

Simulation of photon-assisted streamers

(a short description of the phenomenon in detectors... with no conclusions)

Diego Gonzalez Diaz (Uludag University)

Materials used

Nuclear Instruments and Methods in Physics Research A305 (1991) 91–110
North-Holland

Feedback and breakdown in parallel-plate chambers

P. Fonte ^{a,b}, V. Peskov ^{b,c} and F. Sauli ^b

^a *LIP-Coimbra, Coimbra, Portugal*

^b *CERN, Geneva, Switzerland*

^c *WorldLab, Lausanne, Switzerland*

IEEE TRANSACTIONS ON NUCLEAR SCIENCE, VOL. 43, NO. 3, JUNE 1996

A Model of Breakdown in Parallel-Plate Detectors

P. Fonte

Nuclear Instruments and Methods in Physics Research A 661 (2012) S168–S171

RPC simulation in avalanche and streamer modes using transport equations
for electrons and ions

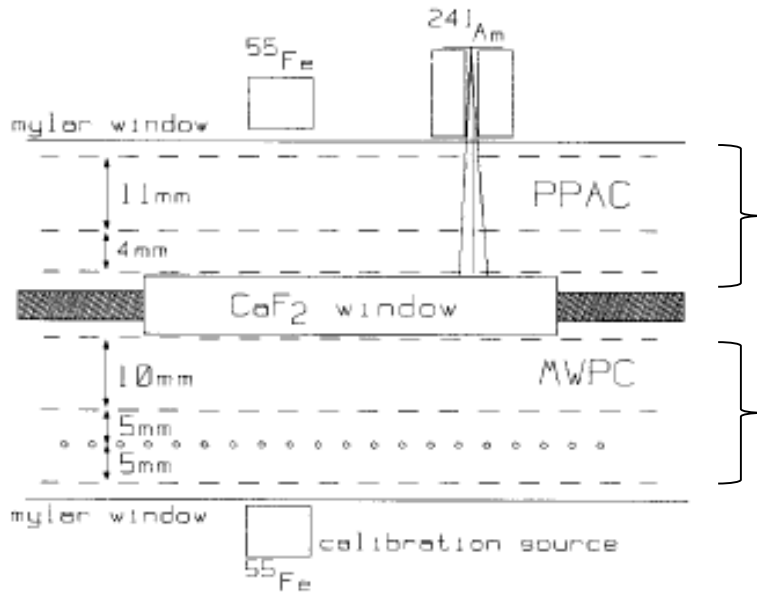
Ahmad Moshaii *, Larousse Khosravi Khorashad, Mehdi Eskandari, Somayeh Hosseini

Department of Physics, Tarbiat Modares University, P.O. Box 14115-175, Tehran, Iran

Calculation of streamer development in MPGDs in an axisymmetric hydrodynamic model

P. Fonte: <https://indico.cern.ch/event/89325/session/0/contribution/16>

Some 'bread-and-butter' phenomenology



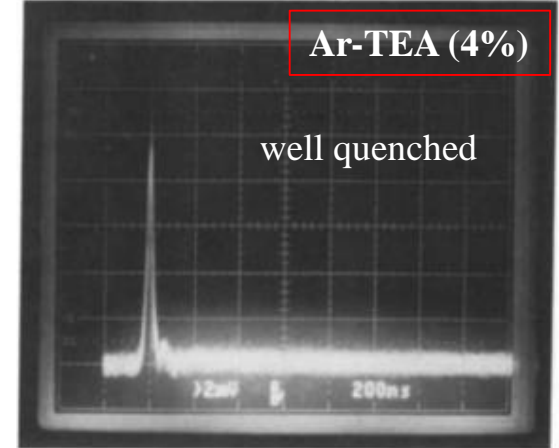
used to study
breakdown

used to measure
light emitted

Fig. 1. Schematic view of the experimental setup.

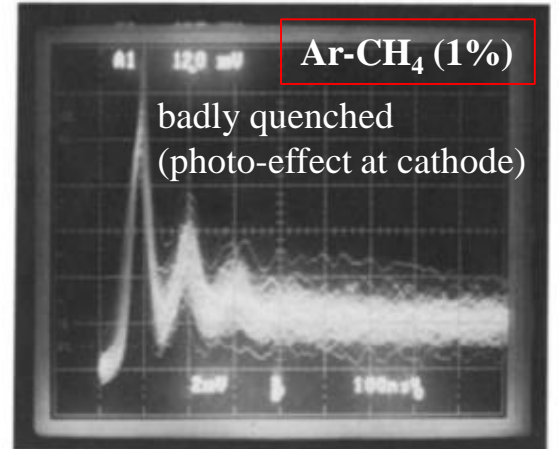
Ar-TEA (4%)

well quenched



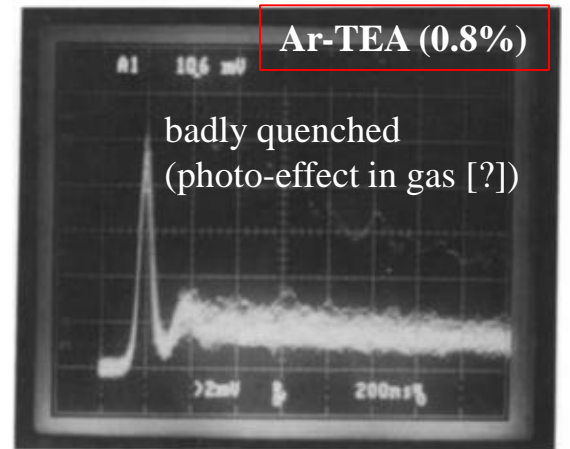
Ar-CH₄ (1%)

badly quenched
(photo-effect at cathode)



Ar-TEA (0.8%)

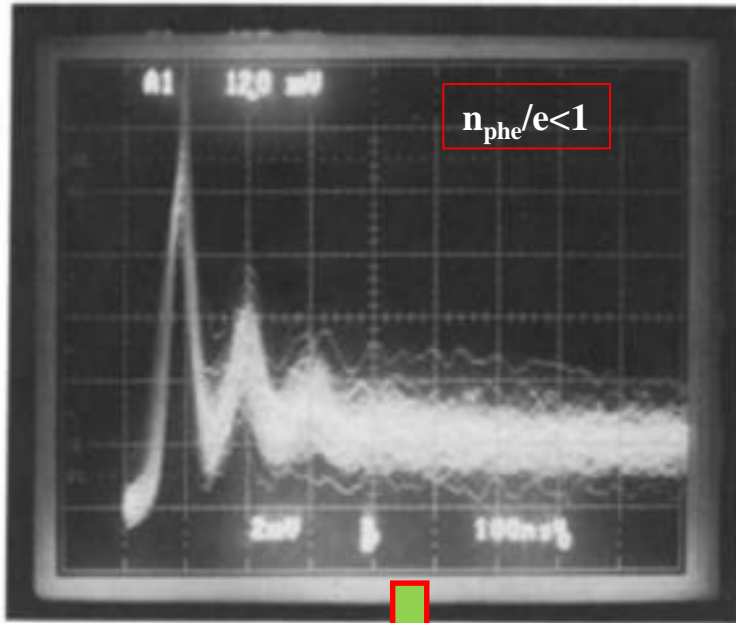
badly quenched
(photo-effect in gas [?])



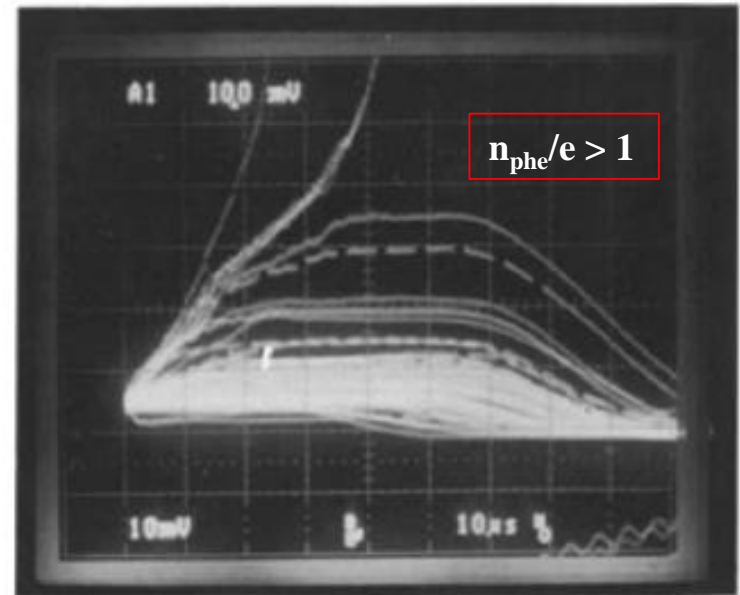
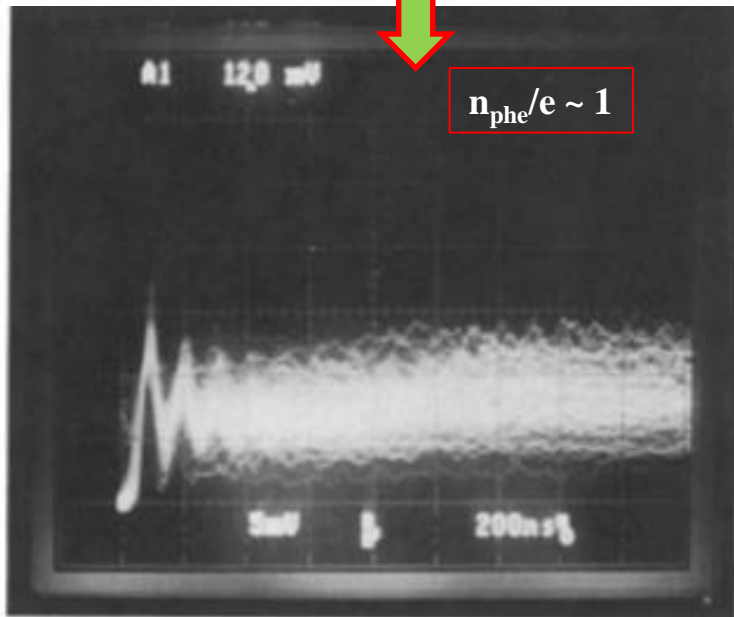
'Slow' breakdown, I

popularly known as:

- 'Slow' breakdown
- 'Generation' breakdown
- 'Paschen/Townsend' breakdown
- 'Photon-feedback'



Ar-CH₄ (~1%)

