



# HIC: Status of Chemical Equilibrium in Heavy Ion Collisions.

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# Outline

Use of Thermal Concepts in Heavy-Ion Collisions

Comparison of Chemical Freeze-Out Criteria

Disappearance of Maxima in Small Systems

Conclusion

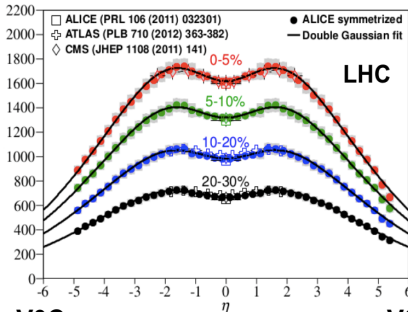
# Particle Multiplicity in Heavy Ion Collisions



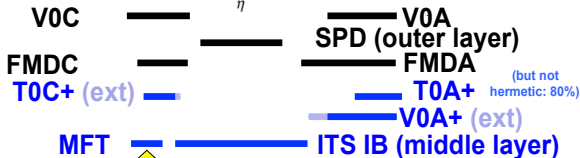
## Acceptance for charged particles



$\eta$  coverages  
for  $z_{\text{vtx}}=0$   
(shown at last  
AW)



Now:  
(T0 now shown)



This is (-3.6,-2.5), i.e. the MFT+MUON acc.



# Particle Multiplicity in Heavy Ion Collisions

About 24 000 particles are produced in a heavy ion collision at the LHC.

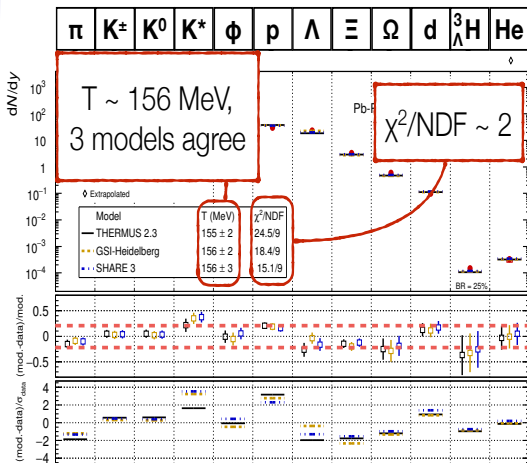
Hence: Use Concepts from Statistical Mechanics to analyze the final state

e.g. use Energy Density, Particle Density, Pressure, Temperature, Chemical Composition, ...

These concepts turn out to be useful at all energies, RHIC, SPS, GSI ...



AL



ALI-PREL-94600

N.B.

RHIC (STAR)

