
Neutralino spin measurement with ATLAS



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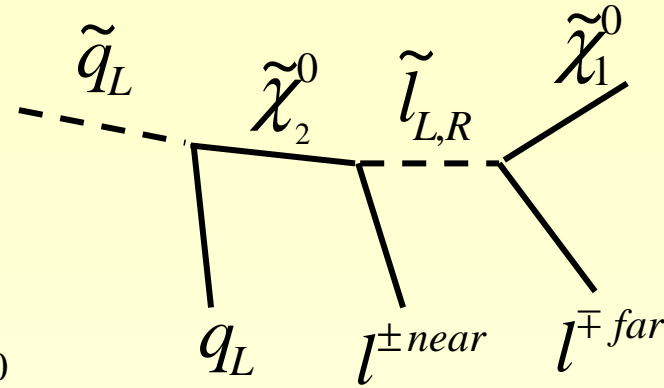


Flavour in the era of the LHC , 4th meeting – CERN, 9 Oct 2006

Contents

- Spin measurement
- SU1 and SU3 points: kinematics
- Event selection
- Study of background
- Charge asymmetries
- Diluting effects on asymmetry
- Summary and conclusions

Spin measurement



■ Left squark cascade decay

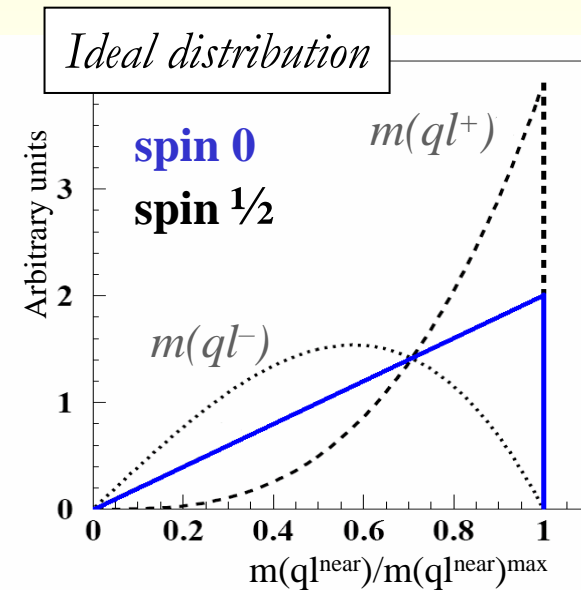
$$\tilde{q}_L \rightarrow q \tilde{\chi}_2^0 \rightarrow q \tilde{l}_{L,R}^{\mp} l^{\pm} \rightarrow ql^{\pm} l^{\mp} \tilde{\chi}_1^0$$

Only electrons and muons considered, from $\tilde{l} = \tilde{e}, \tilde{\mu}$

If neutralino spin is $1/2$, angular distribution of slepton is not spherically symmetric, therefore invariant mass $m(q\tilde{l}^{\pm})$ is charge asymmetric:

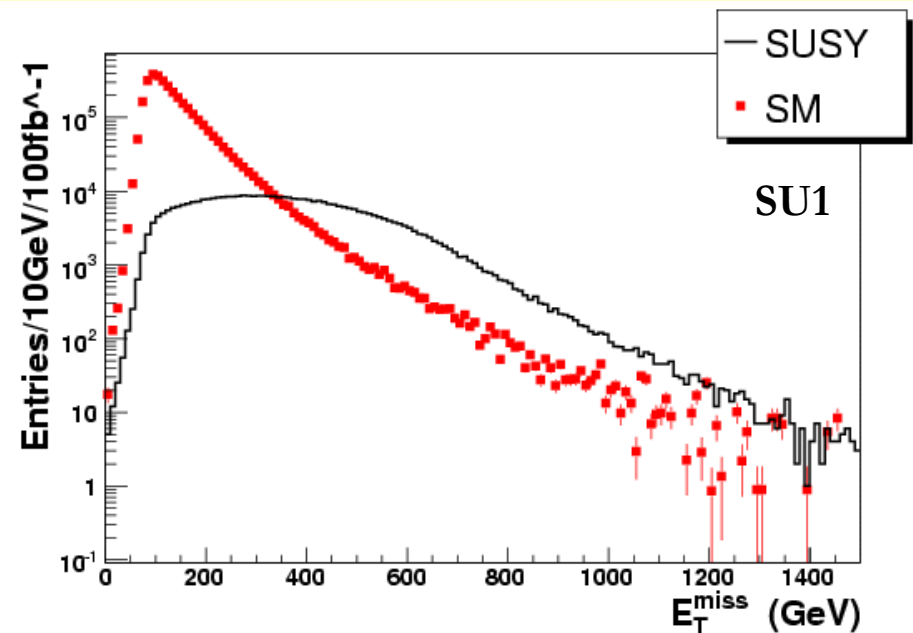
$$A_{ql} = \frac{s^+ - s^-}{s^+ + s^-} \quad s^{\pm} = \frac{d\sigma}{d(m_{ql^{\pm}})}$$

- Studies for **LHCC5** point by A. J. Barr
Phys.Lett.B596 (2004) 205



Event topology

- In **mSUGRA** events, strongly interacting sparticles (\tilde{q} , \tilde{g}) dominate LHC production
- Cascade decays to the stable, weakly interacting lightest neutralino (**L**ightest **S**upersymmetric **P**article) follow, *e.g.* left squark cascade decay.
- **Typical final state signature:**
 - **Large E_T^{miss}** (from LSP)
 - **High p_T jets** (from squark/gluino decay)
 - **Two same-flavour opposite-sign (SFOS) leptons**
- Charge asymmetry $A_{q\ell}$ is suppressed by \tilde{q}/\tilde{q} cancelation (anyway at the LHC much more squarks than antisquarks are expected)
- Asymmetry can be also diluted by the experimental impossibility to distinguish **near** from **far** lepton



mSUGRA points considered

- **Point SU1** – stau coannihilation point

$$m_0 = 70 \text{ GeV} , m_{1/2} = 350 \text{ GeV} , A_0 = 0 \text{ GeV}$$

$$\tan(\beta) = 10 , \text{sign}(\mu) = +$$

$$\sigma_{\text{LO}} = 7.8 \text{ pb} \quad (100 \text{ fb}^{-1} \text{ analysed sample})$$

- **Point SU3** – bulk region point

$$m_0 = 100 \text{ GeV} , m_{1/2} = 300 \text{ GeV} , A_0 = -300 \text{ GeV}$$

$$\tan(\beta) = 6 , \text{sign}(\mu) = +$$

$$\sigma_{\text{LO}} = 19.5 \text{ pb} \quad (30 \text{ fb}^{-1} \text{ analysed sample})$$

ATLAS official Fast Simulation samples + dedicated production in Napoli ATLAS-Tier2

SU1 point kinematics

- Both $\tilde{\chi}_2^0 \rightarrow \tilde{l}_L^\pm l^\mp$ and $\tilde{\chi}_2^0 \rightarrow \tilde{l}_R^\pm l^\mp$ decays are allowed in SU1

(masses are: $m_{\tilde{\chi}_2^0} = 264 \text{ GeV}$, $m_{\tilde{l}_L} = 255 \text{ GeV}$, $m_{\tilde{l}_R} = 154 \text{ GeV}$)

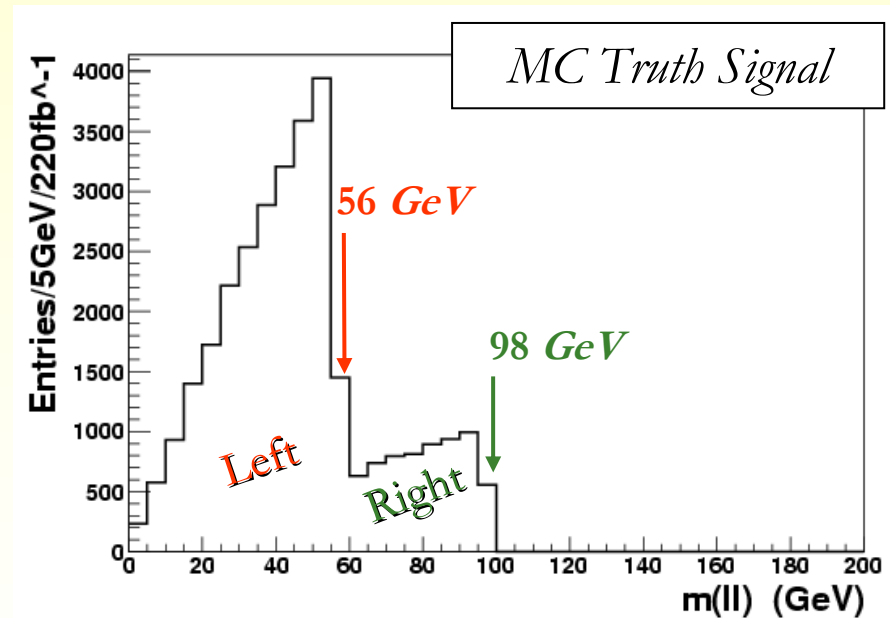
$$\tilde{q}_L \rightarrow q \underset{264}{\tilde{\chi}_2^0} \rightarrow q \underset{255}{\tilde{l}_L} \underbrace{l^\pm}_{\text{soft}} \rightarrow q l^\pm \underbrace{l^\mp}_{\text{hard}} + \underset{137}{\tilde{\chi}_1^0}$$

$$\tilde{q}_L \rightarrow q \underset{264}{\tilde{\chi}_2^0} \rightarrow q \underset{154}{\tilde{l}_R} \underbrace{l^\pm}_{\text{hard}} \rightarrow q l^\pm \underbrace{l^\mp}_{\text{soft}} + \underset{137}{\tilde{\chi}_1^0}$$

