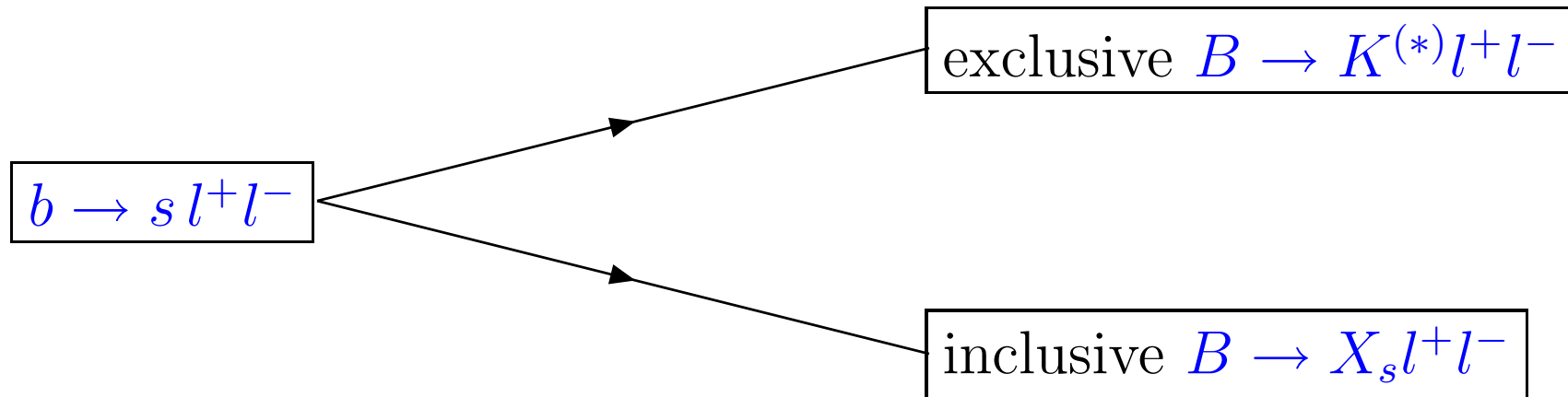


Charm contributions to $b \rightarrow sl^+l^-$ transitions

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- a very promising FCNC transition,
observables: $q^2 = (p_{l^+} + p_{l^-})^2$ distribution, CP- , forward-backward asymm.,...
- exclusive channels: measured by BABAR, Belle, good perspectives for LHC, hadronic matrix elements (form factors): effective theories, LCSR, lattice QCD
- inclusive: more difficult to measure,
theory: OPE, (heavy quark expansion) at NNLO

The problem of $c\bar{c}$ states in $b \rightarrow sl^+l^-$

* a chain of “routine” electroweak decays mimics $b \rightarrow sl^+l^-$, e.g.,

$$\begin{aligned} \text{BR}(B^0 \rightarrow J/\psi K^0) \times \text{BR}(J/\psi \rightarrow l^+l^-) &= \\ = [(8.5 \pm 0.5) \times 10^{-4}] \times 5.9\% &\simeq 5 \times 10^{-5} \end{aligned}$$

$$\begin{aligned} \text{BR}(B^0 \rightarrow \psi(2S)K^0) \times \text{BR}(\psi(2S) \rightarrow l^+l^-) &= \\ = [(6.2 \pm 0.7) \times 10^{-4}] \times [7.4 \times 10^{-3}] &\simeq 4.5 \times 10^{-6} \end{aligned}$$

