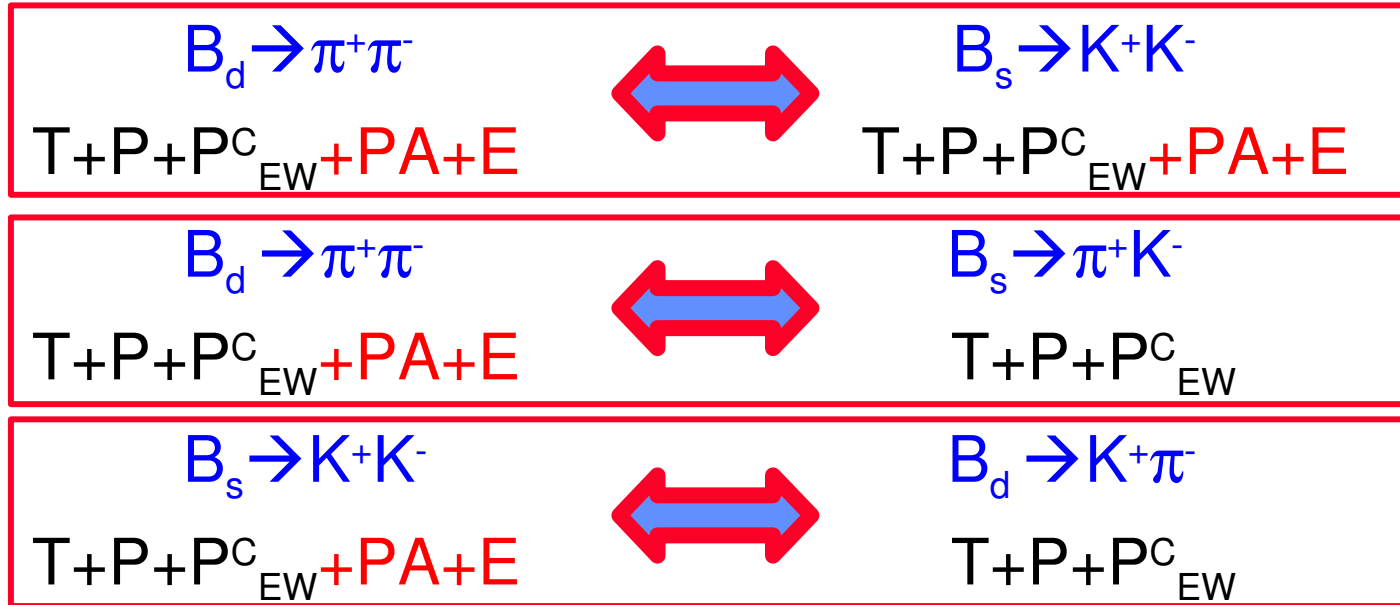


**$B_{s,d} \rightarrow h^+ h^- @ \text{LHCb}$**

A. Sarti  
LNF - INFN

- ◆ Since the original CLEO measurement [ $\text{BR}(B_d \rightarrow K^+\pi^-) \sim 4 \text{BR}(B_d \rightarrow \pi^+\pi^-)$ ], it is well known that **penguin diagrams cannot be neglected** in the two-body B-decay amplitudes
  - This fact complicates considerably the extraction of CKM phases from these decays
  - However, a possibility to eliminate unknown hadronic quantities relies in **exploiting flavour symmetries**, e.g. combining the measurements from U-spin related decay modes
- ◆ The presence of penguins can be viewed as an **opportunity** rather than a **problem**
  - **New Physics contributions might show up inside the loops** of the penguin diagrams, and CKM quantities extracted from these modes can differ from the ones calculated **from tree-level modes**

# U-spin symmetries



T: tree  
P: penguin  
P<sup>C<sub>EW</sub></sup>: colour suppressed electroweak penguin  
PA: penguin annihilation  
E: exchange

- Not all exactly U-spin symmetric, **E** and **PA** contributions missing from flavour specific decays
- **E** and **PA** contributions expected to be relatively small, and can be experimentally probed by measuring the still unobserved  $B_s \rightarrow \pi^+\pi^-$  and  $B_d \rightarrow K^+K^-$  branching ratios  $BR \sim O(10^{-8})$

# B → h<sup>+</sup>h<sup>-</sup> events @ LHCb

Details on LHCb detector can be found in F. Muheim's talk



- Large beauty production cross section expected at 14 TeV p-p collisions: all b-hadron species produced ( $B^\pm, B_d^0, B_s^0, B_c, \Lambda_b \dots \rightarrow 40\%, 40\%, 10\%, 10\%$ )
- B hadrons are produced very likely with a small relative angle and in the very forward (backward) region: **LHCb uses just the forward direction  $1.8 < \eta < 4.9$**   
**@ nominal luminosity ( $L=2 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$ ) about  $10^{12} \text{ } b\bar{b}$  events in  $10^7 \text{ s}$**

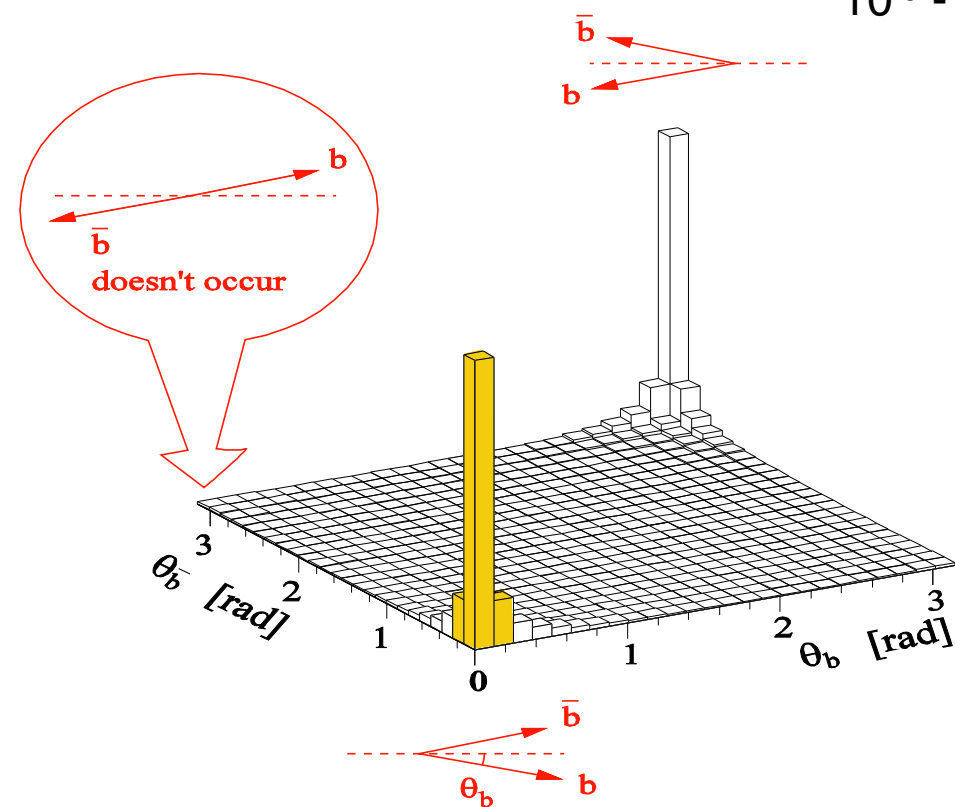
■  $B \rightarrow h^+h^-$  channels have typical branching ratios in the range  $10^{-5} - 10^{-6}$

