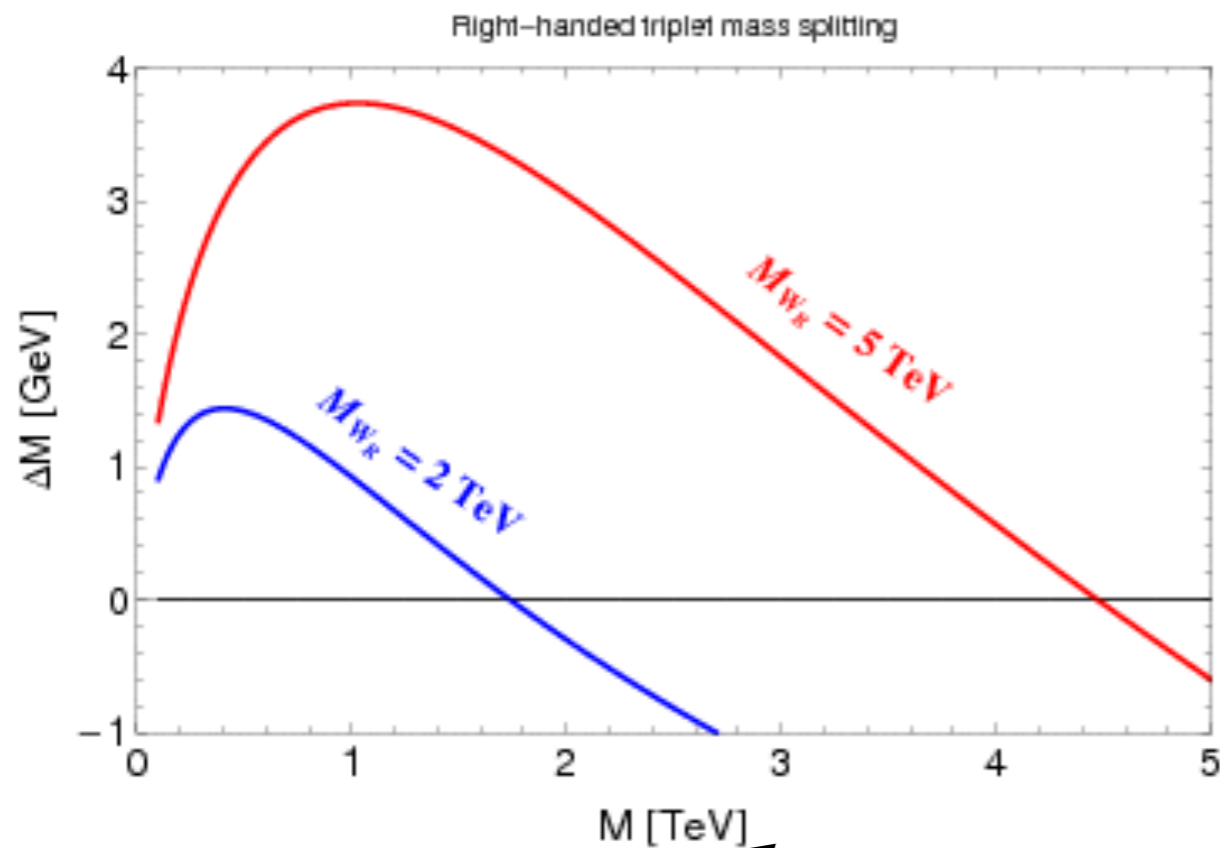


# Neutrality...

Point in parameter space where “chargino” lighter than lightest “neutralino”.

Then we would have charged DM

Loop corrections can provide large contributions to mass splitting, making “neutralino” heavier than “chargino”.



$$\Delta M = M_+ - M_0$$

Minimal Left-Right symmetric DM  
Heeck & Patra  
arXiv:1507.01584

DM mass

# Dark Matter Stability

**Stability** - breaking of quantum numbers under unbroken  $Z_2$  in  $U(1)_{B-L}$   
 $Z_2$  subgroup remains when

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$



$$SU(2)_L \otimes U(1)_Y$$



Fermions odd under  $Z_2$   
Bosons even under  $Z_2$

Lightest Fermion is stable if it carries even B-L charge.

$$O = X \times (\# \text{ of bosons}) \times (\text{odd } \# \text{ of SM fermions})$$

Dark matter gauge eigenstate with lightest neutral component  $\chi$  and B-L charge  $Q$

- B-L gauge invariance
- Lorentz invariance

DM unstable only if  $Q = \pm(1, 3, 5, \dots)/3$

# Conclusions

Why is this study important ???

LHC run 1 might have seen first hints of new physics.

Naturally incorporated in L-R models.

