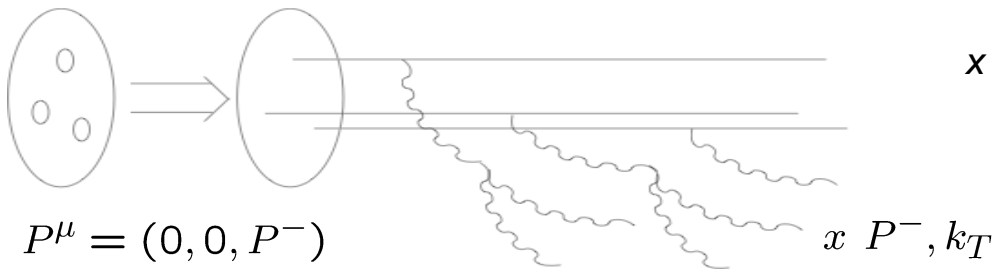


Confirming RHIC saturation signals at the LHC

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Map of parton evolution in QCD



x : parton longitudinal momentum fraction

k_T : parton transverse momentum

the distribution of partons as a function of x and k_T :

QCD linear evolutions: $k_T \gg Q_s$

DGLAP evolution to larger k_T (and a more dilute hadron)

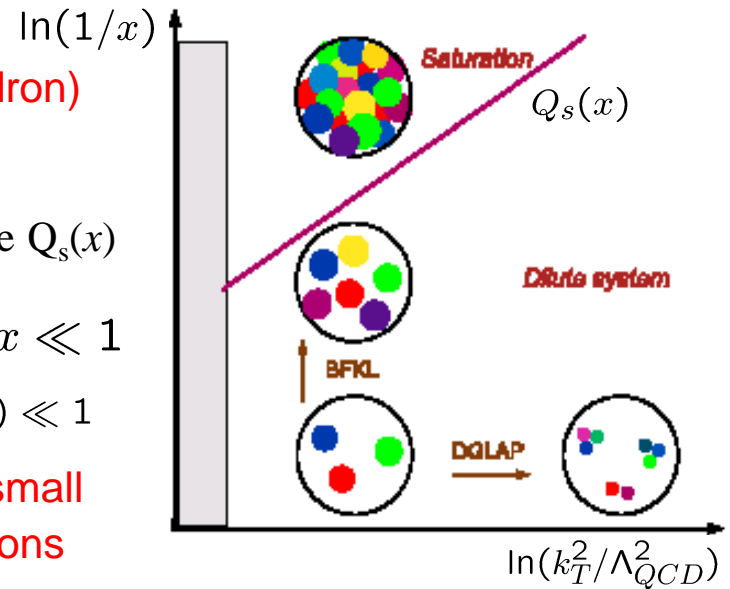
BFKL evolution to smaller x (and denser hadron)

dilute/dense separation characterized by the saturation scale $Q_s(x)$

QCD non-linear evolution: $k_T \sim Q_s$ meaning $x \ll 1$

this regime is non-linear yet weakly coupled: $\alpha_s(Q_s^2) \ll 1$

collinear factorization does not apply when x is too small and the hadron has become a dense system of partons



$$S_{DIS}(x_{Bj}, Q^2) = \int_0^1 dx \int_{partons\ a\ x_{Bj}} j_{a/p}(x, Q^2) \hat{S}_a(x_{Bj}/x, Q^2) + O(Q_0^2/Q^2)$$

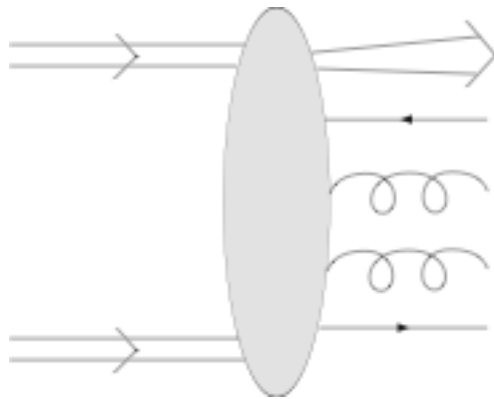
parton density
partonic cross-section
higher twist
 $\frac{(A/x)^{1/3}}{Q^2}$

Saturation signal #1:

forward rapidity suppression
of the nuclear modification
factor in $p+A$ vs $p+p$