 $\sqrt{s} = 200$  GeV $\chi^2/\text{NDF} \sim 1$ 

Better fit in

60-80%,

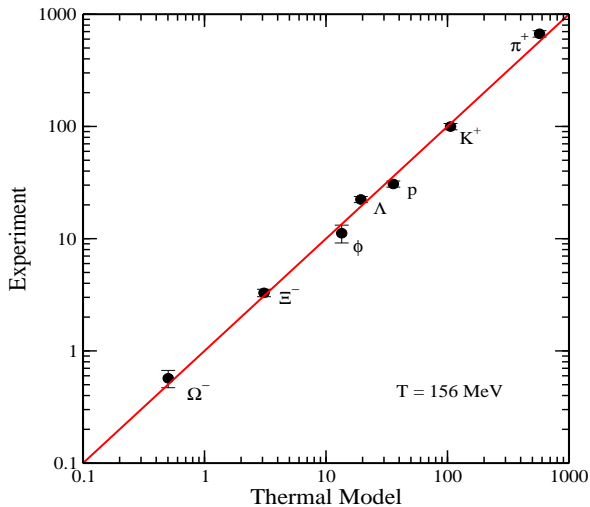
Petran et al, arXiv:1310.5108

Wheaton et al,

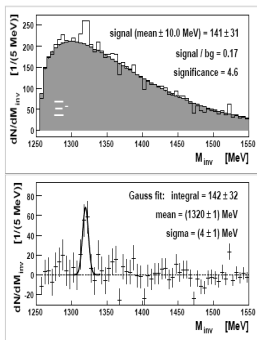
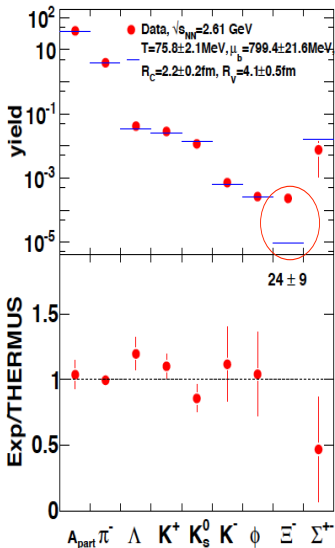
Comput.Phys.Commun, 180 84

Andronic et al, PLB 673 142

## ALICE



# Hadrons in Ar+KCl@1.76A GeV



Strong excess of the  $\Xi^-$

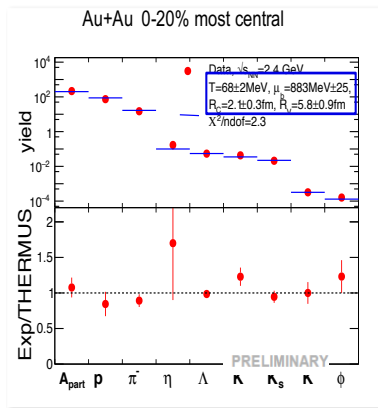
NN-threshold:

$$E_{beam} = 3.74 \text{ GeV} \rightarrow \sqrt{s} - \sqrt{s_{th}} = 630 \text{ MeV!}$$

# Comparing Au+Au data with a statistical model

Macroscopic description based on:

- Grand canonical ensemble  
( $T, \mu = \mu_B, \mu_S, \mu_Q, V$  and sometimes  $\gamma_S$ , usually  $\mu_s$  and  $\mu_Q$  are constrained)
- Strangeness canonical ensemble  
( $T, \mu = \mu_B, \mu_Q, R_c = R_v$ )
- Strangeness canonically suppressed at low temperatures, but not enough to explain data  $\rightarrow$  needs additional parameter:  $R_c < R_v$
- $\phi$  meson (hidden strangeness) not suppressed by  $R_c$  but strongly by  $\gamma_S$



## Implementation: THERMUS V2.3

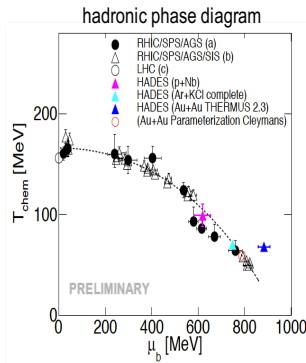
Wheaton & Cleymans Comp. Phys. Com. 180 (2009)

Hadron yields described by  $T, \mu_B, R_v$ , and  $R_c$

- $\rightarrow$  rather large values for  $T$  and  $\mu_B$
- $\rightarrow \gamma_S$  instead of  $R_c$  delivers similar results, but undershoots the  $\phi$  yield



# HADES in the phase diagram



$T$  and  $\mu_b$  higher than expected from parameterization and universal freeze-out line ( $E/B=1\text{GeV}$ )

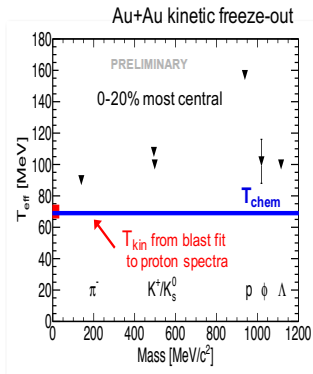
### Systematics of freeze-out points:

Andronic, PBM, J. Stachel, NPA 772 (2006)

Cleymans, Oeschler et al., PRC 73 (2006)

