

Watching Galaxies Fall:

structure formation in the universe as a probe of gravity

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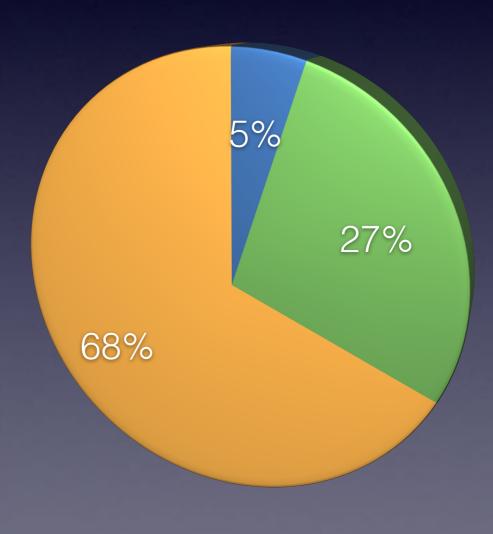
Summary

- Introduction
 - Modified gravity why change a good thing?
- Theory of Structure formation
 - Motion of matter, galaxies as test particles
 - Motion of light, photons as test particles
- Cosmological data
 - Redshift-Space distortions (WiggleZ, BOSS, 2dFLenS)
 - Weak gravitational lensing (CFHTLens, KiDS, DES)
- Beyond potentials fifth forces
 - Deterministic and stochastic velocity bias
- Conclusions

Testing gravity

- The expansion of the Universe is accelerating
- The simplest explanation of a cosmological constant is problematic
- Vacuum energy calculations imply cosmological constant is 10¹²⁰ times larger than its measured value - too small
- Coincidence problem why is density of matter (1/a³) so close to density of dark energy (~constant) today?

- Baryons
- Dark Matter
- Dark Energy



Cosmological Constant Problem

- Why is the energy density of the vacuum so small?
- Alternatively we can ask, why does the vacuum energy gravitate so little?
 - "The effective Newton constant becomes very small at large length scales, so that sources with immense wavelengths and periods -- such as the vacuum energy-- produce minuscule curvature" (Arkani-Hamed, Dimopoulos, Dvali, Gabadadze)
 - Similar to the manner in which long wavelength excitations beyond the Debye sphere are screened by the effective photon mass in a plasma.

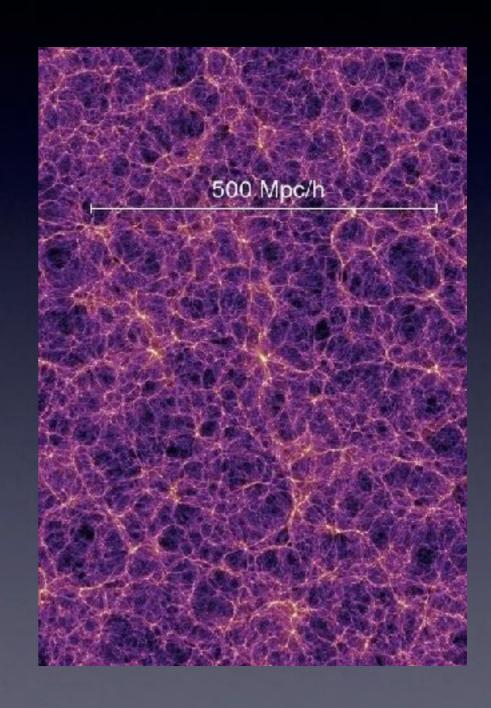
Modified Gravity

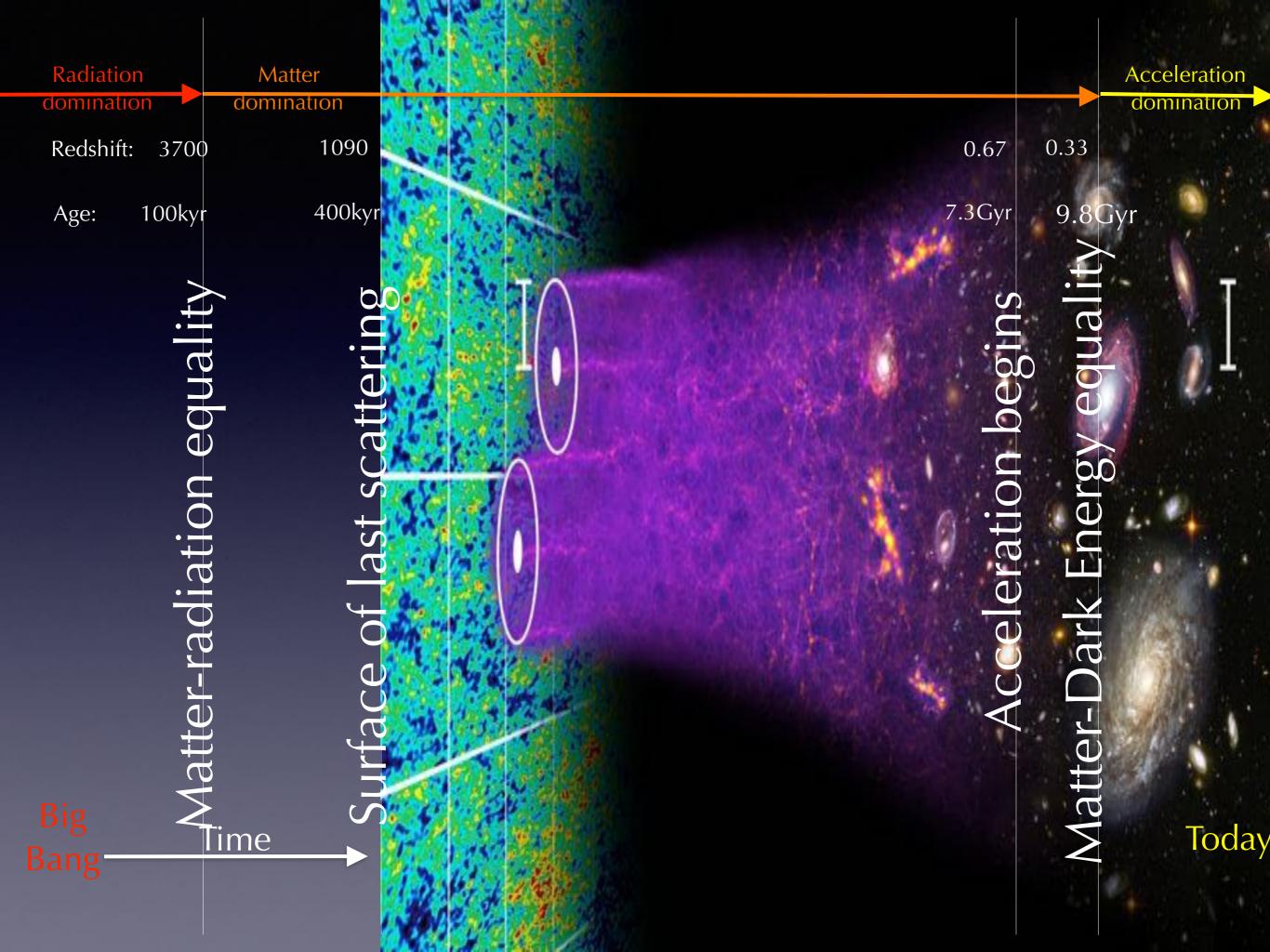
- Consider this as a change in the theory of gravity at large scales
- Can be either:
 - gravity gets weaker on large scales, owing to extradimension effects (Dvali-Gabadadze-Porrati model)
 - graviton has (induced) mass, meaning it does not propagate in the expected manner on large scales (massive gravity)
- Theories like this predict existence of extra degrees of freedom
 - The scalar degrees of freedom will affect the generation and propagation of gravitational instabilities