Single inclusive hadron production

forward rapidities probe small values of x

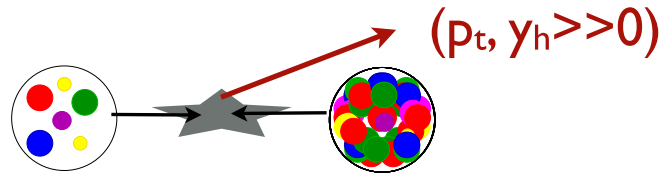


k_T, y transverse momentum k_T , rapidity $y > 0$

values of x probed in the process:

$$x_1 = M_T e^y / \sqrt{s} \quad x_2 = M_T e^{-y} / \sqrt{s}$$

$$M_T^2 = (k_T/z)^2 + m_h^2$$



large- x parton from proj. (pdf)

small- x glue from target (CGC)

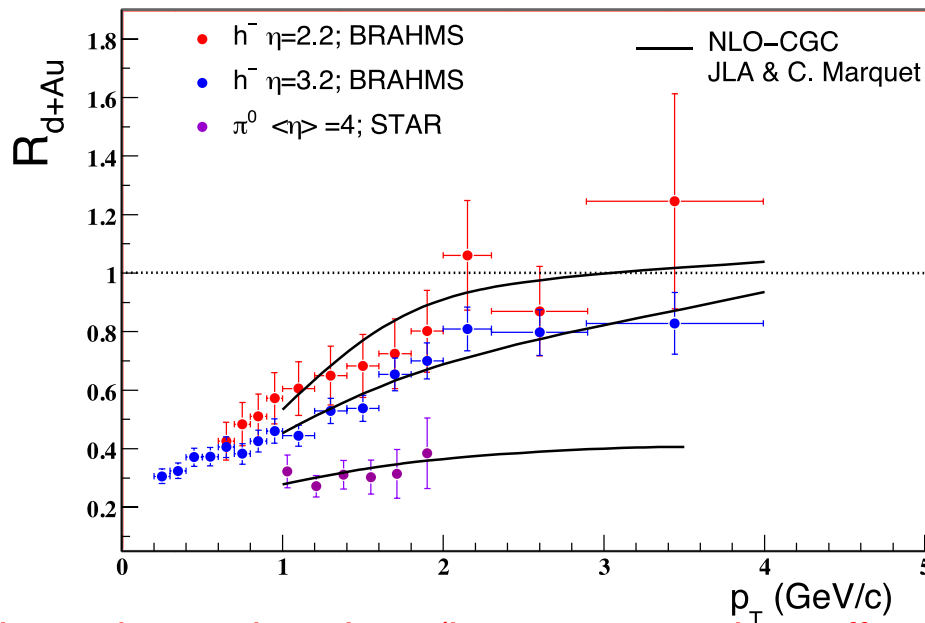
$$\frac{dN_h}{dy_h d^2p_t} = \frac{K}{(2\pi)^2} \sum_q \int_{x_F}^1 \frac{dz}{z^2} \left[x_1 f_{q/p}(x_1, p_t^2) \tilde{N}_F \left(x_2, \frac{p_t}{z} \right) D_{h/q}(z, p_t^2) \right. \\ \left. + x_1 f_{g/p}(x_1, p_t^2) \tilde{N}_A \left(x_2, \frac{p_t}{z} \right) D_{h/g}(z, p_t^2) \right] \xrightarrow{\text{fragmentation}}$$

Nuclear modification factor

$R_{dA} = 1$ in the absence of nuclear effects, i.e. if the gluons in the nucleus interact incoherently as in A protons

$$R_{dA} = \frac{1}{N_{coll}} \frac{\frac{dN^{dA \rightarrow hX}}{d^2kdy}}{\frac{dN^{pp \rightarrow hX}}{d^2kdy}}$$

the suppressed production ($R_{dA} < 1$) was predicted in the Color Glass Condensate picture, along with the rapidity dependence