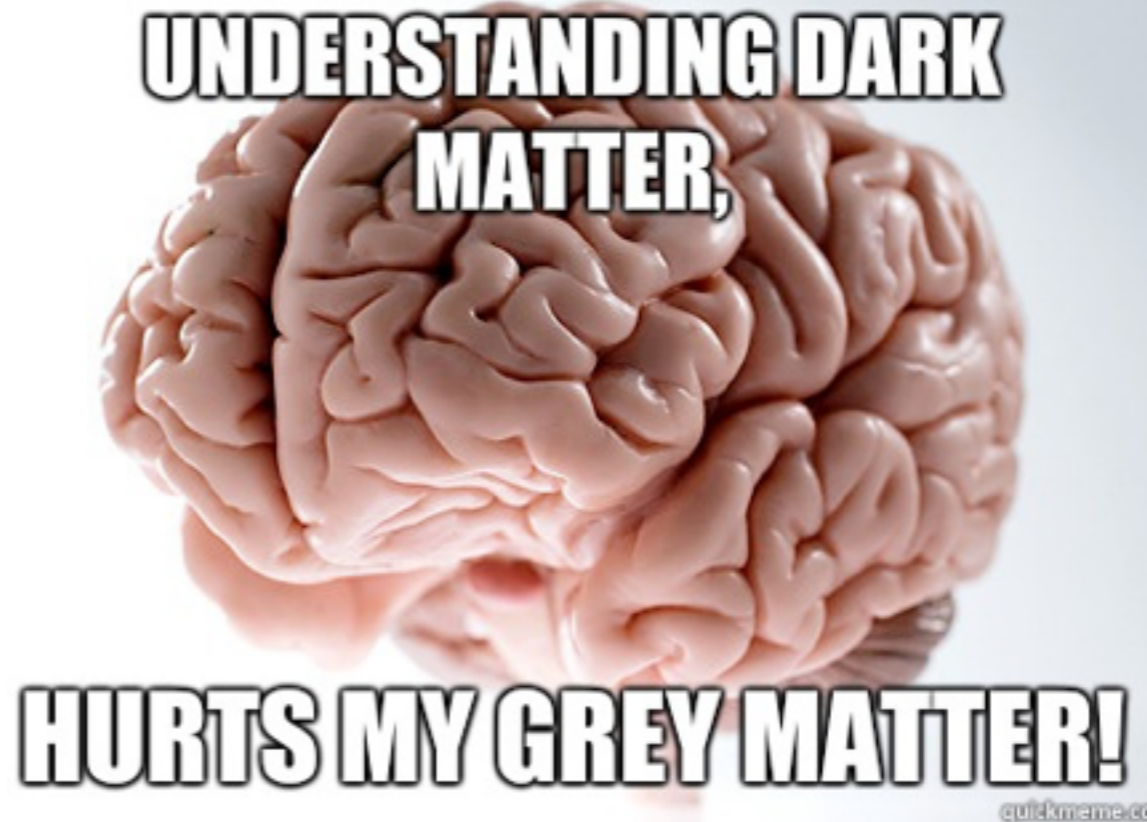
Naturally has DM, stability guaranteed by B-L gauge invariance.

Mixed DM mass splittings are not fixed unlike in pure DM.

Within reach of next generation DD experiments.