- Two edges in the dilepton mass (for Left/Right slepton chains)

$$m(ll)_{\max} = \left[\frac{(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_{L,R}}^2)(m_{\tilde{l}_{L,R}}^2 - m_{\tilde{\chi}_1^0}^2)}{m_{\tilde{l}_{L,R}}^2} \right]^{1/2}$$

- After **Left/Right** slepton decay chain identification, *near* and *far* leptons are distinguishable with high probability from their momenta.

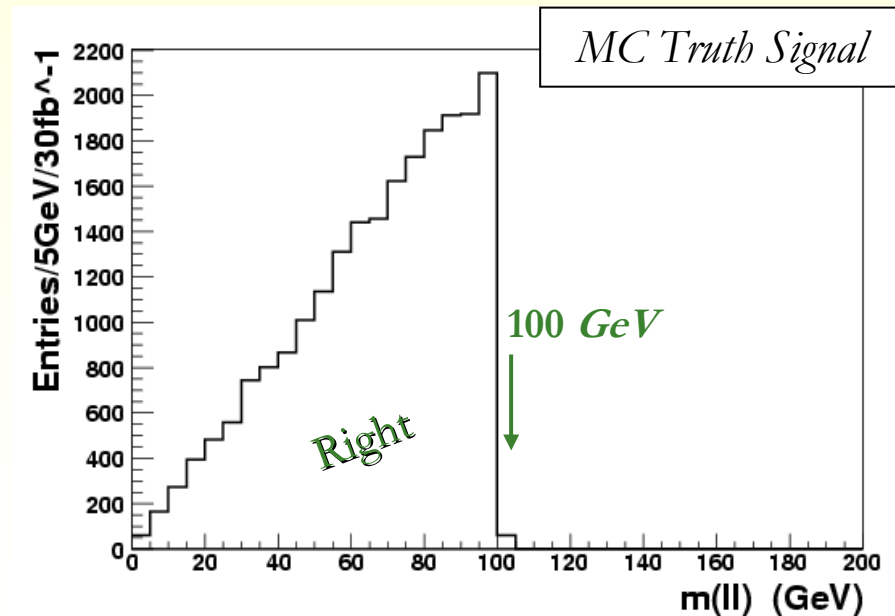


SU3 point kinematics

- Only the right slepton decay $\tilde{\chi}_2^0 \rightarrow \tilde{l}_R^\pm l^\mp$ is allowed in SU3
(masses are: $m_{\tilde{\chi}_2^0} = 219 \text{ GeV}$, $m_{\tilde{l}_L} = 230 \text{ GeV}$, $m_{\tilde{l}_R} = 155 \text{ GeV}$)

$$\tilde{q}_L \rightarrow q \underset{219}{\tilde{\chi}_2^0} \rightarrow q \underset{155}{\tilde{l}_R} (l^\pm) \rightarrow ql^\pm (l^\mp) + \underset{118}{\tilde{\chi}_1^0}$$

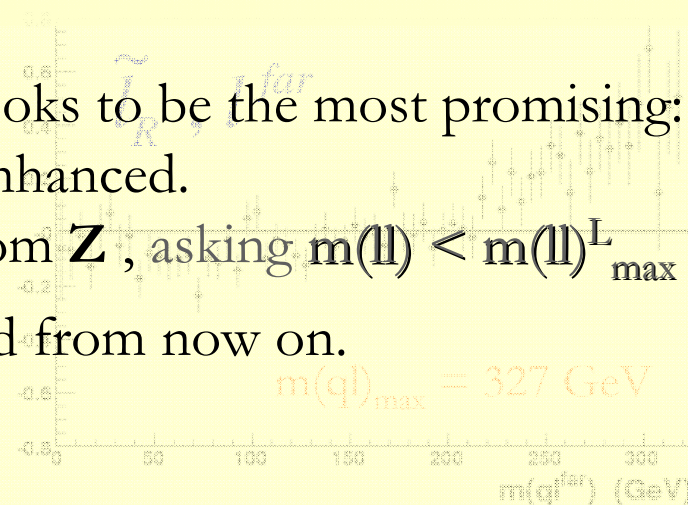
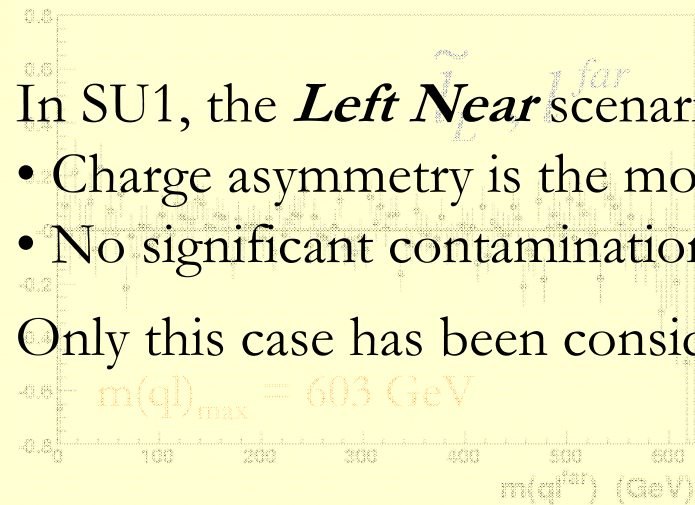
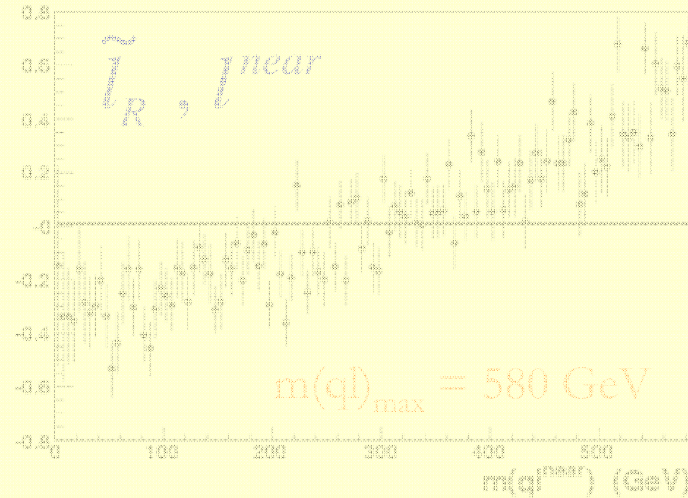
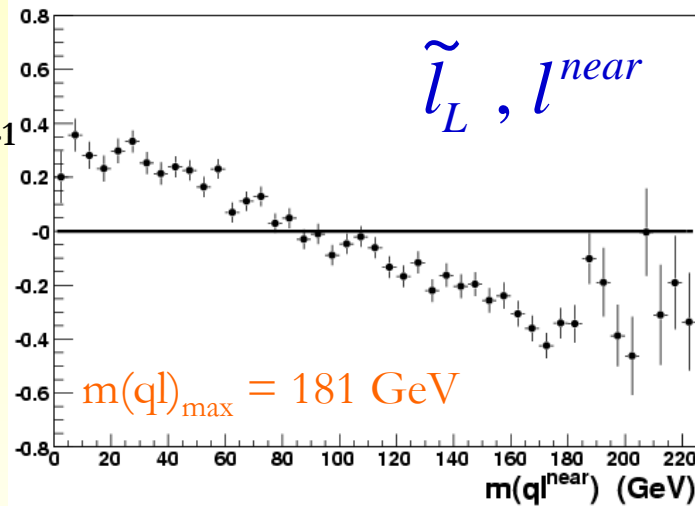
- One edge in the dilepton mass (for **R**ight slepton chain).
- Here *near* and *far* leptons are undistinguishable.



