



Electroweak baryogenesis in the \mathbb{Z}_3 -invariant NMSSM

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Electroweak Baryogenesis

Sakharov Conditions

- B number violation
- departure from equilibrium
- C and CP violation

The Electroweak mechanism

At high temperature EW symmetry is restored before the EWPT.

Result in terms of Sakharov conditions

- EW sphalerons lead to unsuppressed B number violation in symmetric phase
- The EWPT proceeds by bubble nucleation (SFOEWPT)
- C and CP violating interactions with the bubble wall lead bias the sphalerons

Shortcomings of SM

- EWPT is crossover for $m_h \sim 125$
- not enough CP violation

Shortcomings of MSSM

Pheno shortcomings

- little hierarchy problem
- μ problem

EWBG shortcomings

- stops need to be light enough for SFOEWPT
- either stops need to be near degenerate for BAU or Higgsino and a gaugino need to be light and near degenerate
- EDMs and LHC constraints suffocate the MSSM

Phase transitions in NMSSM

Two types of phase transitions

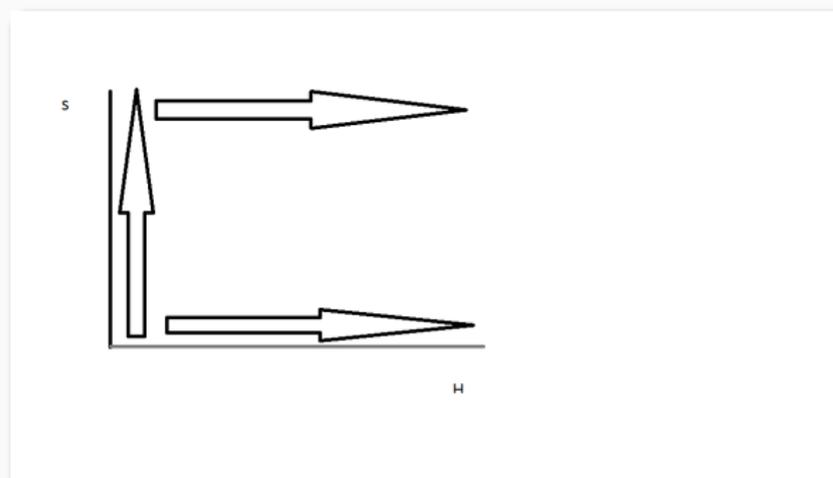


Figure 1: (Type I vs Type II)

- stops can be heavy if singlet catalyze SFOPT
- $\Delta\beta \propto \text{BAU}$ can be an order of magnitude larger.

New CP violating Source

New rephasing invariant: Higgsino-Singlino interactions with space time varying vevs

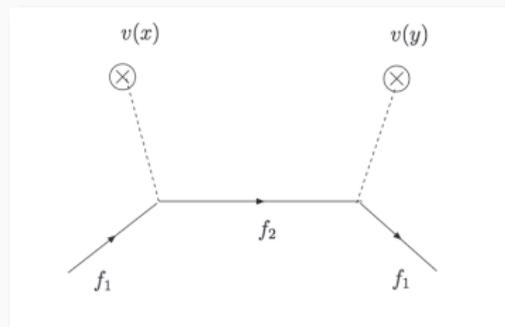


Figure 2: (Higgsino-Singlino interactions)

