

Higher-order corrections to decays and masses of charged Higgs bosons

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Oct. 5, 2016

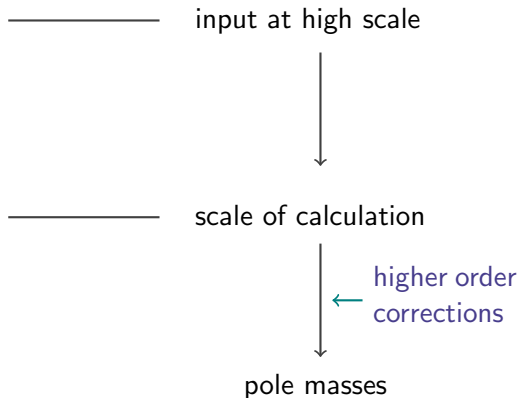
Corrections to the mass of the charged Higgs boson

Mass of the charged Higgs boson:

- Pole/on-shell mass is input in calculations of
 - * production cross sections
 - * partial decay widths
- In perturbative calculations:
Renormalization schemes dependences

Choice of renormalization schemes

- Running parameters (\overline{MS} or \overline{DR})



- On-shell parameters (pole masses, etc.)

input parameters \longleftrightarrow physical observables

Two-Higgs-Doublet Model

Higgs potential:

$$\begin{aligned} V = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 + m_{12}^2 (\Phi_1^\dagger \Phi_2 + \text{h.c.}) \\ & + \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) \\ & + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \lambda_5 \left[(\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2 \right] \end{aligned}$$

- CP conserving
 - invariant under $\Phi_1 \rightarrow -\Phi_1$
 - m_{11}^2, m_{22}^2 fixed by minimum condition
 $m_{12}^2, \lambda_1, \dots, \lambda_5$ free parameters
- \Rightarrow enough free parameters to define all Higgs masses independently
- \Rightarrow all Higgs masses can be chosen as pole masses (on-shell)

Higgs potential:

$$\begin{aligned} V = & (m_1^2 + |\mu|^2)\Phi_1^\dagger\Phi_1 + (m_2^2 + |\mu|^2)\Phi_2^\dagger\Phi_2 + m_{12}^2(\Phi_1^\dagger\Phi_2 + \text{h.c.}) \\ & + \frac{g^2 + g'^2}{8}(\Phi_1^\dagger\Phi_1)^2 + \frac{g^2 + g'^2}{8}(\Phi_2^\dagger\Phi_2)^2 - \frac{g^2 + g'^2}{4}(\Phi_1^\dagger\Phi_1)(\Phi_2^\dagger\Phi_2) \\ & + \frac{g^2}{2}(\Phi_1^\dagger\Phi_2)(\Phi_2^\dagger\Phi_1) \end{aligned}$$

- $m_1^2 + |\mu|^2$, $m_2^2 + |\mu|^2$ fixed by minimum condition

g, g' gauge couplings (fixed in gauge sector)

m_{12}^2 free parameter

⇒ **too** few free parameters to define all Higgs masses independently

⇒ **only one** mass can be chosen independently

(and $\tan\beta =$ ratio of vac. expect. values)

Choices:

- CP-conserving MSSM:

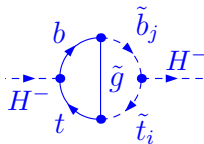
often CP-odd Higgs boson mass as free parameter

⇒ all other masses can be calculated

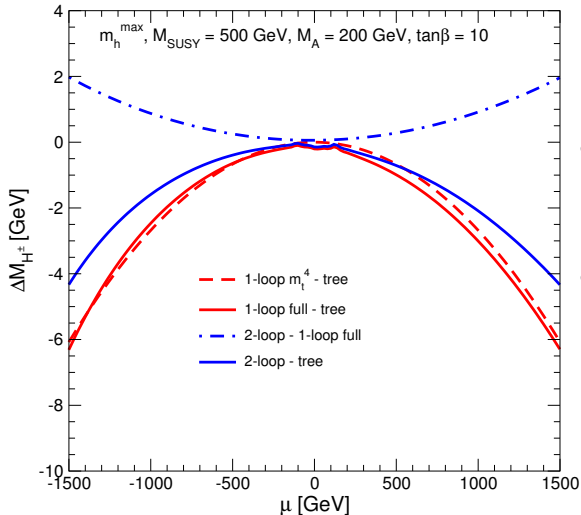
relations change due to
quantum corrections

