Electromagnetic dissociation of relativistic heavy-ions

W.J. Llope and P. Braun-Munzinger & the BNL-AGS 814 collaboration

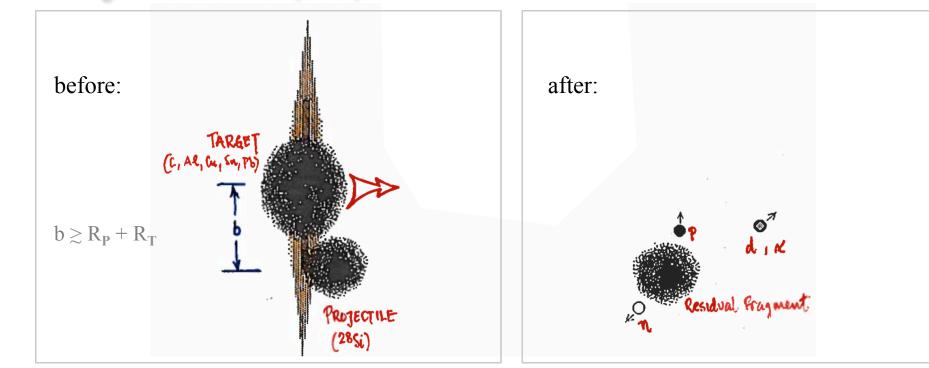
EMD in general Unknowns in ~1980s: Multiple excitations? Exotic decays? BNL-AGS Experiment 814 "folding model" Results from E814 Results from other experiments Summary



QCD thermodynamics – pressure and passion, PBM70, 24-26 Aug 2016, Schloβ Waldthausen, Mainz, Germany

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ElectroMagnetic Dissociation (EMD) as of the late-1980's...

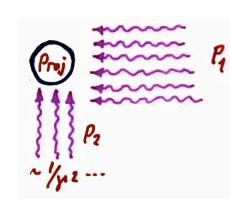


target nucleus is a strong source of (equivalent) photons *a la* Weizsäcker & Williams:

$$\mathcal{N}(E_{\gamma},b) = \frac{1}{E_{\gamma}} \left[\frac{Z_T^2 \alpha}{\pi^2 \beta_P^2 b^2} \right] \left[x^2 K_1^2(x) + \frac{x^2}{\gamma_P^2} K_0^2(x) \right]$$

where

$$x = \frac{E_{\gamma}b}{\gamma_P\beta_P\hbar c}$$



projectile absorbs photon(s) and then decays by some # of N or α emissions...

At the time, all EMD results consistent with single γ absorption then statistical decay...

C. Brechtmann *et al.*, PRC 39, 2222 (1989),
Z. Phys. A 330, 407 (1988), Z. Phys. A 331, 463 (1988)
J.C. Hill *et al.*, PRL 60, 999 (1988)
M.T. Mercier *et al.*, PRL 52, 989 (1984), PRC 33, 1655 (1986)
H.H. Heckman *et al.*, PRL 37, 56 (1976)
J. Barrette *et al.* (E814), PRC 41, 1512 (1990)



EMD open questions (mid-1980s)

EMD has been investigated experimentally with a variety of techniques. - using real photons - using heavy targets & radiochemical techniques
- foil detectors
- complete spectroscopy" Heavy tangets 2 relativistic light ions (160, 28 Si, 325, ...) allows the possibility of MULTIPLE EXCITATION. - resultant decay mechanisms are Not understood. Giant Dipole Resonance: dominates the γ absorption cross-section... p^(N) ~ 0.51 <u>Aproj</u> 2N+1 "Average Nproj Zproj Collective Coordinan Ph. Chomaz et al., Phys. Rep. 252 (1995) 275 for 28 Si: g(1)~ 0.6 pm , p(4)~ 1.2 pm Radius (281i) ~ 3.6 fm - an efficient production mechanism for final states of many neutrons at Zero Temperature?

Wayne StatE UNIVERSITY

BNL-AGS E814

Beam: C1 Status: In progress Approved Hours: 3600 Experiment 814 - Study of Extreme Peripheral Collisions and of the Transition from Peripheral to Central Collisions in Reactions Induced by Relativistic Heavy Ions

Brookhaven National Laboratory	M. Fatyga, R. Hogue, D. Lissauer, T. Ludlam, D. Makowiecki, E. O'Brien, V. Polychronakos, H. Takai, T. Throwe, C. Woody	
CERN	W.J. Willis	
Los Alamos National Laboratory	J. Boissevain, D. Fox, W. Sondheim, J. Sunier, H. van Hecke	
McGill University	J. Barrette, S.K. Mark, C. Pruneau	
University of New Mexico	B. Bassalleck, J. Hall, D. Wolfe	
University of Pittsburgh	W. Cleland, K. Jayananda, D. Kraus, U. Sonnadara, M. Takagui, X. Zhang	
State University of New York, Stony Brook	R. Beliwied, P. Braun-Munzinger, G. David, N. Herrmann, W. Llope, M. Muthuswami, J. Stachel, L. Waters, N. Xu, Y.C. Zhang	
University of Tel-Aviv	R. Heifetz	
Texas A&M University	J.A. Shoemaker, J. Simon, K. Wolf	
Yale University (Nuclear Physics) (High Energy Physics)	S.V. Greene, T. Hemmick, R. Majka, J. Mitcheli, F. Rotondo, J. Sandweiss, B. Shivakumar	
Spokesman:	P. Braun-Munzinger	

Experiment 814 combines 4π calorimeter coverage with a high resolution forward spectrometer. Extreme peripheral collisions are investigated by making an exclusive study of projectile fragmentation processes induced by Coulomb excitation. Central collision, including transverse energy and charged particle multiplicity production and a search for hydrodynamic flow effects are also studied. In the forward spectrometer, spectra for protons, neutrons and eventually, pions and kaons are measured to provide a detailed picture of the stopping process. A search is made for neutral particles, possibly produced in central collisions. The possible production of strange matter in such collisions is studied. The target area calorimetry provides a reaction trigger while fragments are identified in the forward spectrometer by their momentum/charge, Time-Of-Flight, specific energy loss and kinetic energy. This includes a measurement of anti-proton production and a search for anti-nucleus production in such collisions.



