

MSSM dark matter and the muon $g-2$

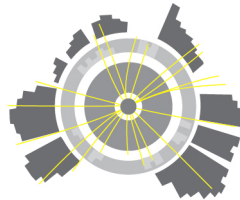
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November 29th 2016



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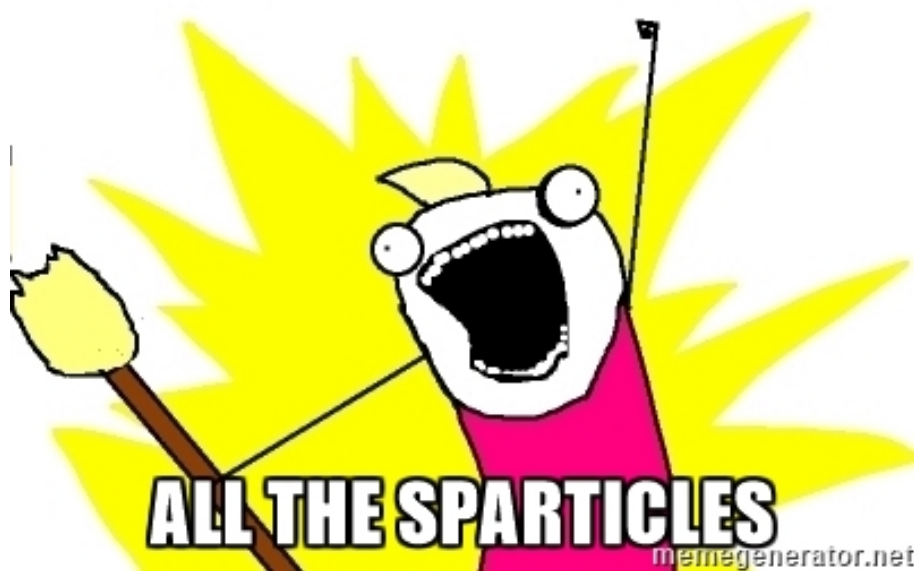
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Particle Physics at the Terascale

Outline

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- 2 The muon $g-2$ in SUSY
- 3 MSSM Parameter Scan
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Introduction

The LHC has resumed its search for Supersymmetry (SUSY) in run 2, but direct evidence is still yet to be found. The search continues on!

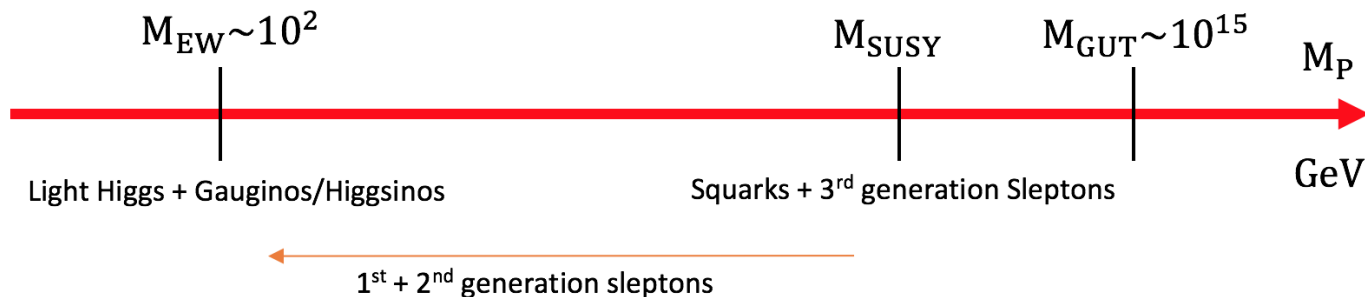


The observation of a Higgs at 125 GeV has strengthened the need for SUSY to appear at the weak-scale.

- Tree-level higgs mass $\sim m_Z$.
- Existence of electroweakinos (partners of EW gauge bosons).
- A light neutralino - great for DM!
- Composition of DM is important (Wino, Bino, Higgsino).

