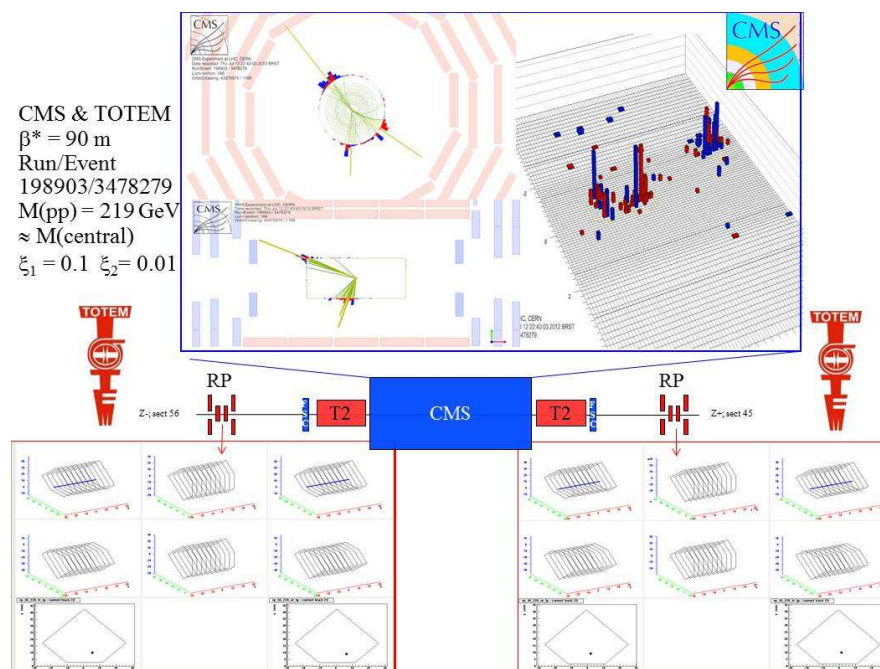


# Search for a di-photon resonance using CT-PPS: Introduction

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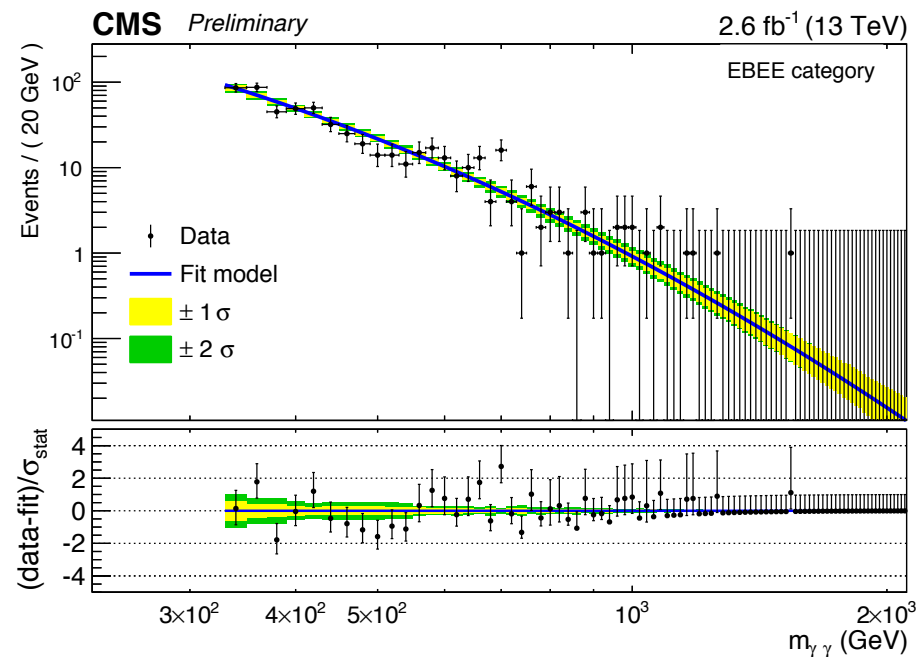
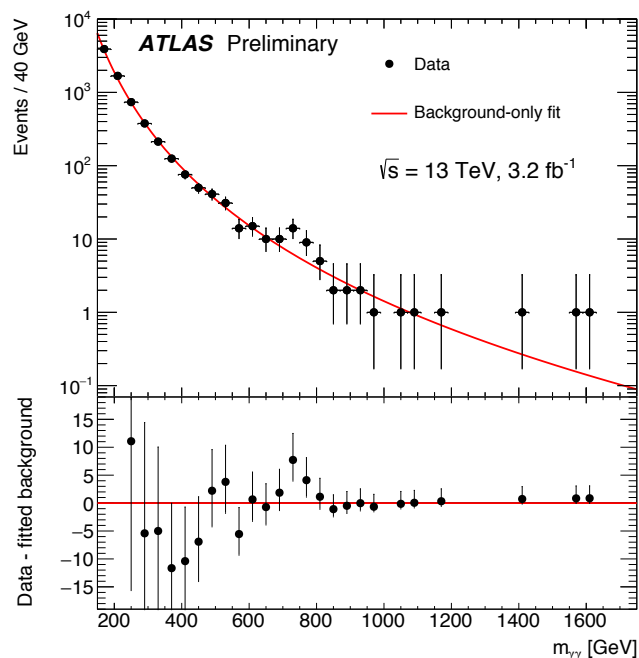
