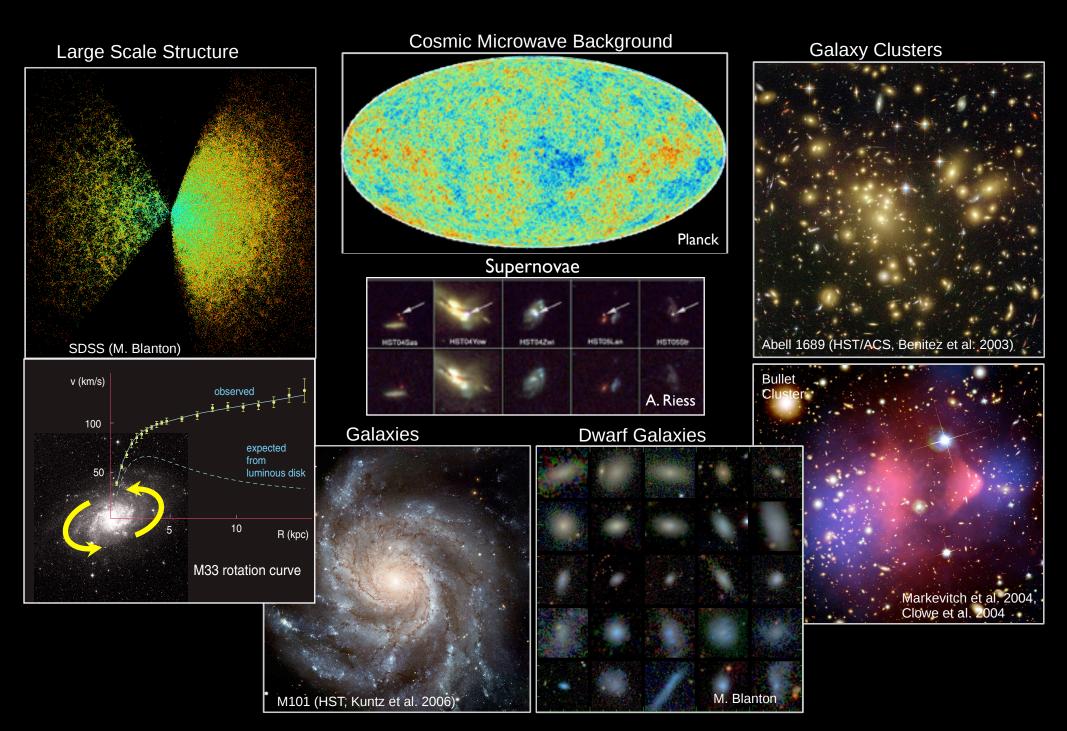
Dark matter theory

Paolo Gondolo University of Utah



Dwarf galaxies

Dwarf galaxies are dominated by dark matter.

40 20

-20

-40

-60

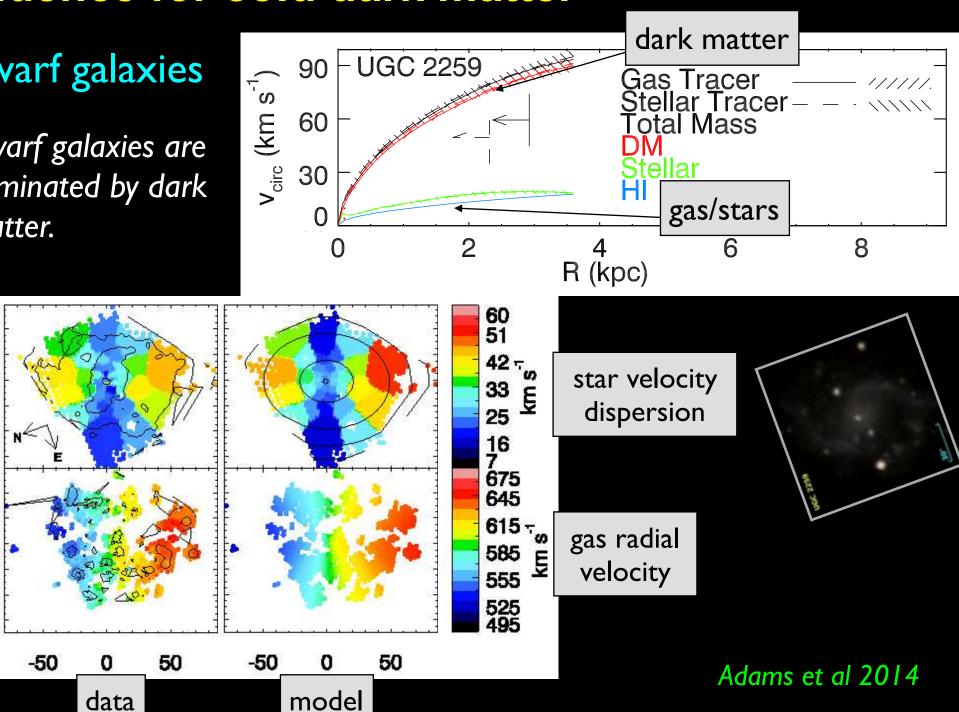
40

20

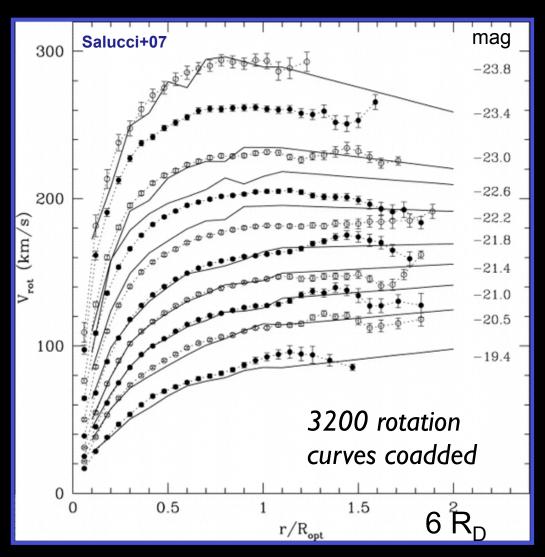
-20

-40

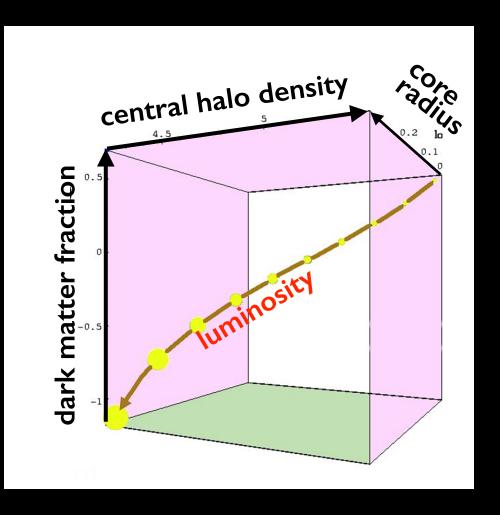
-60



Spiral galaxies



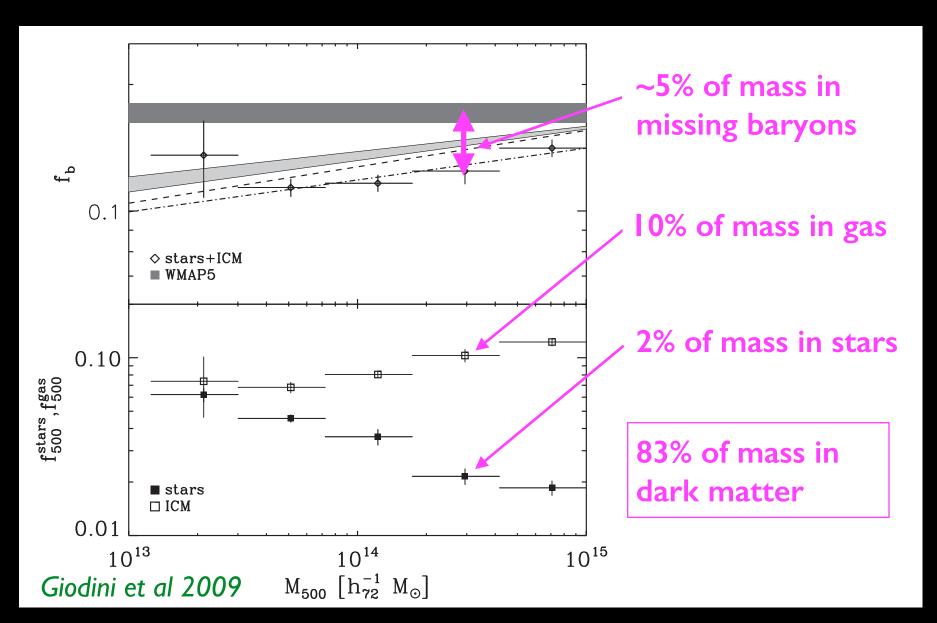
Empirical correlations found from thousands of spiral galaxy rotation curves



Salucci et al 2007

Galaxy clusters

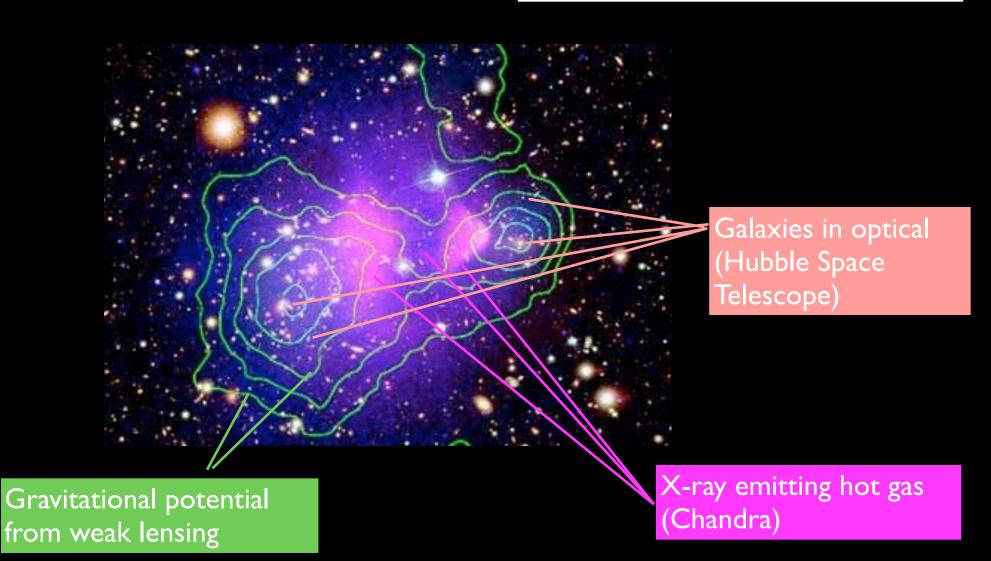
Galaxy clusters are mostly dark matter with some gas and a sprinkle of galaxies



Cold dark matter, not modified gravity

The Bullet Cluster

Symmetry argument: gas is at center, but potential has two wells.

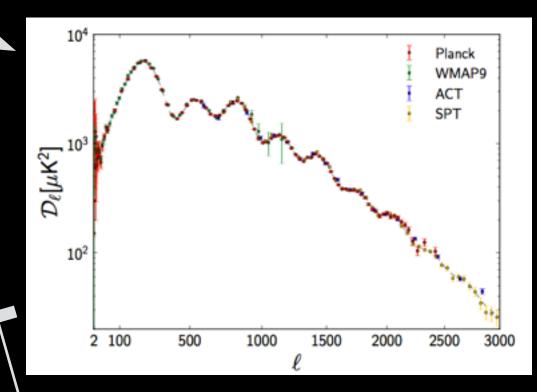


Planck (2013)

Evidence for cold dark matter

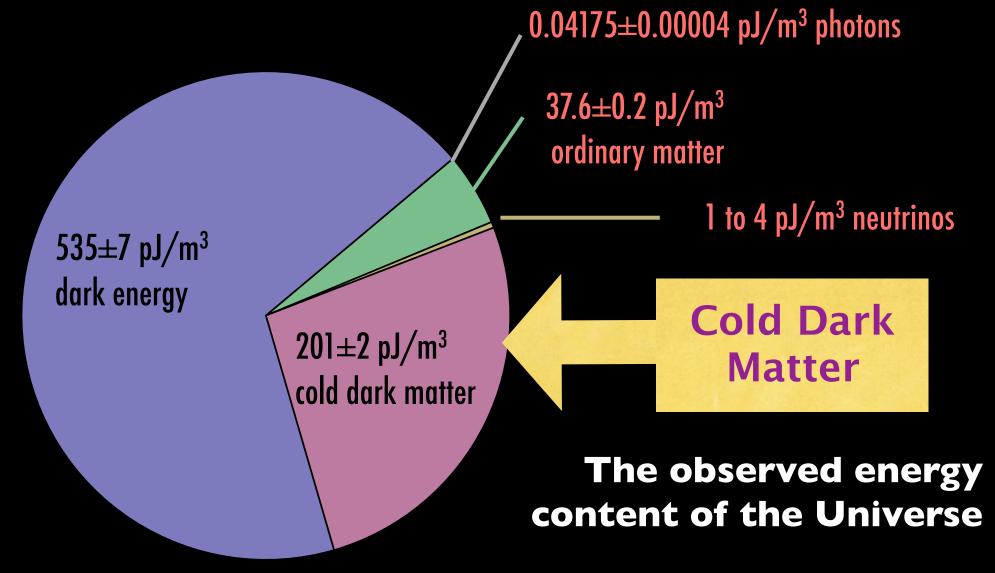
Cosmic Microwave Background fluctuations