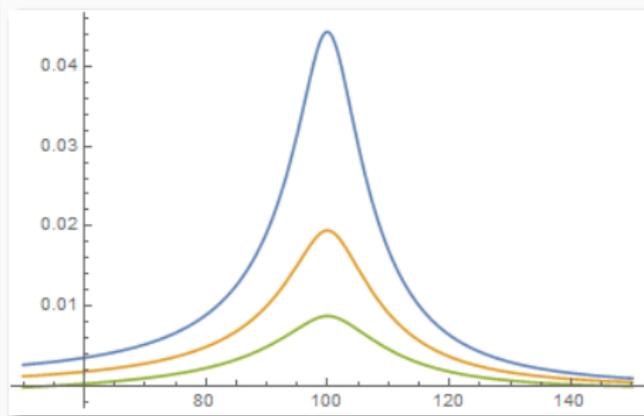
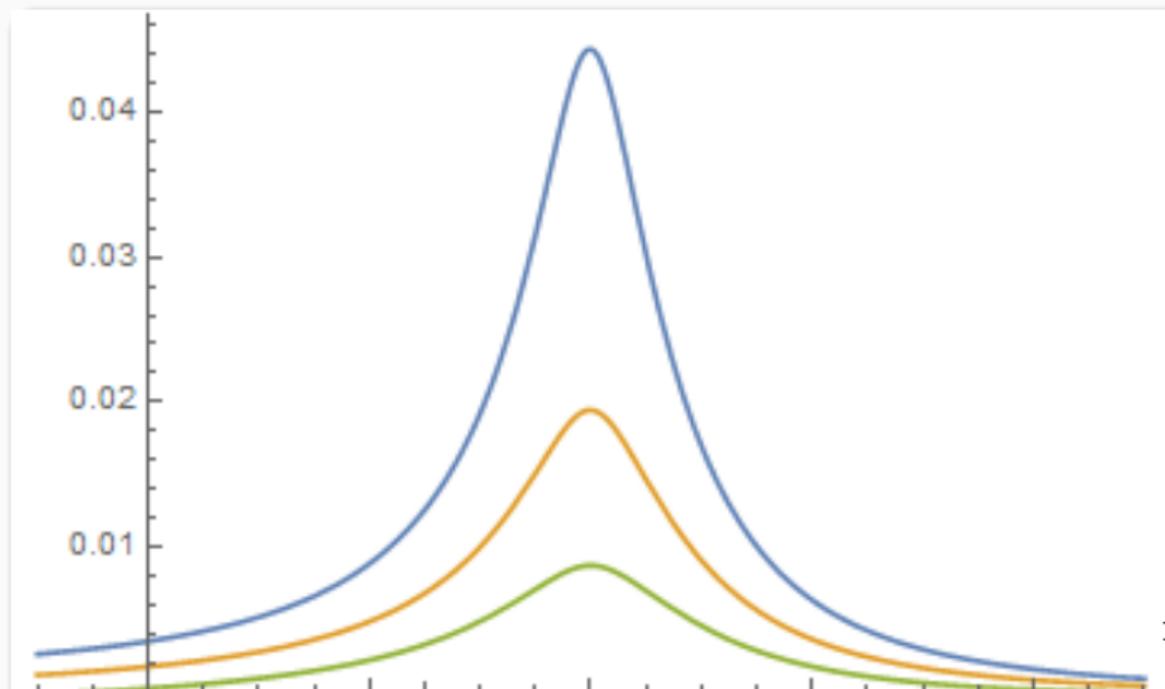


Figure 3: (Singlino source strength)

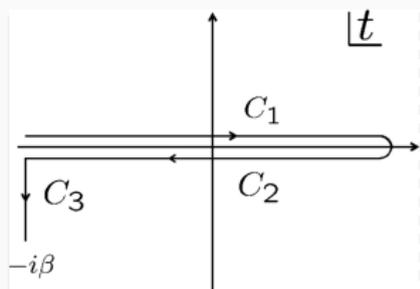
Thermal widths

$$\Gamma_{\tilde{g}} = b(m_H^*)0.01(\lambda^2 + \kappa^2) \quad (1)$$



Baryon Asymmetry in the NMSSM

Use CTP Formalism



- Doubles the number of Greens functions
- Use techniques from

C. Lee, V. Cirigliano, and M. J. Ramsey-Musolf, Phys. Rev. D **71** no 7 075010 (2005)

G. A. White, Phys. Rev. D **93**, no. 4, 043504 (2016)
[arXiv:1510.03901 [hep-ph]].

VEV Insertion Approximation

“The VIA assumes that the particle-antiparticle asymmetry generation is dominated by the region near the phase boundary, where the vevs are small compared to both T and the (mass) difference”

S. Inoue, G. Ovanesyan and M. J. Ramsey-Musolf, Phys. Rev. D **93** .1 015013 (2016).

- Calculating vev insertion diagrams using the VIA is invalid if the mass difference is large

Other Approximations and assumptions

- local equilibrium between quarks and squarks
- Ficks diffusion law $\vec{j} \rightarrow \vec{\nabla} n$
- ignore bubble wall curvature and solve system in bubble wall rest frame z
 - This means $\partial_t \rightarrow v_w \partial_z$
 - the transport equations are inhomogenous coupled ODEs in a single space time variable

Transport equations

Type I phase transitions have no chemical potential for the fluctuations around the singlet vev.