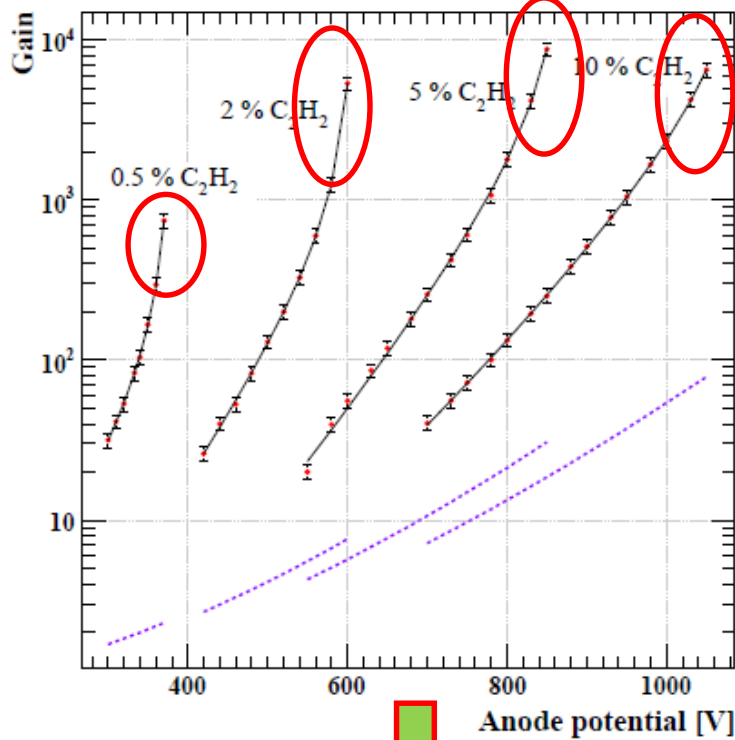


In the same paper... the parameter $n_{\text{phe}}/e = \beta$ is obtained from the gain deviations from an ~exponential law:

$$G(E) \cong \exp(A + BE)$$

$$G := G / (1 - \beta G)$$

'Slow' breakdown, II

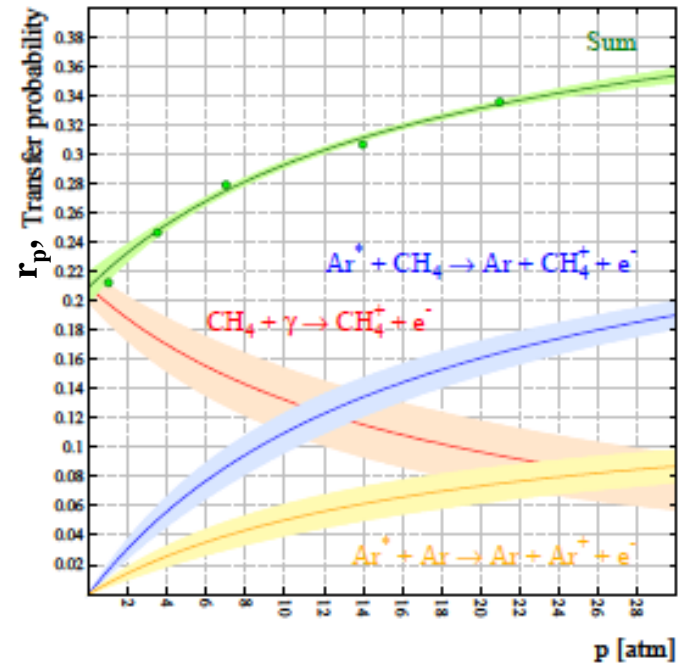
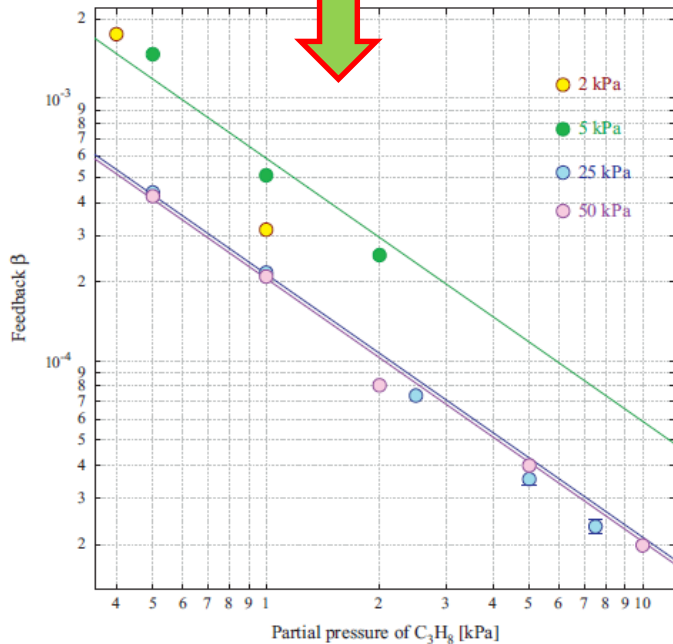
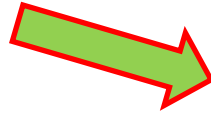


Modern way:

Number of fit parameters reduced from 3 to 2, and they $[r_p, \beta]$ can be (usually) interpreted.

$$G = \exp \int_{\text{tube}}^{\text{anode}} dr \alpha(E(r)) \frac{\sum v_i^{\text{ion}}(E(r)) + \sum r_i v_i^{\text{exc}}(E(r))}{\sum v_i^{\text{ion}}(E(r))}$$

$$G := G / (1 - \beta G)$$



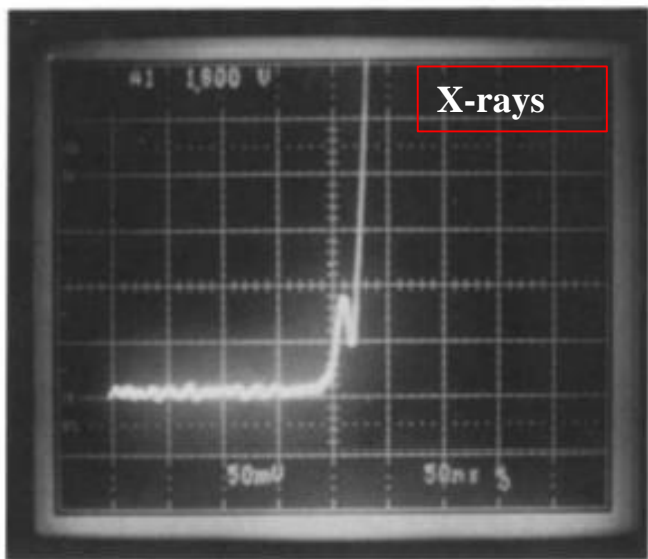
O. Sahin et al., JINST(2010)P05002

O. Sahin et al., NIM A 718(2013)432

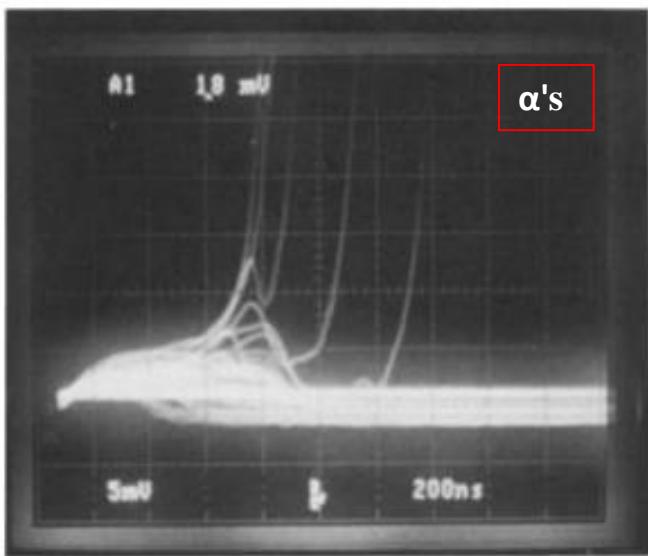
'Fast' breakdown

popularly known as:

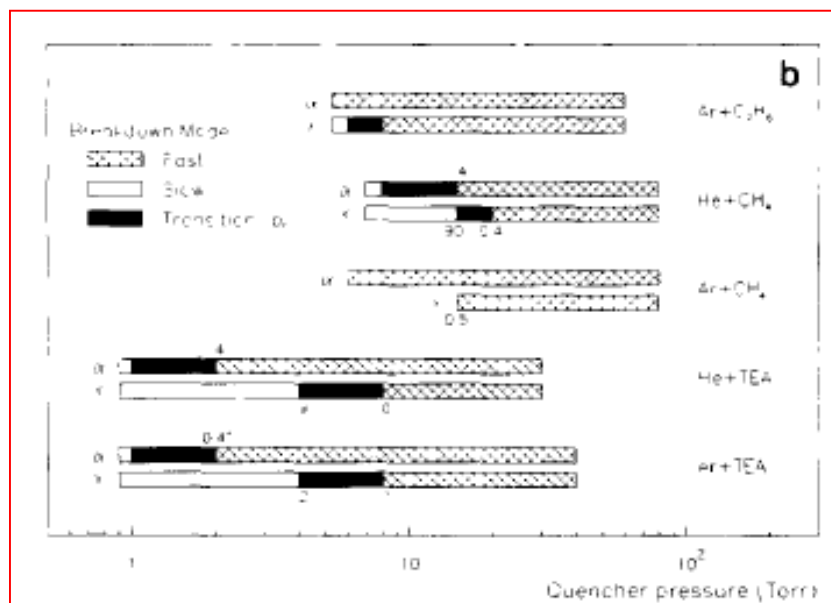
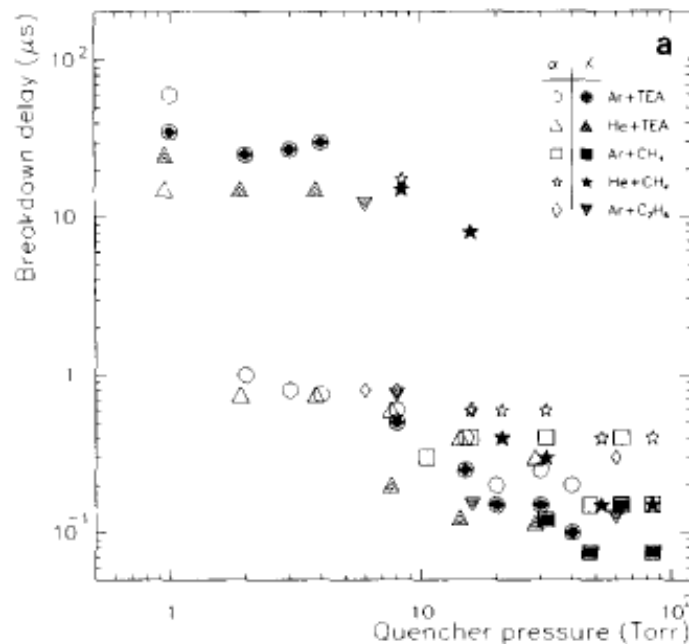
- 'Streamer' mechanism



X-rays



α 's



'Fast' breakdown - experimental evidence in detectors

Very fast process featuring a "precursor" pulse

[RAE64]