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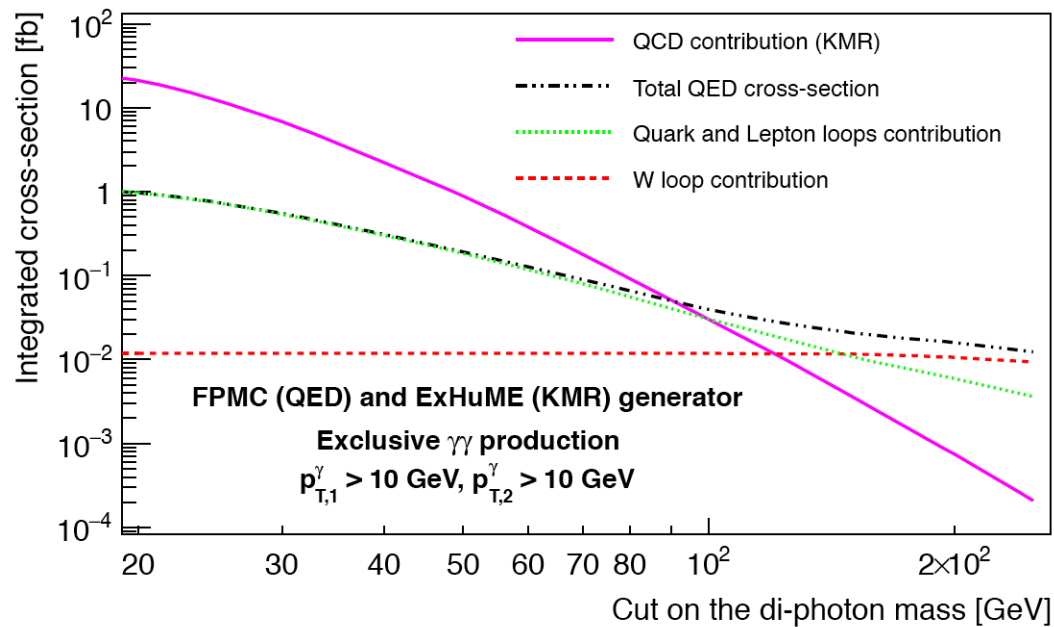
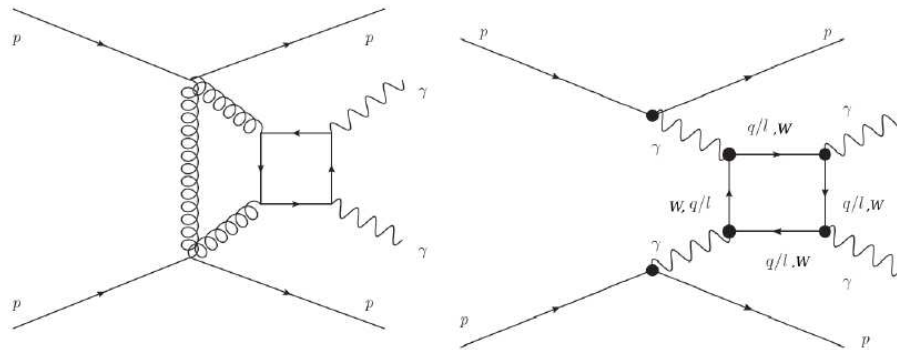
- Selection of  $\gamma$ -induced processes
- Analysis flow
- Existence of a new particle?

