



INSTITUTO DE FÍSICA  
Universidade Federal Fluminense



# Understanding the bulk viscosity of QCD matter

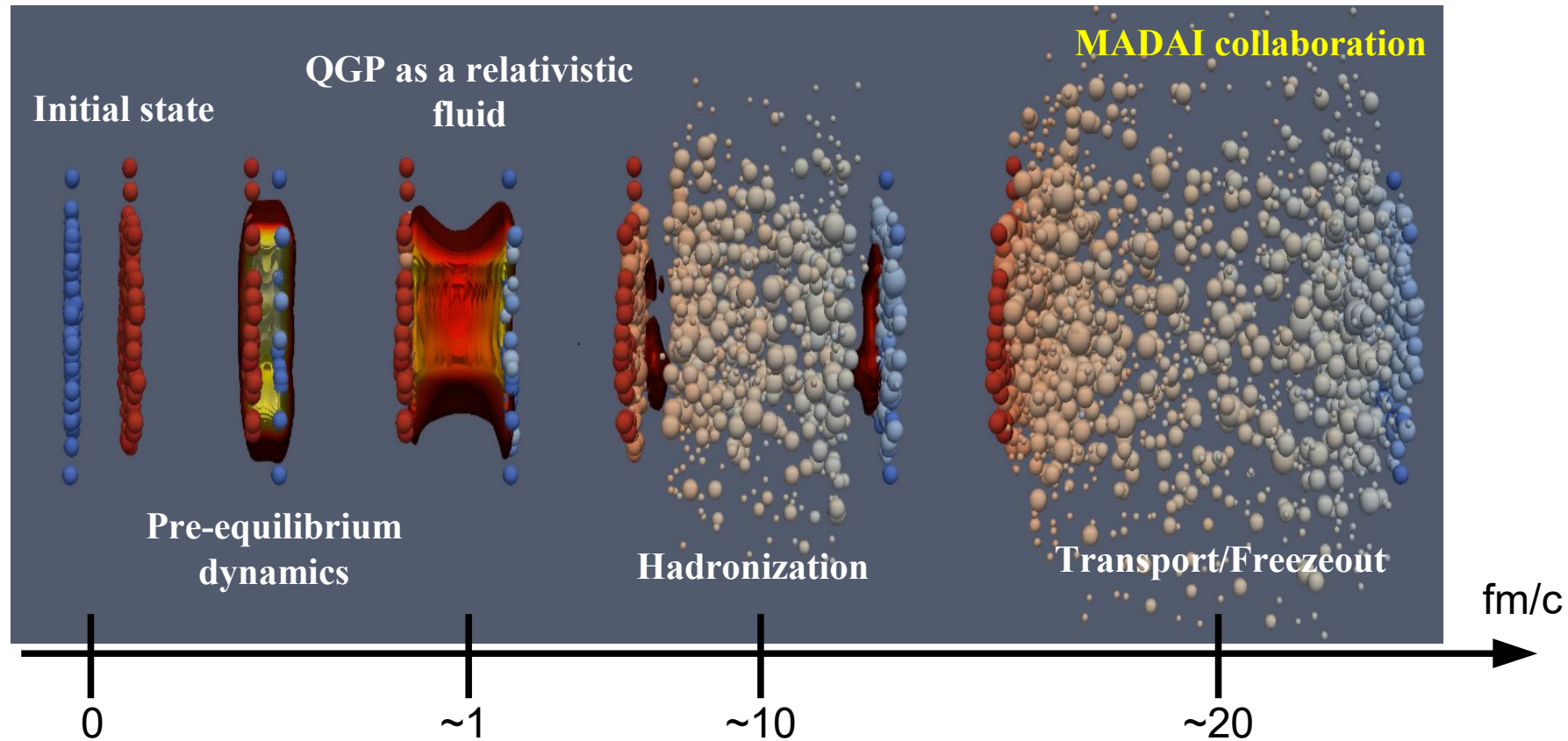
Gabriel S. Denicol (UFF)

**UL**tra-**Rel**ativistic **CH** **HE**avy **IoNZ**

18-20.July.2016

# Ultra-relativistic Heavy Ion Collisions

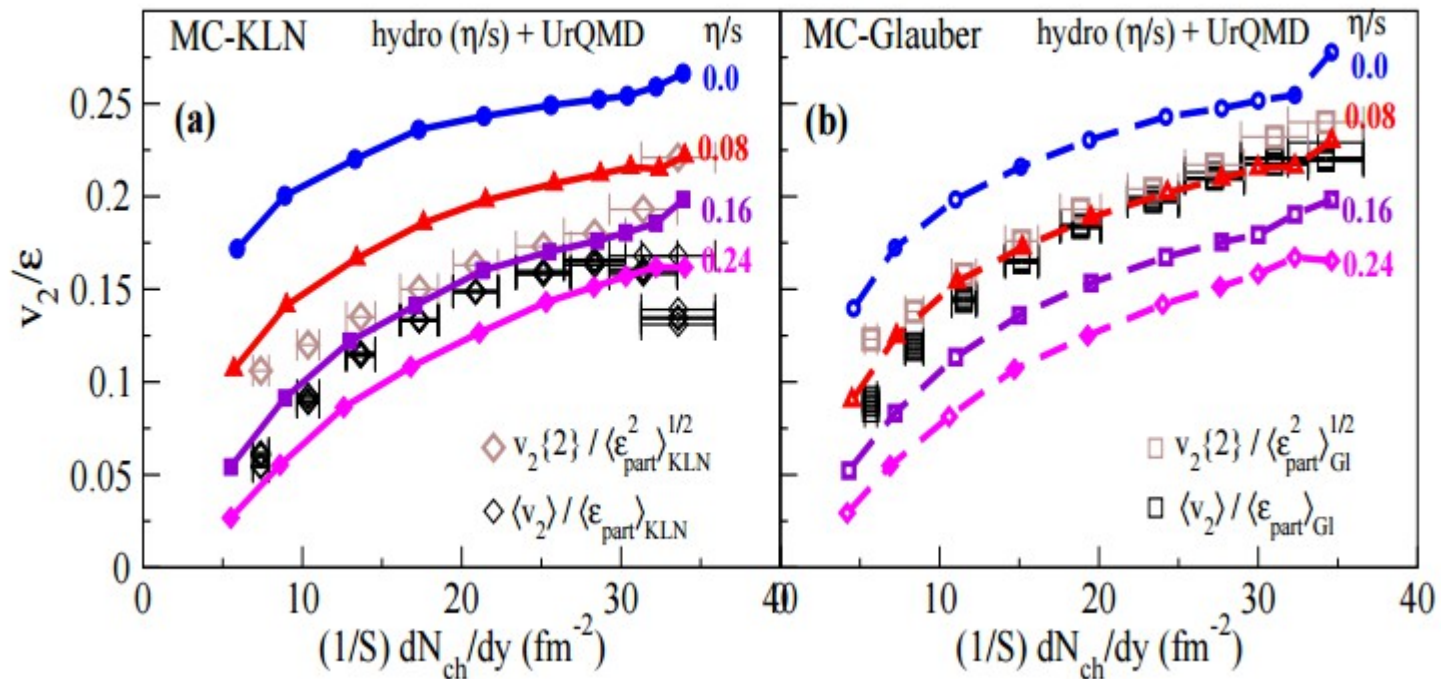
Best (and often only) option to understand certain properties of bulk nuclear matter



Bulk QCD matter is only created transiently  
Need to reverse-engineer its properties

# Extraction of shear viscosity

Shear viscosity can be estimated using elliptic flow data (IC uncertainty)



Song *et al*, PRL 106, 192301 (2011)

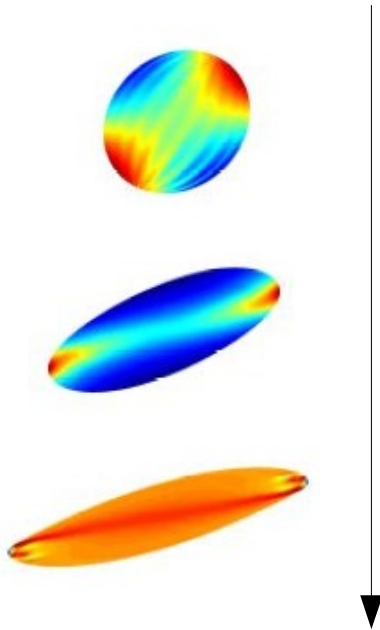
Major contribution from Heinz's group

# What about other transport properties?

## Shear viscosity

Resistance to deformation

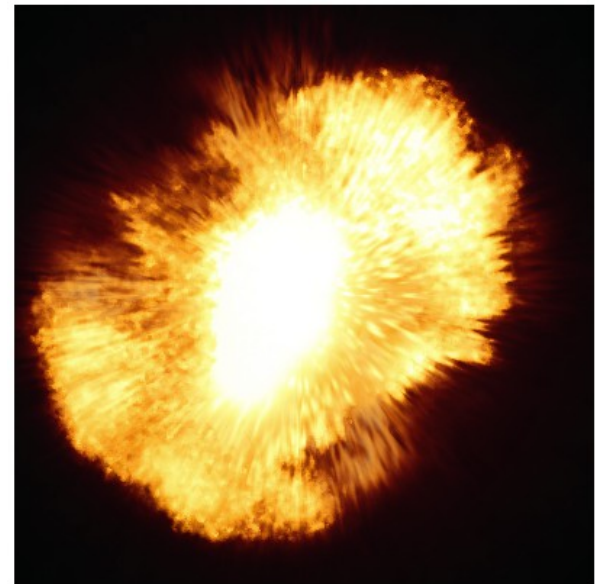
$$\pi^{\mu\nu} = 2\eta \nabla^{\langle\mu} u^{\nu\rangle}$$