### Parameterization of $T$ and $\mu_b$ :

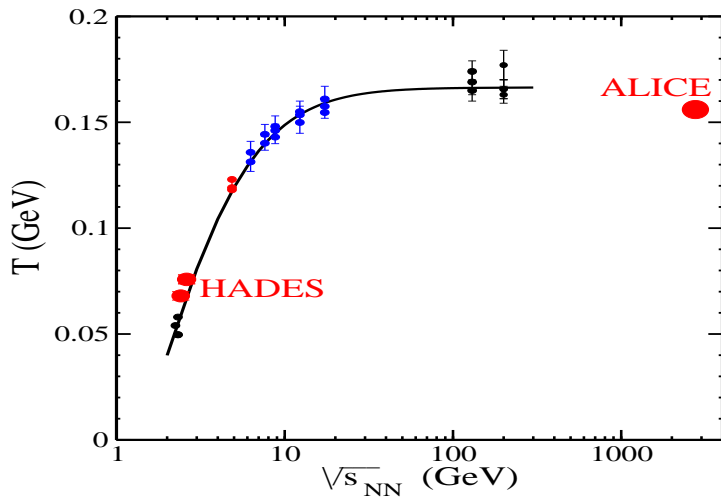
Cleymans, Oeschler et al., PRC 73 (2006)

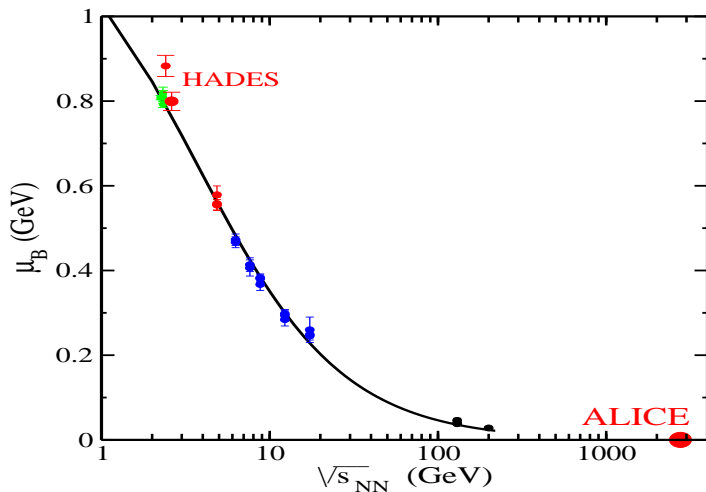


We find  $T_{\text{kin}} \approx T_{\text{chem}} = 68 \text{ MeV}$

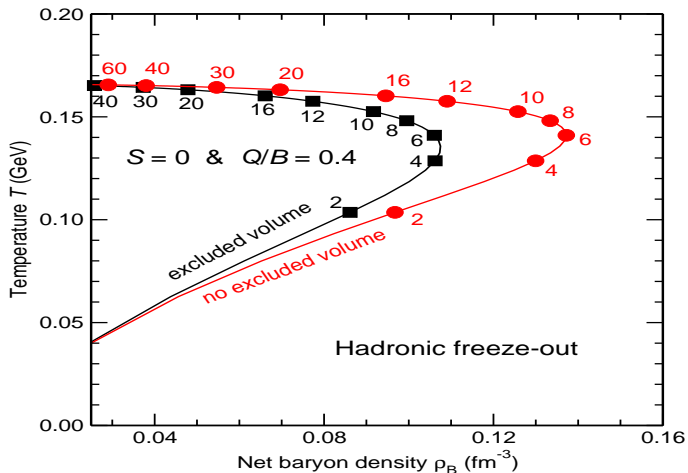
Todo: add fragments (d,t,He) to further constrain  $T_{\text{kin}}$