## CMS and ATLAS observation



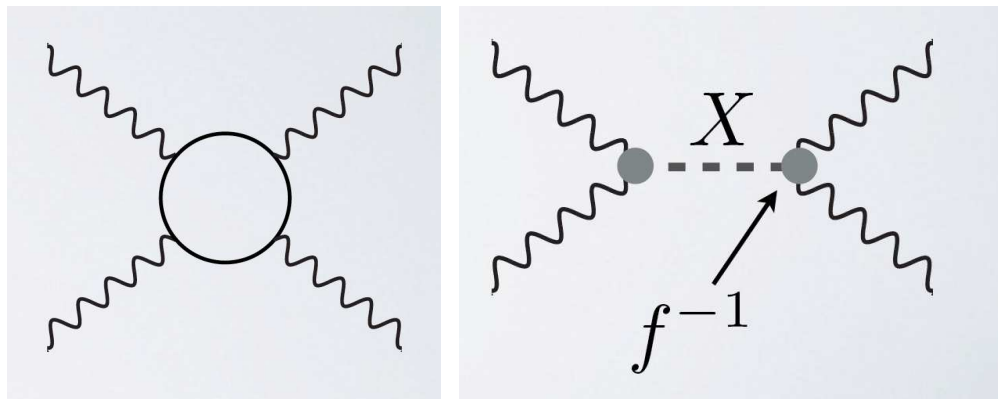
- Potential excess observed by ATLAS and CMS in the diphoton spectrum at 13 TeV
- Can CT-PPS give additional information about this excess if confirmed with more data?
- Specificities of CT-PPS with respect to standard CMS/ATLAS standard searches without tagging the protons
- Do we have a natural explanation for such an excess if confirmed?  
Many publications on ArXiv...

## SM $\gamma\gamma$ exclusive production



- QCD production dominates at low  $m_{\gamma\gamma}$ , QED at high  $m_{\gamma\gamma}$
- Important to consider  $W$  loops at high  $m_{\gamma\gamma}$
- At high masses ( $\sim 750 \text{ GeV}$ ), the photon induced processes are dominant

## Looking for quartic $\gamma\gamma$ anomalous couplings



- Two effective operators at low energies

$$\mathcal{L}_{4\gamma} = \zeta_1^\gamma F_{\mu\nu} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \zeta_2^\gamma F_{\mu\nu} F^{\nu\rho} F_{\rho\lambda} F^{\lambda\mu}$$

- $\gamma\gamma\gamma\gamma$  couplings can be modified in a model independent way by loops of heavy charge particles

$$\zeta_1 = \alpha_{em}^2 Q^4 m^{-4} N c_{1,s}$$

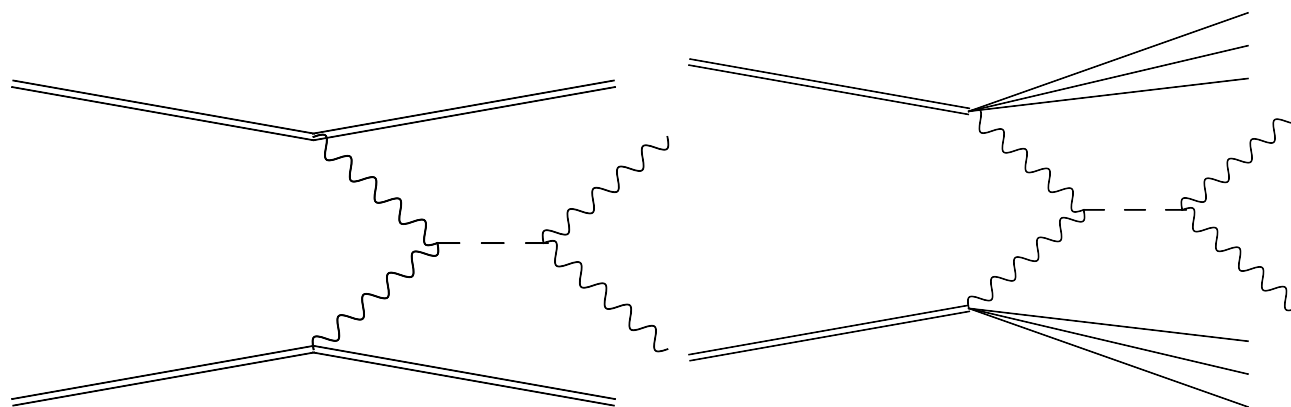
where the coupling depends only on  $Q^4 m^{-4}$  (charge and mass of the charged particle) and on spin,  $c_{1,s}$  depends on the spin of the particle