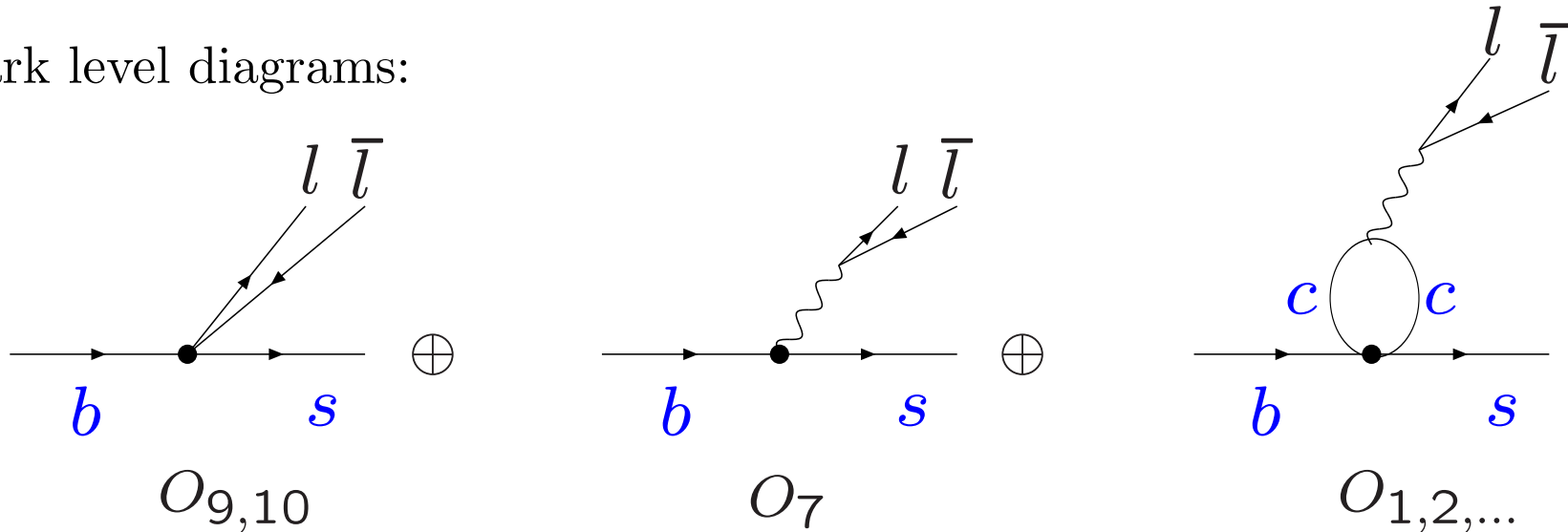
* “direct” decay via $b \rightarrow sl^+l^-$ transition
(subtracting the intervals of q^2 near J/ψ and $\psi(2S)$)

$$\text{BR}(B^0 \rightarrow K^0 l^+ l^-) = (3.12 \pm 1.0) \times 10^{-7} \quad [\text{HFAG average, August 2006}]$$

how large is the effect of $c\bar{c}$ states outside $J/\psi, \psi(2S)$ regions ?

“Anatomy” of $b \rightarrow sl^+l^-$ transitions

Quark level diagrams:



Effective Hamiltonian:
$$H_{eff} = -\frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_{i=1}^{10} C_i O_i$$

$$O_{9(10)} = (\bar{s}b)_{V-A}(\bar{l}l)_{V(A)}, \quad O_7 = (\bar{s}\Sigma_{\mu\nu}b)F^{\mu\nu}, \quad O_1^c = (\bar{s}c)_{V-A}(\bar{c}b)_{V-A}, \quad O_2^c, \dots$$

Effects of 4-quark $O_{1,\dots,6}$ operators in $b \rightarrow sl^+l^-$

a simplified setup:

- neglecting u -quark loops from $O_{1,2}^u$ (CKM suppressed for $b \rightarrow s$),
caution: light-quark resonances at small masses, $B \rightarrow \pi(\rho)l^+l^-$

- neglecting $O_{3,4,5,6}$ with small Wilson coeff. $C_{3,4,5,6}$
include them later, analogous to $O_{1,2}$