Theory of Structure Formation

Tracing structure

- Our observable universe is filled with structure, on all scales
- It's only visible through galaxies
 - The relation between distribution of galaxies and matter is given by the 'bias'
- Galaxies here are functioning as test particles
 tracing out the gravitational field





Post-recombination: Perturbation theory

Metric:

$$ds^{2} = -(1 + 2\Psi)dt^{2} + a^{2}(1 + 2\Phi)d\mathbf{x}^{2}$$

The force equation for matter is

$$\frac{d\theta_m}{dt} = -2H\theta_m - \frac{\partial^2}{a^2}\Psi$$

The perturbation equation is

$$\frac{d\delta_m}{dt} = -\theta_m - \dot{\Phi}$$

From these we derive the growth equation

$$\ddot{\delta} = 2H\theta - \frac{k^2}{a^2}\Psi$$

Einstein Gravity

Here the Newtonian potential φ is given by the Poisson equation

$$-k^2\Phi = 8\pi G \frac{a^2}{2} \left(\rho \delta + 3H(\rho + p) \frac{a^2}{k^2} \theta \right) ,$$

• If we assume no anisotropic stress, $\phi = -\Psi$ and so we can complete the system (assuming matter domination, and sub-horizon scales)

$$\ddot{\delta} = -2H\dot{\delta} + 4\pi G\rho\delta$$

- In the completely matter dominated limit, δ~a
- In the quasi-static limit, the growth of matter fluctuations can be written as

$$\frac{d\ln\delta}{d\ln a} \equiv f = -\frac{\theta_m}{H\delta_m}$$

Modified Gravities

Single fluctuation generated by Newtonian potential Ψ

$$\delta_m(k) = T_{\Phi}^{\delta_m}(k)\Phi_*(k)$$

$$\theta_m(k) = T_{\Phi}^{\theta_m}(k)\Phi_*(k)$$

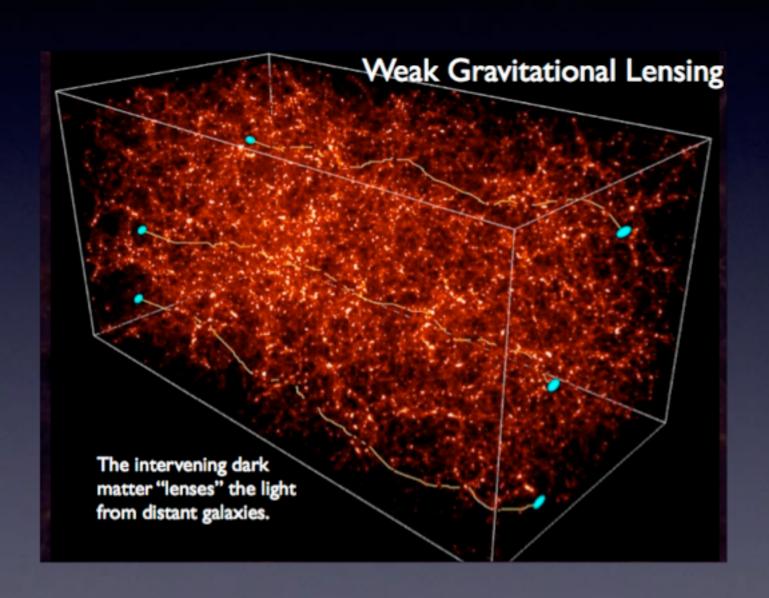
- transfer function T(k): describes how initial metric fluctuation is reprocessed into a late-time configuration of the species
- Growth rate now has a more general form

$$f_{\text{eff}}(k) \equiv -\frac{T_{\Phi}^{\theta_m}(k)}{T_{\Phi}^{\delta_m}(k)}$$

• Deterministic bias $\langle \delta_m \theta_m \rangle = -f_{\rm eff} \langle \delta_m \delta_m \rangle$ $\langle \theta_m \theta_m \rangle = f_{\rm eff}^2 \langle \delta_m \delta_m \rangle$.

