# LIGHT BY LIGHT SCATTERING AND THE 750 GEV DIPHOTON EXCESS

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Based on work with Sylvain Fichet & Christophe Royon 1512.05751 and 1601.01712







International Centre for Theoretical Physics<br>South American Institute for Fundamental Research

INTRODUCTION

#### THE EXCESS AT 750 GEV - ATLAS

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#### $ATLAS 3.2 fb^{-1} (13 TeV)$

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- Global: 2.3 σ $\bullet$

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#### Compare with the Higgs (roughly 3x the statistics…)



#### THE EXCESS AT 750 GEV - CMS  $\mathbf{I}$ -2 0 2



 $\sim$  CMS: 20 fb<sup>-1</sup> (8 TeV) + 2.6 fb<sup>-1</sup> (13 TeV) photon candidates are matched to those selected in the analysis using a k-nearest-neighbourse  $\sim$ algorithm, <mark>● Local: 3.0 σ</mark> **Figure 4 shows, in an analysis of the measured contributions of the different background compositions of the di** 

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- The cross section is roughly  $\sigma(pp \to \phi)$   $BR(\phi \to \gamma \gamma) \sim 10$  fb
- No electric charge *Q=0*

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# PART I: PHOTON FUSION PRODUCTION

#### THE PRODUCTION MECHANISM



1  $f_\gamma$  $\phi F_{\mu\nu}F^{\mu\nu}$  + 1  $f_g$  $\overline{\phi}$   $G_{\mu\nu}$   $G^{\mu\nu}$ *γ γ*

#### THE PRODUCTION MECHANISM







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1  $f_\gamma$  $\phi$   $F_{\mu\nu}F^{\mu\nu}$  + 1  $f_{q}$  $q \sum_{\mu} \oint_{-\infty} q^2 y$   $\mathcal{L}_{eff} = \left[ \frac{1}{c} \phi F_{\mu\nu} F^{\mu\nu} \right] + \left[ \frac{1}{c} \phi \bar{q} q \right]$ 

**Photon Fusion** 



1  $f_\gamma$  $\overline{\phi}$   $F_{\mu\nu}$   $F^{\mu\nu}$ 

Let's assume:

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- 1. What is the cross section in terms of effective coupling *f<sup>γ</sup> -1* ?
	- Model independent!
	- Infer *fγ* from data!



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- 2. What is the effective coupling in terms of fundamental parameters?
	- Model-dependent!
	- Perturbativity?



Fichet, GG, Royon 1512.05751 see also:

Csaki, et al : 1601.00638 Harland-Lang et al, 1601.07187

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Computation of cross section gives:  $\bullet$ 

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\sigma_{pp \to \gamma \gamma X} = 5.9 \text{ fb} \left( \frac{5 \text{ TeV}}{f_{\gamma}} \right)^4 \frac{45 \text{ GeV}}{\Gamma_{\phi}} \times \frac{r_{fs}}{0.8} \times \frac{r_{inel}}{20} \qquad \frac{r_{fs} \sim 0.6...1.0}{r_{inel} \sim 15...25}
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*f<sub>γ</sub>* = 3.9 - 4.9 TeV (68 % CL)  $\left[\right. \overline{o(pp \rightarrow \gamma \gamma X)}$  = 8.6 - 20 fb ] **•**  $f_y = 3.6 - 6.8$  TeV (95 % CL)  $\sigma(pp \to \gamma \gamma X) = 2.5 - 33$  fb ]

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Determination of *fγ* fairly accurate due to *fγ ~ σ -1/4*