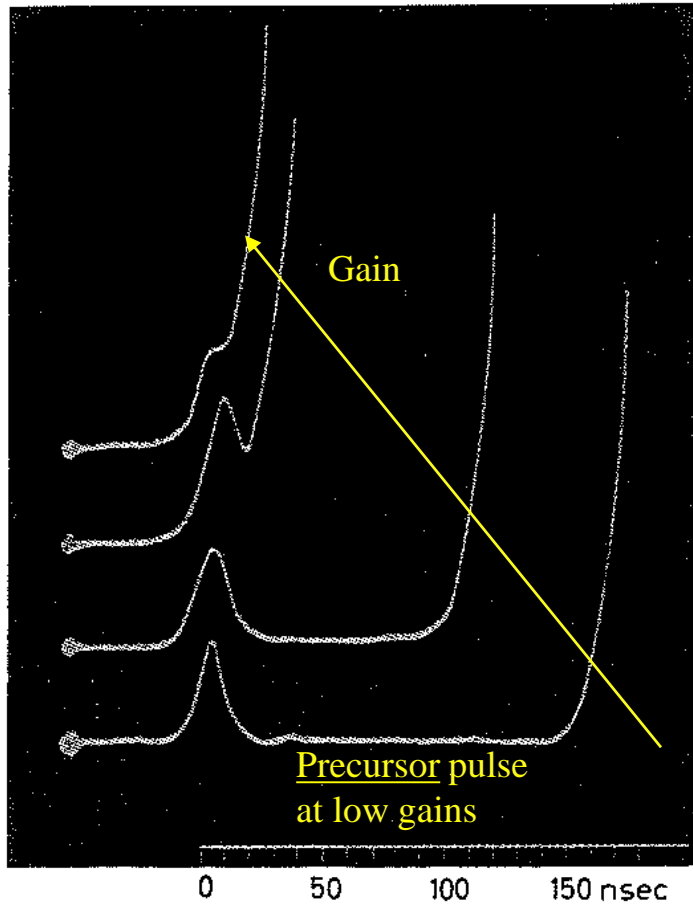
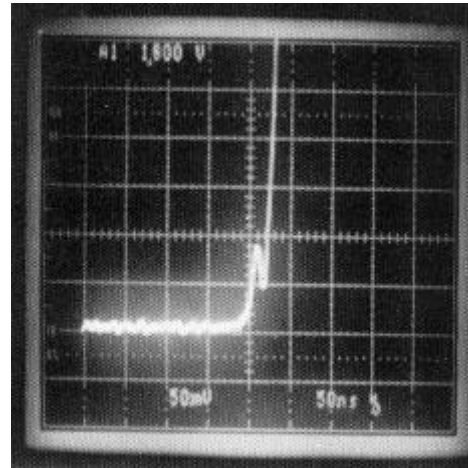


Figure 5.14. Current oscillograms of static breakdown in methylal. Optical method. $E/p = 64.4$, $pd = 230$ Torr cm, $d = 0.8$ cm, $T_- = 90$ nsec $RC = 5$ nsec³⁶

PPAC

[FON91]

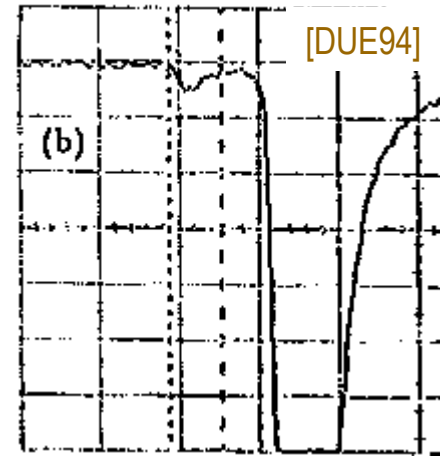


A signature of low-gain cathode streamer-only breakdown

It certainly exists, but it may not be the only mechanism, or it may be the result of different physical processes.

RPC

[DUE94]



single-wire (SQS mode)

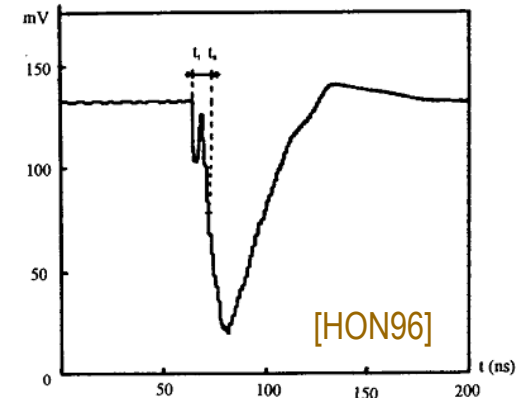


Fig. 1. The pulse shape of the SQS electrical signal $V = 2.45$ kV. Methylal/(Methylal + Ar) = 16.6%.

'Fast' breakdown ('classical interpretation')

$$\frac{\partial n_e(x,t)}{\partial t} + W_e \frac{\partial n_e}{\partial x} = S + \left(\alpha |W_e| - \frac{\partial W_e}{\partial x} \right) n_e; \quad (1)$$

$$\frac{\partial n_i(x,t)}{\partial t} = S + \alpha |W_e| n_i; \quad (2)$$

$$\frac{\partial n_{ph}(x,t)}{\partial t} = \delta |W_e| n_e; \quad (3)$$

$$S(x,t) = \frac{Q}{2\lambda} \int_{-\infty}^{\infty} \frac{\partial n_{ph}(x',t)}{\partial t} \Omega(x-x') \cdot e^{-|x-x'|/\lambda} dx'; \quad (4)$$

$$\Omega(x-x') = \frac{1}{2} \left(1 - \frac{(x-x')}{\sqrt{R_0^2 + (x-x')^2}} \right)$$

photon-feedback term

1D hydro equations for electrons and ions with photon-feedback term

p	20 Torr	
A, B, k	Corresponding to the mixture Ar + TEA	
W_{max}	5 cm/ μ s	
E_v	300 V/cm	
M	10^{-6}	?
λ	500 μ m	?
N_{e0}	100 electrons	
R_0	300 μ m	?
σ_0	300 μ m	
X_0	-900 μ m	
G_0	4×10^6	
d	4 mm.	

$$\frac{E(x)}{2\pi} = \int_{-x}^0 \rho(x+x')(-1-g(x')) dx' + \int_0^{d-x} \rho(x+x')(1-g(x')) dx';$$

$$g(x') = \frac{x'}{(x'^2 + R_0^2)^{1/2}}; \quad \rho(x) = e(n_i(x) - n_e(x))$$

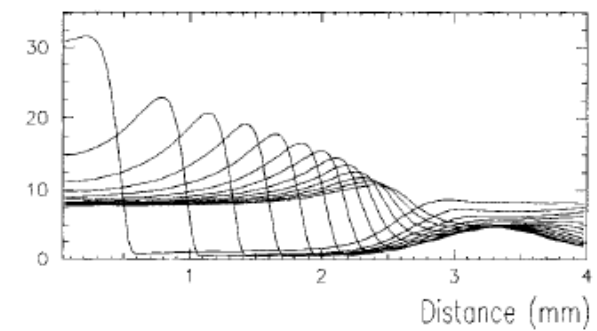
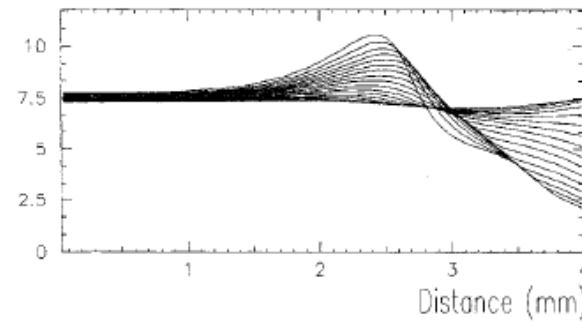
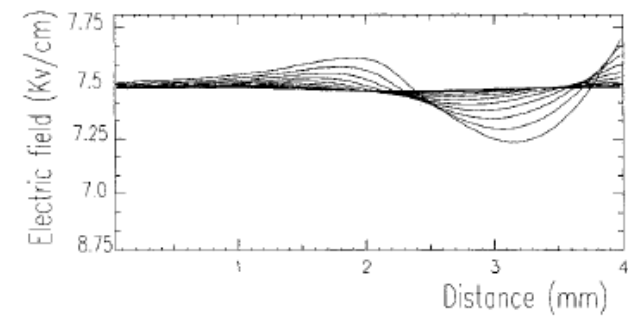
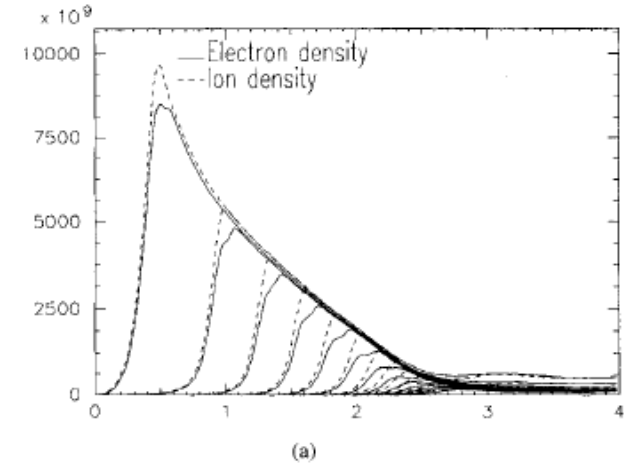
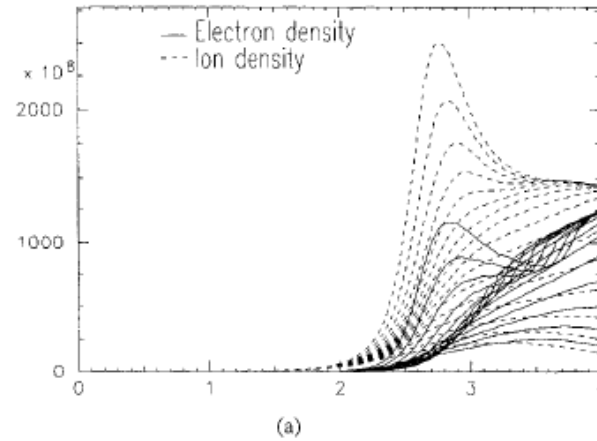
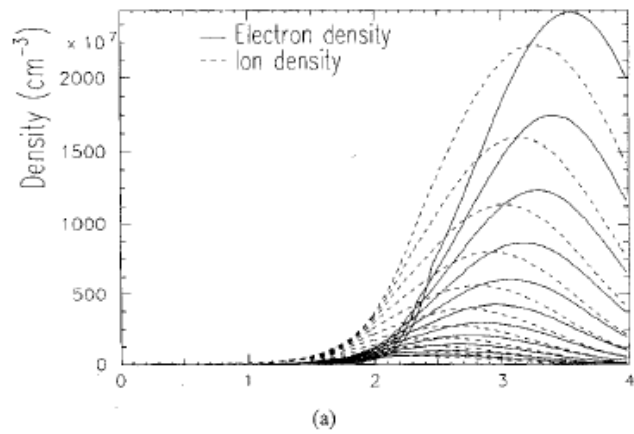
$$\frac{E(x,t)}{2\pi} = - \int_{-(d+x)}^{-x} \rho(x+x')(-1-g(x')) \cdot dx' + \int_{-x}^0 \rho(x+x')(-1-g(x')) dx' + \int_0^{d-x} \rho(x+x')(1-g(x')) dx' - \int_{d-x}^{2d-x} \rho(x+x')(1-g(x')) dx'$$