Massless particle motion: gravitational lensing

- The motions of photons are also perturbed by the local gravitational potential
- This is is manifested as gravitational lensing
- The ellipticities of galaxy shapes become correlated with the matter density, integrated over the whole photon trajectory



Lensing potential

Null condition states

$$k^{\mu}k_{\mu}=0$$

Thin lens approximation gives

$$\frac{d^2x^i}{d\lambda_s} + 2\mathcal{H}\frac{d\eta}{d\lambda_s}\frac{dx^i}{d\lambda_s} - (\Phi_{,x^i} - \Psi_{,x^i})\left(\frac{d\eta}{d\lambda_s}\right)^2$$

Finally we compute deflection equation

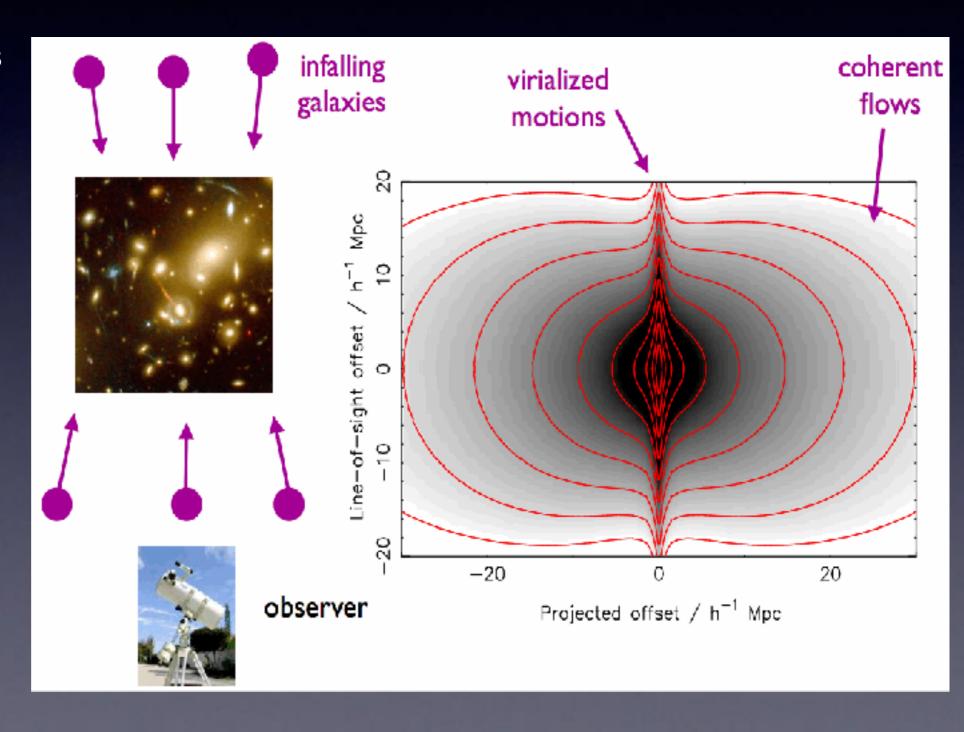
$$\frac{d^2x^i}{dr^2} = \Phi_{,x^i} - \Psi_{,x^i}$$

 Difference between potentials is lensing potential, and deflection is sourced by spatial gradient of lensing potential

Cosmological Data

Redshift-space distortions

- The motions of galaxies are perturbed by the local gravitational field
- The Power spectrum/ correlation function in the line of sight is distorted relative to the transverse direction
- Assuming these motions are generated by matter perturbations, we can measure the growth of structure