This leads to  $\zeta_1$  of the order of  $10^{-14}$ - $10^{-13}$

- $\zeta_1$  can also be modified by neutral particles at tree level (extensions of the SM including scalar, pseudo-scalar, and spin-2 resonances that couple to the photon)  $\zeta_1 = (f_s m)^{-2} d_{1,s}$  where  $f_s$  is the  $\gamma\gamma X$  coupling of the new particle to the photon, and  $d_{1,s}$  depends on the spin of the particle; for instance, 2 TeV dilatons lead to  $\zeta_1 \sim 10^{-13}$

## Possible explanation of the di-photon excess

- S. Fichtel, G von Gersdorff, C. Royon, <http://arxiv.org/pdf/1512.05751.pdf>
- The ATLAS and CMS collaborations measured the sum of the elastic and inelastic contributions



- Two important points:
  - We can select photon-induced processes in a background free mode by tagging the protons in CT-PPS: model independent
  - We can try to interpret present results in terms of the existence of a new particle
- Let us explore the possibility of a spin 0/ spin 2 resonance

## Possible explanation of the di-photon excess

- The explanation must explain two experimental facts: The excess was not observed at 8 TeV, and is also not seen in dijets at 13 TeV with the limited statistics
- If processes are gluon-induced, we would expect the dijet cross section to be of the order of 1 pb (the ratio of partial widths is  $\Gamma_{gg}/\Gamma_{\gamma\gamma} = \alpha_S^2/\alpha^2 \sim 200$ )  $\rightarrow$  natural to consider  $\gamma$ -exchange processes
- When protons are not tagged (the present case), one is dominated by inelastic events (inelastic production of diphotons via photon fusion): the ratio between the total (inelastic, inelastic-elastic, elastic) and the elastic contribution is about a factor 20

$$\sigma_{pp \rightarrow \gamma\gamma X} = (7.3 \text{ fb}) \left( \frac{5 \text{ TeV}}{f_\gamma} \right)^4 \left( \frac{45 \text{ GeV}}{\Gamma_{\text{tot}}} \right) \left( \frac{r_{\text{inel}}}{20} \right) r_{\text{fs}}$$

- Why do not we observe anything at 8 TeV? This is due to the probability to emit a quasi-real photon. The  $\xi$  of the photon has to be much higher at 8 TeV than at 13 TeV and the production is much suppressed (factor estimated to be between 2.4 and 3.9).
- We follow the same strategy as the search for quartic anomalous couplings: same final state
- The resonance production matrix element has been fully implemented in FPMC

## A cleaner measurement: $\gamma\gamma$ exclusive production

- The idea is now to consider elastic production:
  - The theoretical calculation is in better control (QED processes with intact protons), not sensitive to the photon structure function
  - This is a “background-free” experiment (see following slides) and any observed event is signal
  - The survival probability is in better control than in the QCD (gluon) case
  - This is complementary to the search in the jet channel (see following slides)

- Using FPMC and the “resonance” production parameters:

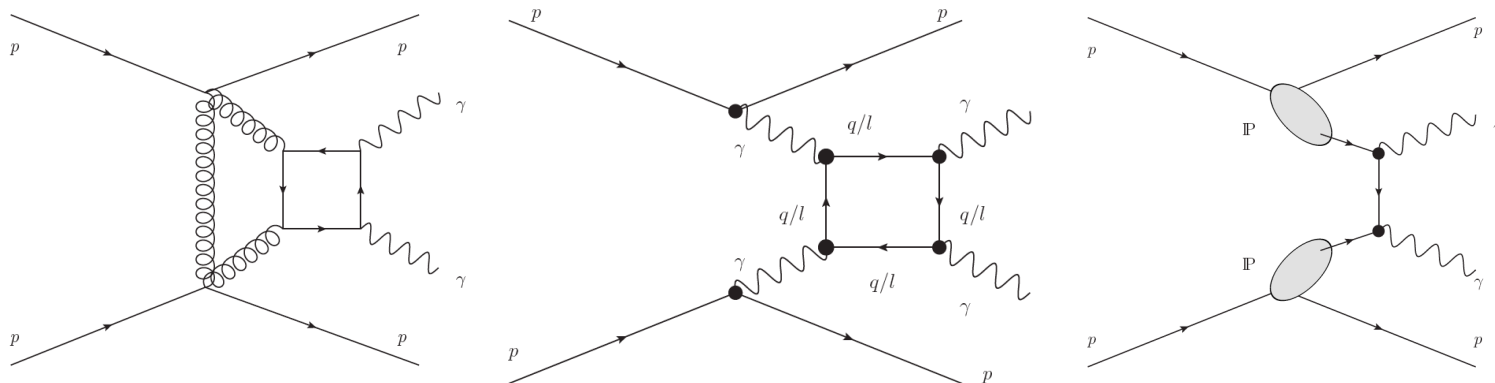
$$\sigma_{pp \rightarrow \gamma\gamma pp} = (0.23 \text{ fb}) \left( \frac{5 \text{ TeV}}{f_\gamma} \right)^4 \left( \frac{45 \text{ GeV}}{\Gamma_{\text{tot}}} \right) r_{fs}$$

with  $f_\gamma \sim 4.6 \text{ TeV}$

- Since we do not have background in this channel, observing 6 events is a  $5 \sigma$  discovery:  $21 \text{ fb}^{-1}$  is needed to reach this number of events

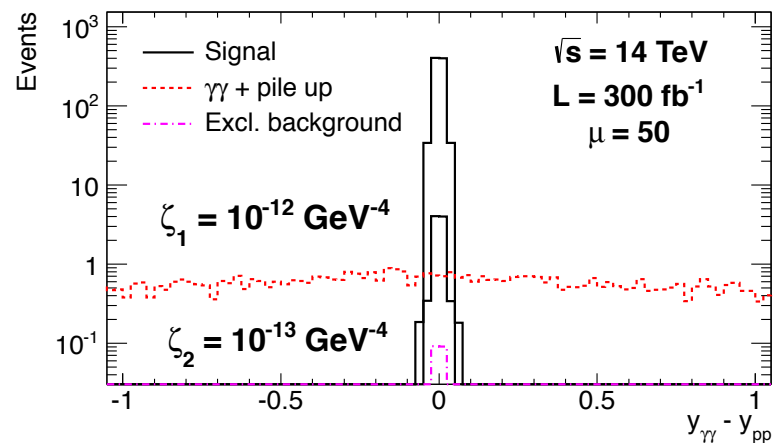
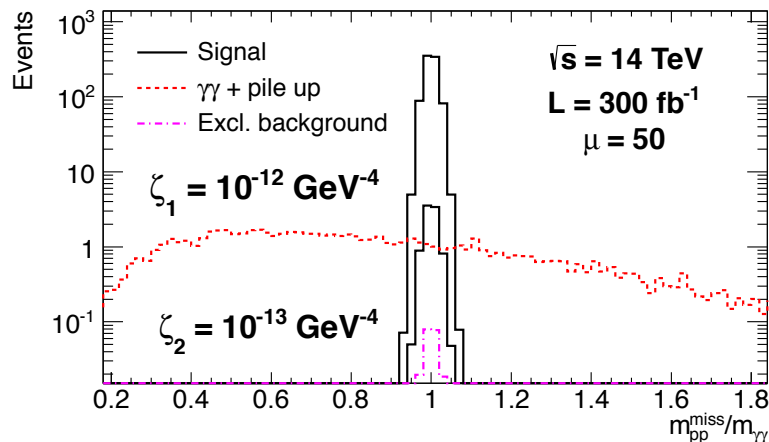
## Search for resonance/ $\gamma\gamma\gamma\gamma$ anomalous couplings: Analysis flow