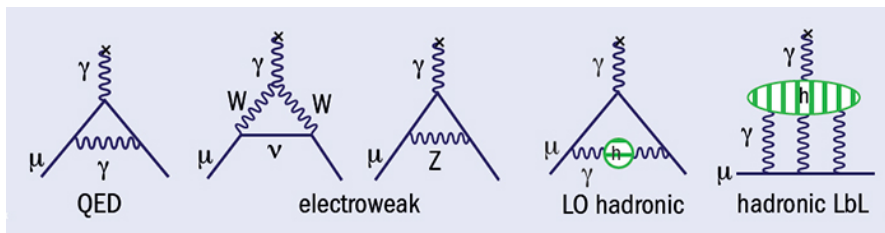
Minimal SUSY mass hierarchy

- Universal squark and 3rd gen slepton masses decoupled
- 125 GeV higgs finely-tuned
- Gauginos/higgsinos at weak scale, protected by chiral symmetry
- Light 1st and 2nd generation sleptons allowed by FCNC constraints
→ **muon g-2**



The muon $g - 2$

Contributions to the SM:



- Main theoretical uncertainty comes from LO Hadronic loop contributions (quarks and gluons)

$$20.6 \times 10^{-10} < \Delta a_\mu < 36.6 \times 10^{-10} \quad (1\sigma)$$

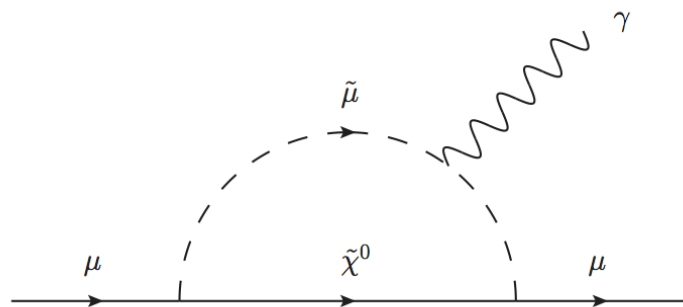
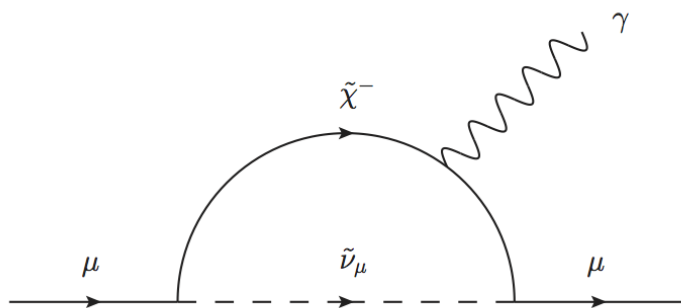
$$12.6 \times 10^{-10} < \Delta a_\mu < 44.6 \times 10^{-10} \quad (2\sigma)$$

where

$$\Delta a_\mu \equiv a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$$

The muon $g - 2$ in SUSY

Contributions come from sneutrino-chargino and smuon-neutralino loop diagrams



The muon $g - 2$ in SUSY

Contribution from the MSSM:

$$\Delta a_\mu = \frac{\alpha m_\mu^2 \mu \tan(\beta)}{4\pi} \left[\frac{M_2}{\sin^2 \theta_W m_{\tilde{\mu}_L}^2} \left(\frac{f_\chi(M_2^2/m_{\tilde{\mu}_L}^2) - f_\chi(\mu^2/m_{\tilde{\mu}_L}^2)}{M_2^2 - \mu^2} \right) \right. \\ \left. + \frac{M_1}{\cos^2 \theta_W (m_{\tilde{\mu}_R}^2 - m_{\tilde{\mu}_L}^2)} \left(\frac{f_N(M_1^2/m_{\tilde{\mu}_R}^2)}{m_{\tilde{\mu}_R}^2} - \frac{f_N(M_1^2/m_{\tilde{\mu}_L}^2)}{m_{\tilde{\mu}_L}^2} \right) \right]$$

f_χ and f_N are loop functions:

$$f_\chi(x) = \frac{x^2 - 4x + 3 + 2 \ln(x)}{(1-x)^3}, \quad f_\chi(1) = -2/3$$

$$f_N(x) = \frac{x^2 - 1 - 2x \ln(x)}{(1-x)^3}, \quad f_N(1) = -1/3$$

Explaining the muon $g - 2$ in the MSSM

- The following particles are important in analyzing the $(g - 2)_\mu$ in the MSSM:

$$\tilde{\mu}, \tilde{\nu}_\mu, \tilde{\chi}^0, \tilde{\chi}^\pm \quad (1)$$

- We can heavily constrain the muon $g - 2$ through slepton and chargino searches at colliders
- Smuons should be kept light (less than around 500 GeV) to increase contribution to the $(g - 2)_\mu$
- Dark Matter (Direct/Indirect) searches can constrain neutralino LSPs in R-Parity conserving SUSY

Constraints from Experiment

- LEP constraints on chargino and Slepton masses:

$$\begin{aligned} m_{\tilde{l}_L}, m_{\tilde{l}_R} &> 100 \text{ GeV} \quad (l = e, \mu) \\ m_{\tilde{\chi}_1^\pm} &> 105 \text{ GeV} \end{aligned}$$

- Constraints on neutralino LSP as a DM candidate:

$$m_{\tilde{\chi}_1^0} > 30 \text{ GeV}$$

- Higgs mass from ATLAS/CMS:

$$123 < m_{h^0} < 127 \text{ GeV}$$

- Higgs precision constraints (LEP, Tevatron and LHC)
- Dark matter relic density (PLANCK 2013)

$$\Omega h^2 = 0.112 \pm 0.006 \quad (1\sigma)$$

- WIMP-nucleon Spin-Independent Cross Section (LUX 2016)

MSSM Parameter Scan

We calculate the $(g - 2)_\mu$ and mass spectrum in the MSSM using FeynHiggs-1.12.0:

- Decoupled Squarks at 5 TeV (Ignore B -Physics constraints)
- Stau sleptons $m_{\tilde{\tau}_L} = m_{\tilde{\tau}_R} = 5 \text{ TeV}$
- Gluino mass $M_3 \sim 3 \text{ TeV}$
- Trilinear coupling A_t in range $|A_t| < 5 \text{ TeV}$ (We keep $|X_t/M_S| < 2$ to avoid charge/colour-breaking minima)
- Rest of higgs sector decoupled by setting $m_{A^0} = 2 \text{ TeV}$

Parameter scan range:

$$10 < \tan(\beta) < 50,$$

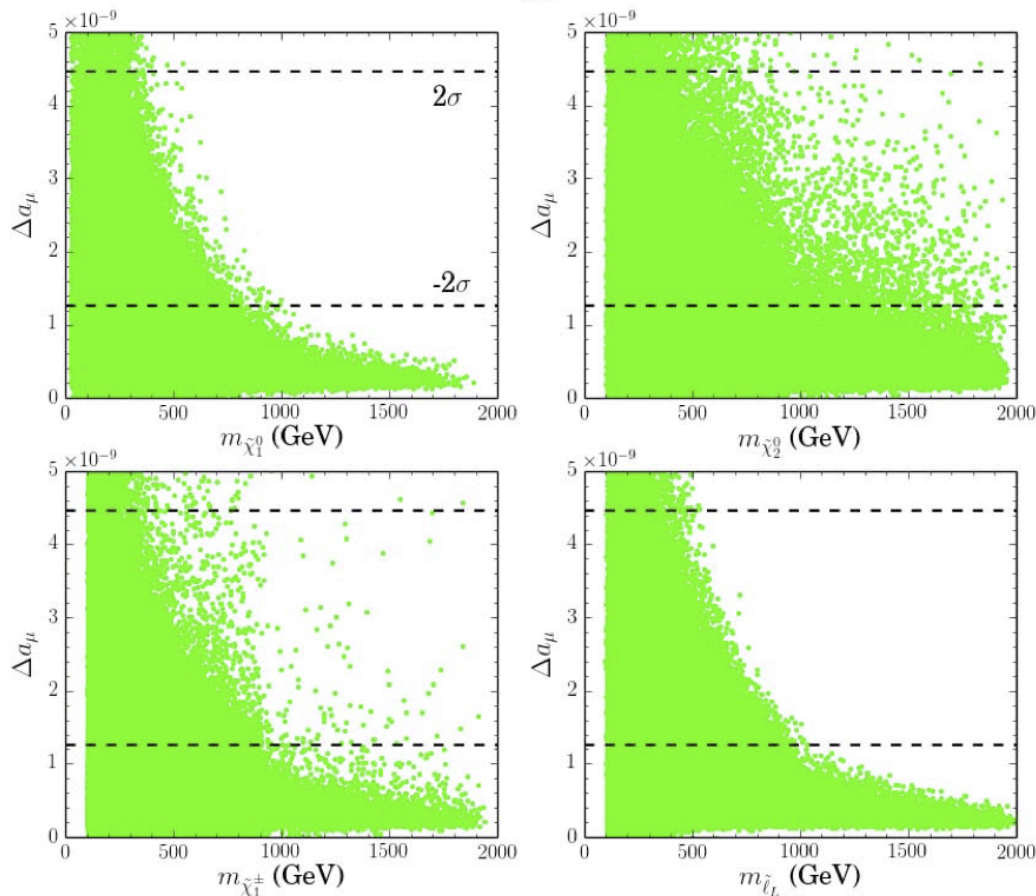
$$|M_1|, |M_2|, |\mu| < 2 \text{ TeV},$$