- Millions of such events produced per year!

$$\sigma_{bb} = 500 \mu\text{b}; \sigma_{\text{inel}} = 80 \text{mb}; \sigma_{bb}/\sigma_{\text{inel}} \sim 6\%$$

- Too many min bias MC events needed to assess S/B: assumption that main bkg comes from b events reduces the events by factor 100

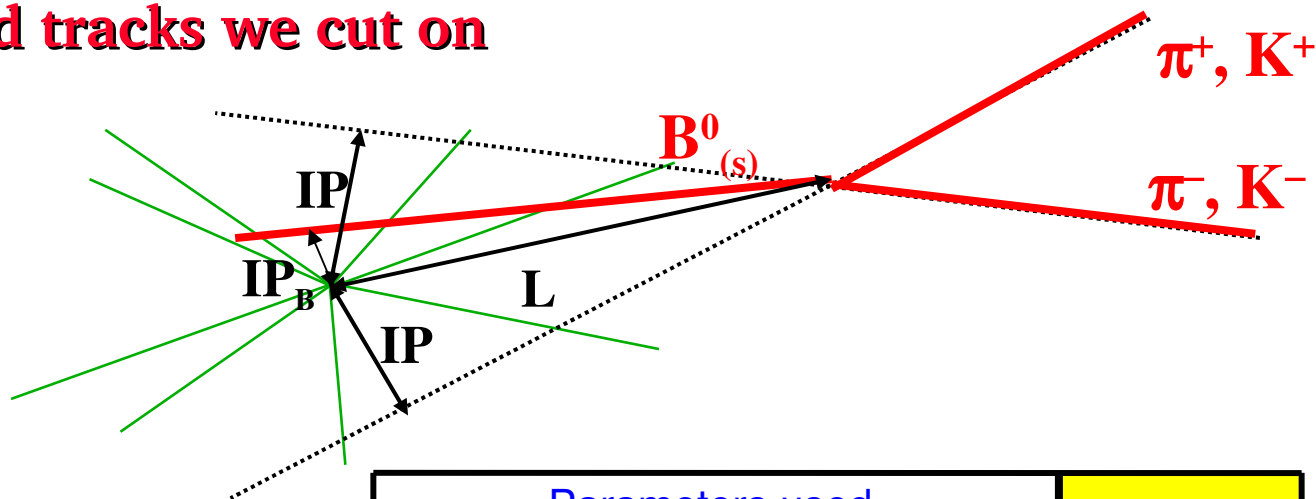
For this study about 300k signal events and 27M inclusive beauty events have been used



# B → h<sup>+</sup>h<sup>'-</sup> selection cuts

## ◆ For each pair of charged tracks we cut on

- max ( $p_{T1}, p_{T2}$ )
- min ( $p_{T1}, p_{T2}$ )
- max ( $IP1/\sigma_{IP1}, IP2/\sigma_{IP2}$ )
- min ( $IP1/\sigma_{IP1}, IP2/\sigma_{IP2}$ )
- $\chi^2$  of common vertex



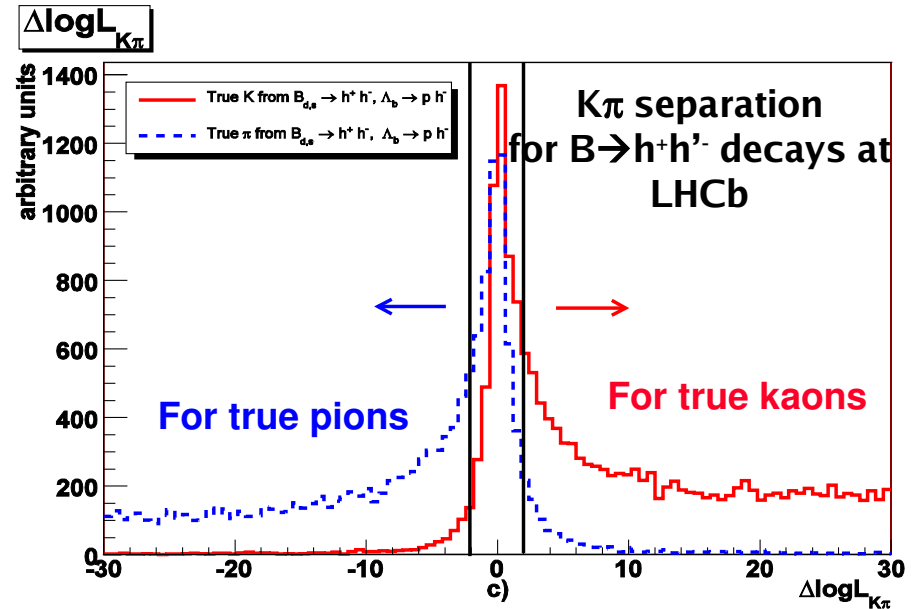
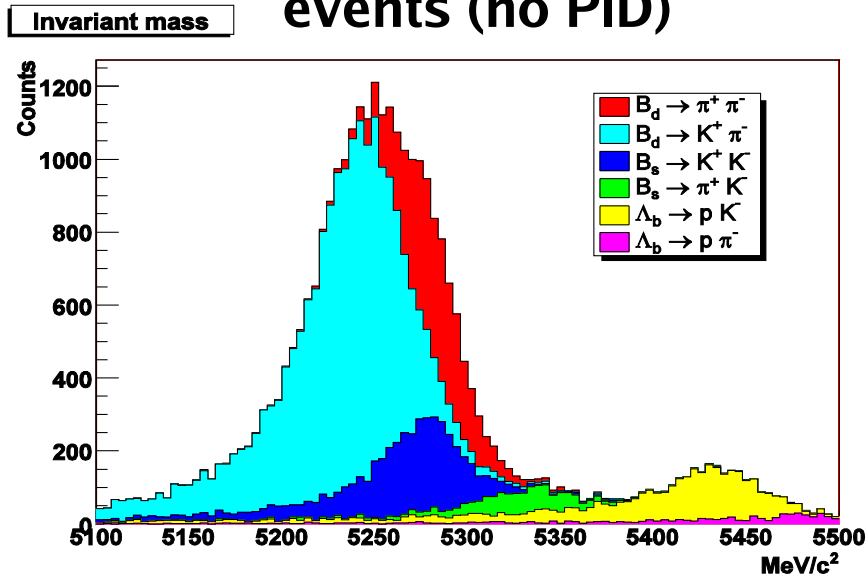
## ◆ Then, the B candidate is selected with cuts on

