

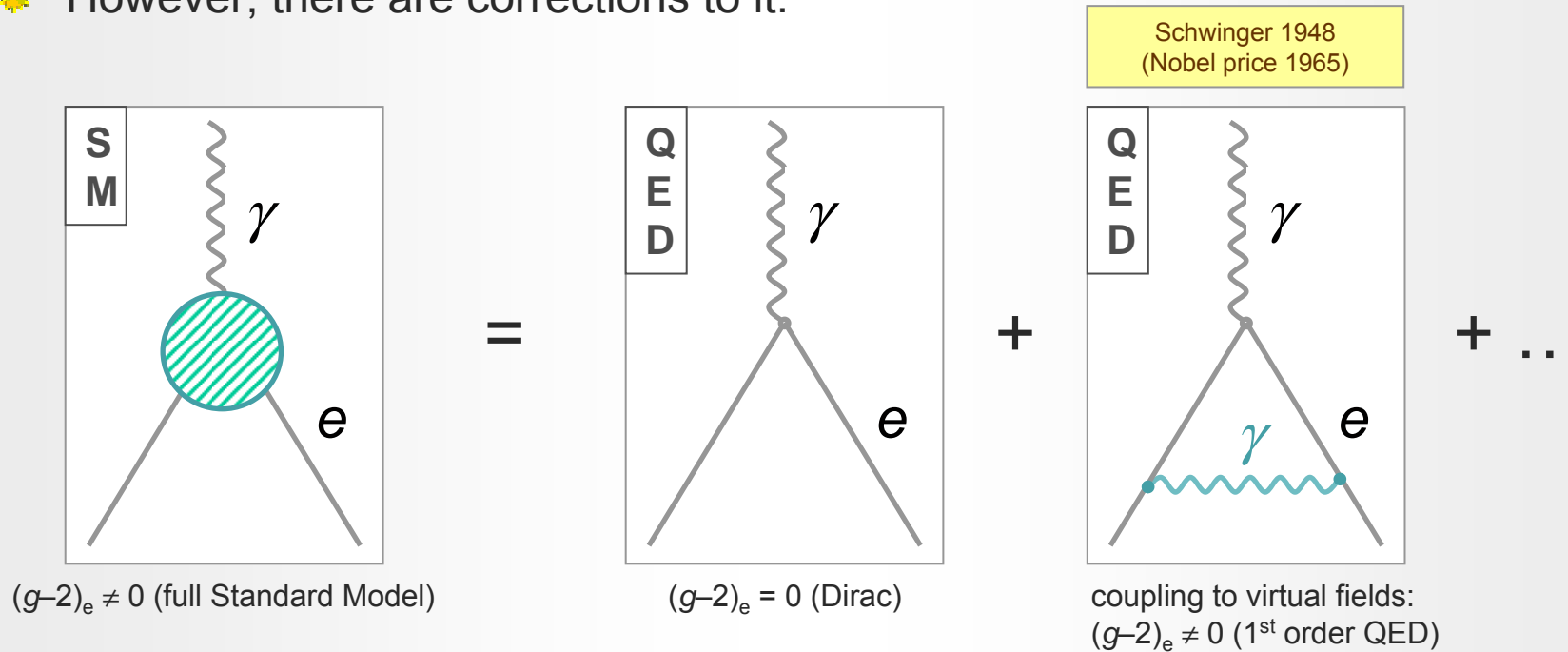
# Evaluation of the Hadronic Vacuum Polarisation Contribution to $(g - 2)_\mu$

Michel Davier (LAL) & Andreas Höcker (CERN)

Flavour in the era of the LHC workshop, October 9–11, 2006, CERN, Switzerland

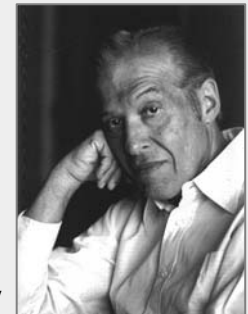
# The *Anomalous* Magnetic Moment ...of the electron

- Dirac's gyromagnetic factor corresponds to the lowest order QED graph
- However, there are corrections to it:



- Quantum fluctuations shift the gyromagnetic ratio

$$a_e \equiv \frac{g_e - 2}{2} = \frac{\alpha}{2\pi} + \dots = \frac{1}{2\pi} \underbrace{\frac{e^2}{4\pi}}_{1/137.036\dots} + \dots = 0.001161\dots$$



Schwinger,  
finding QED easy

# The *Muon* Anomalous Magnetic Moment

☀ Diagrams contributing to the magnetic moment

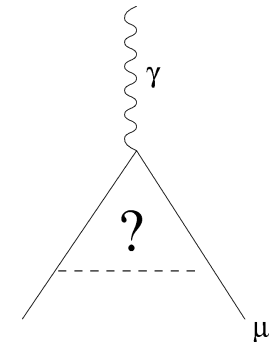
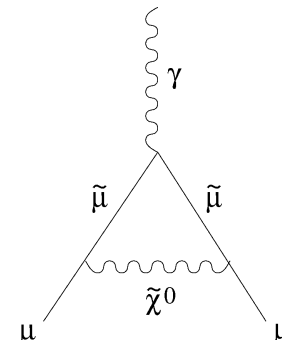
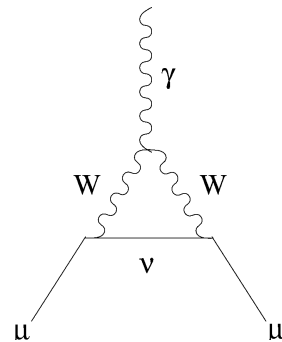
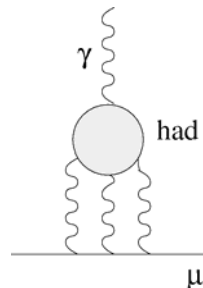
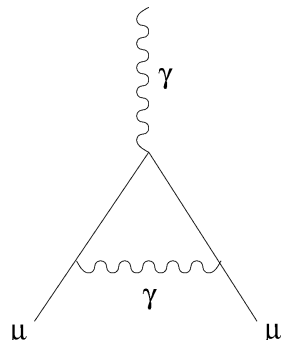
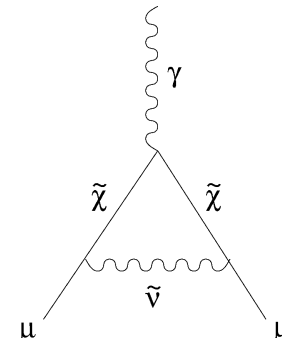
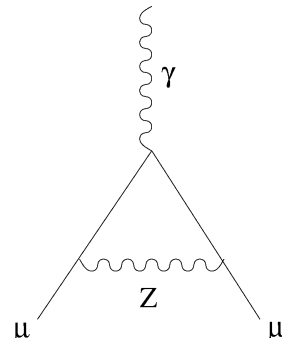
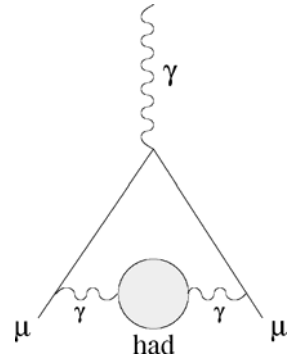
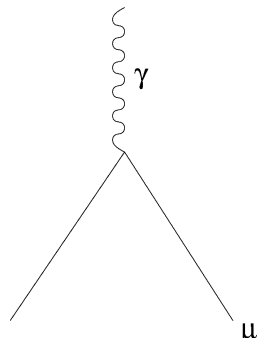
QED

Hadronic

Weak

SUSY...

... or some unknown type of new physics ?



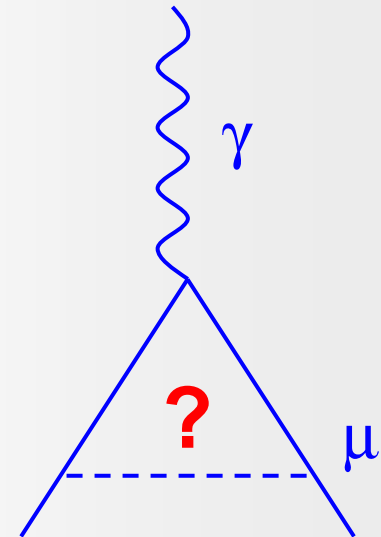
# The Quest for “New Physics”

☀ The experimental precision for  $a_\mu$  will be worse than for  $a_e$ , so why do it ?

🖥 From chiral symmetry expect the New Physics (NP) effects to scale  $\sim m^2(e/\mu)$ :

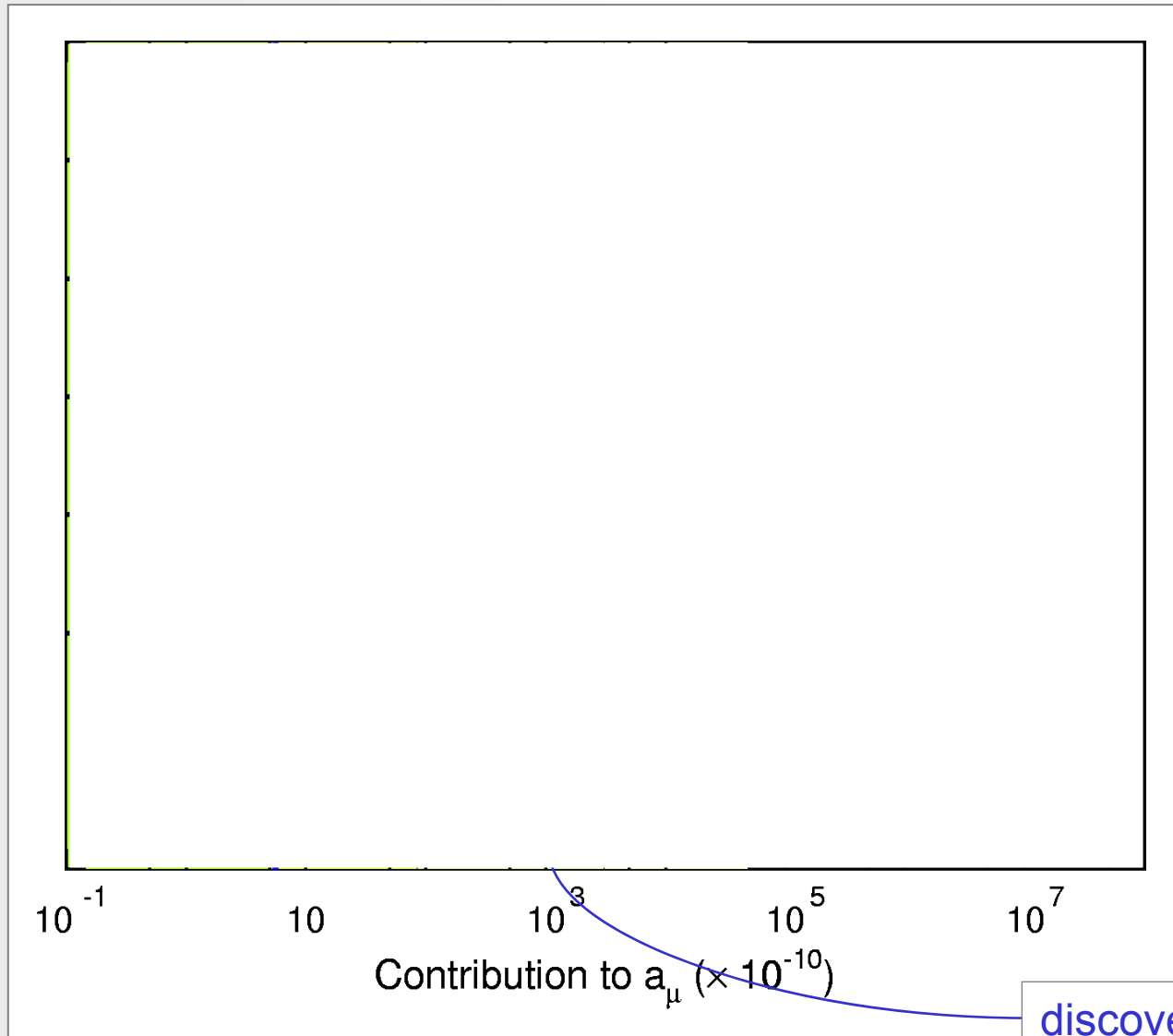
$$a_\ell^{\text{NP}}(\Lambda_{\text{NP}}) \propto \mathcal{O}\left(\frac{m_\ell^2}{\Lambda_{\text{NP}}^2}\right) \implies \frac{a_\mu^{\text{NP}}}{a_e^{\text{NP}}} \propto \mathcal{O}\left(\frac{m_\mu^2}{m_e^2}\right) \approx 42,000$$

🖥 Expect to loose about a factor of 200 in experimental precision



$\implies$   $a_\mu$  should be roughly 200 times more sensitive to NP than  $a_e$  !