- Studies performed at hadron level but taking into account the main detector/pile-up effects
- By default,  $> 1\gamma$  converted is requested (1 mm resolution), but all  $\gamma$  are also considered
- pile-up simulated in AFP/CT-PPS: 50, 100, 200...
- Main detector effects are included (from ATLAS ECFA studies ATL-PHYS-PUB-2013-009), for instance:
  - Photon conversion probability: 15% in barrel, 30% in the end-caps;  $\gamma$  rapidity,  $\Phi$ , and  $p_T$  resolutions taken into account as well as the reconstruction efficiency
  - Misidentification of electron as a  $\gamma$ : 1%
  - Misidentification of jet as a  $\gamma$ : 1/4000,
- All backgrounds were considered: DPE diphoton production, Higgs decaying into photons, exclusive production of diphoton, dilepton, dijet with lepton/jet misidentified, pile up (ND production of Drell-Yan, dijet, diphoton...)





## Search for resonance/ quartic $\gamma\gamma$ anomalous couplings



Cut / Process	Signal (full)	Signal with (without) f.f (EFT)	Excl.	DPE	DY, di-jet + pile up	$\gamma\gamma$ + pile up
$[0.015 < \xi_{1,2} < 0.15,$ $p_{T1,(2)} > 200, (100) \text{ GeV}]$	130.8	36.9 (373.9)	0.25	0.2	1.6	2968
$m_{\gamma\gamma} > 600 \text{ GeV}$	128.3	34.9 (371.6)	0.20	0	0.2	1023
$[p_{T2}/p_{T1} > 0.95,$ $ \Delta\phi  > \pi - 0.01]$	128.3	34.9 (371.4)	0.19	0	0	80.2
$\sqrt{\xi_1 \xi_2} s = m_{\gamma\gamma} \pm 3\%$	122.0	32.9 (350.2)	0.18	0	0	2.8
$ y_{\gamma\gamma} - y_{pp}  < 0.03$	119.1	31.8 (338.5)	0.18	0	0	0

- No background after cuts for  $300 \text{ fb}^{-1}$  **without needing timing detector information**
- Exclusivity cuts using proton tagging needed to suppress backgrounds (Without exclusivity cuts using CT-PPS: background of 80.2 for  $300 \text{ fb}^{-1}$ )

## Consequences for CT-PPS analyses and conclusion

- Looking for di-photon production via photon fusion is an ideal case for CT-PPS: this is  $\sim$  background free after exclusivity cuts, we can be sure that they are photon induced processes
- Even if the model discussed above (if the “resonance” is real) is not true, many other models predict a non-zero contribution of photon-induced processes that can be probed using CT-PPS in a clean way
- S. Fichet, G. von Gersdorff, C. Royon, <http://arxiv.org/pdf/1601.01712.pdf>
- CT-PPS allows to probe diphoton production in a model independent way knowing that any observation is a potential signal
- CT-PPS has the possibility to test  $\gamma$ -induced processes in background free experiments:
  - Fundamental if the resonance is photon-induced
  - if this is not completely the case, measure branching ratio
- If one believes the ATLAS/CMS observation, 20-25 fb<sup>-1</sup> of data might be enough for a 5  $\sigma$  discovery in CT-PPS
- We are also going to look into other channels:  $WW$ ,  $ZZ$ ,  $Z\gamma$  (specially interesting)
- See talks by Gero/Valery