- ◆  $p_T$
- ◆  $IP/\sigma_{IP}$
- ◆  $L/\sigma_L$

Parameters used for selection			
h,h	Smallest $p_t$	GeV/c	> 1
h,h	largest $p_t$	GeV/c	> 3
h,h	smallest $IP/\sigma_{IP}$		> 6
h,h	largest $IP/\sigma_{IP}$		> 12
h, h	vertex fit $\chi^2$		< 5
B	$p_t$	GeV/c	> 1
B	$IP/\sigma_{IP}$		< 2.5
B	$L/\sigma_L$		> 18

- ◆ PID discriminants built as difference between log likelihoods of particle hypotheses

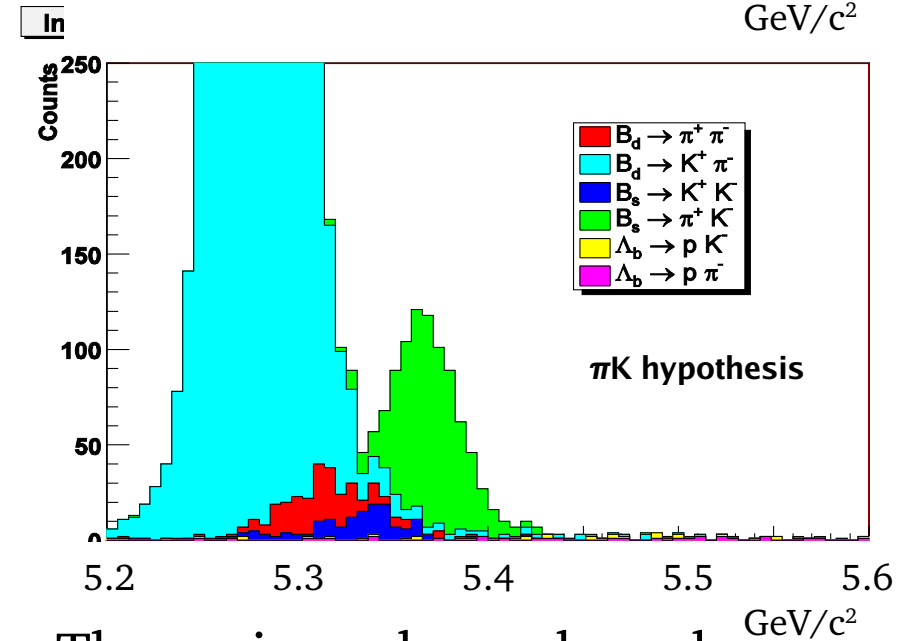
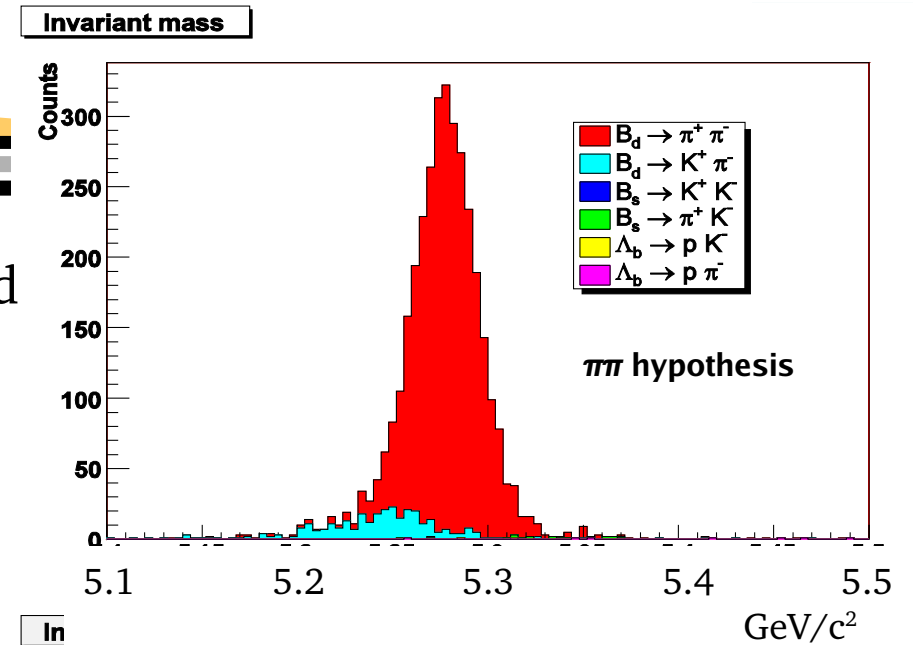
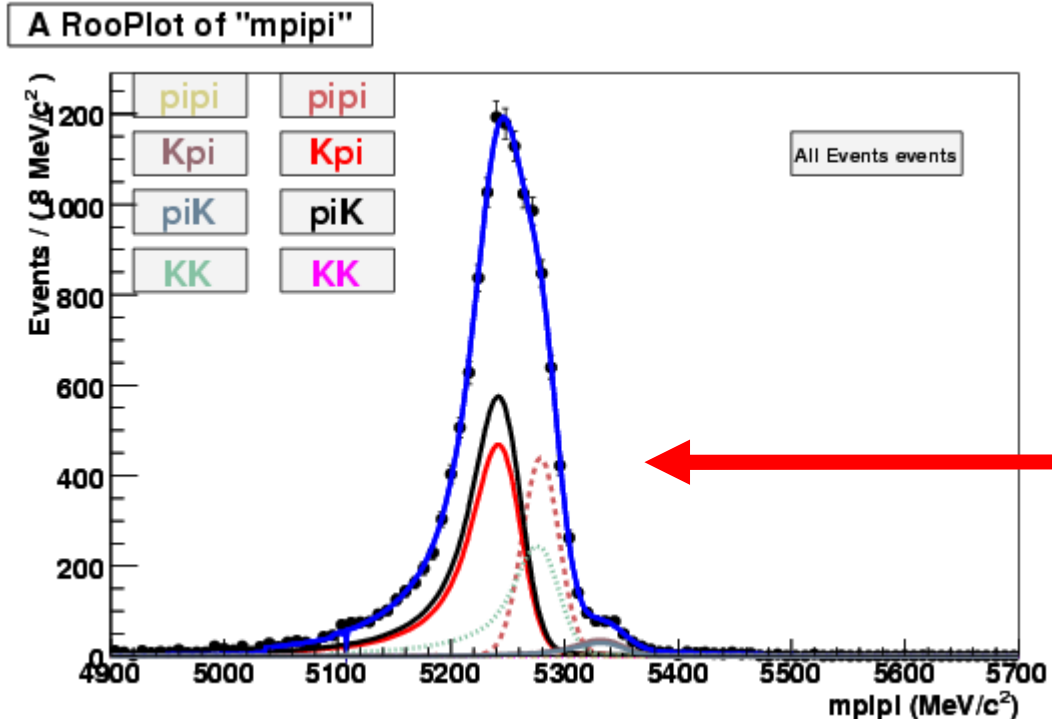
## Invariant mass of selected events (no PID)



