Photon and ion induced breakdown

V. Peskov *CERN,* Johann-Wolfgang-Goethe Univ

I'll start from what is known and generally accepted and then will move to new facts and new points of views

Classical point of view

("fast" and "slow" breakdowns)

A few notes about fast breakdown

Typical time ≤ ns $An_0 \sim 10^8$ electron-Raether limit (A -gas gain, n_0 -primary ionization

Animation of spark development when the total charge in the avalanche reaches the Raethet limit:

- a) filed lines close to the avalanche experience a focusing effect and some secondary avalanches start
- b) moving towards the positive ions "body", b) a thin plasm filament-a streamer- is formed,
- c) c) when the streamer reaches the electrodes, a spark happens

In case of metallic electrodes breakdown disable the detector for some time

Current oscillograms of static breakdown in methylal at various overvoltages increased from the lowest to upper curves (from Raether book) 6 and 6 a Many researches belief in a strong role of avalanche and streamer emission in the detector operation,.. although with some reservations

Linear absorption coefficient for the components of gas mixture used in high-pressure Pestov counters (corresponding to their quantity). The corves 1 and 2 correspond to propylene (C3H6) and neon respectively 7

In case of dielectric electrodes (in contrast to metallic electrodes) a breakdown disable only some part of the detector and one can observe "the consequence of photon emission on the discharge development"(Pestov)

Details of discharge development in RPC

Y. Inoue et al., NIM, A394, 1997,65

Visual observations

10

Large distance (2-3 mm) between consequent streamers is difficult to explain without photon mechanism

A puzzle…

From the "theoretical" point of view, not in all gases the UV emission of avalanches are capable to ionize its own gas

Some related studies

P. Fonte et al., NUM A310, 1991, 140

Our results allow us to put forward a hypothesis about the origin of the emission. For example, in the case of an Ar + TEA mixture in the PPAC we can assume that the molecular emission of TEA cannot ionize TEA molecules in the MWPC. On the other hand, the measurements with the $CH₄$ gas filter show that the first and second continua of Ar^{*} emission (which exist mainly below 140 nm) do not contribute noticeably, so the observed emission can be either third Ar^{*} emission at $\lambda > 140$ nm [10] or fragment emission of TEA, which leads to the emission of atomic lines: C^* , H^{*}, N^{*}, etc. The third continuum of argon corre-

By using $CH₄$ as a gas filter we could separate the sensitive band into two subbands: 120-140 and 140-170 nm. For mixtures containing hydrocarbons the emission is concentrated in the band 140–170 nm, whilst for pure Ar the emission is concentrated in the 120-140 nm band as should be expected from argon excimer emission. These results suggest that for the mixtures containing hydrocarbons the emission is mainly due to fragment emission and not to noble gas excimer or molecular emission.

A few notes about slow breakdown

Typical time…up to ms or even more

Conditions:

An example of a slow breakdown studies

Fig. 2. Typical chamber signals, as detected with a fast current preamplifier: (a) chamber signal delayed after the PM signal (top) by the avalanche time, (b) \sim 25% feedback – our standard level for feedback measurements, (c) many feedback successors, (d) slow breakdown.

In both cases (fast and slow breakdowns) the main contribution comes from the VUV and UV avalanche photons

IONIZATION POTENTIALS FOR COMMON INDUSTRIAL GASES

Many compounds not appearing in this list with an ionization potential of 11.7 eV or less may also be detectable

Other sources of ionization potential data: CRC Handbook of Chemistry and Physics NIOSH Pocket Guide to Chemical Hazards HYPERLINK http://webbook.nist.gov/

Ion induced breakdowns

 $y_{+} = k_{\text{gas}}$ (E/p) (E_i-2φ)

V. Peskov, Zh.Tekh Fiz, 45, 1975,1268

Beside these "classical" ion feedback mechanism there are also observed "exotic" mechanisms called : breakdown preparation/cathode excitation/memory effect and electron jets

Preparation mechanism

Similar effect is often describes in aging papers

(see for example *Aging Workshop*, *NIM A515, 2003*)

Breakdown *I. Ivanchenkov et al, IEEE, 45,1998, 258*

Usual explanation is via Malter effect..