# Chemical Freeze-Out Temperature

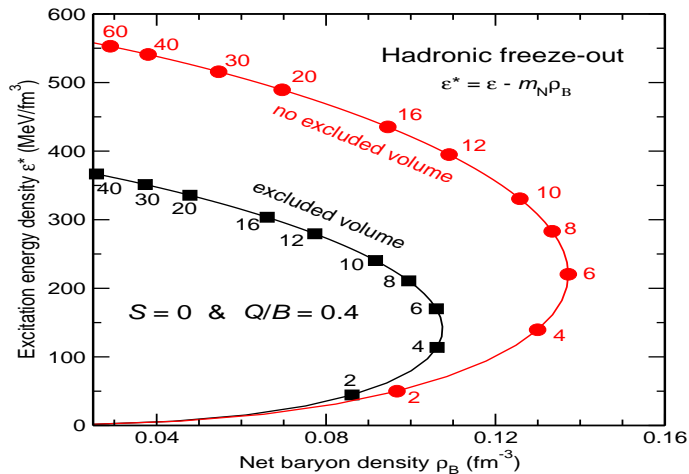


Chemical Freeze-Out  $\mu_B$ 

# Unexpected Results



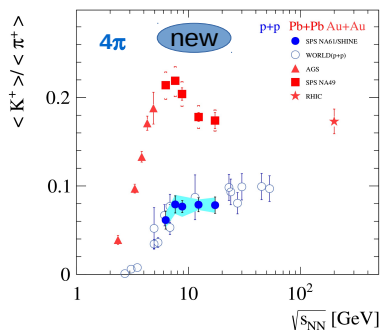
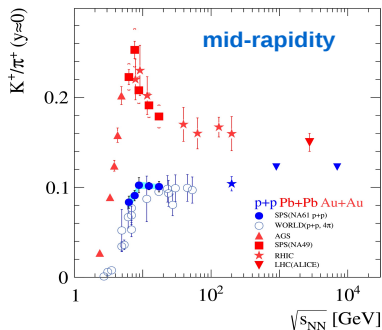
# Unexpected Results



# Unexpected Results

K. Grebieszko (NA61/SHINE) talk at CPOD2016:

Maximum in the  $K^+/\pi^+$  ratio disappears in small systems



To analyze the particle ratios use:

- the Wroblewski factor
- $s/T^3 = 7$  describes chemical freeze-out

## Strangeness in Heavy Ion Collisions vs Strangeness in pp - collisions

Use the Wroblewski factor

$$\lambda_s = \frac{2 \langle s\bar{s} \rangle}{\langle u\bar{u} \rangle + \langle d\bar{d} \rangle}$$

This is determined by the number of **newly** created quark – anti-quark pairs and **before** strong decays, i.e. before  $\rho$ 's and  $\Delta$ 's decay.

Limiting values :

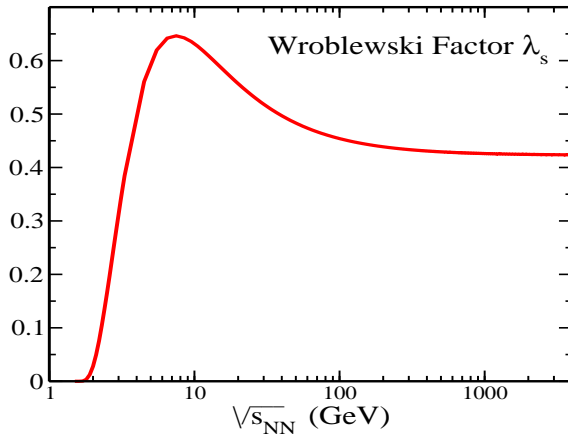
$\lambda_s = 1$  all quark pairs are equally abundant, SU(3) symmetry.

$\lambda_s = 0$  no strange quark pairs.

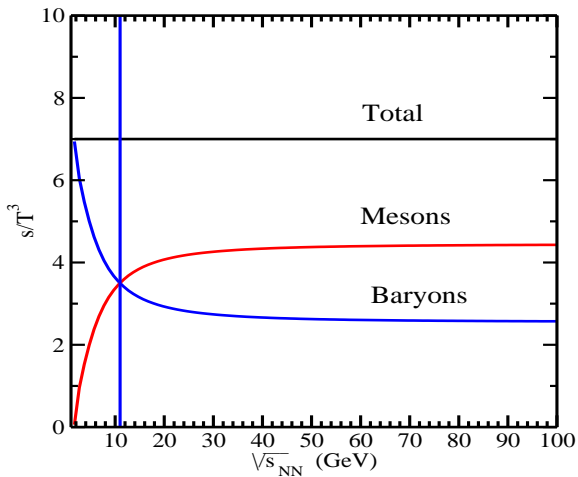




# Wroblewski Factor



$$s/T^3$$



## J.C., H. Oeschler, K. Redlich, S. Wheaton, Phys. Lett. B615 (2005) 50-54

In the statistical model a rapid change is expected as the hadronic gas undergoes a transition from a baryon-dominated to a meson-dominated gas. The transition occurs at a

