

Simplified Collider Limits on New Interactions

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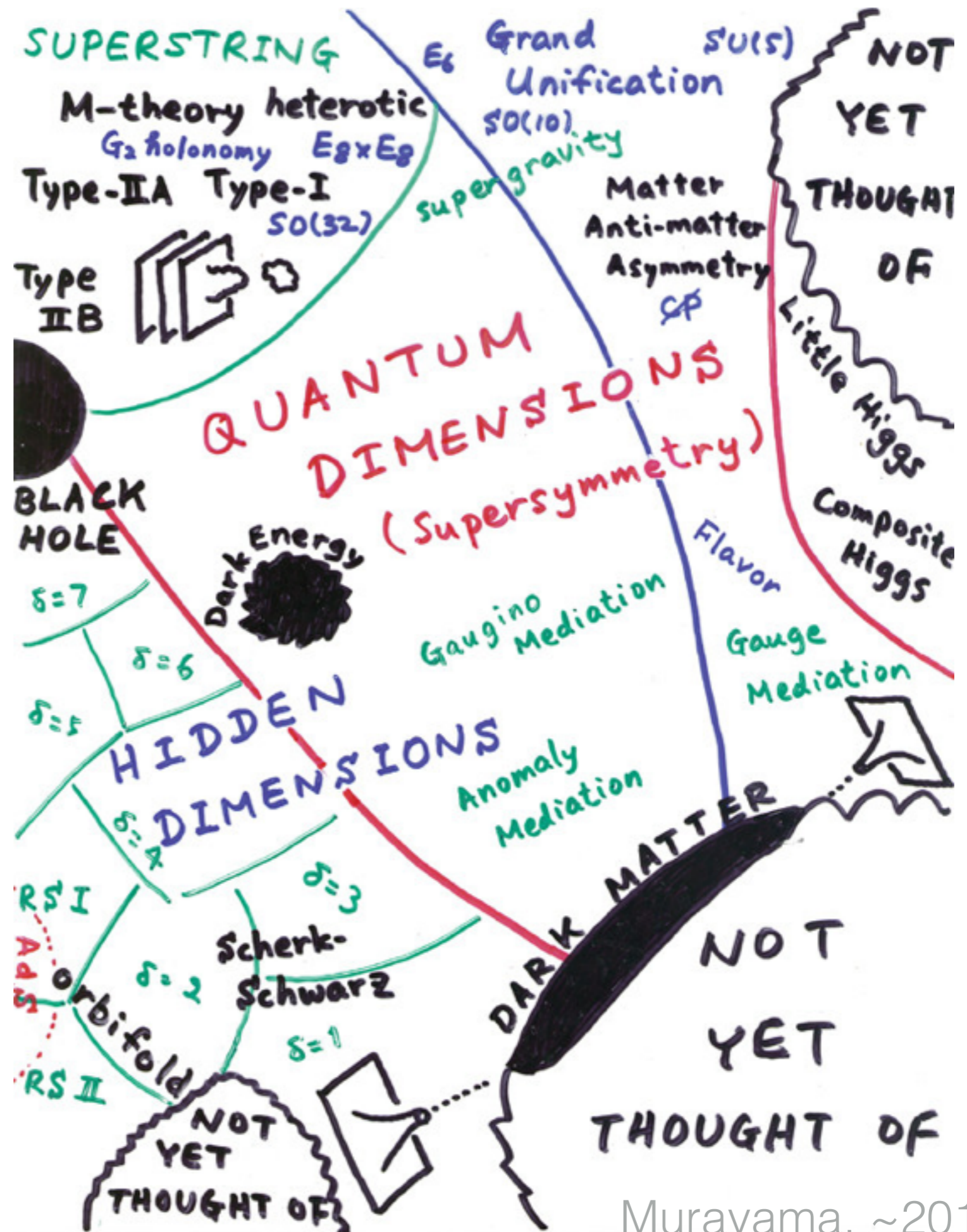
**Phenomenology 2016
Symposium**

9-11 May 2016
University of Pittsburgh
US/Eastern timezone

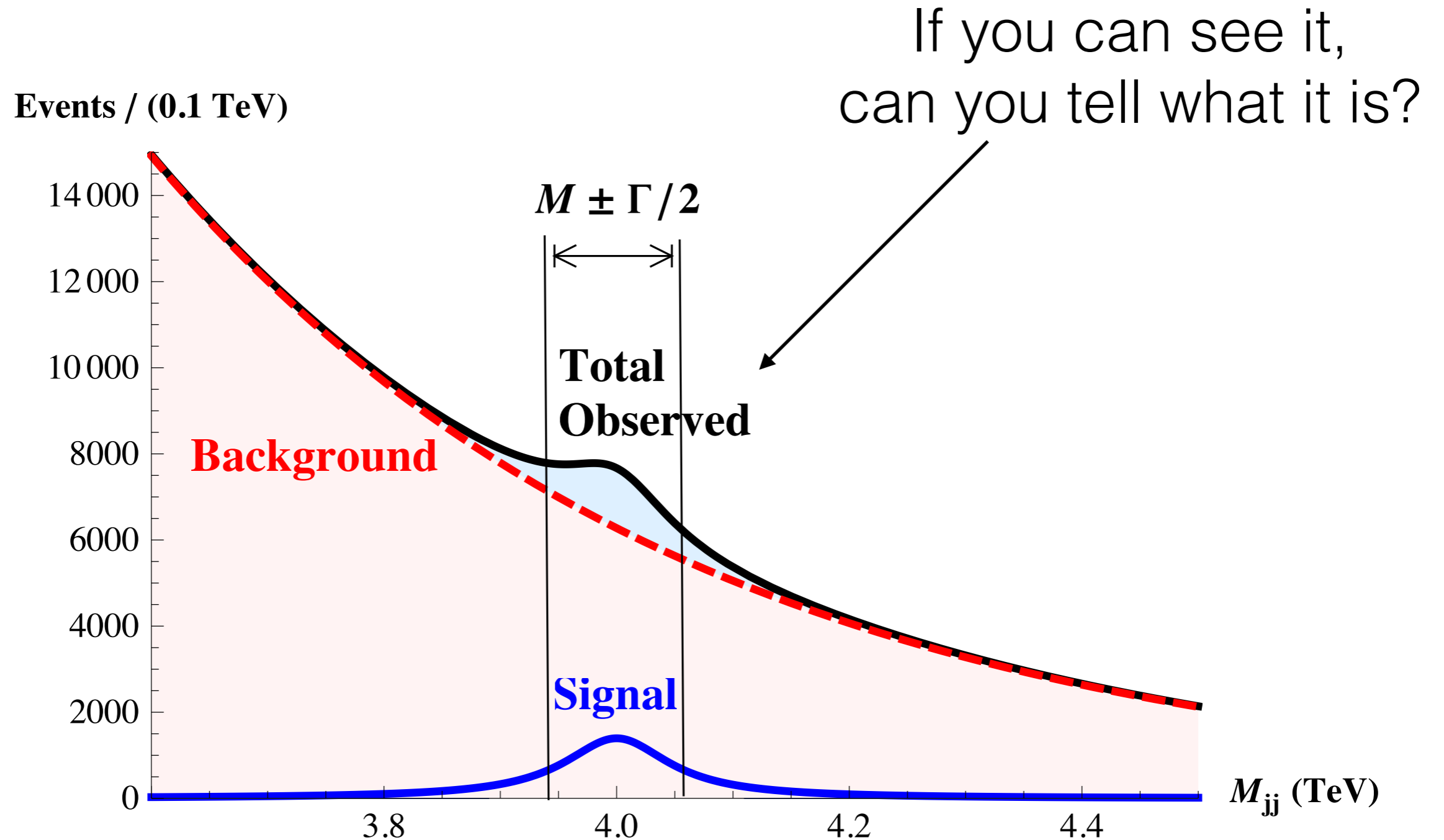
Outline

- New Interactions *typically* imply new resonances
- Traditional “model-independent” Limits
- Simplified Models for s-channel resonances
- “Simplified Limits” from Simplified Models
- Example: Diboson Resonances

New Physics
typically
implies the
existence of
new resonances!