State of the art calculations: Bertulani & Baur

EMD cross-section: $\sigma = \int_{R}^{\infty} 2\pi b P(b) db = \int n(\omega) \sigma_{\gamma}(\omega) \frac{d\omega}{\omega}$

Approximation for the GDR and GQR:

$$\sigma_{\rm C} \simeq \frac{n_{\rm E1}[E_{\rm GR}^{(1)}]}{E_{\rm GR}^{(1)}} \int \sigma_{\gamma}^{\rm E1}(E_{\gamma}) \, \mathrm{d}E_{\gamma} + n_{\rm E2}[E_{\rm GR}^{(2)}] \, E_{\rm GR}^{(2)} \int \frac{\sigma_{\gamma}^{\rm E2}(E_{\gamma}) \, \mathrm{d}E_{\gamma}}{(E_{\gamma})^2}$$

Evaluate prefactors at parameterized E* for GDR,GQR

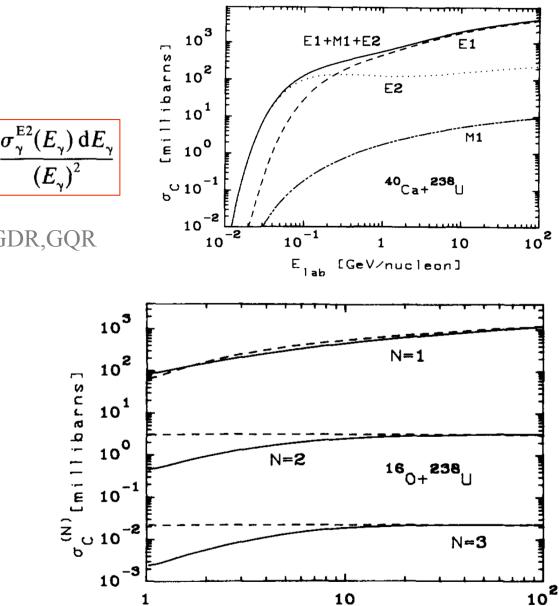
 $E_{GR}^{(1)} = 80/A^{1/3} \text{ MeV}$ <E*> for GDR $E_{GR}^{(2)} = 62/A^{1/3} \text{ MeV}$ <E*> for GQR

GDR and GQR strengths from sum rules: $\int \sigma_{\gamma}^{E1}(E_{\gamma}) dE_{\gamma} \simeq 60NZ/A \text{ MeV mb}$

 $\int \sigma_{\gamma}^{E2} (E_{\gamma}) / (E_{\gamma})^2 dE_{\gamma} \simeq 0.22 Z A^{2/3} \,\mu b \, \text{MeV}^{-1}$

Assumes the photoabsorption cross-section is a delta function at a given E*...

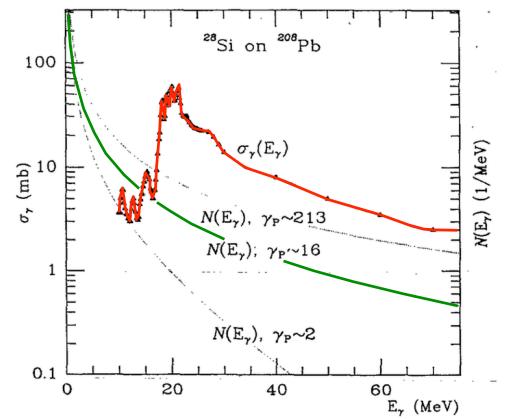
G. Baur and C. Bertulani, Phys. Rep. 163 (1988) 299.





(real) photon absorption cross-sections taken from experiment...

 $10 < E_{\gamma} < \hbar c \gamma_{proj} / b_{min}$ significant width due to quasi-deuteron...



eq. photons are real photons... straight line trajectories... all γ 's absorbed in ground state... numerical integration (requires care)... mean # of 's absorbed at impact parameter b:

$$m_{\gamma}(b) = \int dE_{\gamma} \mathcal{N}(E_{\gamma}, b) \sigma_{\gamma}(E_{\gamma}) = \int dE_{\gamma} \Phi(E_{\gamma}, b)$$

In harmonic oscillator approximation, probability of single γ absorption is Poissonian:

$$\mathcal{P}^{(1)}(b) = m_{\gamma}(b) e^{-m_{\gamma}(b)} \qquad (m_{\gamma} << 1)$$

Properly normalized probability that this γ has an energy E:

$$q^{(1)}(E_{\gamma},b) = \frac{\Phi(E_{\gamma},b)}{m_{\gamma}(b)}$$

Differential cross-section of absorption of exactly one γ of energy E (E*):

$$\frac{d\sigma^{(1)}}{dE^{*}} = \int db \ 2\pi b \mathcal{P}^{(1)}(b) q^{(1)}(E^{*},b)$$

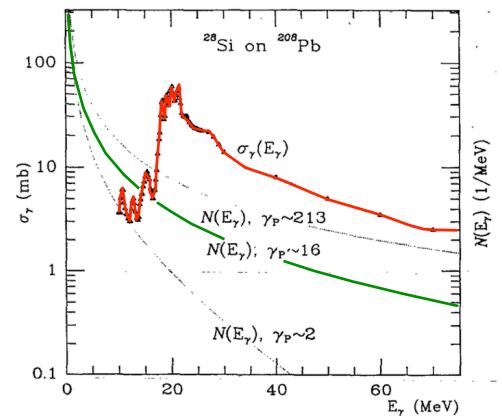
Assume the produced GDR mixes completely into compound nucleus before decaying. Use statistical model (CASCADE) for decay BRs...

single photon cross-sections scale like $\sim Z_T^2$



(real) photon absorption cross-sections taken from experiment...

 $10 < E_{\gamma} < \hbar c \gamma_{proj} / b_{min}$ significant width due to quasi-deuteron...



eq. photons are real photons... straight line trajectories... all γ 's absorbed in ground state... numerical integration (requires care)... mean # of 's absorbed at impact parameter b:

$$m_{\gamma}(b) = \int dE_{\gamma} \mathcal{N}(E_{\gamma}, b) \sigma_{\gamma}(E_{\gamma}) = \int dE_{\gamma} \Phi(E_{\gamma}, b)$$

In harmonic oscillator approximation, probability of two γ absorptions is Poissonian:

$$\mathcal{P}^{(2)}(b) = \frac{m_{\gamma}^{2}(b)}{2!}e^{-m_{\gamma}(b)}$$

Properly normalized probability that these γ 's have result in sum energy E (E*):

$$q^{(2)}(E^*,b) = \frac{\int dE_1 \Phi(E_1,b) \Phi(E^*-E_1,b)}{m_{\gamma}^2(b)}$$

Differential cross-section of absorption of two γ 's resulting in energy E (E*):

$$\frac{d\sigma^{(2)}}{dE^*} = \frac{1}{2!} \int db \ 2\pi b e^{-m_{\gamma}(b)} \\ \times \int dE_1 \Phi(E_1, b) \Phi(E^* - E_1, b)$$