## Bulk viscosity

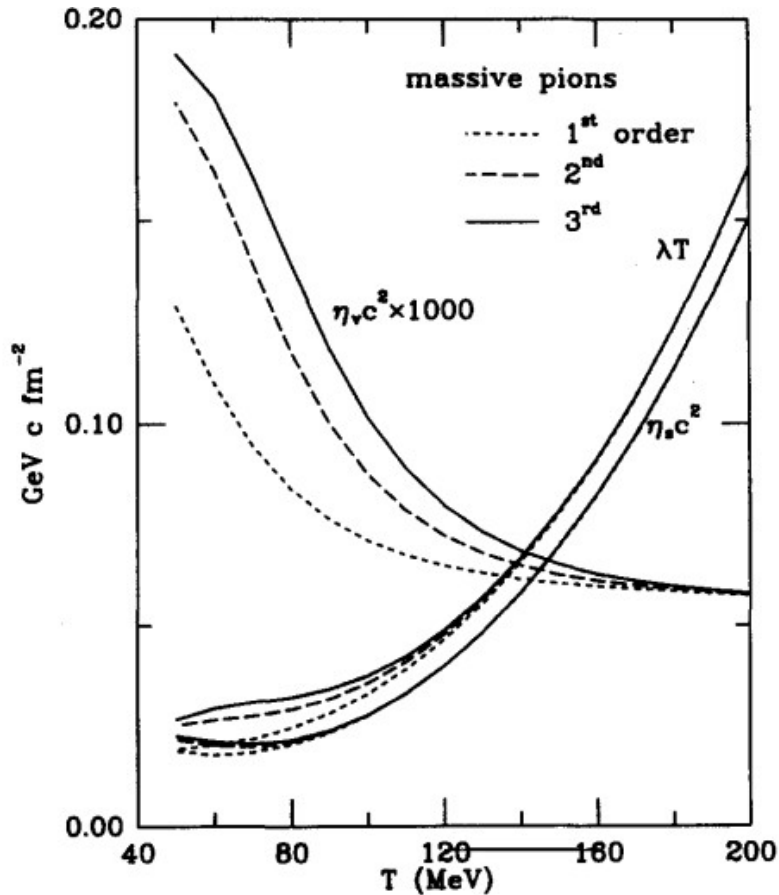
Resistance to volume change

$$\Pi = -\zeta \nabla_{\mu} u^{\mu}$$

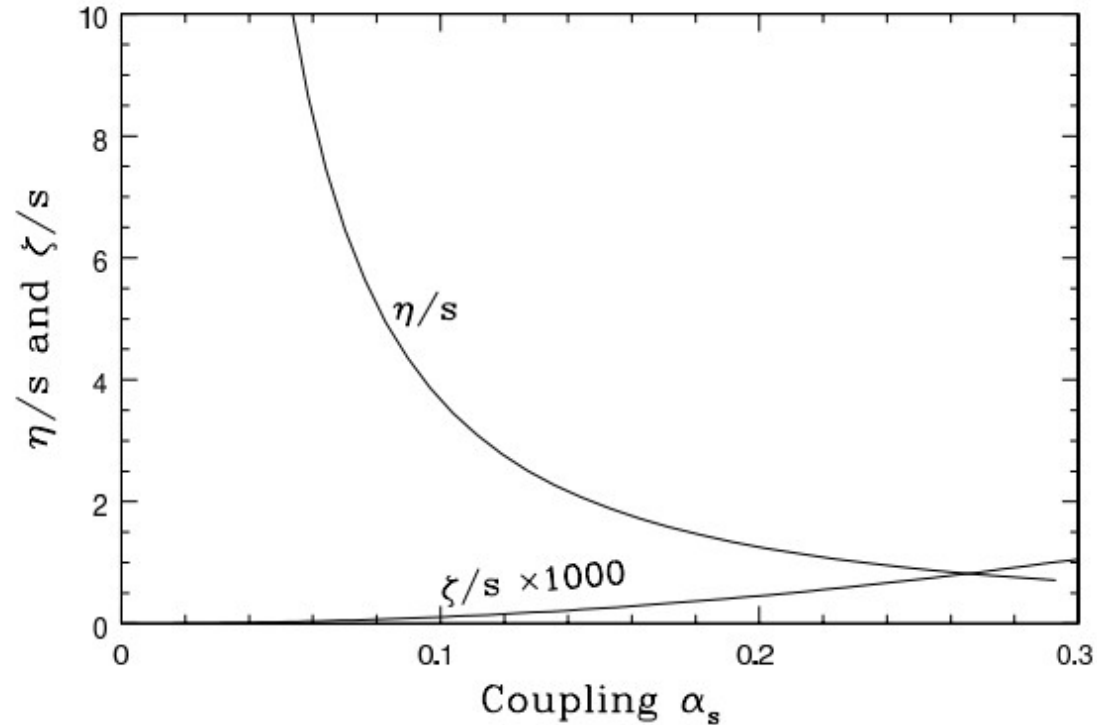


# Some limits

pions gas



pQCD



Arnold *et al*, PRD 74, 085021 (2006)

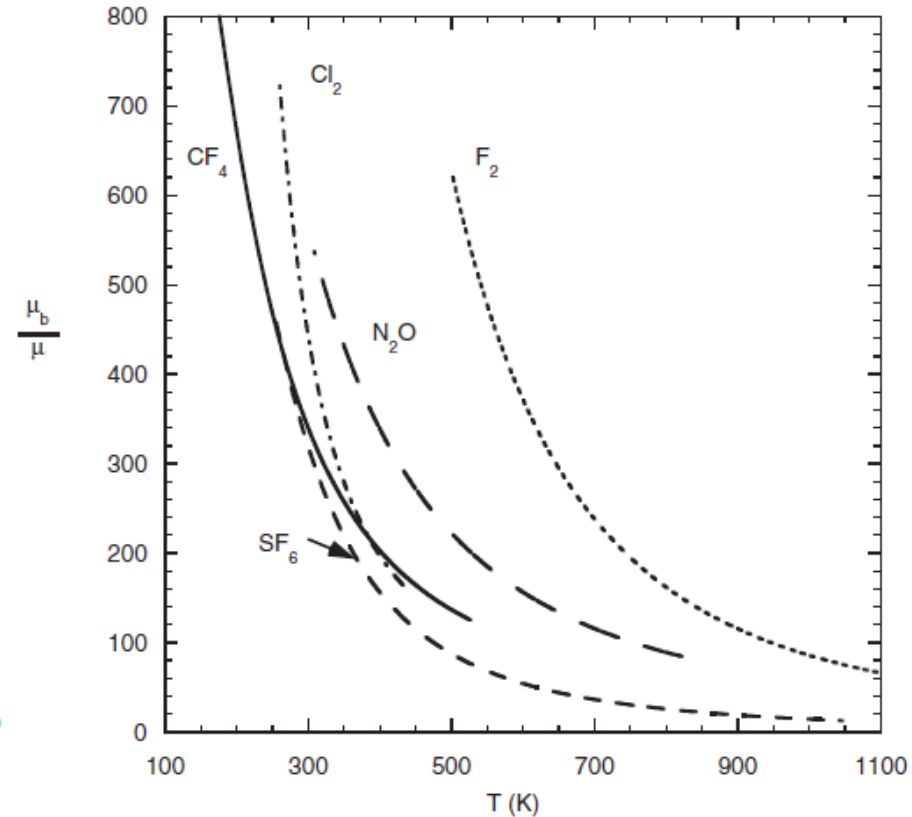
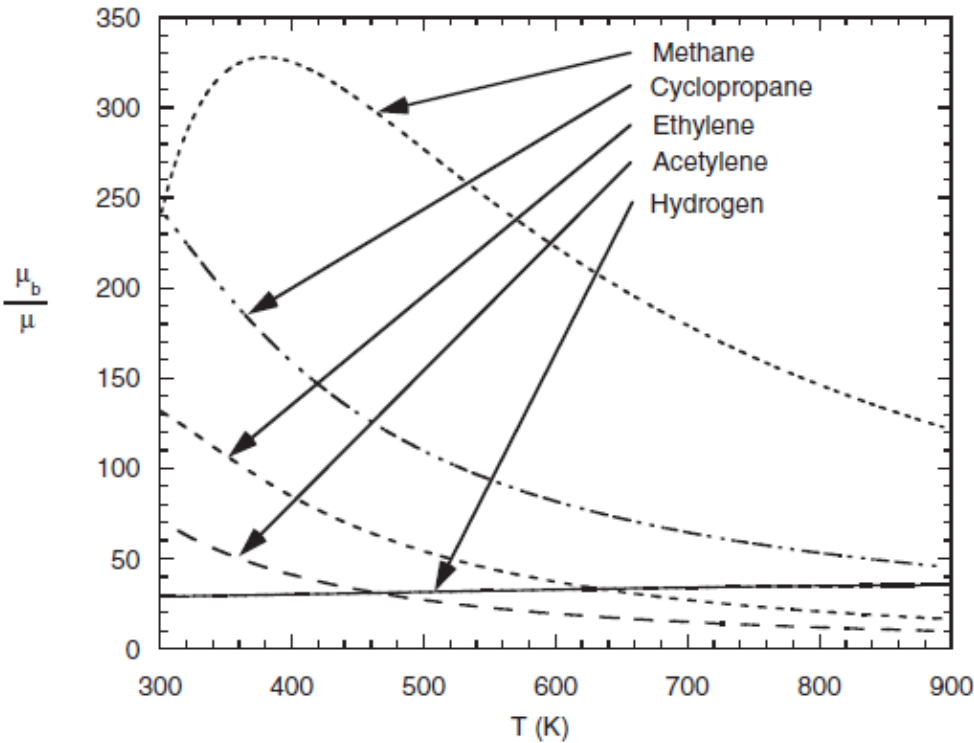
Prakash *et al*, Phys. Rept. 227, 321 (1993)

**Bulk viscosity ~1000 times smaller than shear**

**Near phase transition, we do not know ...<sup>5</sup>**

# Some known examples

## Bulk to shear viscosity ratio for several gases

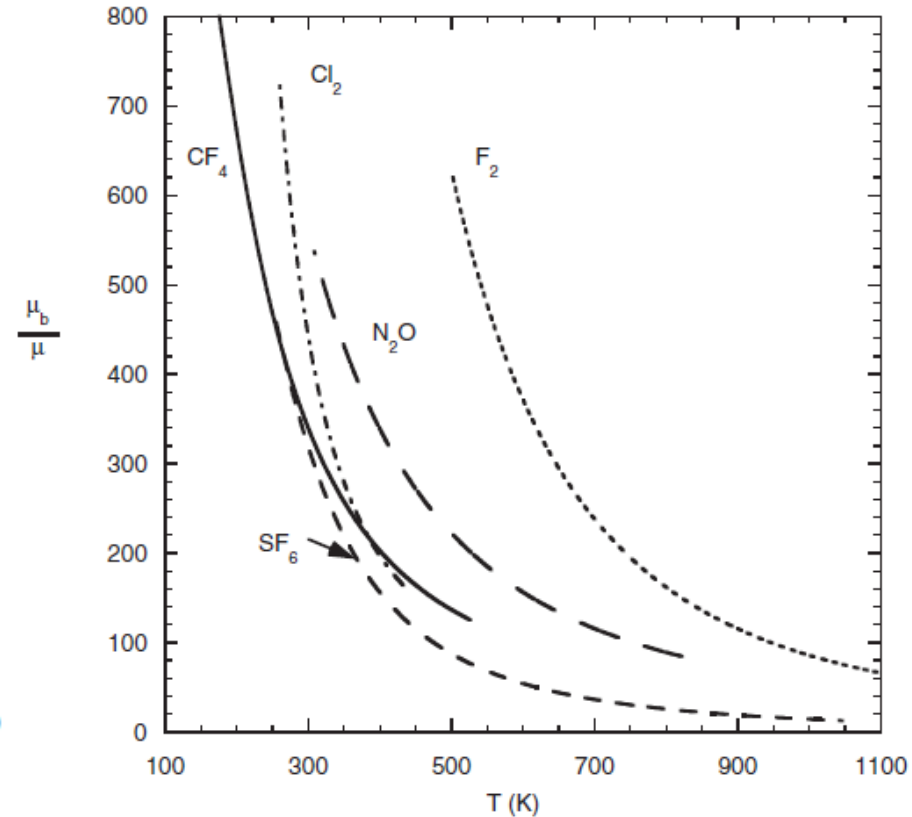
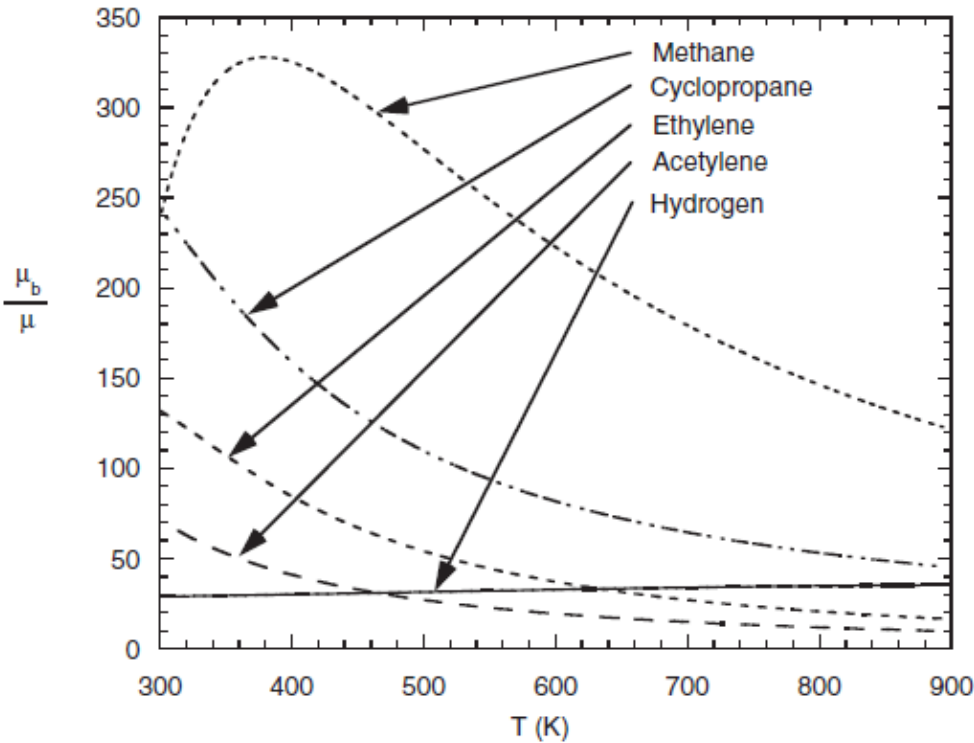