**High lumi: Search for quartic  $\gamma\gamma$  anomalous couplings:**  
**Results from effective theory**

Luminosity	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	3000 fb <sup>-1</sup>
pile-up ( $\mu$ )	50	50	50	200
coupling (GeV <sup>-4</sup> )	$\geq 1$ conv. $\gamma$ 5 $\sigma$	$\geq 1$ conv. $\gamma$ 95% CL	all $\gamma$ 95% CL	all $\gamma$ 95% CL
$\zeta_1$ f.f.	$8 \cdot 10^{-14}$	$5 \cdot 10^{-14}$	$3 \cdot 10^{-14}$	$2.5 \cdot 10^{-14}$
$\zeta_1$ no f.f.	$2.5 \cdot 10^{-14}$	$1.5 \cdot 10^{-14}$	$9 \cdot 10^{-15}$	$7 \cdot 10^{-15}$
$\zeta_2$ f.f.	$2 \cdot 10^{-13}$	$1 \cdot 10^{-13}$	$6 \cdot 10^{-14}$	$4.5 \cdot 10^{-14}$
$\zeta_2$ no f.f.	$5 \cdot 10^{-14}$	$4 \cdot 10^{-14}$	$2 \cdot 10^{-14}$	$1.5 \cdot 10^{-14}$

- **Unprecedented sensitivities at hadronic colliders:** no limit exists presently on  $\gamma\gamma\gamma\gamma$  anomalous couplings
- Reaches the values predicted by extra-dim or composite Higgs models
- Pile up background rejected using exclusivity cuts: timing detectors not used in this analysis
- Introducing form factors to avoid quadratical divergences of scattering amplitudes due to anomalous couplings in conventional way:

$$a \rightarrow \frac{a}{(1+W\gamma\gamma/\Lambda_{cutoff})^2} \text{ with } \Lambda_{cutoff} \sim 2 \text{ TeV, scale of new physics}$$

## Full amplitude calculation

- Effective field theory valid if  $S \ll 4m^2$ ,  $S$  smaller than the threshold production of real particles
- Since the maximum proton missing mass is  $\sim 2$  TeV at the 14 TeV LHC, the effective theory needs to be corrected for masses of particles below  $\sim 1$  TeV  $\rightarrow$  use of form factor which creates an uncertainty on the results (depends on the exact value of form factors)
- Solution: compute the full momentum dependence of the 4 photon amplitudes: computed for fermions and bosons
- Full amplitude calculation for generic heavy charged fermion/vector contribution
- Existence of new heavy charged particles enhances the  $\gamma\gamma\gamma\gamma$  couplings in a model independent way
- Enhancement parametrised with particle mass and effective charge  $Q_{eff} = QN^{1/4}$  where  $N$  is the multiplicity

## Search for quartic $\gamma\gamma$ anomalous couplings: Results from full theory

Cut / Process	Signal (full)	Signal with (without) f.f (EFT)	Excl.	DPE	DY, di-jet + pile up	$\gamma\gamma$ + pile up
$[0.015 < \xi_{1,2} < 0.15,$ $p_{T1,(2)} > 200, (100) \text{ GeV}]$	130.8	36.9 (373.9)	0.25	0.2	1.6	2968
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$ y_{\gamma\gamma} - y_{pp}  < 0.03$	119.1	31.8 (338.5)	0.18	0	0	0

- No background after cuts for  $300 \text{ fb}^{-1}$  **without needing timing detector information**
- For signal: 119.1 events for  $Q_{eff} = 4, m = 340 \text{ GeV}$
- Results for full calculation lay between the effective field result with/without form factor as expected since effective calculation not valid in the region of  $S \sim m^2$

## Full amplitude calculation

- 5  $\sigma$  discovery sensitivity on the effective charge of new charged fermions and vector boson for various mass scenarii for 300  $fb^{-1}$  and  $\mu = 50$

Mass (GeV)	300	600	900	1200	1500
$Q_{\text{eff}}$ (vector)	2.2	3.4	4.9	7.2	8.9
$Q_{\text{eff}}$ (fermion)	3.6	5.7	8.6	-	-

- Unprecedented sensitivities at hadronic colliders reaching the values predicted by extra-dim models - For reference, we also display the result of effective field theory (without form factor) which deviates at low masses from the full calculation
- For  $Q_{\text{eff}} = 4$ , we are sensitive to new vectors (fermions) up to 700 (370) GeV for a luminosity of 300  $fb^{-1}$