	Planck+WP+highL+BAO	
Parameter	Best fit	68% limits
$\Omega_{\mathrm{b}}h^{2}$	0.022161	0.02214 ± 0.00024
$\Omega_{\rm c}h^2$	0.11889	0.1187 ± 0.0017
$100\theta_{\mathrm{MC}}$	1.04148	1.04147 ± 0.00056
τ	0.0952	0.092 ± 0.013
$n_{\rm s}$	0.9611	0.9608 ± 0.0054
$\ln(10^{10}A_{\rm s})\ldots\ldots$	3.0973	3.091 ± 0.025
$\overline{\Omega_{\Lambda} \ldots \ldots \ldots \ldots}$	0.6914	0.692 ± 0.010
σ_8	0.8288	0.826 ± 0.012
Zre	11.52	11.3 ± 1.1
H_0	67.77	67.80 ± 0.77
Age/Gyr	13.7965	13.798 ± 0.037
$100\theta_*$	1.04163	1.04162 ± 0.00056
$r_{\rm drag}$	147.611	147.68 ± 0.45



linear perturbation theory

general relativity and statistical mechanics at 10^4 K ~ 1 eV/k



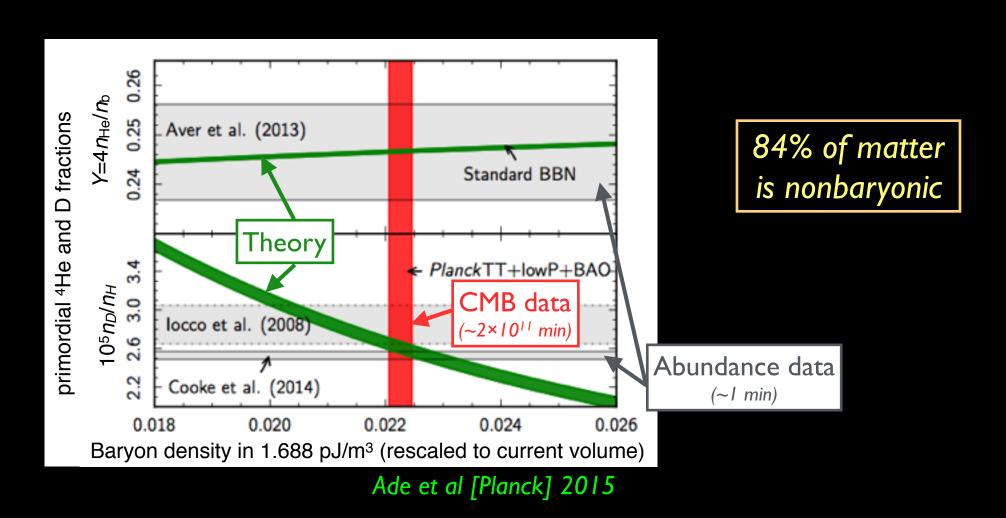
matter $p \ll \rho$ radiation $p = \rho/3$ vacuum $p = -\rho$

Planck (2015)
TT,TE,EE+lowP+lensing+ext

 $1 \text{ pJ} = 10^{-12} \text{ J}$ $\rho_{\text{crit}} = 1688.29 \ h^2 \text{ pJ/m}^3$

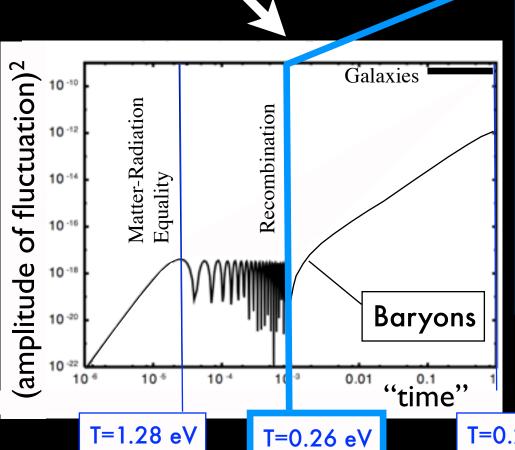
BIG BANG NUCLEOSYNTHESIS

The baryon-to-photon ratio has been the same since ~1 minute after the Big Bang. Baryons are ≤15.7% of the mass in matter.

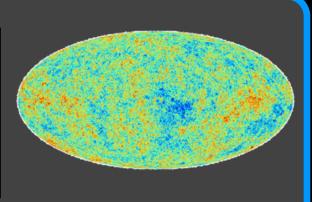


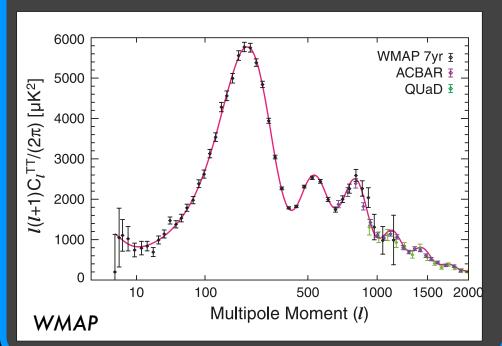


13 billion years ago



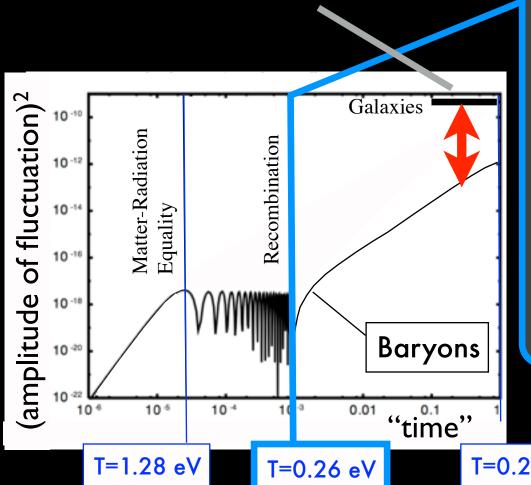
Cosmic Microwave Background fluctuations

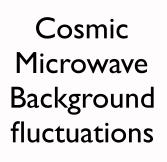


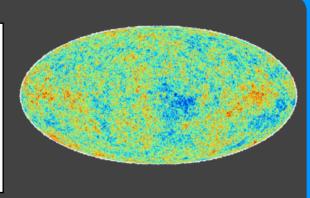


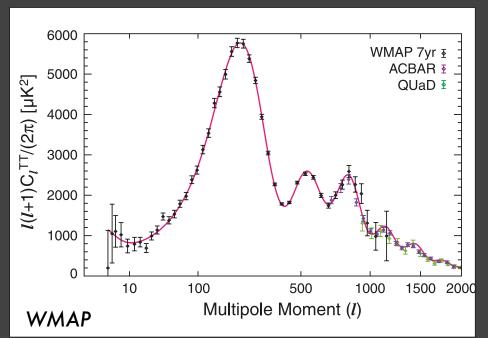
T=0.2348 meV

Without dark matter, fluctuations are too small to gravitationally grow into galaxies in the given 13 billion years.



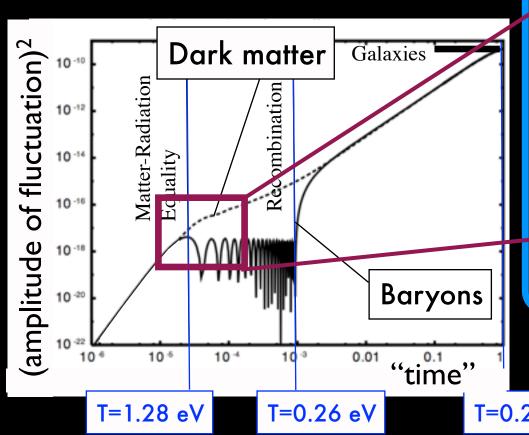






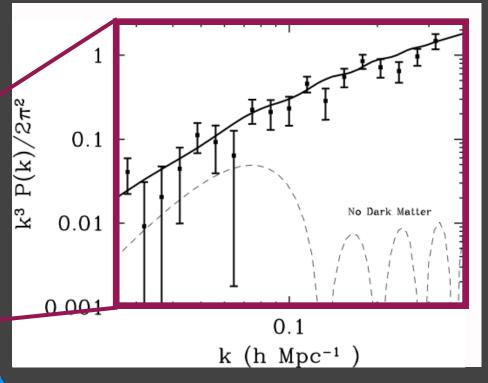
T=0.2348 meV

Dark matter fluctuations, uncoupled to the plasma, start growing early and have enough time to grow into galaxies

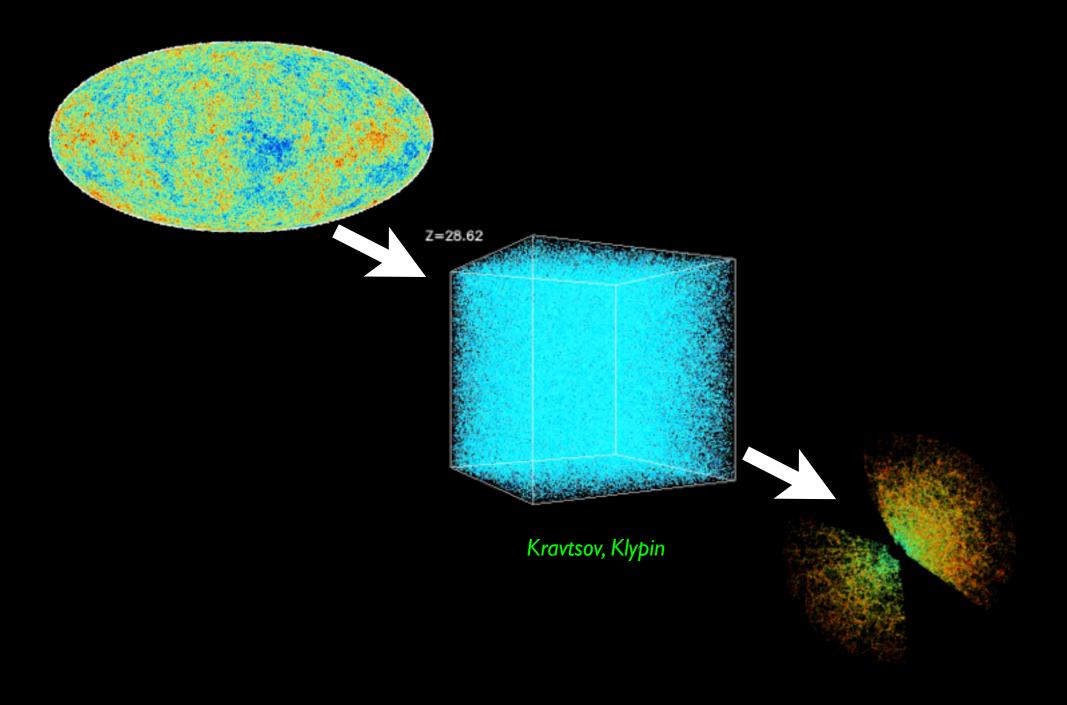


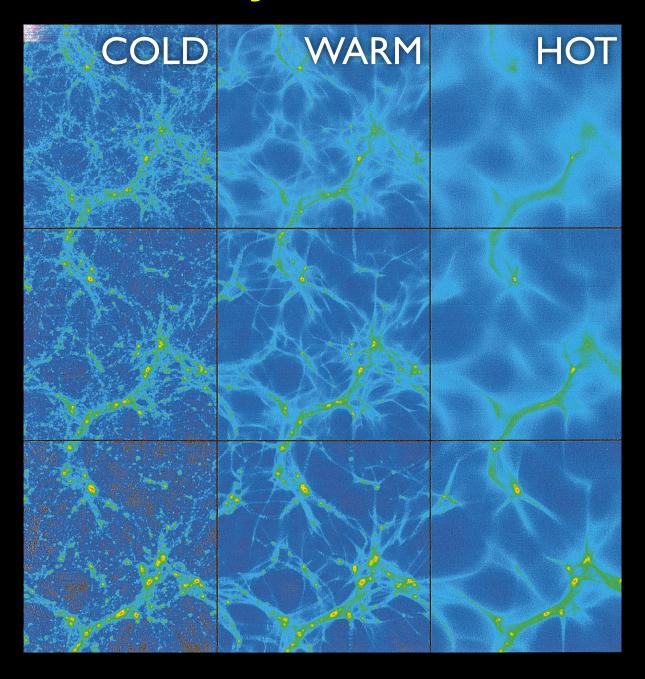
More than 80% of all matter does not couple to the primordial plasma!

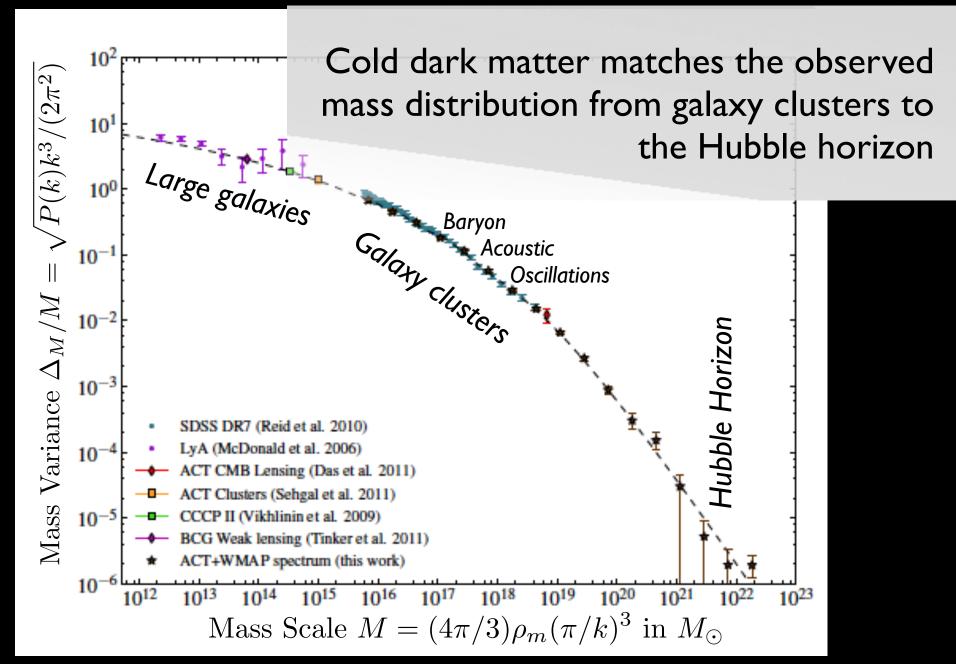
SDSS

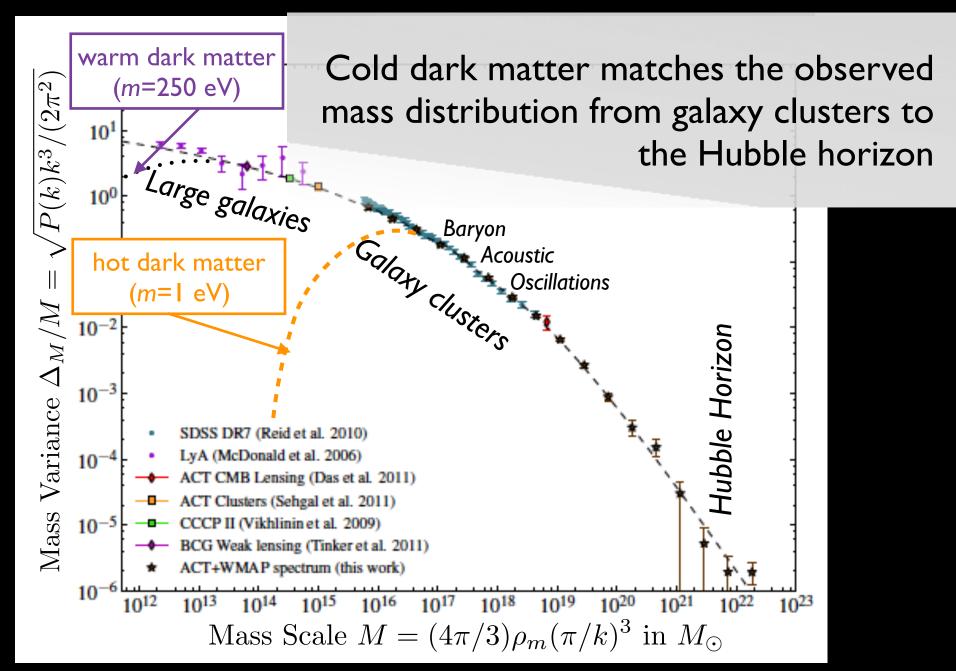


T=0.2348 meV



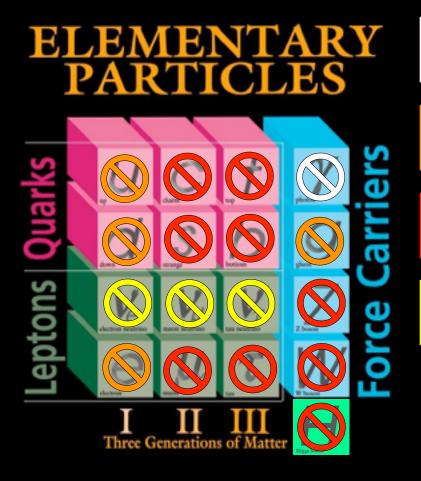








Is dark matter an elementary particle?