M. Cramer, Phys. of fluids 24, 066102 (2012)

In many cases, bulk viscosity can be larger than shear viscosity (failure of Boltzmann eq.) <sup>6</sup>

# Some known examples

## Bulk to shear viscosity ratio for several gases



M. Cramer, Phys. of fluids 24, 066102 (2012)

**bulk viscosity:** vibrational&rotational relax. time  
**shear viscosity:** mean free path

What you will see in this talk

**We investigate the effect of bulk viscosity  
on basic heavy ion collision observables**

**Is bulk viscosity needed to fit the data?**

**Can it be extracted? What observable  
should we use?**



# Fluid-dynamical model

## IP-Glasma initial condition

$$\tau_0 = 0.4 \text{ fm}$$

thermalization “by hand”

## Relativistic Hydrodynamics (MUSIC)

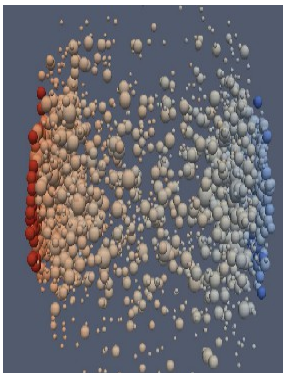
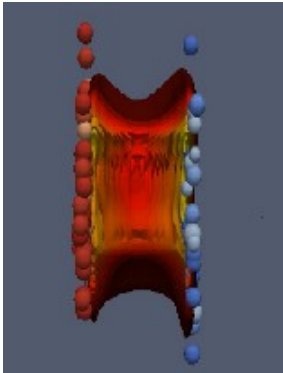
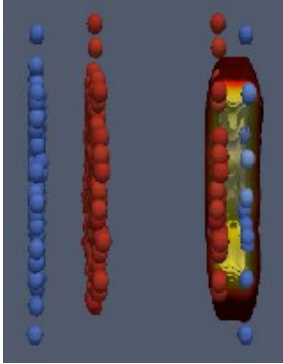
$$\partial_\mu T^{\mu\nu} = 0 \quad \partial_\mu N^\mu = 0$$

+ EoM for dissipative currents

**T=const.**

fluid elements converted  
to particles

## Hadronic transport (UrQMD)



# Fluid-dynamical equations

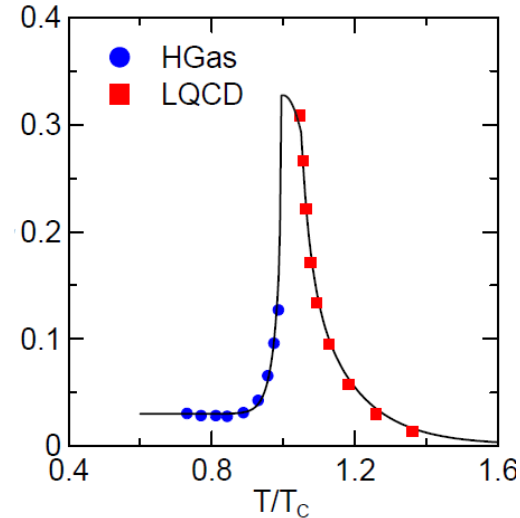
GSD&Niemi&Molnar&Rischke, PRD 85, 114047 (2012)

## Inclusion of bulk viscous pressure, shear-stress tensor, and all couplings

$$\begin{aligned} \dot{\Pi} + \frac{\Pi}{\tau_{\Pi}} &= -\beta_{\Pi}\theta - \delta_{\Pi\Pi}\Pi\theta + \varphi_1\Pi^2 + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu} + \varphi_3\pi^{\mu\nu}\pi_{\mu\nu}, \\ \dot{\pi}^{\langle\mu\nu\rangle} + \frac{\pi^{\mu\nu}}{\tau_{\pi}} &= 2\beta_{\pi}\sigma^{\mu\nu} + 2\pi_{\alpha}^{\langle\mu}\omega^{\nu\rangle\alpha} - \delta_{\pi\pi}\pi^{\mu\nu}\theta + \varphi_7\pi_{\alpha}^{\langle\mu}\pi^{\nu\rangle\alpha} - \tau_{\pi\pi}\pi_{\alpha}^{\langle\mu}\sigma^{\nu\rangle\alpha} \\ &\quad + \lambda_{\pi\Pi}\Pi\sigma^{\mu\nu} + \varphi_6\Pi\pi^{\mu\nu}. \end{aligned}$$

✓  $\eta/s = \text{const}$ ,

✓  $\zeta/s =$

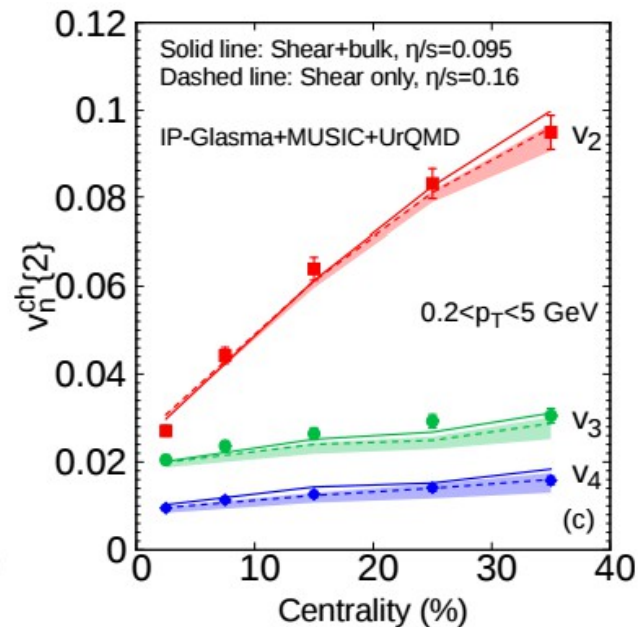
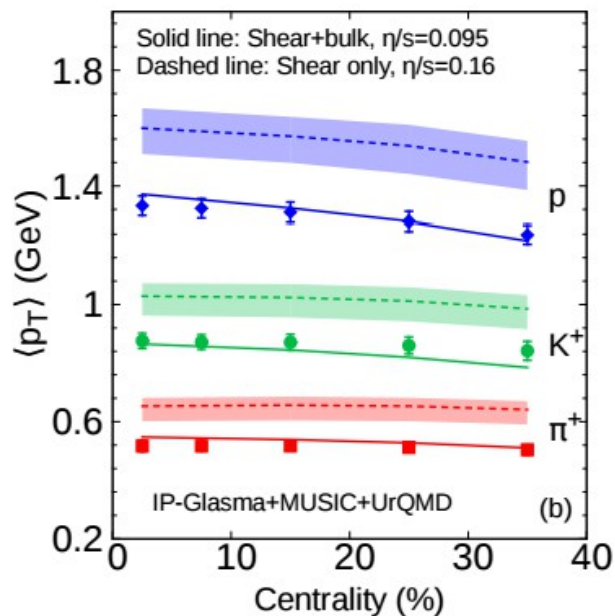
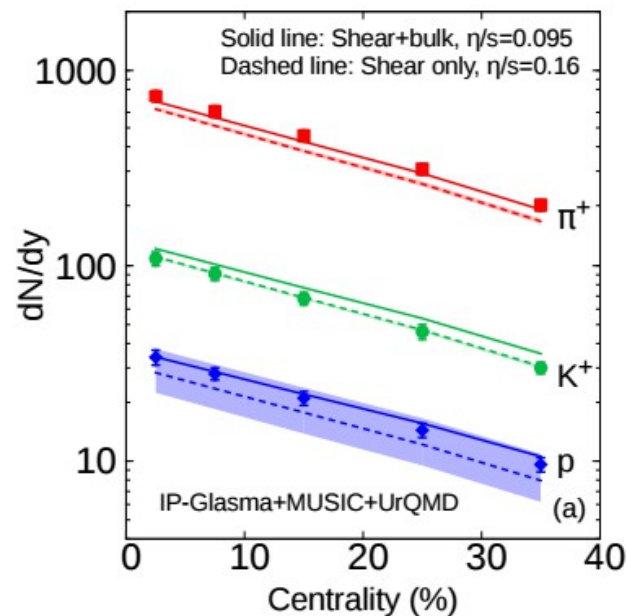


# Integrated observables

IP-Glasma+MUSIC+UrQMD

PRL 115, 132301 (2015)