- ◆ Efficient hadron PID crucial for the  $B \rightarrow h^+ h^-$  channels
  - Invariant masses one on top of the other  $\rightarrow$  kinematic separation possible but weak
- ◆ Calibration of K/ $\pi$  PID on data will be performed using the  $D^{*+} \rightarrow D^0 \pi^+$ ,  $D^0 \rightarrow K^+ \pi^-$  decay chain
  - Large "PID unbiased" samples of such  $D^*$  decays will be acquired with dedicated  $D^*$  trigger chain (300 Hz bandwidth)

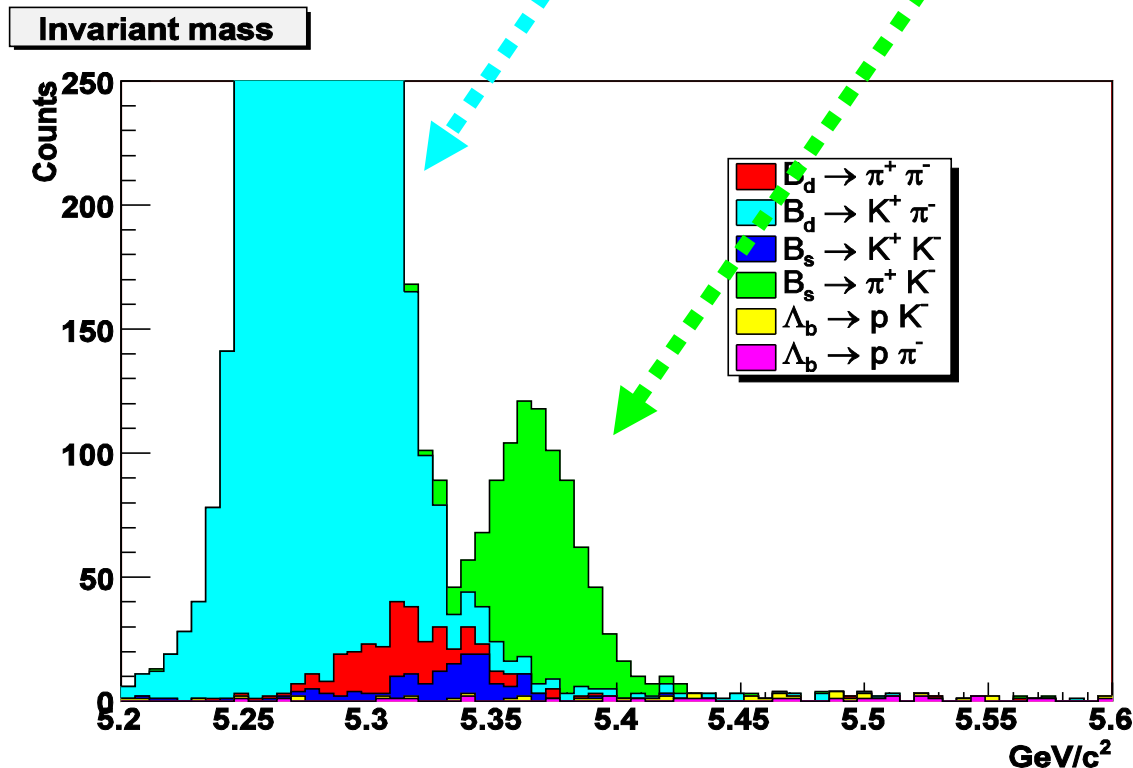
# Using PID information

- ◆ Every  $h^+h^-$  channel is potentially a background for the other channels:
  - The impressive performance of RICH systems allows to select very clean samples cutting on PID information
  - To ease the evaluation of systematics a simultaneous Likelihood fit has been used



The various channels and background are separated in the fit by means of the particle ID observables and invariant mass

- ◆ Discrimination of  $B_d \rightarrow K^- \pi^+$  from  $B_s \rightarrow \pi^+ K^-$  **must rely on mass resolution only** since they share exactly the same signature
  - $B_d \rightarrow K^- \pi^+$  about 16 times more abundant (4 times larger branching fraction, 4 times larger hadronization fraction)



Nominal mass resolution of  $\sigma=16 \text{ MeV}/c^2$  allows for a clean separation between the two signals



# Proper time resolution

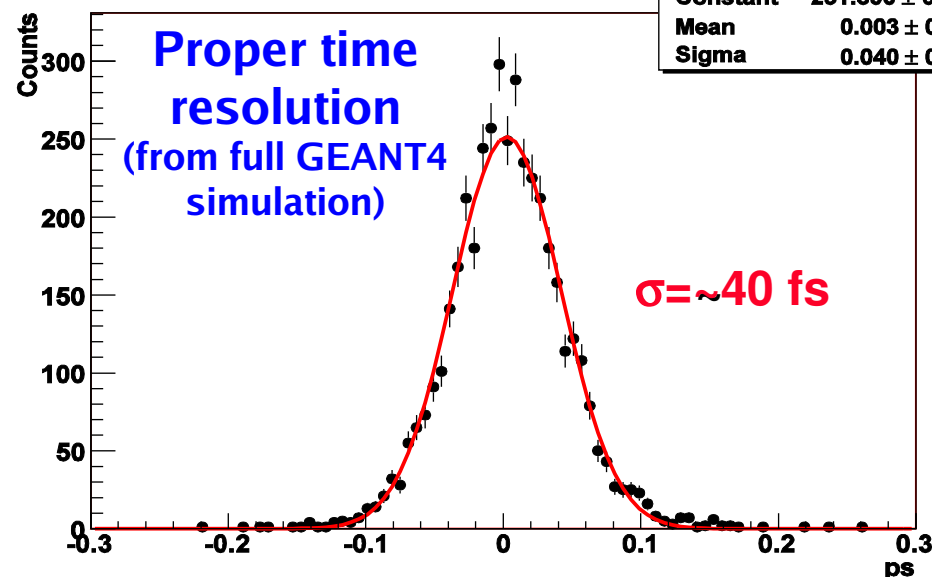
- Excellent proper time resolution allows to perfectly resolve the fast  $B_s$  oscillations

- $\sigma_t = \sim 40$  fs (10 times less than the oscillation period actually)