# Experimental Progress from CERN to BNL



Experimental progress  
on precision of  $(g-2)_\mu$

Outperforms theory  
precision on **hadronic  
contribution**

discovery of magic  $\gamma$

J. Bailey et al.,  
NP B150 1 (1979)



# Confronting Experiment with Theory

The Standard Model prediction of  $a_\mu$  is decomposed in its main contributions:

$$a_\mu^{\text{SM}} \equiv \left( \frac{g-2}{2} \right)_\mu = a_\mu^{\text{QED}} + a_\mu^{\text{had}} + a_\mu^{\text{weak}}$$

of which the **hadronic contribution** has the largest uncertainty

# The Muon Magnetic Anomaly in the Standard Model

## QED contribution

Computed up to 4<sup>th</sup> order  
(5<sup>th</sup> order estimated)

$$a_{\mu}^{\text{QED}} \approx \left( \begin{array}{l} 11,614,098.1 + 41,321.8 \\ + 3,014.2 + 38.1 + \underbrace{0.6}_{\text{estimated}} \end{array} \right) \times 10^{-10} = 11,658,472.1(0.1) \times 10^{-10}$$

Kinoshita-Nio, PRD 70, 113001 (2005)

## Electroweak contribution

Computed up to 2<sup>nd</sup> order

$a_{\mu}^{\text{weak}}$  suppressed by  $\frac{\alpha}{\pi} \frac{m_{\mu}^2}{m_W^2} \sim 10^{-9}$  (!)

$$\underbrace{a_{\mu}^{\text{weak}}}_{\text{1-loop}} = \frac{G_{\mu} m_{\mu}^2}{8\sqrt{2}\pi^2} \left( \frac{5}{3} + \frac{1}{3} (1 - 4 \sin^2 \theta_W) + \mathcal{O}\left(\frac{m_{\mu}^2}{m_W^2}\right) + \mathcal{O}\left(\frac{m_{\mu}^2}{m_H^2}\right) \right) = +19.5 \times 10^{-10}$$

Czarnecki *et al.*,  
PRD 52, 2619 (1995);  
PRL 76, 3267 (1996)

2<sup>nd</sup> order contribution surprisingly large:  
( due to large logs:  $\ln[m_Z/m_{\mu}]$  )

$$\underbrace{a_{\mu}^{\text{weak}}}_{\text{2-loop}} = -4.1(0.2) \times 10^{-10}$$

Note that between  $a_{\mu}$  and  $a_e$ , the same sensitivity factor as for “new physics” applies here

# The Hadronic Contribution to $(g-2)_\mu$

$$a_\mu^{\text{SM}} = a_\mu^{\text{QED}} + a_\mu^{\text{had}} + a_\mu^{\text{weak}}$$

The Situation 1995

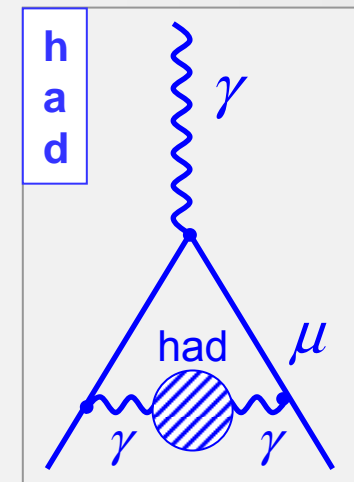
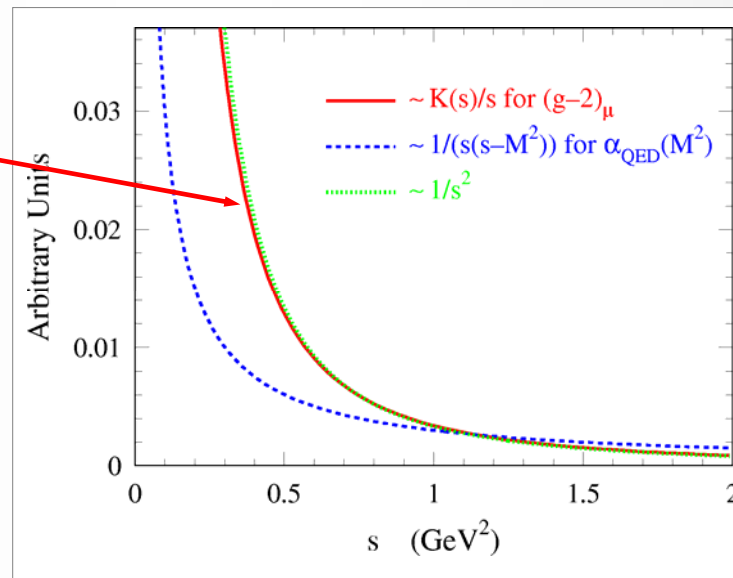
Source	$\alpha(a_\mu)$	Reference
QED	$\sim 0.1 \times 10^{-10}$	[Schwinger '48 & others]
Hadrons	$\sim (15 \oplus 4) \times 10^{-10}$	[Eidelman-Jegerlehner '95 & others]
Z, W exchange	$\sim 0.2 \times 10^{-10}$	[Czarnecki <i>et al.</i> '95 & others]

Dominant uncertainty from lowest order hadronic piece. Cannot be calculated from QCD ("first principles") – but: **we can use experiment (!)**

$$a_\mu^{\text{had,LO}} = \frac{\alpha^2}{3\pi^2} \int_{4m_\pi^2}^{\infty} ds \frac{K(s)}{s} R(s)$$

Dispersion relation, uses unitarity (optical theorem) and analyticity

... (see digression)





# digression: Vacuum polarization ... and the running of $\alpha_{\text{QED}}$

Define: photon vacuum polarization function  $\Pi_\gamma(q^2)$

$$i \int d^4x e^{iqx} \langle 0 | T J_{\text{em}}^\mu(x) (J_{\text{em}}^\nu(0))^\dagger | 0 \rangle = -(g^{\mu\nu} q^2 - q^\mu q^\nu) \Pi_\gamma(q^2)$$

Ward identities: only vacuum polarization modifies electron charge

$$\alpha(s) = \frac{\alpha(0)}{1 - \Delta\alpha(s)} \quad \text{with:} \quad \Delta\alpha(s) = -4\pi\alpha \text{Re} [\Pi_\gamma(s) - \Pi_\gamma(0)]$$

Leptonic  $\Delta\alpha_{\text{lep}}(s)$  calculable in QED. However, quark loops are modified by long-distance hadronic physics, cannot (yet) be calculated within QCD (!)

$$\text{Born: } \sigma^{(0)}(s) = \sigma(s) (\alpha / \alpha(s))^2$$

Way out: Optical Theorem (unitarity) ...

$$12\pi \text{Im} \Pi_\gamma(s) = \frac{\sigma^{(0)}[e^+e^- \rightarrow \text{hadrons}]}{\sigma^{(0)}[e^+e^- \rightarrow \mu^+\mu^-]} \equiv R(s)$$

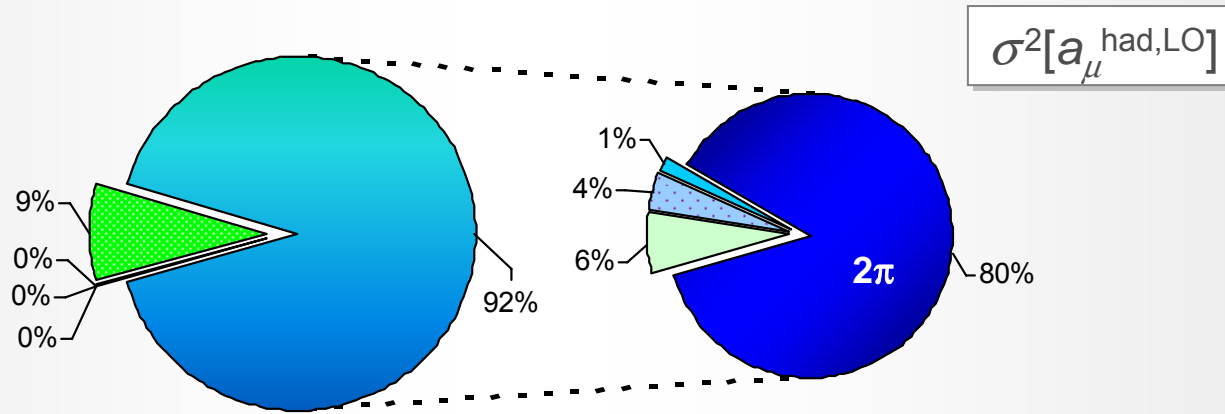
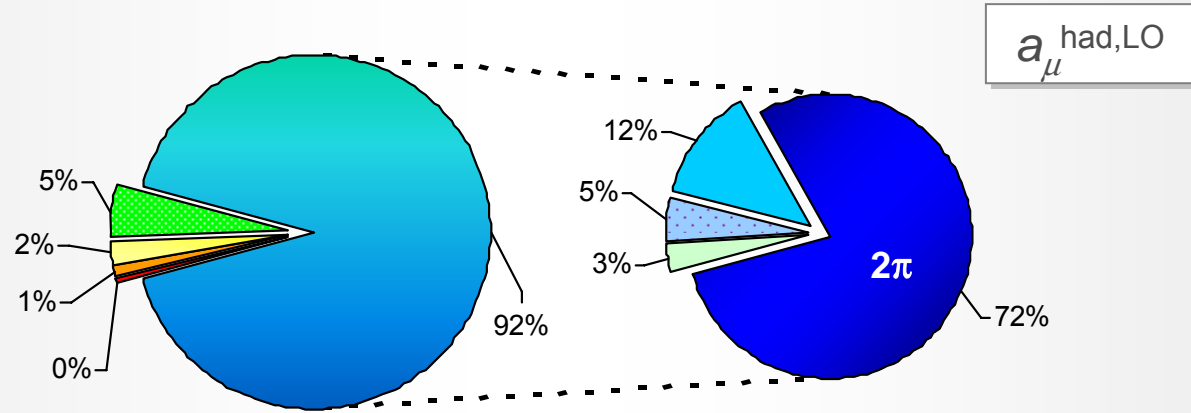
$$\text{Im} [ \text{diagram with shaded blob} ] \propto | \text{diagram with hadron loop} |^2$$

... and the subtracted dispersion relation of  $\Pi_\gamma(q^2)$  (analyticity)

$$\Pi_\gamma(s) - \Pi_\gamma(0) = \frac{s}{\pi} \int_0^\infty ds' \frac{\text{Im} \Pi_\gamma(s')}{s'(s' - s) - i\epsilon} \quad \Rightarrow \quad \Delta\alpha_{\text{had}}(s) = -\frac{\alpha s}{3\pi} \text{Re} \int_0^\infty ds' \frac{R(s')}{s'(s' - s) - i\epsilon}$$

... and equivalently for  $a_\mu$  [had]

# Contributions to the Dispersion Integrals



# Hadronic Contribution to $(g-2)_\mu$ : Roadmap

Eidelman-Jegerlehner, Z.Phys. C67, 585 (1995)

Since 1995, improved determination of the dispersion integral:

- better data
- extended use of QCD

Energy [GeV]	Input 1995	Input after 1998
$2m_\pi - 1.8$	Data	Data ( $e^+e^-$ & $\tau$ ) (+ QCD)
$1.8 - \psi(3770)$	Data	QCD
$J/\psi - \Upsilon$	Data	Data (+ QCD)
$\Upsilon - 40$	Data	QCD
$40 - \infty$	QCD	QCD

Improvement in 4 Steps:

## ■ Inclusion of precise $\tau$ data using SU(2) (CVC)

Alemamy-Davier-H.'97, Narison'01, Trocóniz-Ynduráin'01, + later works

(!)

## ■ Extended use of (dominantly) perturbative QCD

Martin-Zeppenfeld'95, Davier-H.'97, Kühn-Steinhauser'98, Erler'98, + others

## ■ Theoretical constraints from QCD sum rules

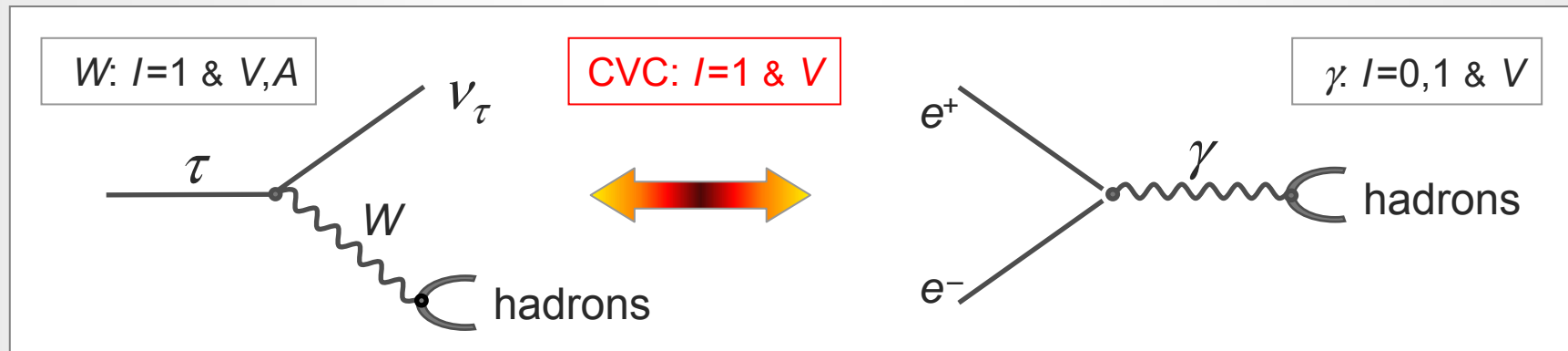
Groote-Körner-Schilcher-Nasrallah'98, Davier-H.'98, Martin-Outhwaite-Ryskin'00, Cvetič-Lee-Schmidt'01, Jegerlehner et al'00, Dorokhov'04 + others

## ■ Better data for the $e^+e^- \rightarrow \pi^+\pi^-$ cross section

CMD-2'02, KLOE'04, SND'05

(!)

