Constraints on shear and rotation with massive galaxy clusters

Ahmad Mehrabi

Bu-Ali Sina university, Hamadan-Iran

CosPA2016,Sydney Nov, 2016

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal CosPA2016, Sydney Nov, 2016 1 / 12

1 Introduction

1 Introduction

- 2 Extended spherical collapse model
 - Evolution of the density contrast
 - Critical density contrast and virial overdensity
 - Massive galaxy cluster number count

1 Introduction

- 2 Extended spherical collapse model
 - Evolution of the density contrast
 - Critical density contrast and virial overdensity
 - Massive galaxy cluster number count
- Observational constrain on the shear and rotation parameter
 Data used in this work

1 Introduction

- 2 Extended spherical collapse model
 - Evolution of the density contrast
 - Critical density contrast and virial overdensity
 - Massive galaxy cluster number count
- Observational constrain on the shear and rotation parameter
 Data used in this work
- 4 Results

Spherical collapse model

 Spherical collapse model (a spherical region with slightly higher density, expand like background and the expansion rate decreases due to gravity and then stop at turn-around and collapse and finally virialize)

- Spherical collapse model (a spherical region with slightly higher density, expand like background and the expansion rate decreases due to gravity and then stop at turn-around and collapse and finally virialize)
- SC provide a useful method to count number of collapsed haloes

- Spherical collapse model (a spherical region with slightly higher density, expand like background and the expansion rate decreases due to gravity and then stop at turn-around and collapse and finally virialize)
- SC provide a useful method to count number of collapsed haloes
- We add the shear and rotation to DM fluid

- Spherical collapse model (a spherical region with slightly higher density, expand like background and the expansion rate decreases due to gravity and then stop at turn-around and collapse and finally virialize)
- SC provide a useful method to count number of collapsed haloes
- We add the shear and rotation to DM fluid
- Study how it affects number of massive clusters

- Spherical collapse model (a spherical region with slightly higher density, expand like background and the expansion rate decreases due to gravity and then stop at turn-around and collapse and finally virialize)
- SC provide a useful method to count number of collapsed haloes
- We add the shear and rotation to DM fluid
- Study how it affects number of massive clusters
- Constrain the cosmological parameters including the shear and rotation parameter using observational data

Related equations

Related equations

The equations describing the evolution of the density contrast of a fluid

$$(\bar{b}_{j} + 3H(c_{\text{eff},j}^{2} - \bar{w}_{j})\delta_{j} + [1 + \bar{w}_{j} + (1 + c_{\text{eff},j}^{2})\delta_{j}]\vec{\nabla} \cdot \vec{u}_{j} = 0$$
 (1)

$$\dot{\vec{u}}_{j} + 2H\vec{u}_{j} + (\vec{u}_{j} \cdot \vec{\nabla}\vec{u}_{j}) + \frac{1}{a^{2}}\vec{\nabla}\phi = 0$$
 (2)

$$\nabla^2 \phi - 4\pi G a^2 \sum_{k} \rho_k \delta_k (1 + 3c_{\text{eff},k}^2) = 0$$
 (3)

Related equations

The equations describing the evolution of the density contrast of a fluid

$$(\bar{b}_{j} + 3H(c_{\text{eff},j}^{2} - \bar{w}_{j})\delta_{j} + [1 + \bar{w}_{j} + (1 + c_{\text{eff},j}^{2})\delta_{j}]\vec{\nabla} \cdot \vec{u}_{j} = 0$$
 (1)

$$\dot{\vec{u}}_{j} + 2H\vec{u}_{j} + (\vec{u}_{j} \cdot \vec{\nabla}\vec{u}_{j}) + \frac{1}{a^{2}}\vec{\nabla}\phi = 0$$
 (2)

$$\nabla^2 \phi - 4\pi G a^2 \sum_{\mathbf{k}} \rho_{\mathbf{k}} \delta_{\mathbf{k}} (1 + 3c_{\text{eff},k}^2) = 0$$
(3)

Add the shear and rotation

$$\dot{\theta} + 2H\theta + \frac{1}{3}\theta^2 + \sigma^2 - \omega^2 + \frac{1}{a^2}\vec{\nabla}^2\phi = 0$$
, (4)

where θ is divergence of peculiar velocity

Related equations

Related equations

The shear and rotation tensor are given by

$$\sigma_{ij} = \frac{1}{2} \left(\frac{\partial u^{i}}{\partial x^{i}} + \frac{\partial u^{i}}{\partial x^{j}} \right) - \frac{1}{3} \theta \delta_{ij}$$
$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u^{i}}{\partial x^{i}} - \frac{\partial u^{i}}{\partial x^{j}} \right)$$

(5)

(6)

Related equations

The shear and rotation tensor are given by

$$\sigma_{ij} = \frac{1}{2} \left(\frac{\partial u^{i}}{\partial x^{i}} + \frac{\partial u^{i}}{\partial x^{j}} \right) - \frac{1}{3} \theta \delta_{ij}$$