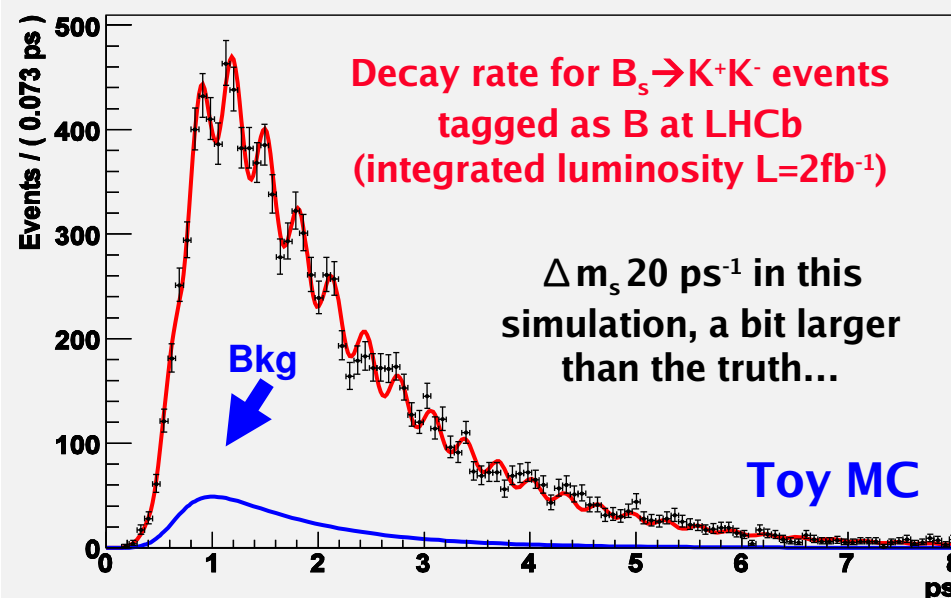
- Calibration of proper time resolution on data is crucial!

- will be done in LHCb by disentangling prompt  $J/\psi \rightarrow \mu^+ \mu^-$  from non-prompt component from B decays
  - The prompt component is affected just by the intrinsic resolution of the detector, hence can be used for calibration
  - “Lifetime unbiased” prompt  $J/\psi$  events will be acquired with dedicated high-mass di-muon trigger (600 Hz bandwidth)

$B_d \rightarrow K^+ \pi^-$  Proper time resolution



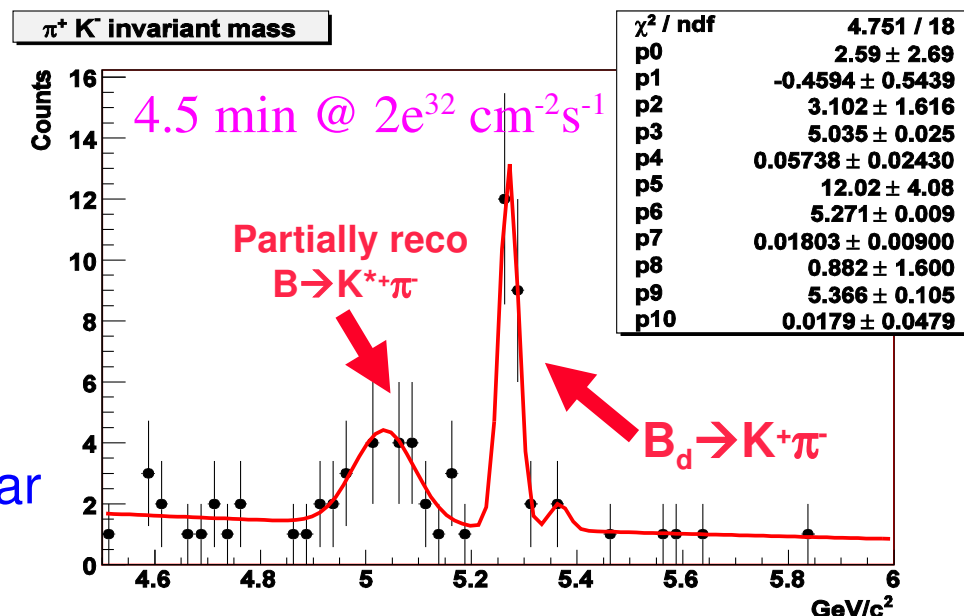
$B_s \rightarrow K^+ K^-$  proper time



# Selection performance

- ◆ Yields are calculated for an integrated luminosity of  $2 \text{ fb}^{-1}$ 
  - $10^7$  seconds at nominal LHCb luminosity
- ◆ Limited background-to-signal ratios
  - Both combinatorial and from the other  $B \rightarrow h^+ h'^-$  modes (due to wrong particle ID)
- ◆ Hundreds of thousands of  $B \rightarrow h^+ h'^-$  decays triggered, reconstructed and selected per year of running with high degree of purity!

$\pi K$  spectrum selected from a sample of 27M  $b\bar{b}$  MC events

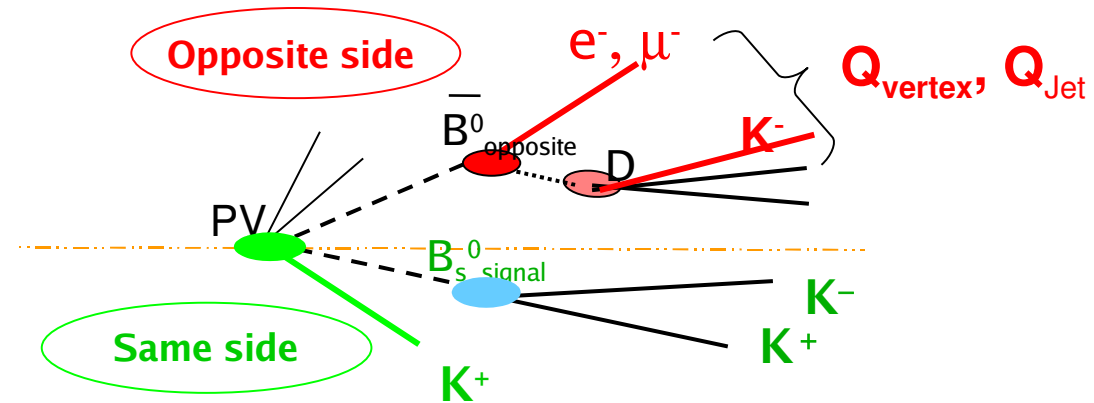


	$\epsilon_{\text{tot}} = \epsilon_{\text{gen}} \times \epsilon_{\text{sel/gen}} \times \epsilon_{\text{trg/sel}}$				BR $10^{-6}$	Yield	B/S bb	B/S spec
	$\epsilon_{\text{gen}}$	$\epsilon_{\text{sel/gen}}$	$\epsilon_{\text{trg/sel}}$	$\epsilon_{\text{tot}}$				
$B_d^0 \rightarrow \pi^+ \pi^-$	$34.9 \pm 0.3$	$7.36 \pm 0.13$	$36.3 \pm 1.1$	$0.93 \pm 0.03$	4.8	35700	0.46	0.08
$B_d^0 \rightarrow K^+ \pi^-$	$34.9 \pm 0.3$	$7.21 \pm 0.07$	$36.8 \pm 0.6$	$0.93 \pm 0.02$	18.5	137600	0.14	0.02
$B_s^0 \rightarrow \pi^+ K^-$	$34.8 \pm 0.3$	$7.25 \pm 0.30$	$40.6 \pm 2.3$	$1.02 \pm 0.06$	4.8	9800	1.92	0.54
$B_s^0 \rightarrow K^+ K^-$	$34.8 \pm 0.3$	$7.11 \pm 0.13$	$39.3 \pm 1.2$	$0.97 \pm 0.03$	18.5	35900	< 0.06	0.08
$\Lambda_b \rightarrow p \pi^-$	$34.6 \pm 0.3$	$7.20 \pm 0.26$	$38.3 \pm 2.3$	$0.95 \pm 0.06$	4.8	9100	1.66	0.11
$\Lambda_b \rightarrow p K^-$	$34.6 \pm 0.3$	$6.80 \pm 0.13$	$36.0 \pm 1.1$	$0.86 \pm 0.03$	18.5	31800	< 0.08	0.02