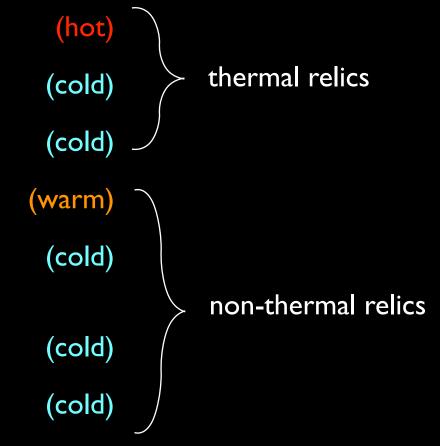


- is the particle of light
- Occuples to the plasma
- disappears too quickly
- is hot dark matter

No known particle can be nonbaryonic cold dark matter!

Particle dark matter

- SM neutrinos
- lightest supersymmetric particle
- lightest Kaluza-Klein particle
- sterile neutrinos, gravitinos
- Bose-Einstein condensates, axions, ultralight scalars, axion clusters
- solitons (Q-balls, B-balls, ...)
- supermassive wimpzillas



Mass range

 $10^{-22} \; {\rm eV} \; (10^{-59} {\rm kg}) \; {\rm B.E.C.s}$ $10^{-8} \; M_{\odot} \; (10^{+22} {\rm kg}) \; {\rm axion \; clusters}$

Interaction strength range

Only gravitational: wimpzillas Strongly interacting: B-balls

Particle dark matter

Hot dark matter

- relativistic at kinetic decoupling (last scattering, start of free streaming)
- big structures form first, then fragment

light neutrinos

Cold dark matter

- non-relativistic at kinetic decoupling
- small structures form first, then merge

neutralinos, axions, WIMPZILLAs, solitons

Warm dark matter

- semi-relativistic at kinetic decoupling
- smallest structures are erased

sterile neutrinos, gravitinos

Particle dark matter

Thermal relics

- in thermal equilibrium with the plasma in the early universe
- produced in collision of plasma particles
- insensitive to initial conditions

```
neutralinos, other WIMPs, ....
```

Non-thermal relics

- not in thermal equilibrium with the plasma in the early universe
- produced in decays of heavier particles or extended structures
- have a memory of initial conditions

```
axions, WIMPZILLAs, solitons, ....
```

QCD axions

QCD axions as dark matter

Hot

Produced thermally in early universe Important for $m_a > 0.1 eV$ ($f_a < 10^8$), mostly excluded by astrophysics

Cold

Produced by coherent field oscillations around mimimum of $V(\theta)$ (Vacuum realignment)

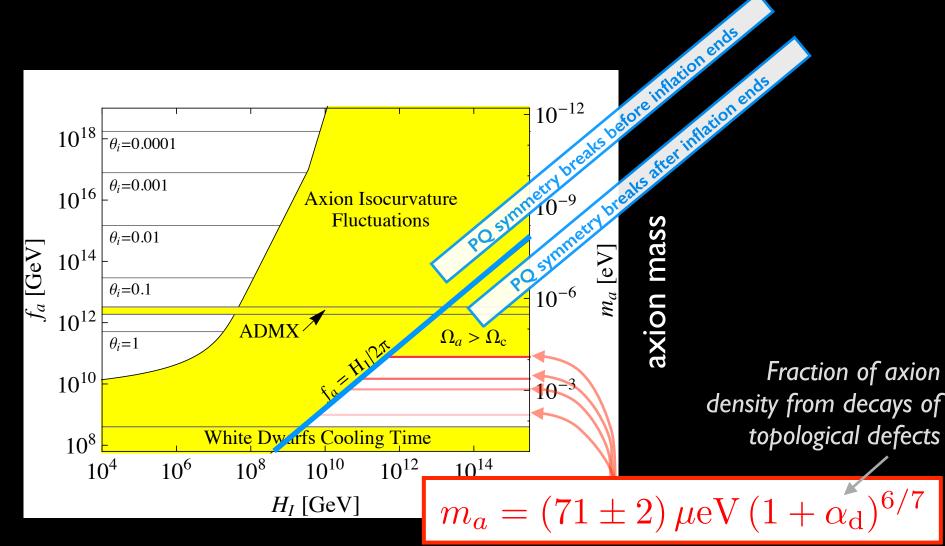
Produced by decay of topological defects

(Axionic string decays)

Still a very complicated and calculation!

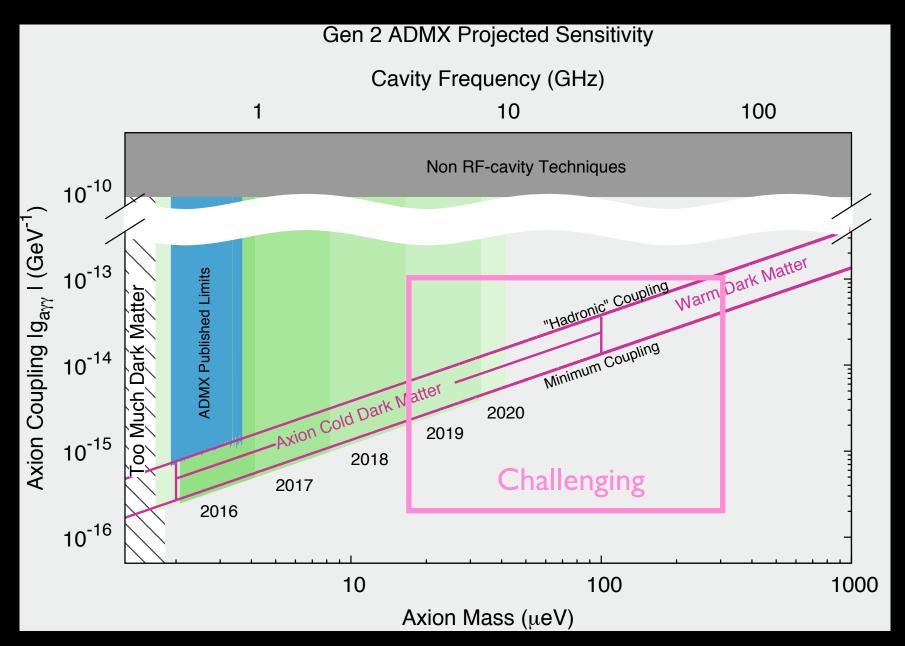
uncertain calculation!

e.g. Hiramatsu et al 2012



Expansion rate at end of inflation

QCD axions as cold dark matter: searches



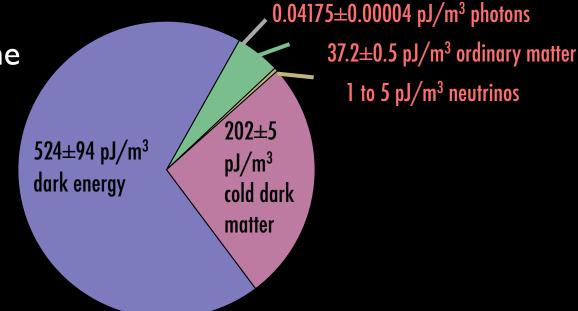
Weakly Interacting Massive Particles

The magnificent WIMP

(Weakly Interacting Massive Particle)

 One naturally obtains the right cosmic density of WIMPs

Thermal production in hot primordial plasma.



One can experimentally test the WIMP hypothesis

The same physical processes that produce the right density of WIMPs make their detection possible

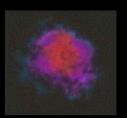
Cosmic density of thermal WIMPs

• At early times, WIMPs are produced in e^+e^- , $\mu^+\mu^-$, etc collisions in the hot primordial soup [thermal production].

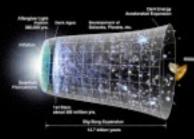
- WIMP production ceases when the production rate becomes smaller than the Hubble expansion rate [freeze-out].
- After freeze-out, there is a constant number of WIMPs in a volume expanding with the universe.

Indirect detection

the WIMP



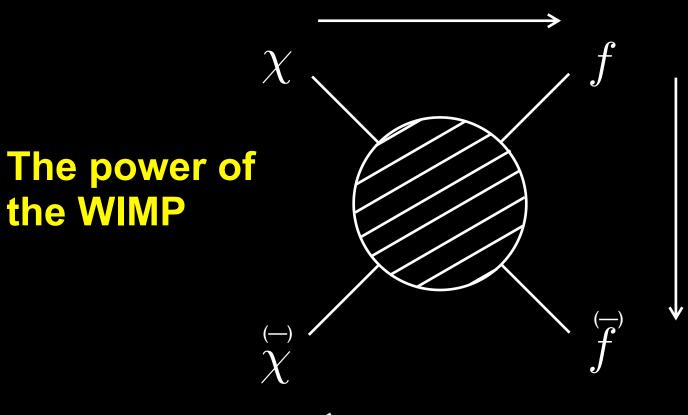




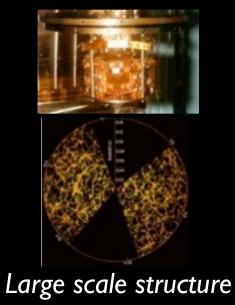
Cosmic density

Scattering

Annihilation



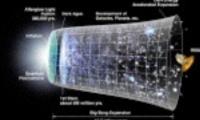
Direct detection



Production

Colliders

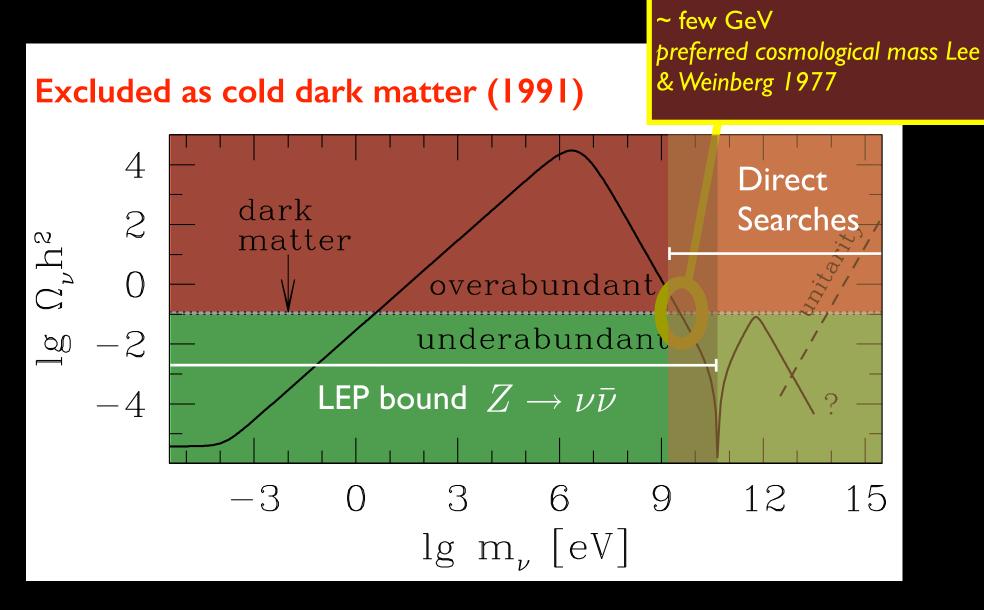




Cosmic density

Massive neutrinos as dark matter

Heavy active neutrino



Sterile neutrino dark matter

Standard model + right-handed neutrinos

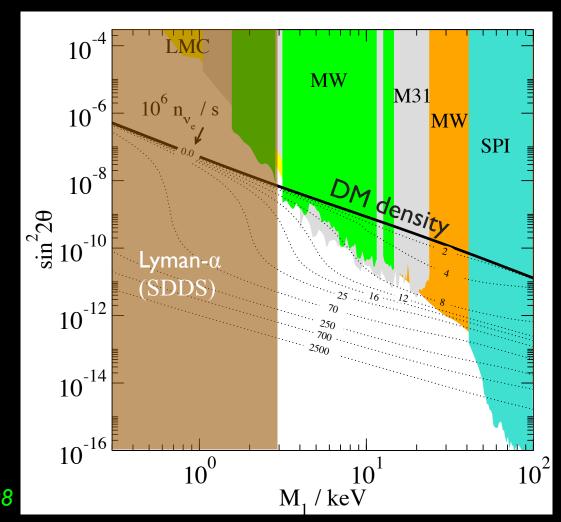
Active and sterile neutrinos oscillate into each other.