$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u^{j}}{\partial x^{i}} - \frac{\partial u^{i}}{\partial x^{j}} \right)$$
(6)

5/12

Considering α as the ratio of the rotational and gravitational terms

$$\delta_{\rm m}^{\prime\prime} + \left(\frac{3}{a} + \frac{E^{\prime}}{E}\right)\delta_{\rm m}^{\prime} - \frac{4}{3}\frac{{\delta^{\prime}}_{\rm m}^2}{1 + \delta_{\rm m}} - \frac{3}{2a^5 E^2}(1 - \alpha)\Omega_{\rm m,0}\delta_{\rm m}(1 + \delta_{\rm m}) = 0$$
(7)

Related equations

The shear and rotation tensor are given by

$$\sigma_{ij} = \frac{1}{2} \left(\frac{\partial u^{i}}{\partial x^{i}} + \frac{\partial u^{i}}{\partial x^{j}} \right) - \frac{1}{3} \theta \delta_{ij}$$

$$\omega_{ij} = \frac{1}{2} \left(\frac{\partial u^{i}}{\partial x^{i}} - \frac{\partial u^{i}}{\partial x^{j}} \right)$$
(6)

Considering α as the ratio of the rotational and gravitational terms

$$\delta_{\rm m}^{\prime\prime} + \left(\frac{3}{a} + \frac{E^{\prime}}{E}\right)\delta_{\rm m}^{\prime} - \frac{4}{3}\frac{{\delta^{\prime}}_{\rm m}^2}{1 + \delta_{\rm m}} - \frac{3}{2a^5 E^2}(1 - \alpha)\Omega_{\rm m,0}\delta_{\rm m}(1 + \delta_{\rm m}) = 0$$
(7)

The linear equation is

$$\delta_{\rm m}^{\prime\prime} + \left(\frac{3}{a} + \frac{E^{\prime}}{E}\right)\delta_{\rm m}^{\prime} - \frac{3}{2a^5 E^2}\Omega_{{\rm m},0}\delta_{\rm m} = 0 \tag{8}$$

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal

Finding the critical density contrast and virial overdensity

Finding the critical density contrast and virial overdensity

First step: initial conditions which lead to collapse in non-linear equation second step: the critical density contrast is the solution of linear equation with the initial conditions

Finding the critical density contrast and virial overdensity

- First step: initial conditions which lead to collapse in non-linear equation second step: the critical density contrast is the solution of linear equation with the initial conditions
- Virial overdensity is given by Δ_{vir} = ζ(x/y)³ where ζ is the overdensity at turn-around, x is the scale factor divided by the turn-around scale factor and y is the ratio between the virialised radius and the turn-around radius

Finding the critical density contrast and virial overdensity

- First step: initial conditions which lead to collapse in non-linear equation second step: the critical density contrast is the solution of linear equation with the initial conditions
- Virial overdensity is given by Δ_{vir} = ζ(x/y)³ where ζ is the overdensity at turn-around, x is the scale factor divided by the turn-around scale factor and y is the ratio between the virialised radius and the turn-around radius



6/12

Finding the critical density contrast and virial overdensity

- First step: initial conditions which lead to collapse in non-linear equation second step: the critical density contrast is the solution of linear equation with the initial conditions
- Virial overdensity is given by Δ_{vir} = ζ(x/y)³ where ζ is the overdensity at turn-around, x is the scale factor divided by the turn-around scale factor and y is the ratio between the virialised radius and the turn-around radius



A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal



Mass function

$$\frac{dn(M,z)}{dM} = \frac{\bar{\rho}_0}{M} \frac{d\nu(M,z)}{dM} f(\nu)$$

(9)

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal CosPA2016, SydneyNov, 2016 7 / 12

Mass function

$$\frac{dn(M,z)}{dM} = \frac{\bar{\rho}_0}{M} \frac{d\nu(M,z)}{dM} f(\nu)$$
(9)

• ST mass function ($\nu = \frac{\delta_c}{\sigma}$)

$$f(\nu) = 0.2709 \sqrt{\frac{2}{\pi}} \left(1 + 1.1096\nu^{0.3} \right) \exp\left(-\frac{0.707\nu^2}{2}\right)$$
(10)

Mass function

$$\frac{dn(M,z)}{dM} = \frac{\bar{\rho}_0}{M} \frac{d\nu(M,z)}{dM} f(\nu)$$
(9)

• ST mass function ($\nu = \frac{\delta_c}{\sigma}$)

$$f(\nu) = 0.2709 \sqrt{\frac{2}{\pi}} \left(1 + 1.1096\nu^{0.3} \right) \exp\left(-\frac{0.707\nu^2}{2}\right)$$
(10)

• The amplitude of mass fluctuations σ

$$\sigma^{2}(R,z) = \frac{1}{2\pi^{2}} \int_{0}^{\infty} k^{2} P(k,z) W^{2}(kR) dk$$
(11)