$m(q\ell^\pm)$ asymmetry for MC truth signal

SU1

L=220 fb⁻¹
generated
sample



In SU1, the **Left Near** scenario looks to be the most promising:

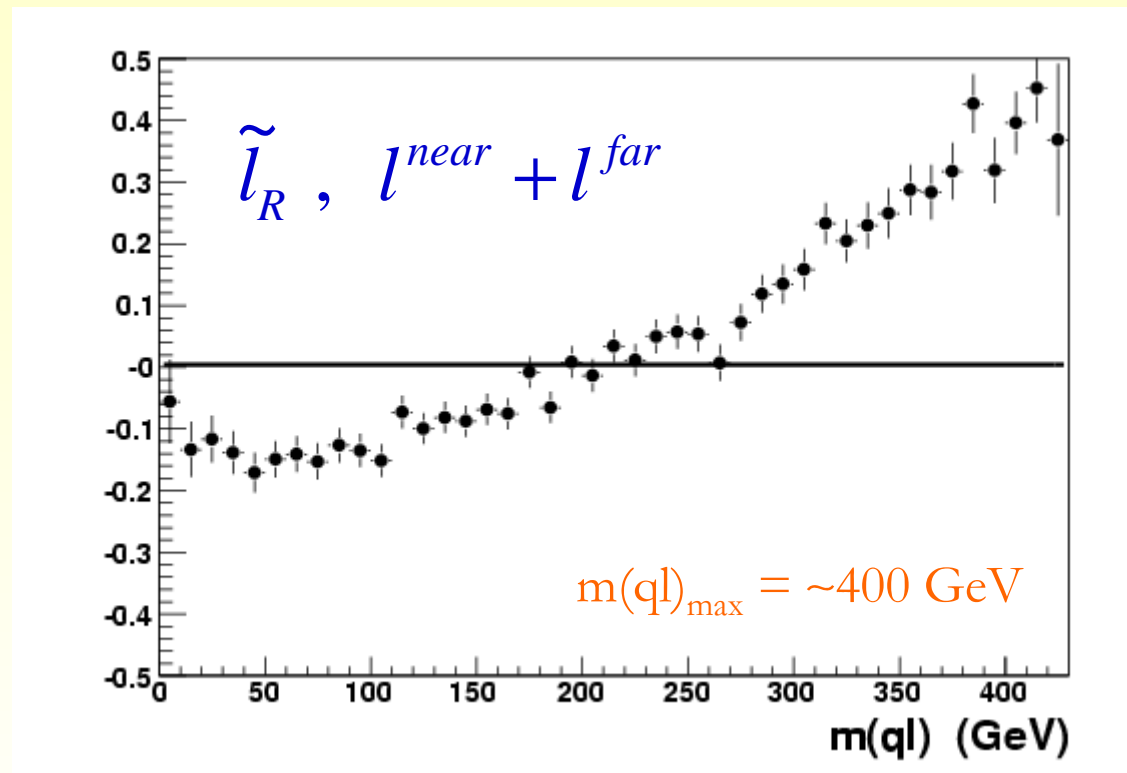
- Charge asymmetry is the most enhanced.
- No significant contamination from **Z**, asking $m(\ell\ell) < m(\ell\ell)_{\max}^L$

Only this case has been considered from now on.

$m(\text{ql}^\pm)$ asymmetry for MC truth signal

SU3

$L=30 \text{ fb}^{-1}$
generated
sample



Even if *near* and *far* leptons are not distinguishable, in SU3 point charge asymmetry is more pronounced than in SU1, owing to the larger cross section and to the higher branching ratios of decays involved in the left squark cascade decay.

Event selection

- $E_T^{\text{miss}} > 100 \text{ GeV}$
- At least 4 jets
 - $p_T(j_1) > 100 \text{ GeV}$,
 - $p_T(j_i) > 50 \text{ GeV}$, $i=2,3,4$
- Two SFOS leptons
 - $p_T(l) > 6 \text{ GeV}$ (for $SU1$), 10 GeV (for $SU3$)
 - $|\eta| < 2.5$, $E_T^{\text{isol}} < 10 \text{ GeV}$ in $\Delta R < 0.2$
- $m(\text{ll}) < 100 \text{ GeV}$
- $m(\text{jll}) < 615 \text{ GeV}$ (for $SU1$)
 $< 500 \text{ GeV}$ (for $SU3$)
- In $SU1$, the decay chain with **Left/Right** slepton is selected according to dilepton mass:
Near/**F**ar leptons are discriminated by their p_T

Preselection cuts

Mostly against
Standard Model

Lepton selection

Cuts on single
isolated leptons

Analysis cuts

Based on model kinematics,
to be applied after observing
invariant mass endpoints

Left : $m(\text{ll}) < 57 \text{ GeV}$

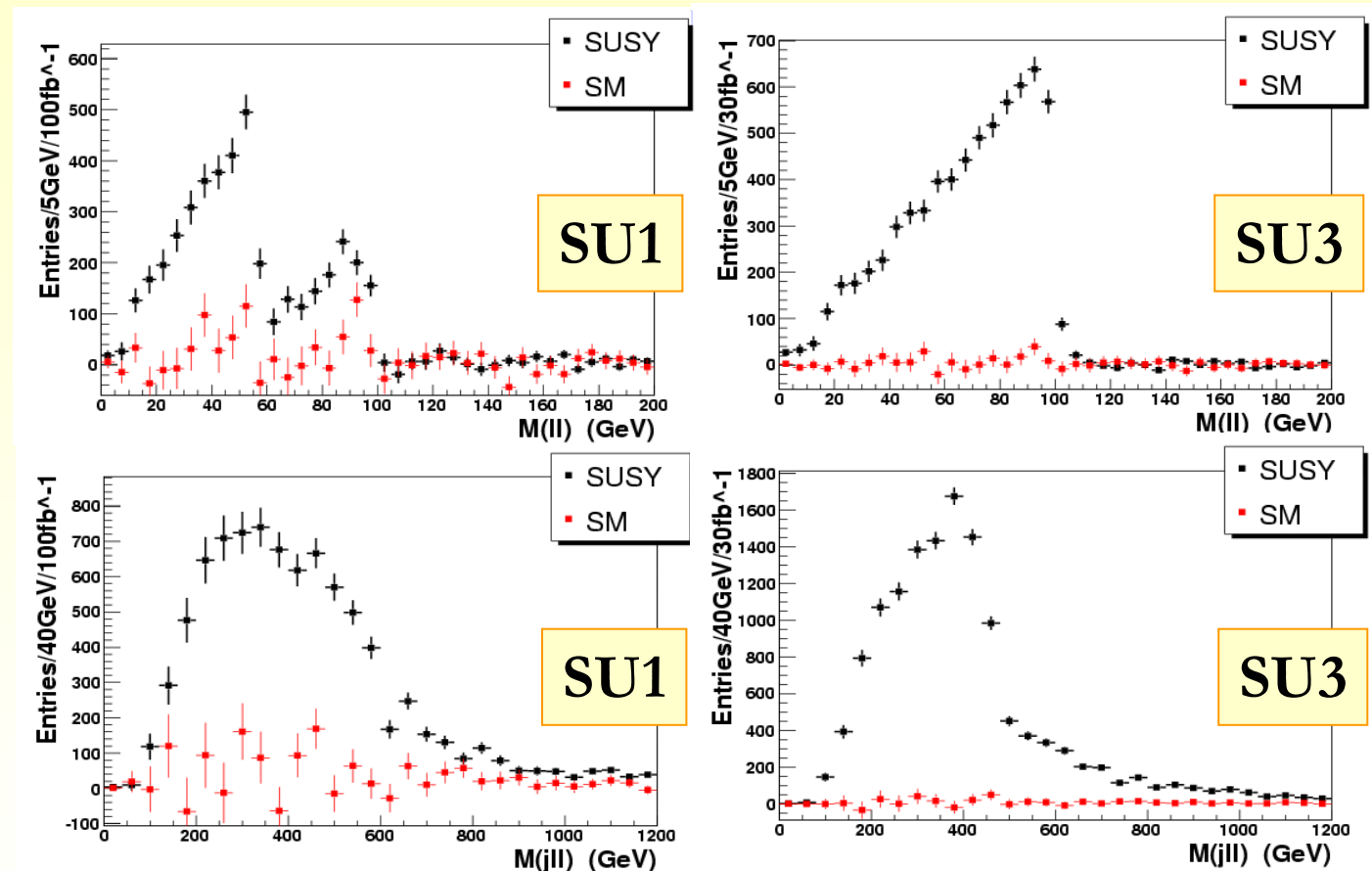
Right : $57 \text{ GeV} < m(\text{ll}) < 100 \text{ GeV}$