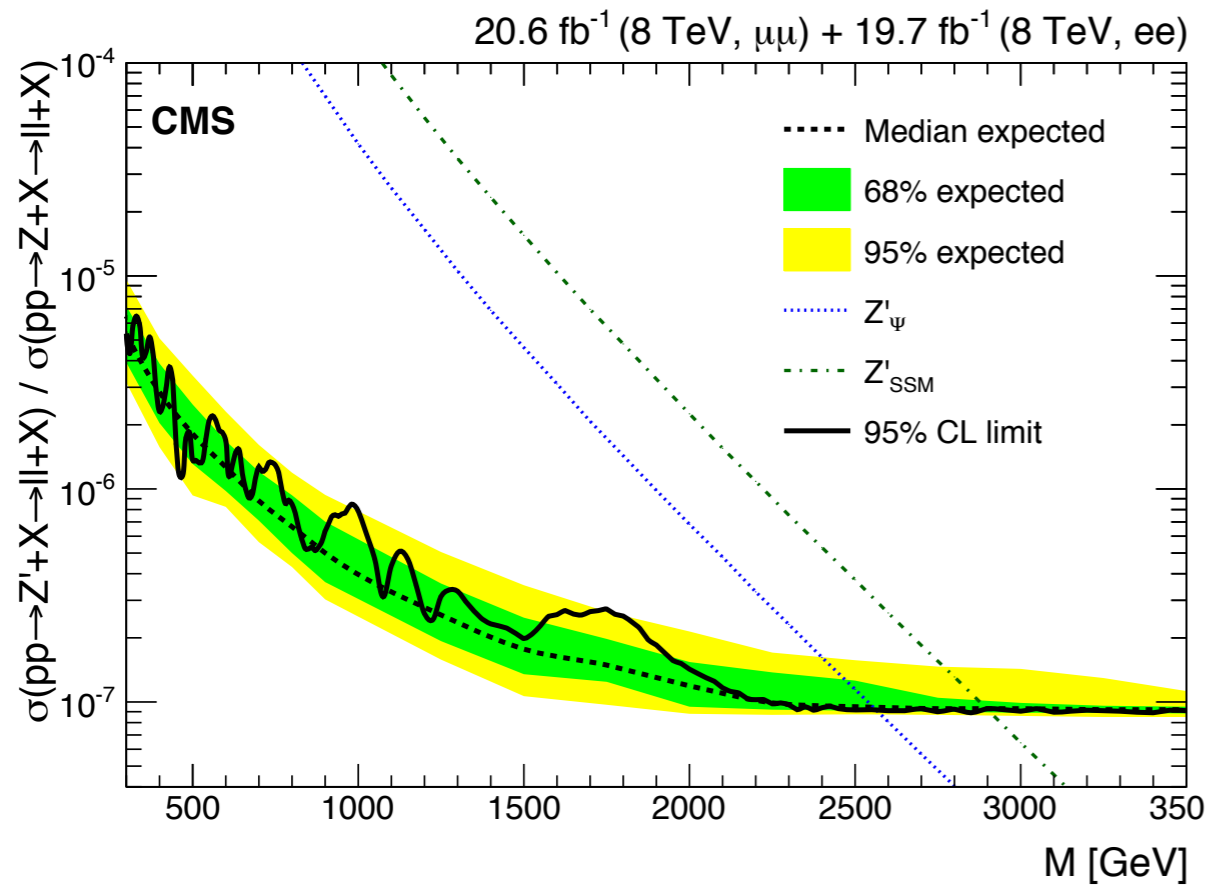


“s-channel” Resonance



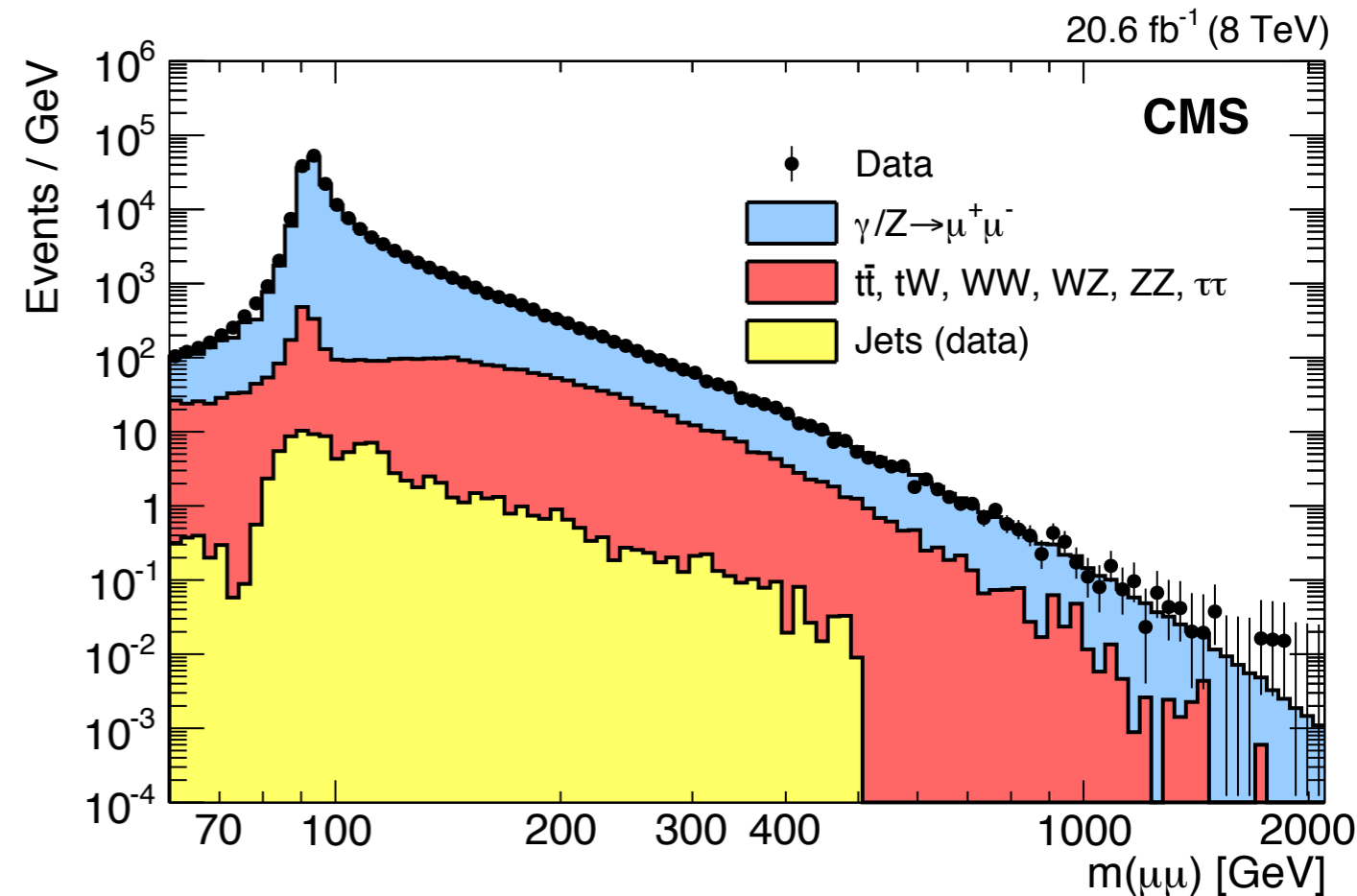