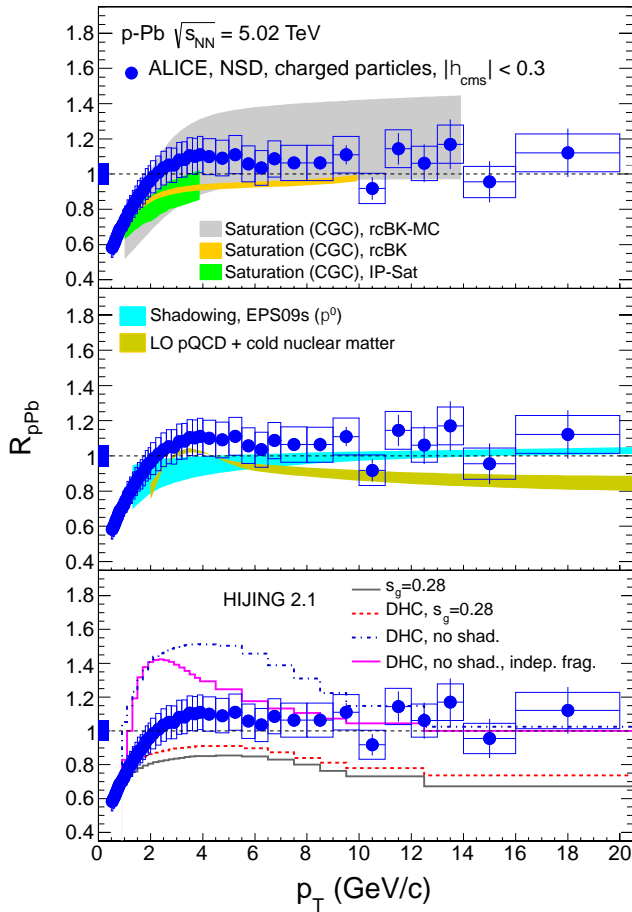
Albacete and CM (2010)

note: alternative explanations (large- x energy loss effects) have been proposed

Kopeliovich et al (2005), Frankfurt et al (2007)

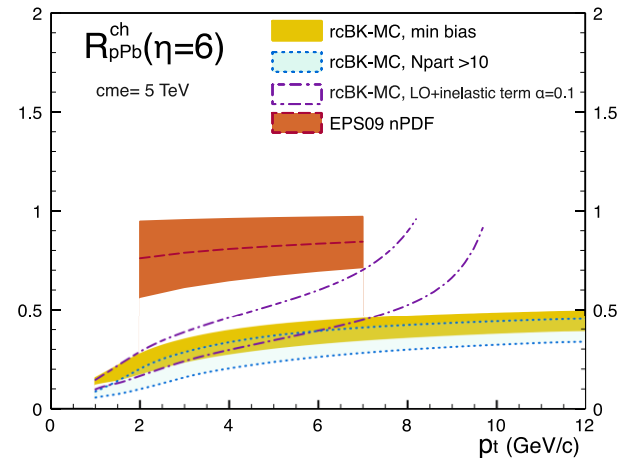
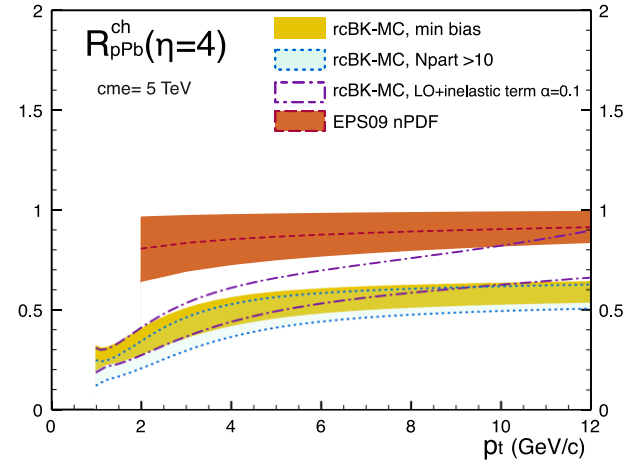
p+Pb @ the LHC