Total list of number densities

$$\begin{aligned}n_{H_1} &= n_{H_u^+} + n_{H_u^0} \\n_{H_2} &= n_{H_d^-} + n_{H_d^0} \\n_{\tilde{H}} &= n_{\tilde{H}_u^+} + n_{\tilde{H}_u^0} - n_{\tilde{H}_d^-} - n_{\tilde{H}_d^0} \\n_t &= n_{t_R} + n_{\tilde{t}_R} \\n_Q &= n_{t_L} + n_{b_L} + n_{\tilde{t}_R} + n_{\tilde{b}_L} \\n_S &= n_S\end{aligned}\tag{2}$$

The baryon number density, ρ_B , satisfies the equation

$$D_Q \rho_B''(z) - v_W \rho_B'(z) - \Theta(-z) \mathcal{R} \rho_B = \Theta(-z) \frac{n_F}{2} \Gamma_{ws} n_L(z)$$

Which can be solved to get

$$Y_B = -\frac{n_F \Gamma_{ws}}{2\kappa_+ D_Q S} \int_{-\infty}^0 e^{-\kappa_- x} n_L(x) dx$$

Results

- Higgsino and Singlino masses in symmetric phase must be $\lesssim 4T_N$
 - this implies zero temperature singlet mass should be $\ll 1$ TeV
- Higgsino and Singlino masses must be within a factor of 2 of each other
- Correct zero-temperature Higgs sector (higgs mass and vev)

Type I vs Type II

- If the Higgs is tachyonic at $S = H = 0$, $\forall T < 200$ it must be a type I phase transition
- Tree level mass terms at $S = H = 0$ tend to be larger than 200 GeV.
- Nearly all points are Type I

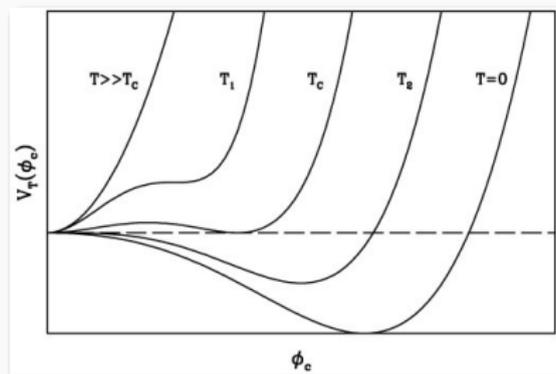


Figure 5

Preliminary (no pheno! Scan running: new results coming soon!)

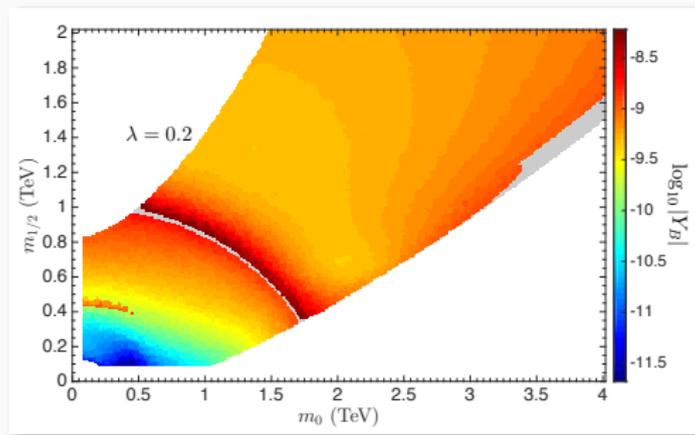


Figure 6: Singlino Higgsino 3 body rates

Effects of three body rates

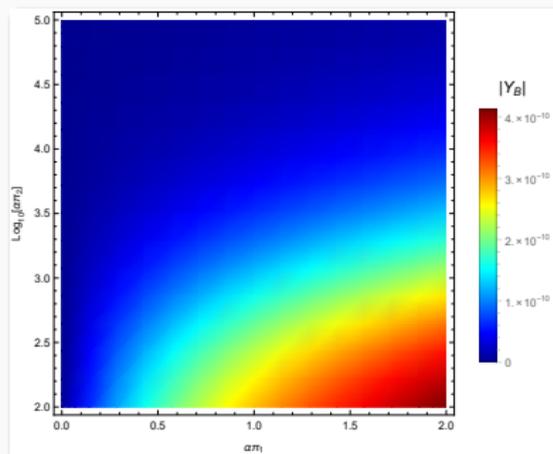
- Three body rates involving Singlinos and Higgsinos:

$$\mu_{H_i} + \mu_{\tilde{H}} \approx 0 \quad (3)$$

whereas supergauge relaxes combination

$$\mu_{H_i} - \mu_{\tilde{H}} \approx 0 \quad (4)$$

- But $\text{BAU} \propto \mu_{\tilde{H}}!$



Summary

Electroweak baryogenesis in the NMSSM overcomes some problems the MSSM has. Some parts of the parameter space can produce the BAU

Questions?