might be detected by LZ

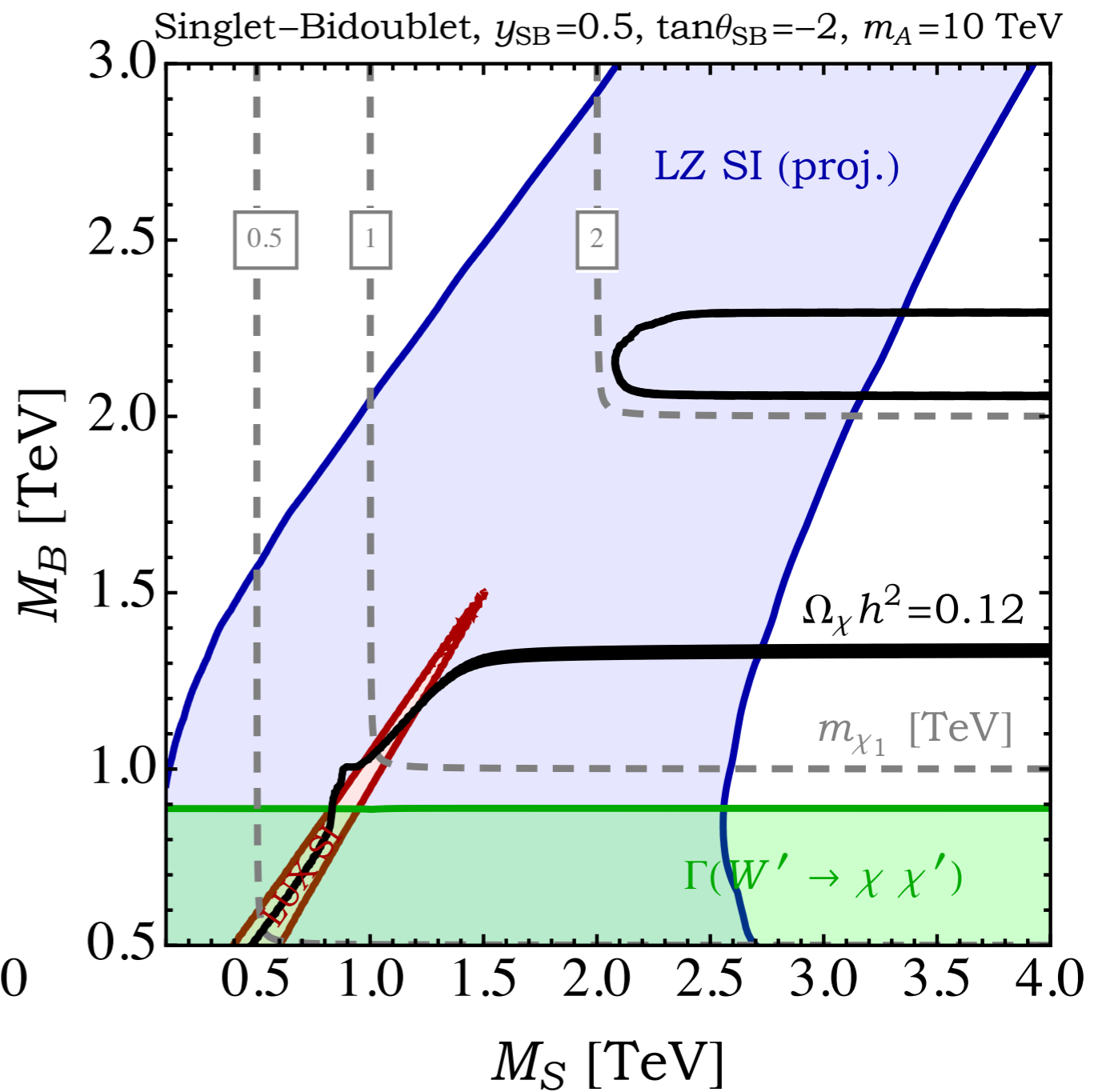
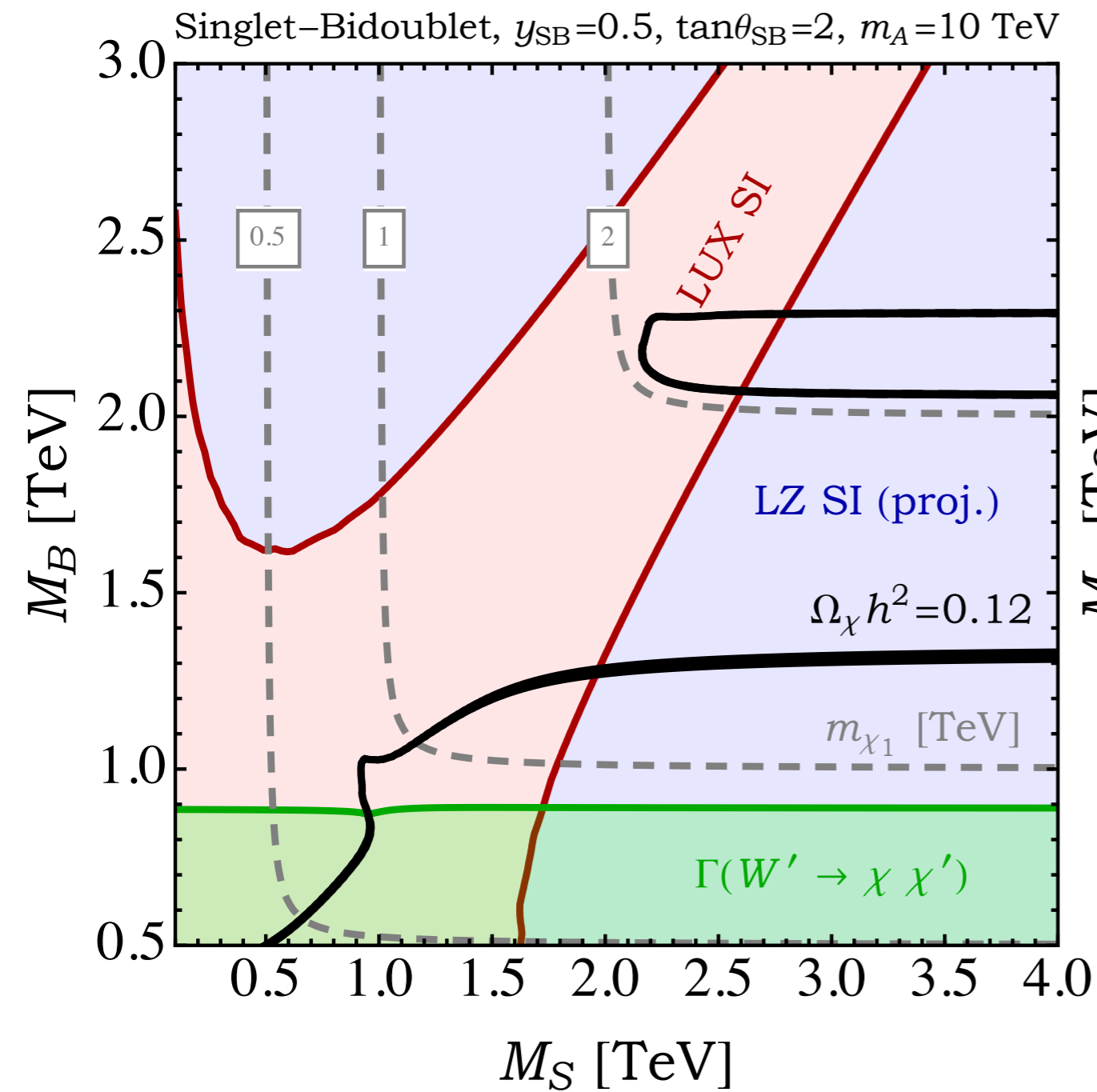