Reconstructed invariant masses

- Simulated (*fast*) SM processes:

- $t\bar{t}$ + jets
- W+jets
- Z+jets

- The first two energetic jets in the event are used to form $m(jl)$: no significant effect on A_{ql} due to wrong jet



- Events with two opposite-flavour opposite-sign (**OFOS**) leptons are kept and statistically subtracted to SFOS events to cancel background from uncorrelated lepton pairs: $\mu^+\mu^- + e^+e^- - \mu^\pm e^\mp$

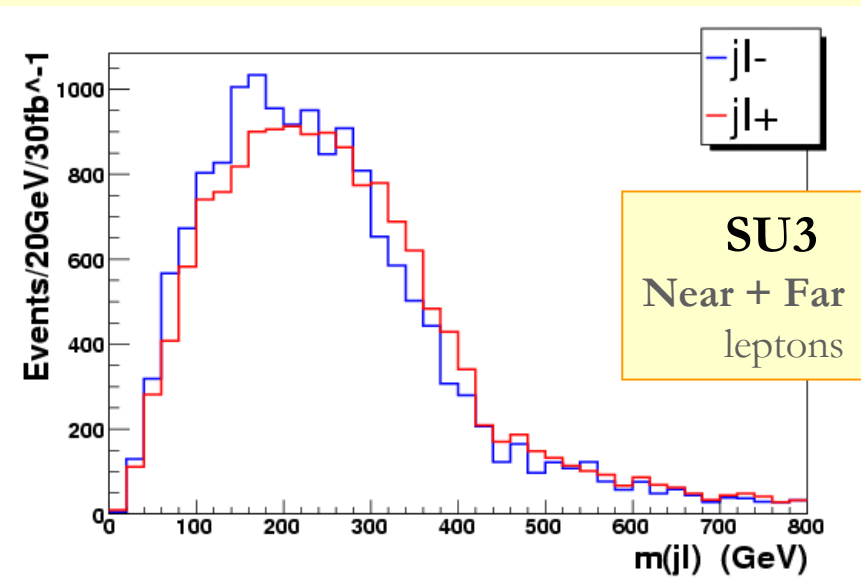
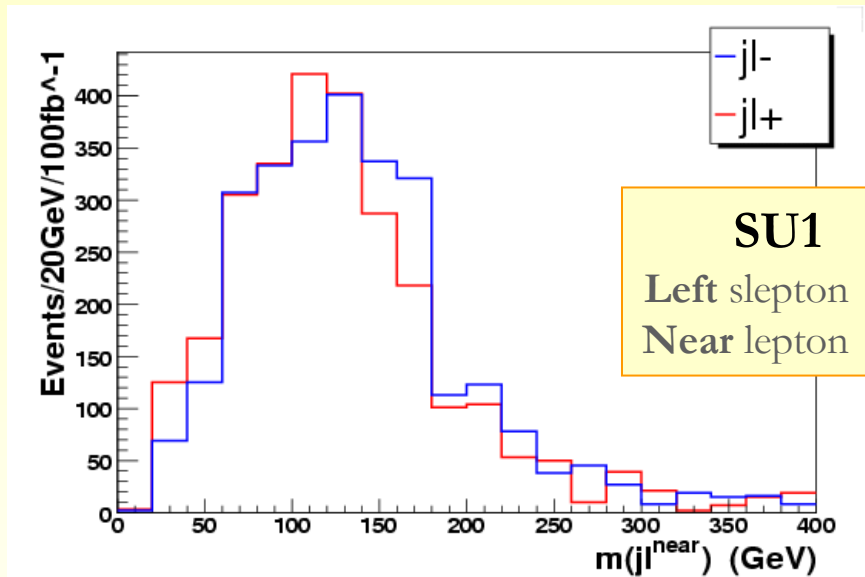
Efficiencies and contaminations

Before <i>OFOS subtraction</i>	Efficiency (SU1)	S/B (SU1)	Efficiency (SU3)	S/B (SU3)
Signal	$(17.0 \pm 0.3) \%$	/	$(20.0 \pm 0.3) \%$	/
SUSY Background	$(0.94 \pm 0.01) \%$	0.33	$(0.75 \pm 0.01) \%$	1
$t\bar{t} + 0$ jets ($lvqq$)	$(1.3 \pm 0.3) 10^{-6}$	~ 80	$(4 \pm 1) 10^{-7}$	~ 1500
$t\bar{t} + 1$ jets ($lvqq$)	$(2.9 \pm 0.2) 10^{-5}$	6.3	$(1.2 \pm 0.1) 10^{-5}$	90
$t\bar{t} + 2$ jets ($lvqq$)	$(1.3 \pm 0.1) 10^{-4}$	3.5	$(5.9 \pm 0.4) 10^{-5}$	45
$t\bar{t} + \geq 3$ jets ($lvqq$)	$(5.3 \pm 0.2) 10^{-4}$	2.6	$(2.3 \pm 0.1) 10^{-4}$	37
$t\bar{t} + 0$ jets ($lv\nu\nu$)	$(1.9 \pm 0.4) 10^{-6}$	~ 200	$(3.0 \pm 0.5) 10^{-6}$	~ 1000
$t\bar{t} + 1$ jets ($lv\nu\nu$)	$(1.5 \pm 0.1) 10^{-4}$	4.8	$(2.2 \pm 0.1) 10^{-4}$	20
$t\bar{t} + 2$ jets ($lv\nu\nu$)	$(2.93 \pm 0.04) 10^{-3}$	0.64	$(3.89 \pm 0.05) 10^{-3}$	2.9
$t\bar{t} + \geq 3$ jets ($lv\nu\nu$)	$(1.65 \pm 0.02) \%$	0.34	$(2.10 \pm 0.02) \%$	1.6
$W + 4$ jets ($lv\nu\nu$)	$(1 \pm 1) 10^{-5}$	24	$(2 \pm 2) 10^{-6}$	~ 1000
$W + \geq 5$ jets ($lv\nu\nu$)	$(2 \pm 2) 10^{-5}$	50	$(1 \pm 1) 10^{-5}$	~ 500
$Z + 3$ jets	$(3.1 \pm 0.4) 10^{-5}$	~ 500	$(3.1 \pm 0.4) 10^{-5}$	~ 2000
$Z + 4$ jets	$(2.0 \pm 0.1) 10^{-3}$	27	$(1.7 \pm 0.1) 10^{-3}$	~ 200
$Z + \geq 5$ jets	$(7.0 \pm 0.2) 10^{-3}$	24	$(5.6 \pm 0.2) 10^{-3}$	~ 200

■ After *OFOS subtraction* :

- most of irreducible background comes from **SUSY** (e.g.: *correlated leptons from direct/indirect χ_2^0 production*), with $S/B > 1$ (for SU1)
- **SM** background from $t\bar{t}bar/W$ +jets becomes compatible with 0. Small Z+jets fraction survives.

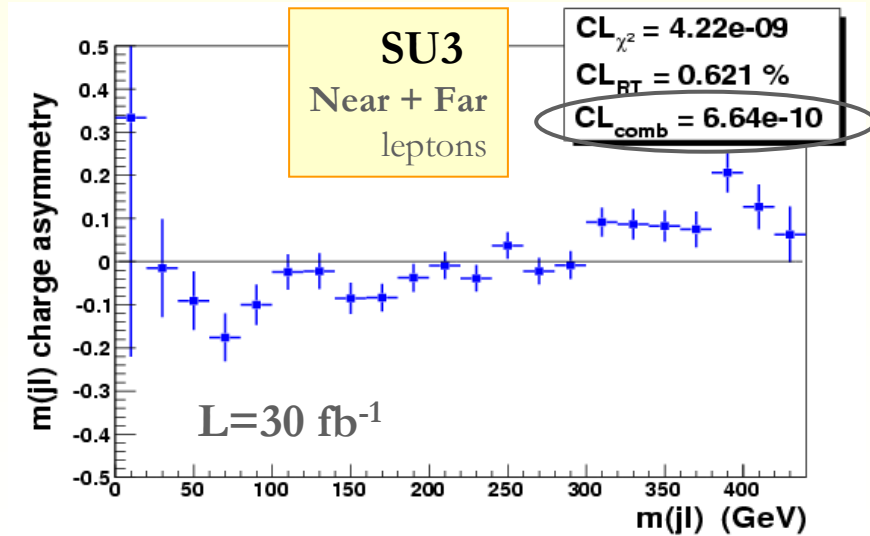
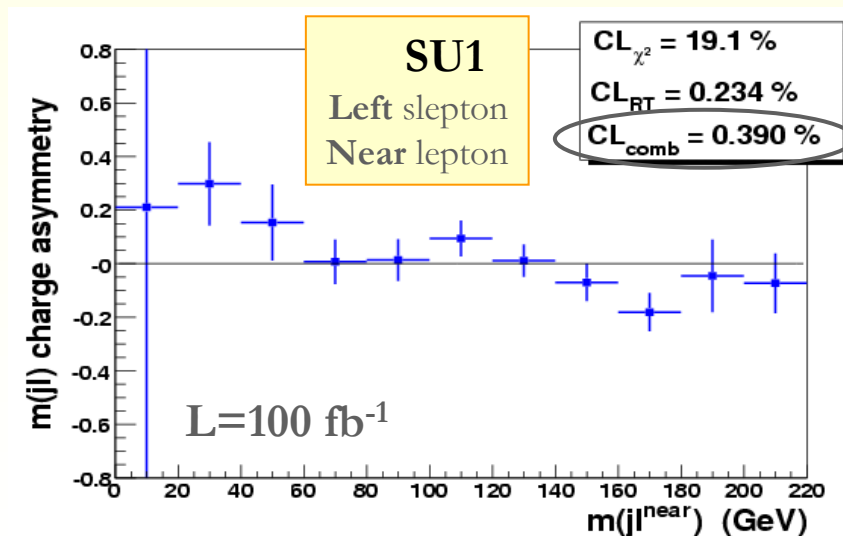
$m(jl^{\pm})$ reconstructed distributions