# Flavour Tagging

## Opposite side

- ◆ High  $p_T$  leptons
- ◆  $K^\pm$  from  $b \rightarrow c \rightarrow s$
- ◆ Vertex charge
- ◆ Jet charge



## Same side

- ◆ Fragmentation  $K^\pm$  accompanying  $B_s$
- ◆  $\pi^\pm$  from  $B^{**} \rightarrow B^{(*)}\pi^\pm$

$\epsilon D^2 = \epsilon(1-2\omega)^2$ : tagging power

$\epsilon$ : tagging efficiency

$\omega$ : mistag probability

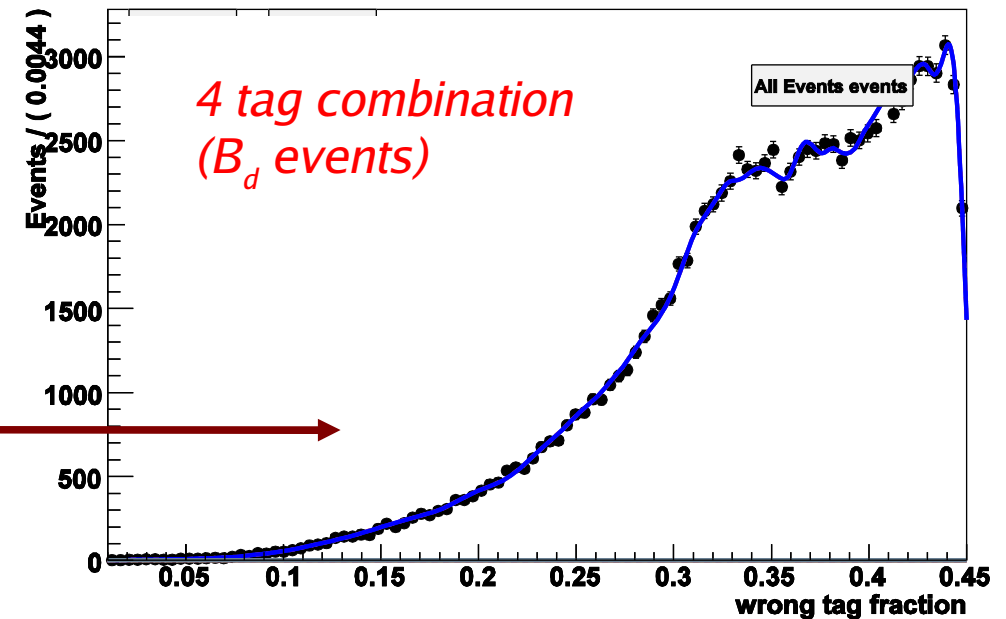
## Tagging power in %

Tag	$B_d$	$B_s$
Muon	1.1	1.5
Electron	0.39	0.69
Kaon opp.side	2.1	2.3
Jet/ Vertex Charge	1.0	0.97
Same side $\pi / K$	0.73 ( $\pi$ )	3.5 (K)
Combined (NNET)	5.05	9.5

Work in continuous progress.

- Strategy I:
  - Define N different tagging categories and use various fixed  $\omega_N$
- Strategy II:
  - Extract from kinematic variables the per event  $\omega$  :  $P(\omega) = P(\omega | p, \dots)P(p)$

A RooPlot of "wrong tag fraction"



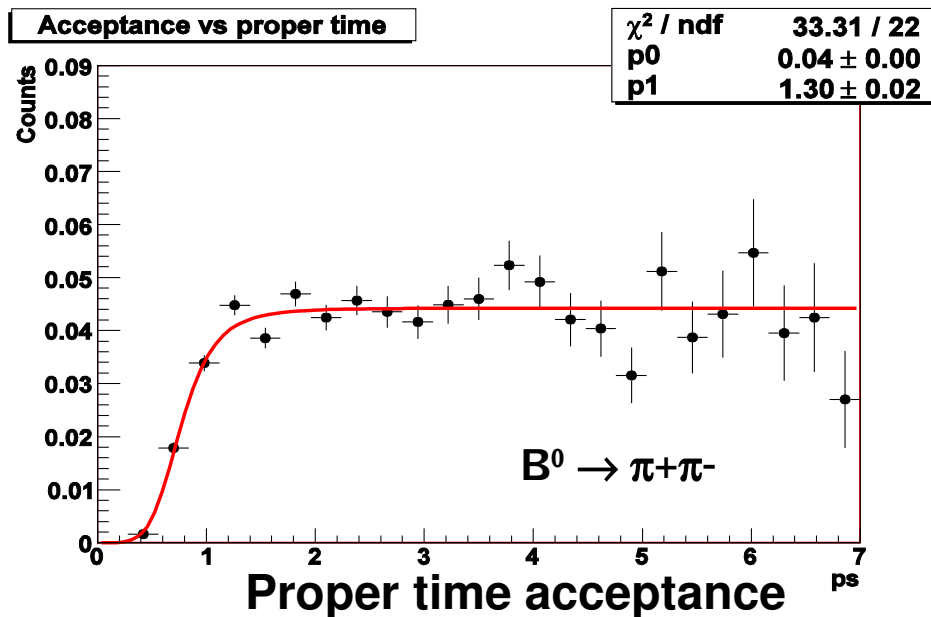
In both cases, control samples are needed in order to evaluate  $\omega$

$\omega_{(N)}$  can be extracted from data using  $B^+$  (ex.  $B^+ \rightarrow J/\psi K^+$ ,  $B^+ \rightarrow D^0 \mu \nu$ ) or  $B^0$  flavour-specific ( $B^0 \rightarrow K^+ \pi^-$ ,  $B^0 \rightarrow J/\psi K^{*0}$ ,  $B^0 \rightarrow D^{*-} \mu^+ \nu$ ) decays