Charged Higgs mass:

$$M_{H^\pm}^2 = \underbrace{M_A^2 + M_W^2}_{=M_{H^\pm}^2{}_{\text{tree}}} - \hat{\Sigma}_{H^+H^-}$$



Corrections to the mass of the charged Higgs boson up to $\mathcal{O}(\alpha_t \alpha_s)$:



- comparison at one-loop:
 - $\frac{m_t^4}{M_W^2}$ terms: good approx.
- 2-loop corrections:
 - * relevant part
 - * \sim up to 2 GeV in example

[Frank, Galeta, Hahn, Heinemeyer, Hollik, H.R., Weiglein, arXiv:1306.1156]

Choices:

- CP-violating MSSM:

neutral Higgs bosons: Loop-corrected mass matrix:

$$\mathbf{M}(p^2) = \begin{pmatrix} M_{H_{\text{tree}}^0}^2 - \hat{\Sigma}_{H^0 H^0}(p^2) & -\hat{\Sigma}_{H^0 h^0}(p^2) & -\hat{\Sigma}_{H^0 A^0}(p^2) \\ -\hat{\Sigma}_{H^0 h^0}(p^2) & M_{h_{\text{tree}}^0}^2 - \hat{\Sigma}_{h^0 h^0}(p^2) & -\hat{\Sigma}_{h^0 A^0}(p^2) \\ -\hat{\Sigma}_{H^0 A^0}(p^2) & -\hat{\Sigma}_{h^0 A^0}(p^2) & M_{A_{\text{tree}}^0}^2 - \hat{\Sigma}_{A^0 A^0}(p^2) \end{pmatrix}$$

all neutral Higgs bosons **mix**

⇒ choose charged Higgs boson as input parameter

⇒ can be fixed as pole mass (no quantum corrections)

(CP-odd Higgs boson mass will receive corrections)

- Additional Higgs singlet superfield
 - ⇒ an additional CP-even and CP-odd Higgs boson
- Additional parameters (+ 3 to 4 complex parameters)
- Even in CP-conserving case:
 - mixing effects for CP-even and CP-odd Higgs bosons
 - ⇒ choose at least charged Higgs boson as free parameter
 - ⇒ can be fixed as pole mass (no quantum corrections)

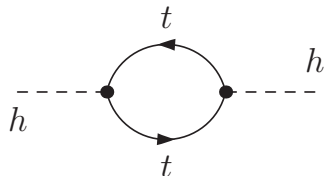
Side remark: Higgs boson masses at higher orders

Two-point-function:

$$-i\hat{\Gamma}(p^2) = p^2 - \mathbf{M}(p^2)$$

with the matrix:

$$\mathbf{M}(p^2) = \begin{pmatrix} M_{h_{\text{Born}}}^2 & -\hat{\Sigma}_{hh}(p^2) & -\hat{\Sigma}_{Hh}(p^2) \\ -\hat{\Sigma}_{Hh}(p^2) & M_{H_{\text{Born}}}^2 & -\hat{\Sigma}_{HH}(p^2) \end{pmatrix}$$



(CP-conserving
MSSM case:
mixing only
betw. CP-even
Higgs bosons h, H)

Calculate the zeros of the determinant of $\hat{\Gamma}$

\Rightarrow loop-corrected Higgs masses

Higgs boson mass for large stop masses

Prediction obtained via Feynman diagrammatic approach:

- + all **log** and **non-log** terms are taken into account
at a **certain order** of perturbation theory
- possible appearance of **large logs**:

$$\Delta M_h \sim \log \frac{M_S}{m_t} \quad M_S = \sqrt{m_{\tilde{t}_1} m_{\tilde{t}_2}}: \text{SUSY particle mass scale}$$

$m_{\tilde{t}_1}, m_{\tilde{t}_2}$: scalar top quark masses
 m_t : top quark mass

- ⇒
- **good** prediction for **lower** SUSY mass scales
 - **no** reliable prediction for **large** SUSY mass scales

Higgs boson mass for large stop masses

Other approach: Renormalization Group Equation (RGE) approach:

see e.g. [Draper, Lee, Wagner, 1312.5743; Vega, Villadoro, arXiv:1504.05200;

Lee, Wagner, arXiv:1508.00576]

★ assume: all SUSY particles are heavy of order $\sim M_S$:

above M_S : MSSM

match at scale M_S :

below M_S : Standard Model

(as effective theory)

$$\lambda^{\text{MSSM}}(M_S) = \lambda^{\text{Standard Model}}(M_S)$$

quartic Higgs coupling

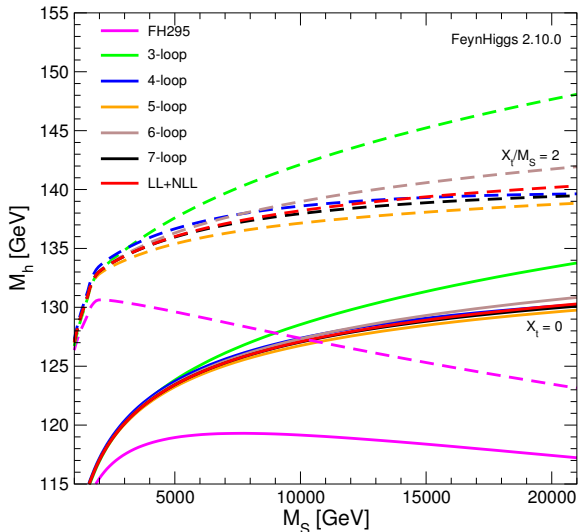
★ evolve λ to lower scale using Standard Model running (RGE)

★ the Higgs mass² is then $M_h^2(m_t) = 2\lambda(m_t)v^2$ $v \approx 174$ GeV

⇒ logs resummed to all orders: good prediction for large SUSY masses

→ Combine both approaches

Higgs boson mass for large stop masses



Comparison of:

★ old FeynHiggs

reliable up to
 $M_s = \mathcal{O}(1\text{TeV})$

★ analyt. solution of RGE:

3-loop ... 7-loop level:

Logs of order

$\mathcal{O}(\alpha_t \alpha_s^2, \alpha_t^2 \alpha_s, \alpha_t^3) \dots$

★ numerical solution:

LL+NLL:

logs resummed

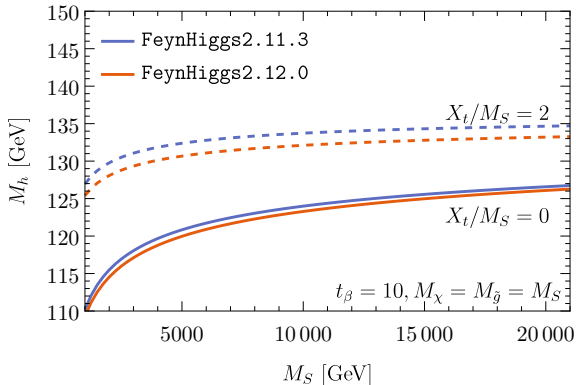
to all orders

$M_A = M_2 = \mu = 1 \text{ TeV}$, $m_{\tilde{g}} = 1.6 \text{ TeV}$, $\tan \beta = 10$

[Hahn, Heinemeyer, Hollik,
 H.R., Weiglein, arXiv:1312.4937]

Higgs boson mass for large stop masses

Improvements:



[Bahl, Hollik, arXiv:1608.01880]

- EW LL, NLL contributions
- $\mathcal{O}(\alpha_s, \alpha_t)$ NNLL contributions