(5)

Space-Charge cylinder

parameters

'Fast' breakdown ('classical interpretation')

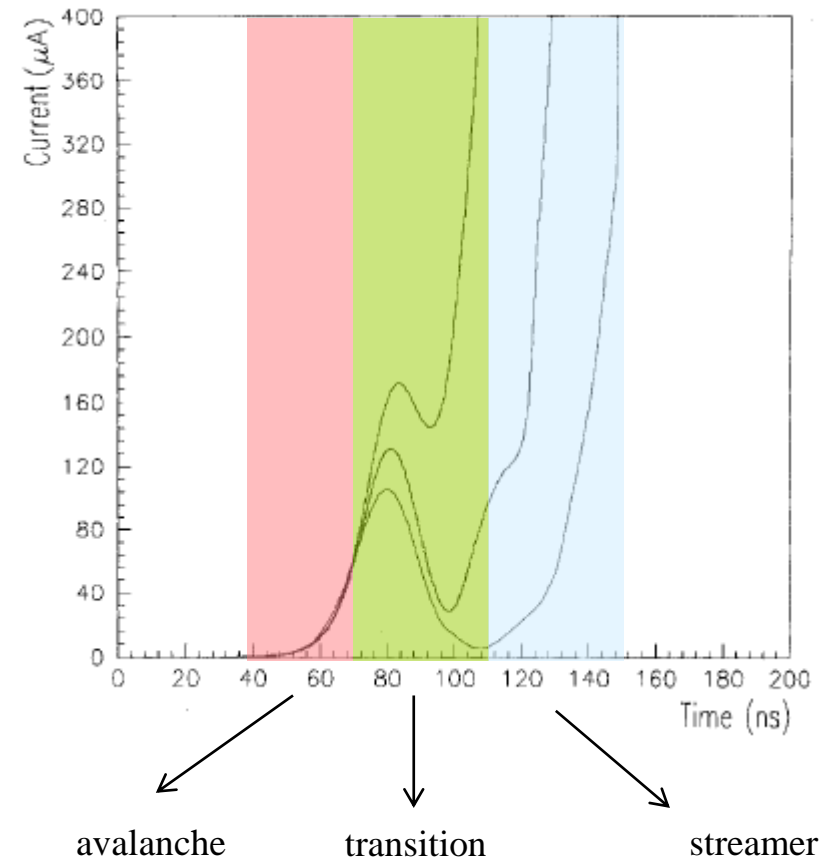


'proportional region' [56ns-76ns]

'transition region' [76ns-106ns]

'propagation region' [106ns-130ns]

'Fast' breakdown ('classical interpretation')



Main (technical) objections:

- Many free parameters, some difficult to experimentally access (spectrum of emission, photo-ionization x-section)
->getting better these days.
- Hydro solutions neglect avalanche or ionization fluctuations, but those will likely trigger breakdown earlier than the average solution to the equations.
- Approximate: transverse dynamics becomes a parameter.
->Probably solvable with present computing power.

Other objections:

- No quantitative comparison with data.
- Does not seem to reconcile well with the fact that improvements on maximum gain with quencher concentrations above some ~1% are very modest.
- Contrary to common wisdom, it does not need of Space-Charge to progress, except if invoking enhanced charge-recombination at low fields, making the process even more difficult to describe.
- It does not seem to be general enough to explain most known systematics on MPGD detectors.

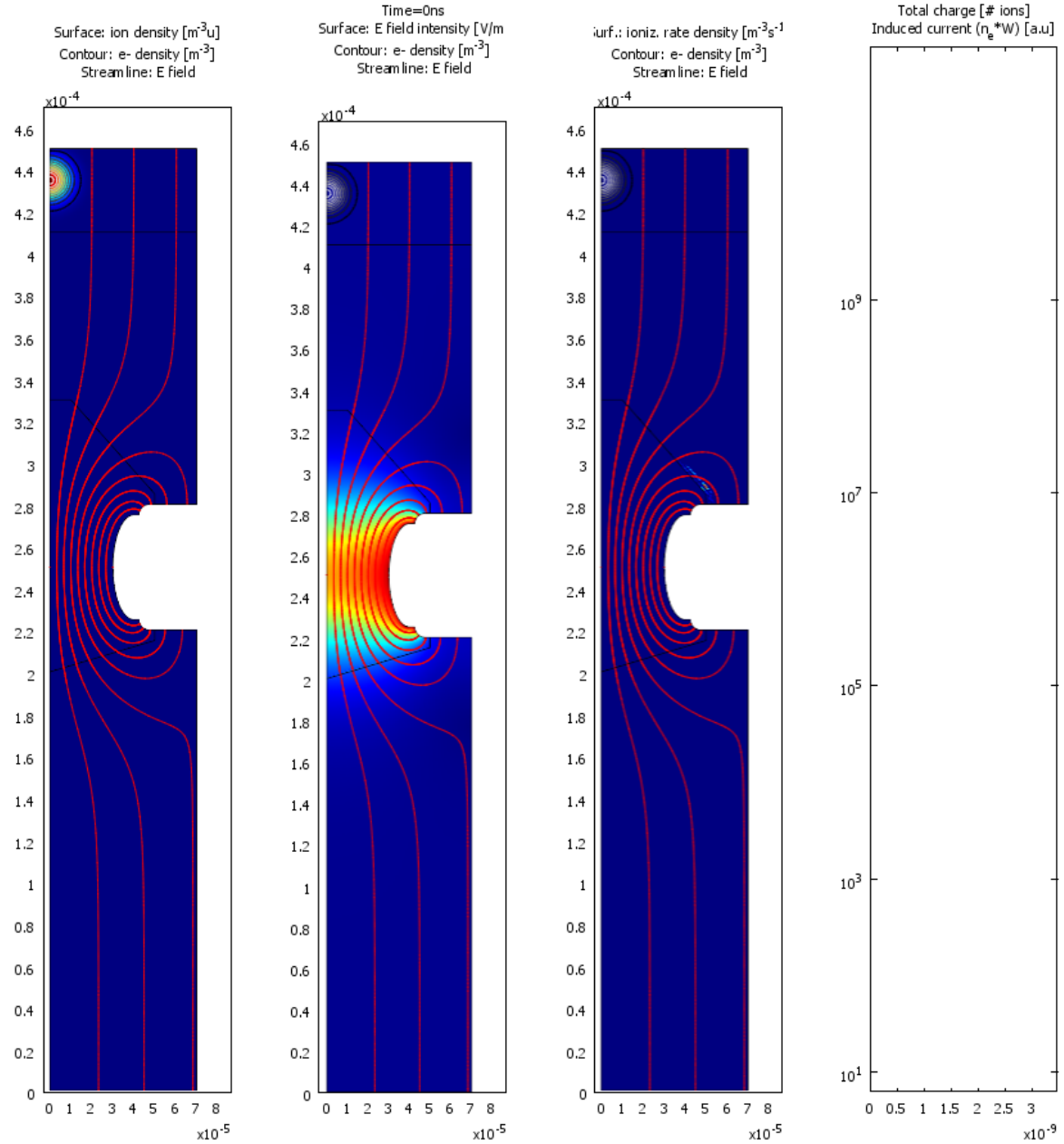


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor is not an unambiguous fingerprint of the underlying physics mechanism!



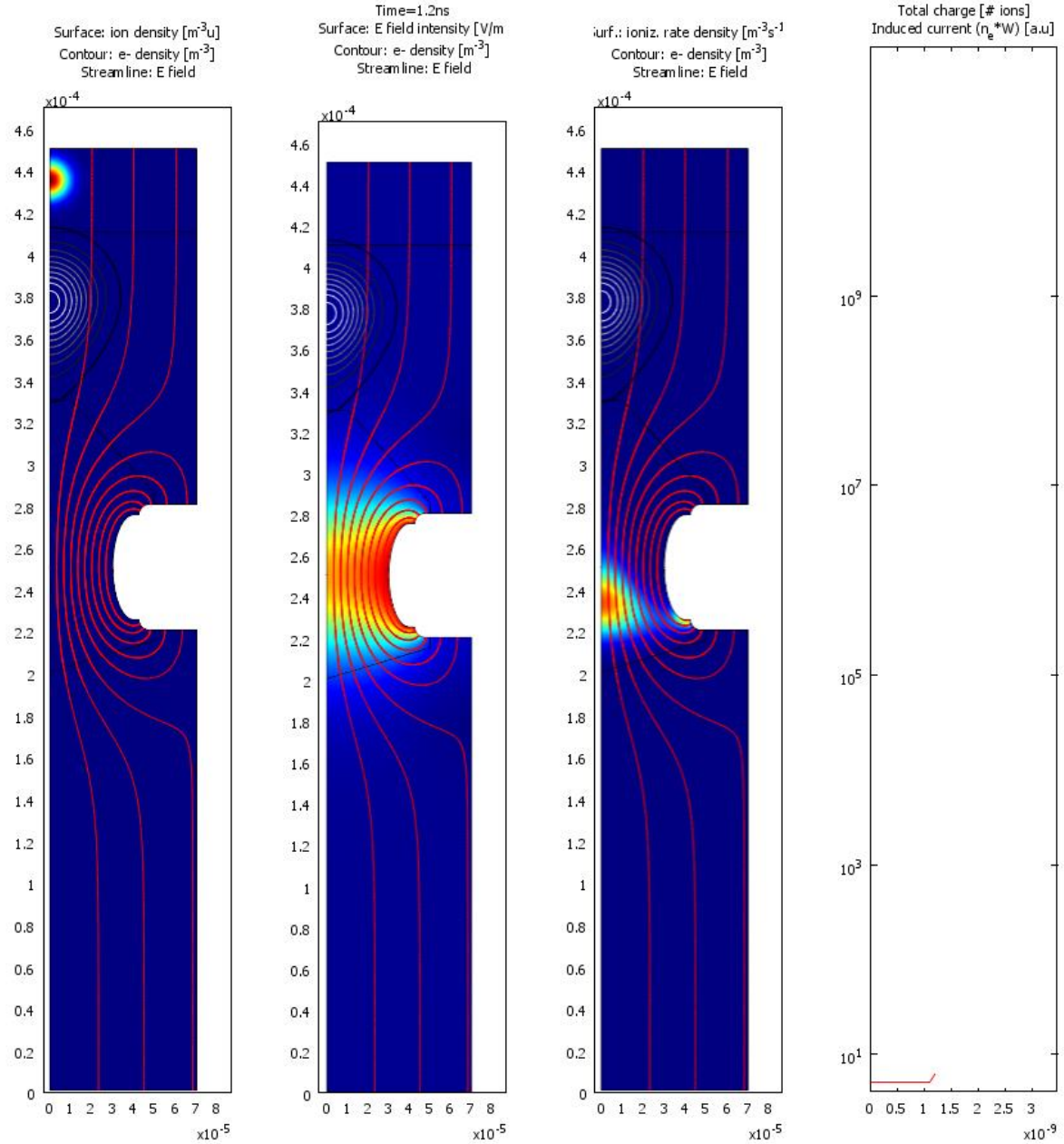


GEM

hole: 60 μm
gap: 100 μm
N₀=100 e⁻
V=1250V

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor is not an unambiguous fingerprint of the underlying physics mechanism!