$$0.1 < m_{\tilde{l}_L}, m_{\tilde{l}_R} < 2 \text{ TeV}, \quad (l = e, \mu)$$

Higgs mass calculated in FeynHiggs, precision constraints in HiggsBounds-4.2.1. SUSY spectrum calculated in SPheno, MicrOmegas to calculate DM relic density and SI WIMP-nucleon CS.

Limits on neutralinos, charginos and smuons

LEP+Higgs data



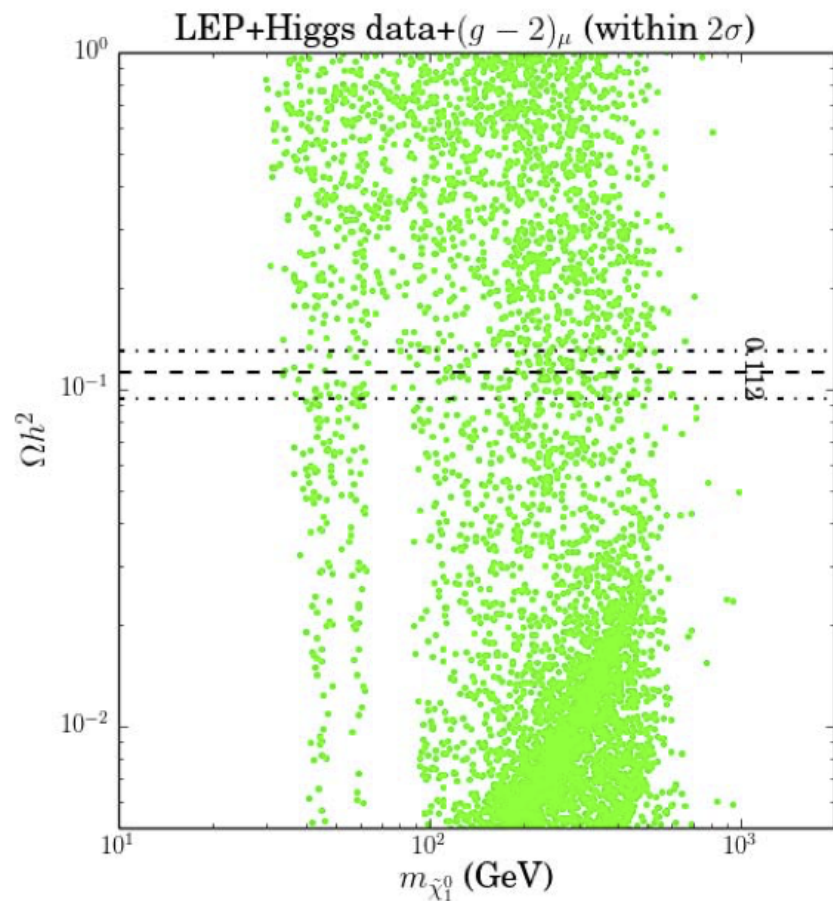
(N)LSP component

Parameters	LSP	NLSP
$M_1 > M_2 > \mu$	Higgsino	Wino
$M_1 > \mu > M_2$	Wino	Higgsino
$M_2 > \mu > M_1$	Bino	Higgsino
$\mu > M_2 > M_1$	Bino	Wino