- temperature  $T = 151$  MeV,
- baryon chemical potential  $\mu_B = 327$  MeV,
- energy  $\sqrt{s_{NN}} = 11$  GeV.

In this region the interplay between temperature and baryon chemical potential leads to peaks in the  $\Lambda / \langle \pi \rangle$ ,  $K^+ / \pi^+$ ,  $\Xi^- / \pi^+$  and  $\Omega^- / \pi^+$  ratios **which occur at different beam energies.**

P. Braun-Munzinger, J.C., H. Oeschler, K. Redlich, Nucl. Phys. A697 (2002) 902.



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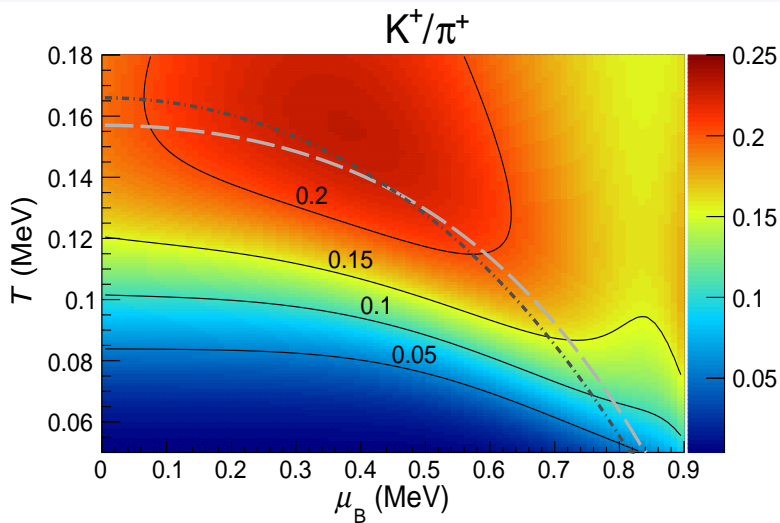
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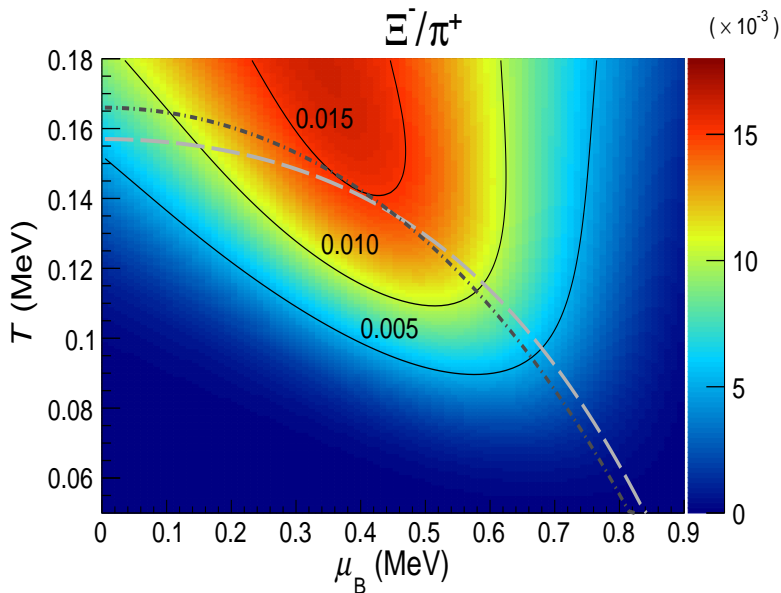
P. Braun-Munzinger, J.C., H. Oeschler, K. Redlich, Nucl. Phys. A697 (2002) 902.