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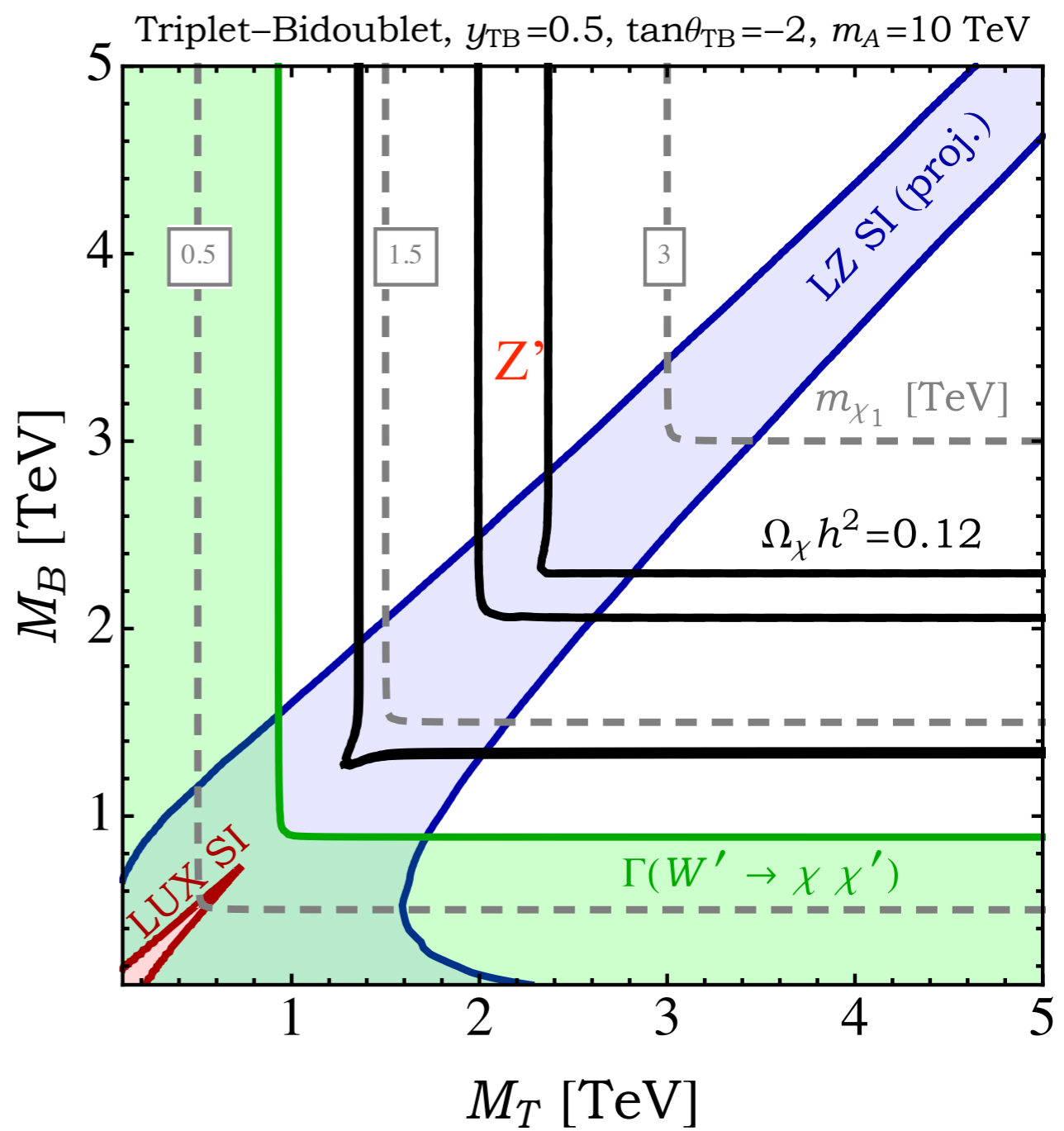
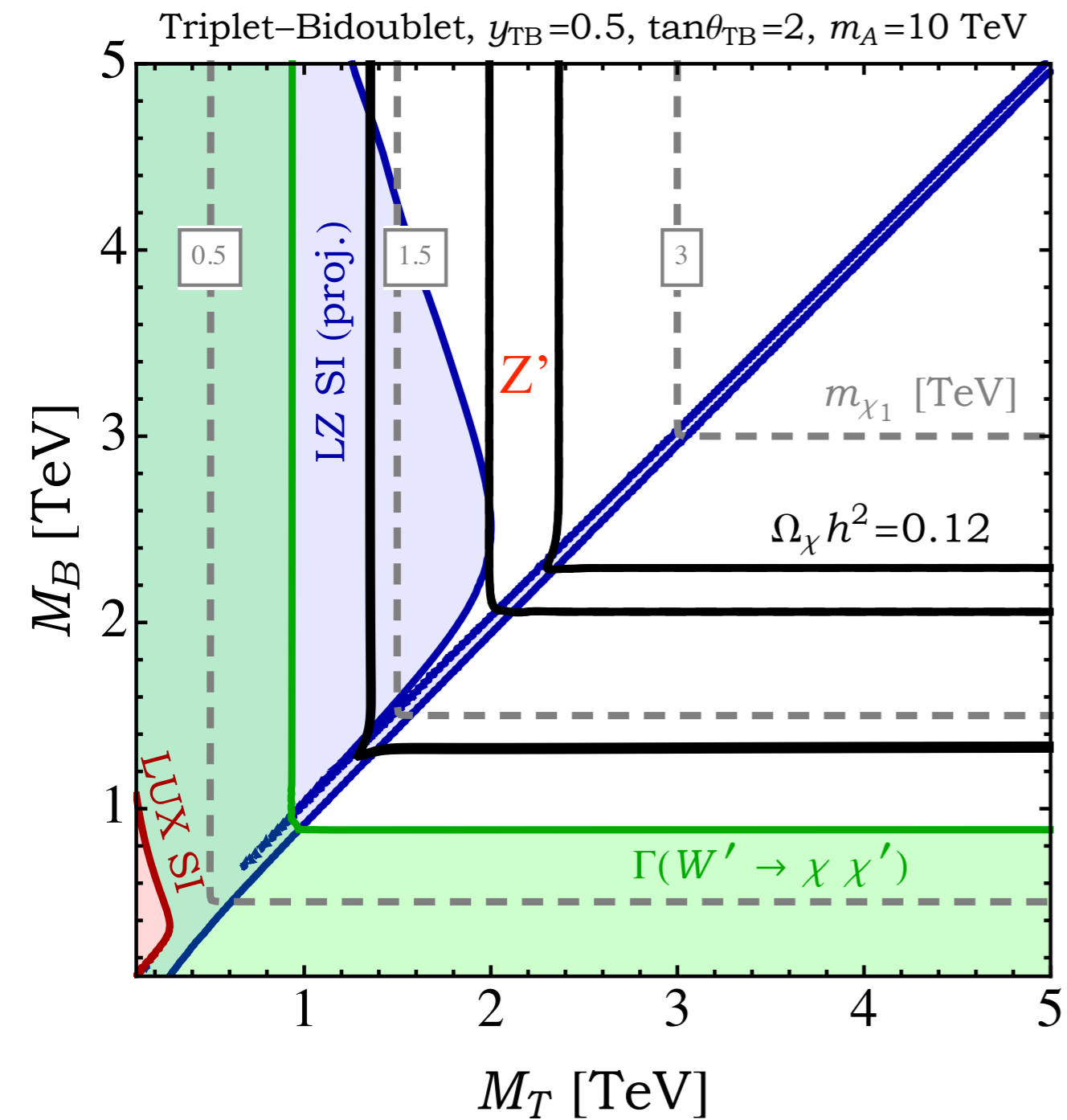
Questions? Comments?

