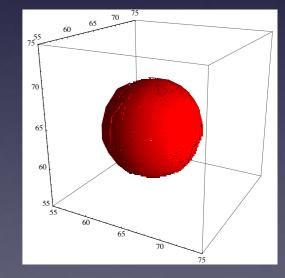
Adiabatic Invariance of I-balls/Oscillons

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Naoyuki Takeda: Keio University (JPN)

Collaboration with Masahiro Kawasaki: ICRR

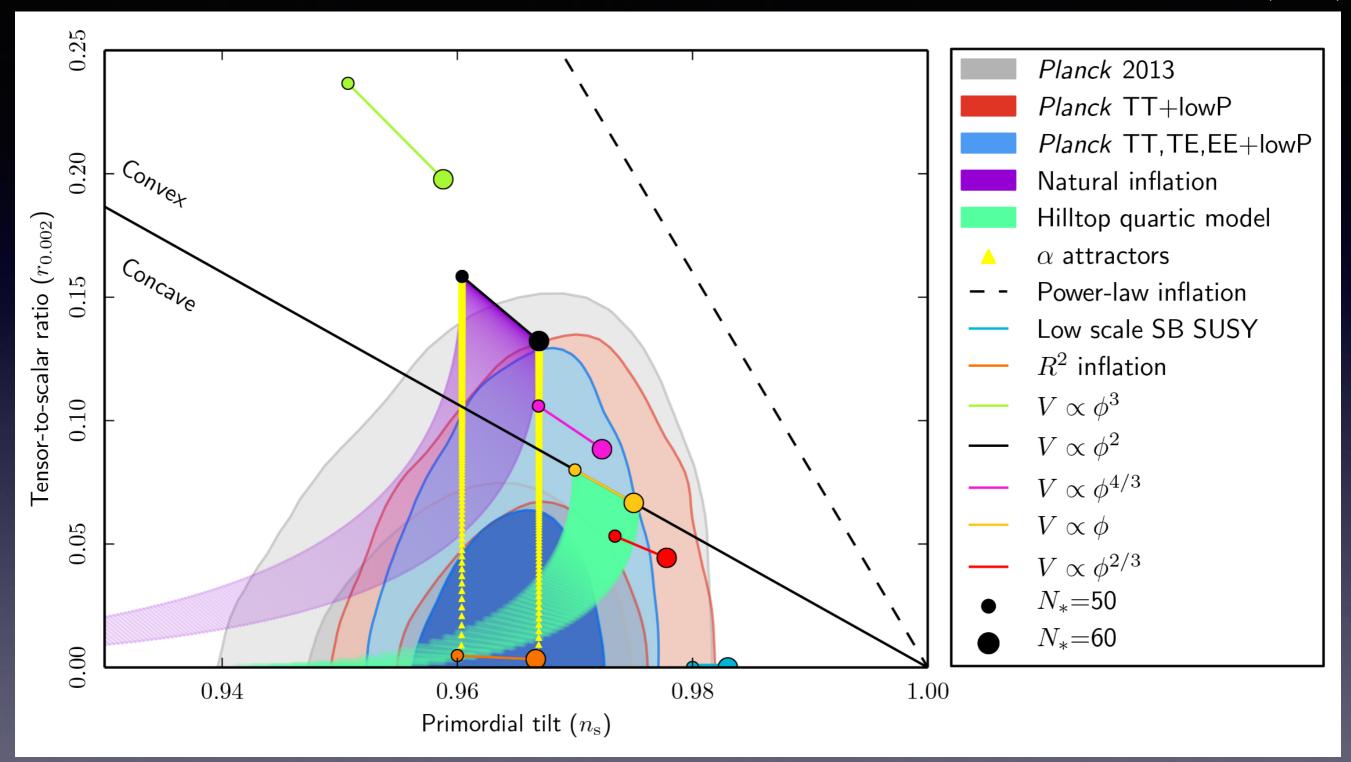
Fuminobu Takahashi: Tohoku University



CosPA2016(Sydney Univ), 1/12/2016

Contents

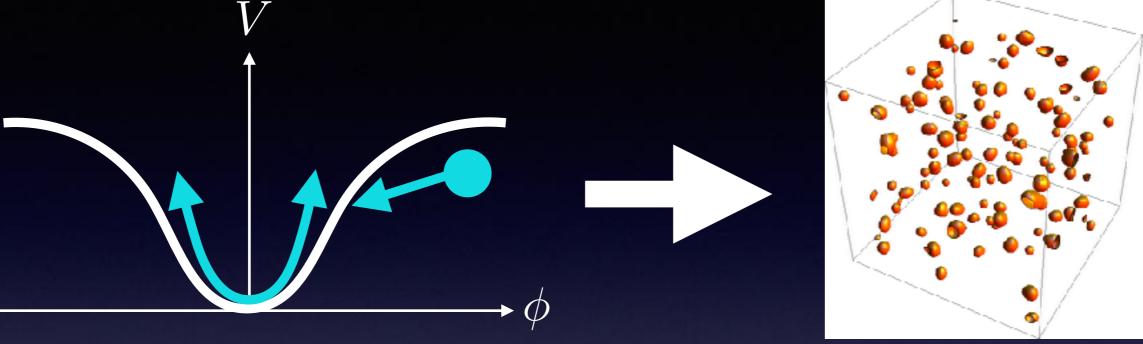
- Introduction
- ▶ I-ball/oscillon
- Adiabatic Invariance
- Work
- Conserv. of Adiabatic Invariance
- Adiabatic deformation of I-ball
- Conclusion



Inflaton potential is flatter than quadratic

Introduction: I-ball/Oscillons