- After event selection and OFOS subtraction, reconstructed $m(jl^+)$ and $m(jl^-)$ are taken bin-by-bin to get asymmetry:
$$A_{ql,i} = \frac{m_i(jl^+) - m_i(jl^-)}{m_i(jl^+) + m_i(jl^-)}$$
- Suitable binning to get a compromise between:
 - good granularity in $m(jl)$,
 - proper statistical treatment of asymmetry and errors.
- $m(jl)$ ranges fixed according to endpoints and expected resolutions: **[0,220] GeV (for SU1)** and **[0,420] GeV (for SU3)**

Evaluation of $m(jl^\pm)$ asymmetry

- Two statistical methods are used to detect a non-zero charge asymmetry:
 - ▣ a non-parametric χ^2 test to a zero constant (= hypothesis of symmetry),
 - ▣ the *Run Test** method, based on the number of observed (vs. expected) fluctuations of bin contents with respect to zero.
- The two methods provide independent **confidence levels** : CL_{χ^2} and CL_{RT} , which can be combined to get CL_{comb} .
- Low confidence level \Rightarrow good evidence of a non-zero asymmetry.



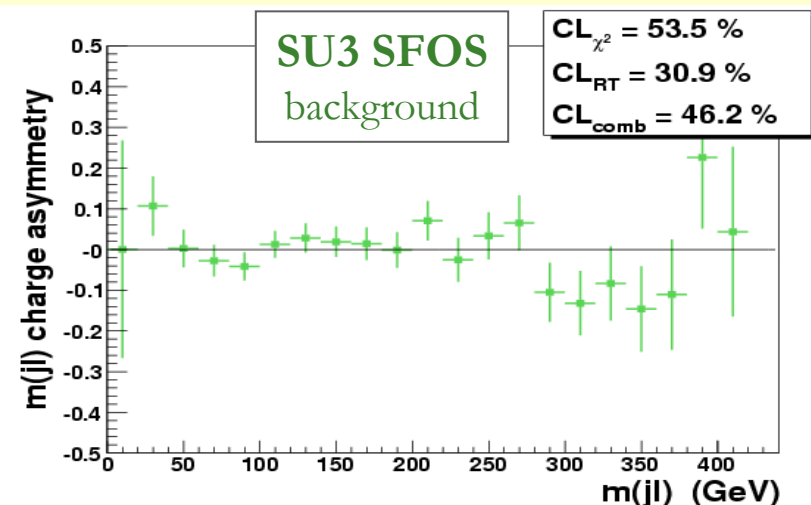
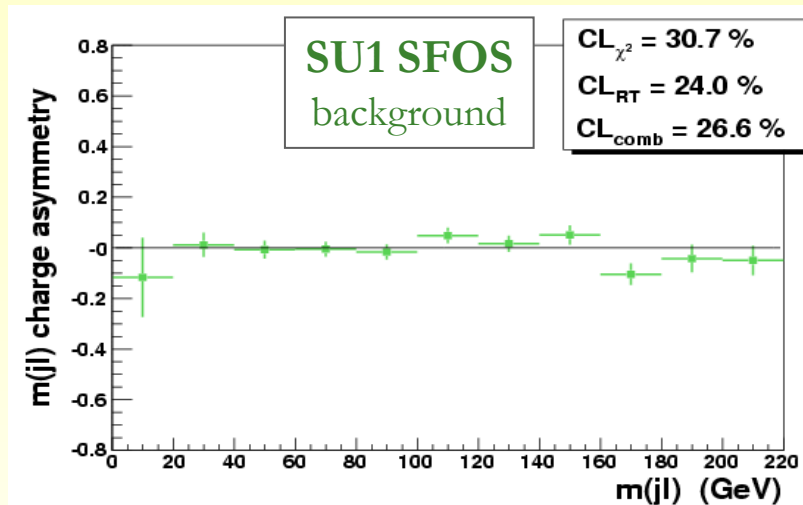
Diluting effects on $m(jl^\pm)$ asymmetry

- a) **SUSY** SFOS background contribution
- b) **SUSY** OFOS contribution
- c) **SM** SFOS and OFOS contributions
- d) Selection of *wrong jet* in $m(jl^\pm)$

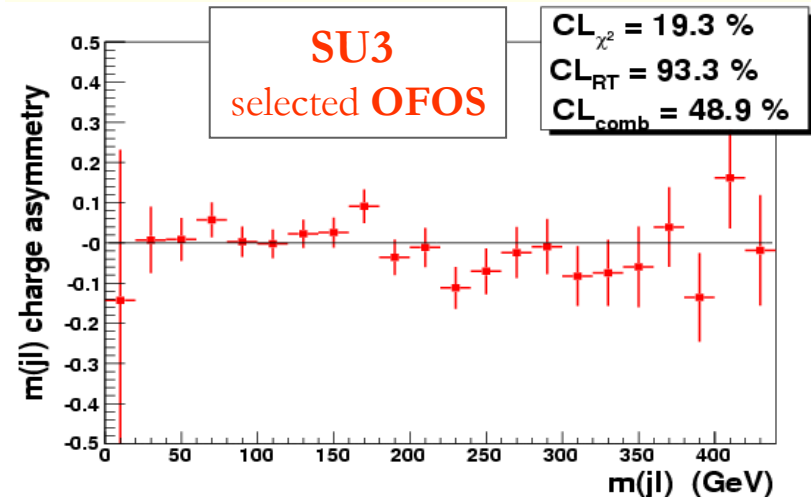
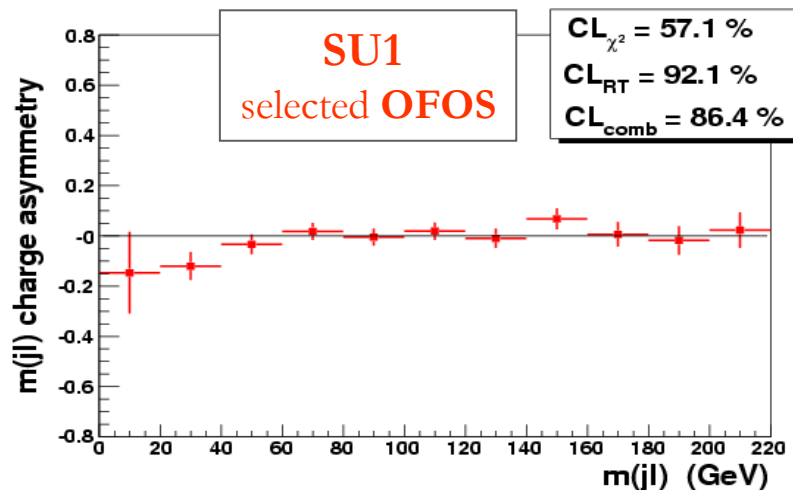
Although not definitively suppressed in the final selection, all of these **background** contributions to \mathbf{A}_{q1} are proved to have a flat behaviour vs. $m(jl^\pm)$ and not to spoil significantly the contribution given by **signal** events.