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$$
\frac{1}{f_{\gamma}} = \alpha \frac{\lambda}{4\pi} Q^2 N \frac{2}{m_{\phi}} B(\tau), \quad B(\tau) = \sqrt{\tau} \left[ 1 + (1 - \tau) \arcsin^2 \left( \tau^{-\frac{1}{2}} \right) \right], \quad \tau = \frac{4m_{\psi}^2}{m_{\phi}}
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- Width and cross section fix two combinations of *Q, N, m<sup>ψ</sup> ,λ*
- For instance:  $Q = 5/2$ ,  $N = 3$ ,  $m_\psi = 360$  GeV,  $\lambda = 5$
- Still perturbative:  $\lambda N^{1/2} \sim 8.6 < 4\pi$

1

 $f_\gamma$ 

#### CONSTRAINING THE UV COMPLETION

- 
- 
- 
- 
- 
- -

Drell-Yan (dilepton), constrains *NQ2*



Gross et al 1602.03877 Goertz et al 1602.04801



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- and find that the 8 TeV Drell-Yan data already exclude Direct production of *ψ* at LHC…
- Landau Poles + Vacuum stability





 $\mathbb{P}^1(\mathbb{R}^d)$  is the second in terms of mass and e $\mathbb{P}^1(\mathbb{R}^d)$  and vectors. in terms and vectors. in

T. Melia, M. Papucci and K. Zurek, arXiv:1512.04928

the case of no requirements of photon conversion at the analysis stage and full integrated luminosity stage and at the medium-luminosity LHC (300 fb<sup>1</sup> [hep-ph]; D. Buttazzo, A. Greljo and D. Marzocca,  $T1$  . The physical distribution  $\mathcal{L}$  $2016$  1, 015017  $\pm151$   $\pm151$ 

# PART II: MEASURING THE *φγγ* COUPLING

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#### Now (part II):

- No assumptions on couplings or production mode ( 100% model independent )
- Is there a way to measure the *φγγ* coupling?  $\bullet$

• Inelastic production (proton destroyed, dominant)



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Elastic production (protons intact, subdominant )



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#### $~\rm{Suppression:}~~~\sim10^{-5}~~~\sim10^{-1}$

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‣ All inelastic events can be completely rejected

- ‣ Essentially background-free (pile up under control)
- ‣ Installed in CMS, planned in ATLAS

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Allows precision measurement of diphoton coupling!

### CROSS SECTION

- Elastic cross section under excellent theoretical control
- With realistic cuts the cross section is

$$
\sigma_{pp\to\gamma\gamma pp} = 0.23 \text{ fb} \left( \frac{5 \text{ TeV}}{f_{\gamma}} \right)^4 \frac{45 \text{ GeV}}{\Gamma_{\phi}} \quad \text{Fichet, GG, Royon} \quad 1601.01712
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Exclusion power at 68% (95%)

$$
f_{\gamma} > 14 \text{ (11) TeV}
$$
\n
$$
f_{\gamma} > 25 \text{ (19) TeV}
$$





Region preferred by diphoton excess at  $\Gamma_{\phi}$  = 45 GeV



We would like to stress that our proposal to measure Region excluded  $\frac{3}{\ln 2}$  in particular does not make an in particular does not make a  $\frac{3}{\ln 2}$  $\sum_{i=1}^n$  by  $\frac{1}{n}$  and  $\frac{1}{n}$ **Fig. 1** dijet searches by Run - I (8 TeV)





- EXPERIMENTAL CONSTRAINTS FIG. 3. Bounds and sensitivities in the *f f<sup>g</sup>* plane, in case  $\blacktriangleright$  Dijet searches and elastic γγ fusion are complementary
- $\longrightarrow$  More dat and tegron b and 95% C.L. credible regions corresponding to the observed  $\blacktriangleright$  More data will improve both bounds and can cover the entire Limit of the region above which *EW* <sup>+</sup>*gg* tot = 45 GeV. Dotted (dashed) lines correspond to *EW /* = 1*.*64 (53.9) region predicted by the diphoton excess

## QUARK VS PHOTON COUPLING

Region excluded by  $Run - I$   $(8 TeV)$ dijet searches



in proferred  $T_{\text{preretred}}$ by diphoton excess at  $\Gamma_\phi$  = 45 GeV We will first argue that *elastic gluon fusion* (EGGF), Region preferred