Oscillating real scalar field forms I-balls/Oscillons

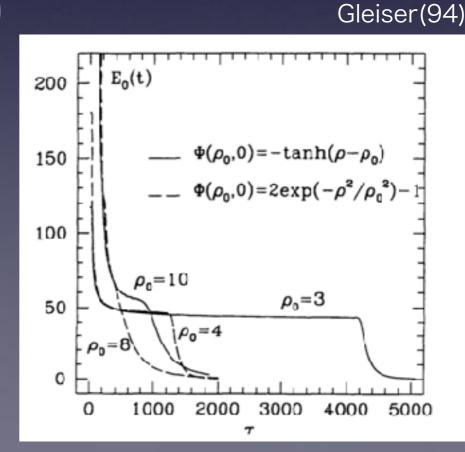


flatter than $m^2\phi^2$

ex):
$$\lambda(\phi^2 - \Lambda^2)$$
, $\frac{m^2}{2}\phi^2 - \lambda\phi^4 + g\frac{\phi^6}{M^2}$, $\Lambda^4\left(1 + \frac{\phi^2}{M^2}\right)$

Amin et al' 12

- Long life time and sudden decay
- Effect on the dynamics of ϕ
 - · decay rate, equation of state etc.
- Try to understand the stability focusing on a conserved charge: adiabatic invariance

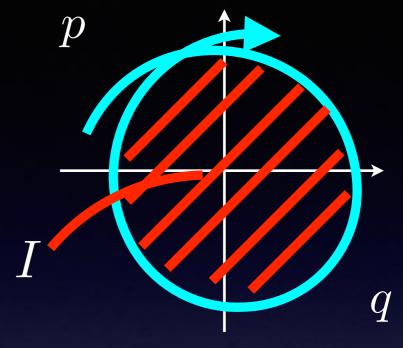


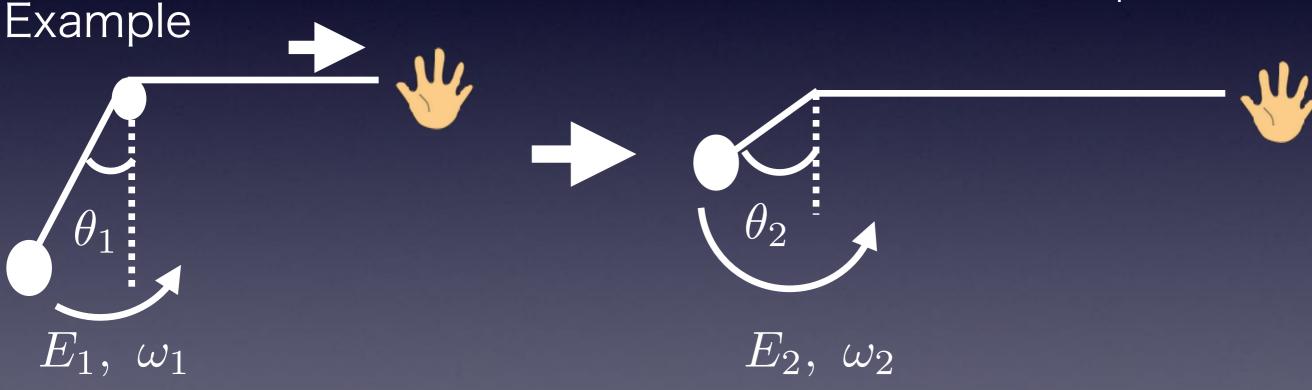
Intro: Adiabatic Invariance (Classical mechanics)