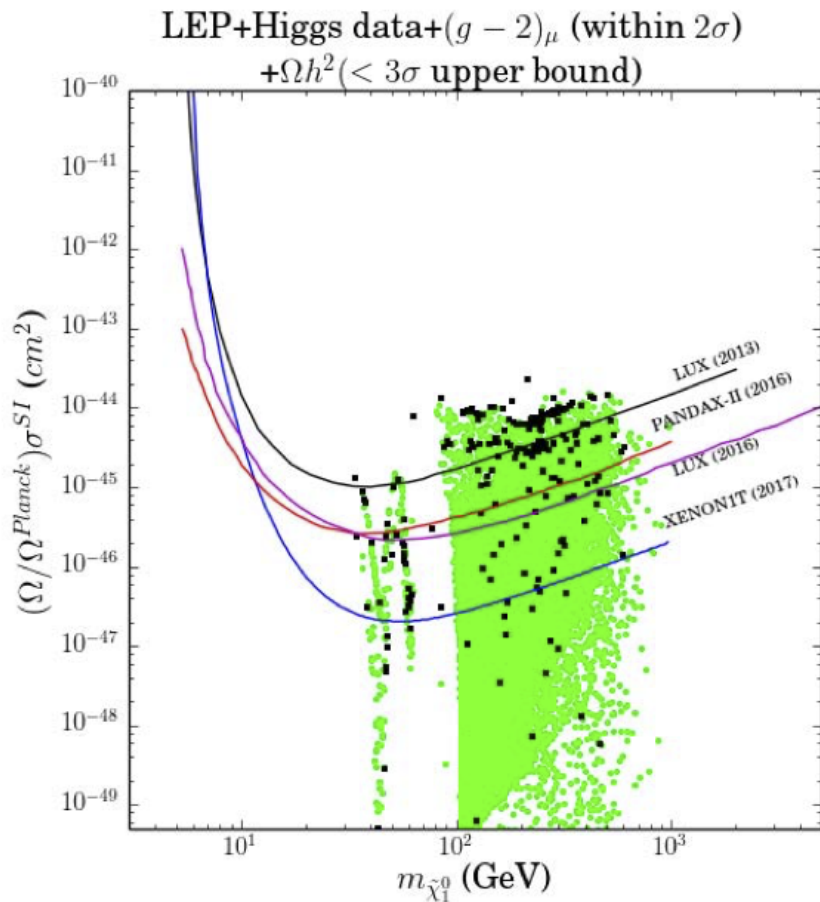
Constraints on χ_1^0 vary for different compositions of Bino, Wino and Higgsinos

- It is well known that pure Bino-like DM relics are typically overabundant (suppressed annihilation cross section), except in the case where the bino co-annihilates with other sparticles
- We can enhance the annihilation rate with a wino or higgsino component in χ_1^0
- To avoid significant constraint, for any LSP abundance less than the relic density, we assume additional DM component (possibly axion-like DM)

Relic Density, Ωh^2



WIMP-nucleon SI Cross Section



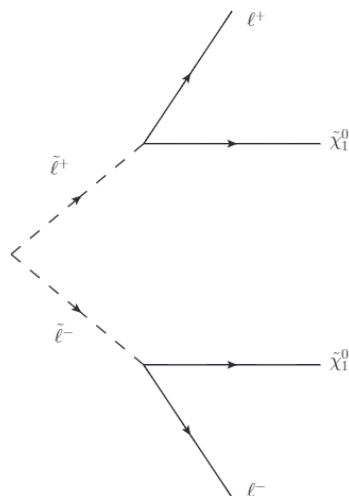
We study constraints from multilepton + MET searches at the LHC.