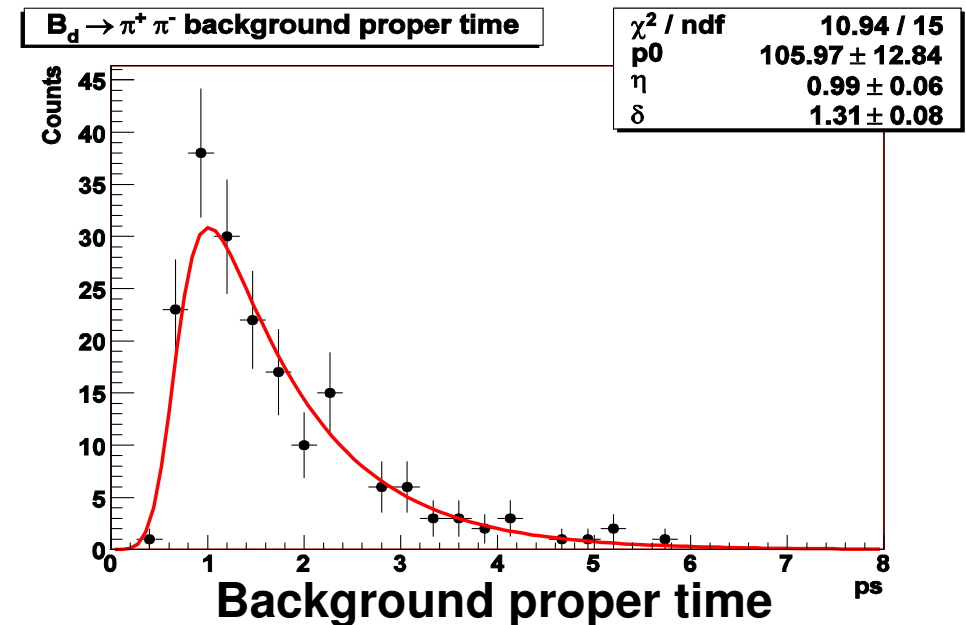
For the  $K_{SS}$ ,  $BR(B_s \rightarrow \pi K)$  is too low.  $\omega_{SS}$  can be fixed to the value extracted from  $B_s \rightarrow D_s \pi$  or  $B_s \rightarrow D_s \mu \nu$  events.

# Toy MC generation

- A toy MC is used in order to generate large samples of signal and background events for estimating sensitivities on CP parameters from time dependent decay rate fits
- Signal samples generated simulating CP violation, proper time acceptance, proper time resolution and tagging
  - For each event the toy MC generates mass, proper time, tagging response, PID response according to the results from the full GEANT4 simulations
- Background mass and proper time spectra generated according to result from the full simulation



will be calibrated on real data using event by event acceptance functions – studies going on in LHCb



will be well known from real data using mass sidebands

# Sensitivity on CP asymmetries

- CP sensitivities are then estimated with an extended unbinned maximum likelihood fit to the toy MC sample
- The likelihood fit is performed simultaneously to all the  $B \rightarrow hh'$  channels, including tagged and untagged samples
- C and S coefficients for  $B_d \rightarrow \pi^+\pi^-$  and  $B_s \rightarrow K^+K^-$  events tagged as B ( $q=+1$ ) or  $\bar{B}$  ( $q=-1$ ) are given by

$$\lambda_f = \frac{q}{p} \frac{\bar{A}_f}{A_f}; \quad f = \pi^+\pi^-, K^+K^-; \quad C_f = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2}; \quad S_f = \frac{2\text{Im}\lambda_f}{1 + |\lambda_f|^2}$$

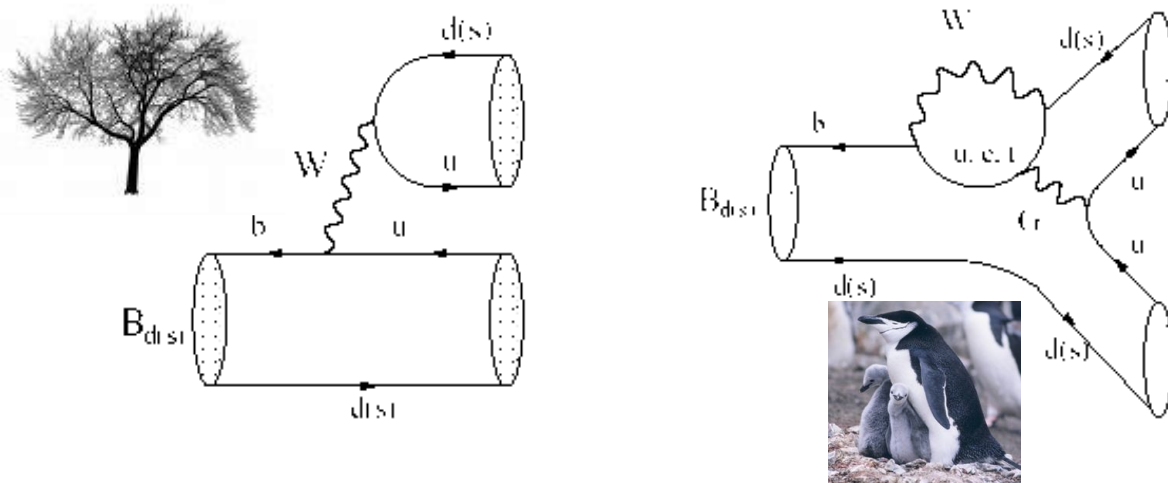
LHCb can reach a sensitivity, on  $B_d$  modes, about 3 times better than the current world average in  $10^7$  s of running at nominal luminosity

Fit results corresponding to Int.  $L=2\text{fb}^{-1}$   
( $10^7$  seconds at nominal LHCb luminosity)

	$B_d \rightarrow \pi^+\pi^-$	$B_s \rightarrow K^+K^-$
$\sigma(C)$	0.043 (0.10*)	0.042
$\sigma(S)$	0.037 (0.12*)	0.044
	$B_d \rightarrow K^+\pi^-$	$B_s \rightarrow \pi^+K^-$
$\sigma(A_{CP})$	0.003 (0.015*)	0.02

\*Current world average

# Trees, penguins and gamma



Sensitivity to  $\gamma$  from the **interference of T and P** amplitudes

- $d, d'$ : penguin-to-tree ratios
- $\theta, \theta'$ : penguin-tree strong phase differences

$$A(B^0 \rightarrow \pi^+ \pi^-) = K(e^{i\gamma} - de^{i\theta})$$

$$A(B_s^0 \rightarrow K^+ K^-) = \frac{\lambda}{1 - \lambda^2/2} K' \left( e^{i\gamma} + \frac{1 - \lambda^2}{\lambda^2} d' e^{i\theta'} \right)$$

← Sensitivity to  $\gamma$  doubly Cabibbo suppressed in this mode ☹

Using U-spin symmetry one gets  $d=d'$  and  $\theta=\theta'$

Method and parametrization from  
**R. Fleischer, PLB 459 (1999) 306**

# Extraction of gamma

$$C(B_d^0 \rightarrow \pi^+ \pi^-) = f_1(d, \theta, \gamma)$$

$$S(B_d^0 \rightarrow \pi^+ \pi^-) = f_2(d, \theta, \gamma, \phi_d)$$

$$C(B_s^0 \rightarrow K^+ K^-) = f_3(d', \theta', \gamma)$$

$$S(B_s^0 \rightarrow K^+ K^-) = f_4(d', \theta', \gamma, \phi_s)$$

$$A_{CP}^{th}(\tau) = \frac{C \cdot \cos(\Delta M \cdot \tau) - S \cdot \sin(\Delta M \cdot \tau)}{\cosh(\frac{\Delta\Gamma}{2} \cdot \tau) - A_{\Delta\Gamma} \cdot \sinh(\frac{\Delta\Gamma}{2} \cdot \tau)}$$