Mass function

$$\frac{dn(M,z)}{dM} = \frac{\bar{\rho}_0}{M} \frac{d\nu(M,z)}{dM} f(\nu)$$
(9)

ST mass function $(\nu = \frac{\delta_c}{\sigma})$

$$f(\nu) = 0.2709 \sqrt{\frac{2}{\pi}} \left(1 + 1.1096\nu^{0.3} \right) \exp\left(-\frac{0.707\nu^2}{2}\right)$$
(10)

The amplitude of mass fluctuations σ

$$\sigma^{2}(R,z) = \frac{1}{2\pi^{2}} \int_{0}^{\infty} k^{2} P(k,z) W^{2}(kR) dk$$
(11)

CosPA2016,SydneyNov, 2016

The linear matter power spectrum at redshift z

$$P(k,z) = P_0(k)T^2(k)D^2(z)$$

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal



■ The comoving number density of clusters above a certain mass *M*₀

$$n(>M_0,z) = \int_{M_0}^{\infty} \frac{dn(M',z)}{dM'} dM'$$
(13)

■ The comoving number density of clusters above a certain mass *M*₀

$$n(>M_0,z) = \int_{M_0}^{\infty} \frac{dn(M',z)}{dM'} dM'$$
(13)

The comoving number of clusters per unit redshift

$$N_{\rm bin} = n(M > M_0, z) \frac{dV}{dz} \tag{14}$$

The comoving number density of clusters above a certain mass M₀

$$n(>M_0,z) = \int_{M_0}^{\infty} \frac{dn(M',z)}{dM'} dM'$$
(13)

The comoving number of clusters per unit redshift

$$N_{\rm bin} = n(M > M_0, z) \frac{dV}{dz} \tag{14}$$

Total number count of massive clusters

$$N = \int_0^z N_{\rm bin} dz' = \int_0^z n(M > M_0, z') \frac{dV}{dz'} dz'$$
(15)

The comoving number density of clusters above a certain mass M₀

$$n(>M_0,z) = \int_{M_0}^{\infty} \frac{dn(M',z)}{dM'} dM'$$
(13)

The comoving number of clusters per unit redshift

$$N_{\rm bin} = n(M > M_0, z) \frac{dV}{dz} \tag{14}$$

Total number count of massive clusters

$$N = \int_{0}^{z} N_{\rm bin} dz' = \int_{0}^{z} n(M > M_{0}, z') \frac{dV}{dz'} dz'$$
(15)

• We examine a logarithmic relation for $\alpha \Rightarrow \alpha = -\beta \log_{10} \frac{M}{M_{e}}$





Data sets

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal CosPA2016, SydneyNov, 2016 10 / 12

Data sets

To fix the background

- To fix the background
 - 1 Snla: Union 2.1 sample (Suzuki et al. 2012) includes 580 Sn

- To fix the background
 - 1 Snla: Union 2.1 sample (Suzuki et al. 2012) includes 580 Sn
 - 2 Baryon Acoustic Oscillation

- To fix the background
 - 1 Snla: Union 2.1 sample (Suzuki et al. 2012) includes 580 Sn
 - 2 Baryon Acoustic Oscillation
 - 3 Cosmic Microwave Background

- To fix the background
 - 1 Snla: Union 2.1 sample (Suzuki et al. 2012) includes 580 Sn
 - 2 Baryon Acoustic Oscillation
 - 3 Cosmic Microwave Background
- To constrain the shear and rotation

- To fix the background
 - 1 Snla: Union 2.1 sample (Suzuki et al. 2012) includes 580 Sn
 - 2 Baryon Acoustic Oscillation
 - 3 Cosmic Microwave Background
- To constrain the shear and rotation
 - 1 Number of massive clusters collected by (Campanelli et al. 2012) The data is number of massive clusters with $M > 8 \times 10^{14} h^{-1} M_{\odot}$ up to redshift 0.9

Results

Summary of results

Parameters	Best fit value	$1 - \sigma$	$2 - \sigma$	$3 - \sigma$
$\Omega_{\rm m}$	0.284	± 0.0064	$+0.013 \\ -0.012$	$+0.017 \\ -0.016$
h	0.678	± 0.017	$+0.032 \\ -0.033$	$+0.038 \\ -0.043$
β	0.0019	$+0.0008 \\ -0.0015$	< 0.0043	< 0.0054

Results

Summary of results

Parameters	Best fit value	$1 - \sigma$	$2 - \sigma$	$3 - \sigma$
$\Omega_{\rm m}$	0.284	± 0.0064	$+0.013 \\ -0.012$	$+0.017 \\ -0.016$
h	0.678	± 0.017	$+0.032 \\ -0.033$	$+0.038 \\ -0.043$
β	0.0019	$+0.0008 \\ -0.0015$	< 0.0043	< 0.0054



A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal

Thanks for your attention

A. Mehrabi (Bu-Ali Sina university, Hamadan-Iran) Constraints on shear and rotation with massive gal CosPA2016, SydneyNov, 2016 12 / 12