Sterile neutrinos can be warm dark matter (mass > 0.3 keV)

Dodelson, Widrow 1994; Shi, Fuller 1999; Laine, Shaposhnikov 2008

vMSM

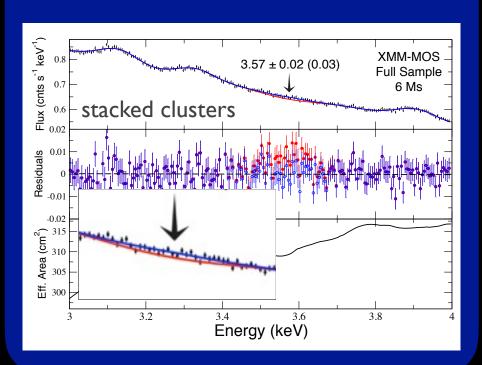
Laine, Shaposhnikov 2008



Sterile neutrino dark matter

An unidentified 3.5-keV X-ray line has been reported in galaxy clusters and the Andromeda galaxy.

Bulbul et al 2014; Boyarski et al 2014; lakubovskyi et al 2015

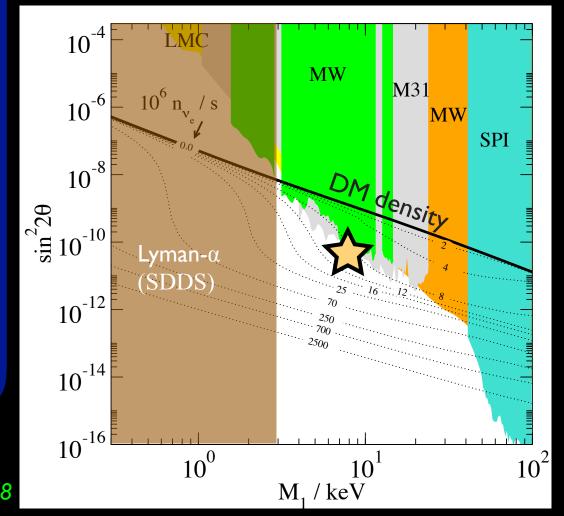


vMSM

Laine, Shaposhnikov 2008

Radiative decay of sterile neutrinos

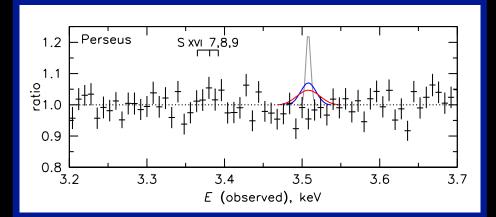
$$u_s o \gamma \nu_a \qquad E_\gamma = m_s/2$$
 $m_v = 7.1 \text{ keV} \qquad \sin^2(2\theta) = 7 \times 10^{-11}$



Sterile neutrino dark matter

The HITOMI data on the Perseus galaxy cluster do not show an X-ray line with the expected strength

Aharonian et al (HITOMI Collab.) 2016

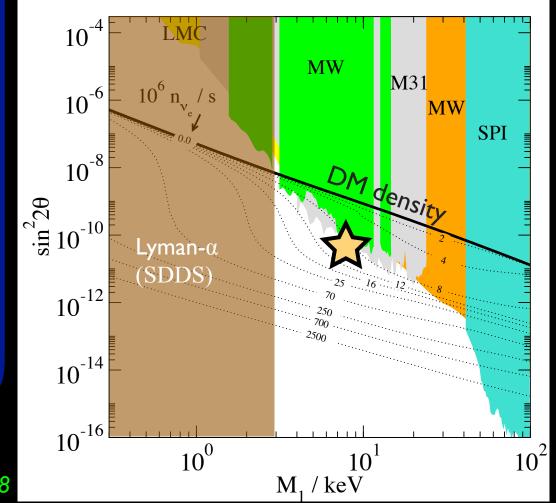


vMSM

Laine, Shaposhnikov 2008

Radiative decay of sterile neutrinos

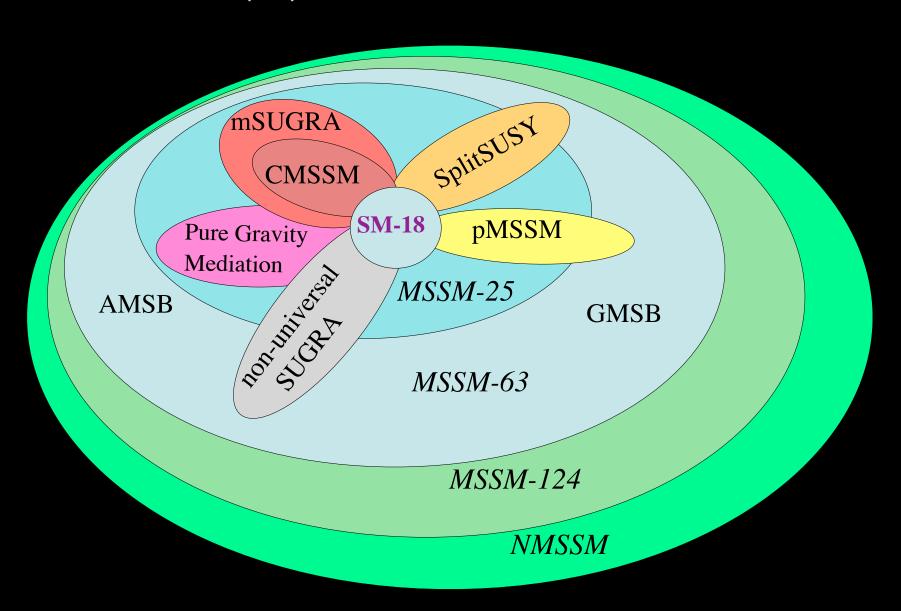
$$u_s o \gamma
u_a \qquad E_\gamma = m_s/2$$
 $m_v = 7.1 \text{ keV} \qquad \sin^2(2\theta) = 7 \times 10^{-11}$



Supersymmetric dark matter

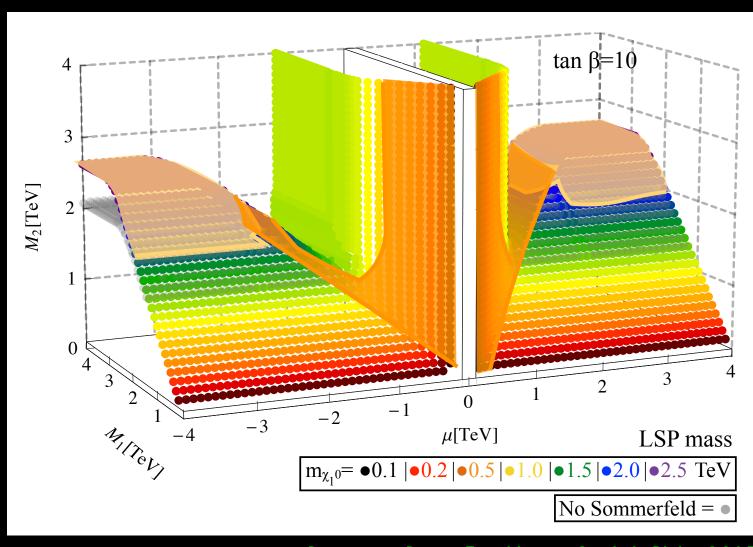
The CMSSM* is in dire straights, but there are many supersymmetric models

*Constrained Minimal Supersymmetric Standard Model



Neutralino dark matter

Neutralino dark matter with decoupled (heavy) sfermions



Excluded by LEP, HESS, LUX

All can be tested by LZ, CTA, and a I 00-TeV pp collider

Bramante, Desai, Fox, Martin, Ostdiek, Plehn 2015

Scalar phantom dark matter

"Gauge singlet scalar dark matter"

"Singlet scalar dark matter"

"Scalar singlet dark matter"

"Scalar Higgs-portal dark matter"

"The minimal model of dark matter"

Minimalist dark matter

do not confuse with minimal dark matter

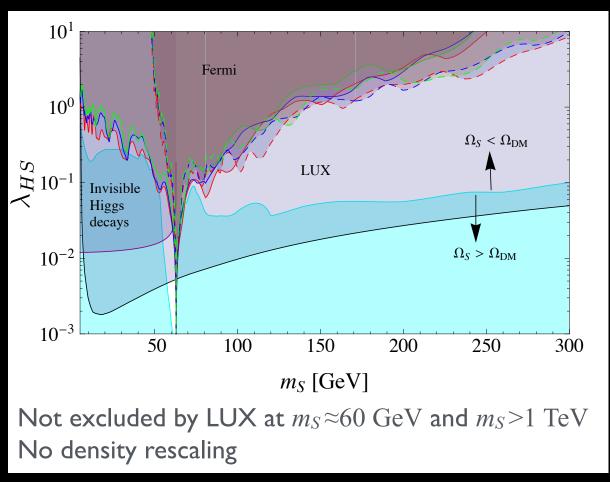
Gauge singlet scalar field S stabilized by a \mathbb{Z}_2 symmetry $(S \rightarrow -S)$

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} S \partial_{\mu} S + \frac{1}{2} \mu_S^2 S^2 - \frac{\lambda_S}{4} S^4 - \lambda_{HS} H^{\dagger} H S^2$$

Silveira, Zee 1985 Andreas, Hambye, Tytgat 2008 Djouadi, Falkowksi, Mambrini, Quevillon 2012 Cline, Scott, Kainulainen, Weniger 2013 Hamada, Kawana 2015

"Scalar phantom" is the original 1985 name

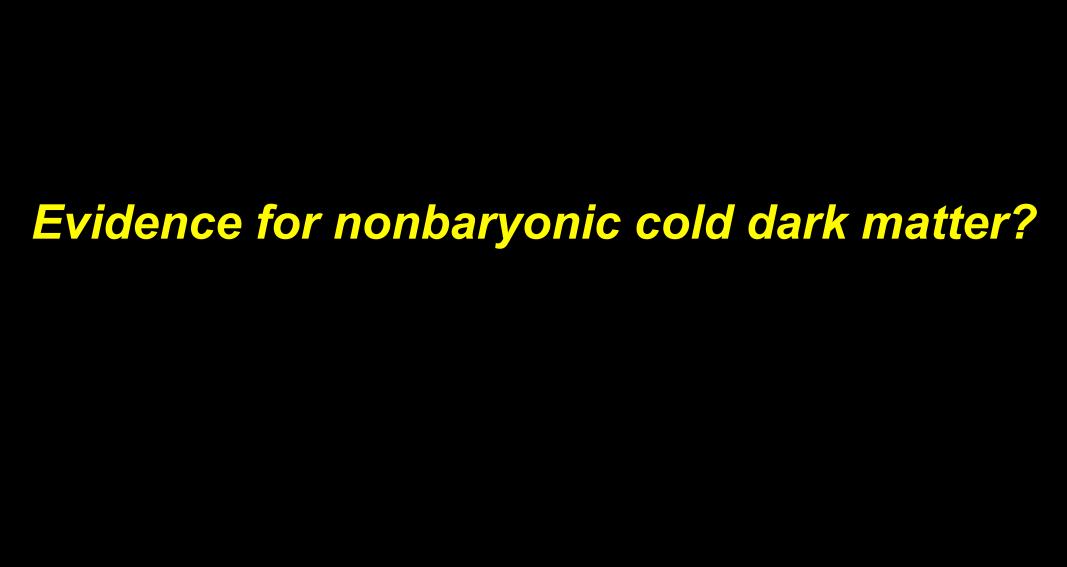
Scalar phantom dark matter



Feng, Profumo, Ubaldi 2015

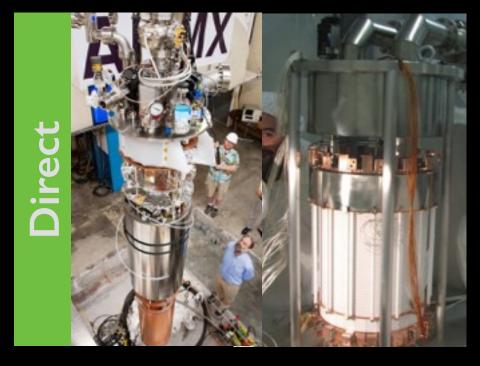
If density is rescaled according to Ω_S , LUX and FERMI exclusion regions are very different

Cline, Scott, Kainulainen, Weniger 2013

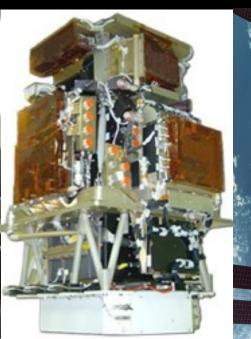


Searches for particle dark matter











Indirect

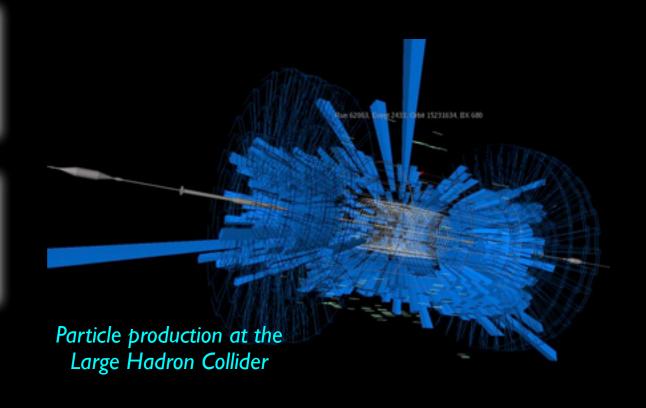
Searches for particle dark matter: colliders

Dark matter particles (or their "cousins") are produced in high-energy collisions

Dark matter particles are produced and escape detection (missing energy)

Charged/colored "cousins" of the dark matter particle are produced

LEP ALEPH, DELPHI, OPAL, ... Tevatron CDF, D0, ... LHC ATLAS, CMS, ...