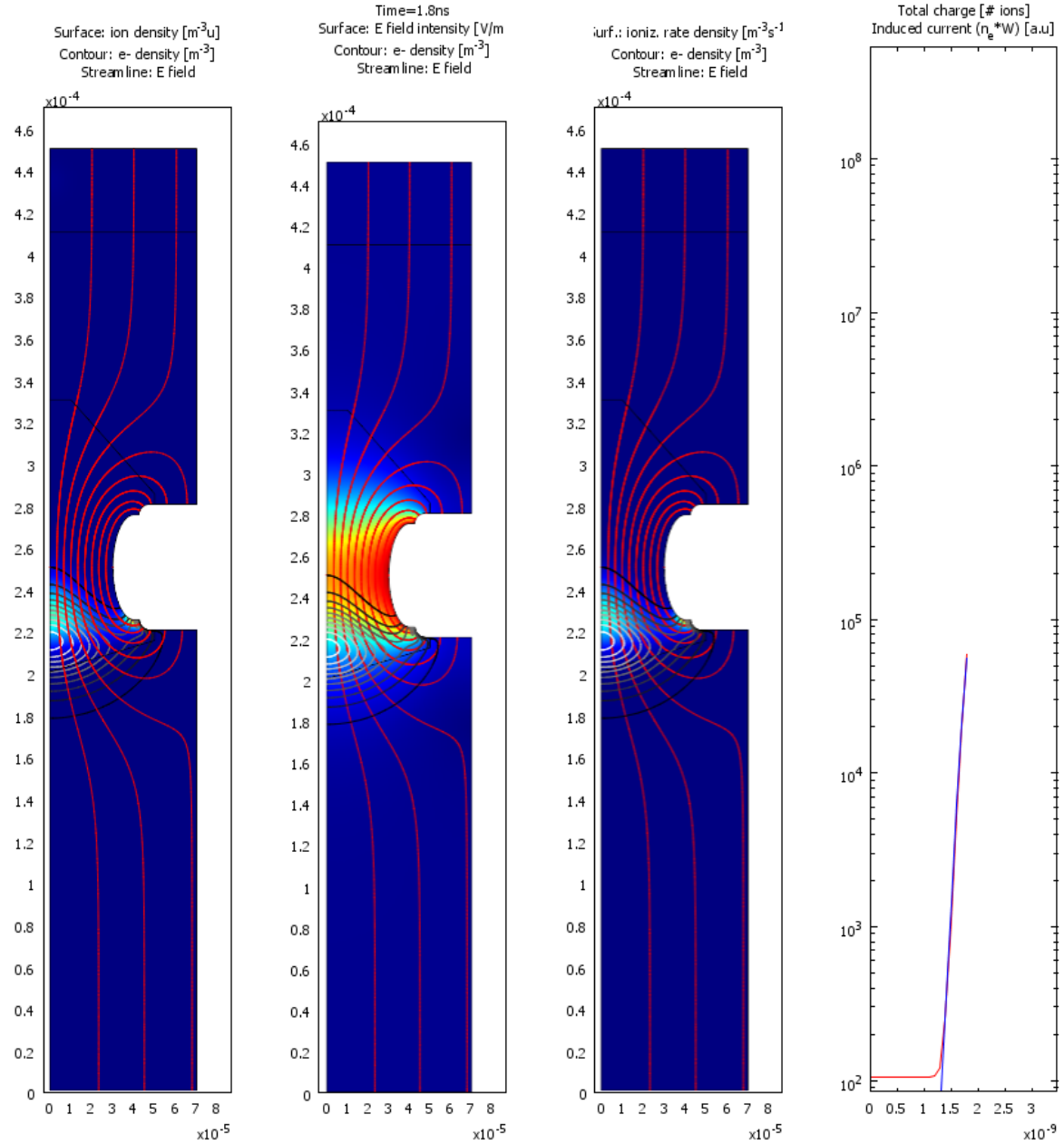


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!



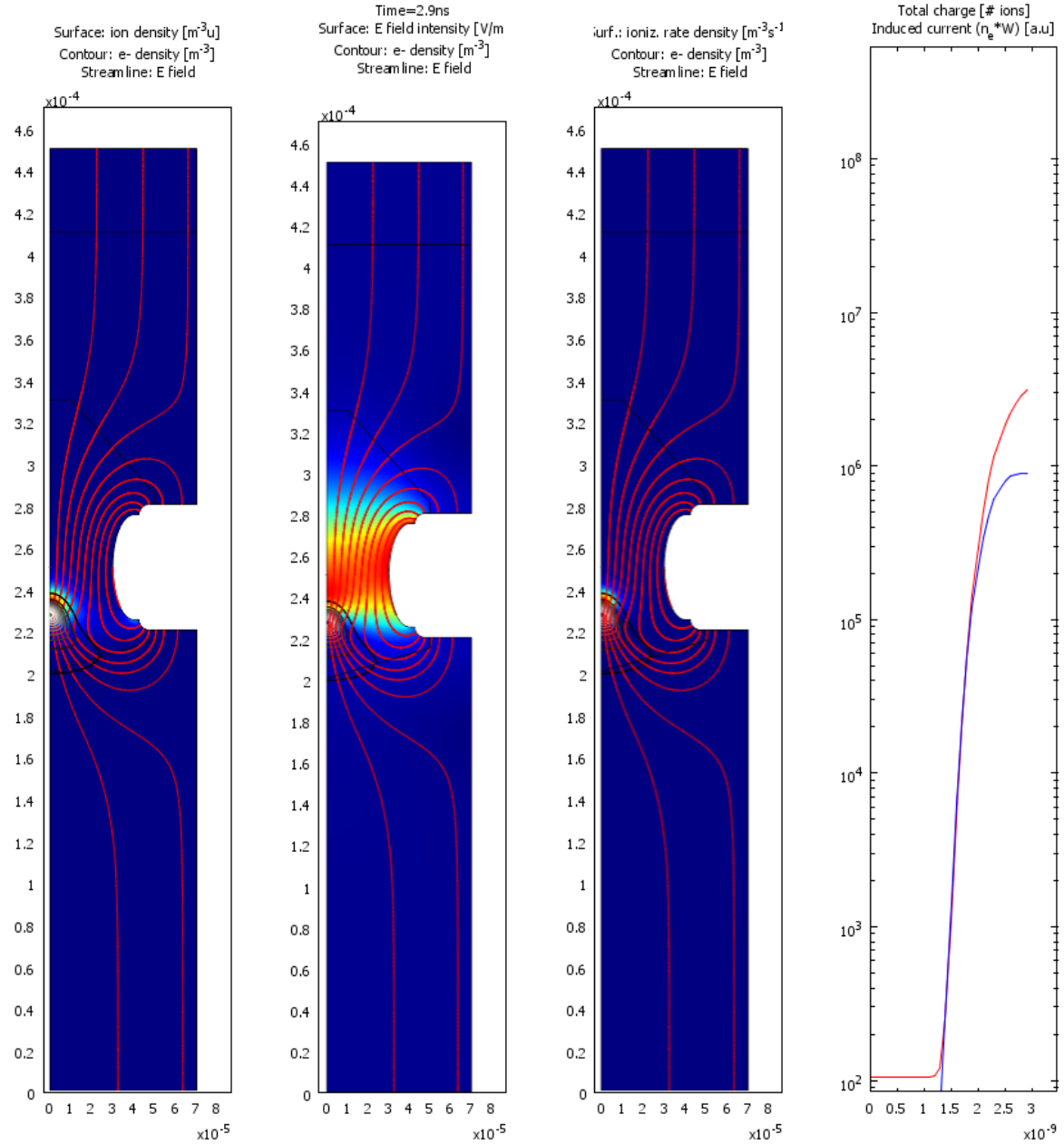


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!



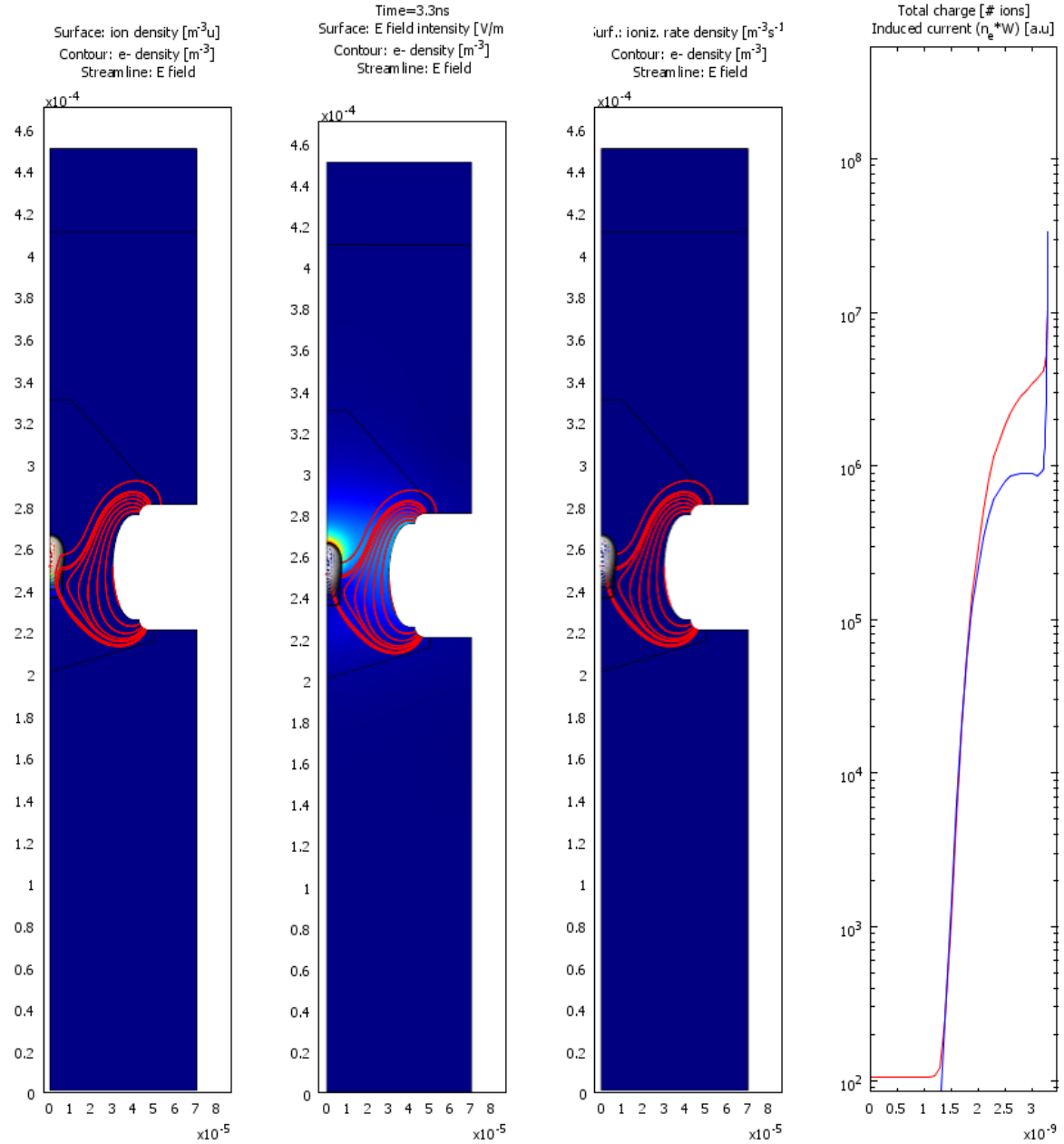


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!