Elastic γγ fusion:  $T_{\text{rad}}$  $protonance$ region at 300 fb<sup>-1</sup> 95% excludable

emission in the gluon ladder has to be suppressed in order

#### CONSTRAINTS FROM GAUGE INVARIANCE

$$
c_{\phi\gamma\gamma} = \boxed{\frac{1}{f_\gamma}\phi\,(F_{\mu\nu})^2}
$$

 $\bullet$  Up to now:  $\qquad \mathcal{L}_{\phi\gamma\gamma} = \frac{1}{f} \phi \left( F_{\mu\nu} \right)^2$  (not SU(2)xU(1) invariant!)

#### CONSTRAINTS FROM GAUGE INVARIANCE

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\mathcal{L}_{\phi\gamma\gamma} = \frac{\left(\frac{1}{f_{\gamma}}\phi(F_{\mu\nu})^2\right)}{\sum_{\substack{f_{\gamma} \\ f_{\gamma}}} \frac{1}{f_{\gamma}} = \frac{s_w}{f_W} + \frac{c_w}{f_B}} \left(\frac{1}{f_B}\phi(H_{\mu\nu})^2\right)} \quad \text{(not SU(2)xU(1) invariant!)}
$$
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 $\bullet\ \textbf{Up to now:}\quad \quad \mathcal{L}_{\phi\gamma\gamma}=\boxed{\frac{\tau}{f}\phi\left(F_{\mu\nu}\right)^2} \quad \quad \textbf{(not SU(2)xU(1) invariant!)}$ 

*f<sup>B</sup>*

$$
\mathcal{L}_{\phi VV} = \underbrace{\left(\frac{1}{f_W} \phi \left(W_{\mu\nu}\right)^2\right)}_{\text{FW}} + \underbrace{\frac{1}{f_B} \phi \left(B_{\mu\nu}\right)^2}_{\text{FB}} + \underbrace{\frac{c_w}{f_B}}_{\text{FW}}
$$

- Diphoton coupling comes with other couplings
- Two parameters, express in terms of  $f_\gamma$ ,  $r = f_W / f_B$ :

$$
\mathcal{L}_{eff} = \frac{1}{f_{\gamma}} \phi \left( F_{\mu\nu}^2 + \frac{2s_w c_w (r-1)}{c_w^2 r + s_w^2} Z_{\mu\nu} F^{\mu\nu} + \frac{s_w^2 r + c_w^2}{c_w^2 r + s_w^2} Z_{\mu\nu}^2 + \frac{2}{c_w^2 r + s_w^2} |W_{\mu\nu}|^2 \right)
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$$

Constraints on these couplings exist from run-1 diboson measurements

Strongest constraint:  $Z\gamma$ 

$$
\sigma_{Z\gamma}^8\, {\rm TeV}\, = \sigma_{\gamma\gamma}^{13} \; {\rm TeV}\;\; \frac{\sigma_{\gamma\gamma}^8\, {\rm TeV}}{\sigma_{\gamma\gamma}^{13}\, {\rm TeV}}\;\, \frac{\Gamma_{Z\gamma}}{\Gamma_{\gamma\gamma}}
$$

 $\bullet$  Strongest constraint: Zγ

$$
\gamma \qquad \qquad \boxed{\sigma_{Z\gamma}^8 \text{TeV}} = \sigma_{\gamma\gamma}^{13} \text{ TeV} \quad \frac{\sigma_{\gamma\gamma}^8 \text{TeV}}{\sigma_{\gamma\gamma}^{13} \text{ TeV}} \, \frac{\Gamma_{Z\gamma}}{\Gamma_{\gamma\gamma}}
$$
\n
$$
\text{upper limit (run I)}
$$









CONCLUSIONS





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