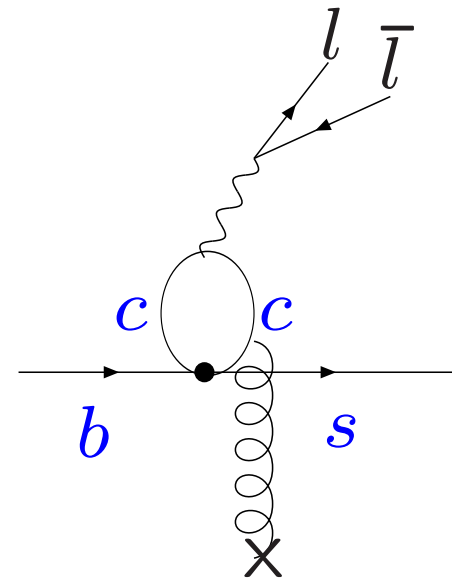
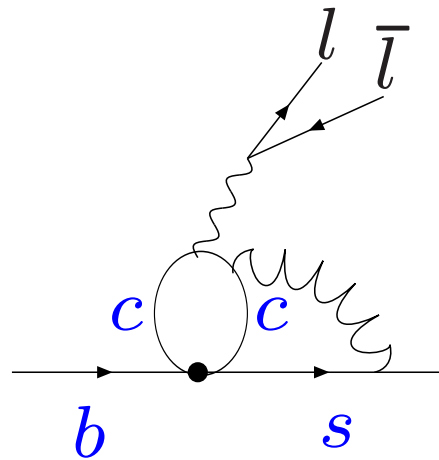
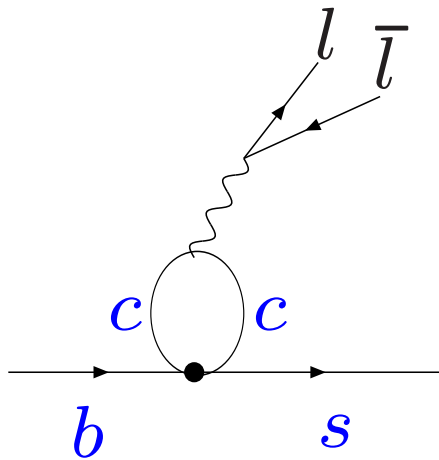
- main effect from c -quark loops, due to

$$C_1 O_1^c + C_2 O_2^c = \left(\frac{C_1}{3} + C_2 \right) (\bar{c}c)_{V-A} (\bar{s}b)_{V-A} + 2C_1 (\bar{c}t^a c)_{V-A} (\bar{s}t^a b)_{V-A}$$

$$\text{at } \mu \sim m_b \quad C_1 \sim 1, C_2 \sim -0.3, \quad C_1/3 + C_2 \ll C_1$$

$c\bar{c}$ -loop effects in $b \rightarrow sl^+l^-$

$q^2 \ll 4m_c^2$, c -quarks form a highly virtual loop ($2m_c \gg \Lambda_{QCD}$)



LO simple loop,

NLO $O(\alpha_s)$,

soft gluons (nonfactoriz.)

factorizable to $C_9^{eff} O_9$

[Bobeth, Misiak, Urban(2000); Asatryan, Greub, Walker(2001)]

Short-distance OPE in $\alpha_s(2m_c), 1/m_c^2$

c -quark loop effects at $q^2 \ll 4m_c^2$ in terms of OPE

$$\begin{aligned}
 & i \int d^4x e^{iqx} T\{\bar{c}\gamma_\nu c(x)[C_1 O_1^c(0) + C_2 O_2^c(0)]\} \\
 &= [(C_1/3 + C_2)g(q^2, m_c^2) + C_1\alpha_s g'(q^2, m_c^2)](\bar{s}\gamma_\nu b) \\
 &+ C_1 h_{\nu\rho\alpha\beta}(q^2, m_c^2)(\bar{s}\Gamma_\rho g_s G_{\alpha\beta} b) + \dots
 \end{aligned}$$

- g, g', h, \dots -calculable Wilson coefficients, $h \sim 1/m_c^2$,

NLO worked out mainly for inclusive $B \rightarrow X_s \gamma$ decays

- new effective operator $\bar{s}g_s \Gamma_\rho G_{\alpha\beta} b$

[Voloshin (1997); A.K., Rückl, Stoll, Wyler(1997);

Ligeti,Randall,Wise;Grant,Morgan,Nussinov,Peccei(1997); Buchalla,Isidori,Rey (1998)]

- are there other important operators? is the local expansion efficient ?
(infinite tower of operators \rightarrow light-cone expansion)

Hadronic matrix elements in $B \rightarrow K^{(*)}l^+l^-$

quark/gluons \rightarrow hadrons

$$O_i \rightarrow \langle K^{(*)} | O_i | B \rangle$$

- $\langle K^{(*)} | \bar{s} \gamma_\nu b | B \rangle$ $B \rightarrow K^{(*)}$ form factors
accessible with lattice QCD, LCSR, QCDF, SCET etc., uncertainty 10-15 %

- $\langle K^{(*)} | \bar{s} \Gamma_\rho g_s G_{\alpha\beta} b | B \rangle$, estimates at $q^2 = 0$ (for $B \rightarrow K^* \gamma$):