- mid-rapidity data



good description but not much non-linear effects

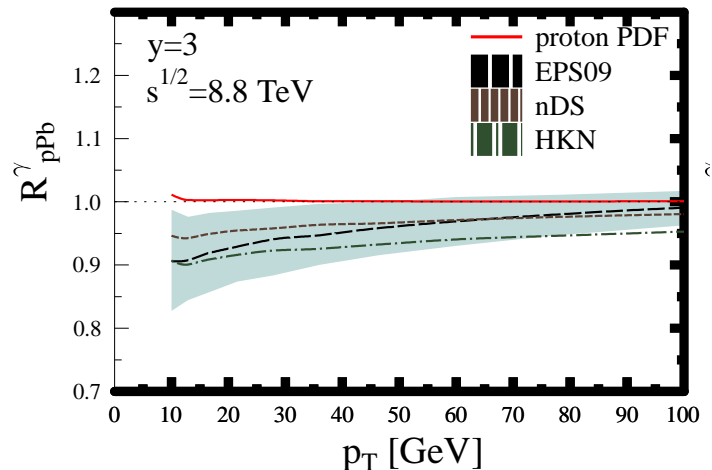
- predictions for forward rapidities



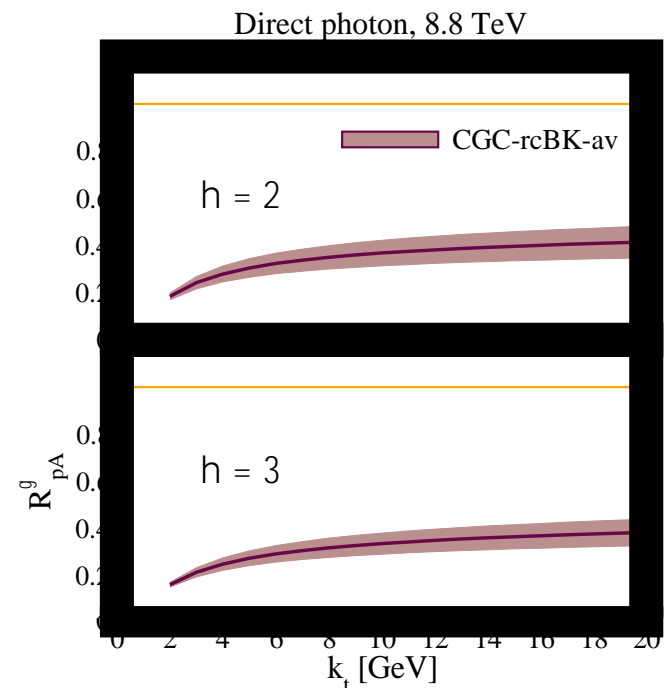
strong non-linear effects

Best way to confirm R_{pA} suppression at the LHC

- isolated photons at forward rapidities
 - no isospin effects in p+Pb vs p+p (contrary to d+Au vs p+p at RHIC)
 - smallest possible x reach: no mass, no fragmentation
 - no cold matter final-state effects (E-loss, ...)
 - large EPS09 / CGC difference in forward rapidity predictions



Arleo, Eskola, Paukkunen and Salgado (2011)



Jalilian-Marian and Rezaeian (2012)

Confirmation of forward rapidity

R_{pA} suppression:

we should push for the FOCAL
upgrade of the ALICE detector

(more details in the Yellow Report)

Saturation signal #2:

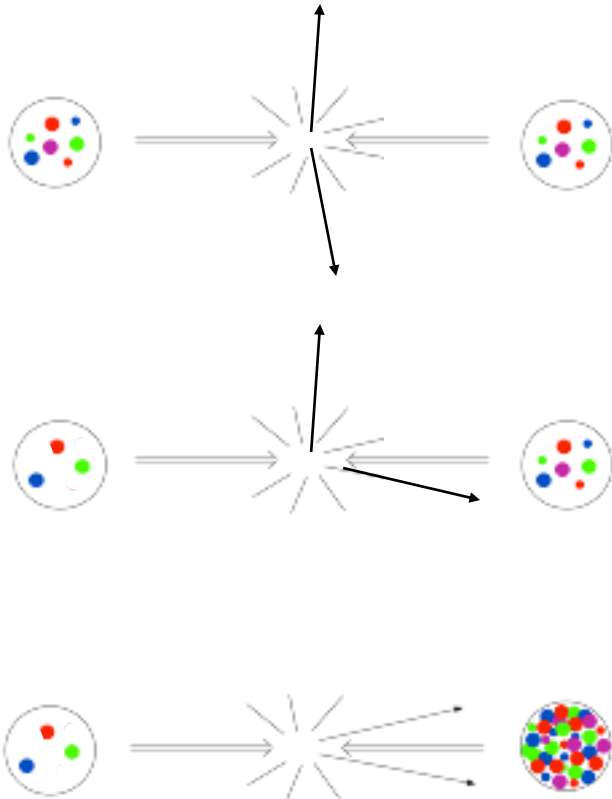
forward rapidity suppression
of di-hadron azimuthal
correlations in $p+A$ vs $p+p$