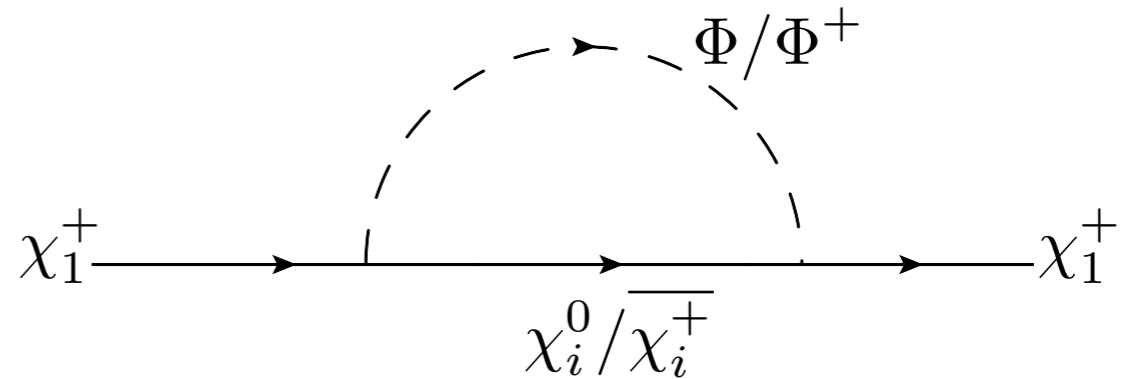
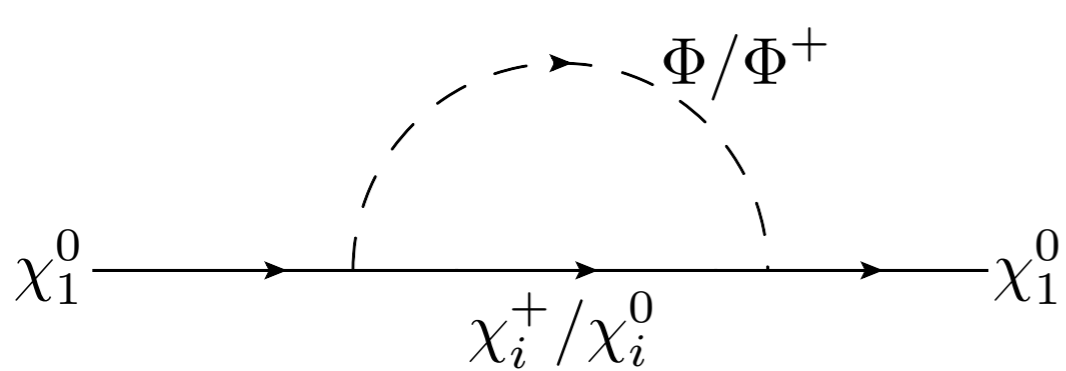
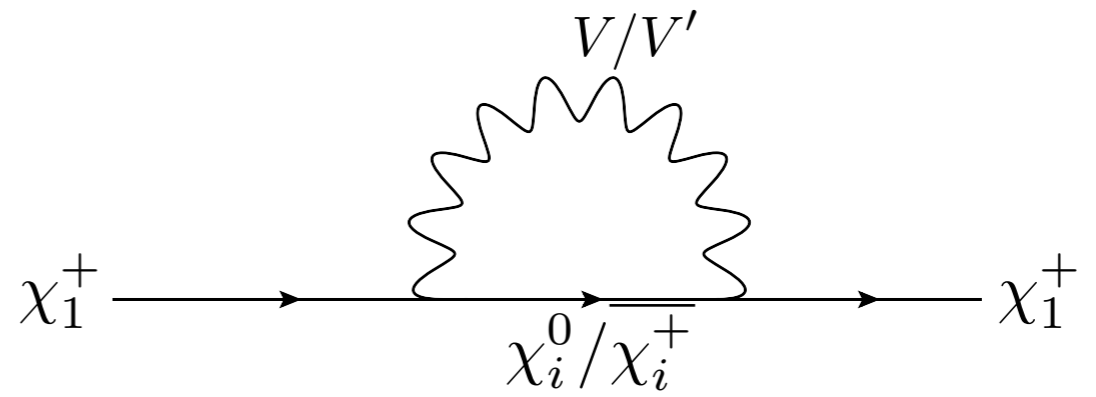
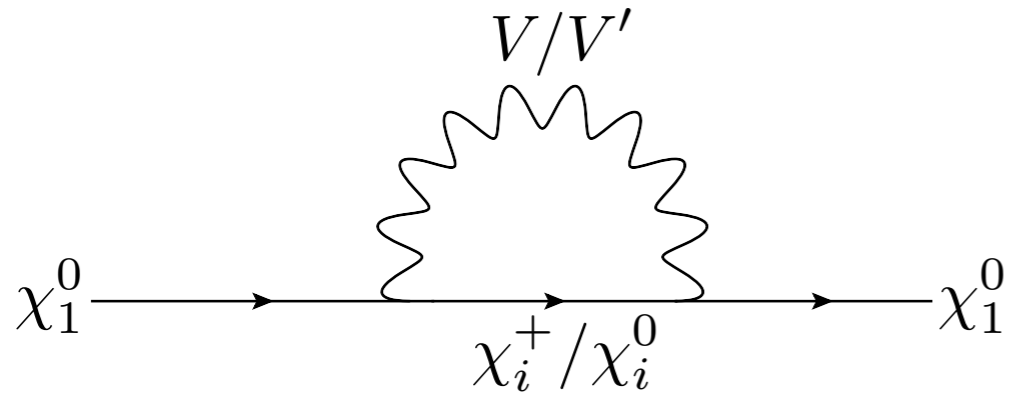