Total cross-section is then the sum over N γ 's:

$$\sigma \coloneqq \int dE^* \left[\frac{d\sigma^{(1)}}{dE^*} + \frac{d\sigma^{(2)}}{dE^*} + \frac{d\sigma^{(3)}}{dE^*} + \cdots \right]$$

n-photon cross-sections scale like $\sim Z_T^{2n}$



QCD thermodynamics – pressure and passion, PBM70, 24-26 Aug 2016, Schloβ Waldthausen, Mainz, Germany

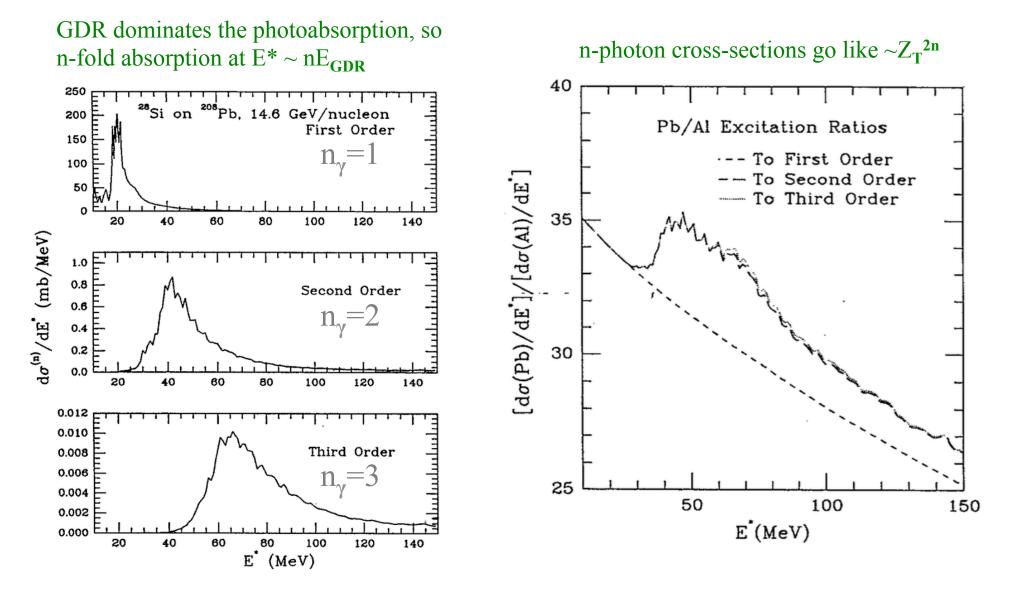
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Importance of realistic photoabsorption cross-section

	Bertulani & Baur	WJL&PBM narrow σ_{γ}	WJL&PBM real σ _γ	
Order	σ (Ref. 20)	$\sigma(\sigma_{\gamma}^{G})$	$\sigma(\sigma_{\gamma}^{expt})$	
	¹⁶ O on ²³⁸ U at	100 GeV/nucleon		
1	N/A	1.03 b	2.5 b	→×2.5
2	3.1 mb	3.16 mb	20.6 mb	→ ×6.5
3	22 µb	22.9 µb	328 <i>µ</i> b	→ ×14
4	160 nb	164.3 nb	4.8 μb	→×30
	³² S on ²³⁸ U at	100 GeV/nucleon		
1	N/A	2.63 b	4.3 b	→×1.7
2	17 mb	17.5 mb	46.3 mb	→×2.6
3	250 µb	273 µb	1.1 mb	→ ×4
4	4 μb	4.3 μb	<u>26.6 μb</u>	→×6.2

calculated cross-sections for EMD are significantly larger when using a realistic photoabsorption cross-section!





study ratios of cross-sections for high- and low-Z targets for decay channels near the peak of the n-fold photoabsorption cross-section! look for enhancements in decay channels w/ Q-values near n*E_{GDR}...



Signatures of multiple photon absorption



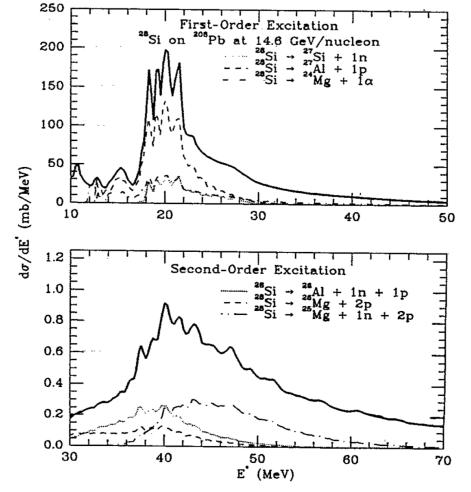


FIG. 8. The differential first-order (top plot, solid line) and second-order (bottom plot, solid line) excitation cross sections versus the excitation energy. The dashed lines indicate the contributions from the important decay channels, labeled in the appropriate figure.

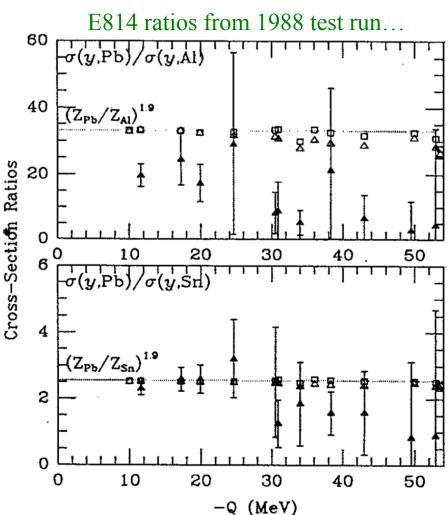
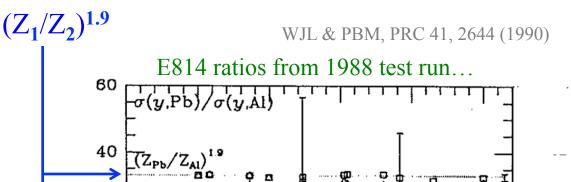


FIG. 10. The ratios of the E814 cross sections, Ref. 10, for the electromagnetic dissociation of 14.6 GeV/nucleon ²⁸Si projectiles (solid points), the calculated ratios to first order (open triangles) and to third order (open squares). The dotted line gives the ratio expected for first-order excitation.

E814 1988 data: significant nuclear contamination... Error bars need to be *much* smaller...