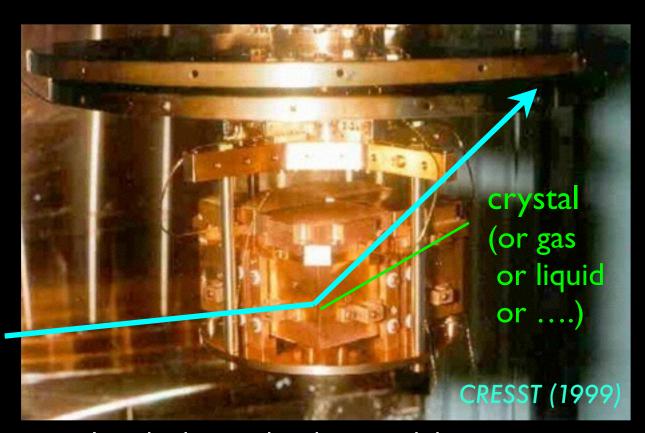


Searches for particle dark matter: direct

Dark matter particles that arrive on Earth scatter off nuclei or electrons in a detector

Goodman, Witten 1985

Dark matter particle

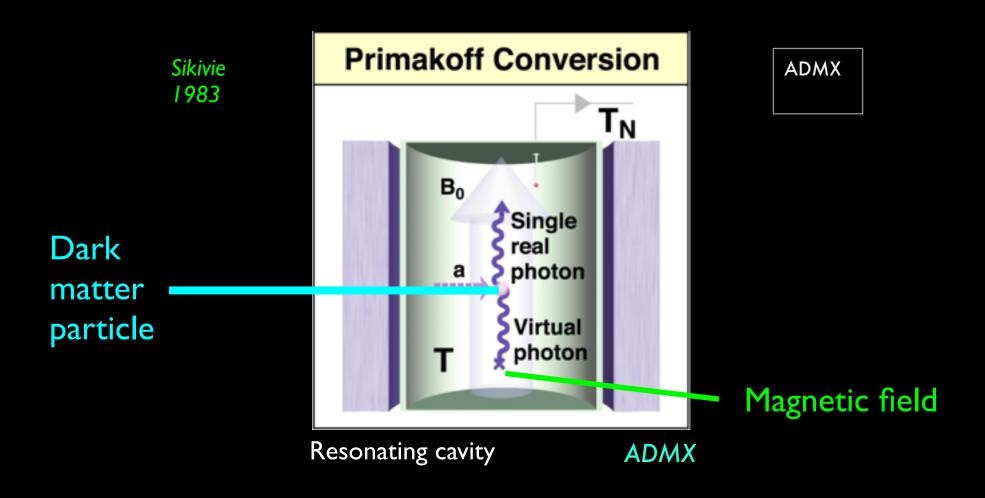


Low-background underground detector

DAMA, SuperCDMS, XENON, LUX, XMASS, PICO, CoGeNT, DEAP, DRIFT, ANAIS, CRESST, LZ, DARWIN, DM-ICE, NEWAGE, ...

Searches for particle dark matter: direct

Dark matter particles that arrive on Earth transform into photons in a detector

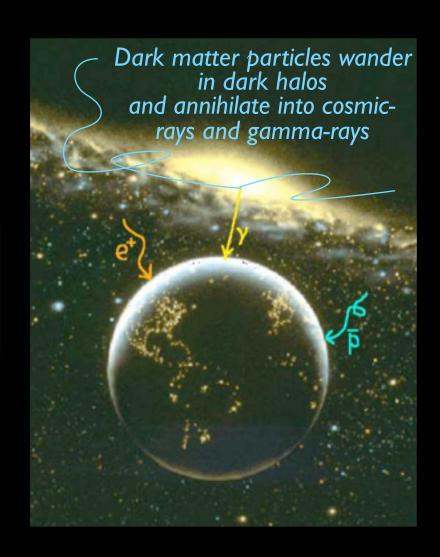


Searches for particle dark matter: indirect

Dark matter particles transform into ordinary particles, which are then detected or inferred

Gunn, Lee, Lerche, Schramm, Steigman 1978; Stecker 1978

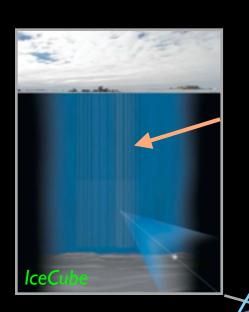
Gamma-rays, positrons, antiprotons from our galaxy and beyond



cosmic-rays
PAMELA
AMS
...
gamma-rays
MAGIC
HESS
VERITAS
Fermi-LAT
HAWK
CTA
...

Searches for particle dark matter: indirect

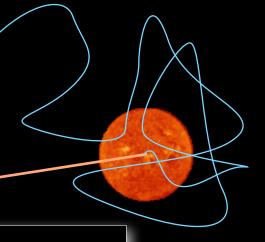
Dark matter particles transform into ordinary particles, which are then detected or inferred



IceCube ANTARES

• • •

Dark matter particles sink into the Sun/Earth where they transform into neutrinos



Neutrinos from the Sun

Press, Spergel 1985; Silk, Olive, Srednicki 1985

Neutrinos from the Earth

Freese 1986; Krauss, Srednicki, Wilczek 1986

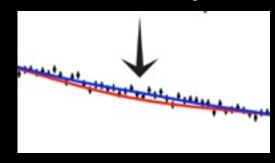
Signals from dark matter?

GeV γ -rays



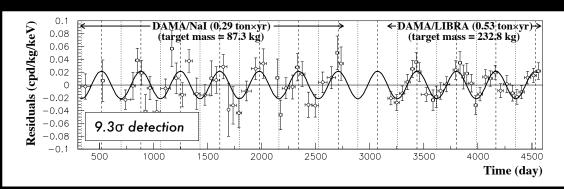
Hooper et al 2009-14

3.5 keV X-ray line



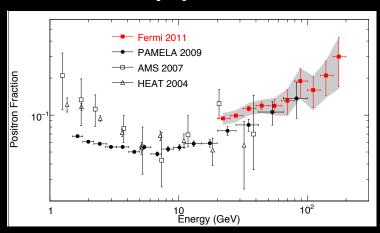
Bulbul et al 2014

Annual modulation in direct detection



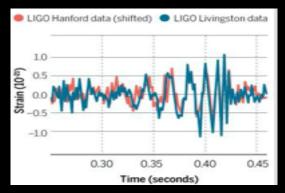
Bernabei et al 1997-now

Cosmic-ray positrons



Adriani et al 2009; Ackerman et al 2011; Aguilar et al 2013

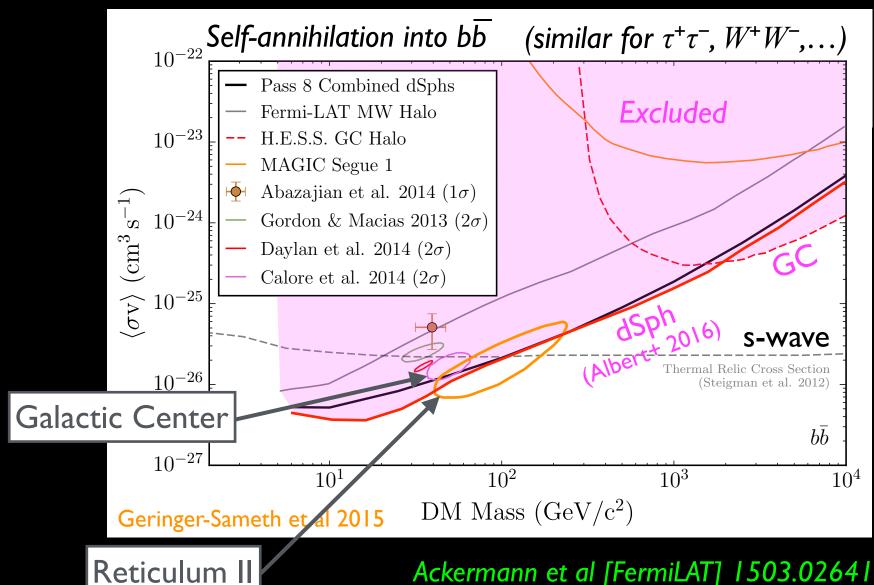
Black-hole mergers



Bird et al, Kashlinsky 2016

Gamma-rays from dark matter

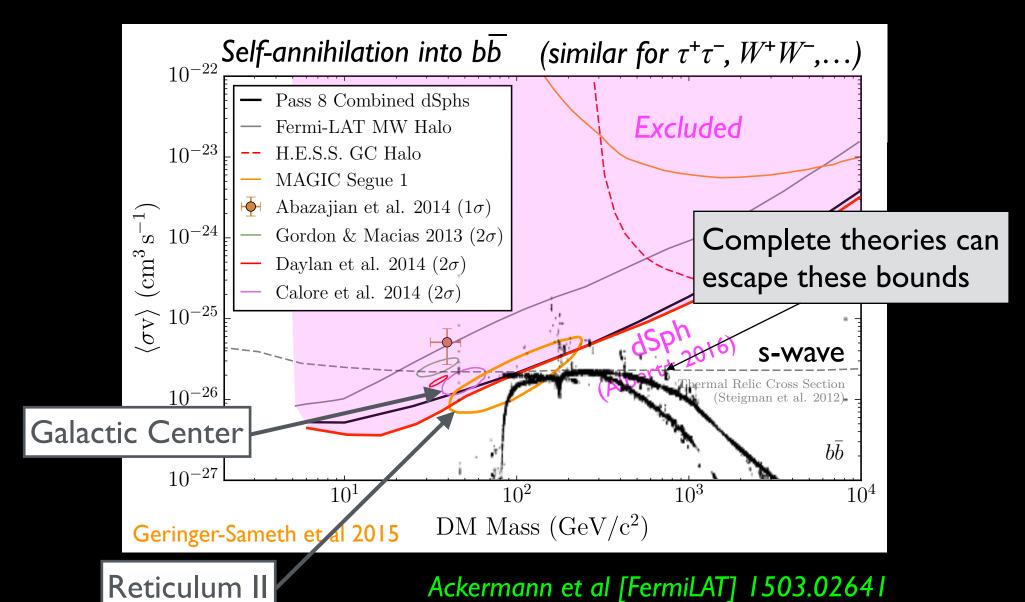
Upper limits on the WIMP annihilation cross section from dwarf spheroidal galaxies and Galactic Center



Ackermann et al [FermiLAT] 1503.02641

Gamma-rays from dark matter

Upper limits on the WIMP annihilation cross section from dwarf spheroidal galaxies and Galactic Center

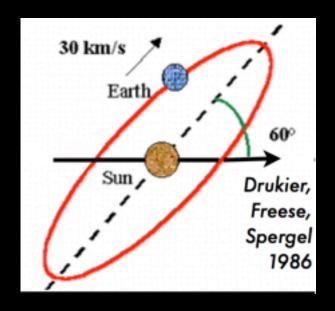


Recent trends: "make no assumptions"

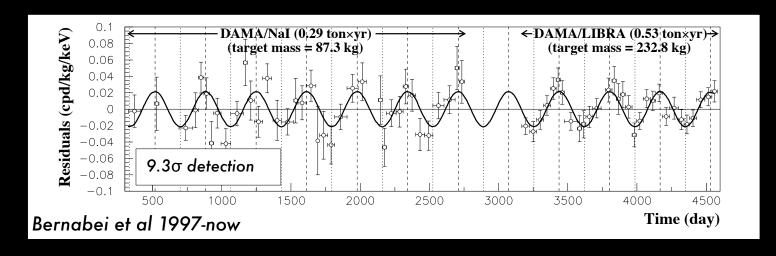
Annual modulation in direct detection

 The revolution of the Earth around the Sun modulates the WIMP event rate

Drukier, Freese, Spergel 1986



DAMA observes such kind of modulation



DAMA annual modulation

Model Independent Annual Modulation Result

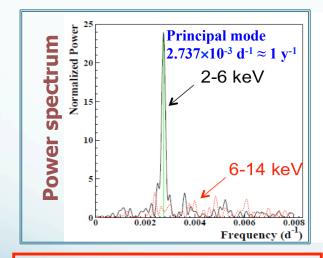
DAMA/NaI + DAMA/LIBRA-phase1 Total exposure: 487526 kg×day = 1.33 ton×yr

EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648

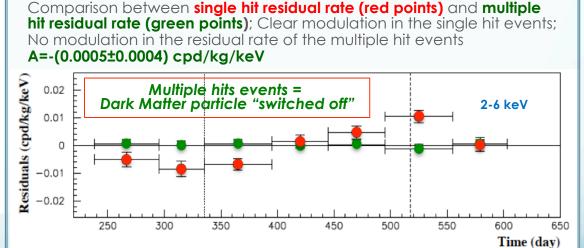
The measured modulation amplitudes (A), period (T) and phase (t_0) from the single-hit residual rate vs time

	A(cpd/kg/keV)	$T=2\pi/\omega$ (yr)	t ₀ (day)	C.L.
DAMA/NaI+DAMA/LIBRA-phase1				
(2-4) keV	0.0190 ±0.0020	0.996 ±0.0002	134 ± 6	9.5σ
(2-5) keV	0.0140 ±0.0015	0.996 ±0.0002	140 ± 6	9.3σ
(2-6) keV	0.0112 ±0.0012	0.998 ±0.0002	144 ± 7	9.3σ

 $A\cos[\omega(t-t_0)]$



No systematics or side reaction able to account for the measured modulation amplitude and to satisfy all the peculiarities of the signature



This result offers an additional strong support for the presence of DM particles in the galactic halo further excluding any side effect either from hardware or from software procedures or from background

The data favor the presence of a modulated behaviour with all the proper features for DM particles in the galactic halo at about 9.2 σ C.L.