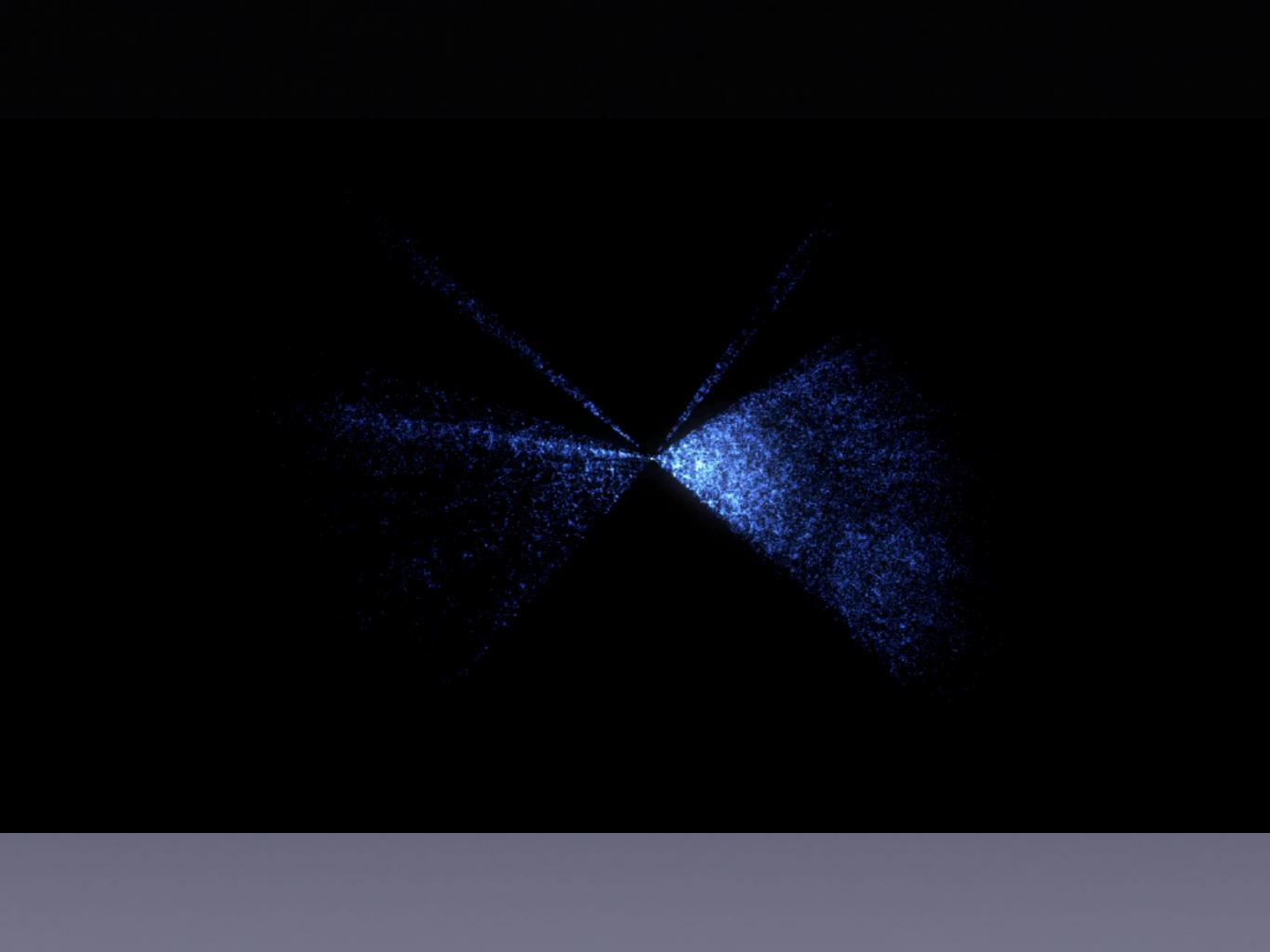
Multipole power spectra

- Density and velocity divergence have different angular dependence
- Use Power spectra decomposed into Legendre polynomials (Cole, Fisher and Weinberg 1994)

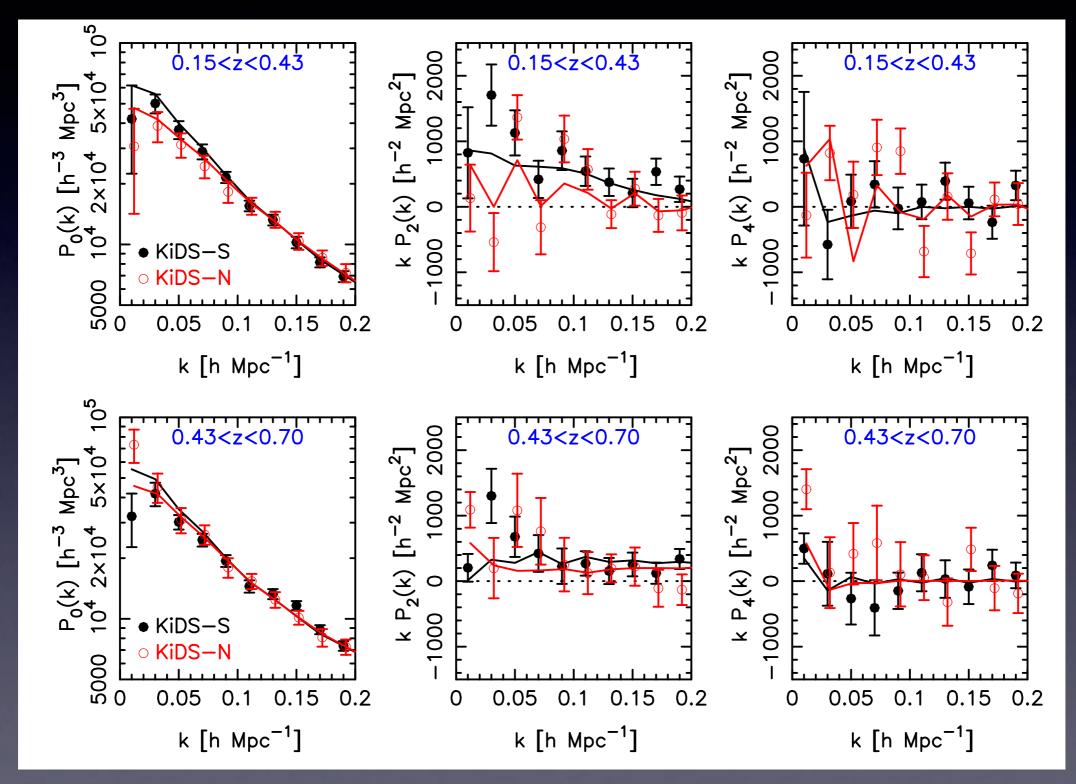
$$P(k,\mu) = \sum_{\ell=0}^{\infty} P_{\ell}(k) L_{\ell}(\mu)$$

Orthogonality of the Legendre polynomials leads to the relation

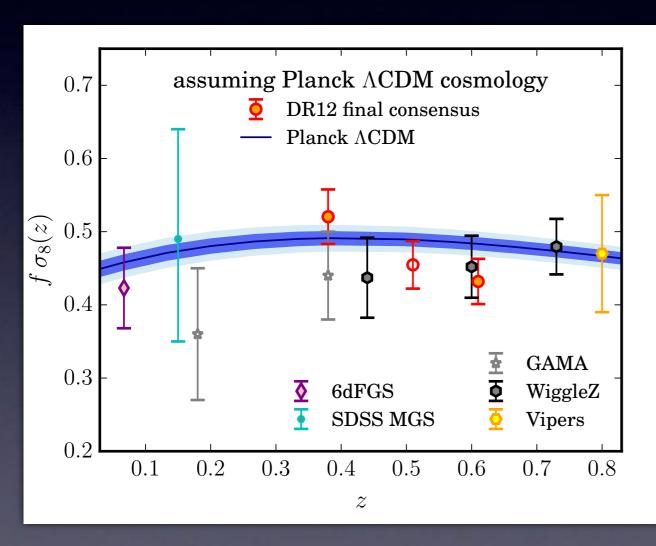
$$P_{\ell} = \frac{2\ell + 1}{2} \int_{-1}^{+1} d\mu P^{s}(k, \mu) L_{\ell}(\mu)$$