Signatures of multiple photon absorption



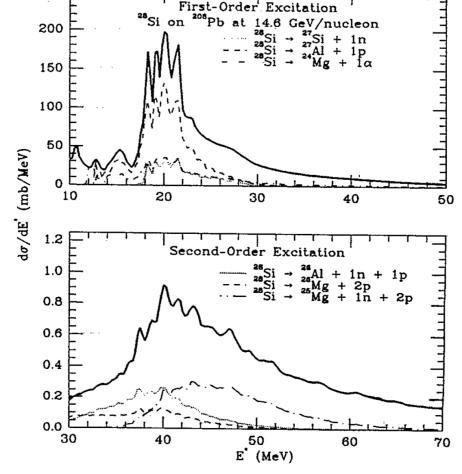


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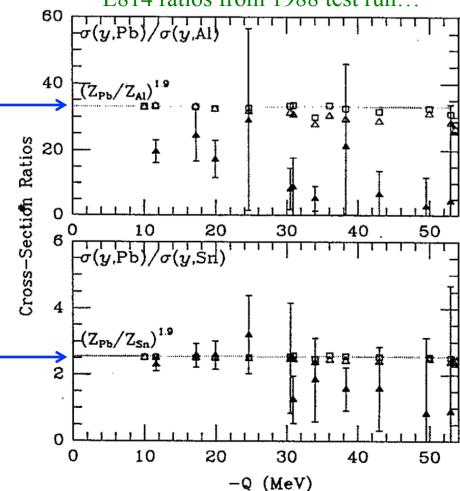


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E814 1988 data: significant nuclear contamination... Error bars need to be *much* smaller...



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Goals of the E814 "extreme peripheral collisions" program

Accurate separation between EM and nuclear excitations in experimental data taken by the E814 experiment for ²⁸Si on C, Al, Cu, Sn, Pb targets at 14.6 GeV/c/N use many detectors to reject nuclear interactions and redundantly measure all FS particles

Compare measured EMD cross-sections with those from a semi-classical Weizsäcker-Williams framework which incorporates the statistical model for the nuclear decay to look for discrepancies possibly related to exotic non-statistical decay mechanisms

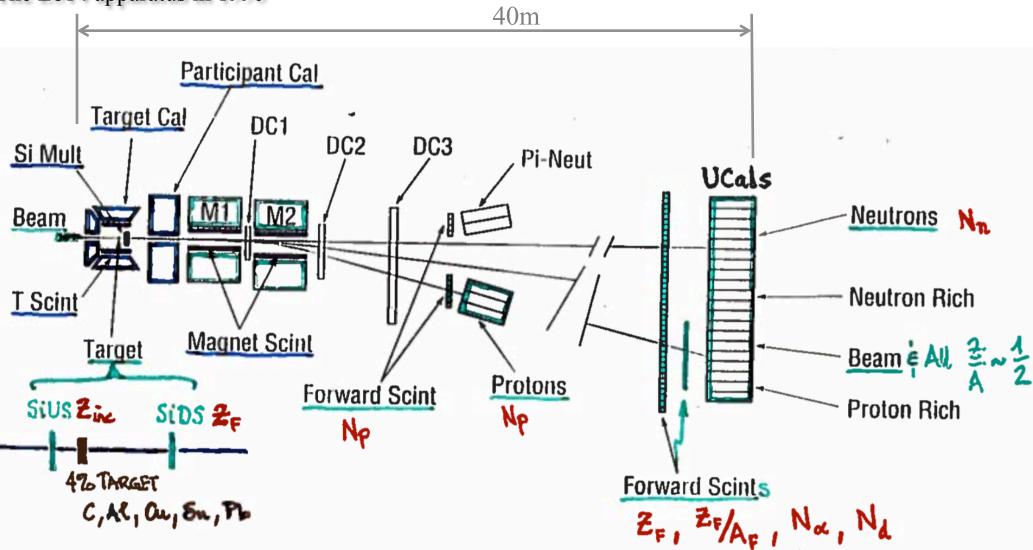
Establish with precision the scaling of the cross-sections for EM and nuclear excitations with the target species

nuclear cross-sections $\sigma \sim (A_P^{\frac{1}{3}} + A_T^{\frac{1}{3}})$ EMD cross-sections $\sigma_n \sim Z_T^{1.9n}$ $n = \# \text{ of } \gamma \text{ 's absorbed}$

Establish or deny that multiple electromagnetic excitations contribute to the measured cross-sections study how the EMD cross-sections scale with the target charge Z_T



The E814 apparatus in 1990



" 4π calorimetry" hermetically surrounding the target... separate EMD from peripheral nuclear reactions

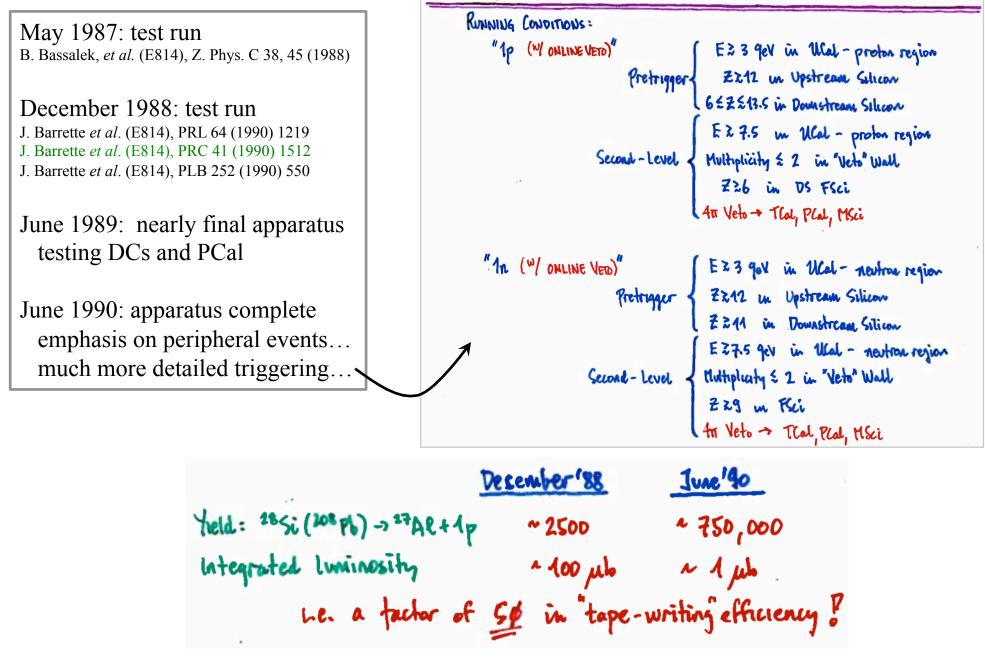
Forward magnetic spectrometer and calorimetry... (redundantly) exclusive measurements of the final state