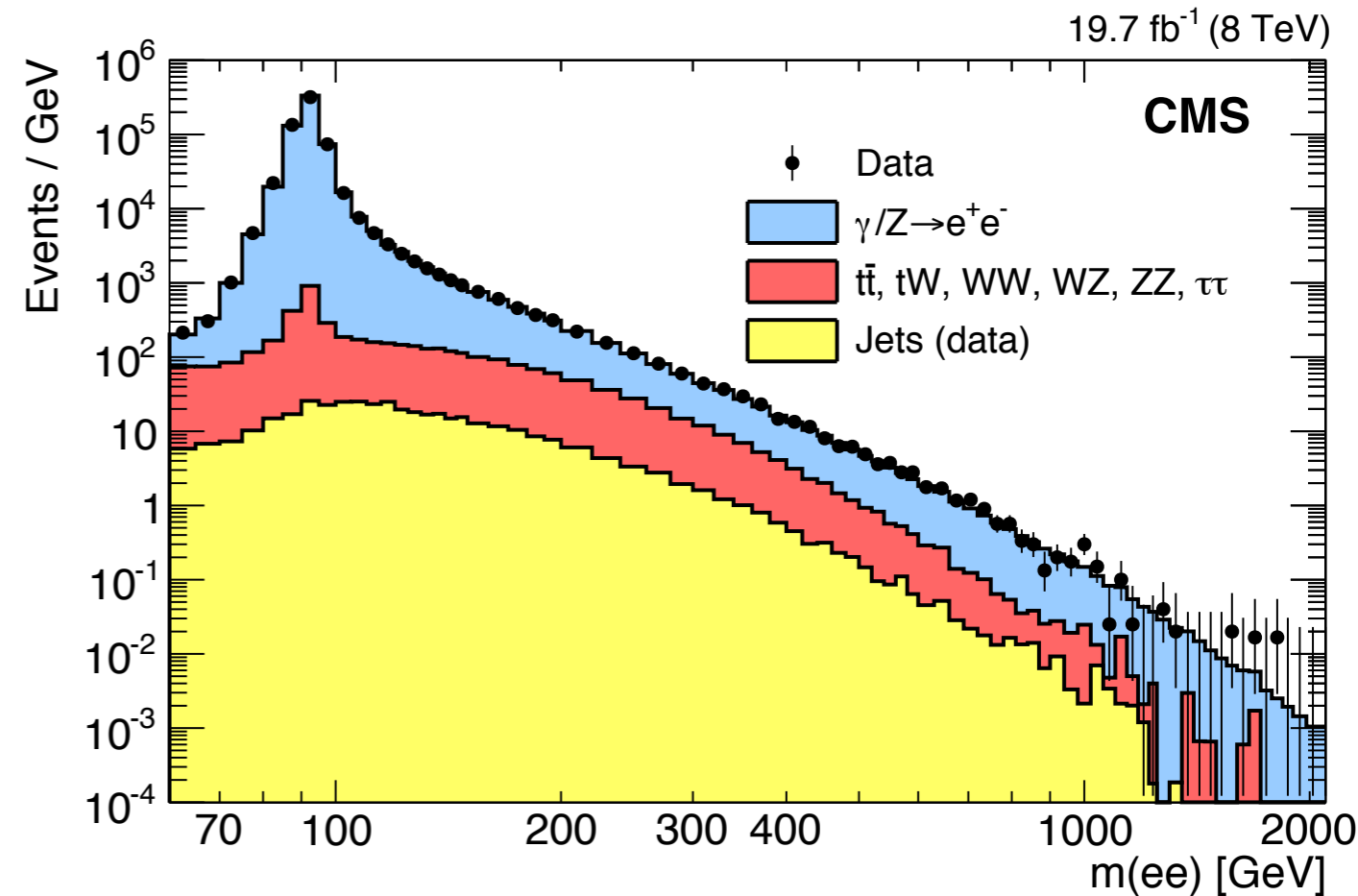
Dilepton Resonances

Flavor-independent Z'