# Using also Tau Data through CVC – SU(2)



Hadronic physics factorizes in [Spectral Functions](#) :

Isospin symmetry connects  $I=1$   $e^+e^-$  cross section to vector  $\tau$  spectral functions:

$$\sigma^{(I=1)}[e^+e^- \rightarrow \pi^+\pi^-] = \frac{4\pi\alpha^2}{s} \nu[\tau^- \rightarrow \pi^-\pi^0\nu_\tau]$$

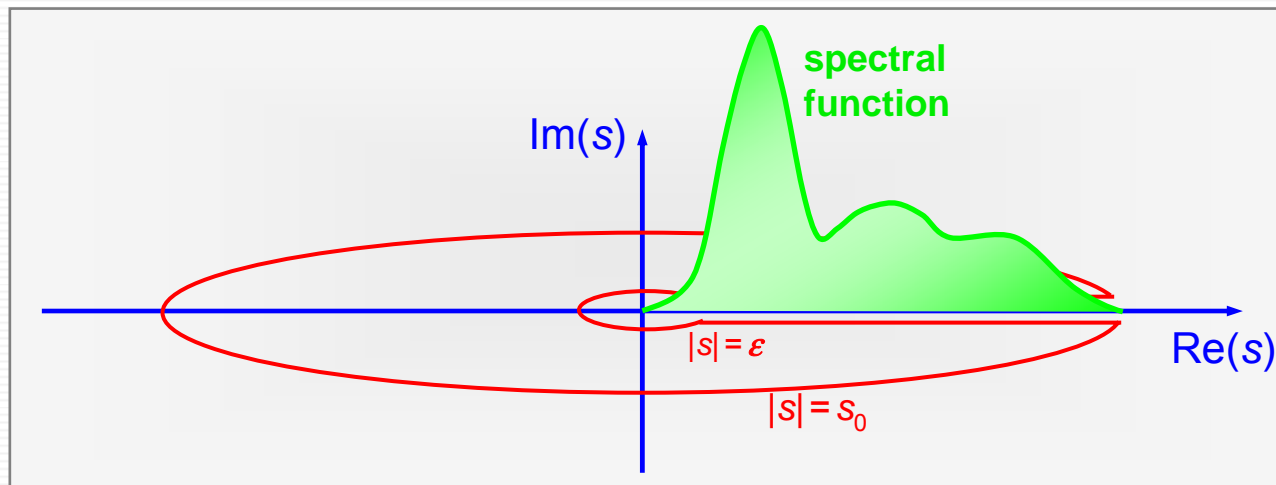
Experimentally:  $\tau$  and  $e^+e^-$  data are complementary with resp. to normalisation and shape uncertainties

$$\nu[\tau^- \rightarrow \pi^-\pi^0\nu_\tau] \propto \underbrace{\frac{\text{BR}[\tau^- \rightarrow \pi^-\pi^0\nu_\tau]}{\text{BR}[\tau^- \rightarrow e^-\bar{\nu}_e\nu_\tau]}}_{\text{branching fractions}} \underbrace{\frac{1}{N_{\pi\pi^0}} \frac{dN_{\pi\pi^0}}{ds}}_{\text{mass spectrum}} \underbrace{\frac{m_\tau^2}{(1-s/m_\tau^2)^2 (1+s/m_\tau^2)}}_{\text{kinematic factor (PS)}}$$

## digression: Tau Spectral Functions and QCD

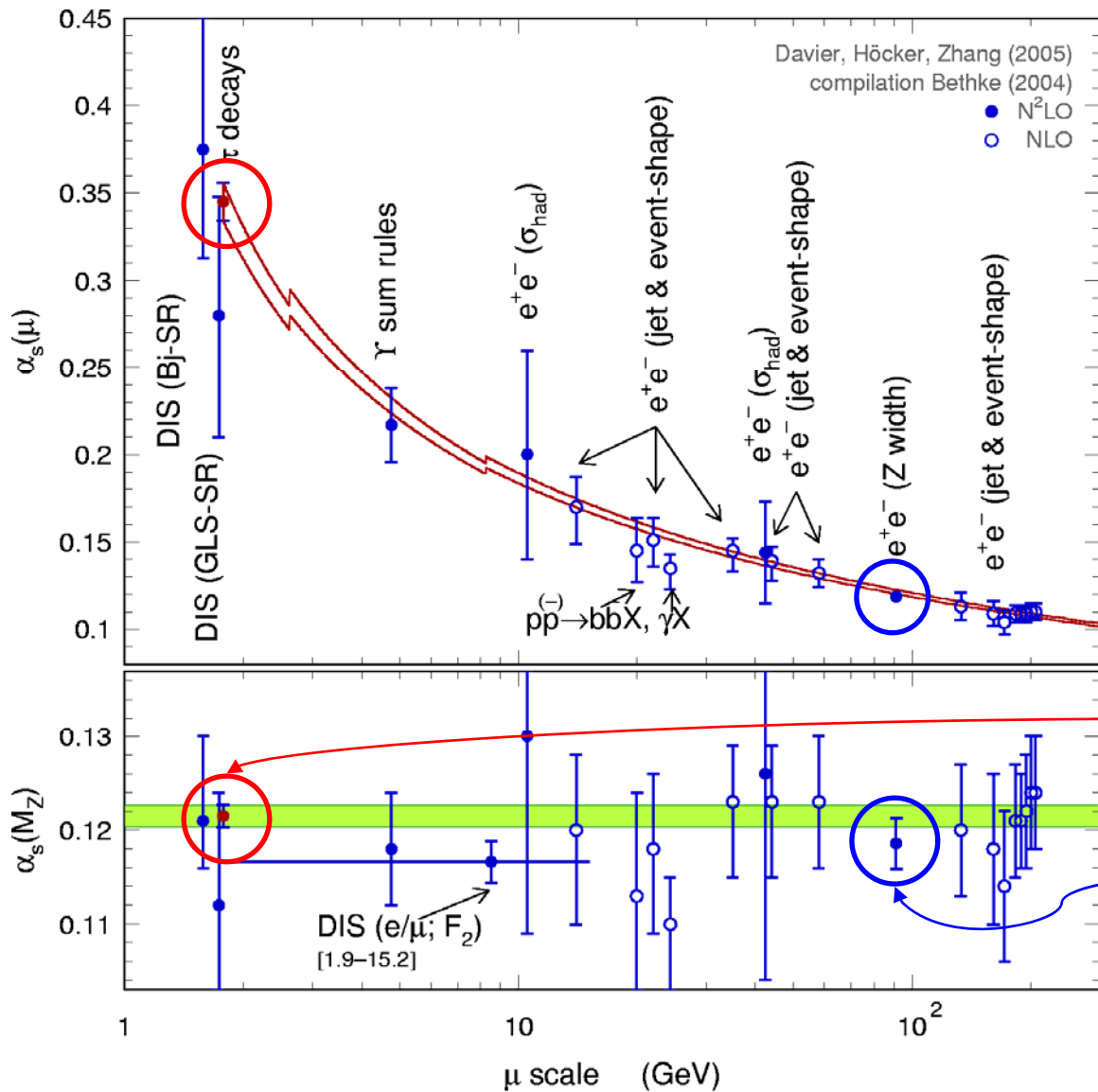
- (1) Optical theorem  $\nu(s) \propto \text{Im} \Pi(s)$
- (2) Apply Cauchy's theorem for "save" (i.e., sufficiently large)  $s_0$ :

$$R(s_0) \propto \int_0^{s_0} ds \underbrace{f(s)}_{\text{kinematic factor}} \text{Im} \Pi(s) \Leftrightarrow -\frac{1}{2i} \oint_{|s|=s_0} ds f(s) \Pi(s)$$



- (3) Use the Adler function to remove unphysical subtractions:  $D(s) = s \frac{d\Pi(s)}{ds}$
- (4) Use global quark-hadron duality in the framework of the Operator Product Expansion (OPE) to predict:  $D(s) \sim D_{\text{pert}}(s) + D_{\text{q-mass}}(s) + D_{\text{non-pert}}(s)$
- (5) Use analytical moments  $f_n(s) = f(s) \cdot \text{poly}_n(s)$  to fix non-perturbative parameters of the OPE ... and then fit  $\alpha_s(m_\tau)$

# digression: QCD Results from Hadronic Tau Decays



Evolution of  $\alpha_s(m_\tau)$ , measured using  $\tau$  decays, to  $M_Z$  using RGE (4-loop QCD  $\beta$ -function & 3-loop quark flavor matching) shows the excellent compatibility of  $\tau$  result with EW fit !

Davier-Höcker-Zhang  
 hep-ph/0507078

$\alpha_s(M_Z) = 0.1215 \pm 0.0012$   
 (DHZ'05, theory dominated)

$\alpha_s(M_Z) = 0.1186 \pm 0.0027$   
 (LEP'00, exp. dominated)

LEP EW WG,  
 LEPEWWG 2002-01

# SU(2) Breaking

Electromagnetism does not respect isospin and hence we have to consider isospin breaking when dealing with an experimental precision of 0.5%

Corrections for SU(2) breaking applied to  $\tau$  data for dominant  $\pi^-\pi^+$  contribution:

## ■ Electroweak radiative corrections:

▶ dominant contribution from short distance correction  $S_{EW}$  to effective 4-fermion coupling  $\propto (1 + 3\alpha(m_\tau)/4\pi)(1+2\langle Q \rangle)\log(M_Z/m_\tau)$

Marciano-Sirlin' 88

▶ subleading corrections calculated and small

Braaten-Li' 90

▶ long distance radiative correction  $G_{EM}(s)$  calculated  
[ add FSR to the bare cross section in order to obtain  $\pi^-\pi^+(\gamma)$  ]

Cirigliano-Ecker-Neufeld' 02

## ■ Charged/neutral mass splitting:

Alemany-Davier-Höcker' 97, Czyż-Kühn' 01

▶  $m_{\pi^-} \neq m_{\pi^0}$  leads to phase space (cross sec.) and width (FF) corrections

▶  $\rho$ - $\omega$  mixing (EM  $\omega \rightarrow \pi^-\pi^+$  decay) corrected using FF model

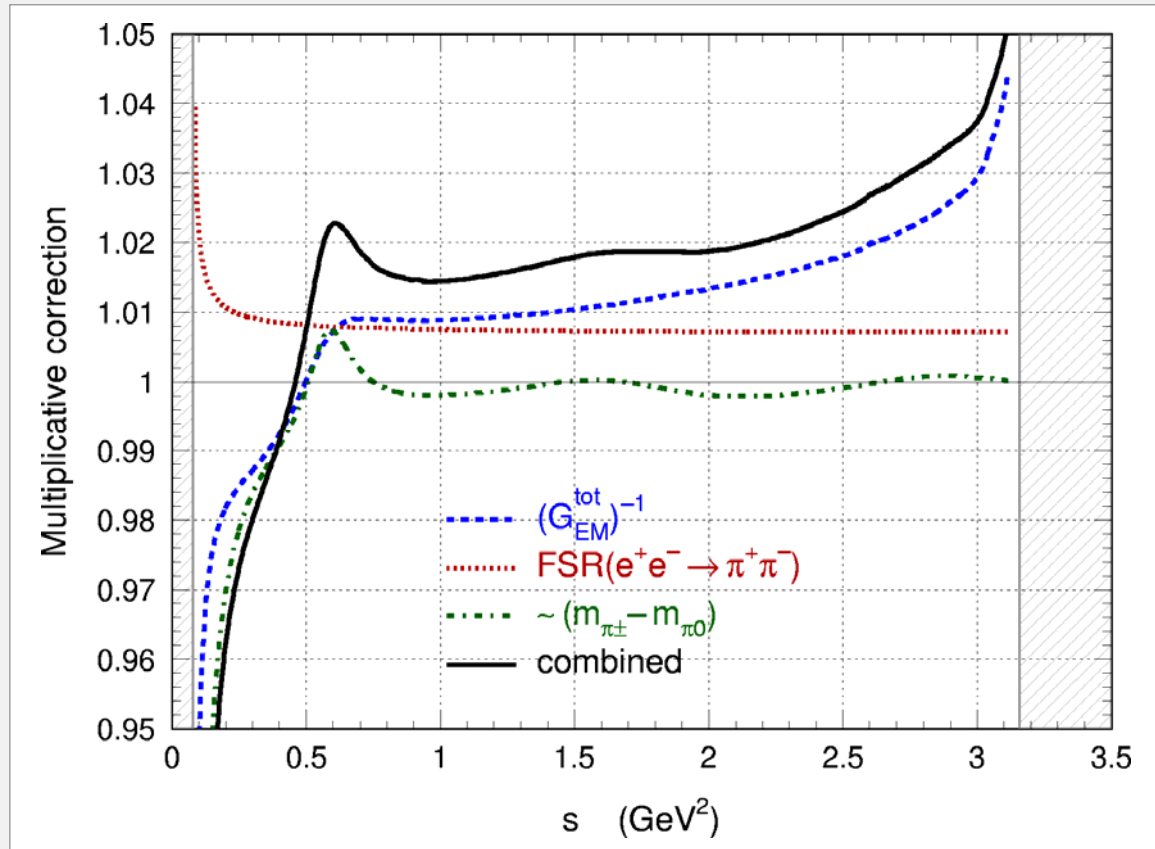
▶  $m_{\rho^-} \neq m_{\rho^0}$  and  $\Gamma_{\rho^-} \neq \Gamma_{\rho^0}$  [not corrected !]

■ Electromagnetic decays, like:  $\rho \rightarrow \pi\pi\gamma$ ,  $\rho \rightarrow \pi\gamma$ ,  $\rho \rightarrow \eta\gamma$ ,  $\rho \rightarrow l^+l^-$

■ Quark mass difference  $m_u \neq m_d$  generating "second class currents" (negligible)

# Mass Dependence of SU(2) Breaking

Multiplicative SU(2) corrections applied to  $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$  spectral function:



Only  $\beta^3$  and EW short-distance corrections applied to  $4\pi$  spectral functions



# $e^+e^-$ Radiative Corrections

Multiple radiative corrections are applied on measured  $e^+e^-$  cross sections

Situation often unclear: whether or not and if - which corrections were applied

- Vacuum polarization (VP) in the photon propagator:

- ▶ leptonic VP in general corrected for
- ▶ hadronic VP correction not applied, but for CMD-2 (in principle: iterative procedure)

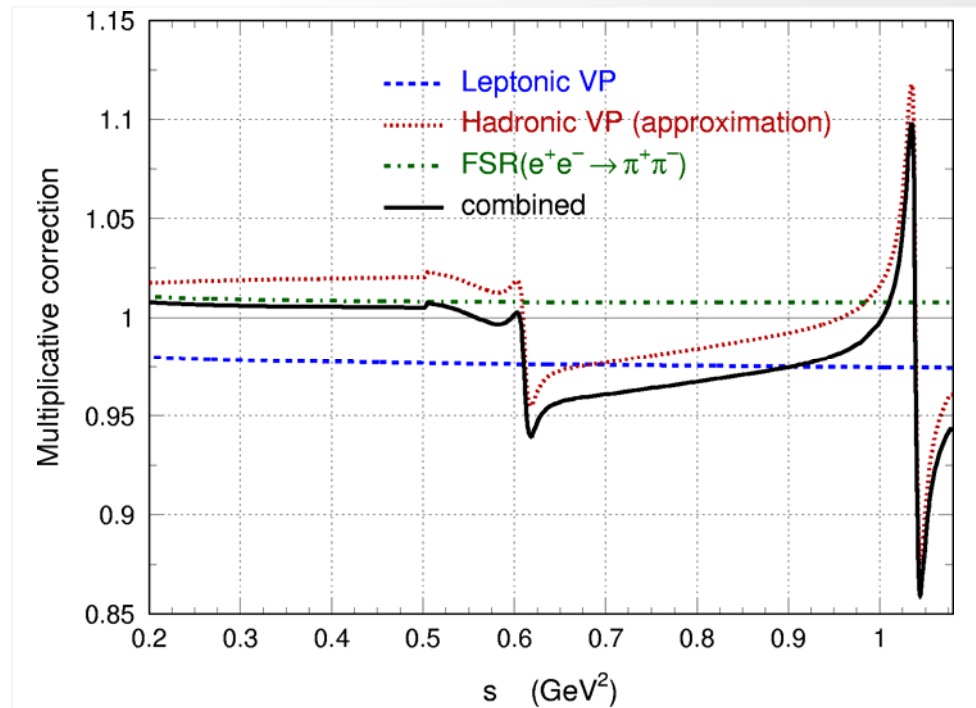
- Initial state radiation (ISR)