SUSY background asymmetry

a)



b)



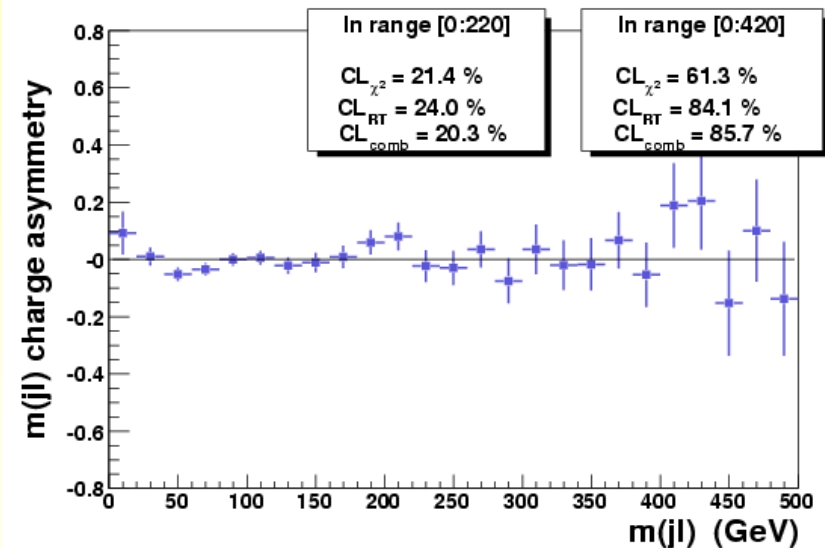
- No relevant impact on the final A_{q_l} asymmetry can be observed

SM background asymmetry

- Both **SFOS** and **OFOS** SM selected lepton pairs show reasonably flat charge asymmetries with high CL_{comb} compared to SUSY (signal+background) events.
- This is tested on both $m(jl)$ ranges used for SU1 and SU3.
- Similar results varying bin-width and range values.

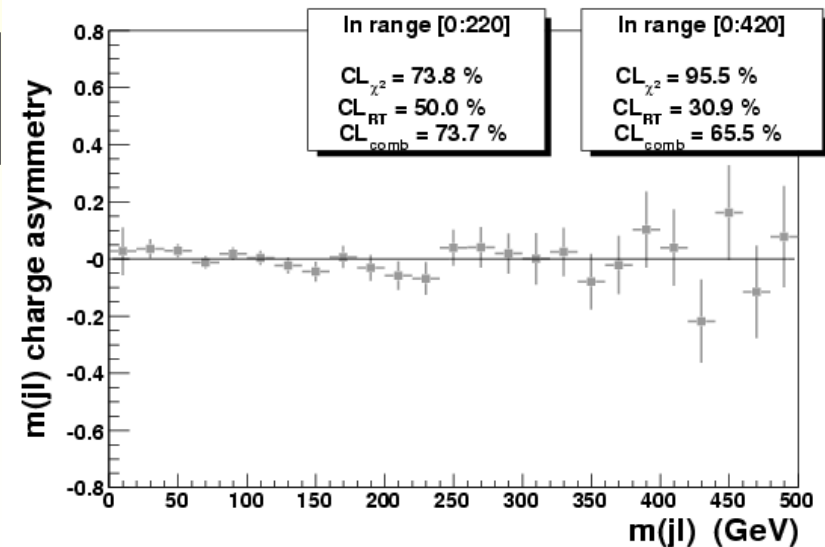
SM SFOS
background

c')



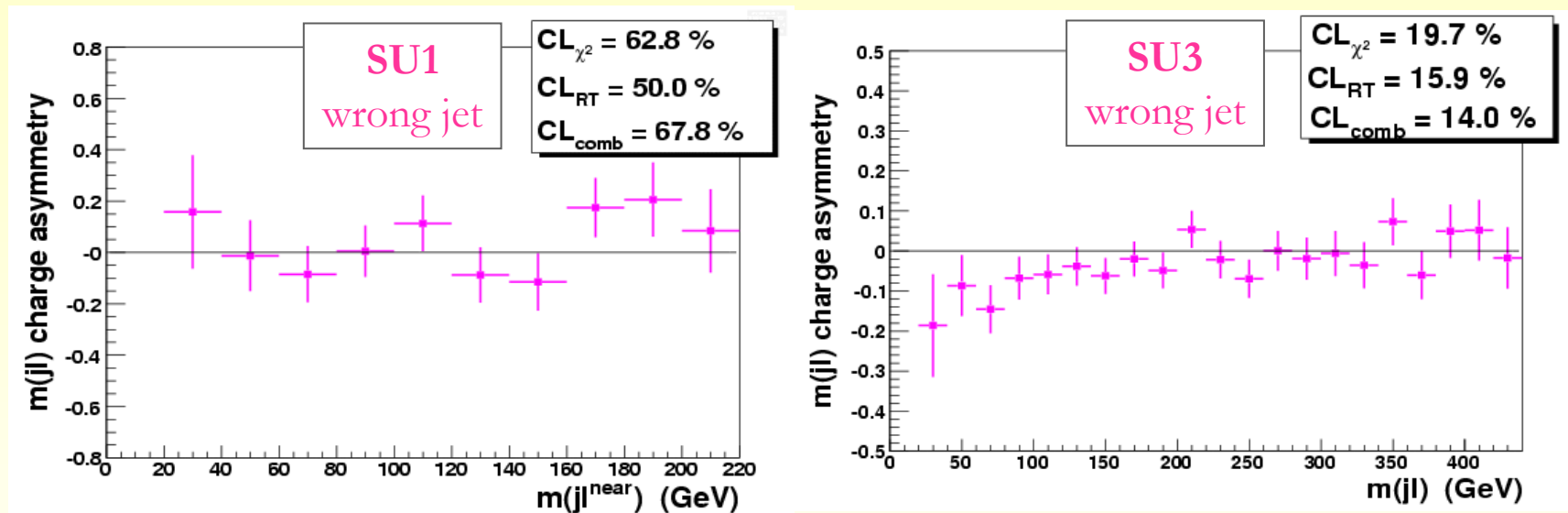
SM OFOS
selection

c'')



Wrong jet selection

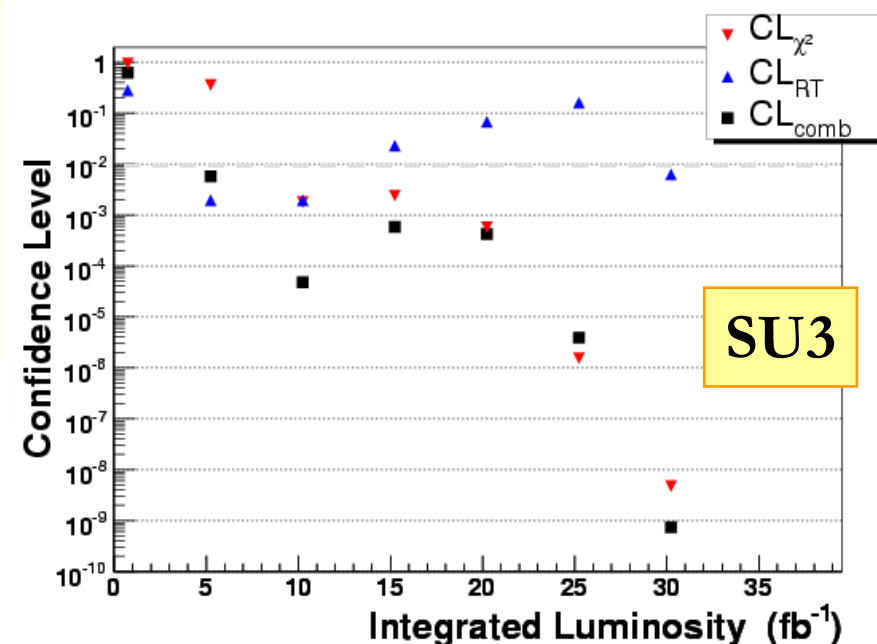
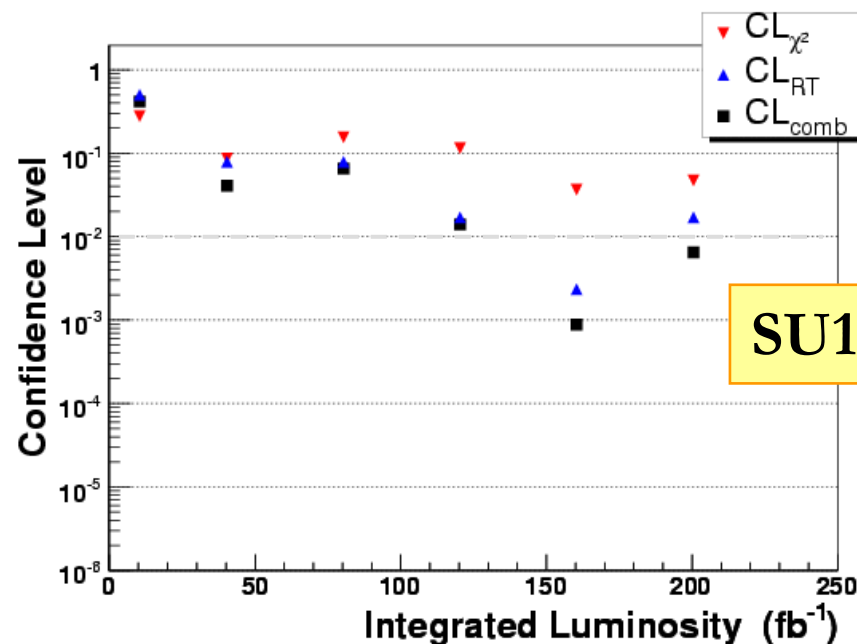
d)



- In both SU1 and SU3 points, **SFOS** events with a wrongly selected jet (wrt the signal quark) provide no significant $m(jl^{\pm})$ asymmetry.
- Selecting the two highest p_T jets in the event systematically brings to retain (at least) one wrong jet: this is an acceptable compromise for keeping the largest possible signal yield without spoiling asymmetry.