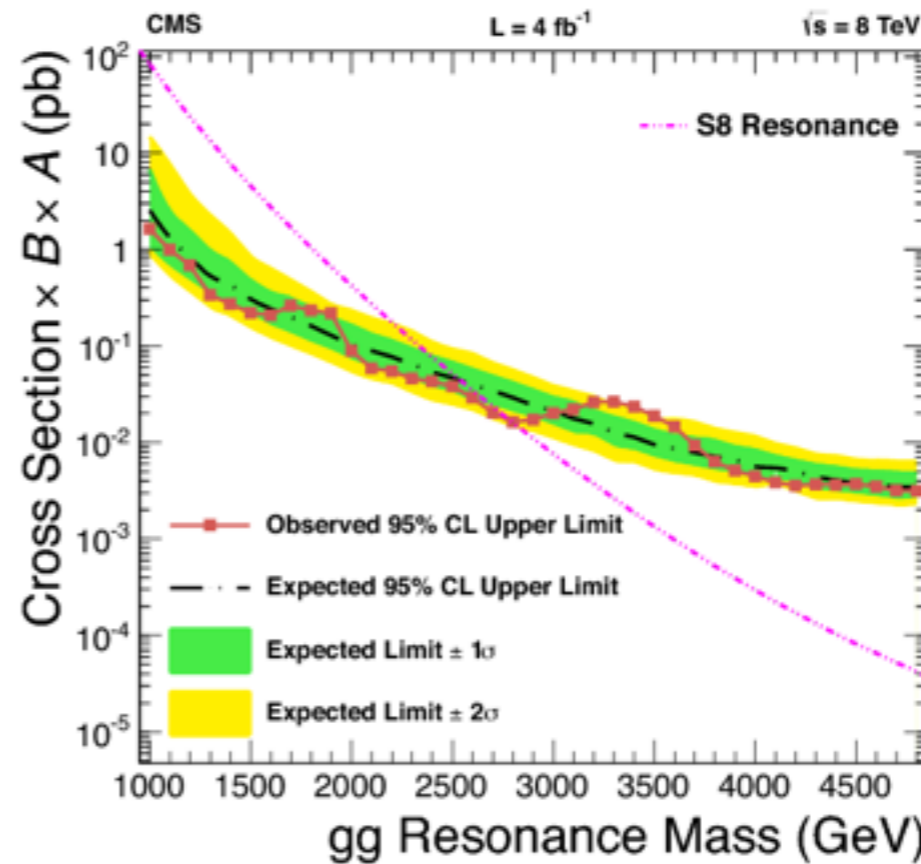
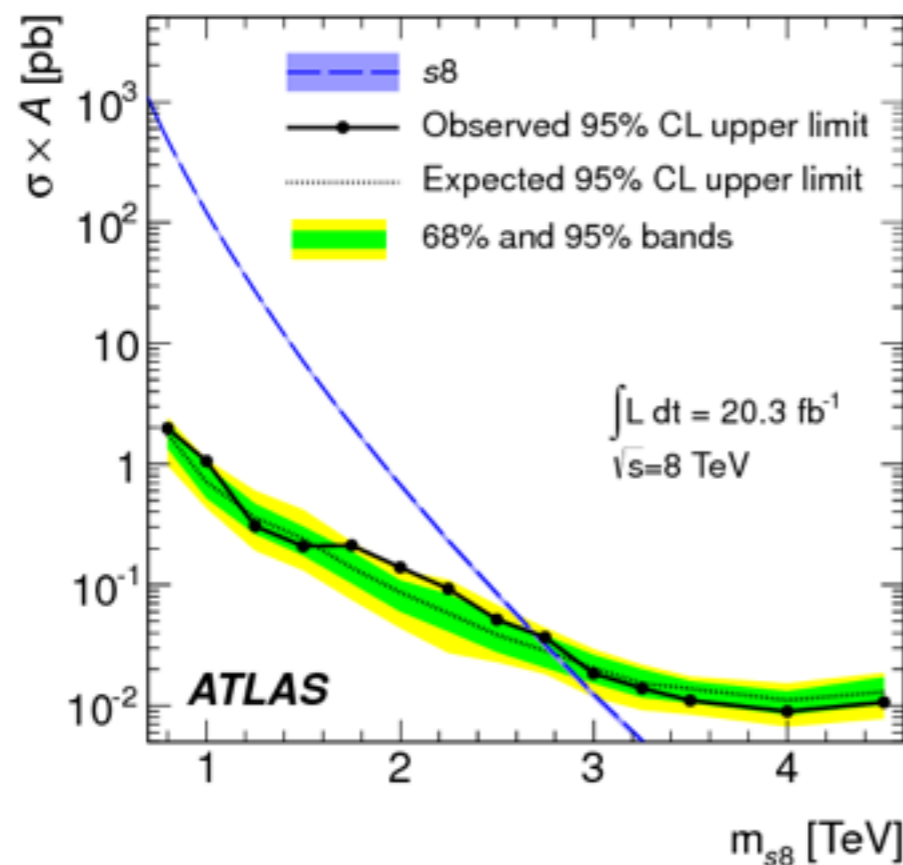
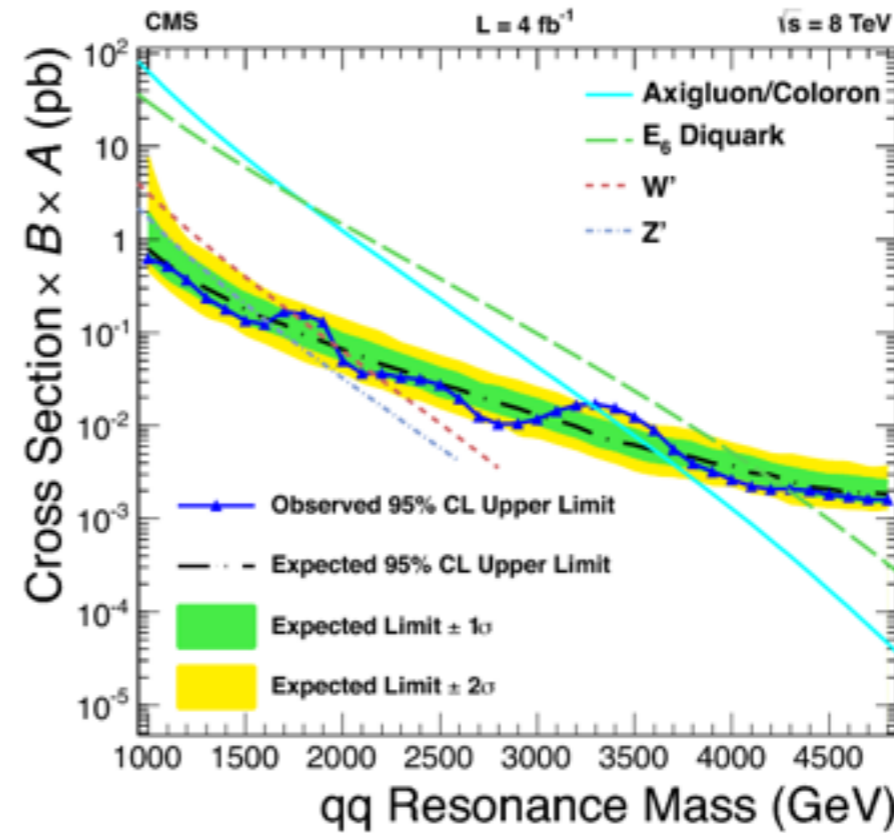
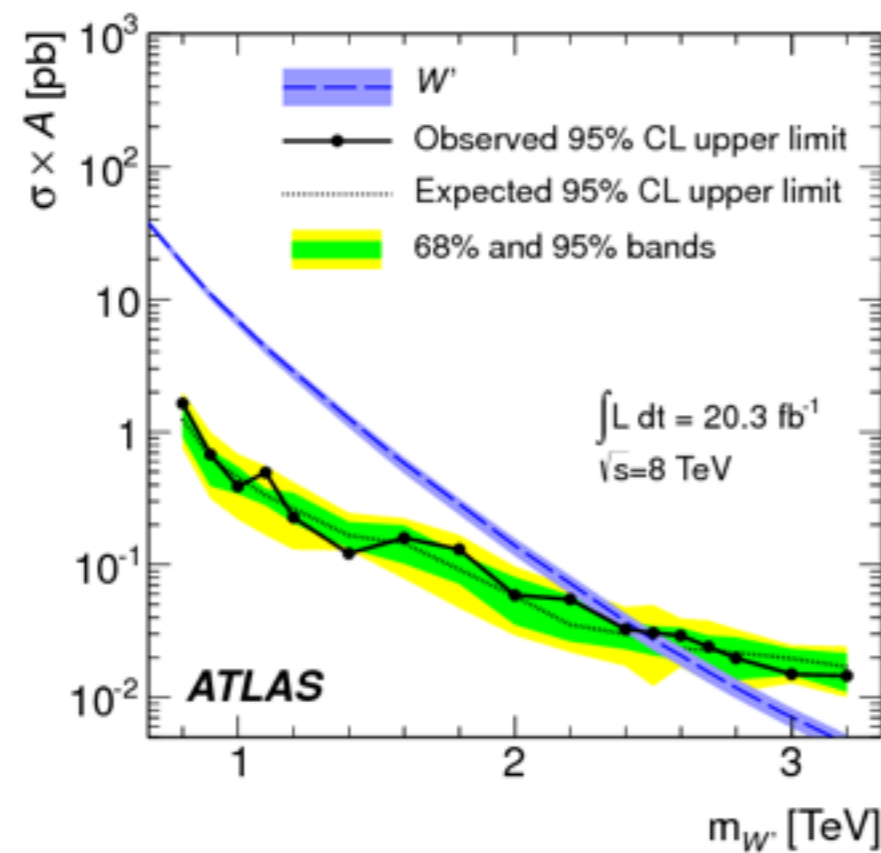
Limits: 2.5-3 TeV