DAMA annual modulation

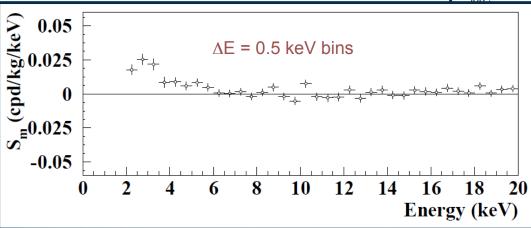
Model Independent Annual Modulation Result

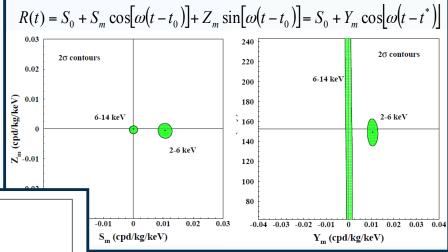
DAMA/NaI + DAMA/LIBRA-phase1 Total exposure: 487526 kg×day = 1.33 ton×yr

EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648

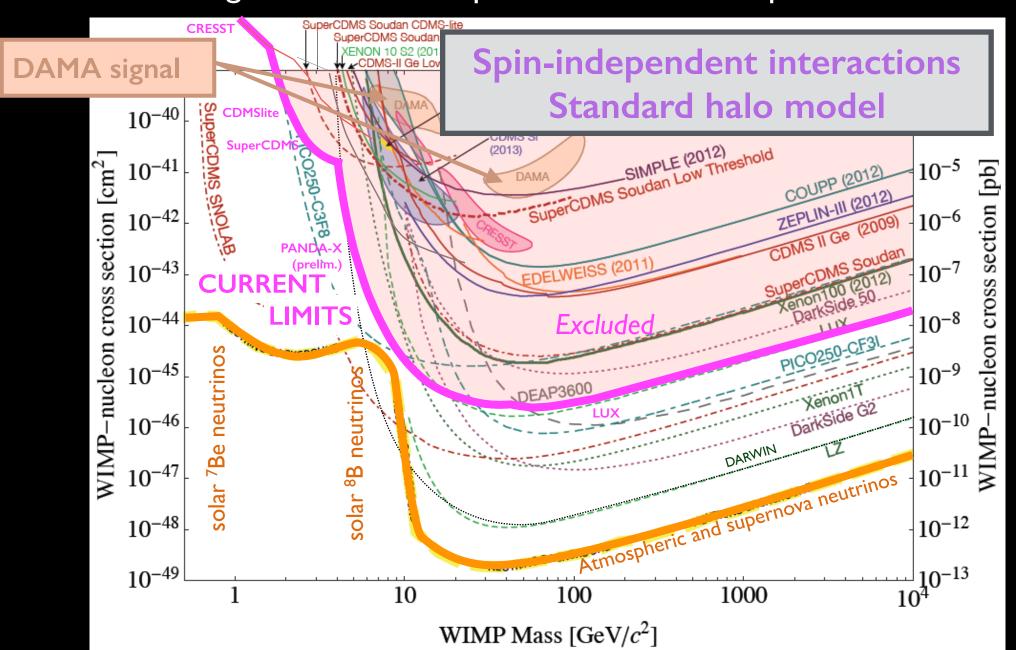
- No modulation above 6 keV
- No modulation in the whole energy spectrum
- No modulation in the 2-6 keV multiple-hit events

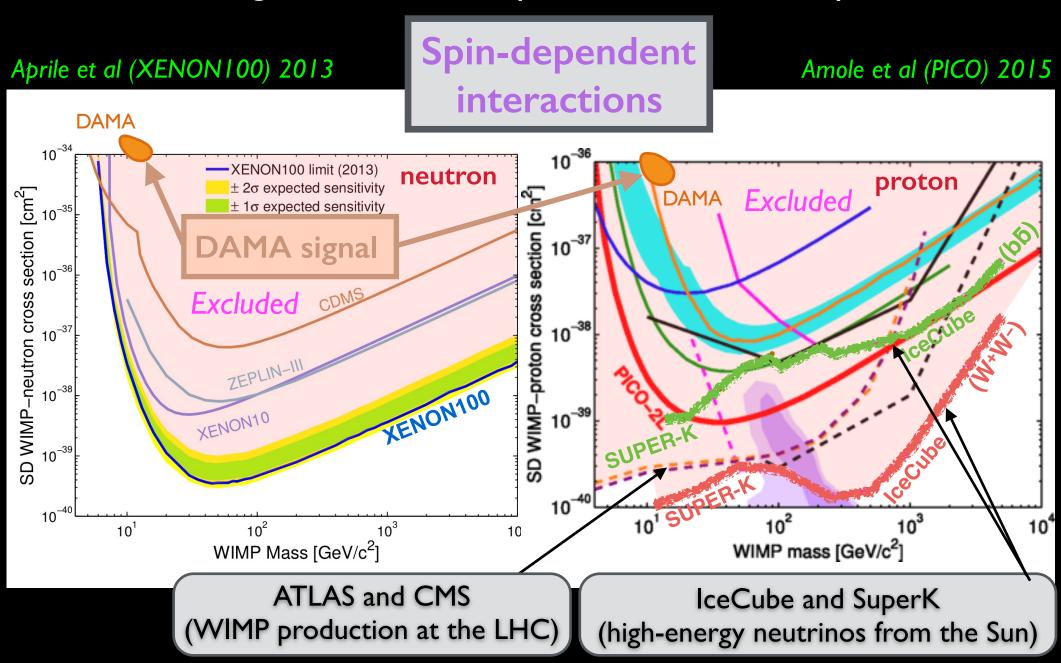
$$R(t) = S_0 + S_m \cos\left[\omega(t - t_0)\right]$$
here $T = 2\pi/\omega = 1$ yr and $t_0 = 152.5$ day

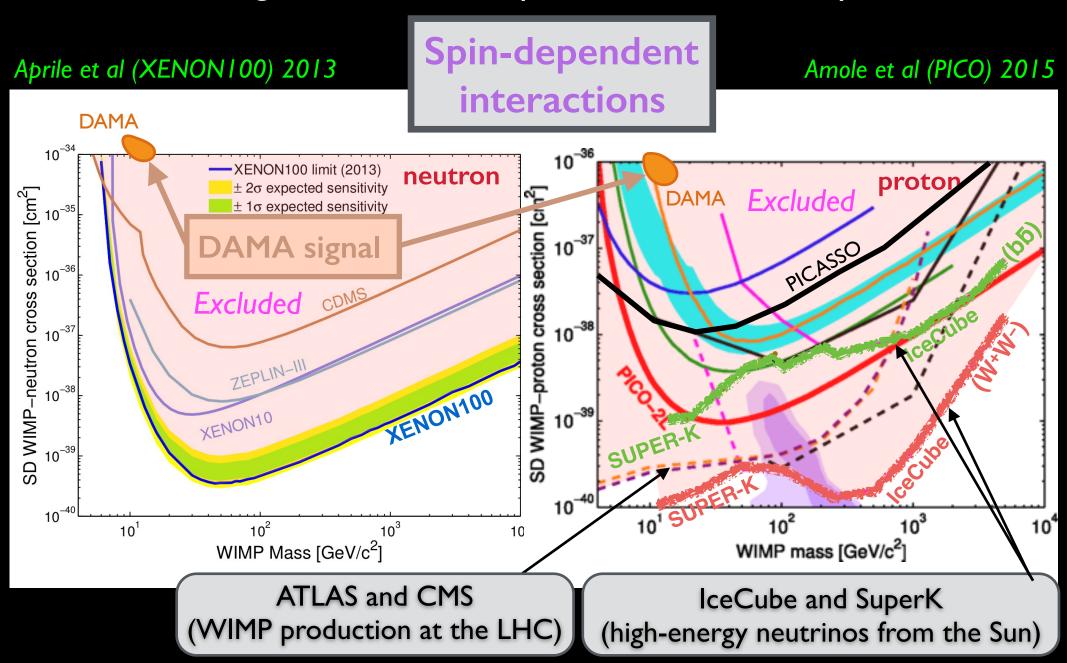


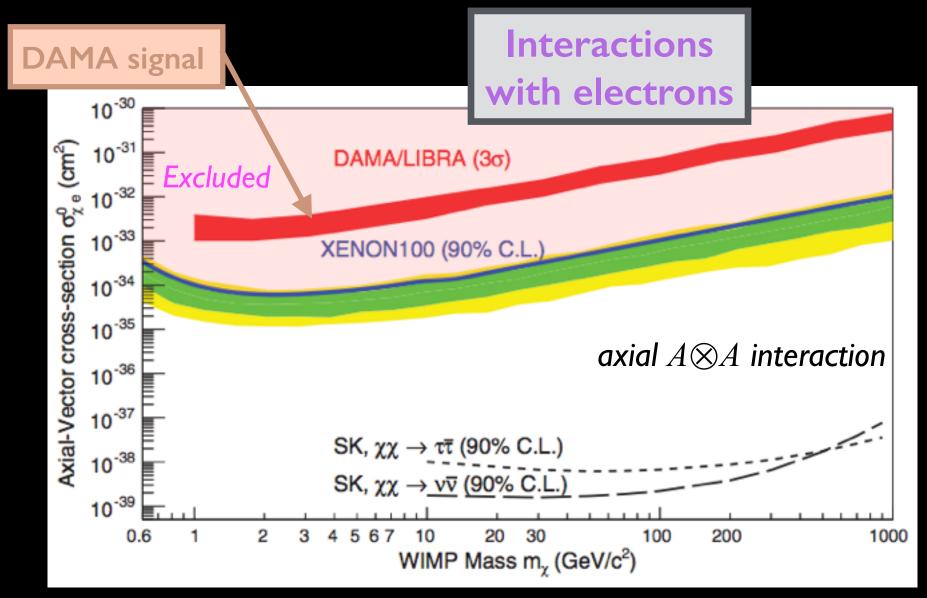


No systematics or side processes able to quantitatively account for the measured modulation amplitude and to simultaneously satisfy the many peculiarities of the signature are available.









Recoil spectrum

The recoil spectrum (scattering rate per unit target mass)

$$\frac{dR}{dE_R} = \frac{1}{m_T} \int_{v > v_{\min}(E_R)} n_{\chi} v \frac{d\sigma}{dE_R} f(\mathbf{v}) d^3 v$$

Traditionally, $v^2 d\sigma/dE_R = \text{const} \times (\text{nuclear form factor})$

In trying to explain the data, modify the cross section

- set different couplings to neutrons and protons ("isospin-violating")
- put additional velocity or energy dependence in $v^2 d\sigma/dE_R$

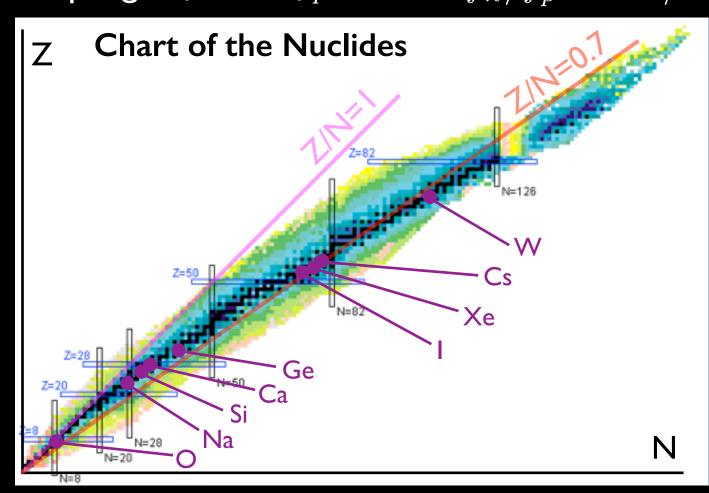
Isospin-violating (nonisoscalar) dark matter

Spin-independent couplings to protons stronger than to neutrons may allow modulation signals compatible with other null searches

Kurylov, Kamionkowski 2003; Giuliani 2005; Cotta et al 2009; Chang et al 2010; Kang et al 2010; Feng et al 2011; Del Nobile et al 2011;

coupling
$$Nf_n + Zf_p \approx 0$$
 for $f_n/f_p \approx -Z/N$

Why $f_n/f_p = -0.7$ suppresses the coupling to Xe



Velocity and energy-dependent scattering

Energy and/or velocity dependent scattering cross sections

nucleus	DM	$v^2 d\sigma/dE_R$		
Hucicus		light mediator	heavy mediator	
"charge"	"charge"	$1/E_R$	$1/M^4$	
"charge"	dipole	$1/E_R$	E_R/M^4	
dipole	dipole	$const + E_R/v^2$	E_R^2/M^4	

All terms may be multiplied by nuclear or DM form factors $F(E_R)$

See e.g. Barger, Keung, Marfatia 2010; Fornengo, Panci, Regis 2011; An et al 2011

"Make no assumptions"

All particle physics models

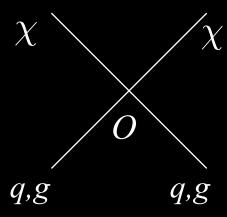
- Consider all possible interactions between dark matter and standard model particles
- This program has been carried out in some limits (e.g., non-relativistic conditions, heavy mediators)

All astrophysical models

- Halo-independent methods of analysis have been developed
- Ideally they require no assumption on the astrophysical density and velocity distributions of dark matter particles

Contact operators

if mediator mass ≫ exchanged energy



Four-particle contact operator

There are many possible operators.

Interference is important although often, but not always, neglected.

Long(ish) distance interactions are not included.

Nonrelativistic contact operators

Nonrelativistic WIMP-nucleon contact operators classified

Fitzpatrick et al 2012, Dent et al 2015

$$\mathcal{O}_{1} = 1_{\chi} 1_{N} \qquad \qquad \mathcal{O}_{7} = \vec{S}_{N} \cdot \vec{v}_{\chi N}^{\perp}$$

$$\mathcal{O}_{3} = -i \vec{S}_{N} \cdot \left(\frac{\vec{q}}{m_{N}} \times \vec{v}_{\chi N}^{\perp}\right) \qquad \mathcal{O}_{8} = \vec{S}_{\chi} \cdot \vec{v}_{\chi N}^{\perp}$$

$$\mathcal{O}_{4} = \vec{S}_{\chi} \cdot \vec{S}_{N} \qquad \qquad \mathcal{O}_{9} = -i \vec{S}_{\chi} \cdot \left(\vec{S}_{N} \times \frac{\vec{q}}{m_{N}}\right)$$

$$\mathcal{O}_{5} = -i \vec{S}_{\chi} \cdot \left(\frac{\vec{q}}{m_{N}} \times \vec{v}_{\chi N}^{\perp}\right) \qquad \qquad \mathcal{O}_{10} = -i \vec{S}_{N} \cdot \frac{\vec{q}}{m_{N}}$$

$$\mathcal{O}_{6} = \left(\vec{S}_{\chi} \cdot \frac{\vec{q}}{m_{N}}\right) \left(\vec{S}_{N} \cdot \frac{\vec{q}}{m_{N}}\right) \qquad \qquad \mathcal{O}_{11} = -i \vec{S}_{\chi} \cdot \frac{\vec{q}}{m_{N}}$$

and more in Dent et al 2015

At leading order in q and v, only O_1 and O_4 appear, which are the spin-independent and spin-dependent terms, respectively.

Nonrelativistic contact operators

Operators appearing in "simplified" WIMP models Dent et al 2015

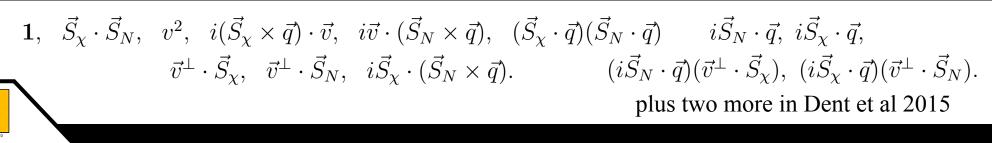
	$s_{\chi}=0$ X^{0}	s _χ =0 X ⁺	$ \begin{array}{c c} s_{\chi}=1/2 \\ X^{0} \end{array} $	$s_{\chi}=1/2$ X^{+}	$s_{\chi}=1$ X^{0}	$s_{\chi}=1$ X^{+}
S	1	1	1,11	1,11	1	1
V	_		1,8,9	1,8,9	4,5,6,8,9,11,17	
T		_	_	4,10,11,12		4,11,12,18
A	10		4,7,9	4,7,9	4,6,9,10,14,18	
P	10	10	6,10	6,10	10	10

Entries denote operator index. X^0 (X^+) is neutral (charged) mediator.