- We study electroweakinos at $\sqrt{s} = 8$ TeV LHC from slepton/sneutrino and W/Z decays
- Parameter sets that pass the previous collider and direct/indirect dark matter searches are considered
- Points are considered within the 2σ limit of Δa_μ
- We also present the prospects for electroweakino searches with a 100 TeV collider
- NLO events are simulated using MadGraph 5 interfaced with Pythia 6
- These are passed to CheckMATE-1.2.2 to check exclusion limits at 95% CL

Electroweakinos and sleptons at colliders

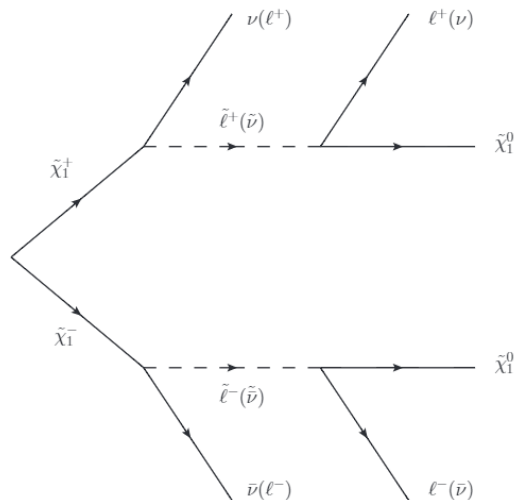
$2\ell + \cancel{E}_T$ (2 leptons + missing energy) ¹

(a)



(a) via direct slepton decays

(b)



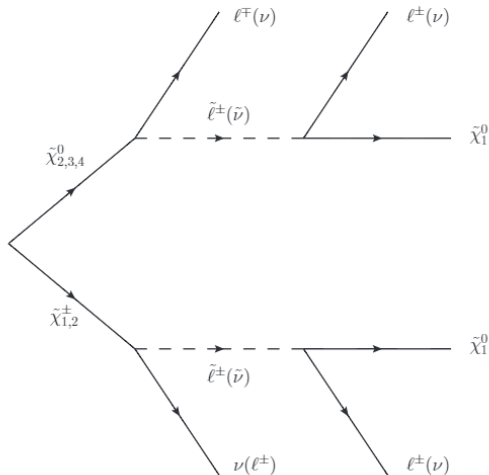
(b) via sleptons/sneutrinos

¹atlas-conf-2013-049

Electroweakinos and sleptons at colliders

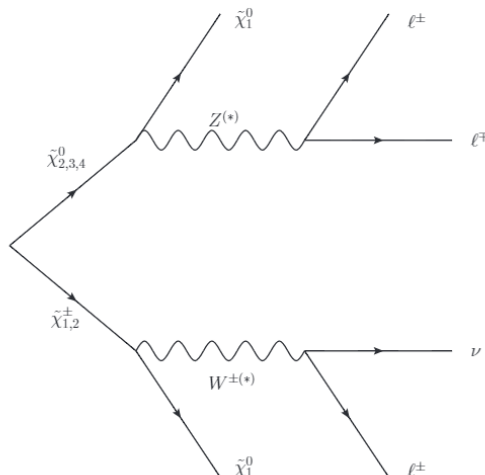
$$3\ell + \cancel{E}_T \text{ (3 leptons + missing energy) }^2$$

(a)