- ▶ corrected by experiments

- Final state radiation (FSR)

[need  $e^+e^- \rightarrow \text{hadrons} (\gamma)$  in disp. integral]

- ▶ usually, experiments obtain bare cross section so that FSR has to be added “by hand”; done for CMD-2, (supposedly) not done for others



# Inputs to the Analyses

Motivation for published 2002-2003 re-analyses:

- **New high precision  $e^+e^-$  results** (0.6% sys. error) around  $\rho$  from CMD-2 (Novosibirsk)
- **New  $\tau$  results** from ALEPH using full LEP1 statistics
- **New  $R$  results** from BES between 2 and 5 GeV
- **New theoretical analysis of SU(2) breaking**
- **discrepancy between  $e^+e^-$  and  $\tau$  input !**

CMD-2 PL B527, 161 (2002)

ALEPH Phys. Rep. 2005

BES PRL 84 594 (2000); PRL 88, 101802 (2002)

Cirigliano-Ecker-Neufeld  
JHEP 0208 (2002) 002

Davier-Eidelman-Höcker-Zhang  
Eur.Phys.J. C27 (2003) 497; C31 (2003) 503

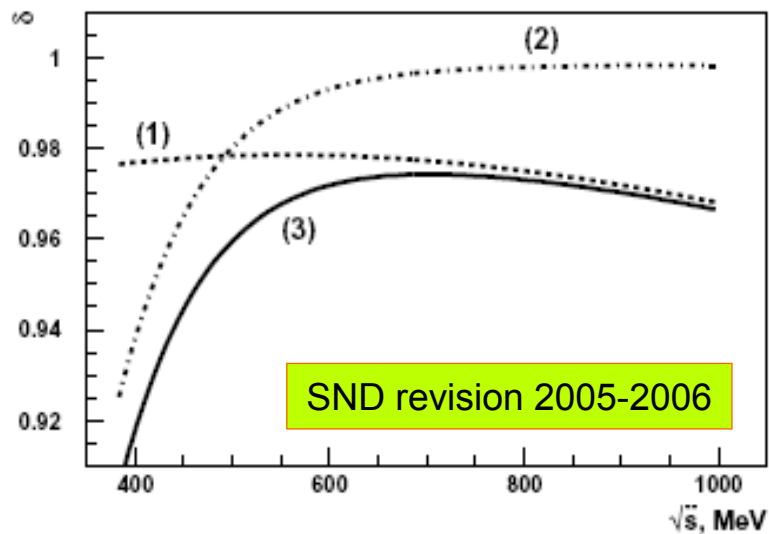
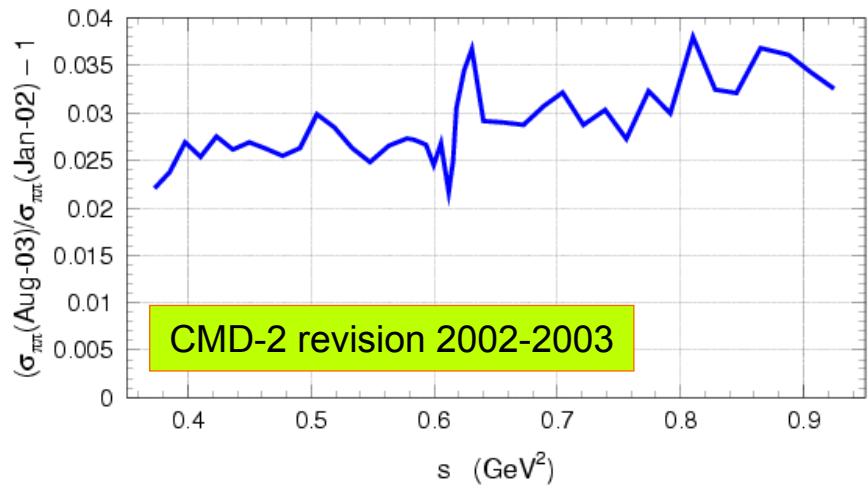
Hagiwara-Martin-Nomura-Teubner,  
Phys.Rev. D69 (2004) 093003 (no  $\tau$  data)

Jegerlehner,  
hep-ph/0312372 (no  $\tau$  data)

New input:

- **KLOE 2004, SND 2005** (revised 2006), **CMD-2 2006** (still preliminary)
- reduction of VP+FSR additional uncertainties
- **BABAR** exclusive channels  $2\pi\pi^0$ ,  $4\pi$ ,  $6\pi$ ,  $4\pi 2\pi^0$
- **QCD** in 1.8–3.7 GeV

# $e^+e^-$ Radiative Corrections Revised



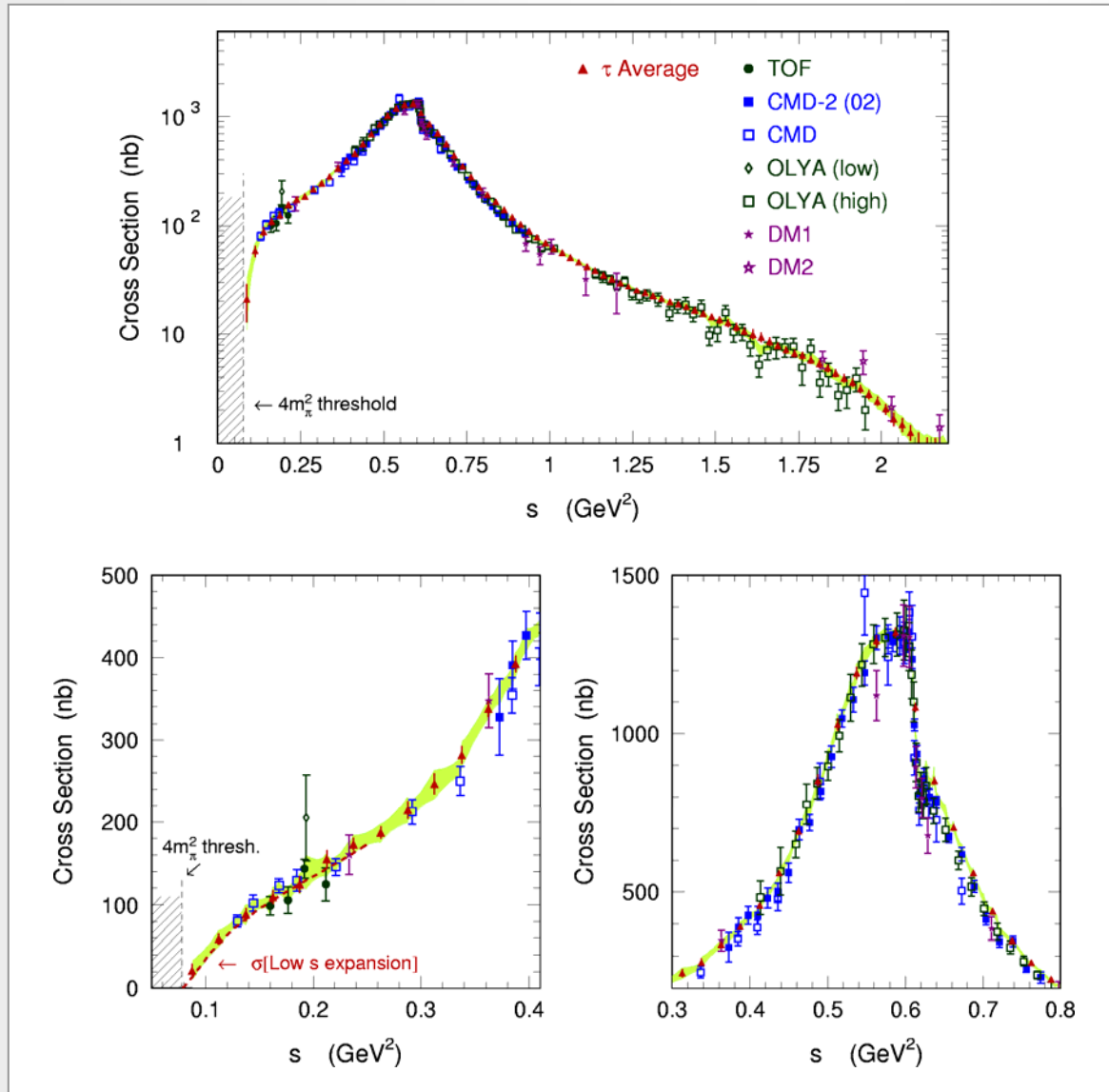
- ▶ 2.2-2.7% luminosity correction from change in  $\sigma_{\text{Bhabha}}$
- ▶ 1.2-1.4% change in  $\sigma_{\mu\mu}$
- ▶ Both changes affect  $ee / \mu\mu / \pi\pi$  separation

0.4% systematic error unchanged  
(0.6% total)

- ▶ Monte Carlo generators not correct for (1)  $\pi\pi(\gamma)$  and (2)  $\mu\mu(\gamma)$
- ▶ 3% change in  $\sigma_{\pi\pi}$  (3)

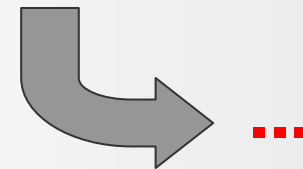
0.2% systematic error unchanged  
(1.3% total)

# Comparing $e^+e^- \rightarrow \pi^+\pi^-$ and $\tau \rightarrow \pi^-\pi^0\nu_\tau$



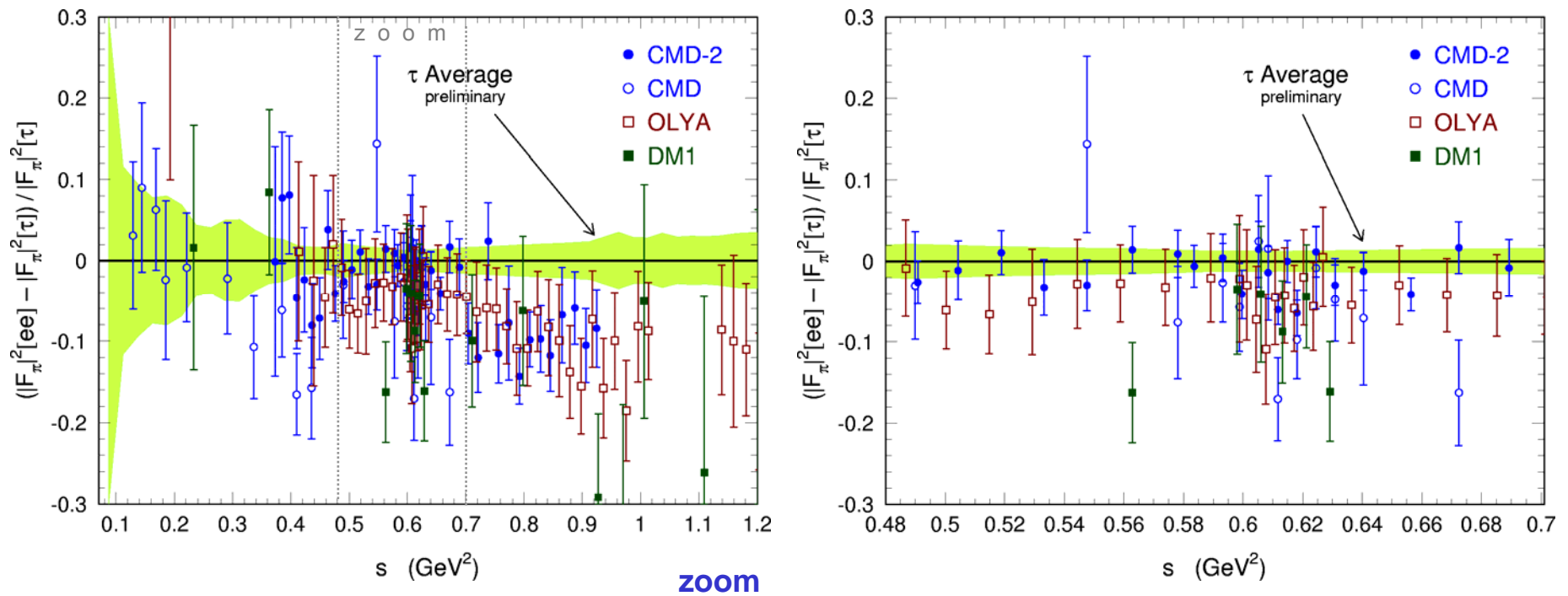
Correct  $\tau$  data for missing  $\rho$ - $\omega$  mixing (taken from BW fit) and all other SU(2)-breaking sources

Remarkable agreement  
But: not good enough...