2dFlens - Results



Growth history



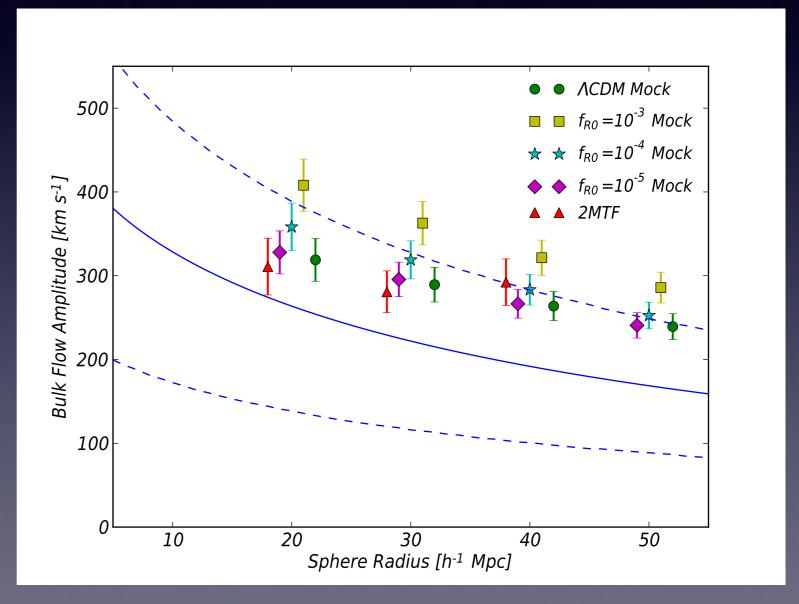
Alam et al 2016

- Growth rate (f) and amplitude of fluctuations (σ₈) sourced by both gravitational force and expansion rate
- Need to fit for both simultaneously, so some degeneracy with BAO signal

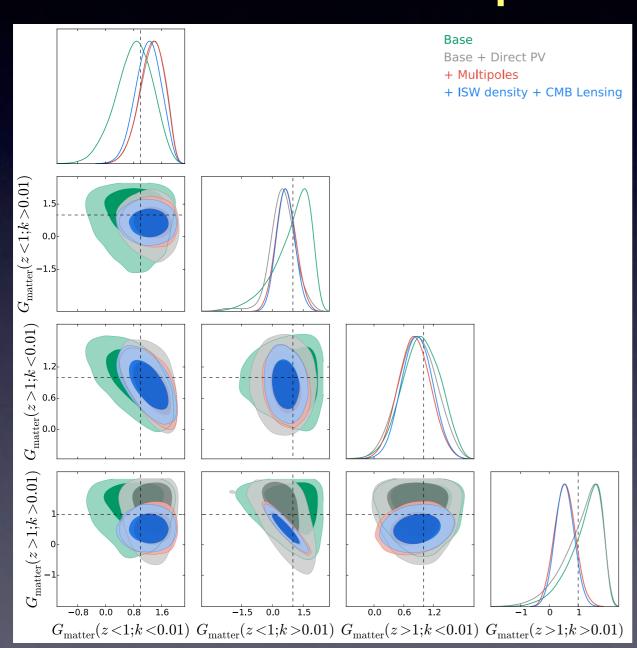
Peculiar velocities

- Data from the Two Micron All-Sky Survey (2MASS; Skrutskie et al. 2006) Tully-Fisher Survey (2MTF; Masters 2008) covers most of the sky, and uses 2018 galaxies to measure the bulk flow.
- Numerical simulation for the same sky, and select from the same redshift distribution
- f(R) gravity predicts a larger bulk flow velocity than ΛCDM

 $z_{\rm observed} \simeq z_{\rm cosmological} + z_{\rm peculiar}$



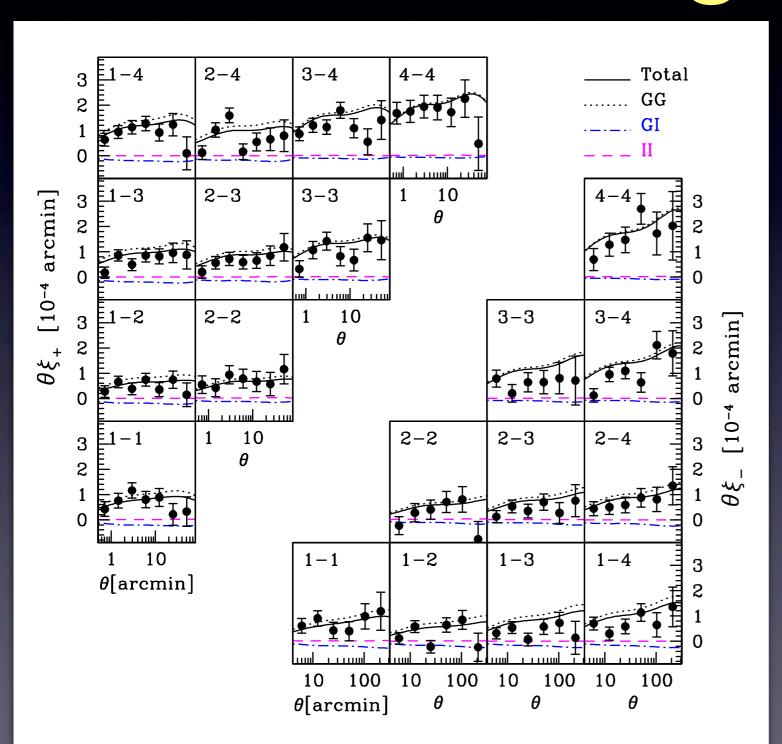
Peculiar velocity power spectra



Johnson, Blake, Dossett, Koda, Parkinson, Joudaki, MNRAS 458 (2016), 2725-2744

- Can measure power spectra of velocities
 - No galaxy bias
- Use to measure deviations from gravitational force law on different scales