Classical Malter effect predicts single electron emission *(see L. Malter , Phys. Rev, 49,1936, 478*)

However, in most cases a slow current increase is just an integral of high amplitude pulses *I. Ivanchenkov et al, IEEE, 45,1998, 258*

This strongly contradict to classical Ma;ter effect

More detailed studies reveal that the preparation mechanism may exhibit not only as current pulses but also as a short-term current growth

This was called a cathode excitation effect

Figure 2a), b): Two typical oscillograms showing a preparation

P. Fonte etal., IEEE Nucl Sci,46,1999,321 ²⁷

Measurements with single wire counters:

It is well known that In single wire counters: $Ay=1$

V max

This curve is typical for many gaseous detectors, including MPGDs (check with your experience!)

Early studies of the cathode excitation effect (Karabadjak Thesis)

In self-quenchd mode ion emission does dot contribute to the γ_{+} studies

Fig. 8. Sketch of a multisection gas counter for γ_+ measurements, 1 - Cylindrical anode; 2 - multisectional cylindrical cathode; $3 -$ glow-discharge; $4 -$ cathode spot of the glow-discharge; 5 - external VUV radiation source; 6 - collimating system for the VUV radiation.

More detailed studies

Glow discharge at elevated pressures

G. F. Karabadzhak and V. D. Peskov

S. I. Vavilov Institute of Physics Problems, Academy of Sciences of the USSR, Moscow (Submitted July 22, 1985; resubmitted January 15, 1986) Zh. Tekh. Fiz. 57, 891-895 (May 1987)

Fig. 9. Values of γ , as a function of the angle Θ between the upper section (see fig. 8) and investigated one and for different time delays τ after turning off the glow-discharge, $1 - \tau = 9 \mu s$; $2 - \tau = 10$ ms.

Changes in QE after intense ion bombardment

P. Fonte , V. Peskov, ArXive

So it was clearly observed that after intense ion bombardment the QE temporally increase as well as γph and γ+

Therefore , the feedback loop Aγ=1 will appear at lower A

Further studies

PUBLISHED BY IOP PUBLISHING FOR SISSA

RECEIVED: August 1, 2010 ACCEPTED: October 4, 2010 PUBLISHED: November 22, 2010

Further evaluation of a THGEM UV-photon detector for RICH - comparison with MWPC

V. Peskov, $a,b,1$ M. Cortesi, a R. Chechik a and A. Breskin a

['] inst

Figure 10. Counting rate of spurious pulses and visible-light induced pulses vs. time after induction of the cathode excitation effect; CsI-coated triple THGEM; gas mixture: Ne+10%CH₄; gas gain \sim 10⁴.

Figure 11. Counting rate of spurious pulses and visible-light pulses vs. time after cathode excitation induction at gain 10⁴, followed by a 10-fold gain increase. CsI-coated triple THGEM, Ne+10%CH₄ mixture, gas

35

We also succeed to reproduced COMPASS problems

Jets

What is the origin of these gigantic pulses?

38

Explosive field emission

Besides classical field emission calculated by Zommerfeld and others there is another phenomena -explosive field emission

(from [17]). Enlargement shows pulses due to the explosive field emission.

R. Latham, "High voltage vacuum insulation", new Yoork,1995

See: G.A. Lubimov, V.I. Rahovski, Uspekh. Phys. Nauk, 125, 1978, 665, V. Peskov Journ, de Physique Coll. C7, suppl#7, 1979,C7-333

A proposed mode of electrons jets in gaseous detectors:

Figure 3: Schematic illustration of a two-step process which leads to emission of jets and bursts from thin dielectric films.

rate-induced spark at low gas gain. A linear increase in current is visible for about one second before breakdown, followed by several large pulses that coincide with a spark.