Confidence level vs. luminosity

- Results have been obtained for different integrated luminosities.
 - In **SU1** point at least 100 fb^{-1} are needed to observe some non-zero charge asymmetry (e.g. with $CL_{\text{comb}} < 1\%$).
 - For the case of **SU3** point, $5 \div 10 \text{ fb}^{-1}$ could be already enough.



Summary and conclusions

- If **SUSY** will be discovered at **LHC**, measuring properties of new particles (*e.g.* spin) is needed for demonstrating that they are indeed the predicted super-partners.
- The method used here on two selected **mSUGRA** points shows that **m(q1[±])** charge asymmetry is detectable with few fb⁻¹ (in SU3) or at least 100fb⁻¹ (in SU1).
- The study must be also performed with more realistic (*full*) ATLAS simulation and for other SUSY points.
- New techniques to detect charge asymmetry and reduce diluting effects (background, systematics, etc.) are under investigation.

Backup slides

Branching ratios & mass endpoints

Point SU1

$$\begin{aligned}\sigma(\tilde{q}_R\tilde{g}) &= 1.757 \text{ pb} , & \sigma(\tilde{q}_L\tilde{g}) &= 1.620 \text{ pb} , \\ \sigma(\tilde{q}_L\tilde{q}_R) &= 0.885 \text{ pb} , & \sigma(\tilde{q}_L\tilde{q}_L) &= 0.665 \text{ pb} , \\ \sigma(\tilde{g}\tilde{g}) &= 0.554 \text{ pb} , & \sigma(\tilde{q}_L\tilde{\chi}_2^0) &= 0.154 \text{ pb} .\end{aligned}$$

$$\begin{aligned}BR(\tilde{q}_L \rightarrow q\tilde{\chi}_2^0) &= 31.5 \% , & BR(\tilde{\chi}_2^0 \rightarrow \tilde{l}_L l) &= 6 \% , \\ BR(\tilde{\chi}_2^0 \rightarrow \tilde{l}_R l) &= 3 \% , & BR(\tilde{l}_{L,R} \rightarrow \tilde{\chi}_1^0 l) &= 100 \% .\end{aligned}$$

$$\begin{aligned}m(ll)^{max} &= \left[\frac{(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{l}_{L,R}}^2)(M_{\tilde{l}_{L,R}}^2 - M_{\tilde{\chi}_1^0}^2)}{M_{\tilde{l}_{L,R}}^2} \right]^{1/2} = \\ &= 56.05 \text{ GeV (left), } 97.93 \text{ GeV (right)} , \\ m(qll)^{max} &= \left[\frac{(M_{\tilde{q}_L}^2 - M_{\tilde{\chi}_2^0}^2)(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{\chi}_1^0}^2)}{M_{\tilde{\chi}_2^0}^2} \right]^{1/2} = \\ &= 611.64 \text{ GeV (left, right)} , \\ m(ql^{near})^{max} &= \left[\frac{(M_{\tilde{q}_L}^2 - M_{\tilde{\chi}_2^0}^2)(M_{\tilde{\chi}_2^0}^2 - M_{\tilde{l}_{L,R}}^2)}{M_{\tilde{\chi}_2^0}^2} \right]^{1/2} = \\ &= 180.14 \text{ GeV (left), } 580.09 \text{ GeV (right)} , \\ m(ql^{far})^{max} &= \left[\frac{(M_{\tilde{q}_L}^2 - M_{\tilde{\chi}_2^0}^2)(M_{\tilde{l}_{L,R}}^2 - M_{\tilde{\chi}_1^0}^2)}{M_{\tilde{l}_{L,R}}^2} \right]^{1/2} = \\ &= 327.21 \text{ GeV (right), } 603.08 \text{ GeV (left)} .\end{aligned}$$

Point SU3

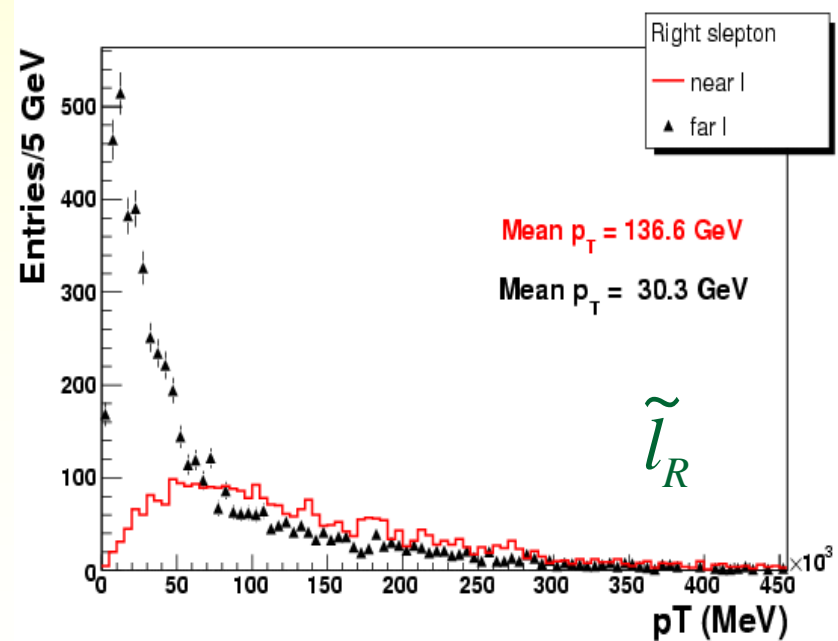
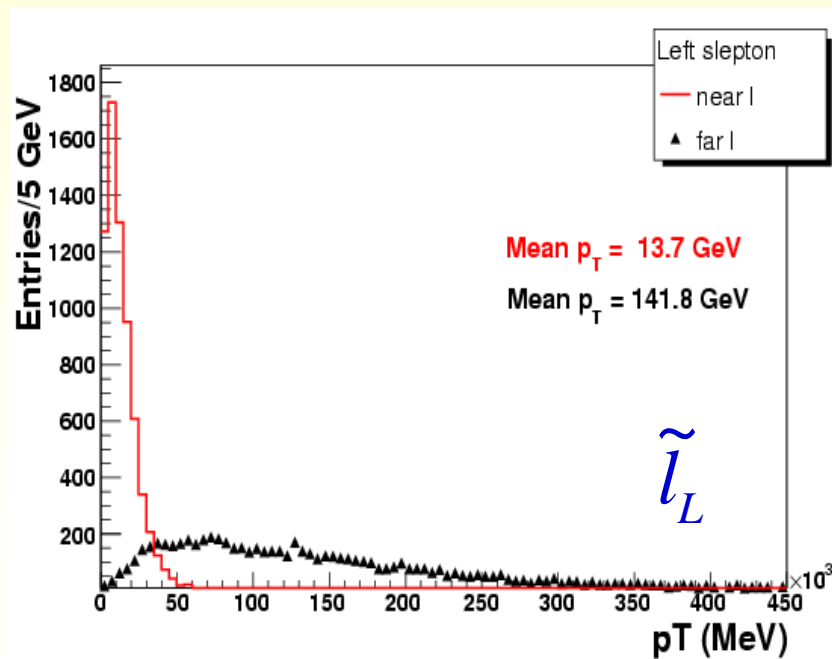
$$\begin{aligned}\sigma(\tilde{q}_R\tilde{g}) &= 4.469 \text{ pb} , & \sigma(\tilde{q}_L\tilde{g}) &= 4.426 \text{ pb} , \\ \sigma(\tilde{q}_L\tilde{q}_R) &= 2.085 \text{ pb} , & \sigma(\tilde{q}_L\tilde{q}_L) &= 1.716 \text{ pb} , \\ \sigma(\tilde{g}\tilde{g}) &= 1.544 \text{ pb} , & \sigma(\tilde{q}_L\tilde{\chi}_2^0) &= 0.203 \text{ pb} .\end{aligned}$$

$$\begin{aligned}BR(\tilde{q}_L \rightarrow q\tilde{\chi}_2^0) &= 32 \% , & BR(\tilde{\chi}_2^0 \rightarrow \tilde{l}_R l) &= 17.6 \% , \\ BR(\tilde{l}_{L,R} \rightarrow \tilde{\chi}_1^0 l) &= 100 \% ,\end{aligned}$$

$$\begin{aligned}m(ll)^{max} &= 100.17 \text{ GeV} , \\ m(qll)^{max} &= 500.34 \text{ GeV} , \\ m(ql^{near})^{max} &= 417.74 \text{ GeV} , \\ m(ql^{far})^{max} &= 385.85 \text{ GeV} .\end{aligned}$$