(a) via sleptons/sneutrinos

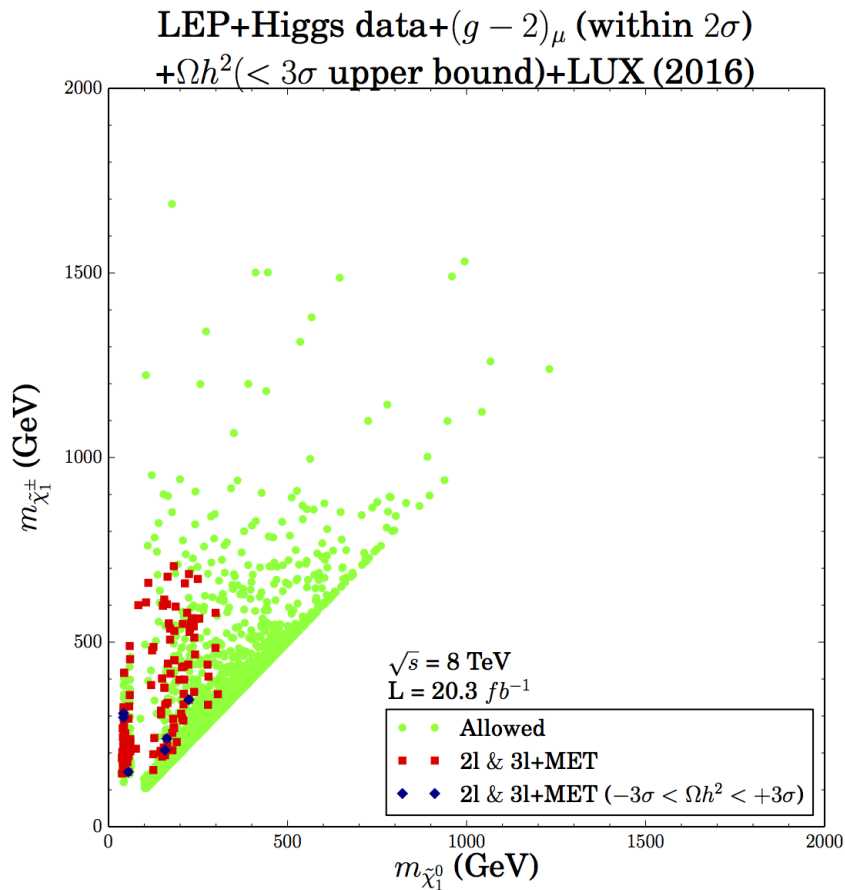
(b)



(b) via gauge bosons

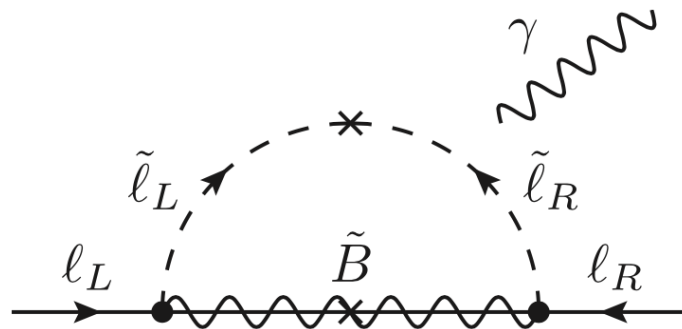
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Results in $m_{\tilde{\chi}_1^0} - m_{\tilde{\chi}_1^\pm}$ plane



Large μ case

It has been noted that one can explain the $(g - 2)_\mu$ can be explained with a dominant bino-smuon loop contribution.



This is enhanced with a large **smuon left-right mixing**.

Too large, and this can spoil the electroweak vacuum stability.

Large μ case

We scan the extended region:

$$\begin{aligned}10 < \tan(\beta) < 50, \\ |M_1|, |M_2| < 3 \text{ TeV}, \\ 10 < \mu < 100 \text{ TeV}, \\ 0.1 < m_{\tilde{l}_L}, m_{\tilde{l}_R} < 2 \text{ TeV}, \quad (l = e, \mu)\end{aligned}$$

with staus decoupled at $m_{\tilde{\tau}_L} = m_{\tilde{\tau}_R} = 10 \text{ TeV}$ and $A_\tau = 0$.

To explain $(g - 2)_\mu$ within 2σ , we find upper limits of $m_{\tilde{\chi}_1^0} < 2.4 \text{ TeV}$ and $m_{\tilde{\ell}_1} < 1.1 \text{ TeV}$.

The previous DM constraints severely limit this case, and so is **not the preferred scenario**.

100 TeV Analysis

The 3 lepton + MET events at 100 TeV are expected to have the largest reach over the MSSM parameter space.