Available online at www.sciencedirect.com

SCIENCE ω Direct*

Nuclear Instruments and Methods in Physics Research A 535 (2004) 632-643

www.elsevier.com.flocate/nima

Model of high-current breakdown from cathode field emission in aged wire chambers \hat{r}

Adam M. Boyarski*

Stanford Linear Accelerator Center, M.S. 95, 2575 Sand Hill Road, Menlo Park, Stanford, CA 94025, USA

Received 18 March 2004; accepted 25 June 2004 Available online 13 August 2004

Some of the results presented in this paper were interpreted via the jets mechanism

Role of adsorbed layers?

If Ei<2φ –no ion recombination

Figure 4: Improvement in the rate characteristics of a PPAC by optimizing the gas mixture.

Observations of jets and cathode excitation effect in operation of some gaseous detectors

Examples:

Glass RPC

Fig. 5. The efficiency and the rate of noise pulses versus the voltage applied on the RPC [4]. The curve with the triangle symbols correspond to measurements done in anti-coincidence with the signals from the scintillators, respectively. Gas mixture Ar/Isobutane/Freon (R134) in the ration 48/4/48.

C. Iacobaeus et al., IEEE Trans Nucl. Sci, 49, 2002, 1622 ⁴⁶

High rate Si and GaAs RPCs

Fig. 11. The rate of the afterpulses for the PPAC (Cu-electrodes) and the RPC (Si). Gas mixture Xe (20%)+Kr (40%)+CO₂ (20%) at 1 atm.

Fig. 6. Pulse-height spectra of signals from RPCs measured in the case of single primary electrons produced from the cathode by (a) UV emission and (b) in the case of noise pulses. The gas mixture Xe (40%)+Kr (40%)+CO₂ (20%) was used (1 atm).

C. Iacobaeus et al., IEEE Trans Nucl. Sci, 49, 2002, 1622 ⁴⁷

GEM at extreme counting rates

Fig. 12. The current from the GEM (at 350 V) recorded directly on a 50- Ω input of the oscilloscope when the GEM was exposed to a pulsed gamma radiation, producing $\sim 10^7$ counts/mm² on the 2.5 cm \times 2.5 cm GEM area. No other resistors (except the 50- Ω input of the scope) were connected. The upper figure shows the current pulse from a racetrack current monitor. The lower figure shows the corresponding current pulse from the GEM readout. The gas mixture $Ar + 20\% CO_2$ was used for the measurement (1 atm).

Fig. 13. The same setup as in Fig. 12, but 420 V applied over the GEM electrodes. The upper oscillogram shows the current pulse from the racetrack current monitor, the lower shows the current from the GEM readout. One can clearly see current pulses of large amplitudes, corresponding to a large number of primary electrons $>10^5$.

C. Iacobaeus et al., IEEE Trans Nucl. Sci, 49, 2002, ⁸1622

Delayed discharge propagation between GEMs

Fig. 41. Two breakdowns following each other: the breakdown in the GEM was followed with some delay by discharge propagation to the collector

See: V. Peskov. "Discharge propagation between GEMs," WG-2 meeting 49

Fig. 42. In a more sensitive scale one can see a steady current increase before the second breakdown happens

This phenomena resembling the current growth before breakdown mentioned earlier

Conclusions

- •Besides well established breakdown mechanisms- streamer and feedback related -it was also observed another one: "memory/cathode excitation-jets"
- This mechanism mainly show up at high counting rates. For example COMPASS RICH already experience in the memory effect and some tests of RPCs at GIF also reveal this effect too
- Very often cathode excitation and jets mechanism are mixed
- It will be important to further study these effects , because they can cause problems at future MPGDs applications in LHC experiments.

Backups