Di-hadron final-state kinematics

final state : k_1, y_1 k_2, y_2

$$x_p = \frac{k_1 e^{y_1} + k_2 e^{y_2}}{\sqrt{s}} \quad x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}}$$

scanning the wave functions:



$$x_p \sim x_A < 1$$

central rapidities probe moderate x

$$x_p \text{ increases} \quad x_A \sim \text{unchanged}$$

$$x_p \sim 1, x_A < 1$$

forward/central doesn't probe much smaller x

$$x_p \sim \text{unchanged} \quad x_A \text{ decreases}$$

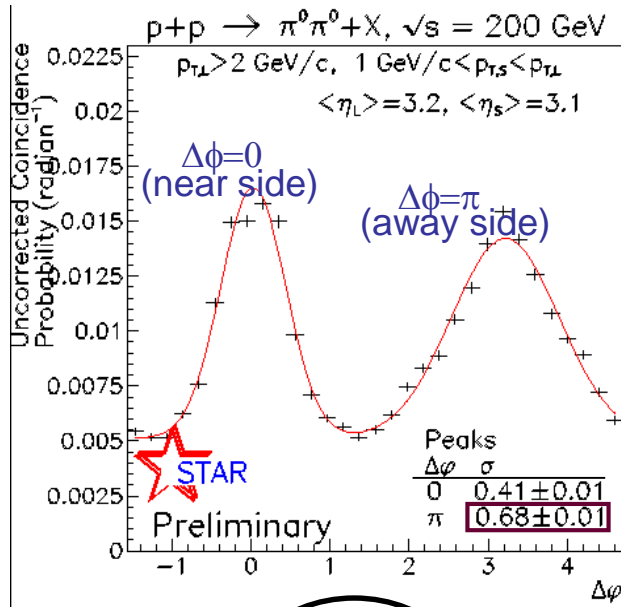
$$x_p \sim 1, x_A \ll 1$$

forward rapidities probe small x

Di-hadron angular correlations

comparisons between d+Au \rightarrow $h_1 h_2 X$ (or p+Au \rightarrow $h_1 h_2 X$) and p+p \rightarrow $h_1 h_2 X$

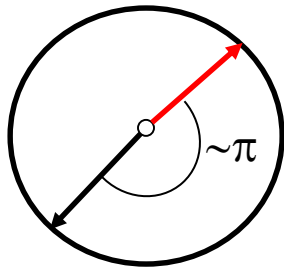
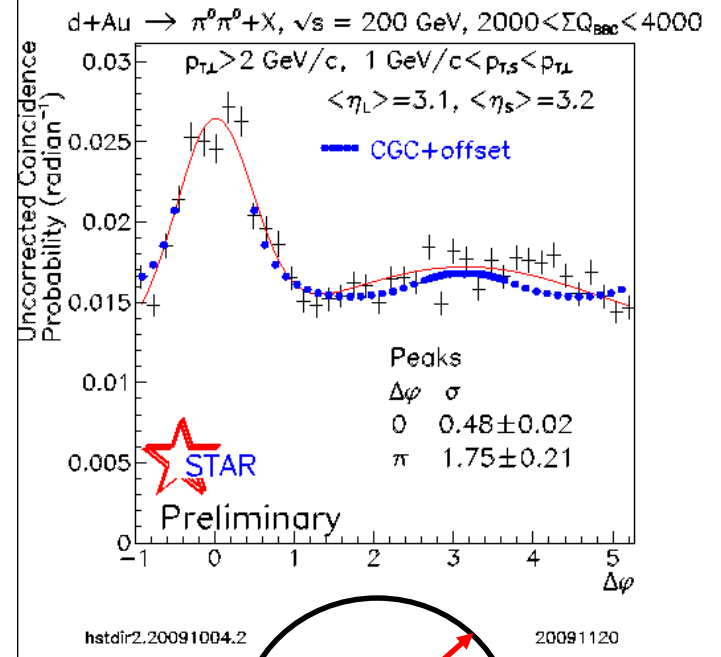
p+p collisions