Lensing Data



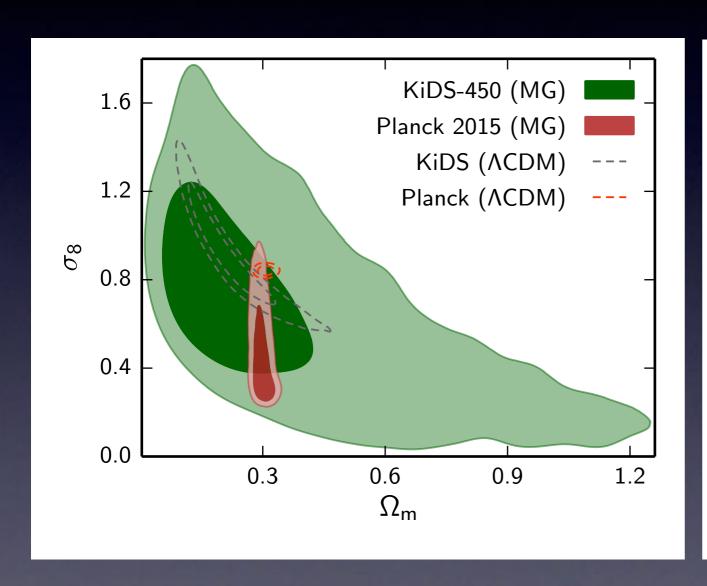
Lensing convergence power spectra

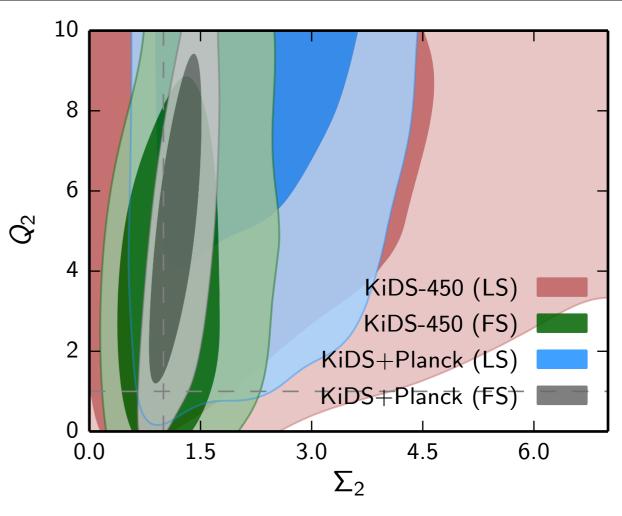
$$P\kappa^{ij}(\ell) = \int_0^{r_H} dr \frac{q_i(r)q_j(r)}{[f_{\kappa}(r)]^2} P_{\delta}\left(\frac{\ell}{f_{\kappa}(r)}, r\right)$$

 tomographic angular two-point shear correlation function

$$\xi_{\pm}^{ij}(\theta) = \frac{1}{2\pi} \int d\ell \ell P_{\kappa}^{ij}(\ell) J_{0,4}(\ell\theta)$$

Parameter constraints





Fifth forces

Fifth forces

- (With Seery and Burrage)
- Metric: $ds^2 = -(1+2\Psi)dt^2 + a^2(1+2\Phi)d\mathbf{x}^2$
- The force equations for matter and radiation are

$$\frac{d\theta_m}{dt} = -2H\theta_m - \frac{\partial^2}{a^2}\Psi + F_{5r}$$

$$\frac{d\theta_r}{dt} = -H\theta_r - \frac{\partial^2}{a^2}\Psi - \frac{1}{4}\frac{\partial^2}{a^2}\delta_r + F_{5m}$$

Perturbation equations

$$\frac{d\delta_m}{dt} = -\theta_m - 3\dot{\Phi} + j_{5m}$$

$$\frac{d\delta_r}{dt} = -\frac{4}{3}\theta_r - 4\dot{\Phi} + j_{5r}$$

Stochastic bias

- Several primordial perturbations:
 - fifth force mediated by scalar field (φ) that also has its own fluctuations

$$\delta_m(k) = T_{\Phi}^{\delta_m}(k)\Phi_*(k) + T_{\phi}^{\delta_m}(k)\delta\phi_*(k)$$

$$\theta_m(k) = T_{\Phi}^{\theta_m}(k)\Phi_*(k) + T_{\phi}^{\theta_m}(k)\delta\phi_*(k).$$

Relation between growth and density has now changed

$$\theta_m = -f_{\text{eff}}\delta_m + \frac{W}{T_{\Phi}^{\delta_m}}$$

 Here W is a Wronskian-like function which measures the correlation of δ and θ

New correlation functions

The correlation functions now satisfy

$$\langle \theta \delta \rangle_{k} = \langle \delta \delta \rangle_{k} \left(-f_{eff}^{\Phi}(k) + \frac{f_{eff}^{\Phi}(k) - f_{eff}^{\phi}(k)}{1 + \rho(k)} \right)$$
$$\langle \theta \theta \rangle_{k} = \langle \delta \theta \rangle_{k} \left(-f_{eff}^{\Phi}(k) + \frac{f_{eff}^{\Phi}(k) - f_{eff}^{\phi}(k)}{1 + \sigma(k)} \right)$$
$$= \langle \delta \delta \rangle_{k} \left(-f_{eff}^{\Phi 2}(k) + \frac{f_{eff}^{\Phi 2}(k) - f_{eff}^{\phi 2}(k)}{1 + \rho(k)} \right)$$

• If we measure $\langle \delta_m \delta_m \rangle$, $\langle \delta_m \theta_m \rangle$ and $\langle \theta_m \theta_m \rangle$, and do not see complete correlation between δ and θ , we have gone beyond f_{eff} as a test of modified gravity