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Datasets

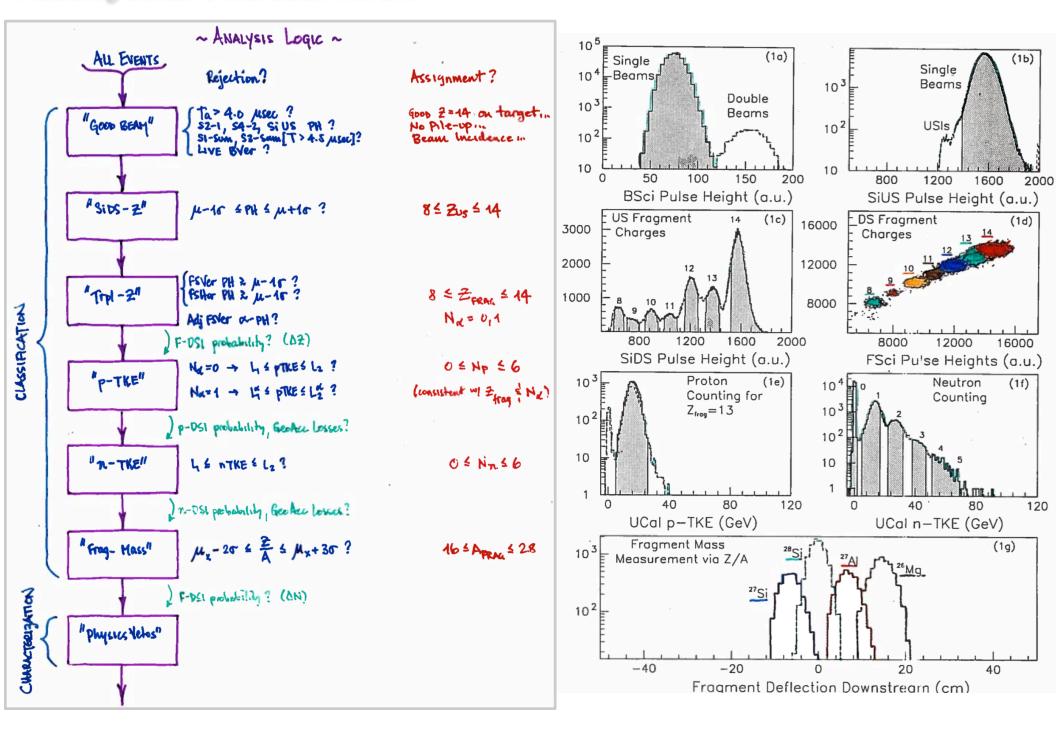
Unit Z (or A) resolution on all nucleons and nuclei in the final state required at the trigger level ... Otherwise E814 only measures down-stream interactions...





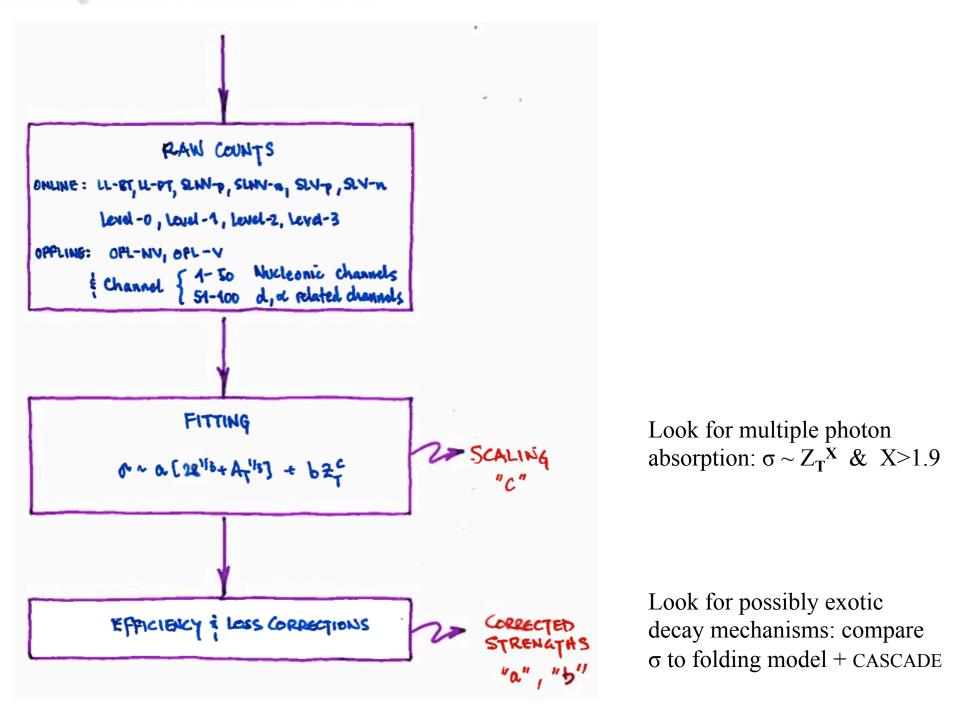
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Measuring exclusive final states with E814





Measuring exclusive final states with E814

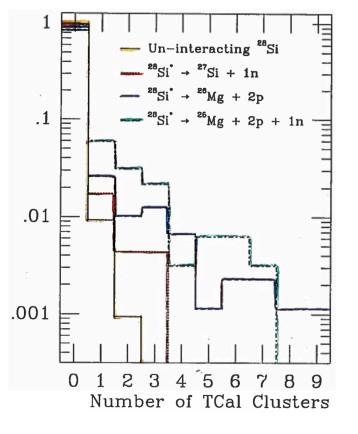


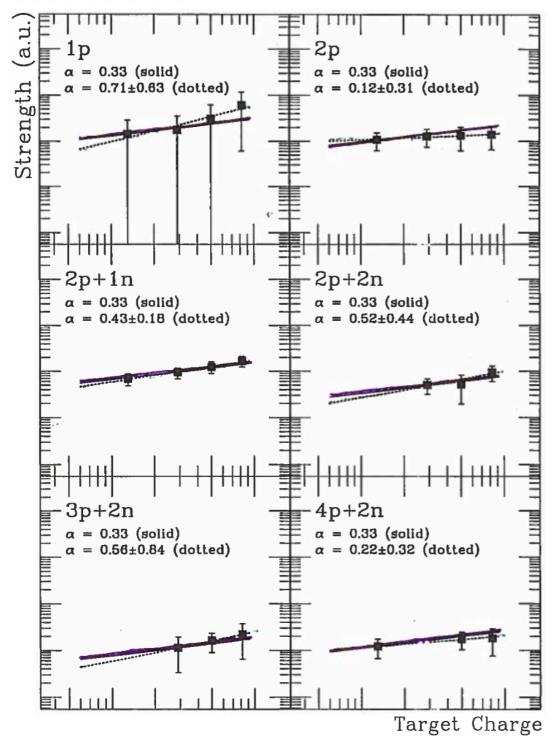


Making sure we understood the obviously nuclear excitations

for nuclear excitations, expect: $\sigma \sim (28^{\frac{1}{3}} + A_{T}^{\frac{1}{3}})$

select exclusive events accompanied by activity in target calorimeter

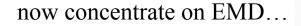






EMD cross-sections

Experimental Data (solid), WW/CASCADE Calculation (red)