Back Up slides



Direct Detection dominated by SM higgs and heavy higgs interaction





$$V = Z/W^- \quad V' = Z'/W'^- \quad \Phi = h/A/H \quad \Phi^+ = H^\pm$$

$\Delta M = M_{\chi^+} - M_{\chi^0}$  could help constrain parameter space further.

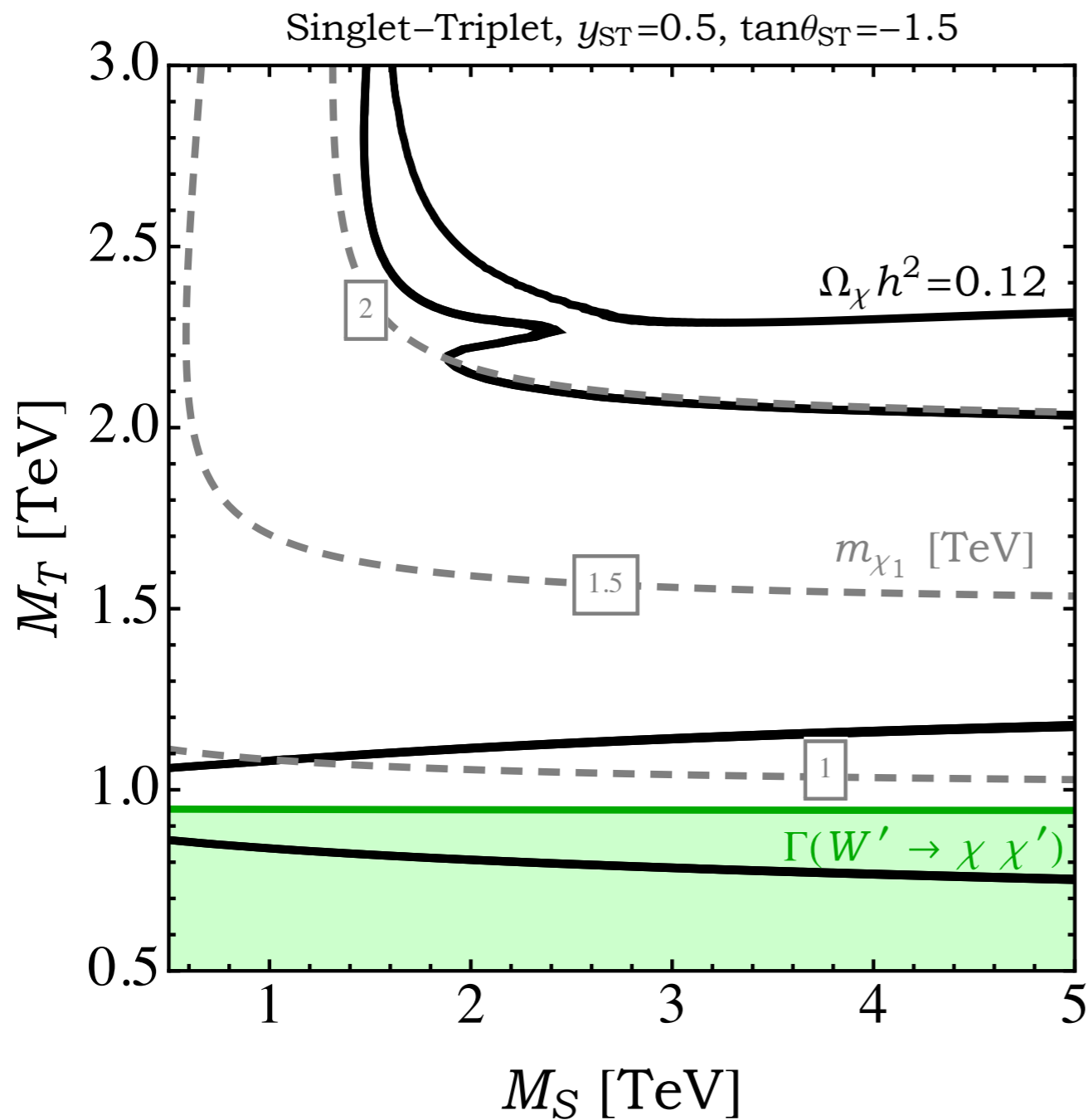
Please see appendix E

Berlin, Fox, Hooper & GM (arXiv:1604.06100)