Once the direct and mixing-induced CP-violating terms are measured, one has a system of **7 unknowns and 4 equations**

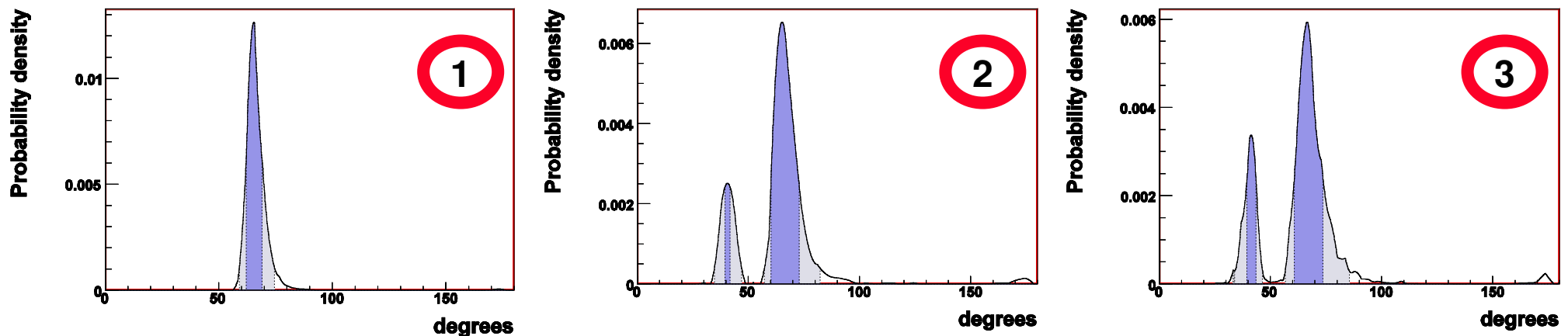
However, the mixing phase  $\phi_d$  ( $\phi_s$ ) is (will be) measured from  $B_d \rightarrow J/\psi K_S$   
( $B_s \rightarrow J/\psi \phi$ )  $\rightarrow$  **5 unknowns**

Relying on U-spin symmetry one eliminates two further unknowns  $\rightarrow$  **3 unknowns, system over-constrained,  $\gamma$  can be extracted unambiguously**



# Sensitivity on gamma ( $L=2\text{fb}^{-1}$ )

- Extraction of  $\gamma$  performed with a Bayesian approach in 3 different scenarios (in this exercise  $\gamma=65^\circ$  is assumed)
  - (1) Perfect U-spin symmetry
    - i.e. using the constraints  $d=d'$  ;  $\theta = \theta'$
  - (2) Weaker assumption: perfect U-spin symmetry just for  $d, d'$ 
    - $d=d'$  ; no constraint on  $\theta, \theta'$
  - (3) Even weaker: given U-spin breaking for  $d, d'$ 
    - $\xi = d'/d = [0.8, 1.2]$  ; no constraint on  $\theta, \theta'$



Sensitivity ranging from  $\sim 4^\circ$  in case of perfect U-spin assumed, to  $7^\circ$ - $10^\circ$  partially releasing U-spin assumptions (but a secondary fake solution appears)

# Exercise: using world averages

(not using LHCb results)

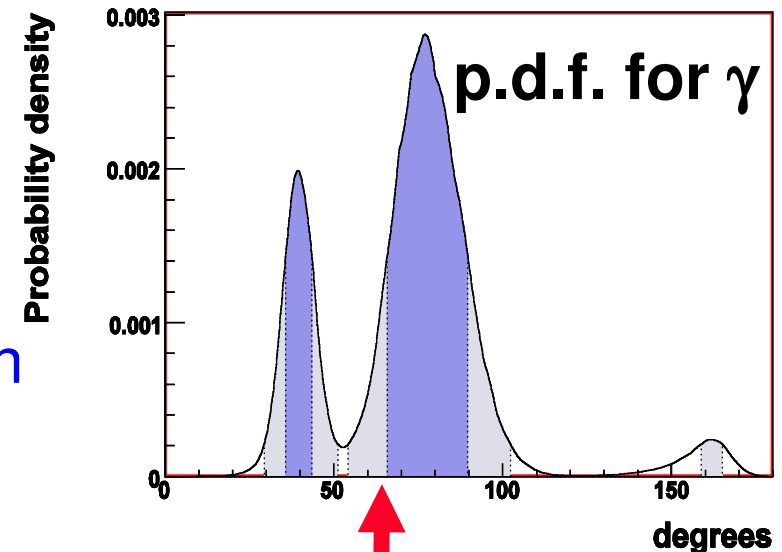
- Instead of using the  $B_s \rightarrow K^+ K^-$  one can make use of the  $B_d \rightarrow K^+ \pi^-$ 
  - It just differs for the spectator quark, but neglecting penguin annihilation and exchange diagrams which are not present in  $B_d \rightarrow K^+ \pi^-$
  - Assuming U-spin symmetry the cosine term of the time dependent asymmetry of  $B_s \rightarrow K^+ K^-$  is equal to the charge asymmetry of  $B_d \rightarrow K^+ \pi^-$ , thus we can replace  $C_{KK}$  with  $A_{K\pi}$
  - One observable is missing ( $S_{KK}$ ), but still can solve system of constraints for  $\gamma$

## ◆ Current world averages from HFAG

- $C_{\pi\pi} = -0.37 \pm 0.10$  (BaBar/Belle)
- $S_{\pi\pi} = -0.50 \pm 0.12$  (BaBar/Belle)
- $A_{K\pi} = -0.093 \pm 0.015$  (BaBar/Belle/CLEO/CDF)

## ◆ p.d.f. for $\gamma$ obtained allowing for a U-spin breaking

- $\Delta\theta = \theta' - \theta = \pm 20^\circ$
- $\xi = d'/d = [0.8, 1.2]$



Unitarity Triangle fits expectation

- LHCb will collect large samples of  $B \rightarrow h^+h^-$  decays
  - Its excellent vertexing and PID capabilities will allow to **collect several hundreds of thousands of  $B \rightarrow h^+h^-$  events with very high purity**
- The CP sensitivity reachable on  $B_d$  modes ( $10^7$  s of data taking,  $L = 2 \cdot 10^{32} \text{ cm}^{-2}\text{s}^{-1}$ ) is  $\sim 3$  times better than the current world averages (B factories, Tevatron)
- Besides the general interest of measuring CP violation for these channels, especially for the still unmeasured  $B_s$  decays, the  $B \rightarrow h^+h^-$  modes can provide useful information to constrain the CKM angle
  - By relating the  $B_d$  and  $B_s$  decay modes by means of the U-spin flavour symmetry (even allowing a symmetry breaking to some extent it is still possible to extract useful information)
- The extraction of  $\gamma$  from these decays can reveal contributions from NP
  - $\gamma$  obtained from these decay modes might thus be different from  $\gamma$  from tree level modes, such as  $B \rightarrow DK$