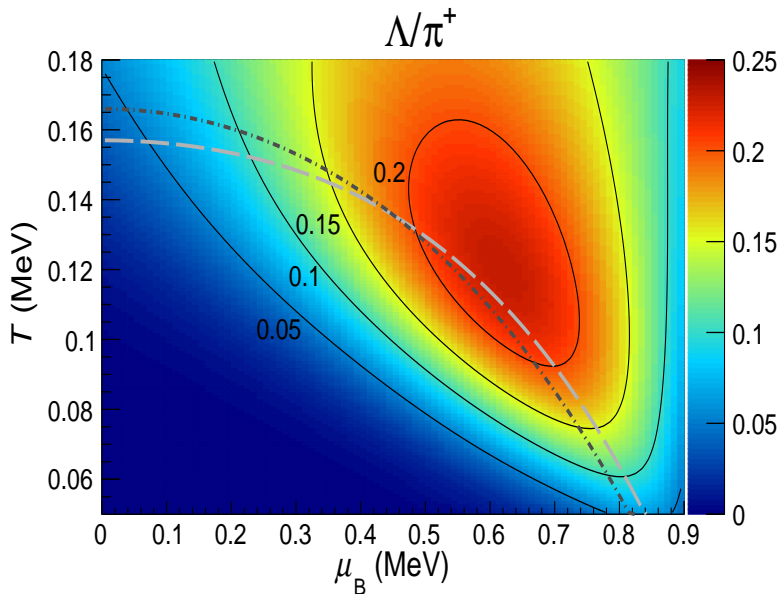




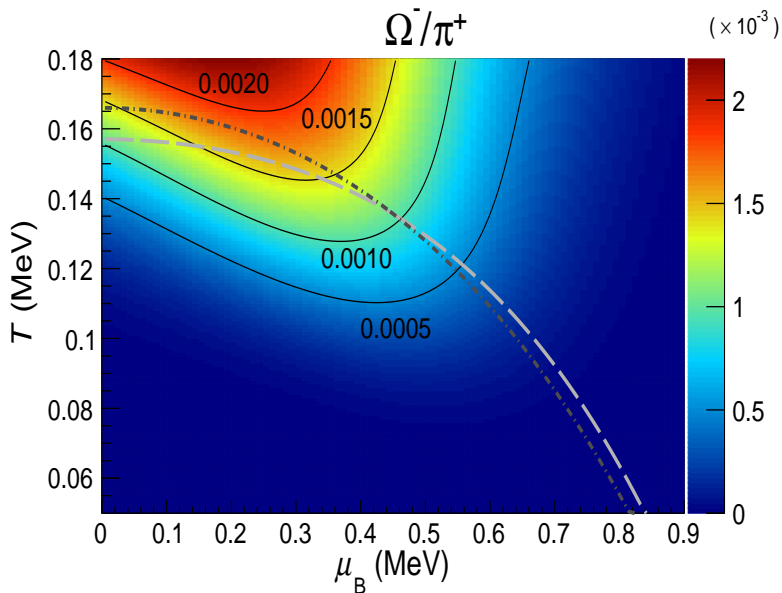
J.C., B. Hippolyte, H. Oeschler, K. Redlich, N. Sharma arXiv:1603.09553

V. Vovchenko, V.V. Begun, M.I. Gorenstein, arXiv:1512.08025[nucl-th]









## Small systems.

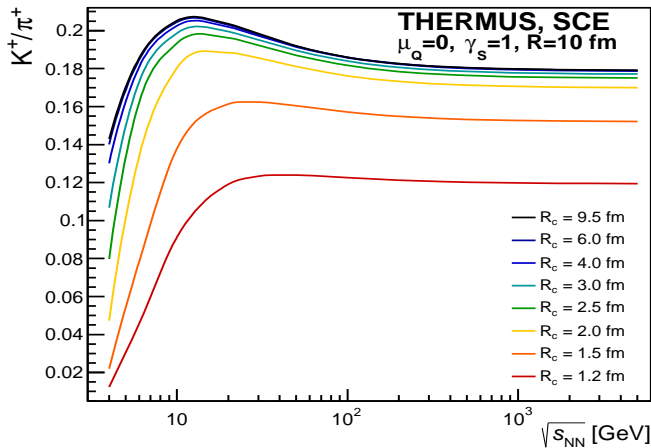
- Use the canonical ensemble with strangeness conservation.
- Introduce two volumes: global volume and a strangeness correlation volume .
- Reduce the strangeness correlation volume to describe small systems.