$T_{\text{switch}} = 145 \text{ MeV}$

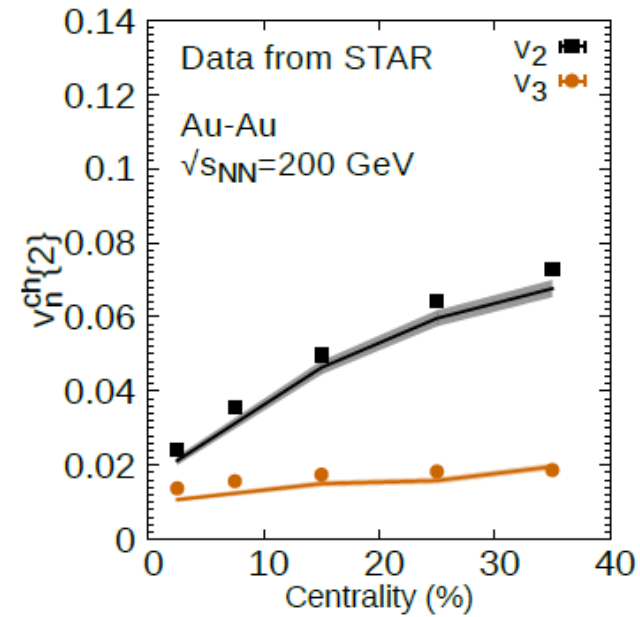
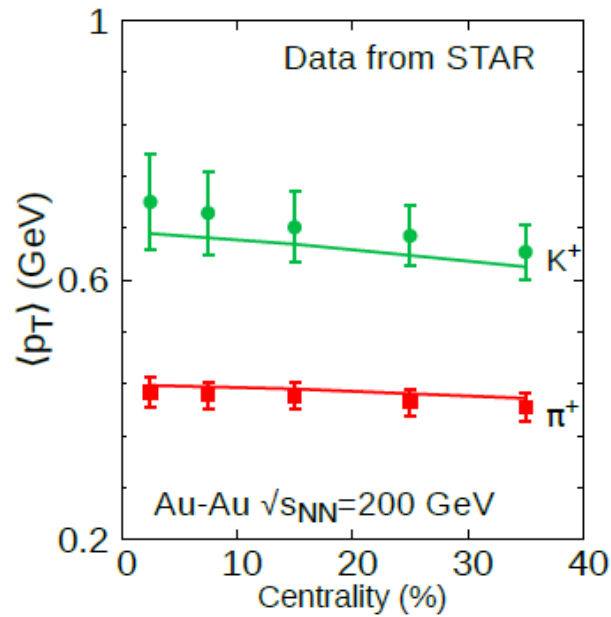
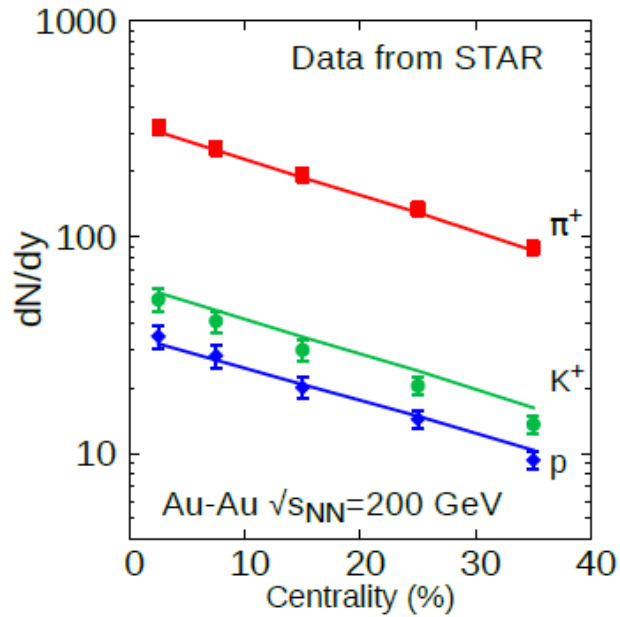


**bulk viscosity increases multiplicity, reduces  $\langle p_T \rangle$ , and reduces  $V_n$**

Value of shear viscosity extracted changes significantly  
 $\eta/s=0.16 \longrightarrow 0.095$

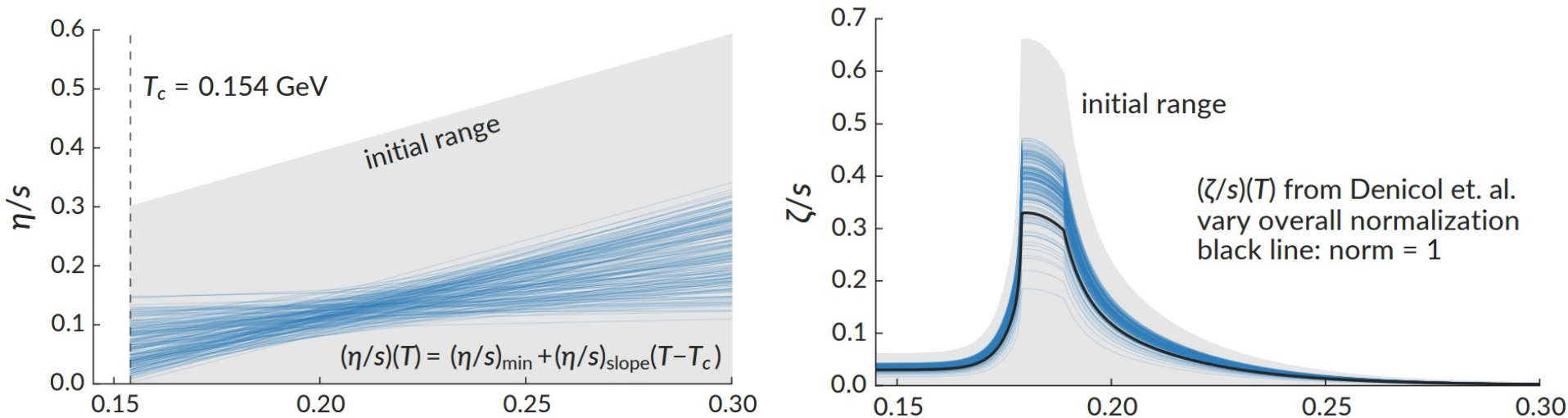
# What about RHIC?

$T_{\text{switch}}=165 \text{ MeV}$   
 $\eta/s=0.06$



**Same bulk viscosity can also describe RHIC data**

Model: Trento initial state + Hydro (bulk&shear) + UrQMD  
Very good description of Multiplicity, mean-pT, and Vn

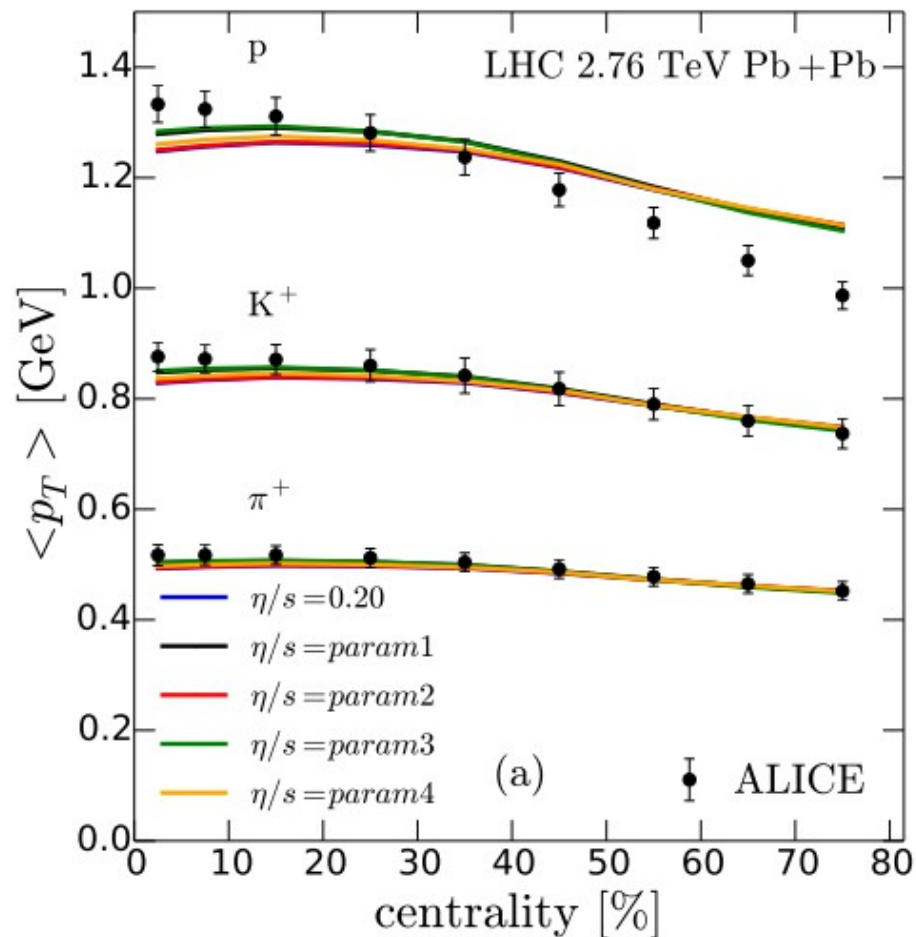
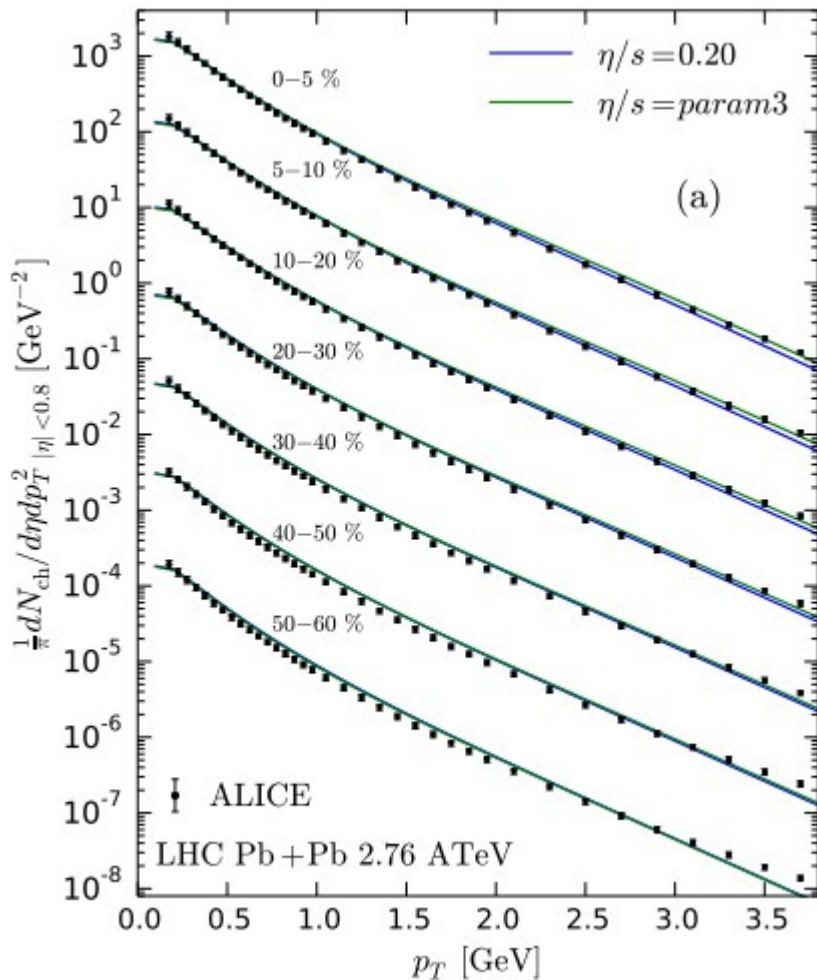


**finite bulk viscosity is favored**

Similar findings also by Schenke&Monnai for  
MC-Glauber with valence quarks

# EKRT model + second order viscous hydro

Niemi *et al*, PRC 93, no. 2, 024907 (2016)



**Good description without including bulk viscosity!**  
**But, chemical freeze-out at 175 MeV**

**Non-equilibrium  
contribution to pressure**

$$\text{Pressure} = P_0 + \Pi$$

“ $\delta P$ ”

Hydrodynamical picture works, if this is a perturbative correction to the thermodynamic pressure

$$P_{\text{PCE}} = P_0 + \delta P \quad \longrightarrow \quad \text{bulk?}$$

Different descriptions of the same effect?