Contact operators: direct detection

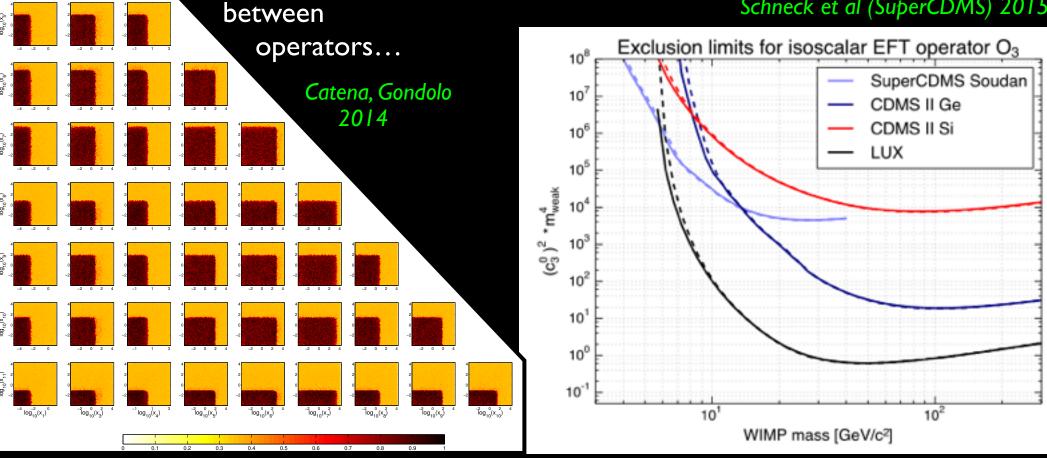
Short-distance WIMP-nucleon operators classified

Fitzpatrick et al 2012

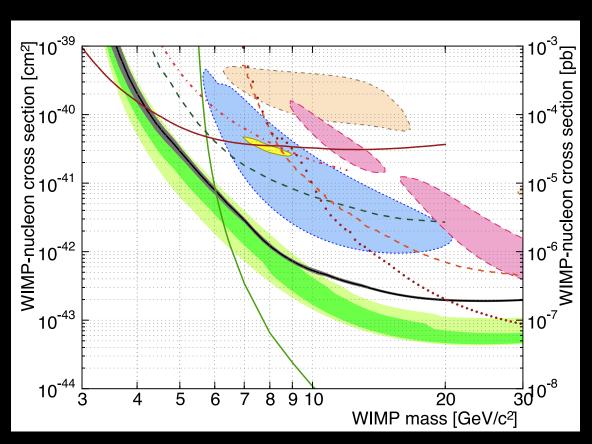


Experimental limits on single operators... **Correlations**

Schneck et al (SuperCDMS) 2015



$$\begin{pmatrix} \text{event} \\ \text{rate} \end{pmatrix} = \begin{pmatrix} \text{detector} \\ \text{response} \end{pmatrix} \times \begin{pmatrix} \text{particle} \\ \text{physics} \end{pmatrix} \times \begin{pmatrix} \text{astrophysics} \end{pmatrix}$$
FIXED



Standard Halo Model

truncated Maxwellian

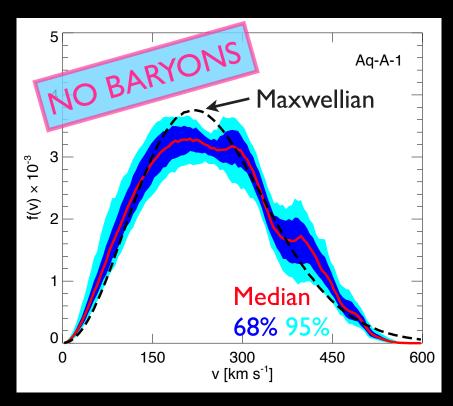
$$f(\vec{v}) = C e^{-|\vec{v} + \vec{v}_{obs}|/\bar{v}_0^2} \Theta(v - v_{esc})$$



The spherical cow of direct WIMP searches

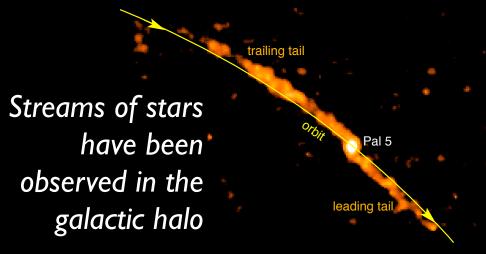
Gelmini

We know very little about the dark matter velocity distribution near the Sun

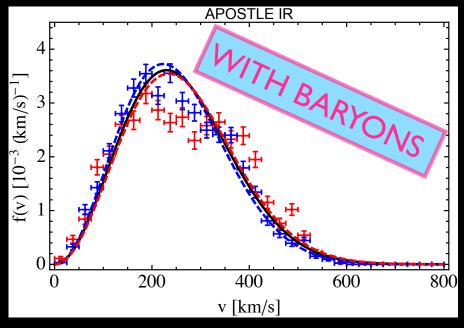


Vogelsberger et al 2009

Cosmological N-Body simulations including baryons are challenging but underway



Odenkirchen et al 2002 (SDSS)



Bozorgnia et al 2016

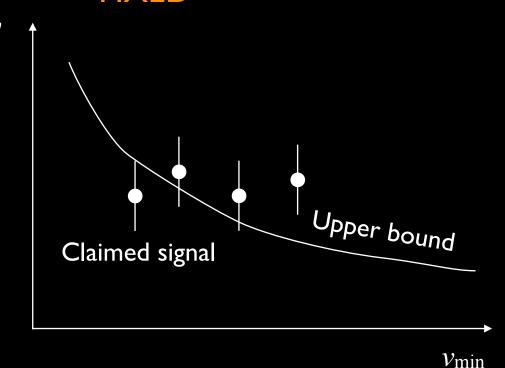
$$\begin{pmatrix} \text{event} \\ \text{rate} \end{pmatrix} = \begin{pmatrix} \text{detector} \\ \text{response} \end{pmatrix} \times \begin{pmatrix} \text{particle} \\ \text{physics} \end{pmatrix} \times \begin{pmatrix} \text{astrophysics} \end{pmatrix}$$

FIXED ARBITRARY

Rescaled astrophysics factor common to all experiments

$$\eta(v_{\min}) = \frac{\rho_{\chi}}{m_{\chi}} \int_{v_{\min}}^{\infty} \frac{f(\mathbf{v})}{v} d^3v$$

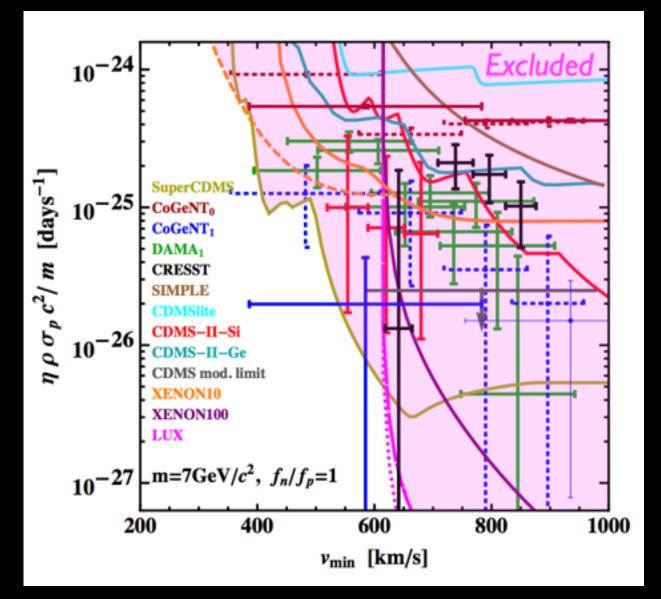
Proxy for dark matter flux



Minimum WIMP speed to impart recoil energy E_R

Spin-independent isoscalar interactions

$$\frac{d\sigma}{dE_R} = \frac{2m}{\pi v^2} A^2 f_p^2 F^2(E_R)$$



Halo modifications alone cannot save the SI signal regions from the Xe and Ge bounds

CDMS-Si event rate is similar to yearly modulated rates

Still depends on particle model

Del Nobile, Gelmini, Gondolo, Huh 2014

Anapole dark matter

The anapole moment is a C and P violating, but CP-conserving, electromagnetic moment

First measured experimentally in Cesium atoms

Zeldovich 1957

Wood et al 1997

Anapole dark matter

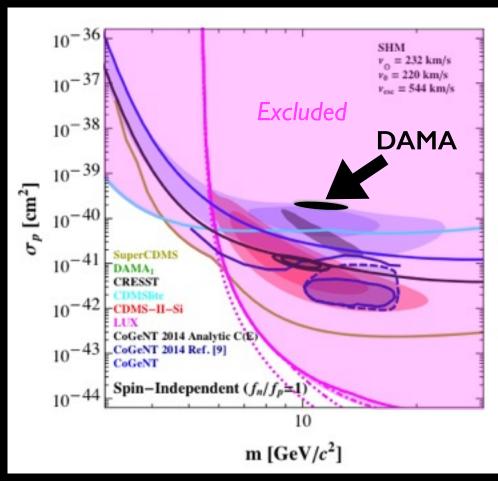
spin-1/2 Majorana fermion

$$\mathcal{L} = \frac{g}{2\Lambda^2} \bar{\chi} \gamma^{\mu} \gamma^5 \chi \, \partial^{\nu} F_{\mu\nu}$$

$$H = -\frac{g}{\Lambda^2} \, \vec{\sigma} \cdot \vec{\nabla} \times \vec{B}$$

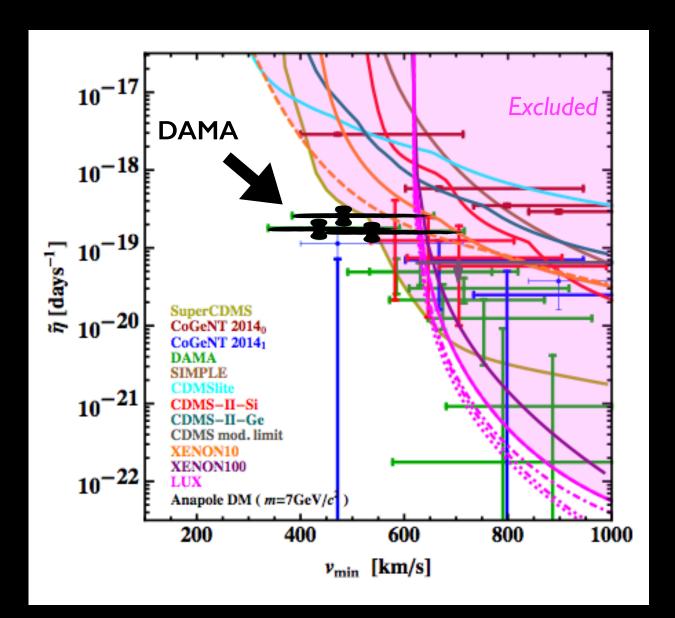
Direct detection limits with standard dark halo

Del Nobile, Gelmini, Gondolo, Huh 2014



Anapole dark matter

$$\frac{d\sigma}{dE_R} = \frac{2m}{\pi v^2} \frac{e^2 g^2}{\Lambda^2} \left[\left(v^2 - v_{\min}^2 \right) F_L^2(E_R) + F_T^2(E_R) \right]$$



spin-1/2 Majorana fermion

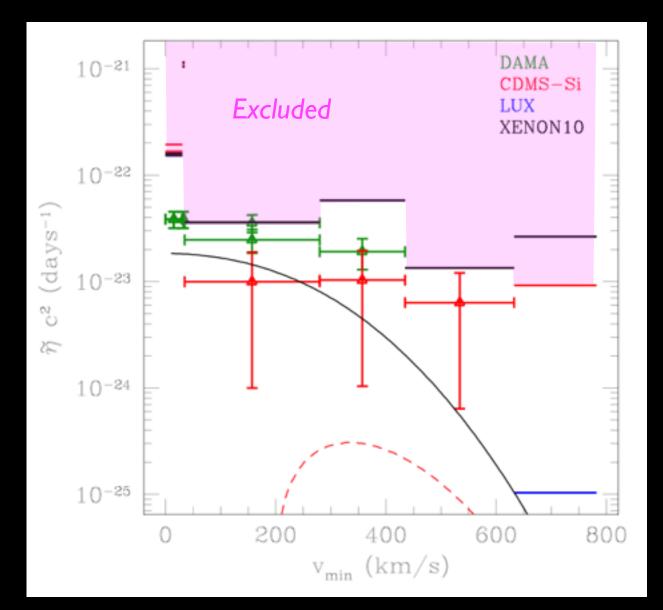
$$\mathcal{L} = \frac{g}{2\Lambda^2} \bar{\chi} \gamma^{\mu} \gamma^5 \chi \, \partial^{\nu} F_{\mu\nu}$$

$$H = -\frac{g}{\Lambda^2} \, \vec{\sigma} \cdot \vec{\nabla} \times \vec{B}$$

For anapole dark matter, the lowest DAMA bins may be compatible with null searches

Exothermic nonisoscalar scattering

$$\frac{d\sigma}{dE_R} = \frac{2m}{\pi v^2} [Z f_p + (A - Z) f_n]^2 F^2(E_R)$$



For light exothermic nonisoscalar scattering, the DAMA modulation may be compatible with other experiments

$$m = 3 \text{ GeV}/c^2$$

$$\delta = -70 \text{ keV}$$

$$f_n/f_p = -0.79$$

$$(A, Z)$$

$$\chi'(m + \delta)$$

$$(A, Z)$$

Still depends on particle model

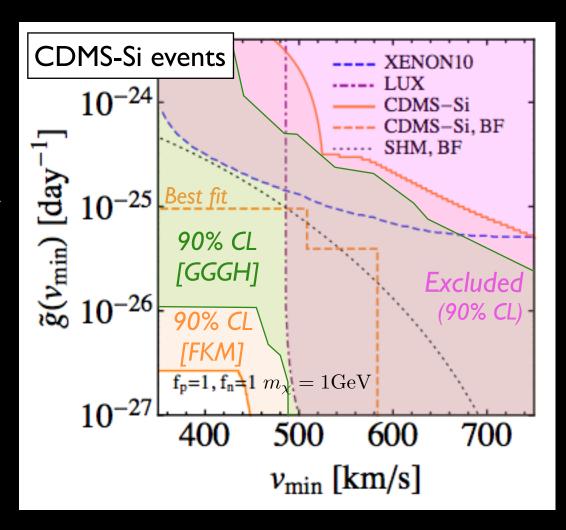
The statistics of the halo-independent approach is being understood.