J.C., B. Hippolyte, H. Oeschler, K. Redlich, N. Sharma  
arXiv:1603.09553

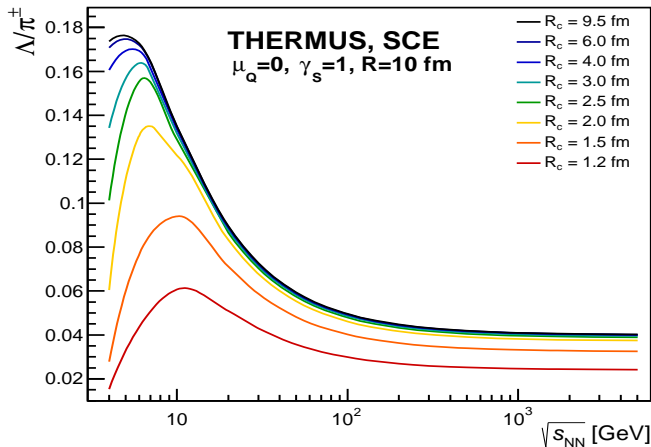
S. Hamieh, K. Redlich and A. Tounsi, Phys. Lett. B486 (2000) 61

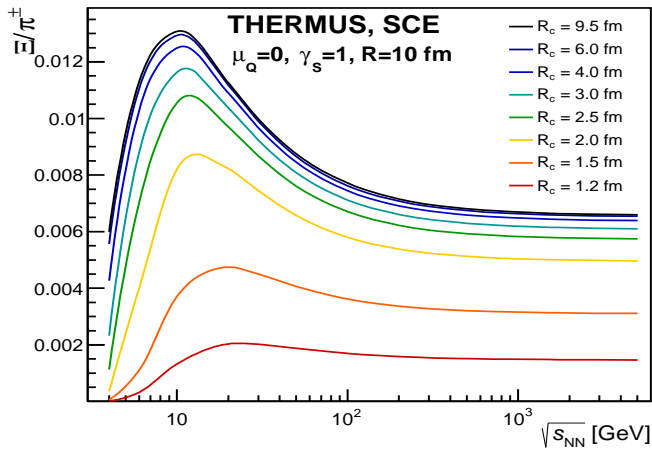


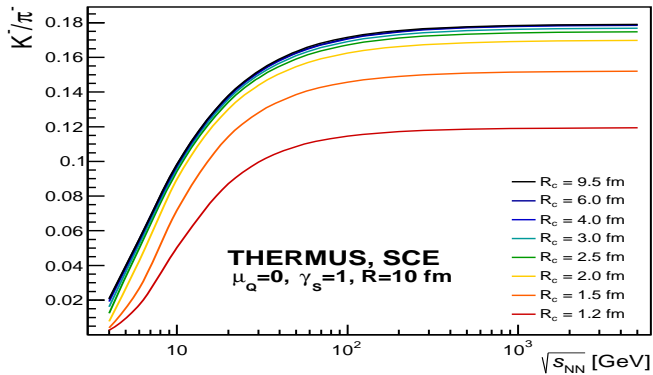
# Maximum in $K^+/\pi^+$ ratio disappears



# Maximum in $\Lambda/\pi^+$ ratio survives







## Conclusions

- Maximum in  $K^+/\pi^+$  ratio disappears for small systems,
- A small maximum in  $\Lambda/\pi$  ratio **SURVIVES** for small systems,

If this is confirmed experimentally then a hadronic scenario explains the behaviour seen in the hadronic ratios and there is no need for other mechanisms.

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**Uli, many happy returns.**