# The $e^+e^- \tau$ Problem

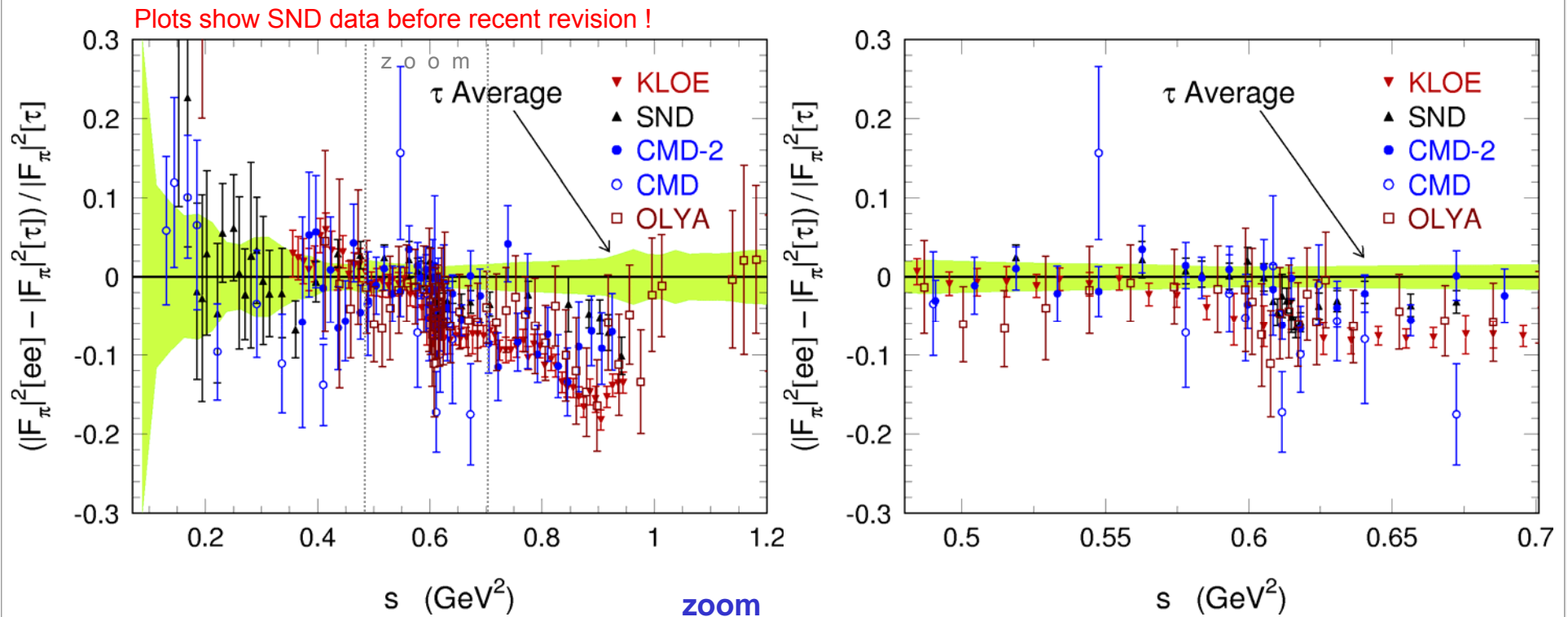
Relative difference between  $\tau$  and  $e^+e^-$  data:



zoom

# New $e^+e^- \rightarrow \pi^+\pi^-$ Data from KLOE (“radiative return“) & SND

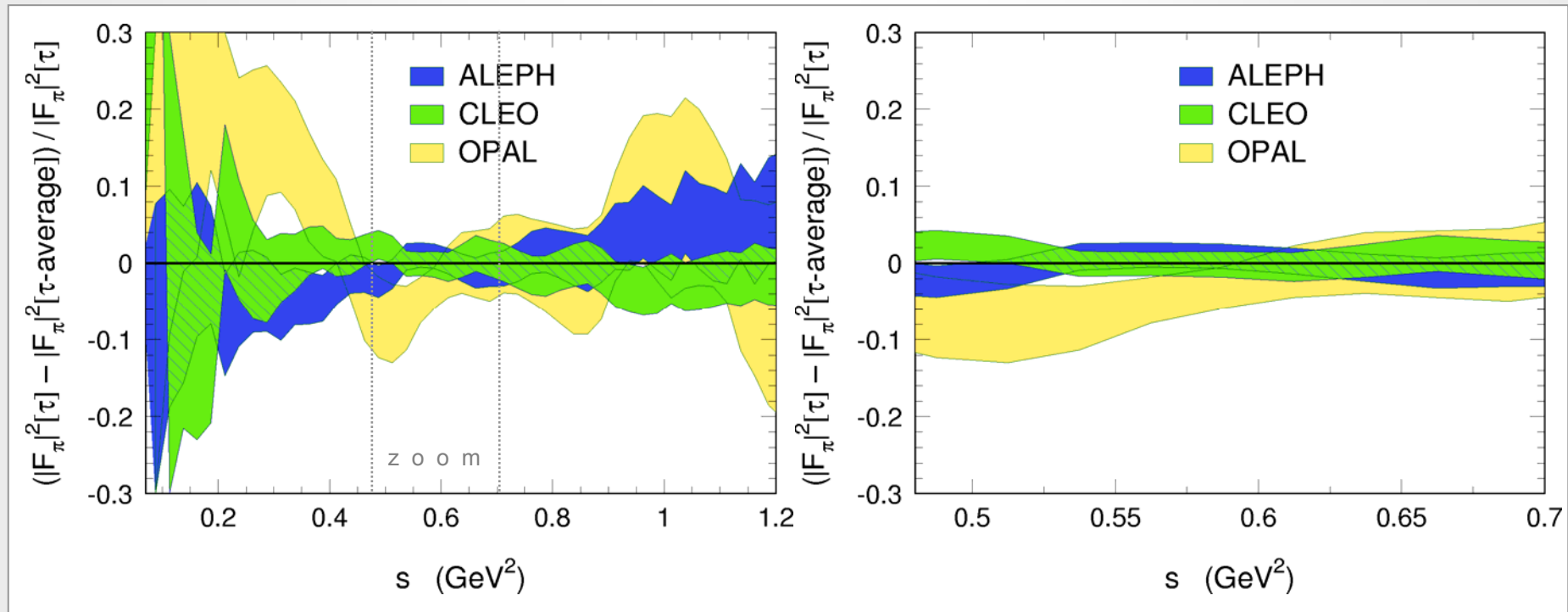
Relative difference between  $\tau$  and  $e^+e^-$  data:



SND,  
hep-ex/0506076 (2005)

KLOE,  
PL B606, 12 (2005)

# $\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$ : Comparing ALEPH, CLEO, OPAL



- ▶ Good agreement observed between ALEPH and CLEO
- ▶ ALEPH more precise at low  $s$
- ▶ CLEO better at high  $s$

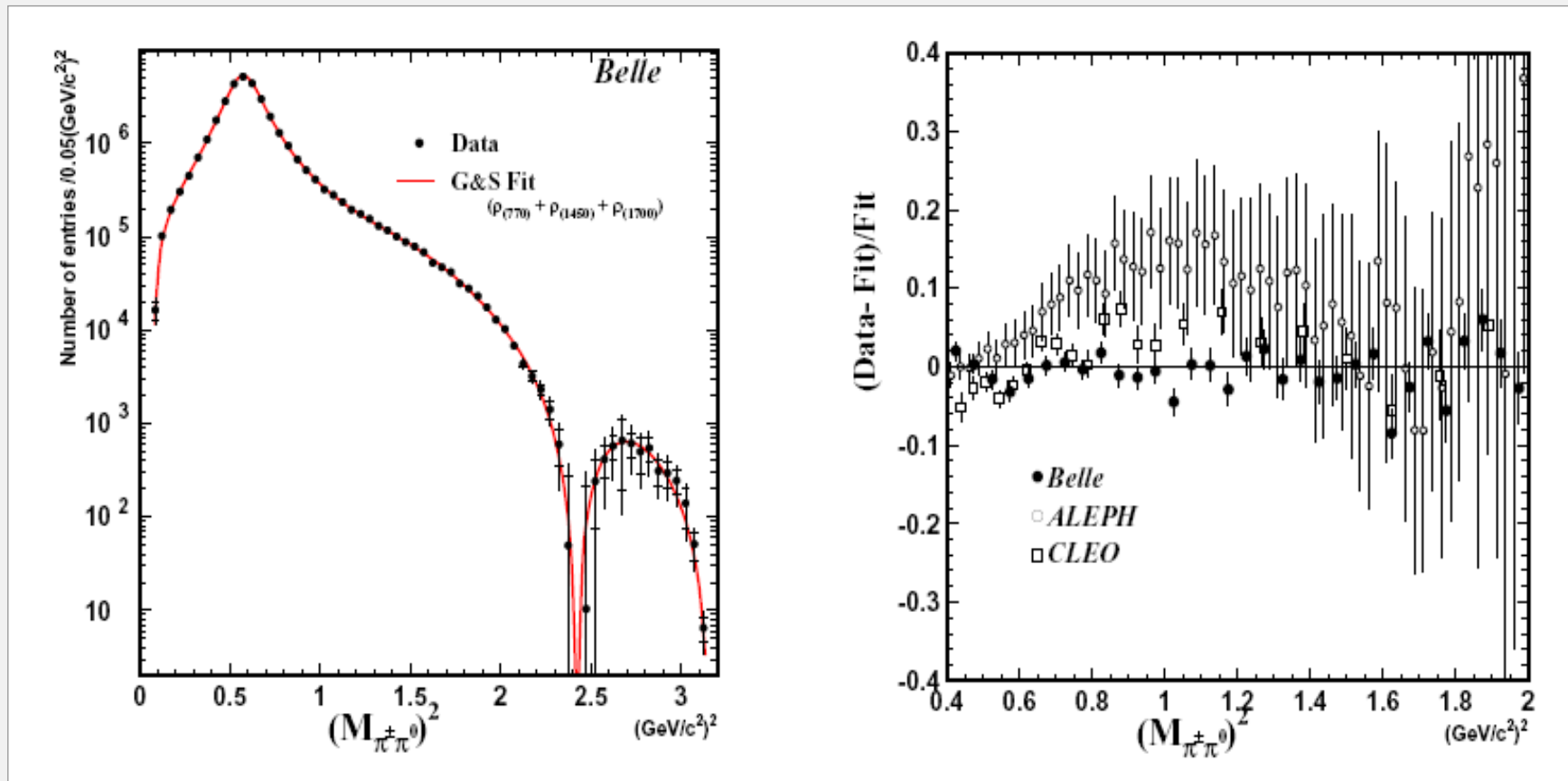
ALEPH,  
Phys. Rept. 421, 191 (2005)

CLEO,  
PRD 61, 112002 (2000)

OPAL,  
EPJ C35, 437 (2004)

# $\tau \rightarrow \pi^- \pi^0 \nu_\tau$ : Preliminary Results from Belle

- ▶ Preliminary spectral function presented by Belle at EPS 2005
- ▶ High statistics: see significant dip at 2.4 GeV<sup>2</sup> for first time in  $\tau$  data !
- ▶ Discrepancies with ALEPH/CLEO at large mass and with  $e^+e^-$  data at low mass

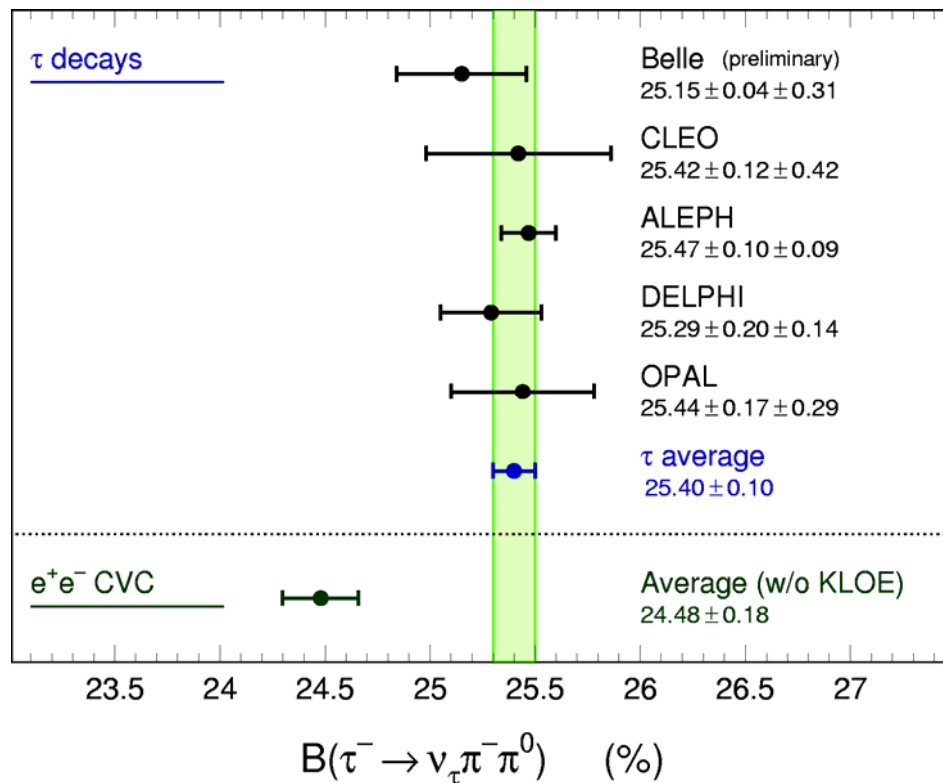




# Another Way to Look at the Data

Infer  $\tau$  branching fractions (more robust than spectral functions) from  $e^+e^-$  data:

$$\text{BR}_{\text{CVC}}(\tau^- \rightarrow \pi^- \pi^0 \nu_\tau) = \frac{6\pi |V_{ud}|^2 S_{EW}}{m_\tau^2} \int_0^{m_\tau^2} ds \text{kin}(s) \cdot v^{\text{SU(2)-corrected}}(s)$$



Difference:  $\text{BR}[\tau] - \text{BR}[e^+e^- (\text{CVC})]$ :

Mode	$\Delta(\tau - e^+e^-)$	“Sigma”
$\tau^- \rightarrow \pi^- \pi^0 \nu_\tau$	$+0.92 \pm 0.21$	4.5
$\tau^- \rightarrow \pi^- 3\pi^0 \nu_\tau$	$-0.08 \pm 0.11$	0.7
$\tau^- \rightarrow 2\pi^- \pi^+ \pi^0 \nu_\tau$	$+0.91 \pm 0.25$	3.6

$e^+e^-$  data on  $\pi^- \pi^+ \pi^0 \pi^0$  not satisfactory

# Final Remarks on Main $\pi^+\pi^-$ Contribution

The problem of the  $\pi^+\pi^-$  contribution:

- **Experimental situation:**

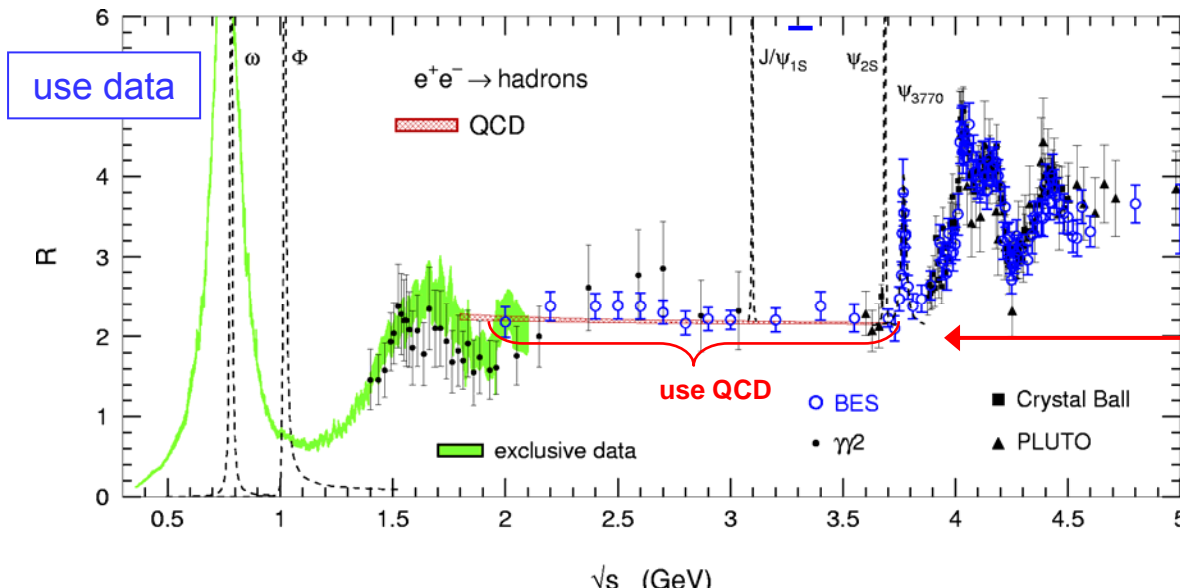
- ▶ revised SND results in agreement with CMD-2
- ▶  $\tau$  data without  $m(\rho)$  and  $\Gamma(\rho)$  corr. in strong disagreement with both data sets
- ▶ ALEPH, CLEO and OPAL  $\tau$  data in ok agreement, preliminary Belle less so
- ▶  $e^+e^-$  spectral functions have now reached the precision of  $\tau$  data

- **Concerning the remaining line shape discrepancy (0.7- 0.9 GeV<sup>2</sup>):**

- ▶ **SU(2) corrections:** basic contributions identified and stable since long; overall correction applied to  $\tau$  is  $(-2.2 \pm 0.5)\%$ , dominated by uncontroversial short distance piece; additional long-distance corrections found to be small
- ▶  **$\rho$  lineshape corrections** can improve the situation, but cannot account for the difference above 0.7 GeV<sup>2</sup>

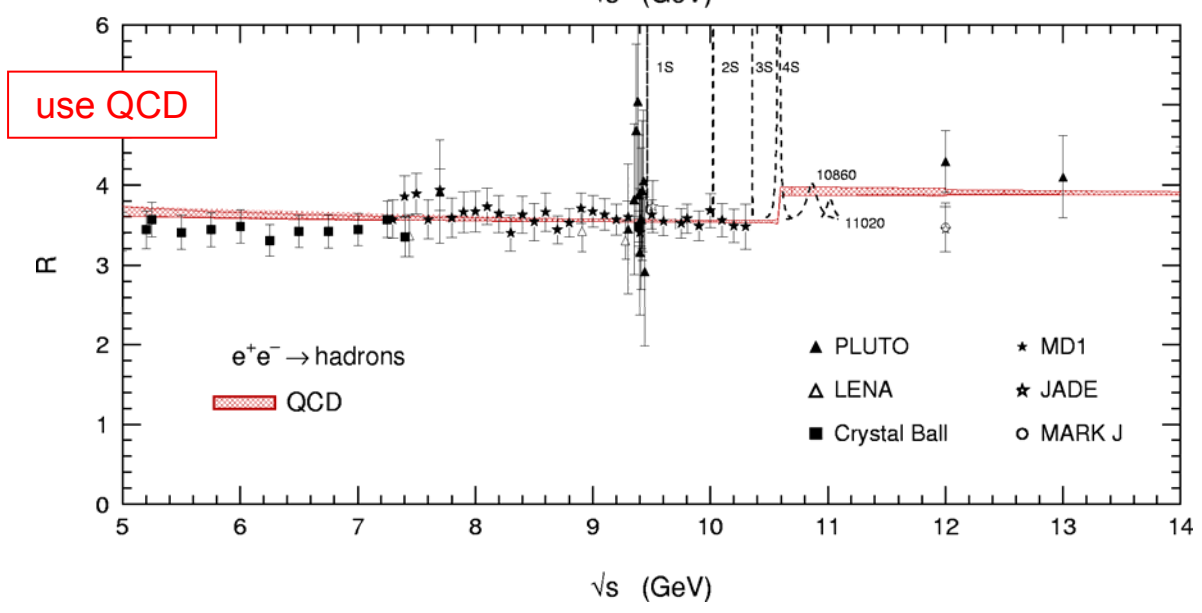
- ▶ The agreement between SND and CMD-2 invalidates the use of  $\tau$  data until a better understanding of the discrepancies is achieved (an interesting question as such)
- ▶ Discrepancy between KLOE and CMD-2/SND results: not safe to take advantage of decreased error when including KLOE

# Evaluating the Dispersion Integral



$$a_{\mu}^{\text{had,LO}} = \frac{\alpha^2}{3\pi^2} \int_{4m_{\pi}^2}^{\infty} ds \frac{K(s)}{s} R(s)$$

Agreement between Data (BES) and pQCD (within *correlated* systematic errors)

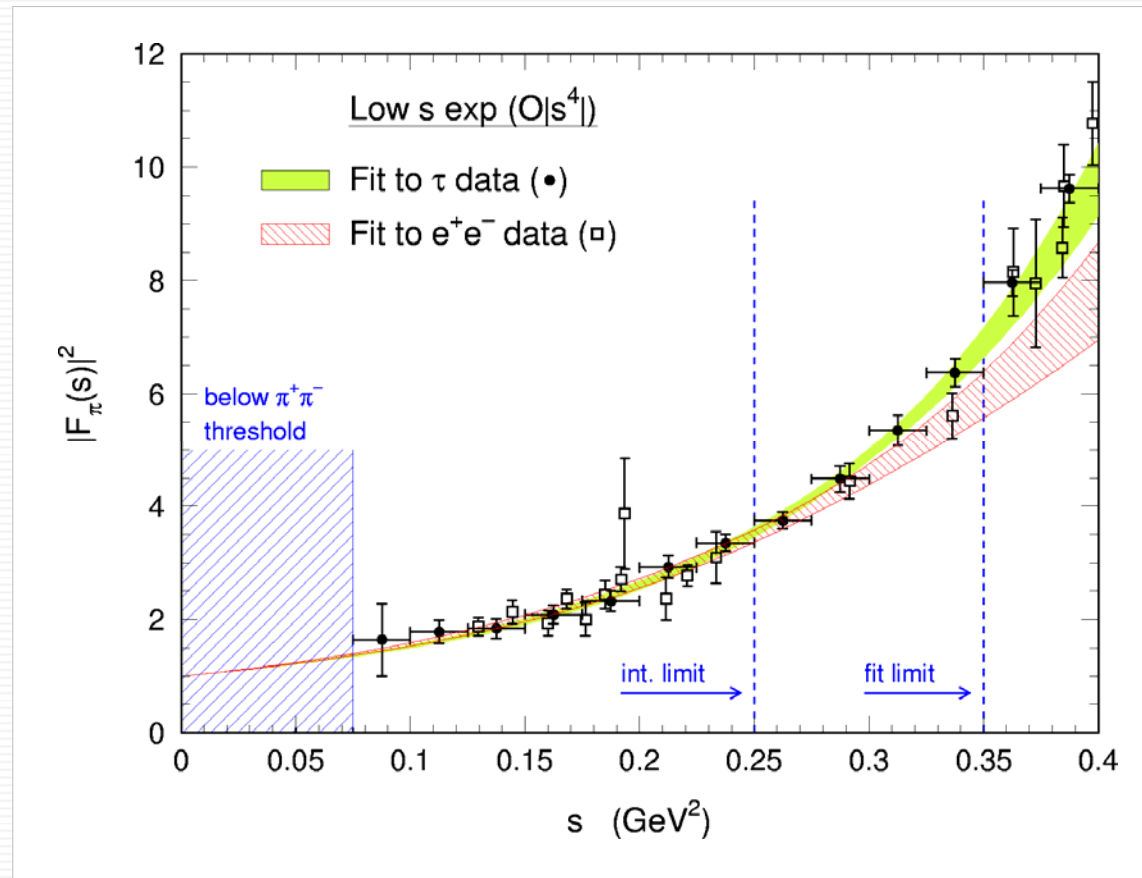


Better agreement between exclusive and inclusive ( $\gamma\gamma$ ) data than in 1997-1998 analyses

# *digression:* Specific Contributions: Threshold

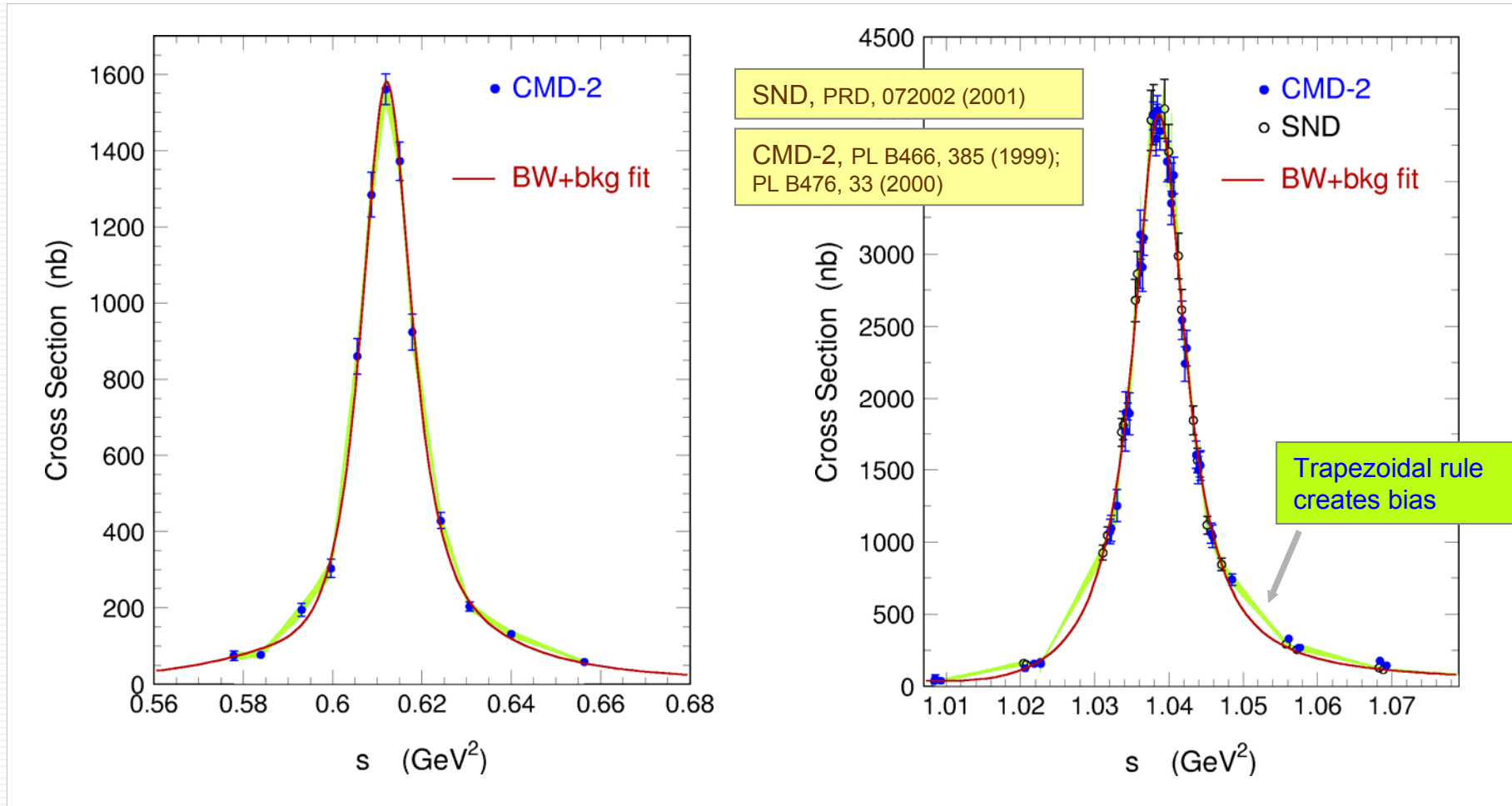
Use expansion for  $\pi^+\pi^-$  threshold inspired by chiral perturbation theory:

$$\sigma_{\pi\pi} = \frac{\pi\alpha^2\beta^3}{3s} |F_\pi|^2 \quad \text{and:} \quad F_\pi = 1 + \frac{1}{6} \langle r^2 \rangle_\pi s + c_1 s^2 + c_2 s^3 + O(s^4)$$