Example: Galileons

 Galileons are scalar fields that are invariant under shifts in the field value

$$\pi \to \pi + c + b_{\mu}x^{\mu}$$

 Only 5 possible Lagrangians that give 2nd order equations of motion and are ghost free

$$\mathcal{L}_{1} = M^{3}\pi$$

$$\mathcal{L}_{2} = (\nabla \pi)^{2}$$

$$\mathcal{L}_{3} = (\Box \pi)(\nabla \pi)^{2}/M^{3}$$

$$\mathcal{L}_{4} = (\nabla \pi)^{2}[2(\Box \pi)^{2} - 2\pi_{;\mu\nu}\pi^{;\mu\nu} - R(\nabla \pi)^{2}/2]/M^{6}$$

$$\mathcal{L}_{5} = (\nabla \pi)^{2}[(\Box \pi)^{3} - 3(\Box \pi)\pi_{;\mu\nu}\pi^{;\mu\nu} + 2\pi^{\nu}_{:\mu}\pi^{\rho}_{:\rho}\pi^{\mu}_{:\rho} - 6\pi_{;\mu}\pi^{;\mu\nu}\pi^{;\rho}G_{\nu\rho}]/M^{9}$$

Coupling to matter

We expect a coupling of the Galileon field (for example) to matter

$$T_{\phi} = -rac{1}{2\Lambda} \left(1 + rac{\phi}{\Lambda}
ight)^{-1} T_m$$

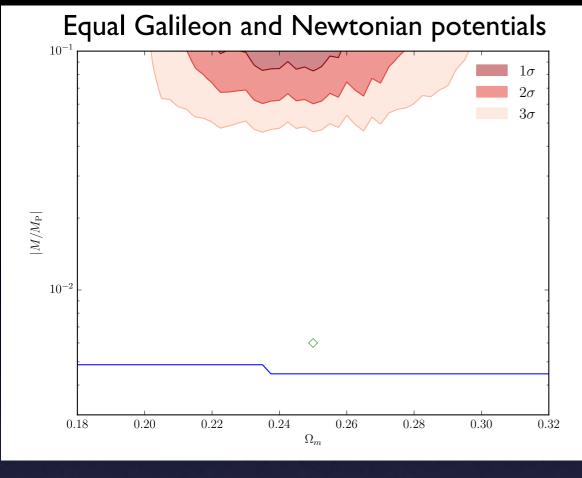
- The field starts small, and remains small during the evolution of the universe
- The perturbation of the field can grow, and will grow quickly at late times through the coupling to the matter perturbation
 - While the Galileon fluctuation will be smaller than the matter perturbation at matter-radiation decoupling, it can grow and become very large today.
- This is stochastic bias, where structures can form without presence of standard `Newtonian' potential generated by the presence of matter

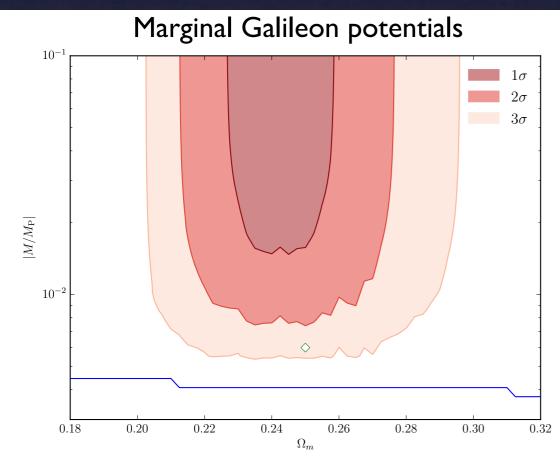
Screening mechanisms

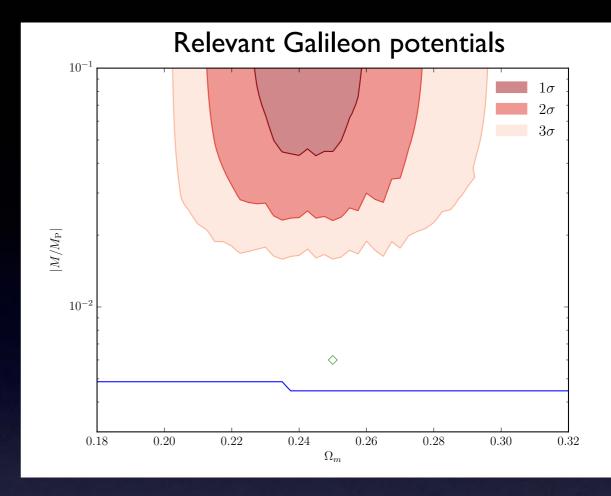
- If the gravity is different, we can test it on lab or solar system system scales
 - e.g. fifth force effect, or scale-dependent G_{Newton}
- Three "screening mechanisms" save the theories

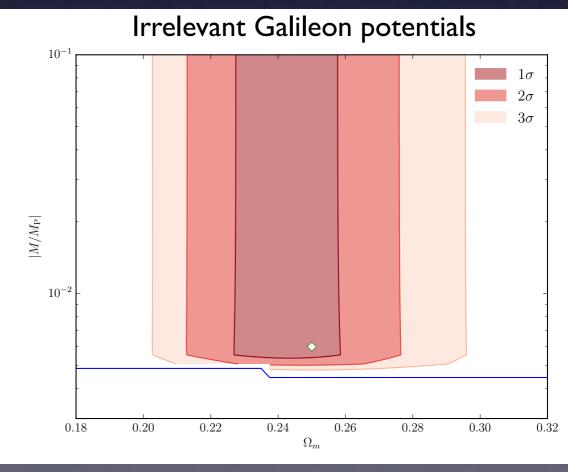
$L \supset -\frac{1}{2}Z(\varphi_0)(\partial \delta \varphi)^2 - \frac{1}{2}m^2(\varphi_0)\delta \varphi^2 + (\beta(\varphi)/m_P)\delta \varphi \delta T$