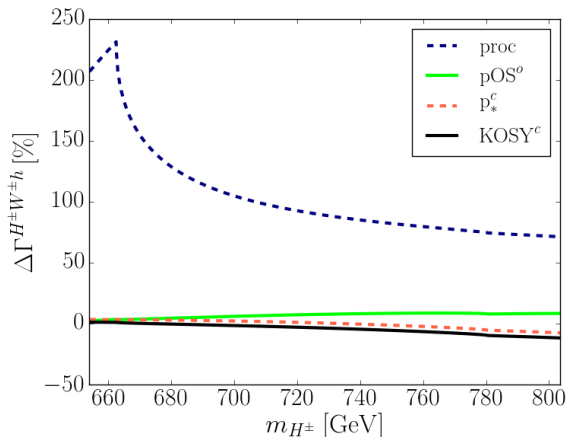
Charged Higgs decays

QCD corrections and dominant electroweak effects:

known for a while, here for $H^+ \rightarrow t\bar{b}$:

- 2HDM: e.g. [Coarasa Perez, Guasch, Hollik, Sola, hep-ph/9808278]
- MSSM: e.g. [Coarasa Perez, Garcia, Guasch, Jimenez, Sola, hep-ph/9711472;
Carena, Garcia, Nierste, Wagner, hep-ph/9912516]

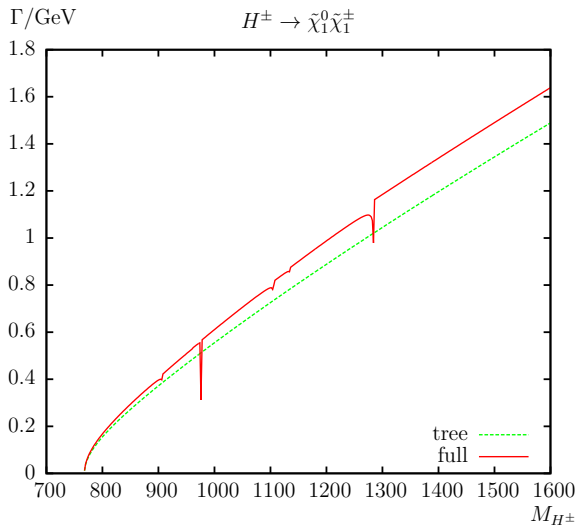
Recent work in 2HDM: $H^+ \rightarrow hW^+$



- complete EW contributions
- different renormalization schemes

[Krause, Lorenz, Mühlleitner, Santos, Ziesche, arXiv:1608.01880]

Recent work in MSSM: $H^+ \rightarrow \chi^0 \chi^+$



[Heinemeyer, Schappacher, arXiv:1503.02996]

- complete EW contributions
- dips = threshold effects
- compare also to HFOLD

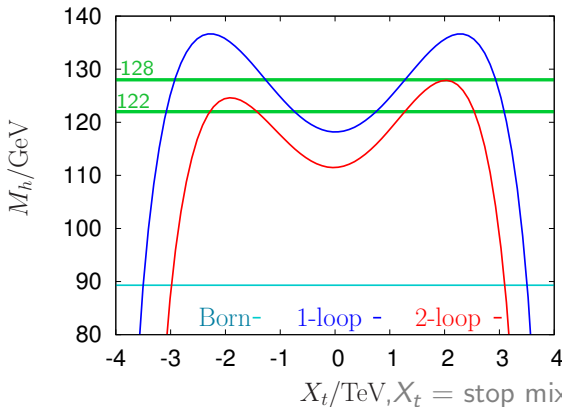
[Eberl, Frisch, Hlucha,
arXiv:1012.5025]

- NMHDECAY (decay part based on HDECAY [Djouadi, Kalinowski, (Mühlleitner,) Spira, hep-ph/9704448])
[Ellwanger, Gunion, Hugonie, hep-ph/0406215]
- NMSSMCALC (decay part based on HDECAY)
[Baglio, Gröber, Mühlleitner, Nhung, H.R., Spira, Streicher, Walz, arXiv:1312.4788]
→ QCD effects and Δb type effects included