CMS EXO-12-061

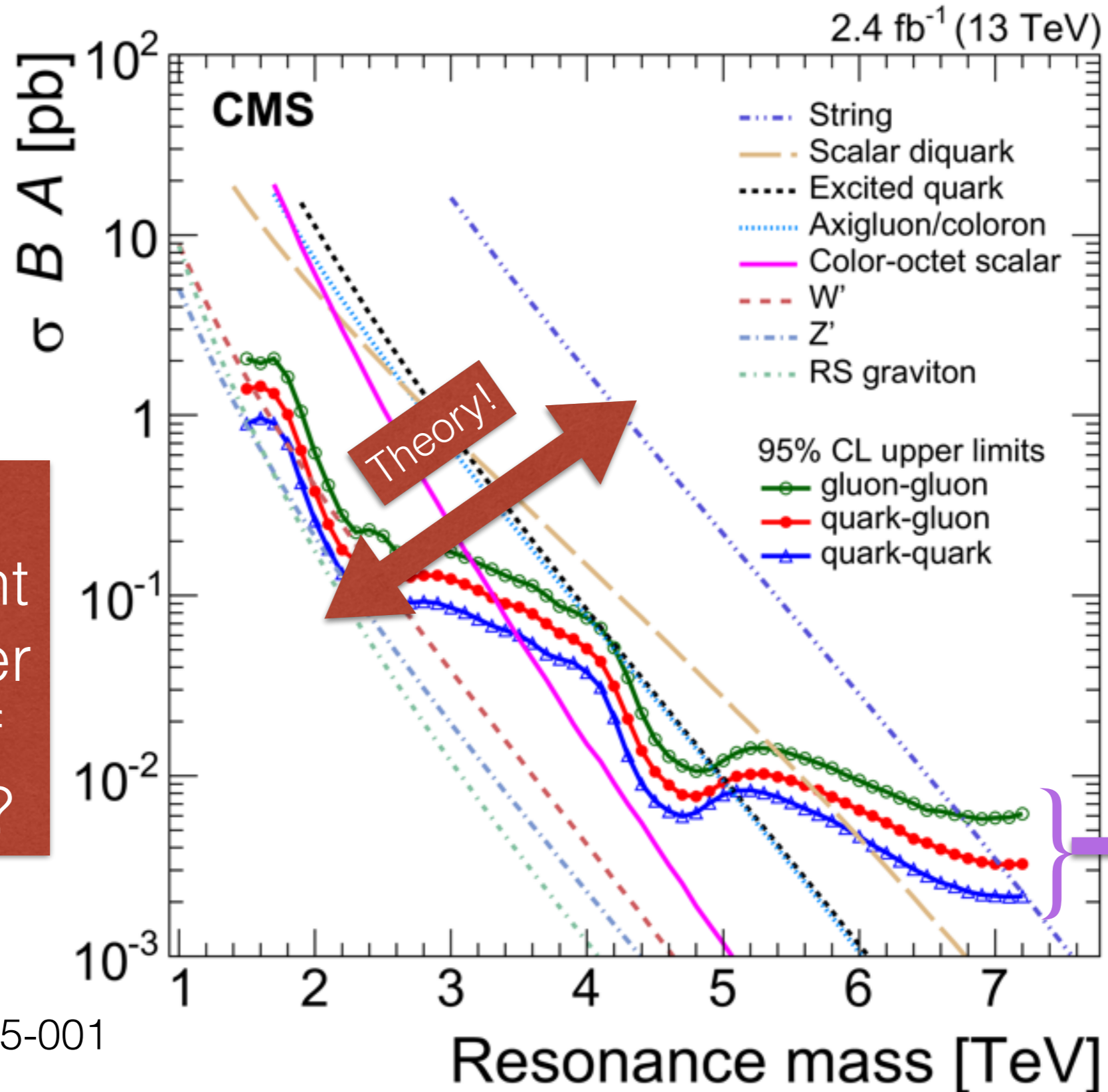


Dijet Resonances

ATLAS CERN-PH-EP-2014-147
CMS EXO-12-016



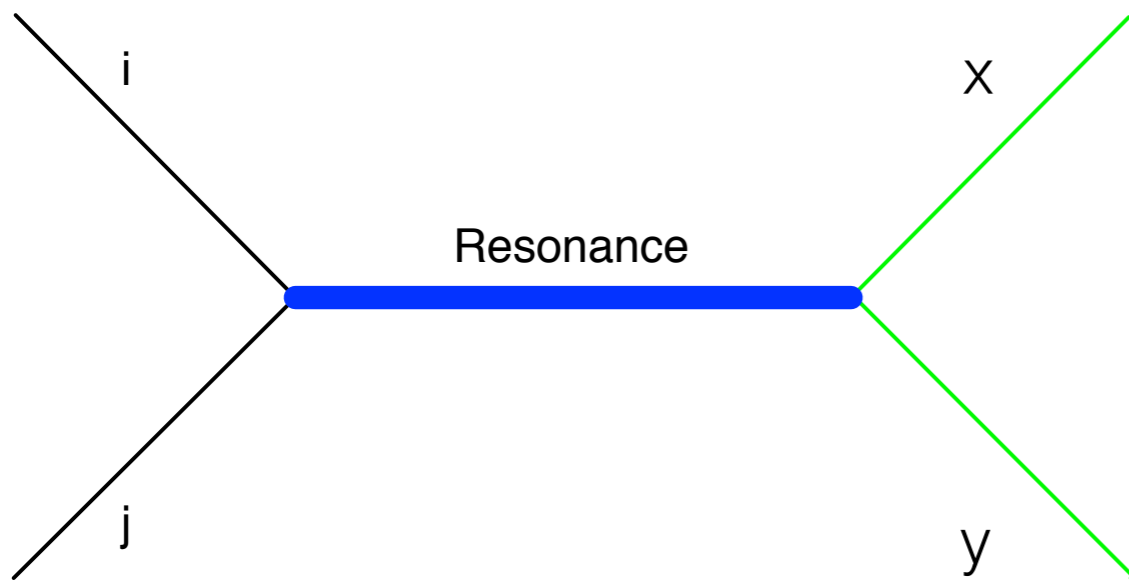
“Model-Independent” (dijet) Limits



Can we represent a broader class of models?

dijet characteristics

Simplified s-channel Model



$i, j = u, d, g, \gamma, W, Z$

$x, y = j, t, b, g, \gamma, W, Z, h$

Characteristics vs. Observables

- i, j : event characteristics
- *Couplings*: $BR, \sigma * BR$
- *Mass and width*: $d\sigma/dm_{ab}$
- *Spin*: $d\sigma/d\cos\theta_{ab}$
- x, y : in each decay channel
 - flavor tagging
 - jet substructure