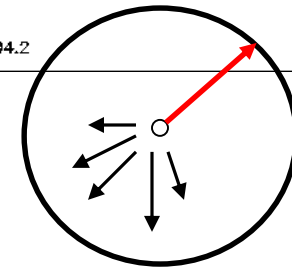
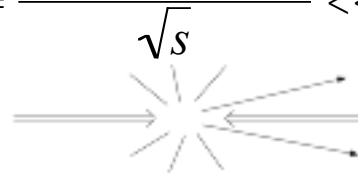
$$\frac{1}{N_{trig}} \frac{dN_{pair}}{d\Delta\phi}$$

Albacete and CM (2010)

central d+Au collisions



$$x_A = \frac{k_1 e^{-y_1} + k_2 e^{-y_2}}{\sqrt{s}} \ll 1$$

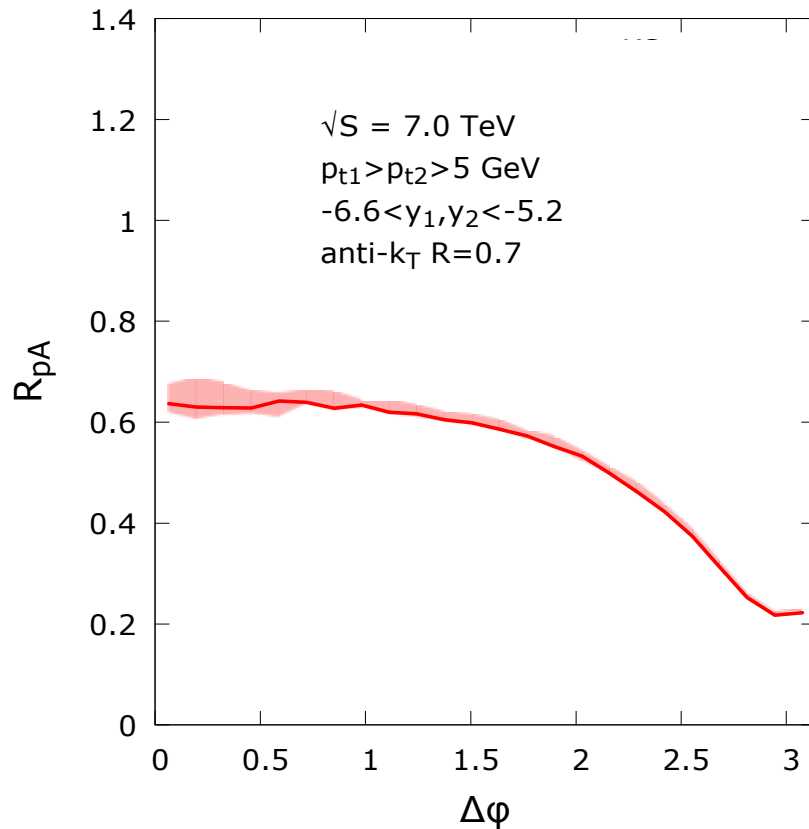


however, when $y_1 \sim y_2 \sim 0$ (and therefore $x_A \sim 0.03$), the p+p and d+Au curves are almost identical

R_{pA} of forward-forward di-jets

- at the LHC this can be done with di-jets!

strong nuclear modification predicted due to strong non-linear effects



Albacete
and CM (2010)

the kinematics for this plot
are chosen assuming that the
di-jets are detected using
the CASTOR detector

for detection in the central part of CMS (i.e. $3 < y < 5$),
 R_{pA} goes from 1 to 0.6 with increasing $\Delta\phi$

Confirmation of forward-forward
azimuthal decorrelation:

we should push for p+Pb data
taking with ZDC+T2+CASTOR

Conclusions

- Fundamental consequence of QCD dynamics:
 - at asymptotically small x /large A , QCD evolution becomes non-linear
- Non-linear evolution of gluon density in Au at RHIC:
 - suppression of single hadron production in d+Au vs p+p
 - suppression of back-to-back correlations of di-hadrons in d+Au vs p+p
- Awaiting more LHC p+Pb forward rapidity data (so far only quarkonia which are not the best probe of saturation physics)
- The most sensitive measurements could be done with FOCAL upgrade of ALICE (isolated photons) and p+Pb data taking with ZDC-T2-CASTOR (forward-forward dijets)