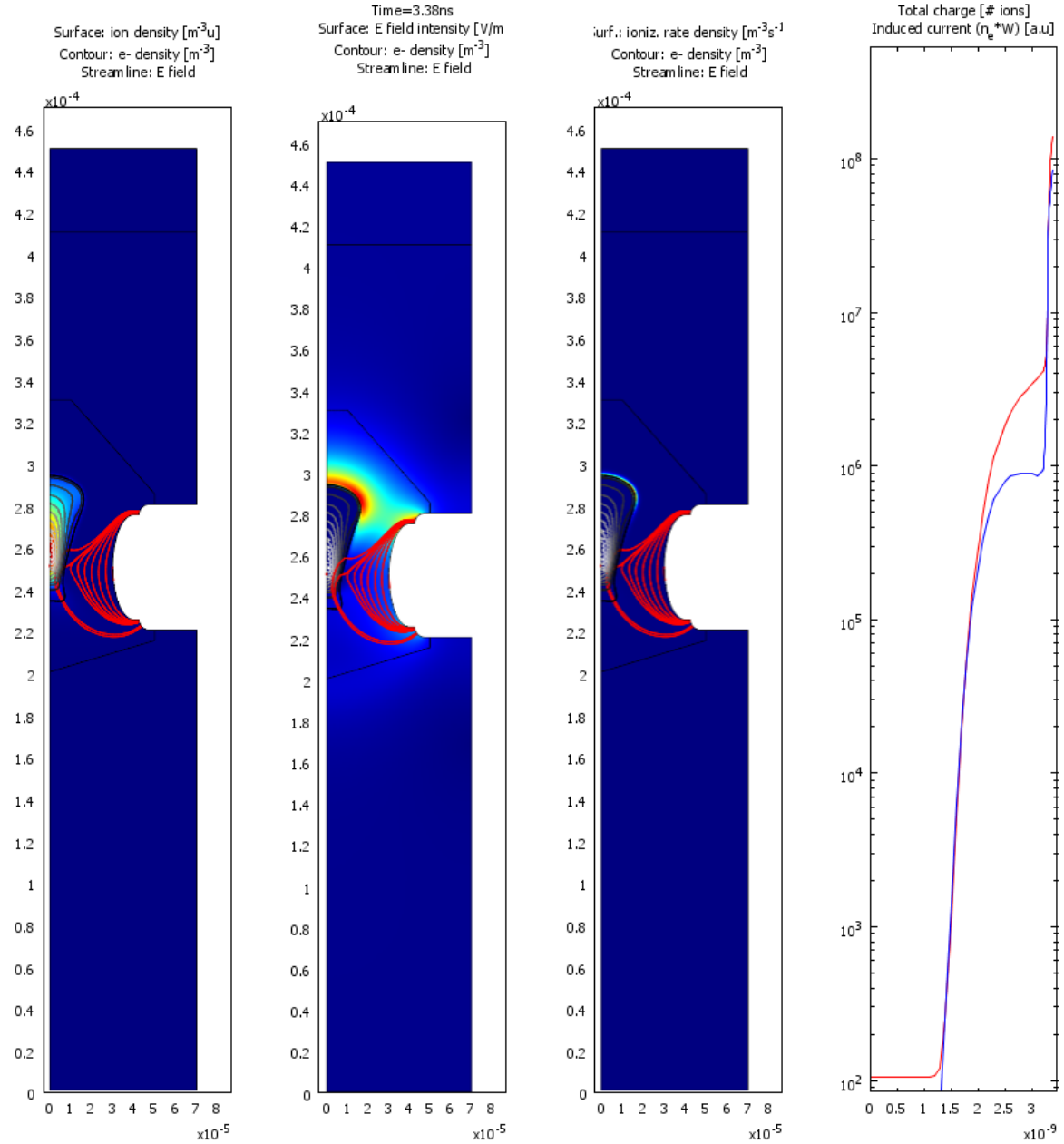


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!



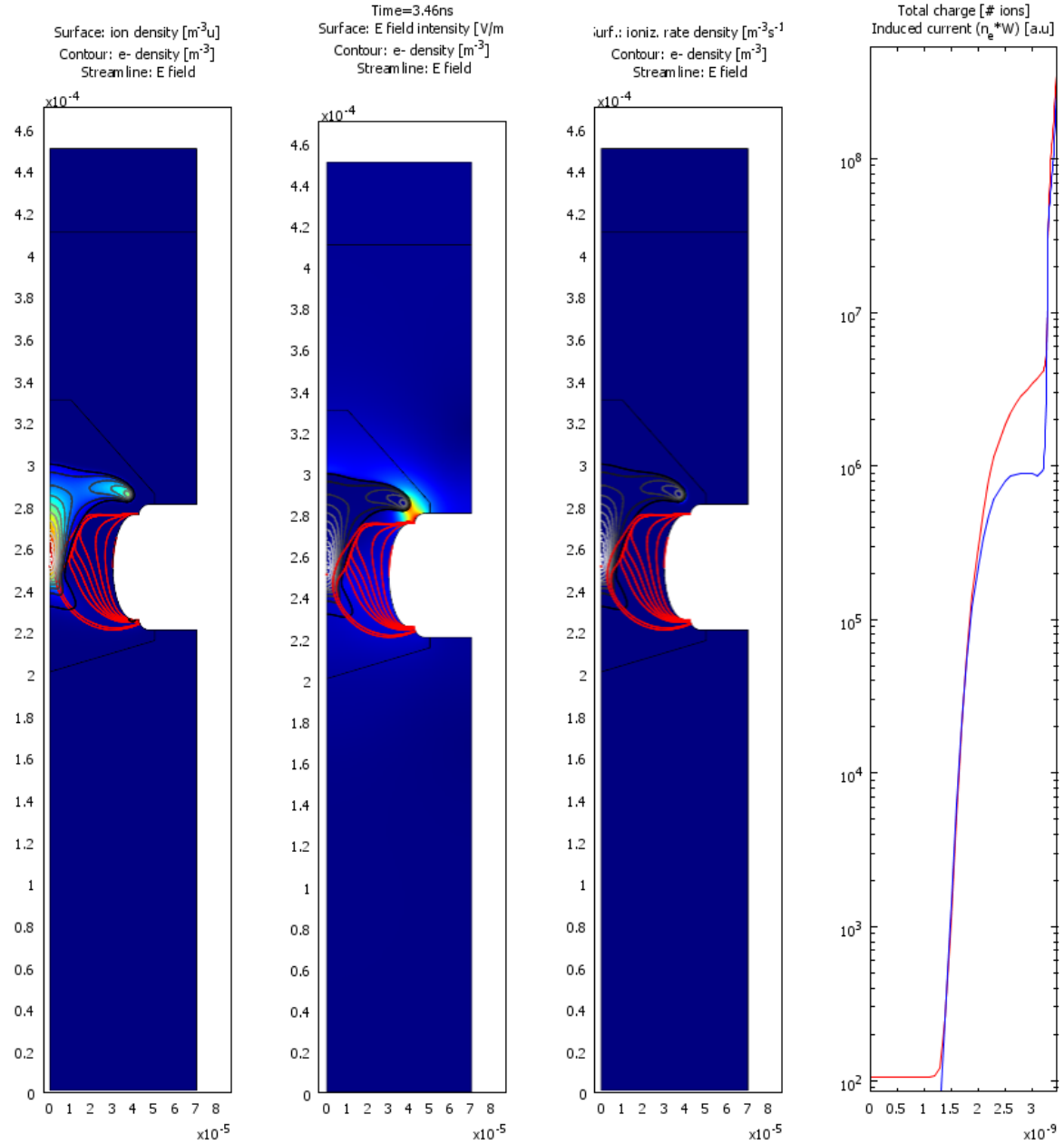


GEM

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!

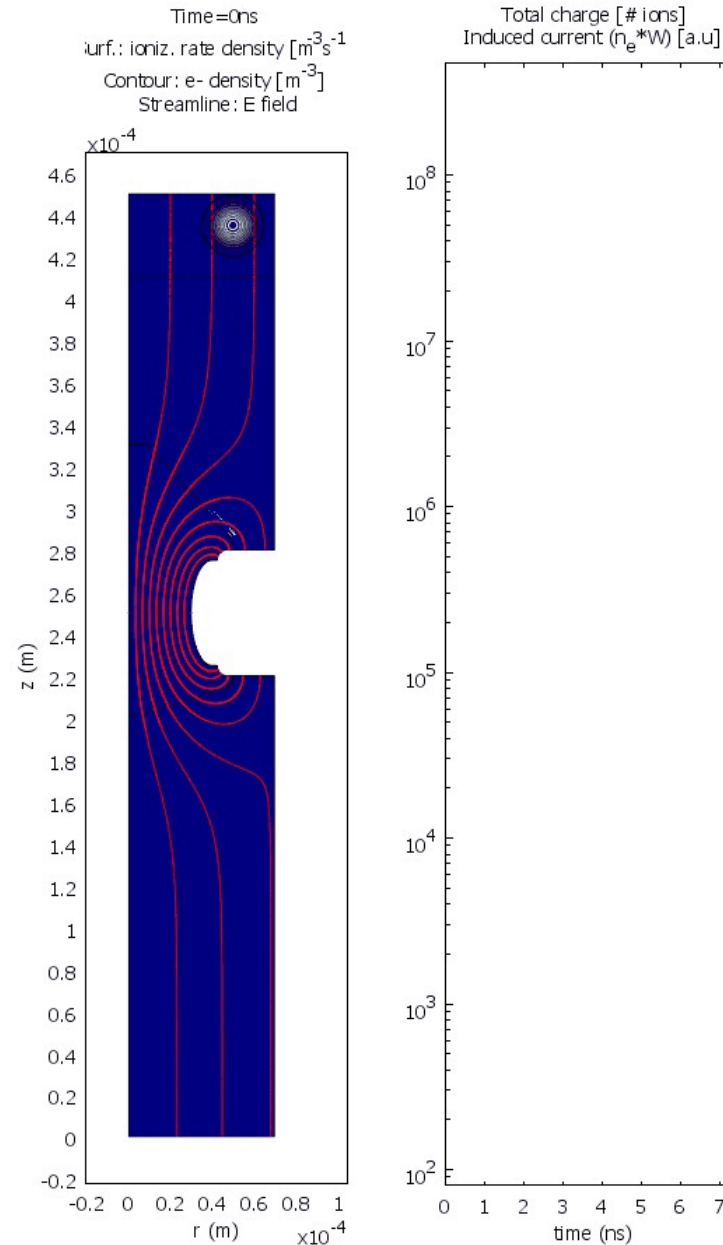


GEM lateral (ring) avalanche

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!