NB: If x, y can be light quarks,
t-channel process may be relevant

Narrow Width Approximation

$$\sigma_R(pp \rightarrow x + y) = \int_{s_{min}}^{s_{max}} d\hat{s} \hat{\sigma}(\hat{s}) \cdot \left[\frac{dL^{ij}}{d\hat{s}} \right]$$

$$\hat{\sigma}_{ij \rightarrow R \rightarrow xy}(\hat{s}) = 16\pi(1 + \delta_{ij}) \cdot \mathcal{N} \cdot \frac{\Gamma(R \rightarrow i + j) \cdot \Gamma(R \rightarrow x + y)}{(\hat{s} - m_R^2)^2 + m_R^2 \Gamma_R^2}, \quad \mathcal{N} = \frac{N_{S_R}}{N_{S_i} N_{S_j}} \cdot \frac{C_R}{C_i C_j}$$

$$\frac{1}{(\hat{s} - m_R^2)^2 + m_R^2 \Gamma_R^2} \approx \frac{\pi}{m_R \Gamma_R} \delta(\hat{s} - m_R^2)$$

$$\sigma_R(pp \rightarrow x + y) = 16\pi^2 \cdot \mathcal{N} \cdot \frac{\Gamma_R}{m_R} \cdot (1 + \delta_{ij}) BR(R \rightarrow ij) \cdot BR(R \rightarrow xy) \left[\frac{1}{s} \frac{dL^{ij}}{d\tau} \right]_{\tau = \frac{m_R^2}{s}}$$

(Note: Can be corrected for K-factor(s) & Acceptance)

Simplified Limits

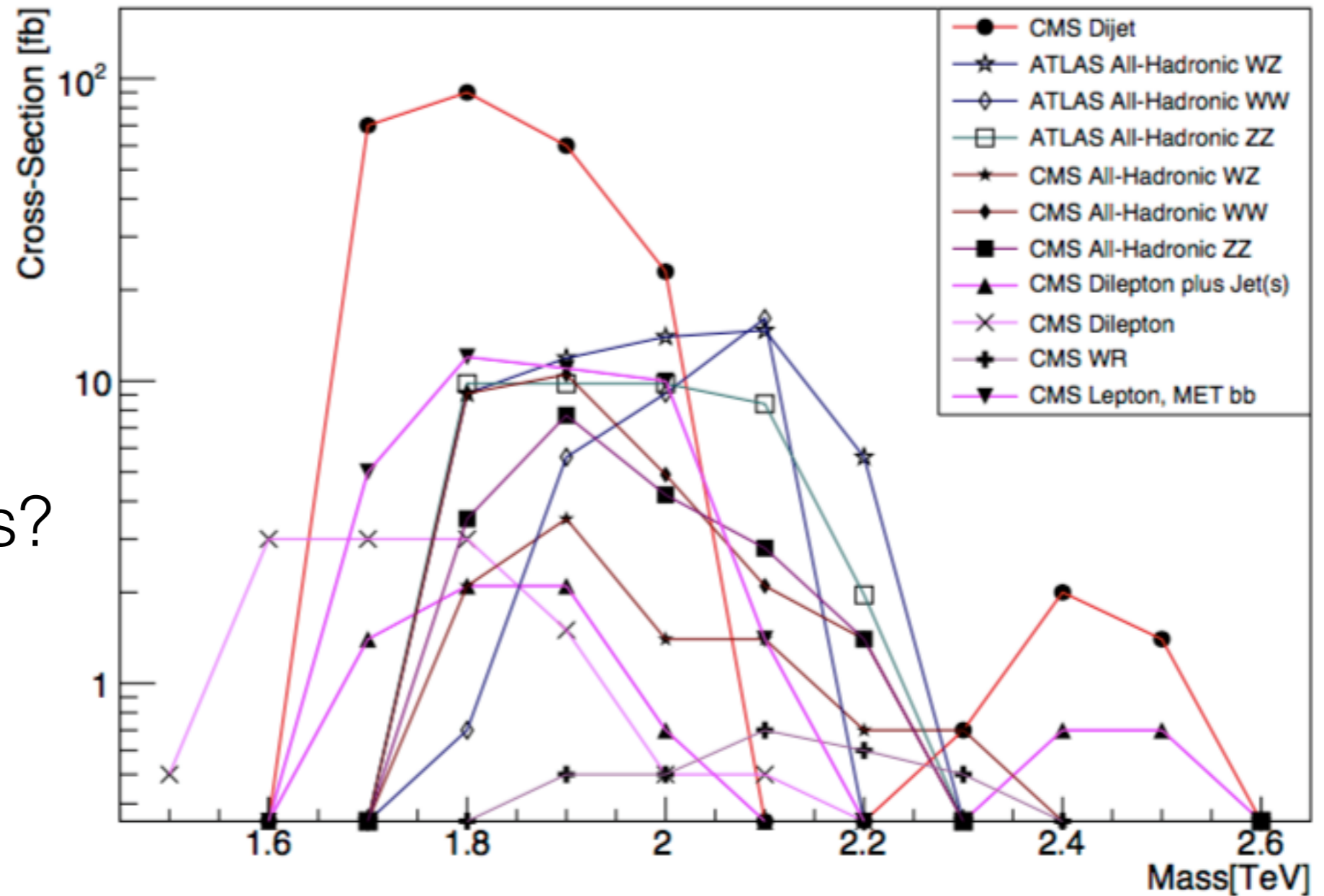
$$\sigma_R(pp \rightarrow x + y) = 16\pi^2 \cdot \mathcal{N} \cdot \frac{\Gamma_R}{m_R} \cdot (1 + \delta_{ij}) BR(R \rightarrow ij) \cdot BR(R \rightarrow xy) \left[\frac{1}{s} \frac{dL^{ij}}{d\tau} \right]_{\tau = \frac{m_R^2}{s}}$$

In the narrow-width approximation, for a given i & j , a bound on σ_R is a limit on