-local QCD sum rules [A.K., G. Stoll, D. Wyler (1997)];

-LCSR [P. Ball, R. Zwicky (2006)], see talk by R. Zwicky

small $O(1\%)$ corrections with large $O(50\%)$ uncertainties

Increasing the lepton pair mass q^2

With the local OPE answer for $\bar{c}c$ - effects
valid at $q^2 \ll 4m_c^2 \sim 6 \text{ GeV}^2$:

$$(\bar{m}_c(m_c) = 1.25 \text{ GeV})$$

- how large is $q_{max}^2 < m_{J/\psi}^2 = 9.6 \text{ GeV}^2$
- are $q^2 > m_{\psi(2S)}^2 = 13.6 \text{ GeV}^2$ accessible ?

Models of “long-distance” contributions in $b \rightarrow sl^+l^-$

- use the LO $\bar{c}c$ loop everywhere (with Im-part above $4m_c^2$):
too crude for the q^2 -distribution, nonfactorizable effects ignored
- add a bunch of ψ resonances to the $\bar{c}c$ loop:
double counting of quark and hadronic d.o.f
- the method avoiding double counting of quark and hadron d.o.f.
[Krüger, Sehgal(1997) (for $B \rightarrow X_s l^+ l^-$)]
dispersion relation (analyticity, unitarity) with ψ resonances \oplus open charm
continuum matched to a LO $\bar{c}c$ loop at small q^2

but: nonfactorizable $O(\alpha_s), O(1/m_c^2)$ contributions not taken into account
(see discussion in [Buchalla, Isidori, 1998])

analogous method used in $K \rightarrow \pi l^+ l^-$ [Okun, Shifman, Vainshtein, Zakharov, (1976)]

Is the region $q^2 < m_{J/\psi}^2$ directly accessible with OPE ?

- series of effective soft-gluon operators diverge at $q^2 \rightarrow 4m_c^2$
- QCD sum rules for charmonium: the OPE is valid at $q^2 = 0$ and matches dispersion relation with ψ 's, but the gluon condensate effects grow very fast at $q^2 \rightarrow 4m_c^2$
- $B \rightarrow J/\psi K^{(*)}$, $B \rightarrow \psi(2S)K^{(*)}$ decay amplitudes extracted from the measured BR's cannot be approximated by factorization, nonfactorizable effects are at O(100%) level

Use of dispersion relation in J/ψ channel

- match OPE to dispersion relation at small q^2 , where **both** are valid:

$$\begin{aligned} & \langle K^{(*)} | i \int d^4x e^{iqx} T \{ c \gamma_\nu \bar{c}(x) (C_1 O_1^c(0) + C_2 O_2^c(0)) | B \rangle \\ &= \sum_{\psi=J/\psi, \dots} \frac{f_\psi A_\nu(B \rightarrow \psi K^{(*)})}{m_\psi^2 - q^2 - im_\psi \Gamma_\psi} + \text{open charm} \end{aligned}$$

- use dispersion relation to predict the q^2 distribution below $\psi(3S)$
- very preliminary crude estimate for $B \rightarrow Kl^+l^-$: (work in progress)

$$C_9^{eff, O_{1,2}}(q^2) = C_9 + \delta_c(q^2)$$

the charm effect not exceeding 5% (10%) of C_9 in at $q^2 \leq 7 \text{ GeV}^2$ ($\leq 8 \text{ GeV}^2$)

- a potential problem with q^2 above $\psi(3S)$:
broad resonances, dependence on the model for the widths

Summary and final comments

- local OPE provides a reliable framework for estimating $\bar{c}c$ effects at small q^2
more work needed to understand $(1/m_c^2)^n$ suppressed contributions
- matching with the dispersion relation helps to avoid double counting and unaccounted nonperturbative effects
- optimistically, charm effects in $B \rightarrow K^{(*)}l^+l^-$ add \leq few % to c_9 below J/ψ
- light-quark loops need a parallel/similar analysis for the region $q^2 > m_\rho^2, m_\phi^2$, (broad resonances?)
- not clear if the same analysis can be simply translated to $B \rightarrow X_s l^+ l^-$
(a specific problem of “hidden” background due to $J/\psi \rightarrow l^+ l^- X$ with a lepton-pair mass $< m_{J/\psi}$)