**Why did bulk  
become necessary?**



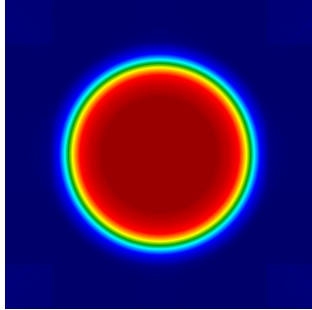
# Evolution of initial state

Coarse-graining  
size

$\varepsilon(x,y)$

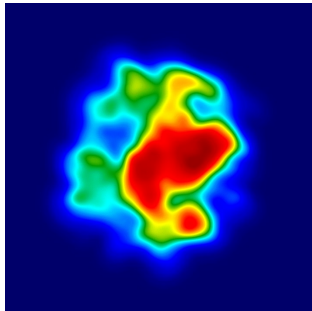
Smooth Initial conditions  
( $v_2=0$  in central collisions)

$\lambda \sim 5$  fm



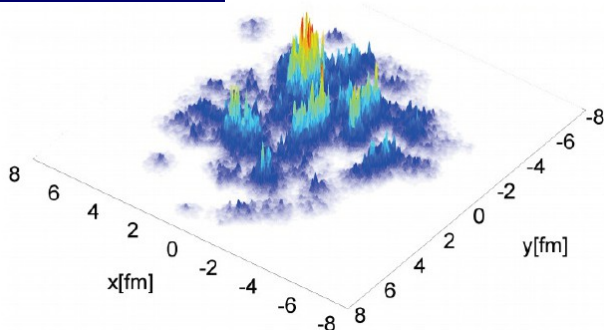
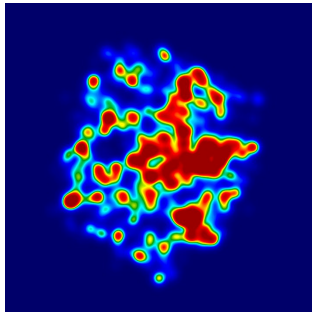
Fluctuations in nucleon positions  
(odd harmonics)

$\lambda \sim 1$  fm



sub-nucleonic fluctuations  
( $v_n$  distributions & correlations)

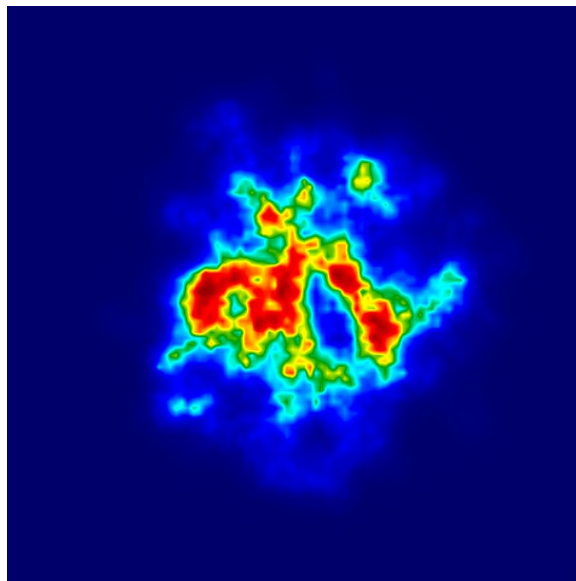
$\lambda \sim 0.4$  fm



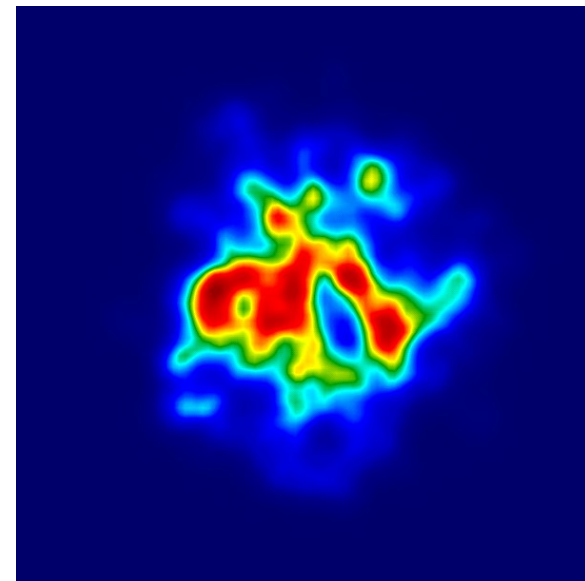
This happens because the IP-Glasma model gives rise to an initial state with large gradients of pressure and the subsequent fluid-dynamic expansion accordingly produces a significant radial flow.

# Effect of smoothing

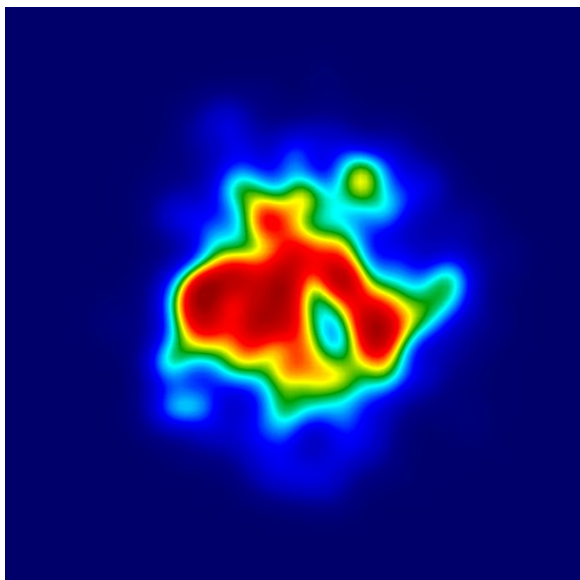
IP-Glasma event  
0-5% class



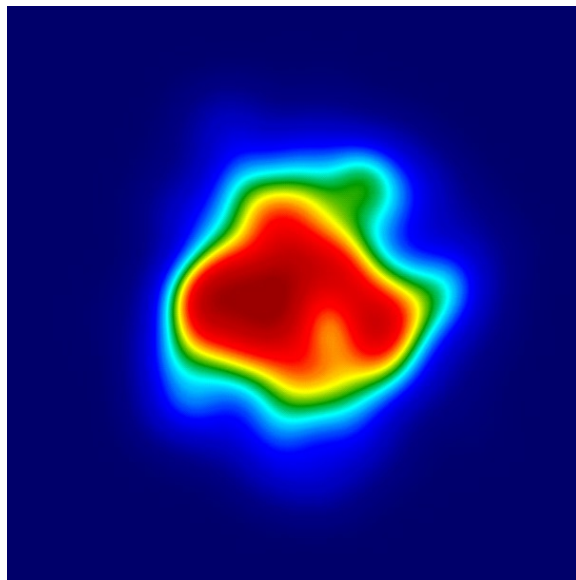
$\lambda = 0.1$  fm



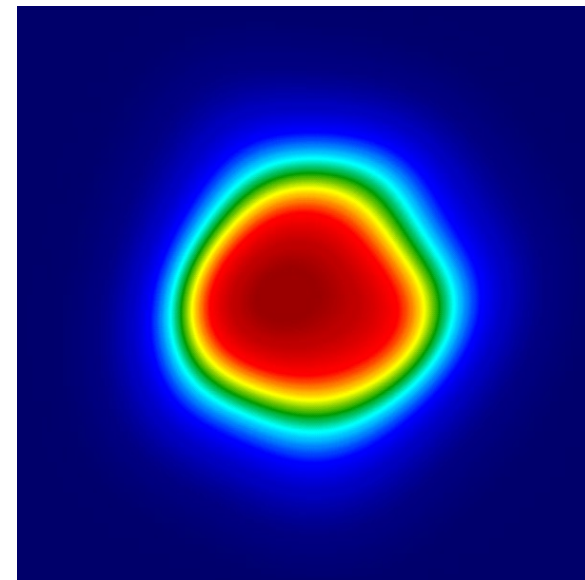
$\lambda = 0.2$  fm



$\lambda = 0.4$  fm

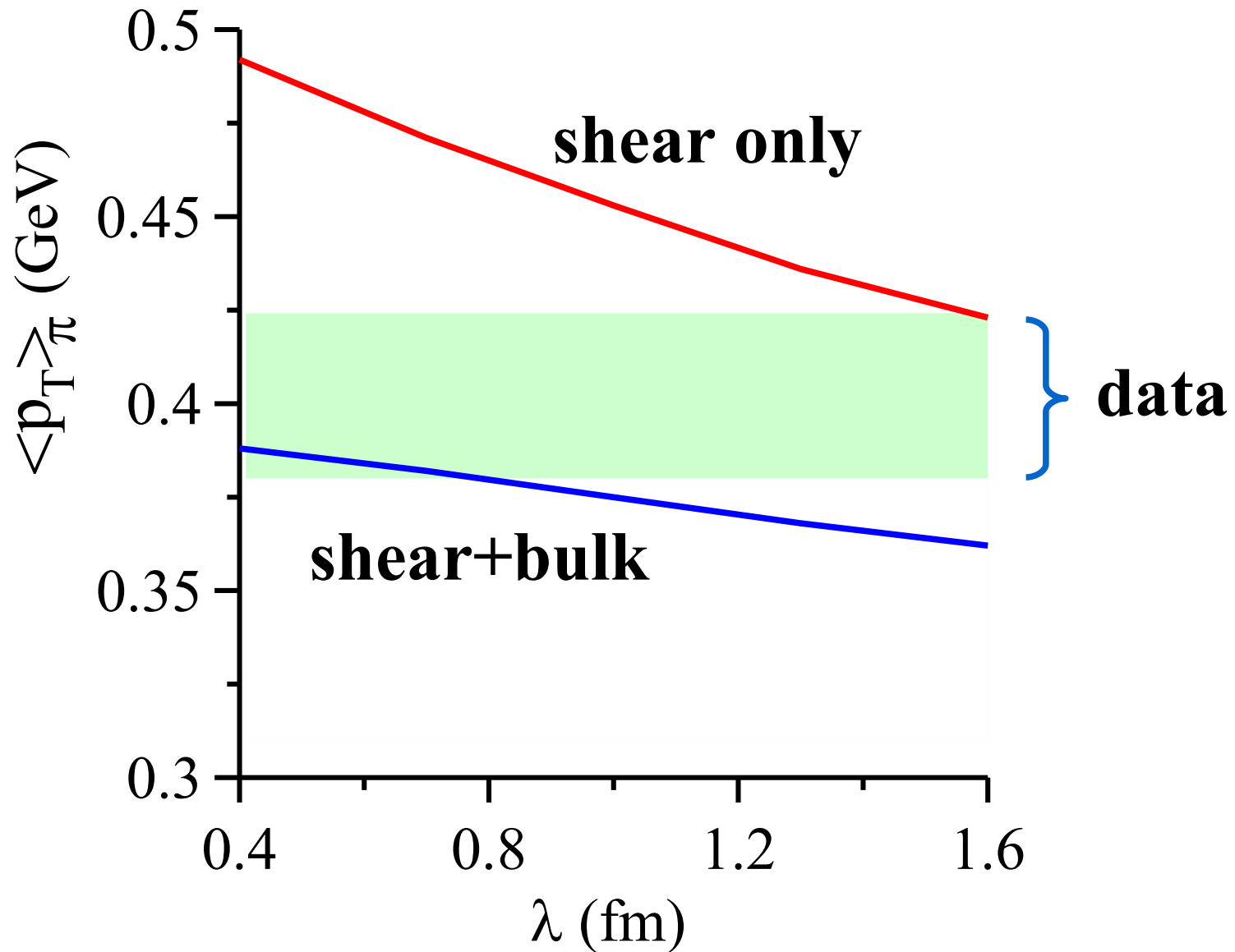


$\lambda = 0.8$  fm



$\lambda = 1.6$  fm

# Effect of initial thermalization scale



# What else can the spectra tell us?

Azimuthal information – Fourier series

Rapidity information – Legendre polynomials

Bzdak&Teaney, measured by ATLAS

Need of a more systematic expansion

# Can we extract more information from spectra?

$$f^{(i)}(x_T) = \frac{1}{\langle N \rangle_i} \frac{dN^{(i)}}{dx_T} \quad x_T = 2p_T / \langle \langle p_T \rangle \rangle$$