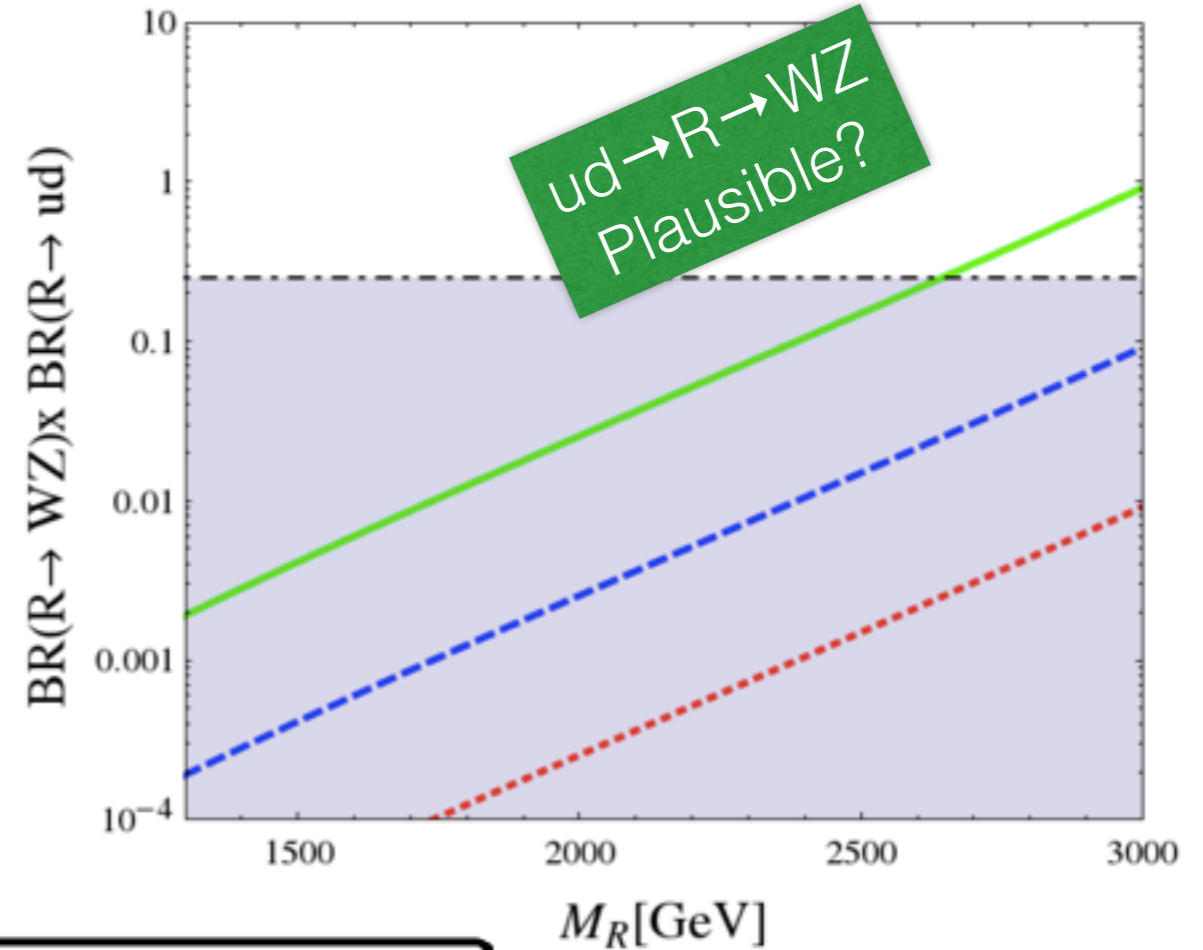
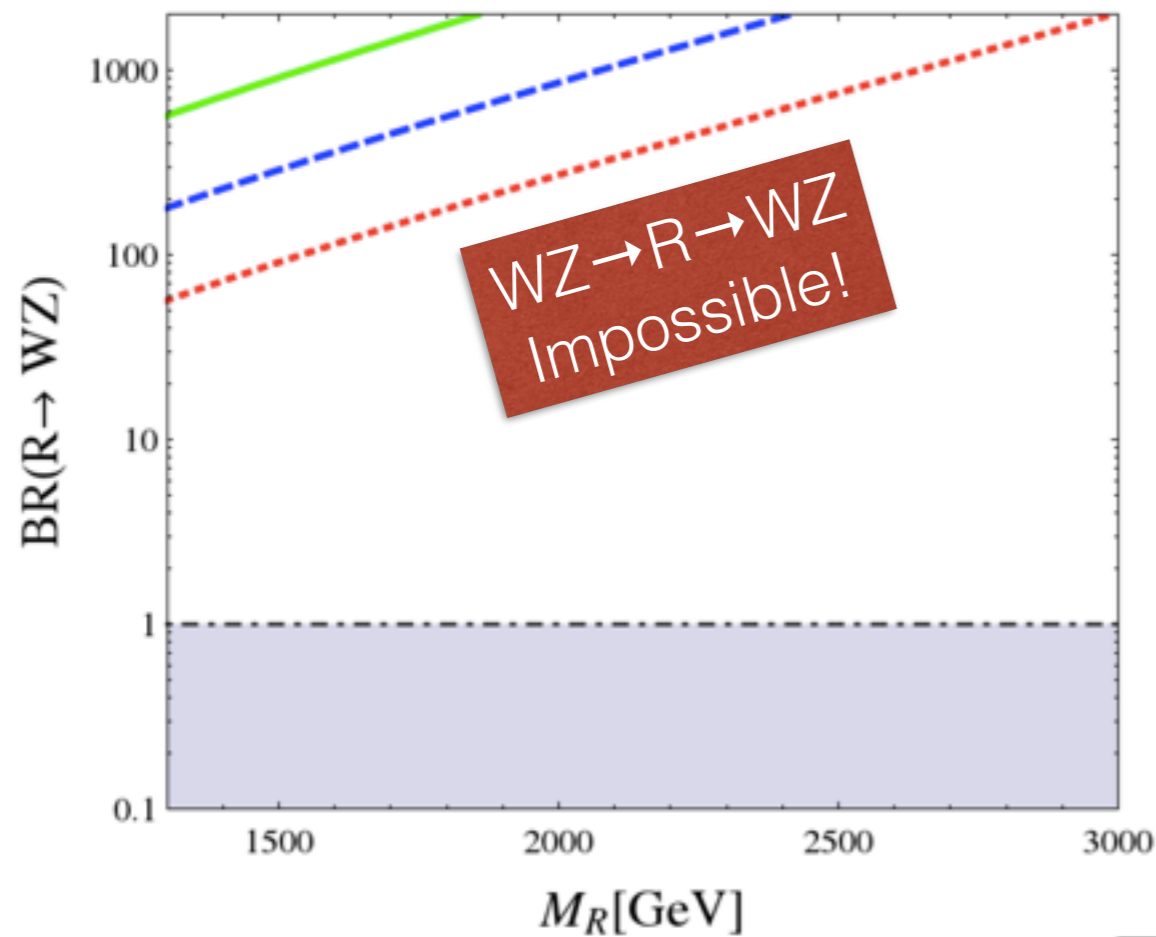
$$\mathcal{N} \cdot \frac{\Gamma_R}{m_R} \cdot BR(R \rightarrow i + j) \cdot BR(R \rightarrow x + y)$$

Diboson Excess

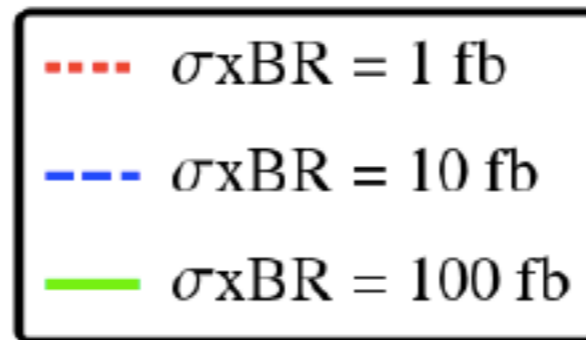
1-100 fb
“WZ” excess?