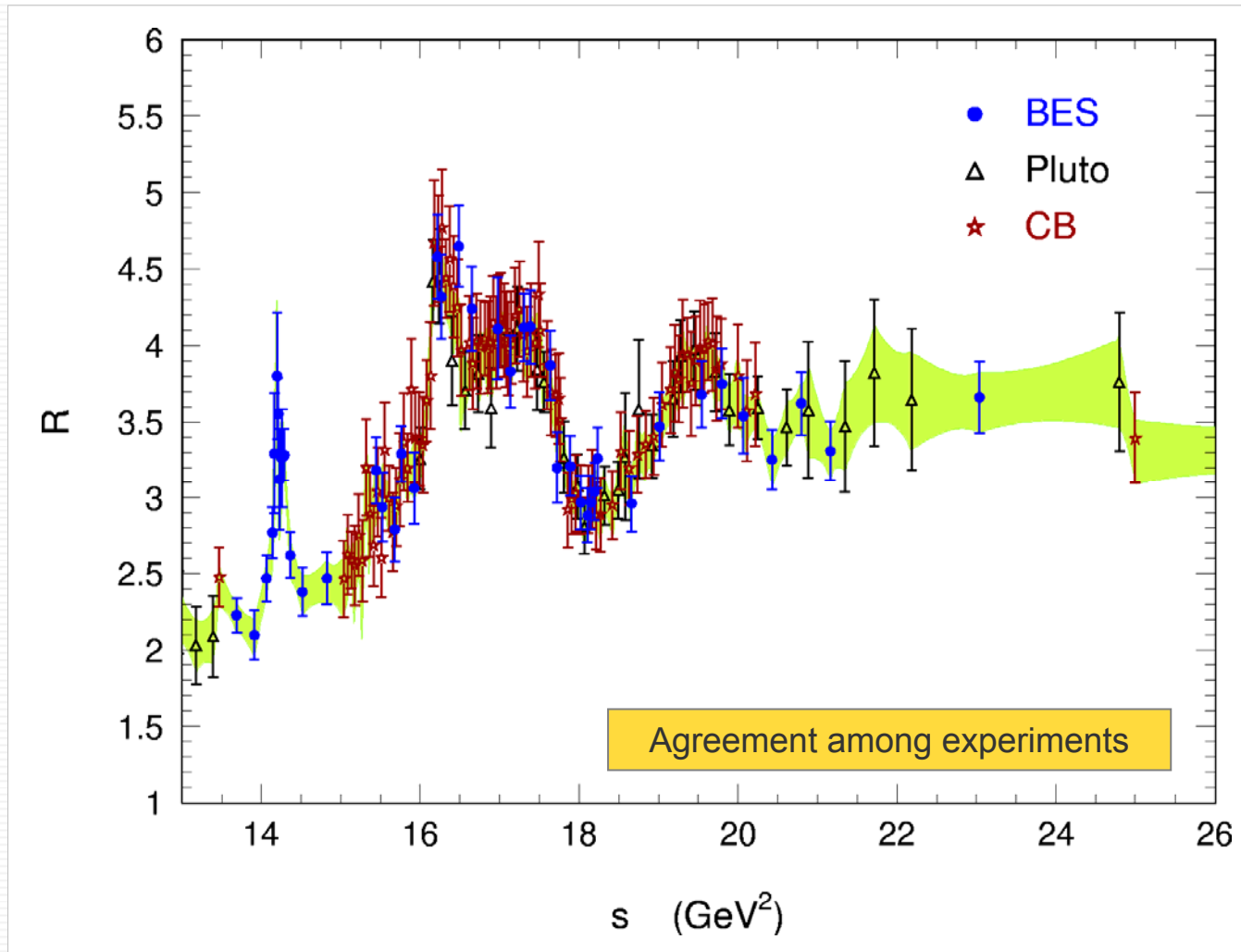
# *digression:* Specific Contributions: Narrow Light Resonances

Use direct data integration for  $\omega(782)$  and  $\phi(1020)$  to account for non-resonant contributions.



# *digression:* Specific Contributions: Charm Threshold

New precise BES data improve  $c\bar{c}$  resonance region:



# Results: the Compilation (including newest data)

Contributions to  $a_{\mu}^{\text{had,LO}}$  [in  $10^{-10}$ ] from the different energy domains:

Modes	Energy [GeV]	$e^+e^-$	$\tau$
<b>Low s expansion</b>	$2m_{\pi} - 0.5$	$55.6 \pm 0.8 \pm 0.1_{\text{rad}}$	$56.0 \pm 1.6 \pm 0.3_{\text{SU}(2)}$
$\pi^+\pi^-$ (+SND+CMD-2)	0.5 – 1.8	$449.0 \pm 3.0 \pm 0.9_{\text{rad}}$	$464.0 \pm 3.0 \pm 2.3_{\text{SU}(2)}$
$\pi^+\pi^-2\pi^0$	$2m_{\pi} - 1.8$	$16.8 \pm 1.3 \pm 0.2_{\text{rad}}$	$21.4 \pm 1.3 \pm 0.6_{\text{SU}(2)}$
$2\pi^+2\pi^-$ (+BABAR)	$2m_{\pi} - 1.8$	$13.1 \pm 0.4 \pm 0.0_{\text{rad}}$	$12.3 \pm 1.0 \pm 0.4_{\text{SU}(2)}$
$\omega(782)$	0.3 – 0.81	$38.0 \pm 1.0 \pm 0.3_{\text{rad}}$	–
$\phi(1020)$	1.0 – 1.055	$35.7 \pm 0.8 \pm 0.2_{\text{rad}}$	–
<b>Other excl.</b> (+BABAR)	$2m_{\pi} - 1.8$	$24.3 \pm 1.3 \pm 0.2_{\text{rad}}$	–
$J/\psi, \psi(2S)$	3.08 – 3.11	$7.4 \pm 0.4 \pm 0.0_{\text{rad}}$	–
<b>R [QCD]</b>	1.8 – 3.7	$33.9 \pm 0.5_{\text{theo}}$	–
<b>R [data]</b>	3.7 – 5.0	$7.2 \pm 0.3 \pm 0.0_{\text{rad}}$	–
<b>R [QCD]</b>	5.0 – $\infty$	$9.9 \pm 0.2_{\text{theo}}$	–
<b>Sum (w/o KLOE)</b>	$2m_{\pi} - \infty$	$690.8 \pm 3.9 \pm 1.9_{\text{rad}} \pm 0.7_{\text{QCD}}$	$710.1 \pm 5.0 \pm 0.7_{\text{rad}} \pm 2.8_{\text{SU}(2)}$

# The Full Hadronic Contribution

## ☀ Hadronic leading order contribution

$$a_{\mu}^{\text{had,LO}}[e^+e^-] = (693.4 \pm 5.3 \pm 3.5_{\text{rad}}) \times 10^{-10}$$

$$\text{Electroweak: } (15.4 \pm 0.2) \times 10^{-10}$$

$$\text{QED: } (11\,658\,471.9 \pm 0.1) \times 10^{-10}$$

## ☀ Hadronic next-to-leading order contributions

### ■ Vacuum polarization (1-loop) + additional photon or VP insertion:

🖼 computed akin to LO part via dispersion integral with modified kernel function

$$a_{\mu}^{\text{had,NLO}} = -9.8(0.1) \times 10^{-10}$$

### ■ Light-by-light scattering :

🖼 dispersion relation approach not possible (4-point function)

🖼 no first-principle calculation yet (e.g., on the lattice)

🖼 model calculations using short dist. quark loops,  $\pi^0$ ,  $\eta^{(\prime)}$ , ... pole insertions and  $\pi^{\pm}$  loops in the large- $N_C$  limit

$$a_{\mu}^{\text{had,LBL}} = +12.0(3.5) \times 10^{-10}$$

Knecht-Nyffeler, Phys.Rev.Lett. 88 (2002) 071802

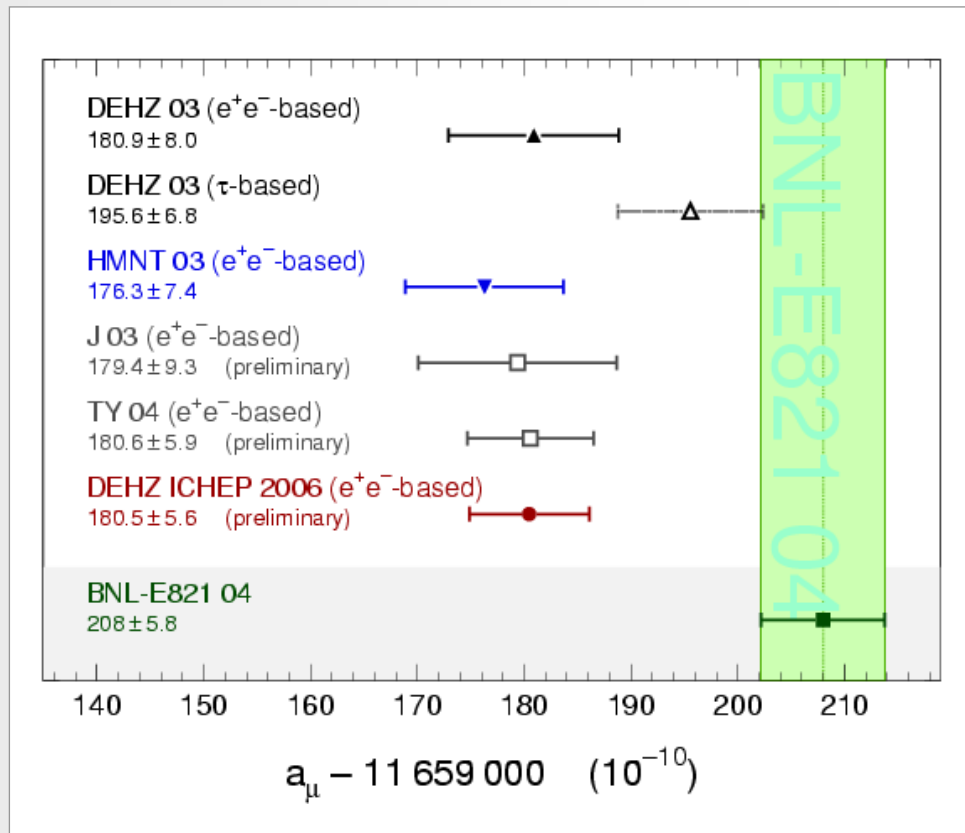
Melnikov-Vainshtein, PRD 70, 113006 (2004)



# And the Complete Result

$$a_{\mu}^{\text{SM}}[e^{+}e^{-}] = (11\,659\,180.5 \pm 4.4_{\text{had,LO}} \pm 3.5_{\text{LBL}} \pm 0.2_{\text{QED+weak}}) \times 10^{-10}$$

DEHZ (Tau 2006)



BNL E821 (2004):

$$a_{\mu}^{\text{exp}} = (11\,659\,208.0 \pm 6.3) 10^{-10}$$

Observed Difference with Experiment:

$$a_{\mu}^{\text{exp}} - a_{\mu}^{\text{SM}} = (27.5 \pm 8.4) \times 10^{-10}$$

➔ 3.3 "standard deviations"

# conclusions

## and perspectives

# conclusions

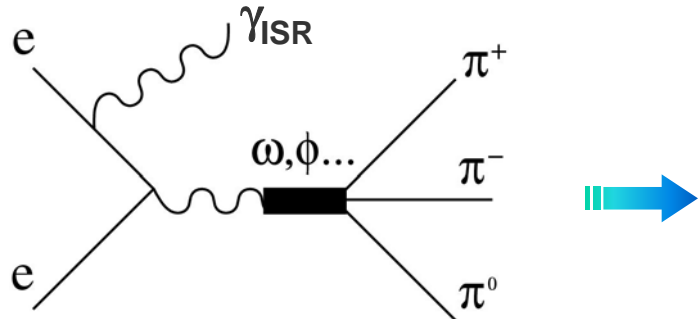
- ▶ Phenomenal experimental progress from BNL (E821)  $g-2$  measurement
- ▶ Improved theory prediction due to new CMD-2 and SND data
- ▶ Hadr. part dominates SM uncertainty (5.6), but more precise than experiment (6.3)
- ▶ Disagreement between SND/CMD-2 and KLOE data sets; so far KLOE not incl.
- ▶ Tau data in agreement (but Belle); revised SND data confirm  $\tau / e^+e^-$  discrepancy
- ▶ What is behind the  $4.5\sigma$   $\tau / e^+e^-$  discrepancy of the CVC BR ?
- ▶ KLOE will publish cross sections based on pion/muon ratios
- ▶ BABAR ISR:  $\pi^+\pi^-$  spectral function over full mass range, multihadron channels
- ▶ Difference between experiment and  $SM_{[e^+e^-]}$  within range of possible New Physics

# appendix

## the BABAR ISR programme

# Radiative Return Cross Section Results from BABAR

- ✦ The *Radiative Return*: benefit from huge luminosities at  $B$  and  $\phi$  Factories to perform continuous cross section measurements



$$\frac{d\sigma(s, x)}{dx d(\cos \theta)} = H(s, x, \theta) \cdot \sigma_0(s(1-x))$$

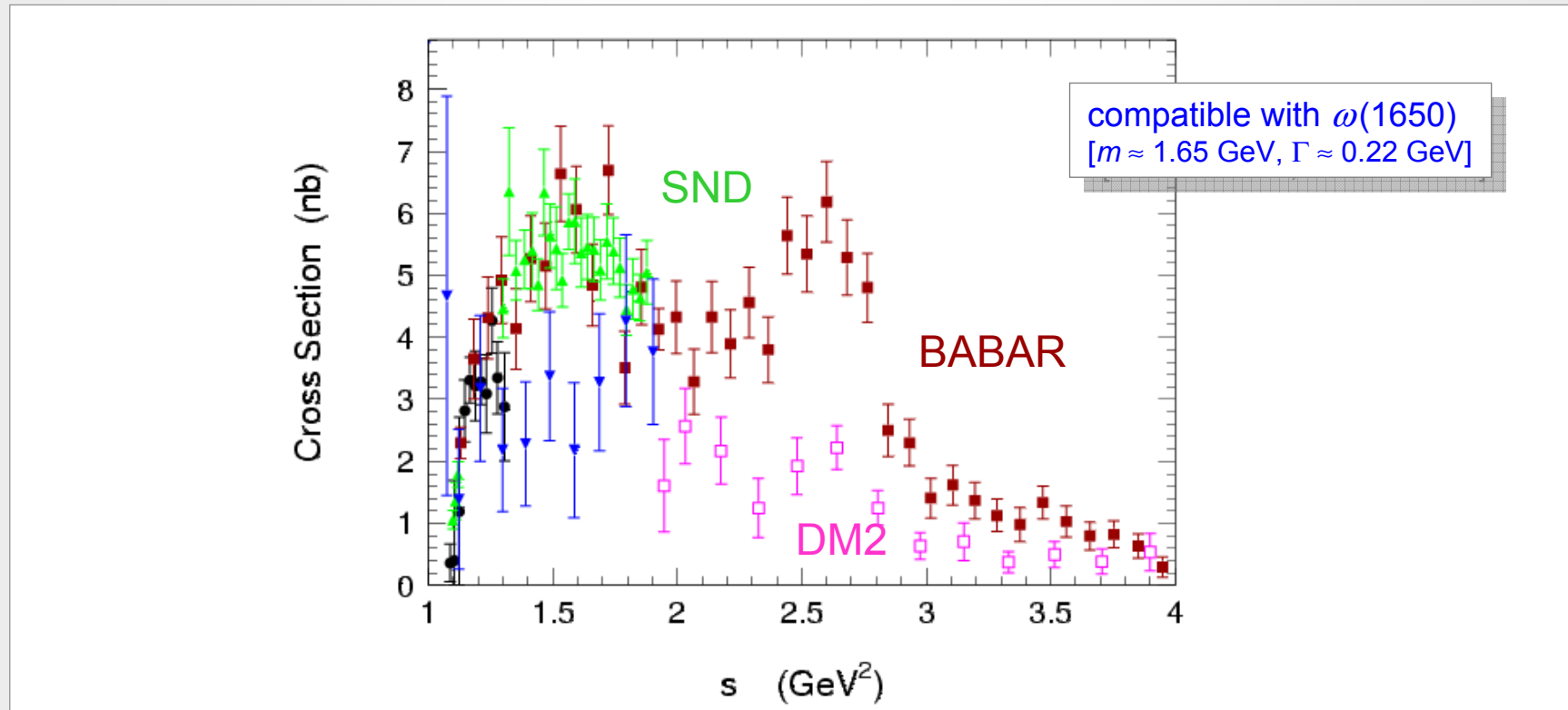
$$H(s, x, \theta) = \frac{\alpha}{\pi x} \left( \frac{2 - 2x + x^2}{\sin^2 \theta} - \frac{x^2}{2} \right), \quad x = \frac{2E_\gamma}{\sqrt{s}}$$

$H$  is radiation function

- High PEP-II luminosity at  $\sqrt{s} = 10.58$  GeV  $\rightarrow$  precise measurement of the  $e^+e^-$  cross section  $\sigma_0$  at **low c.m. energies** with BABAR
- Comprehensive program at BABAR
- Results for  $\pi^+\pi^-\pi^0$ , preliminary results for  $2\pi^+2\pi^-$ ,  $K^+K^-\pi^+\pi^-$ ,  $2K^+2K^-$  from  $89.3 \text{ fb}^{-1}$

# The $e^+e^- \rightarrow \pi^+\pi^-\pi^0$ Cross Section

- ☀ Cross section above  $\phi$  resonance : DM2 missed a resonance !