**Expansion using Laguerre polynomials,  $L_n$**

$$f^{(i)}(x_T) = x_T \exp(-x_T) \sum_{n=0}^{\infty} \frac{\ell_n^{(i)}}{n+1} L_n^{(1)}(x_T)$$

**characterize spectra**

**Can we extract more information from spectra?**  $x_T = 2p_T / \langle\langle p_T \rangle\rangle$

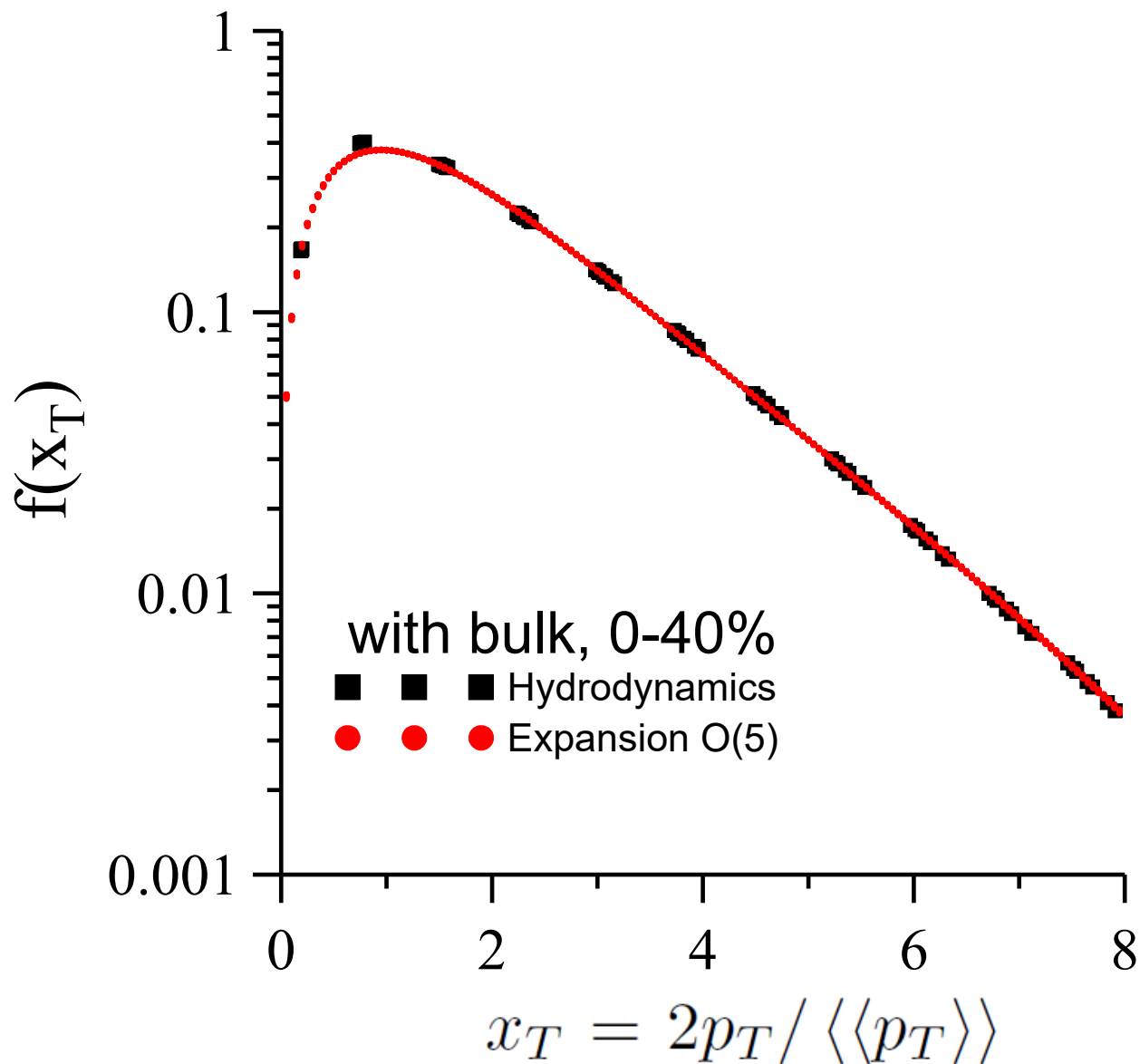
Expansion using Laguerre polynomials,  $L_n$

$$f^{(i)}(x_T) = x_T \exp(-x_T) \sum_{n=0}^{\infty} \frac{\ell_n^{(i)}}{n+1} L_n^{(1)}(x_T)$$

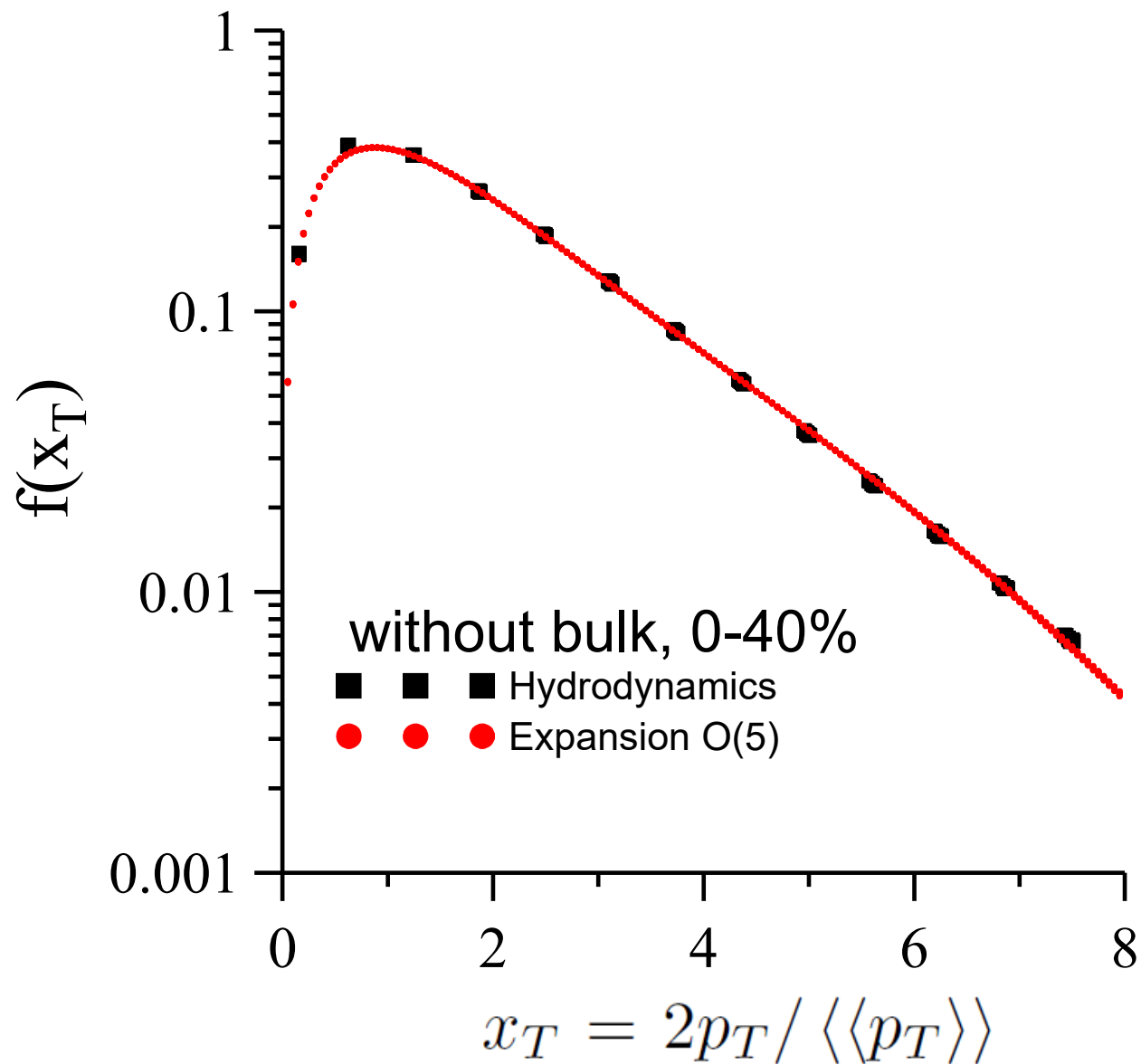
**Coefficients:**  $\ell_n^{(i)} = \langle L_n^{(1)} \rangle_{(i)}$