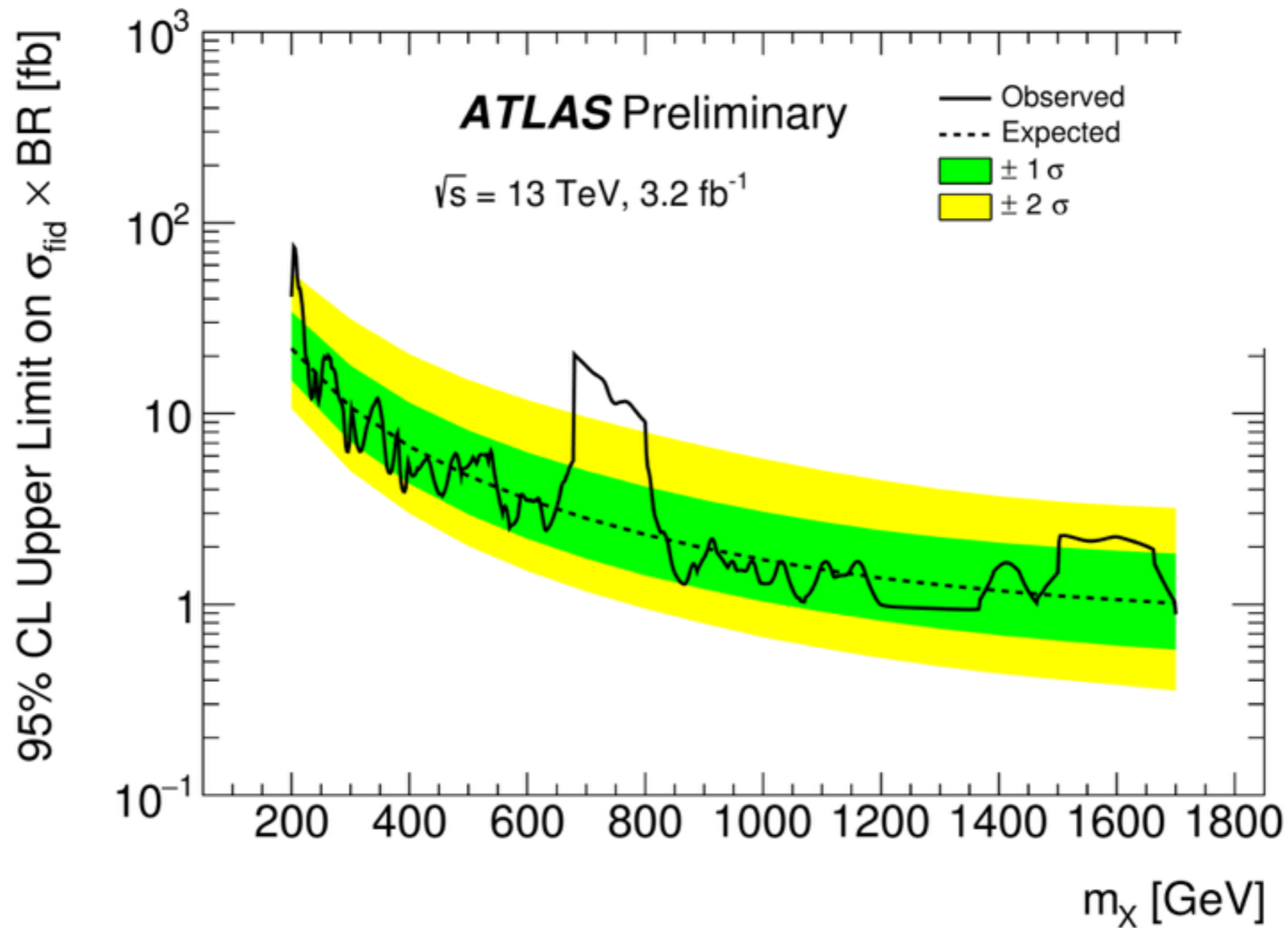
DiBoson Vector Resonances



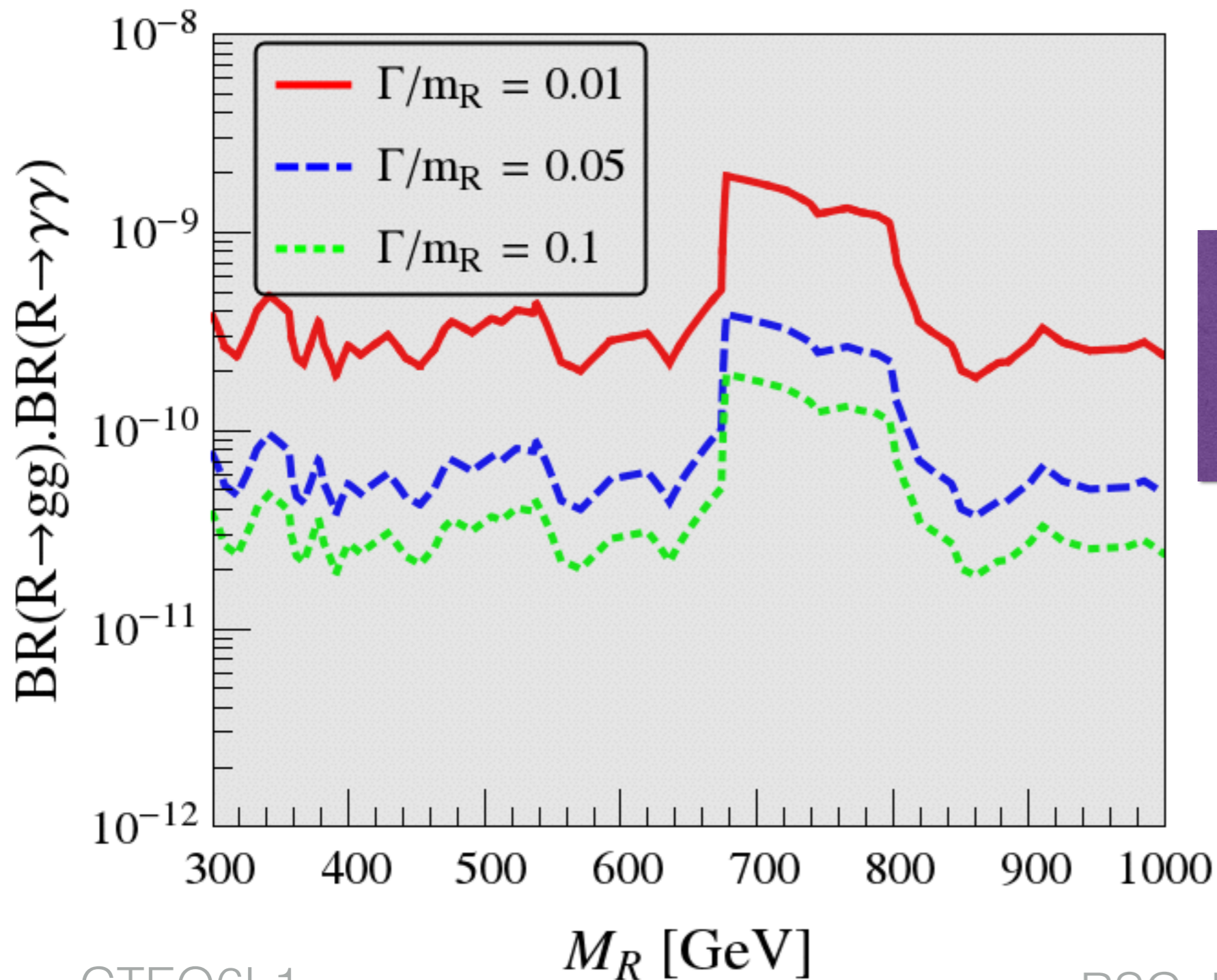
$$\frac{\Gamma_R}{M_R} = 0.1$$



Diphoton Resonance



Simplified Limits on Scalar Diphoton Resonances



$gg \rightarrow S \rightarrow \gamma\gamma$
 $BR(gg) \cdot BR(\gamma\gamma)$
for various Γ/M

Simplified Limits on Scalar Diphoton Resonances: *Diphoton Production*

$$\sigma_R(pp \rightarrow \gamma\gamma) = 32\pi^2 \cdot \left(\frac{1}{4}\right) \cdot \frac{\Gamma_R}{m_R} \cdot BR(R \rightarrow \gamma\gamma)^2 \cdot \left[\frac{1}{s} \frac{dL^{\gamma\gamma}}{d\tau} \right]_{\tau = \frac{m_R^2}{s}}$$

To be consistent with observed signal
a scalar *produced* through photon fusion
(since $BR(s \rightarrow \gamma\gamma) < 1$), $\Gamma/M > 3 \times 10^{-4}$

Using simplified models based on the narrow width approximation we can derive model-universal “simplified limits”:
model-independent limits on branching ratios for different production mechanisms.