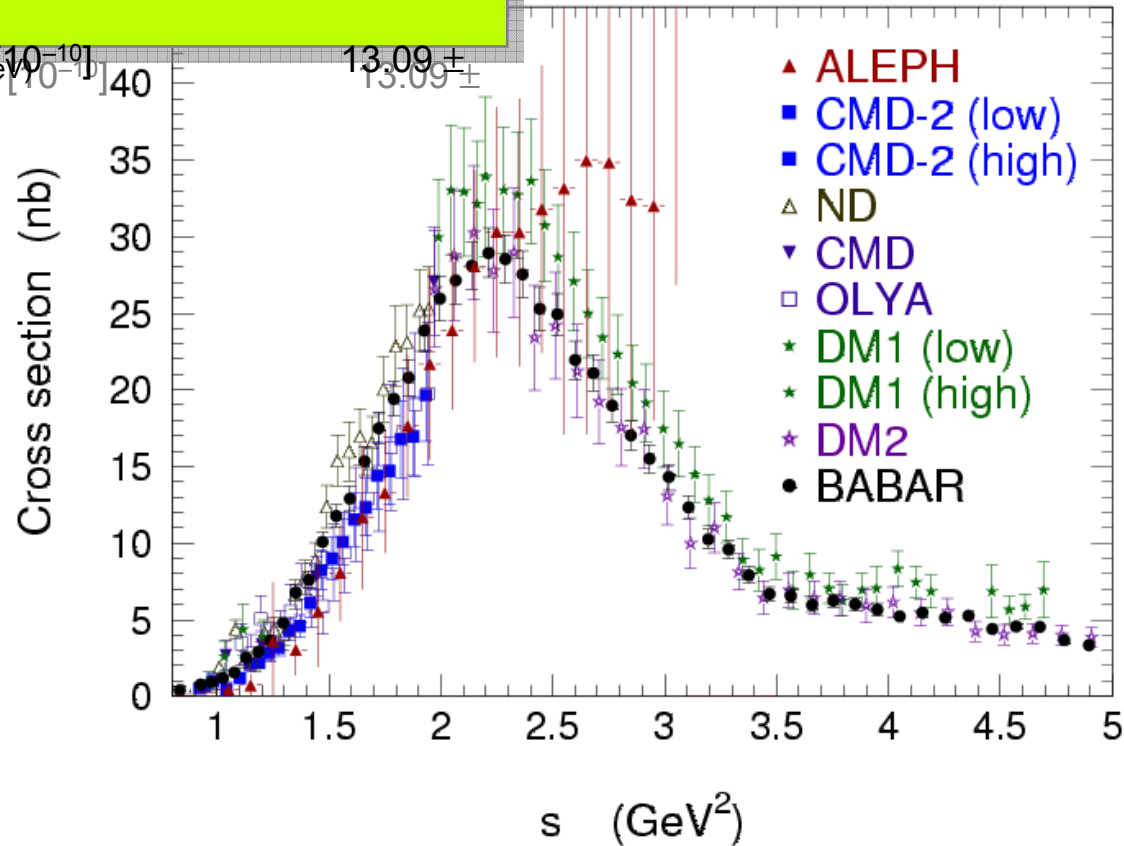
contribution to  $a_\mu^{\text{had}} (<1.8 \text{ GeV})$  :

- all before BABAR [ $10^{-10}$ ]  $2.45 \pm 0.26 \pm 0.03$
- all + BABAR [ $10^{-10}$ ]  $2.79 \pm 0.19 \pm 0.01$
- all + BABAR - DM2 [ $10^{-10}$ ]  $3.25 \pm 0.09 \pm 0.01$

# The $e^+e^- \rightarrow 2\pi^+2\pi^-$ Cross Section

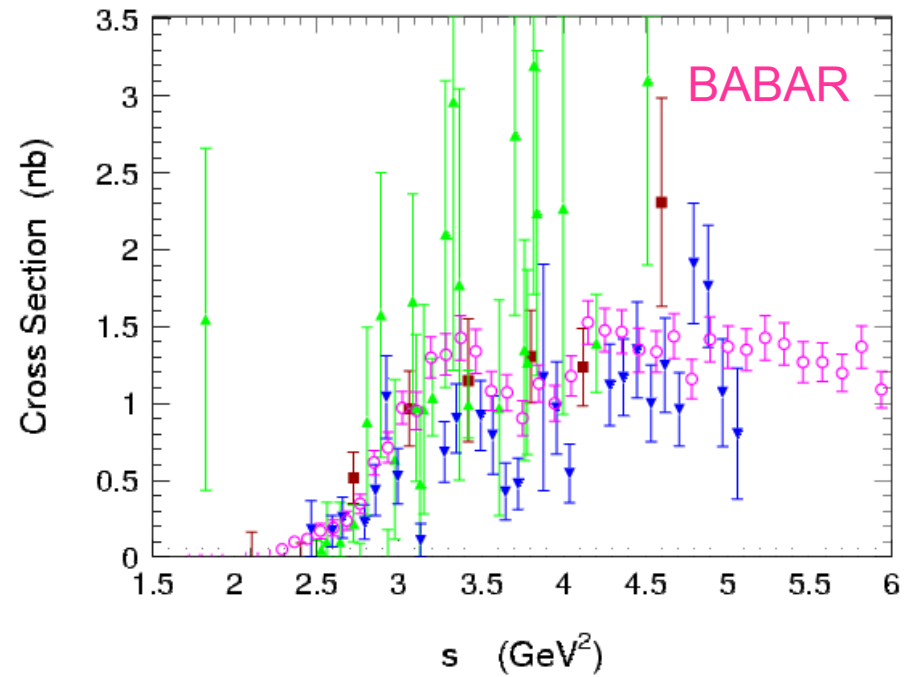
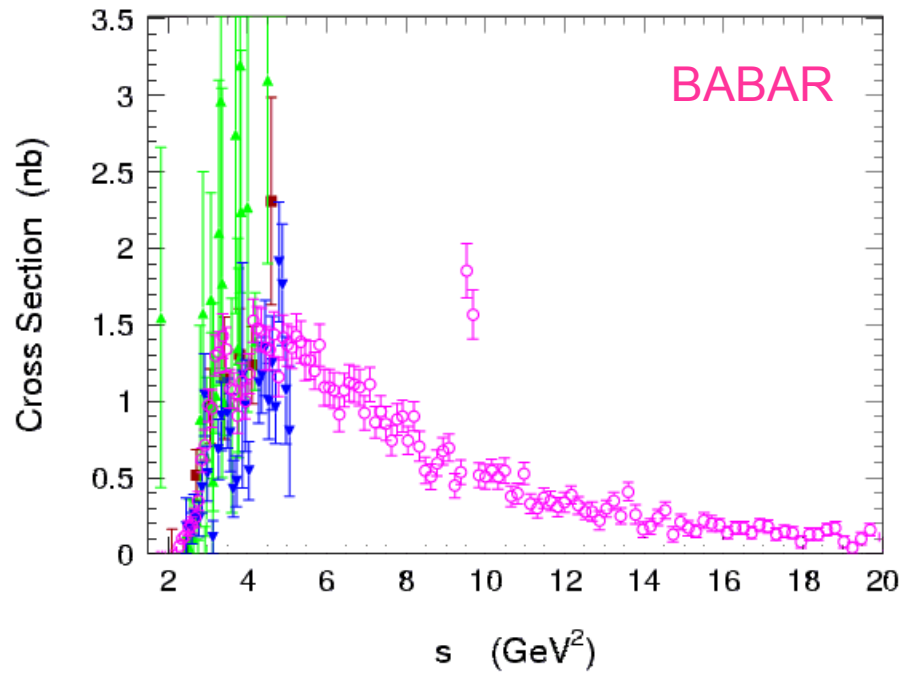
- all before BABAR [ $10^{-10}$ ]  $14.20 \pm 0.87 \pm 0.24$

- all + BABAR [ $10^{-10}$ ]  $13.09 \pm 0.44 \pm 0.00$



- Good agreement with direct  $e^+ e^-$  measurements
- Most precise result above 1.4 GeV

# BaBar ISR: $3\pi^+3\pi^-$

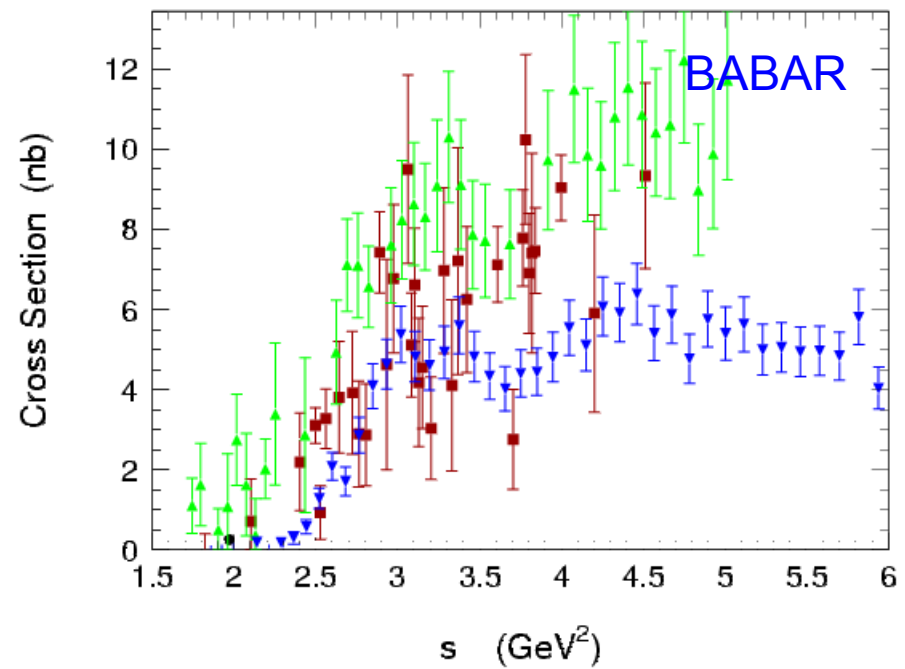
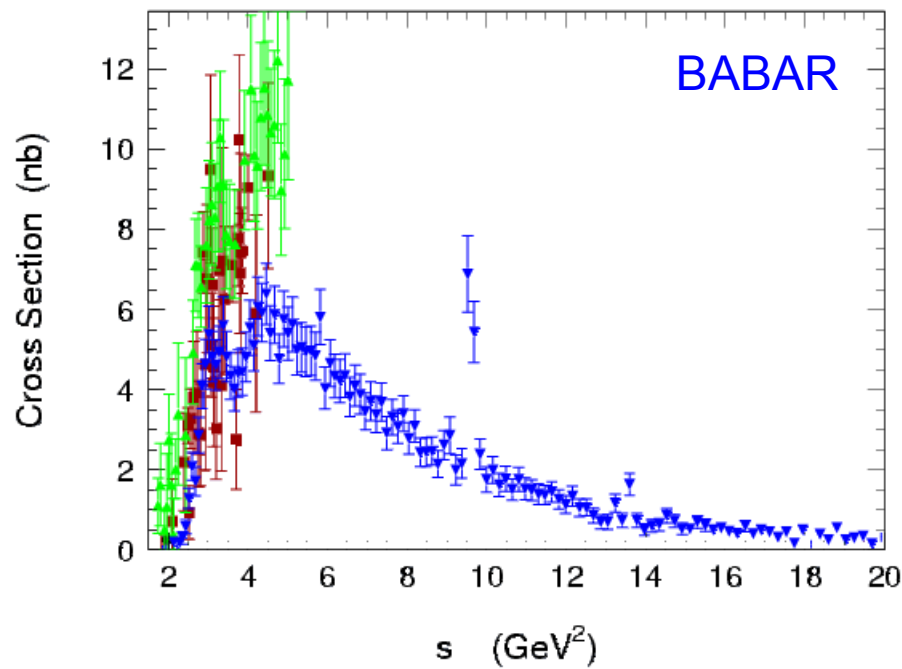


contribution to  $a_\mu^{\text{had}} (<1.8 \text{ GeV})$  :

- all before BABAR [ $10^{-10}$ ]  $0.10 \pm 0.10$
- all + BABAR [ $10^{-10}$ ]  $0.108 \pm 0.016$



# BABAR ISR: $2\pi^+2\pi^-2\pi^0$



contribution to  $a_\mu^{\text{had}} (<1.8 \text{ GeV})$  :

- all before BABAR [ $10^{-10}$ ]  $1.42 \pm 0.30 \pm 0.03$
- all + BABAR [ $10^{-10}$ ]  $0.89 \pm 0.09$