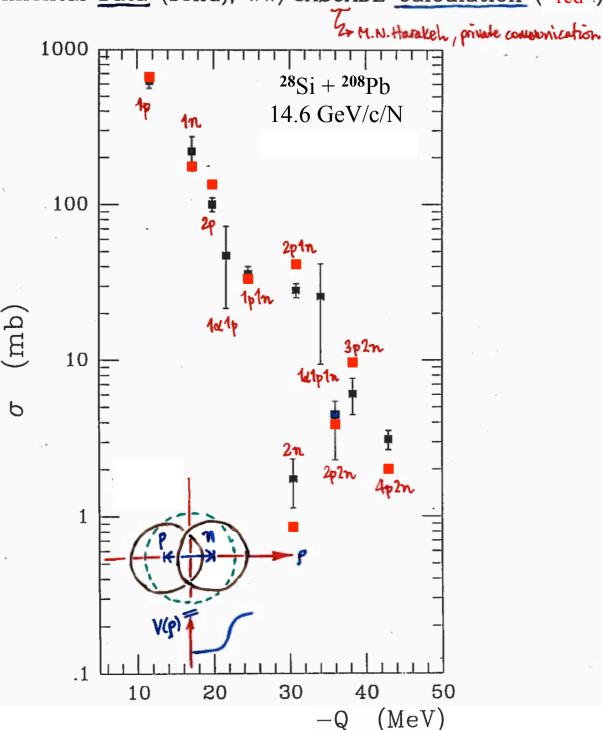
Assume nGDR mixes completely into the compound nucleus before decaying

Use a statistical model (CASCADE) to give the decay branching fractions for a given final state...

F. Pühlhofer, Nucl. Phys. A280, 267 (1977) b

tends to overestimate p-channels and underestimate n-channels.

some factor ~2 discrepancies but nothing considered "exotic"





EMD cross-sections

now concentrate on EMD...

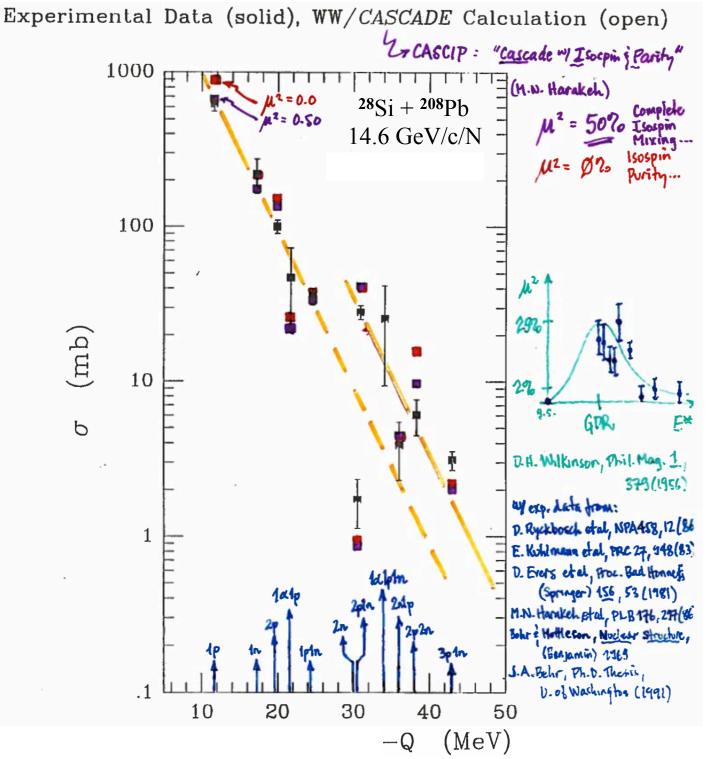
Assume nGDR mixes compl into the compound nucleus b decaying

Use a statistical model (CASCADE) to give the decay branching fractions for a give final state...

F. Pühlhofer, Nucl. Phys. A280, 26

tends to overestimate p-chan and underestimate n-chan

Also tried a version of CASC₄ including isospin & parity (C from M. Harrakeh... ...generally small chang





Z_T scaling

nuclear excitations still contribute even after careful cuts in the target region... "quiet nuclear component"

Rejected what I could, then fit cross-sections $vs. Z_T$ with:

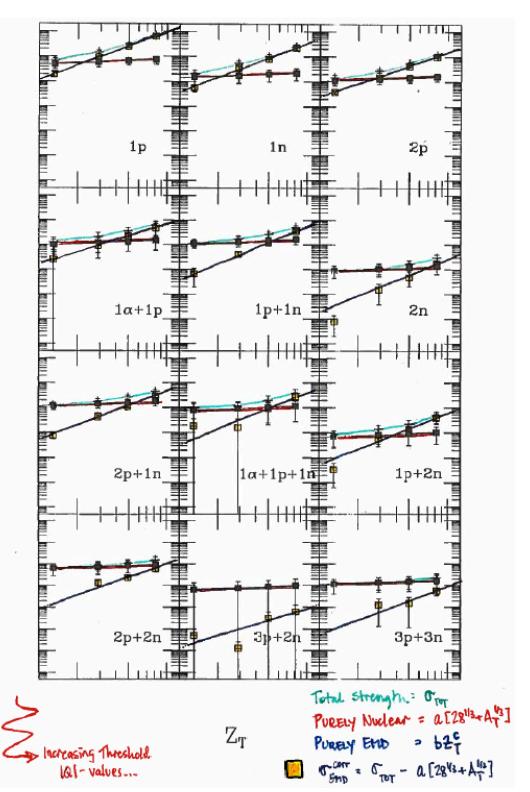
 $\sigma \sim a(28^{\frac{1}{3}}+A_T^{\frac{1}{3}})+bZ_T^{c}$ a,b: cross-sections c: scaling

For single photon absorption: $c \sim 1.9$ due to different ranges of impact parameters sampled for different targets

WJL & PBM, PRC 41, 2644 (1990)