$$\langle A \rangle_{(i)} = \int dx_T f^{(i)}(x_T) A^{(i)}(x_T)$$

# Does it work?

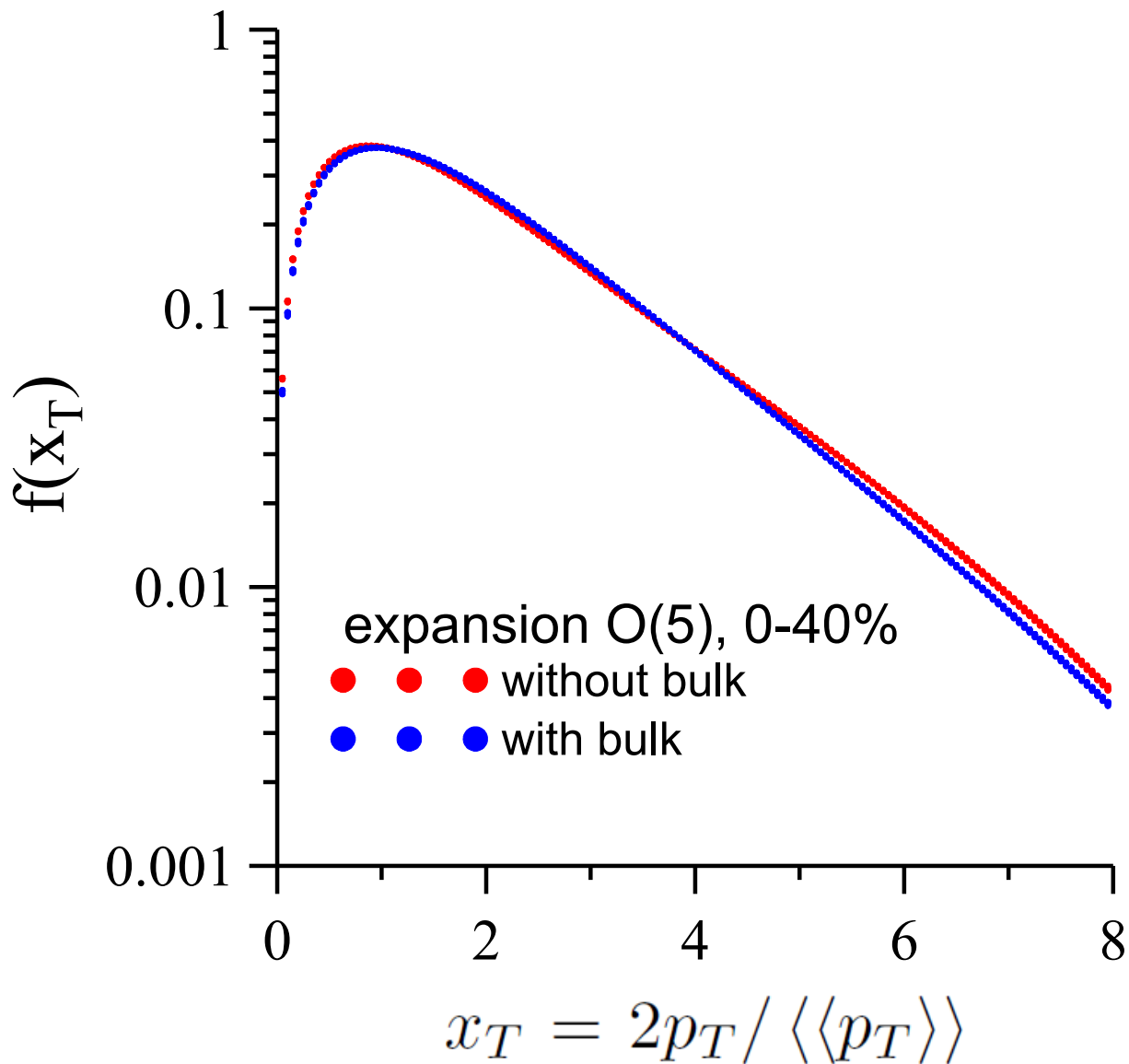


# Does it work?

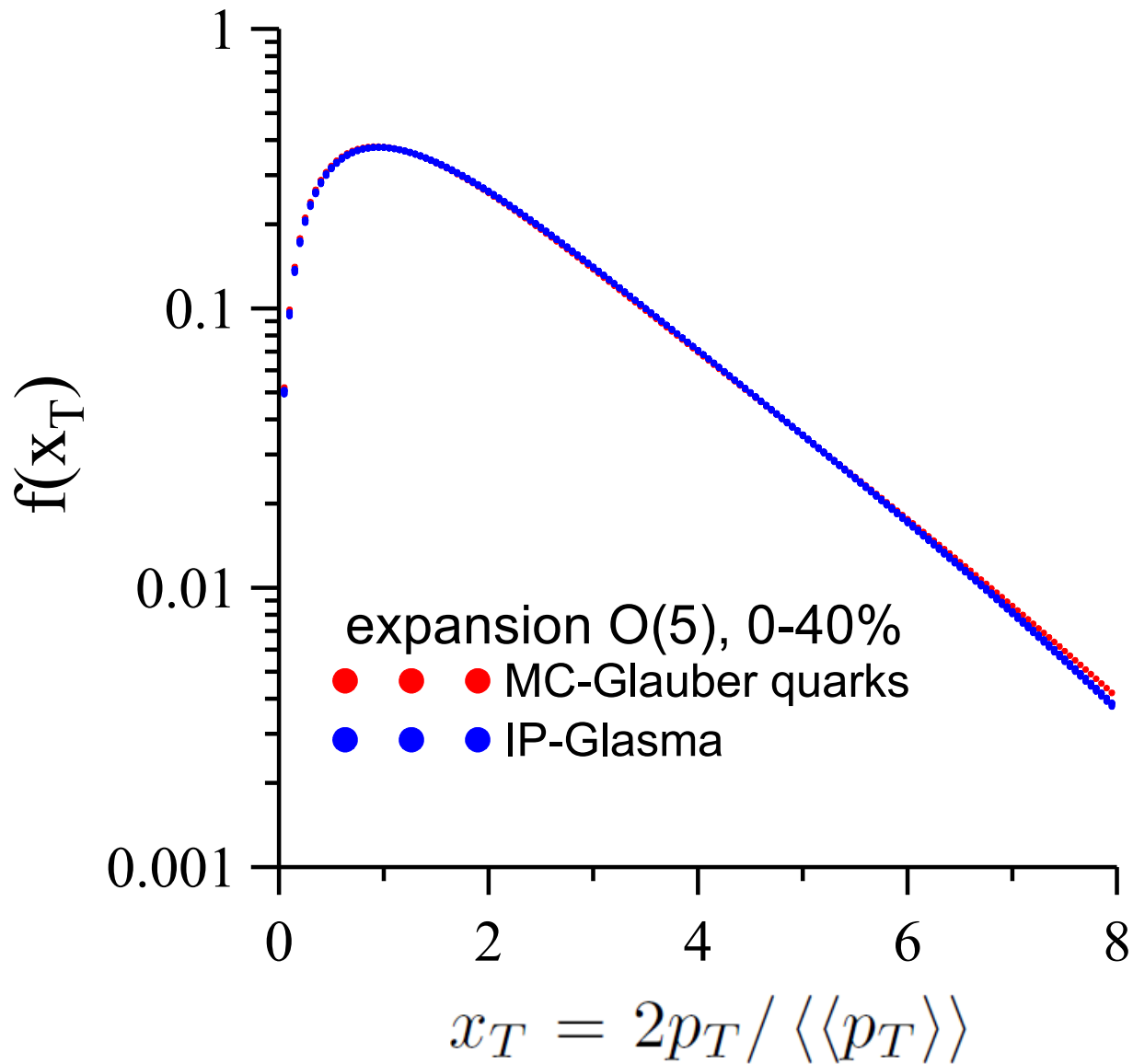




# Bulk viscosity changes the shape of f



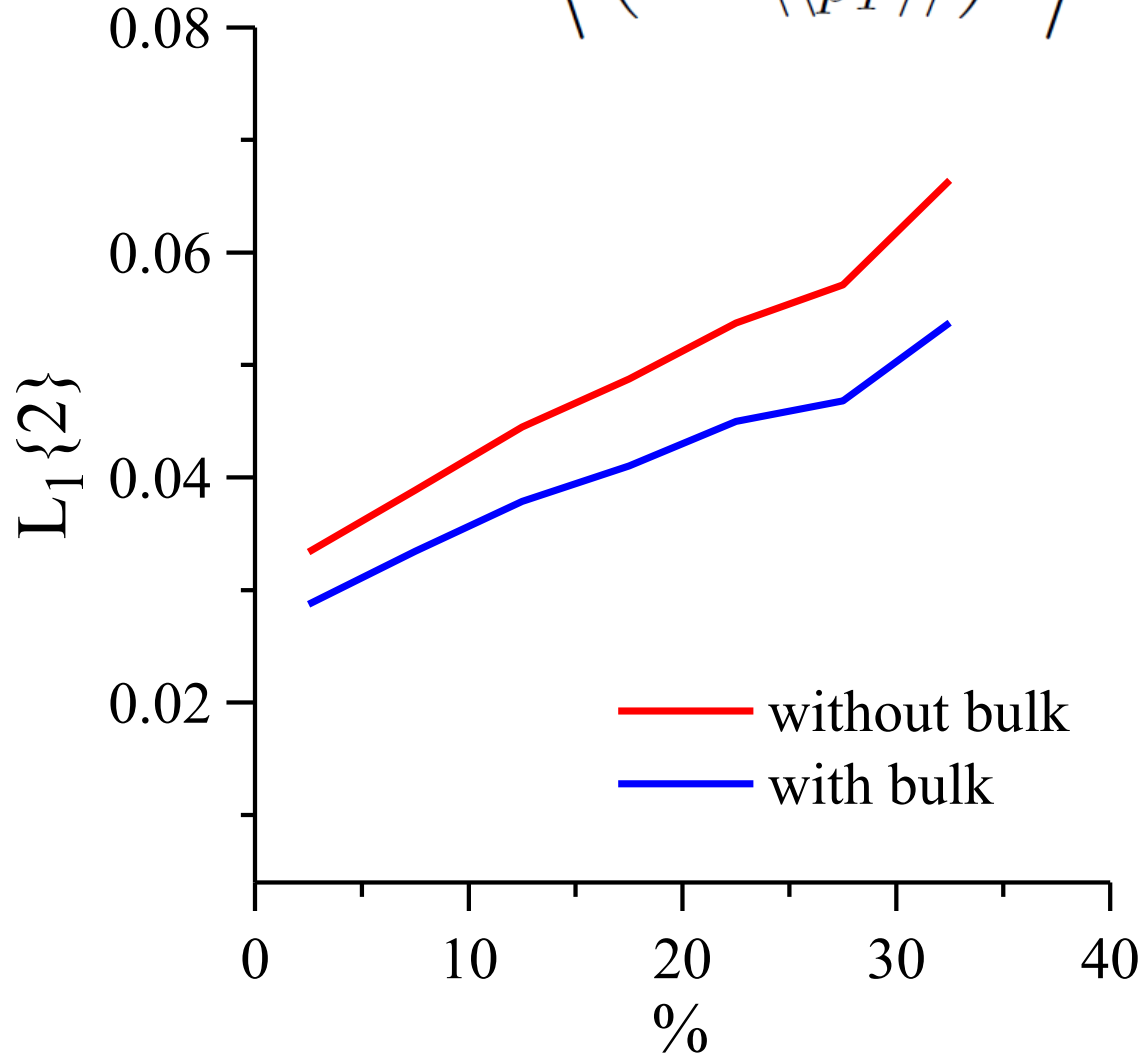
# Initial condition does not



# Expansion coefficients

$$L_n\{2\} = \sqrt{\langle \langle L_n \rangle^2 \rangle}$$

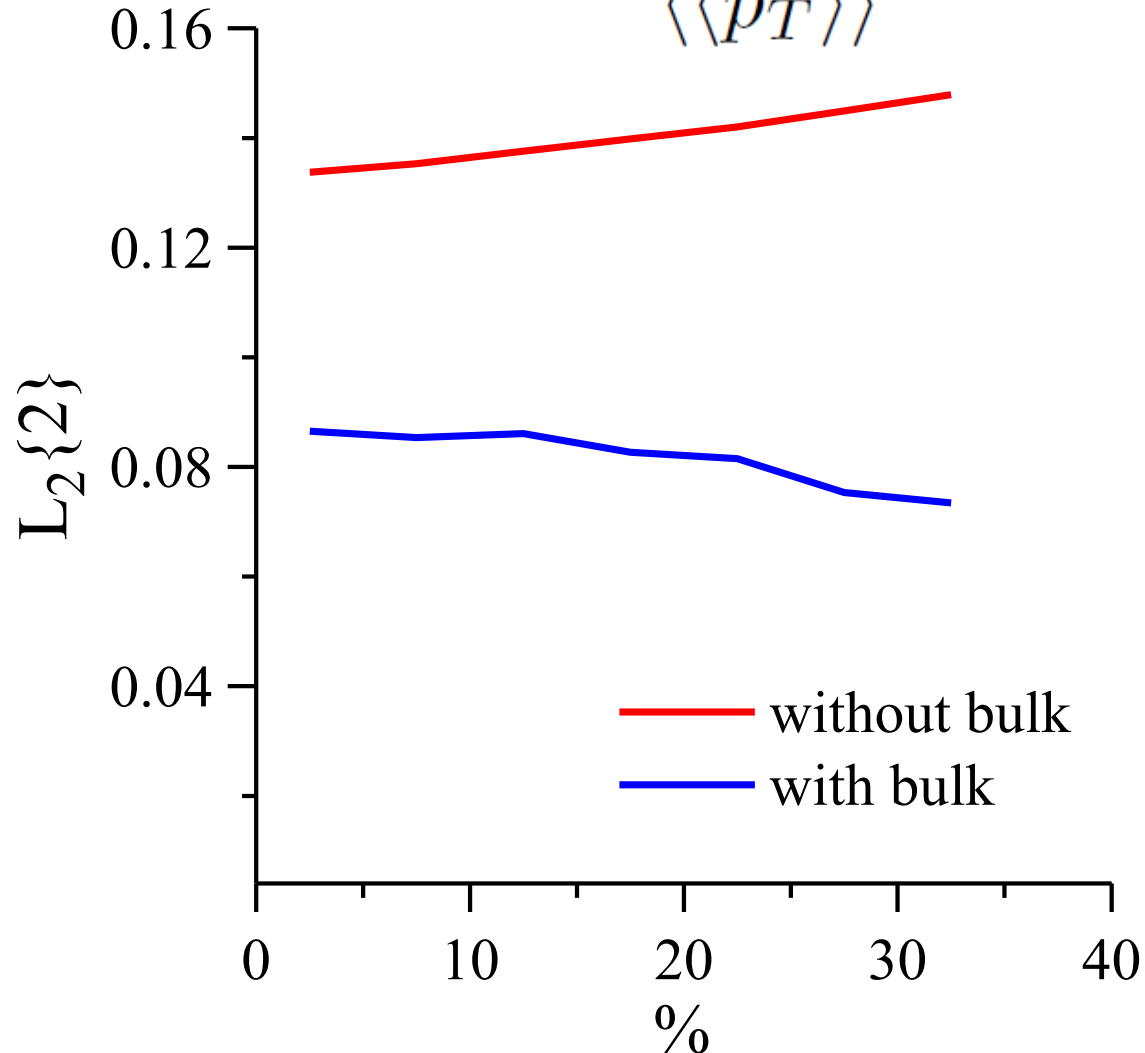
$$\langle \langle L_1 \rangle^2 \rangle = 4 \left\langle \left( 1 - \frac{\langle p_T \rangle}{\langle \langle p_T \rangle \rangle} \right)^2 \right\rangle$$



# Expansion coefficients

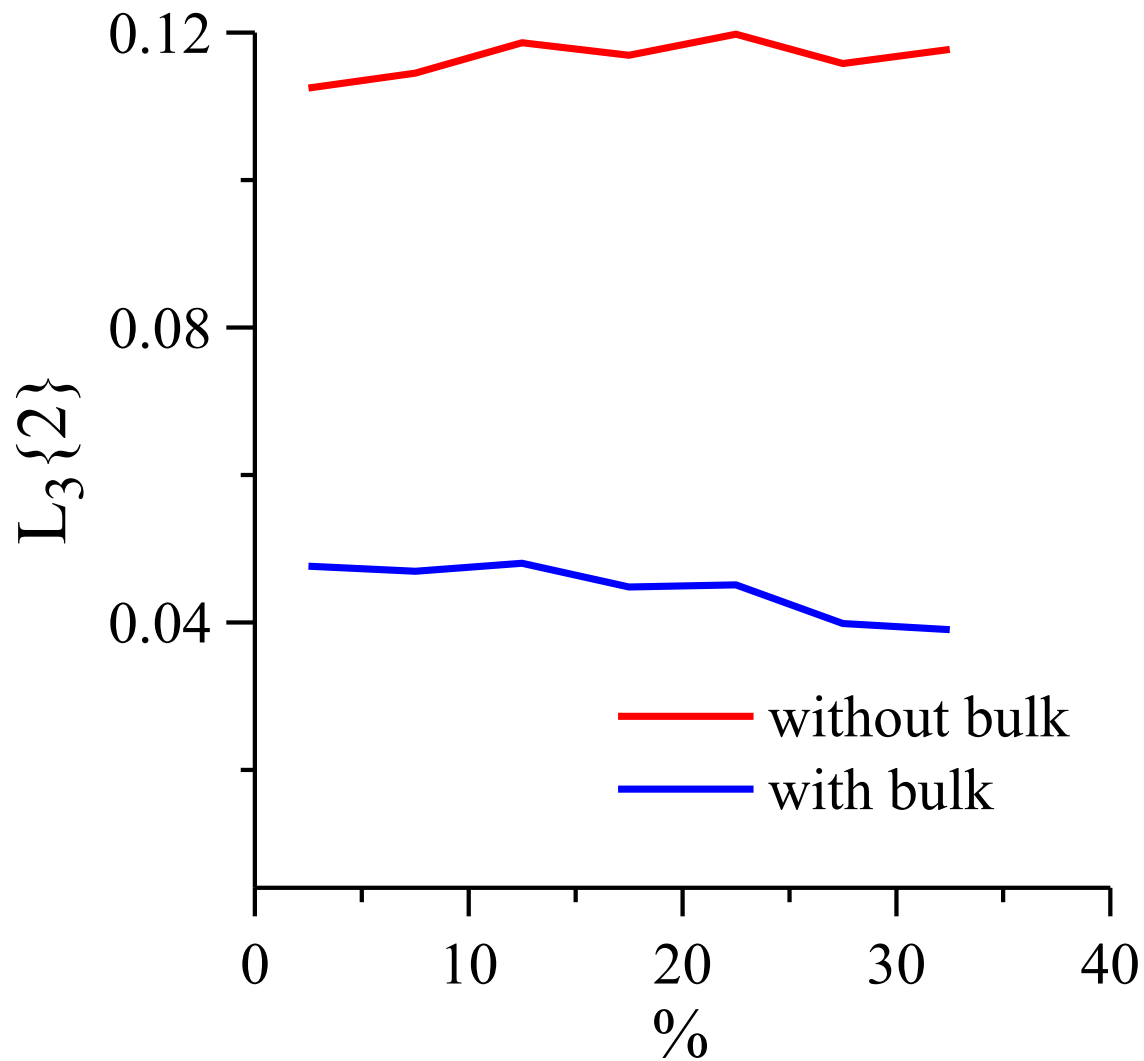
$$L_n\{2\} = \sqrt{\langle \langle L_n \rangle^2 \rangle}$$