# Singlet-Triplet

$$M_{\chi^\pm} < M_{\chi_1} \quad M_S \gtrsim 50 \text{ TeV} \gg M_T$$

$$\tan \theta_{ST} = -1.5$$

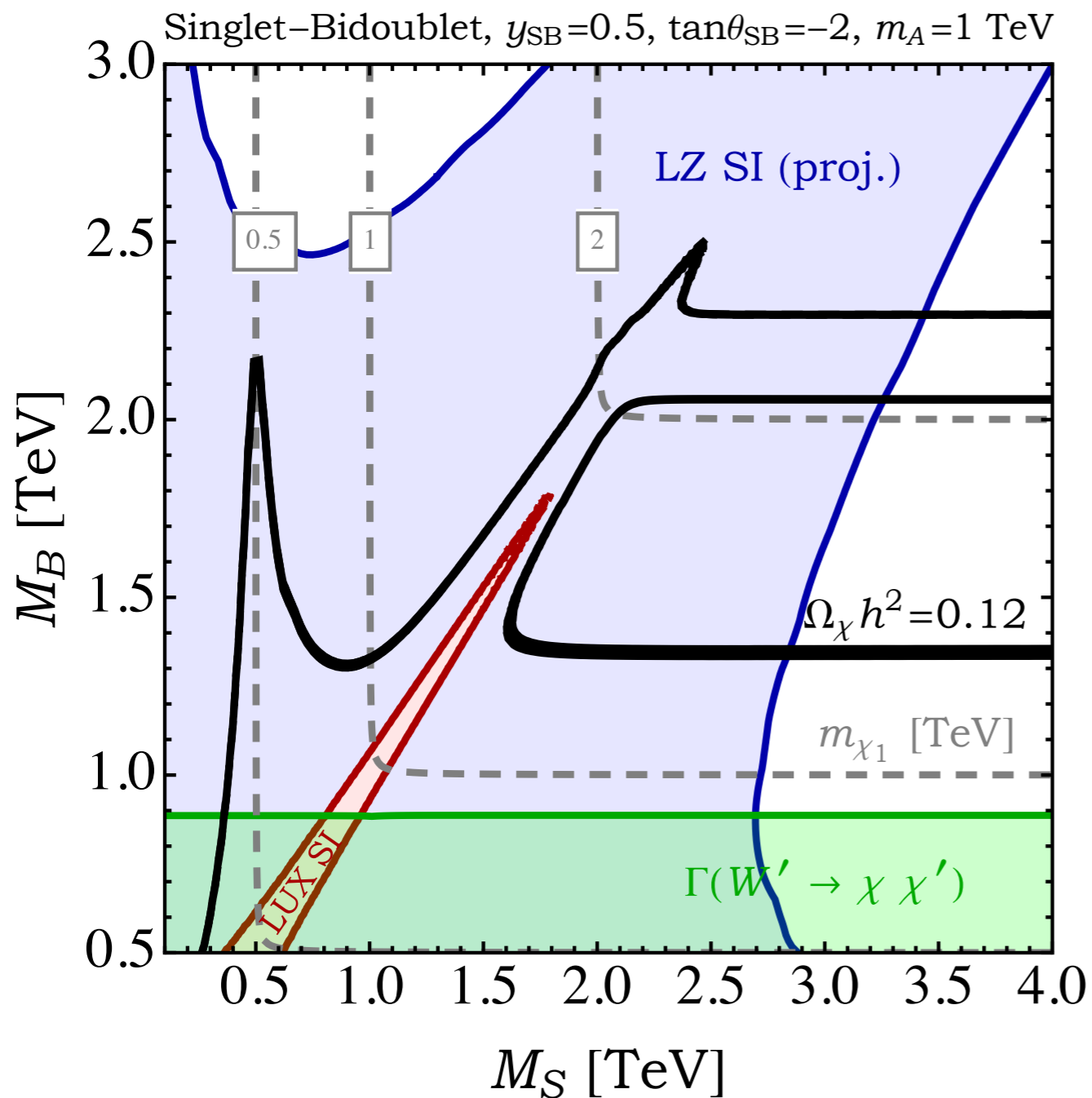


↑  
roughly here

# Singlet-Bidoublet

$$M_{\chi^\pm} < M_{\chi_1} \quad M_S \gtrsim 50 \text{ TeV} \gg M_B$$

$$\tan \theta_{SB} = -2$$



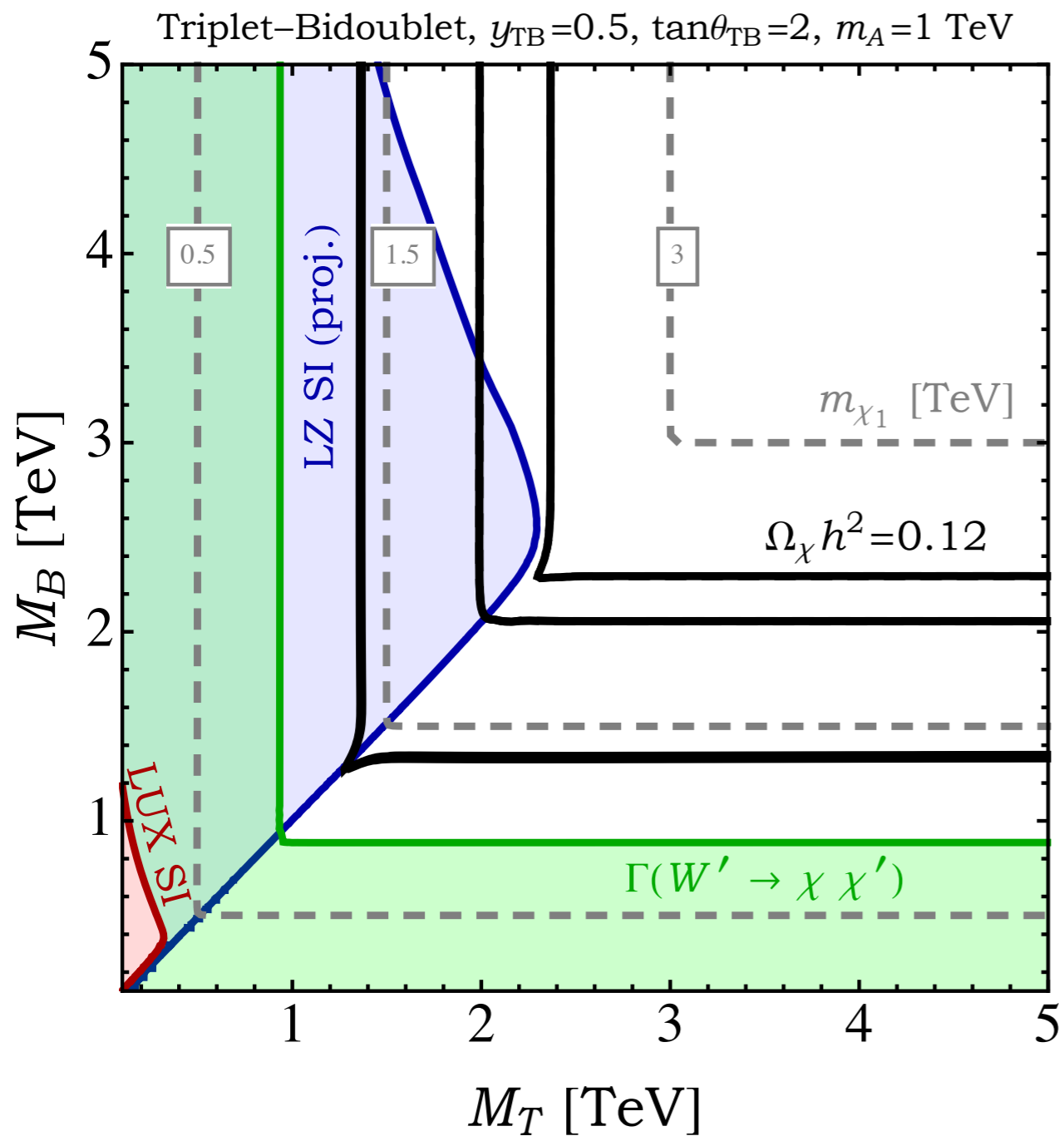
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# Triplet-Bidoublet