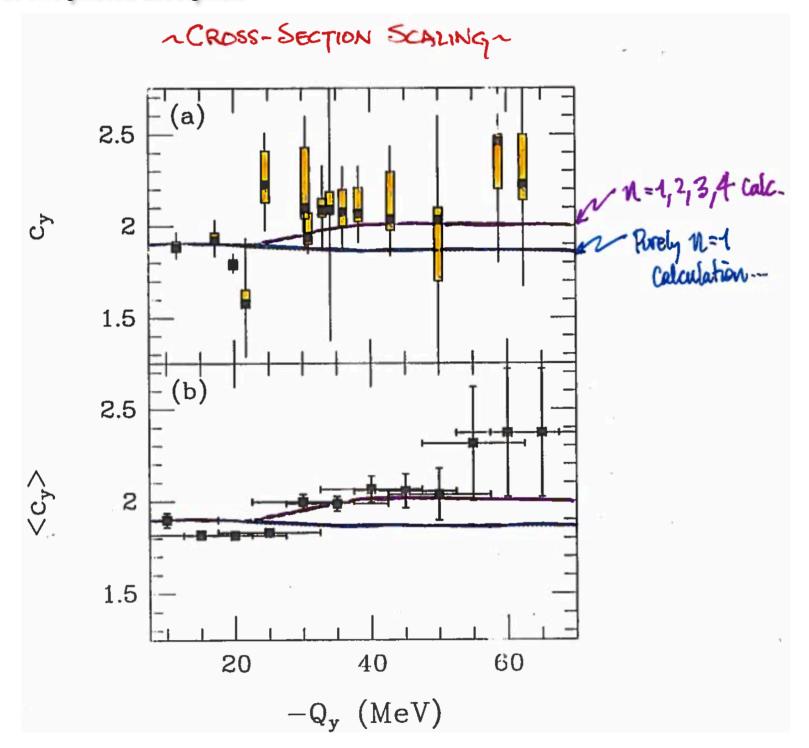
(mb)

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Observation of two photon absorption





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Lower beam energies should make the signal of 2-photon absorption easier to see...

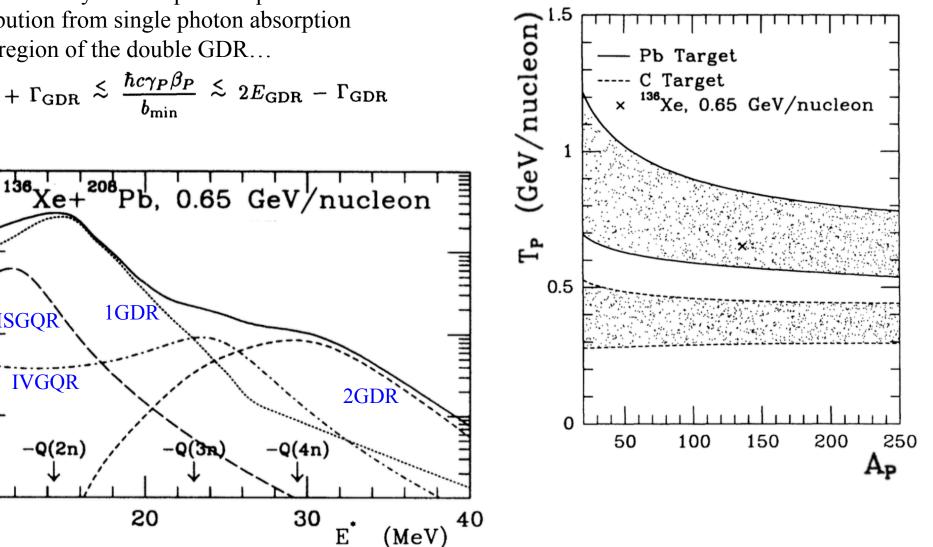
Use adiabaticity in WW photon spectrum to limit contribution from single photon absorption in the region of the double GDR...

 $E_{\rm GDR} + \Gamma_{\rm GDR} \lesssim \frac{\hbar c \gamma_P \beta_P}{b_{\rm min}} \lesssim 2 E_{\rm GDR} - \Gamma_{\rm GDR}$

1GDR

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Beam energies near 0.6-1 GeV/N can emphasize 2-photon w.r.t. 1-photon...



1000

100

10

0.1

10

ISGQR

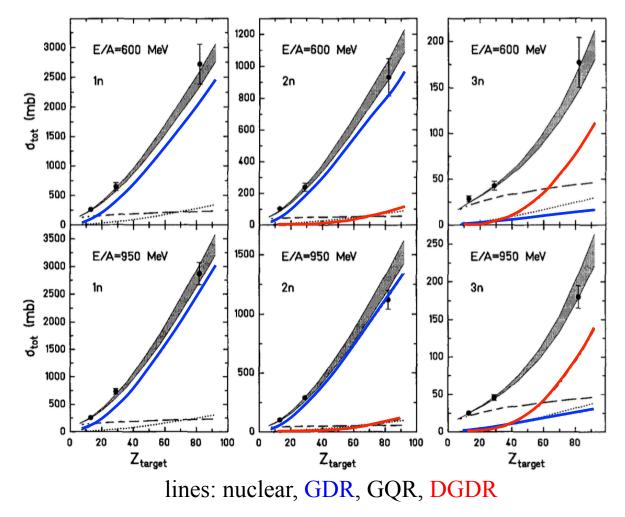
IVGQR

-Q(2n)

do/dE (mb/MeV)

Experimental results at lower beam energies

¹⁹⁷Au+X, T. Aumann *et al.*, Nucl. Phys. A569, 157c (1994)
²³⁸U+X, T. Aumann *et al.*, Nucl. Phys. A599, 329c (1996)
see also K. Boretzky *et al.* (LAND), PLB 384, 30 (1996)



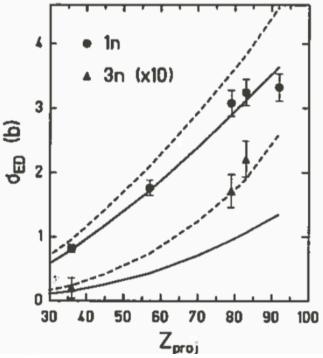


Figure 4. Calculated ED cross sections using different parametrizations for the minimum impact parameter b_{min} . The solid curves are obtained with the Kox parametrization [11], the dashed one with $b_{min} = 1.2(A_{proj}^{1/3} + A_{targ}^{1/3})$ fm.

It is obvious that it is not possible to describe simultaneously the 1n and the 3n cross sections by varying the minimum impact parameter. A better approach would be to replace the fixed value of b_{min} by a smooth transition as suggested by Llope and Braun-Munzinger [4].

at low beam energies, trickiest part is getting b_{min} right...