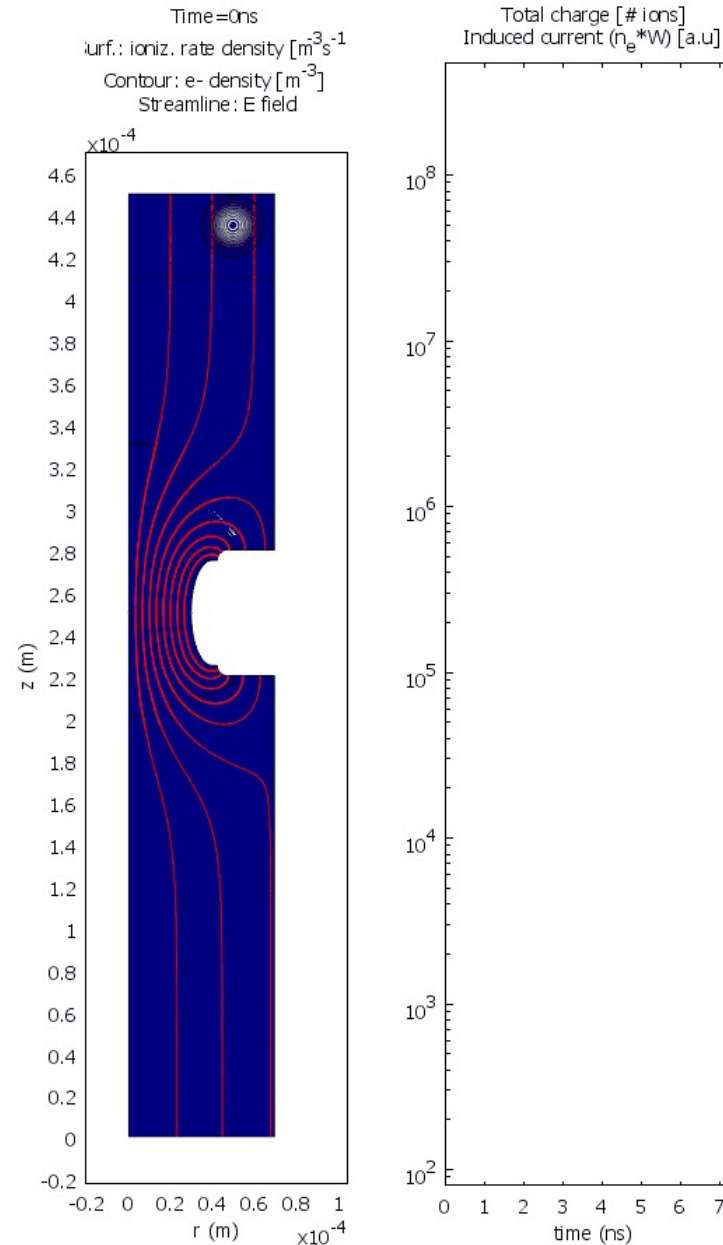


GEM lateral (ring) avalanche

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1250\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor
is not an unambiguous fingerprint of the
underlying physics mechanism!



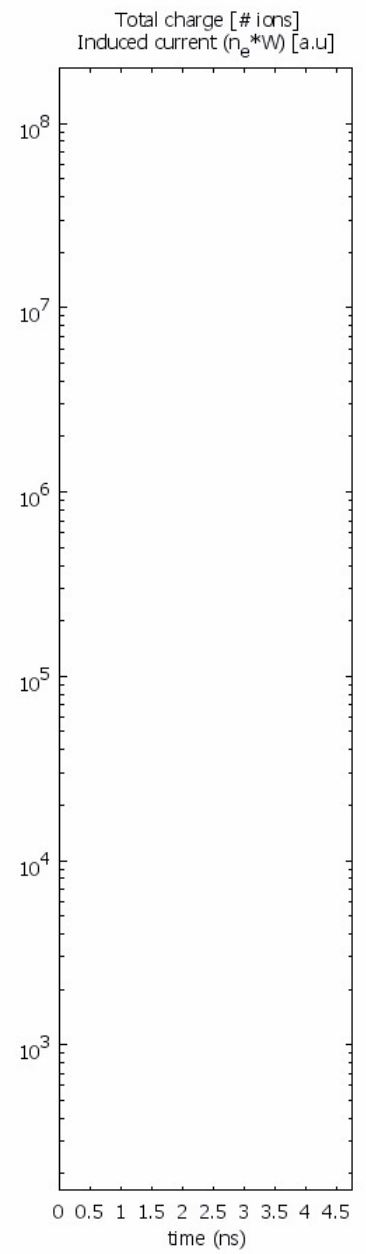
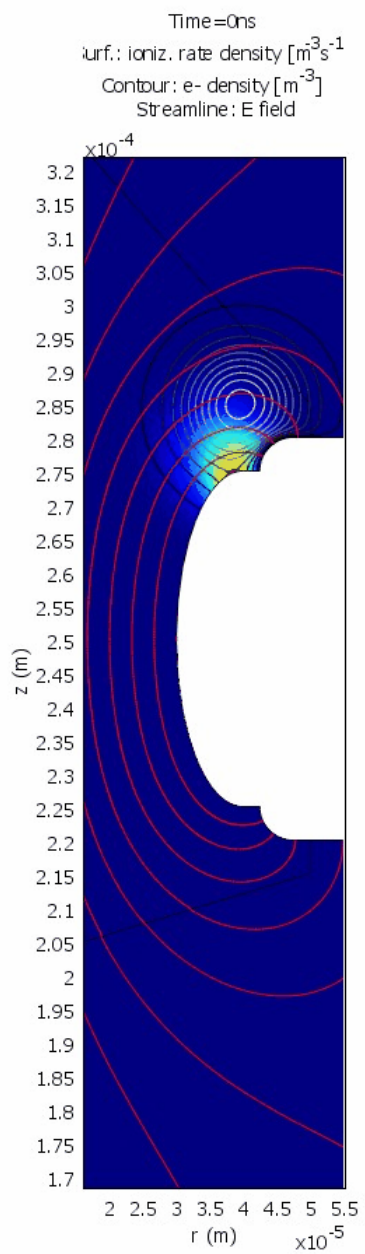


GEM surface avalanche

hole: 60 μm
gap: 100 μm
 $N_0=100 e^-$
 $V=1150\text{V}$

Diffusion-assisted streamer
(no photons)

Unfortunately, the presence of precursor is not an unambiguous fingerprint of the underlying physics mechanism!



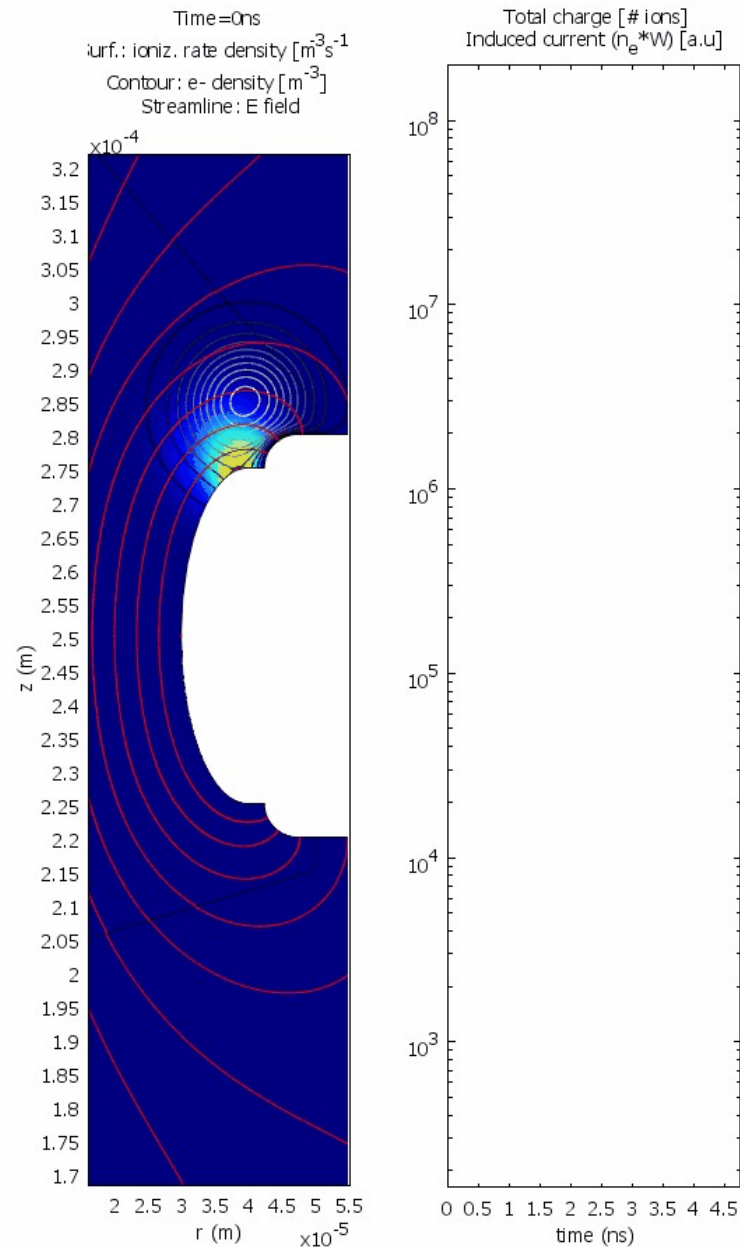
GEM

surface avalanche

hole: 60 μm
 gap: 100 μm
 $N_0=100 e^-$
 $V=1150\text{V}$

This happens 100V below the streamer limit in the space, limiting the practical GEM gain.