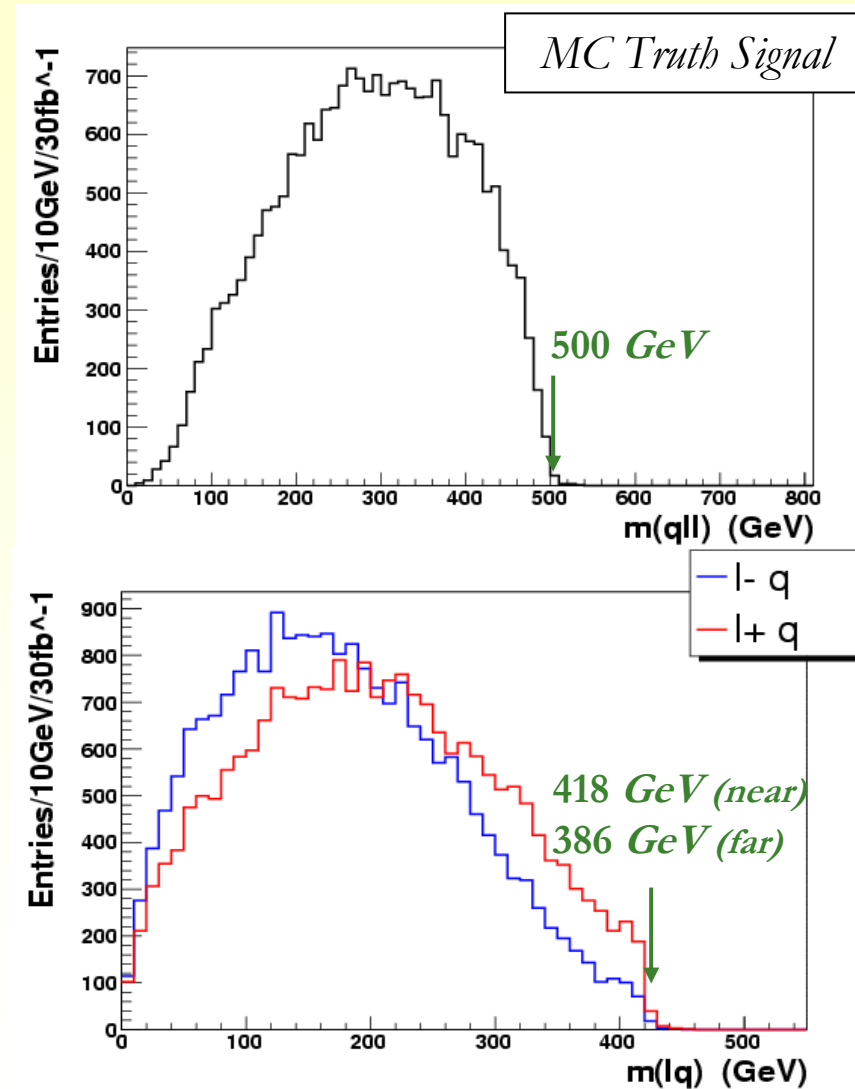
Distinguishability of leptons in SU1

- After **Left/Right** slepton decay chain identification, *near* and *far* leptons are easily distinguishable from their momenta



SU3 mass distributions for truth signal

- Lepton-lepton-quark invariant mass distribution
- Lepton-quark invariant mass distribution



SUSY vs. Universal Extra Dimensions

J. M. Smillie and B. R. Webber, *JHEP* (2005) 069, [[hep-ph/0507170](#)]

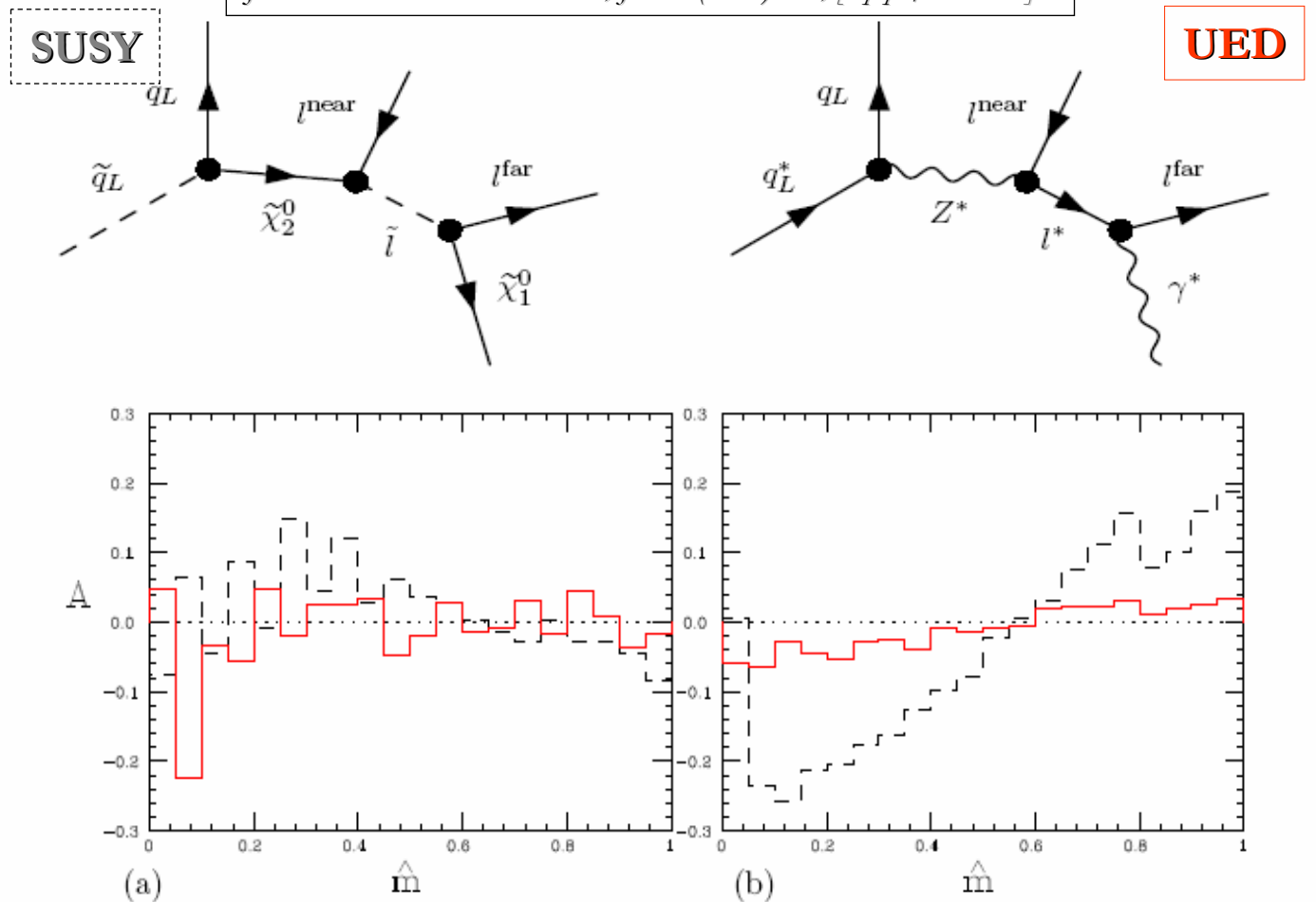


Figure 12: Detector-level charge asymmetries with respect to the jet + lepton rescaled invariant mass, for the (a) UED and (b) SUSY mass spectra given above. Dashed: SUSY. Solid/red: UED.