We scale the signal (S) and background (B) events for the 8 TeV analysis by the ratio:

$$N^{100\text{ TeV}} = (\sigma^{100\text{ TeV}} / \sigma^{8\text{ TeV}}) (3000\text{ fb}^{-1} / 20.3\text{ fb}^{-1}) N^{8\text{ TeV}}$$

Sources of background (B):

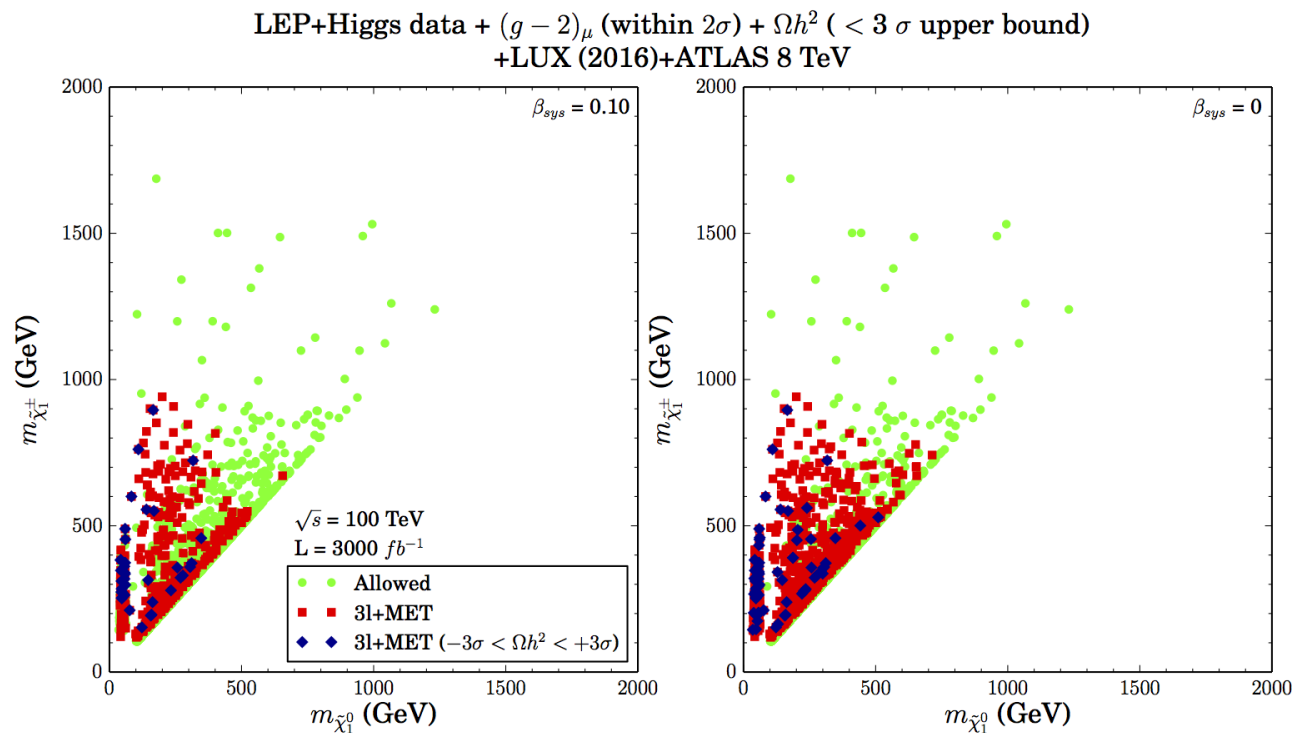
- WZ, ZZ, H
- $ttV + ttZ$
- VVV
- Reducible (t single/pair, WW , single W/Z with jets or photons)

We exclude events corresponding to:

$$\frac{S}{\sqrt{B + (\beta_{\text{sys}} B)^2}} \geq 2$$

where β_{sys} parameterizes the systematic uncertainty.

Results for 100 TeV Analysis



Conclusions

- We studied constraints from direct/indirect measurements on the MSSM with heavy squarks and light sleptons.
- A 100 TeV collider could potentially probe almost the entire mass range for electroweakinos in this model as an explanation for the muon $(g - 2)_\mu$ and dark matter.
- One can further the analysis using monojet-like signals with greater sensitivity to the degenerate mass region.
- Our 100 TeV analysis can be considered a preliminary one, that can be improved once the collider environment details are known (and/or a public code is released).