Summary

- Corrections to the mass of the charged Higgs boson:
 - * non - if mass chosen as input and fixed as pole mass
 - * otherwise: need to take corrections into account
- For 2HDM, MSSM, NMSSM:
framework ready to do complete EW corrections

Implications of a 125 GeV Higgs boson (MSSM)



generated using FeynHiggs

[Hahn, Heinemeyer, Hollik, H.R.,
Weiglein, Williams]

1-loop [Frank, Hahn, Heinemeyer,
Hollik, H.R., Weiglein]

2-loop $\mathcal{O}(\alpha_{\{t,b\}}\alpha_s, \alpha_{\{t,b\}}^2, \alpha_t\alpha_b)$
[Degrassi, Slavich, Zwirner;

Brignole, Degrassi, Slavich, Zwirner;
Heinemeyer, Hollik, H.R., Weiglein;
Dedes, Degrassi, Slavich]

- A 125 ± 3 GeV mass constrains the parameter space but does not exclude the MSSM. (theory uncertainty ≈ 3 GeV)
- here: no known 3-loop contributions included

[Martin; Harlander, Kant,
Mihaila, Steinhauser]

Higgs-boson self energies in the MSSM

	real parameters				complex parameters	
	$\overline{\text{DR}}$ scheme		OS/mixed schemes		OS/mixed schemes	
	$p^2 = 0$	$p^2 \neq 0$	$p^2 = 0$	$p^2 \neq 0$	$p^2 = 0$	$p^2 \neq 0$
one-loop				✓		✓
two-loop						
$\mathcal{O}(\alpha_t \alpha_s)$	✓	✓	✓	✓	✓	
$\mathcal{O}(\alpha_t^2)$			✓		✓	
$\mathcal{O}(\alpha_b \alpha_s)$			✓			
$\mathcal{O}(\alpha_t \alpha_b, \alpha_b^2, \alpha_\tau^2)$			✓			
$\mathcal{O}(\alpha_\tau \alpha_b)$			✓			
full	✓					
1st 5 rows above		✓				
three-loop						
$\mathcal{O}(\alpha_t \alpha_s^2)$			✓			

Higher orders: “only” logarithmic contributions known

Higgs-boson self energies in the MSSM

- [Chankowski, Pokorski, Rosiek, 92; hep-ph/9303309; Dabelstein hep-ph/9409375]
- [Frank, Hahn, Heinemeyer, Hollik, HR, Weiglein, hep-ph/0611326]
- [Zhang, hep-ph/9808299; Degrassi, Slavich, Zwirner, hep-ph/0105096]
- [Degrassi, Di Vita, Slavich, arXiv:1410.3432]
- [Heinemeyer, Hollik, Weiglein, hep-ph/9812472; Degrassi, Slavich, Zwirner, hep-ph/0105096]
- [Borowka, Hahn, Heinemeyer, Heinrich, Hollik, arXiv:1404.7074, arXiv:1505.03133; Degrassi, Di Vita, Slavich, arXiv:1410.3432]
- [Heinemeyer, Hollik, HR, Weiglein, arXiv:0705.0746]
- [Brignole, Degrassi, Slavich, Zwirner, hep-ph/0112177]
- [Hollik, Paßher, arXiv:1401.8275; arXiv:1409.1687]
- [Brignole, Degrassi, Slavich, Zwirner, hep-ph/0206101; Heinemeyer, Hollik, HR, Weiglein, hep-ph/0411114]
- [Dedes, Degrassi, Slavich, hep-ph/0305127]
- [Allanach, Djouadi, Kneur, Porod, Slavich, hep-ph/0406166]
- [Martin, hep-ph/0211366]
- [Martin, hep-ph/0405022]
- [Harlander, Kant, Mihaila, Steinhauser, arXiv:0803.0672; arXiv:1005.5709]