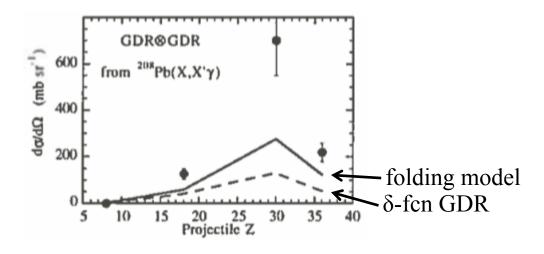
Now even lower! $E/A \le 95 \text{ MeV/N}$

Data taken at GANIL with high resolution magnetic spectrograph "SPEG"

Forward magnetic spectrometer and BaF_2 crystals around the target to measure γ 's

Measure dexcitation of nGDR via γ emission and E* in magnetic spectrometer.

The solid lines on the Fig. 7 are calculations for the excitation cross sections using the <u>simplest possible model</u> of the DGDR [24], which considers the DGDR strength distribution to be determined completely by the single photon E1 photo absorption cross section; additional damping and other broadening effects are ignored.



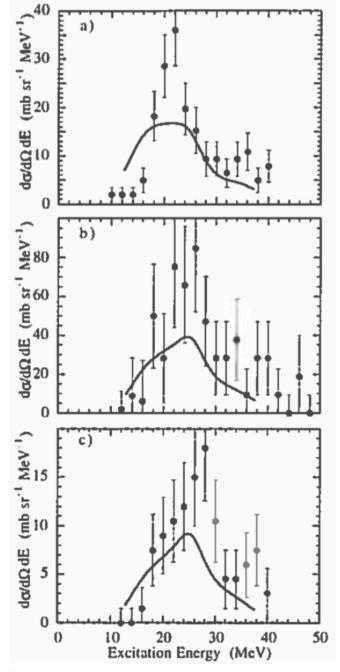


Figure 7. Experimental (points) and theoretical (curves) cross sections for DGDR excitation in target nuclei: a) 60*A MeV ⁸⁶Kr on ²⁰⁸Pb, b) 80*A MeV ⁶⁴Zn on ²⁰⁹Bi, c) 95*A MeV ³⁶Ar on ²⁰⁸Pb.

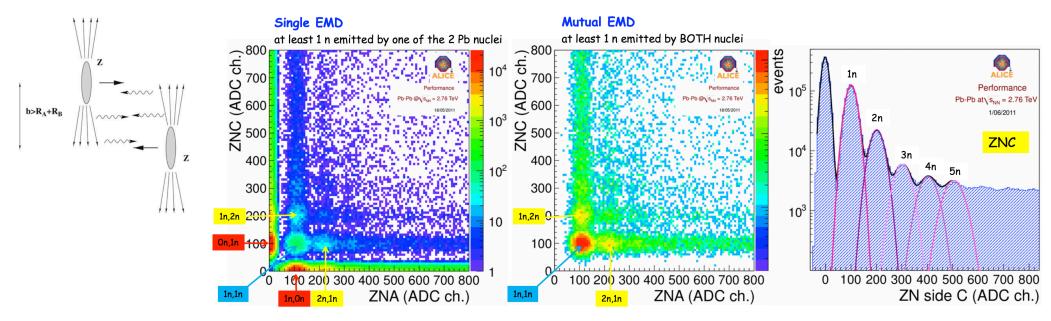


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Modern EMD

N. DeMarco, PLHC2011, Perugia ALICE, PRL 109, 252302 (2012)

EMD at RHIC and LHC is an important contibution to beam lifetime.



Compare to RELDIS model

I.A. Pshenichov et al., Phys. Rev. C 60 044901 (1999)

single and double virtual photon absorption by nuclei "following Llope and Braun-Munzinger" intranuclear cascades of produced hadrons

statistical decay of excited residual nuclei (evaporation, fission and multifragmentation)

	DATA (PRELIMINARY)	RELDIS MODEL
σ^{sEMD} + σ^{had}	(195.6 ± 0.1 stat. ^{+24.2} _{-11.7} syst.) b	(192.9 \pm 9.2 syst.) b
σ^{sEMD} - σ^{mEMD}	(176.9 \pm 0.1 stat. $^{+21.6}_{-10.6}$ syst.) b	$(179.7 \pm 9.2 \text{ syst.}) \text{ b}$
σ^{mEMD}	$(5.7 \pm 0.2 \text{ stat.} \ ^{+0.7}_{-0.3} \text{ syst.}) \text{ b}$	(5.5 \pm 0.6 syst.) b
σ^{sEMD}	(185.7 ± 0.2 stat. +22.6 syst.) b -11.1	(185.2 \pm 9.2 syst.) b



Summary

Peter and I developed a "folding model" for the cross-sections for n-photon EMD, $n\geq 1$ one must use realistic photoabsorption cross-sections in EMD σ calculations

With the E814 collaboration, we looked for any hints of multiple γ absorption look closely at decay channels with Q-values near n*E_{GDR} look closely at the scaling of the cross-sections with the target charge Z_T

Saw a hint of 2-photon EMD in 14.6 GeV/c/N ²⁸Si + A (but did not claim "observation") an independent analysis was consistent, but the scaling enhancement was not quite as strong

We suggested lower beam energies as the adiabatic cutoff reduces the 1γ near the DGDR.

Unambiguous double GDR photoexcitation seen in experiments at GSI and later elsewhere...

At higher beam energies, all data consistent with the folding model and statistical decay no "exotic" decays seen

At lower beam energies, the results depend strongly on the choice of b_{min} , and the total cross-sections are underestimated (many assumptions we used are breaking down).



Summary

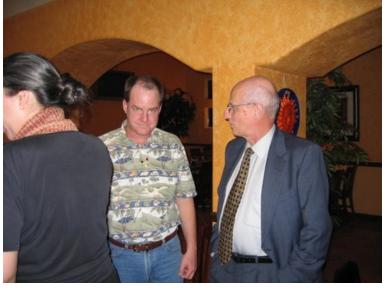
Birthday party for Jimmy Dee (~2007)

Just one night in a quite average Mexican restaurant northwest of Houston TX.

Peter and Johanna flew all the way around the world for this.

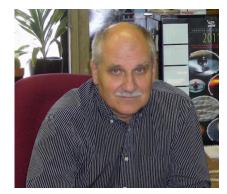








Greetings and best wishes from the Wayne State heavy-ion group!



Sergei Voloshin ALICE & STAR Flow, CME, vorticity Λ polarization



Jörn Putschke ALICE & STAR Jets in HI Jetscape topical group From all of us at WSU: All the best & many thanks for everything Peter!

> +6 postdocs (2 at CERN) +5 graduate students



Claude Pruneau ALICE Flow

Correlations



W.J. Llope ALICE & STAR GEM framing for ALICE TPCU rapidity correlations in RHIC BES



BACKUP



Giant Resonances

M.N. Harakeh, Joliet-Curie School, La Colle Sur Loup, France; 12-17 September 2011