$$\langle \langle L_2 \rangle \rangle = 2 \frac{\langle \langle p_T^2 \rangle \rangle}{\langle \langle p_T \rangle \rangle^2} - 3$$



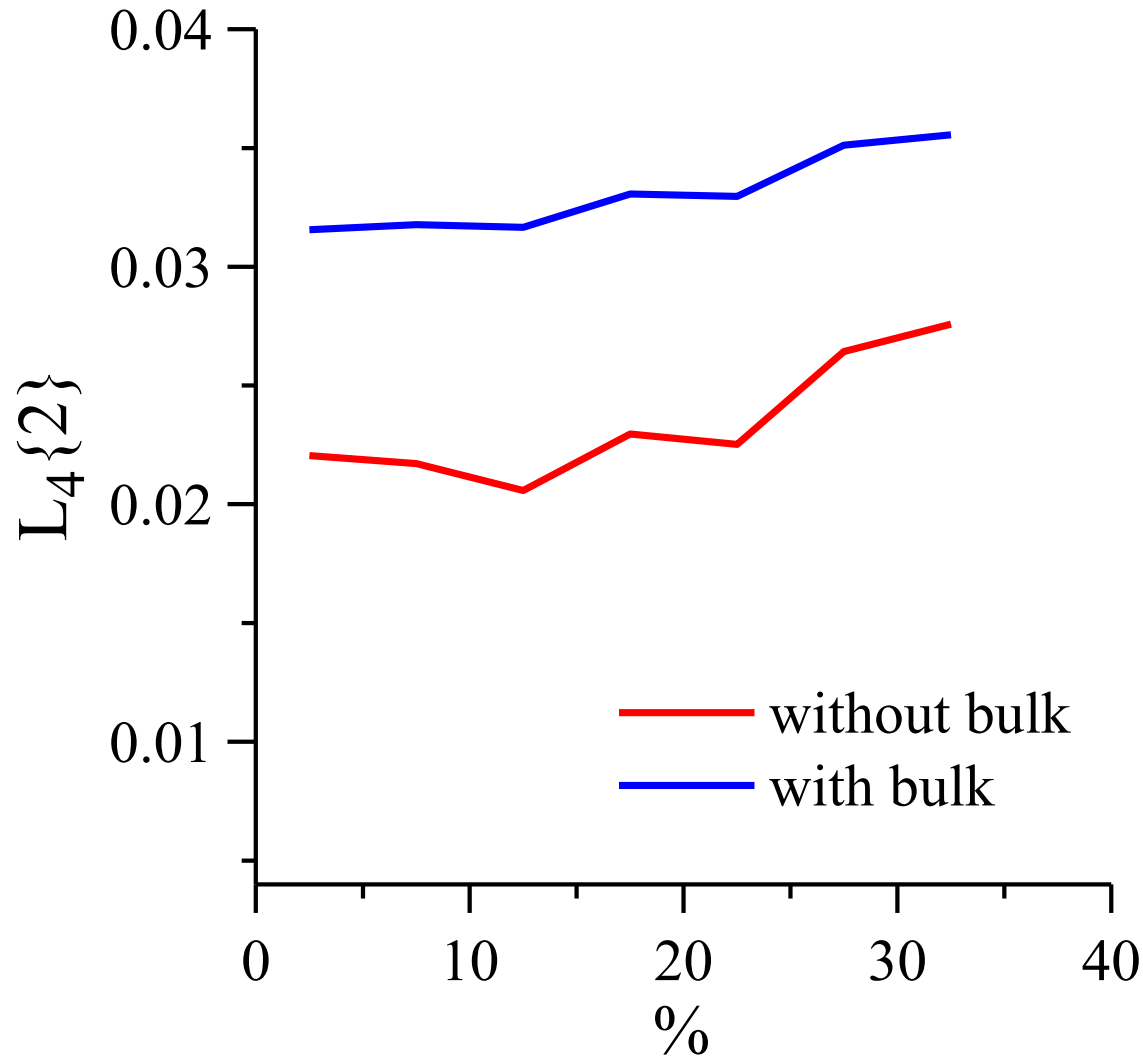
# Expansion coefficients

$$L_n\{2\} = \sqrt{\langle \langle L_n \rangle^2 \rangle}$$



# Expansion coefficients

$$L_n\{2\} = \sqrt{\langle \langle L_n \rangle^2 \rangle}$$



# Summary/conclusions

## **We studied the effect of bulk viscosity on the transverse momentum spectra**

- ✓ Bulk viscosity has a considerable effect on multiplicity, transverse momentum, and flow harmonics
- ✓ The mean transverse momentum is correlated to the initial energy density gradients
- ✓ Effect of bulk viscosity on spectra can be studied using Laguerre moments of the distribution function

**Happy Birthday Uli!**