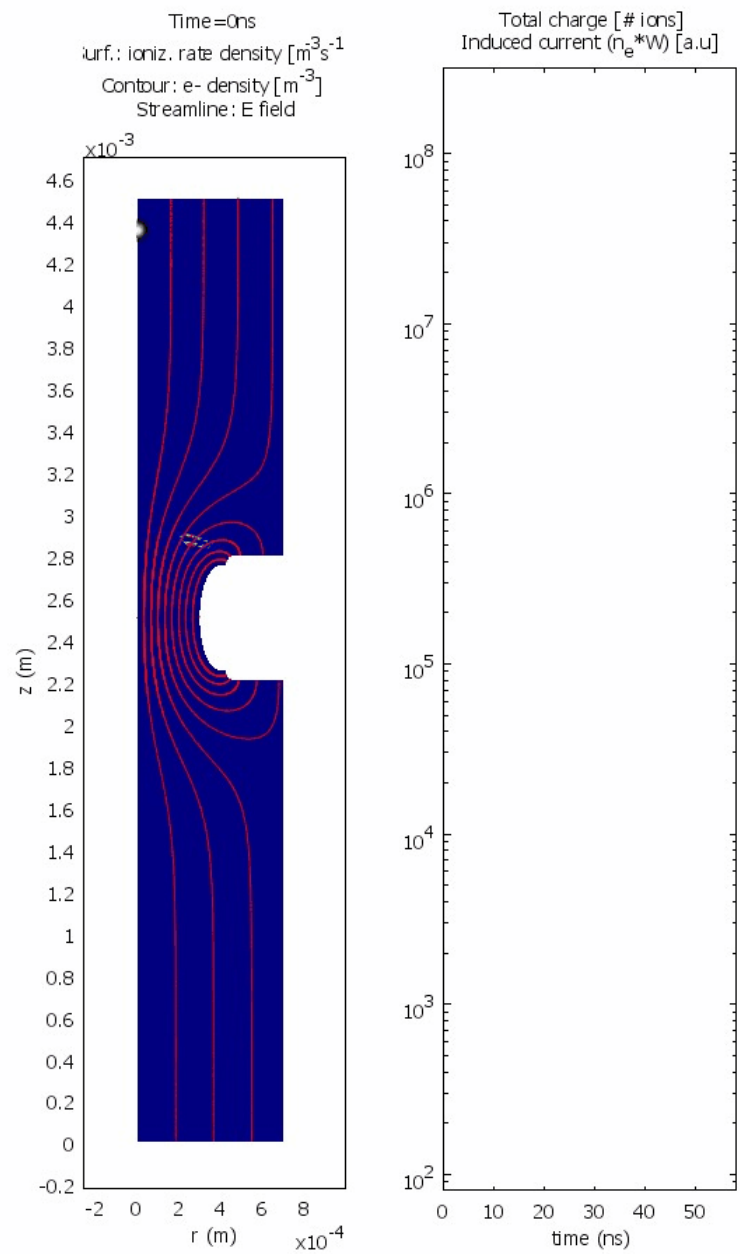
Solved by multistepping.





THGEM (GEMx10)

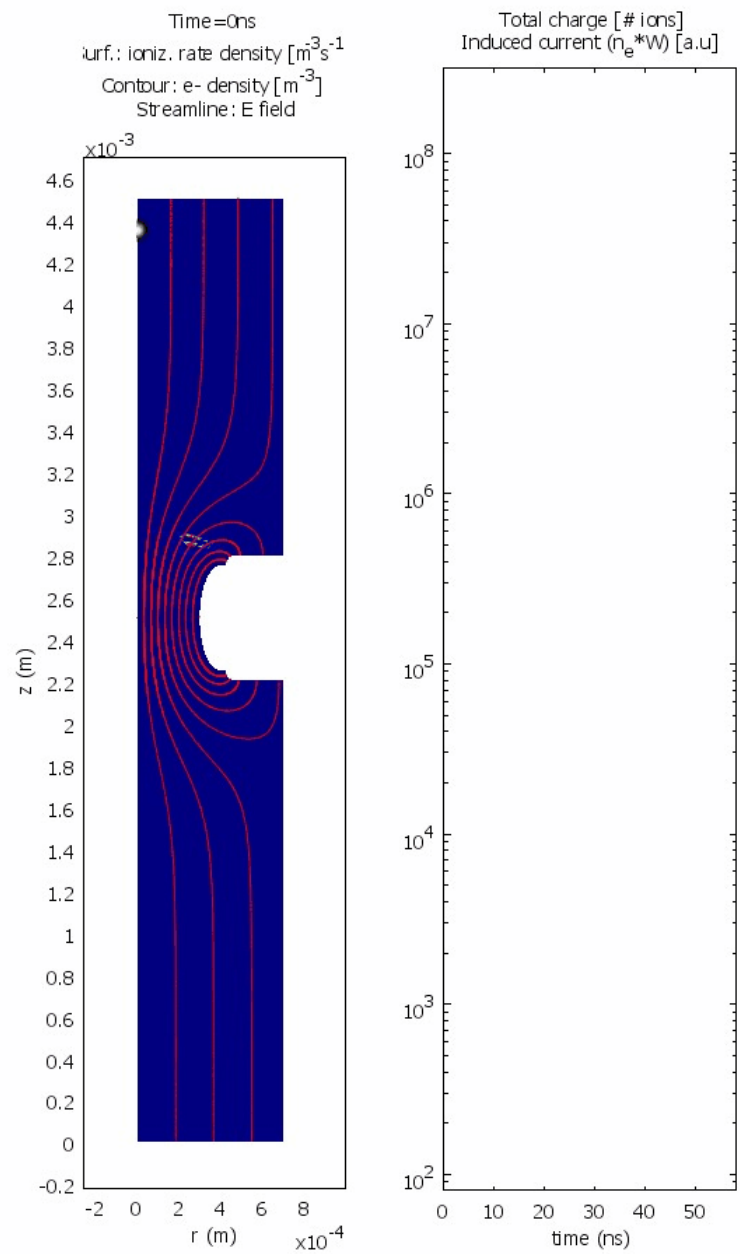
hole: 600 μm
 gap: 1 mm
 $N_0=100 e^-$
 $V=4600\text{V}$





THGEM (GEMx10)

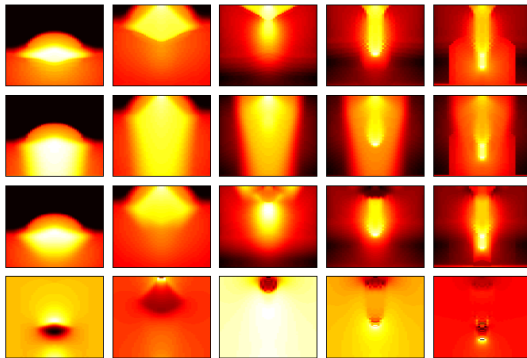
hole: 600 μm
 gap: 1 mm
 $N_0=100 e^-$
 $V=4600\text{V}$



Note: streamer and discharge simulations are ubiquitous these days (e.g.: plasma physics and insulation studies)

Dielectric Breakdown in Insulating Gases

Space Charge Effects and Non-Uniform Fields



Enrique Humberto Radames Gaxiola

Enrique Humberto Radames Gaxiola

Eindhoven, 15 maart 1999

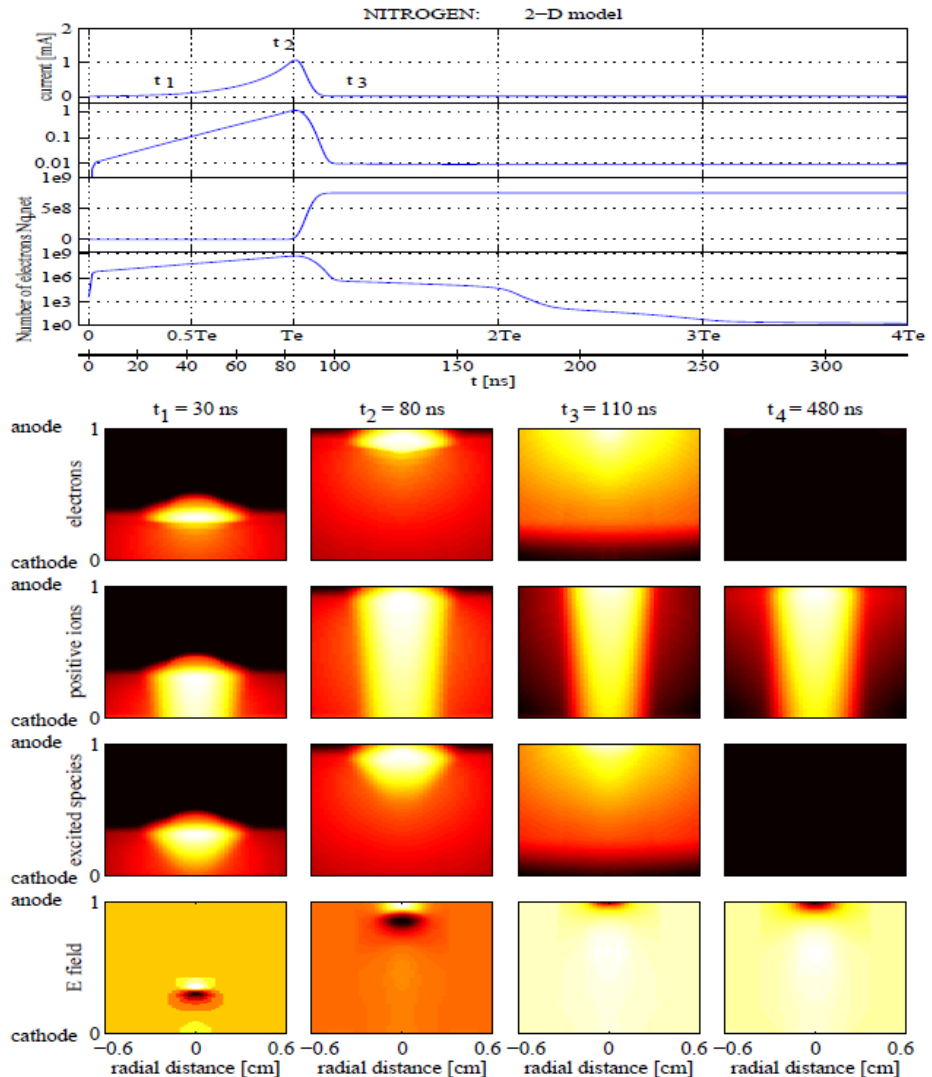


Fig.4.3 Simulated current waveform ($N_0=5.5 \cdot 10^6$, $T_e=83 \text{ ns}$, $A=0.07 \text{ cm}^2$). Electron, positive ion and excited species density and resultant electrical field at various time instances.

Conditions: atmospheric nitrogen $E=30 \text{ kV/cm}$, $gap=1 \text{ cm}$.