Classical mechanics

 Phase space surrounded by quasiperiodic oscillation is conserved.

$$I_{\rm cl} = \oint dq \cdot p = \frac{\overline{\dot{q}^2}}{\omega}$$





• ω, E change: $\omega_1 < \omega_2, \; \theta_1 < \theta_2, \; E_1 < E_2$ Adiabatic invariance is conserved: $E/\omega = {\rm const}$

Work: Adiabatic Invariance (Scalar field)

Adiabatic Charge

$$I = \int d^3x \oint d\phi \cdot \pi = \int d^3x \frac{\overline{\dot{\phi}^2}}{\omega}$$

 ω : frequency

- We proved that Adiabatic Charge is conserved for scalar field theory
- Assumption
- Field is divided into separable form

$$\phi(t, \vec{x}) = \Phi(\vec{x}) f(t)$$

Constant of motion

$$\tilde{\mathcal{H}} = \mathcal{H} - \frac{1}{2}\partial_i \left(\phi \partial^i \phi\right)$$

= au

1st 2no

$$J^0 \equiv \frac{2\pi}{\omega} \overline{\dot{\phi}^2}, \ J^i \equiv -\frac{\pi}{\omega} \left(\dot{\phi} \partial_i \phi - \phi \partial_i \dot{\phi} \right)$$

a: parameter of external force

Kawasaki Takahashi

Work: Potential for separable form

- Separable form: $\phi(t, \vec{x}) = \Phi(\vec{x}) f(t)$
- ▶ E.O.M.

$$\frac{\ddot{f}}{f} - \frac{\nabla^2 \Phi}{\Phi} = -\frac{\partial V(\Phi f)}{\partial (\Phi f)}$$

- lacktriangledown and f need to evolve independently
- V is the function of ϕ
- C satisfies def. of logarithmic function

$$C(\Phi f) = C(\Phi) + C(f) - A(1) - B(1)$$

$$\rightarrow C(\phi) = \ln(\phi)$$

$$V = \frac{m^2}{2}\phi^2 \left[1 - K \ln \frac{\phi^2}{2M^2} \right]$$

Kawasaki, Takahashi, Takeda '15

m, M: mass parameter, 0 < K < 1

Work: I-ball solution

I-ball solution

Language multiplier with fixed I

$$\bar{E}_{\tilde{\omega}} = \bar{E} + \tilde{\omega} \left[I - \int d^3 x \frac{\bar{\phi}^2}{\omega} \right]$$

potential

$$V = \frac{m^2}{2}\phi^2 \left[1 - K \ln \frac{\phi^2}{2M^2} \right]$$

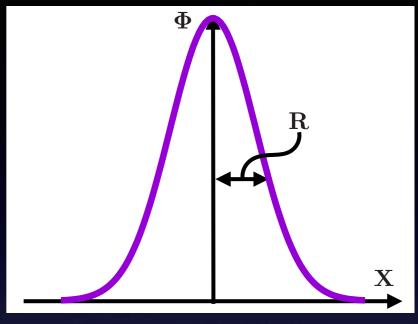
■ Lowest energy condition: $\delta \bar{E}_{\tilde{\omega}}/\delta \Phi = 0$