$$M_{\chi^\pm} < M_{\chi_1} \quad M_B \ll M_T \gtrsim \mathcal{O}(10) \text{ TeV}$$

$$M_T \ll M_B \gtrsim \mathcal{O}(10) \text{ TeV}$$



roughly here

# Gauge Boson Sector

After Symmetry breaking,  $\Delta_R$  and  $\phi$  get vevs, Gauge Boson mass term:

For the Charged Bosons,  $W_L^\pm$  and  $W_R^\pm$

$$\mathcal{L}_{W^\pm} = (W_L^{+\mu} \quad W_R^{+\mu}) \begin{pmatrix} \frac{1}{2}g_L^2 v^2 & -\frac{1}{2}g_L g_R s_{2\beta} v^2 \\ -\frac{1}{2}g_L g_R s_{2\beta} v^2 & g_R^2(v_R^2 + \frac{1}{2}v^2) \end{pmatrix} \begin{pmatrix} W_L^{-\mu} \\ W_R^{-\mu} \end{pmatrix}$$

For the Neutral Bosons,  $W_L^3$ ,  $W_R^3$  and  $B$

$$\mathcal{L}_{W^3, B} = \frac{1}{2} (W_{L\mu}^3 \quad W_{R\mu}^3 \quad B_\mu) \begin{pmatrix} \frac{1}{2}g_L^2 v^2 & -\frac{1}{2}g_L g_R v^2 & 0 \\ -\frac{1}{2}g_L g_R v^2 & g_R^2(v_R^2 + \frac{1}{2}v^2) & -2g_{B-L} g_R v_R^2 \\ 0 & -2g_{B-L} g_R v_R^2 & 2g_{B-L}^2 v_R^2 \end{pmatrix} \begin{pmatrix} W_{L\mu}^3 \\ W_{R\mu}^3 \\ B_\mu \end{pmatrix}$$

Diagonalize both Charged and Neutral Gauge Boson Mass Matrices

Gives masses for  $\gamma$ ,  $Z$  and  $Z'$ ,  $W^\pm$  and  $W'^\pm$



# Important Model Parameters

Parameter space motivated by CMS & ATLAS 8 TeV Diboson excess

To match observed rates in diboson excess, we need:

$$M_{W'} = 1.8 \text{ TeV} - 2.0 \text{ TeV} \quad v_R = 3 \text{ TeV} - 4 \text{ TeV}$$

$$g_R = 0.4 - 0.6 \quad \tan \beta = 0.5 - 2$$

We choose  $M_{W'} = 1.9 \text{ TeV}$  and  $g_R = 0.45$  such that  $v_R = 4.2 \text{ TeV}$

and  $M_{Z'} = 4.41 \text{ TeV}$

We also choose  $\tan \beta = 2$

We choose to not deal with Higgs triplet and bidoublet potential terms since they add more parameters.

We work with  $M_h = 125 \text{ GeV}$  and  $M_{heavy} = 1 \text{ TeV}$

$|\tan \theta_{ST}| \sim 1$  there is an enhanced parity symmetry acting on triplets

singlet mixes with only lin comb of neutral triplet parts

other lin comb of neutrals degenerate with charged states

$$\sigma_{SI} \approx 2 \times 10^{-44} \text{ cm}^2 \times \left( \frac{\lambda_h^{(1)}}{0.1} \right)^2 \quad \text{Singlet-Bidoublet}$$

$$\lambda_h^{(1)} \propto (1 + \sin 2\beta \sin 2\theta_{SB}) m_{\chi_1} + (\sin 2\beta + \sin 2\theta_{SB}) M_B$$

## Triplet-Bidoublet

Destructive interference btw h and H DD diagrams

$$\lambda_h^{(1)} \propto (1 + \sin 2\beta \sin 2\theta_{\text{TB}})m_{\chi_1} - (\sin 2\beta + \sin 2\theta_{\text{TB}})M_B$$

Cancellations between higgs contributions happen  $M_T \gg M_B \sim m_{\chi_1}$

$$\lambda_h^{(1)} \propto (1 - s_{2\beta})(1 - s_{2\theta_{\text{TB}}})$$