- Vainshtein mechanism: higher-order corrections (cubic and above) recover GR on scales smaller than Vainshtein radius (DGP, Galileon)
- e : mass of field large enough to suppress range of fifth force (f(R) theories)
- Symmteron mechnism: direct coupling to stress-energy tensor (T) is small









Lensing

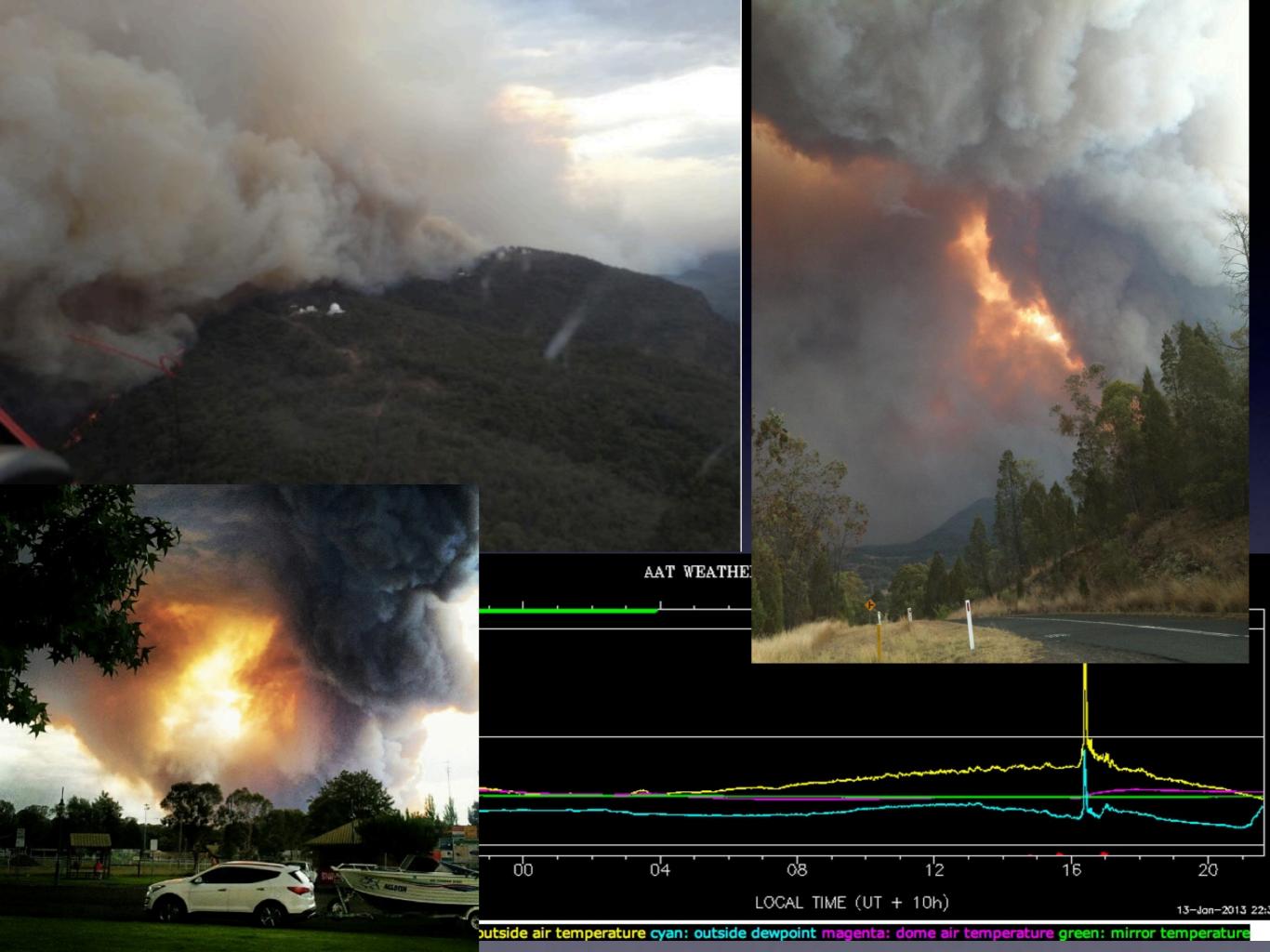
- Fifth forces: cannot distinguish between strength of coupling and size of field fluctuations
- Need to measure size of metric fluctuations independently
- A conformal coupling induces no change in the lensing potential
 - Photons only feel Newtonian potential
- However, matter formation (that sources the lensing potential) still influenced by fifth forces
- Considering only Galileon actions up to L_3 (which we are), there is no anisotropic stress, so $\phi = -\Psi$
- Cross-correlation of lensing with RSD will determine strength of stochastic bias

Conclusions

- Structure formation tests of gravity effectively measure force law on largest scales
- The growth rate f gives the correlation between density and velocity statistics of galaxies
- Lensing power spectrum gives correlation between induced ellipticity and density of matter
- If fifth forces become important at late times, the density and velocity (or velocity and lensing) perturbations may no longer be perfectly correlated
 - Galileons are an example of this behaviour, as the fluctuations in the field grow at late times and become a new source of structure formation, leading to stochastic bias
- Cross-correlation of lensing and RSD data will provide direct measurement of stochastic bias effect

Thank you





2dFLens

Spectroscopic survey, providing follow-up of lensing

galaxies from KiDS survey

 Swinburne: Chris Blake (PI), Karl Glazebrook, Andrew Johnson, Shahab Joudaki, Felipe Marin

- University of Queensland: David Parkinson
- Mount Stromlo, ANU: Mike Childress, Chris Wolf
- Edinburgh: Catherine Heymans, Alexandra Amon
- Bonn: Thomas Erben, Hendrik Hildebrandt, Dominik Klaes
- Konrad Kuijken (Leiden)
- Chris Lidman (AAO)
- Greg Poole (U.Melbourne)