Unbinned likelihood analysis

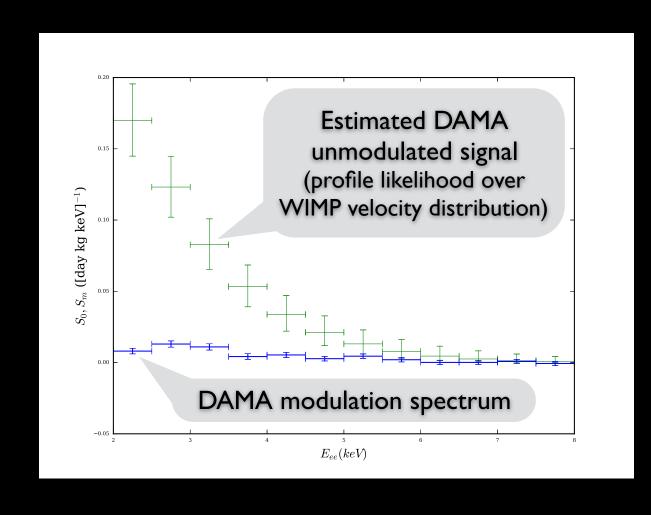
$$\mathcal{L} = \frac{e^{-\int_{E_{\min}}^{E_{\max}} \frac{dR}{dE} dE}}{N!} \prod_{i=1}^{N} \frac{dR}{dE} \Big|_{E=E_i}$$

The extent of the 90% CL region is still unclear

Fox, Kahn, McCullough 2015 Gelmini, Georgescu, Gondolo, Huh 2015



New techniques and proper statistical treatment let the astrophysics-independent approach address questions beyond the comparison of experiments.



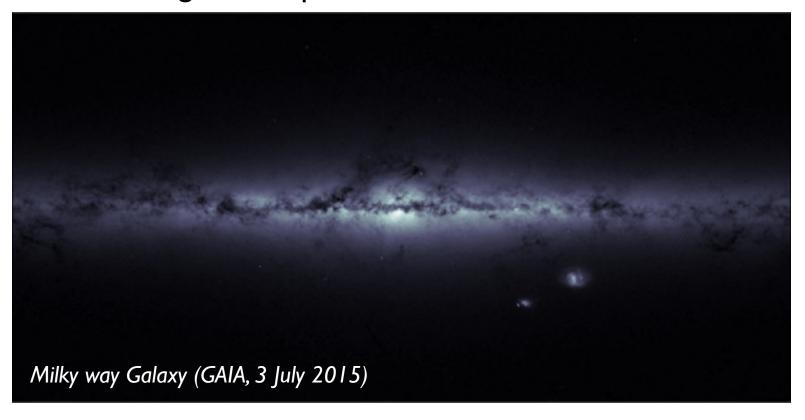
Astrophysics-independent estimate of the DAMA unmodulated signal

Gondolo, Scopel 2016 (in preparation)

In the next future

In the next future..... Small-scale structure

GAIA is measuring the 3D position of about one billion stars.

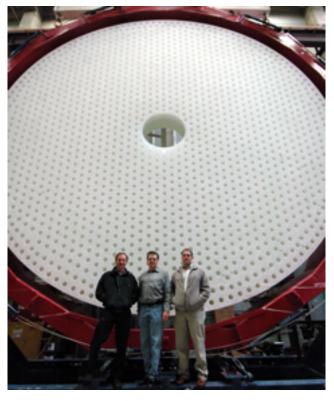


GAIA may detect dark matter substructures in the galactic halo. Feldmann, Spolyar 2014

In the next future..... Large-scale structure

The Dark Energy Survey (DES) and the Large Synoptic Survey Telescope (LSST) will map the large-scale distribution of dark matter and its growth with time.





DES

LSST

In the next future..... DAMA's revenge?

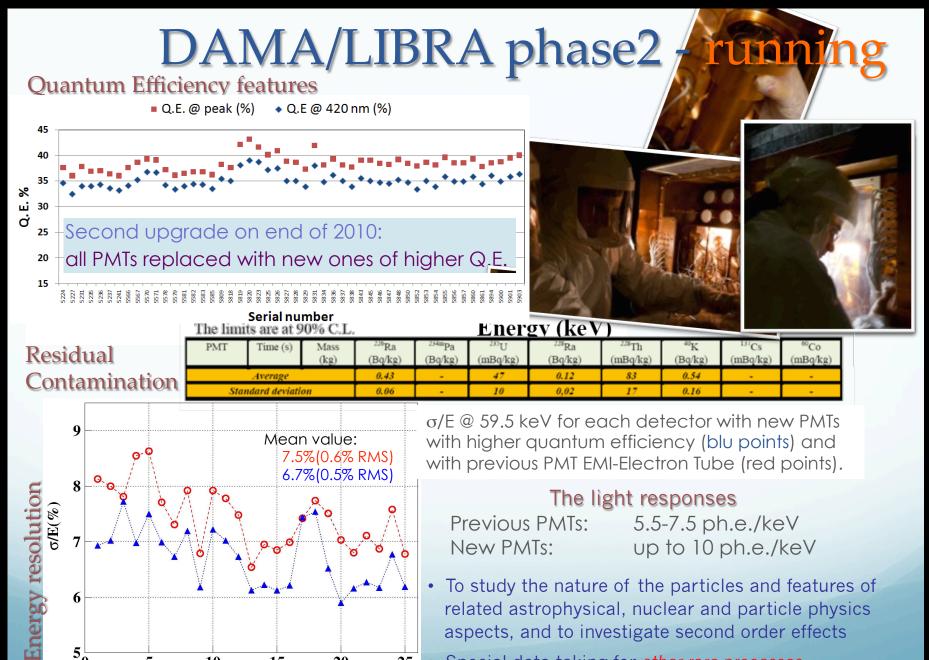
5

10

15

Detector number

20



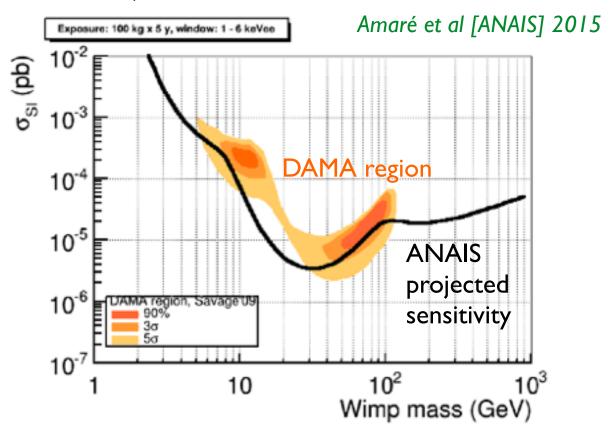
Special data taking for other rare processes

In the next future..... Direct check on DAMA

Experiments have been proposed that can directly check the DAMA modulation using the same target material

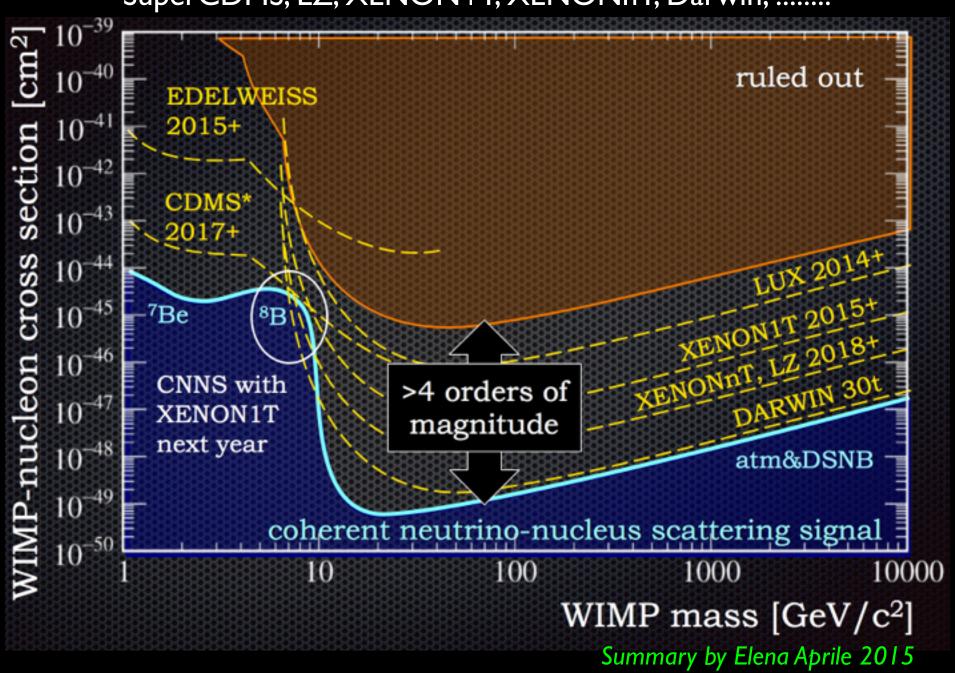
COSINE-100 (DM-ICE+KIMS-Nal)

ANAIS, SABRE, ...

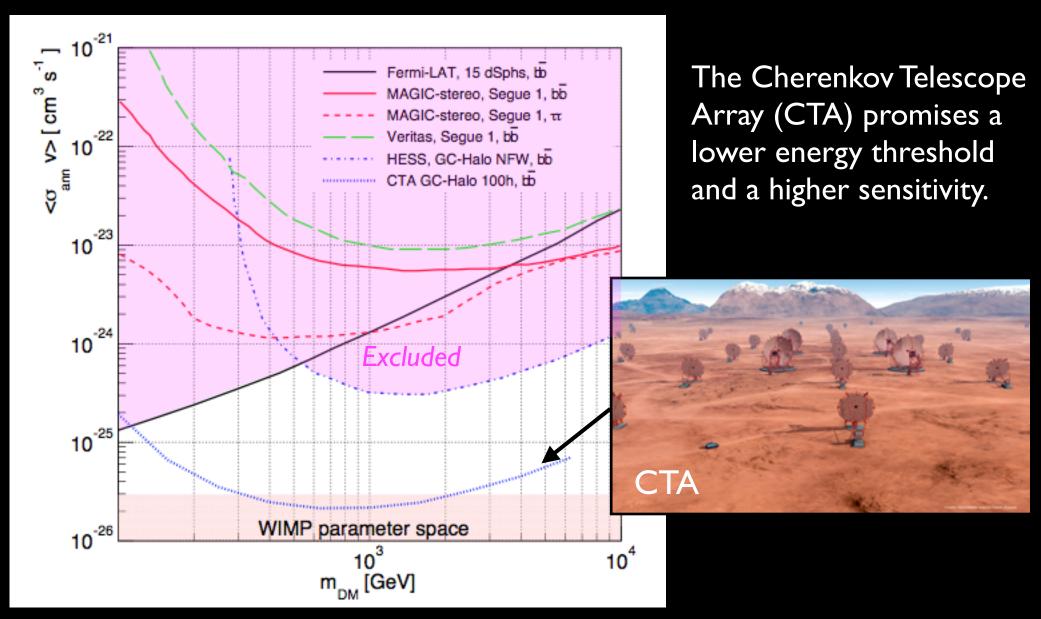


In the next future..... Giant direct detectors

SuperCDMS, LZ, XENON1T, XENONnT, Darwin,



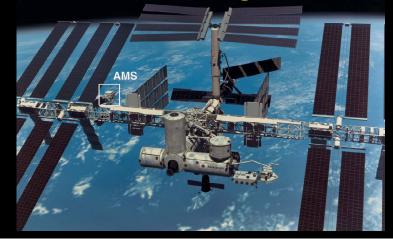
In the next future..... High-energy γ-rays

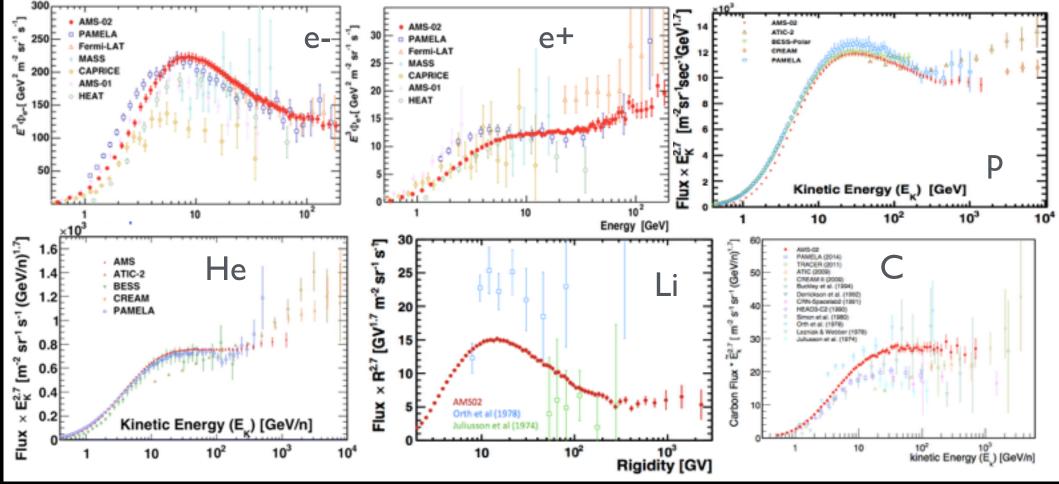


In the next future..... Precision cosmic rays

AMS (Alpha Magnetic Spectrometer)

Isotopic ratios measured to better than 1% precision up to Fe and ~100 GeV/nucleon allow for better Galactic cosmic ray models





Summary

- There is overwhelming astrophysical evidence for nonbaryonic cold dark matter
- The nature of cold dark matter is still unknown, and many candidates have been proposed
- There is some controversial detection of dark matter signals
- The next future will see measurements of the small- and large-scale structure of dark matter, as well as larger and more precise searches for dark matter signals