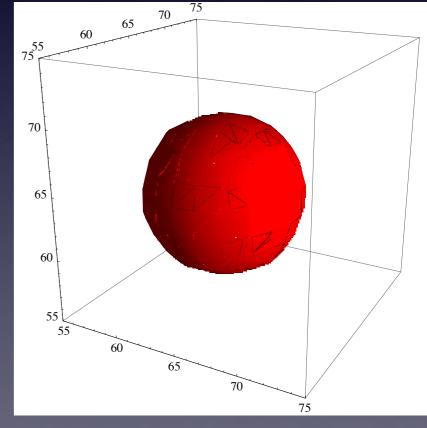
$$\Phi(r) = \Phi_c \exp\left(-r^2/R^2\right), \ R = \sqrt{\frac{2}{K}} \frac{1}{m}$$

- ▶ I-ball is the lowest lowest energy state for a fixed I
- Lemmna

Adiabatic charge is conserved even for $\dot{K} \neq 0$

I-ball should deform following prediction from I





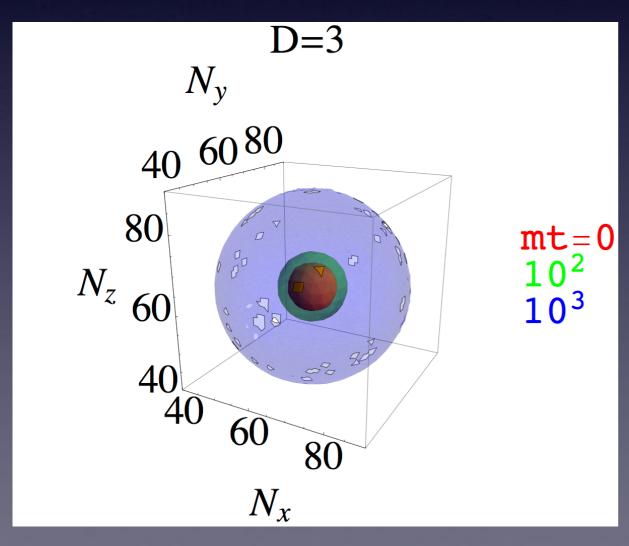
Work: Adiabatic deformation of I-ball

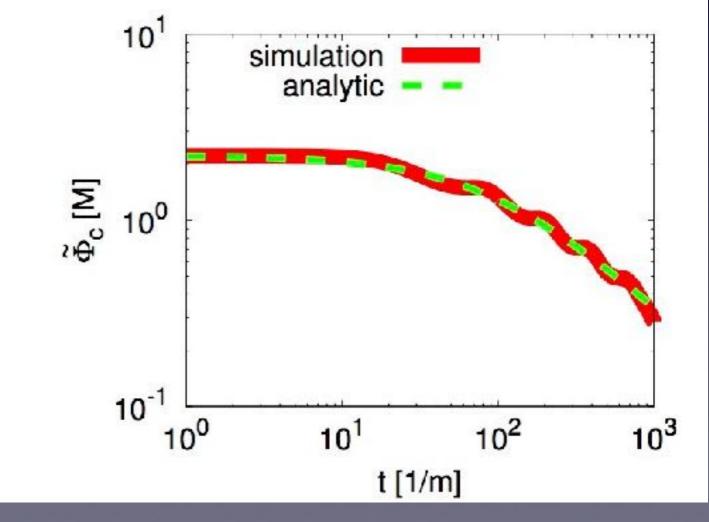
Deformation

$$K(t) = \frac{K_0}{1+\alpha mt}$$
; α , K_0 : dimensionless constant

- Radius: $R(t) = \sqrt{2/K_0}\sqrt{1 + \alpha mt}/m$
- Field amplitude: $\Phi_c = \Phi_{c,0} (1 + \alpha mt)^{-3/4}$
- Numerical simulation

$$\alpha = 10^{-2}$$





Deformation follows the prediction of Adiabatic Charge

Conclusion

- Inflaton potential is suggested to be flatter than quadratic
- Some scalar field with flatter potential forms I-ball
- We focused on the adiabatic charge to verify the I-ball's stability.
- We proved the conservation of the adiabatic charge for the scalar field
- Further revealed that the I-ball deforms adiabatically following the prediction of the adiabatic charge
- Adiabatic Charge is important for the stability of I-ball