

# Weak Charm Decays with lattice QCD: status and prospects



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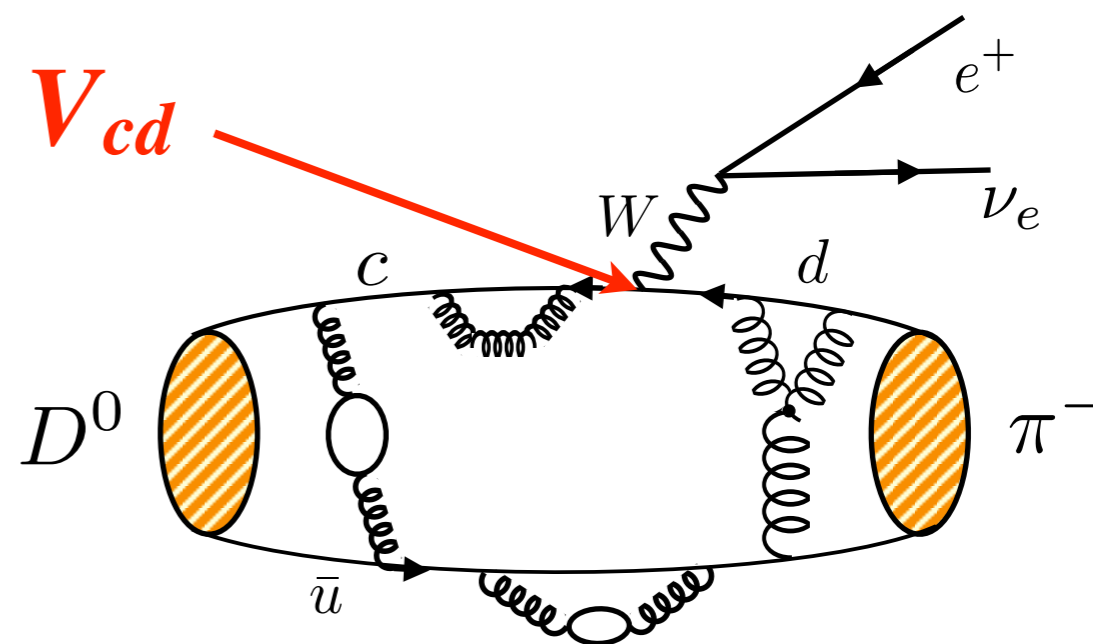
# Outline

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- Motivation and Introduction
  - ◆ lattice QCD
- Results
  - ◆ leptonic  $D$ -meson decays
  - ◆ semileptonic  $D$ -meson decays
  - ◆ neutral  $D$ -meson mixing
- Phenomenology
  - ◆ CKM determinations
  - ◆ second row unitarity test
  - ◆ NP scale
- Summary and Outlook

# Introduction

example:  $D \rightarrow \pi \ell \nu$



generic weak process involving hadrons:

$$(\text{experiment}) = (\text{known}) \times (\text{CKM element}) \times (\text{had. matrix element})$$

$$\Gamma_{K\ell 3}, \Gamma_{K\ell 2}, \dots$$

$$\frac{d\Gamma(B \rightarrow \pi \ell \nu)}{dq^2}, \frac{d\Gamma(D \rightarrow K \ell \nu)}{dq^2}, \dots$$

$$\Delta m_{d(s)}$$

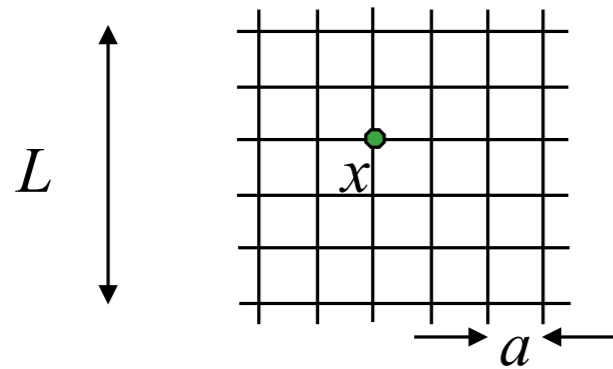
⋮

**Lattice QCD**

parameterize the ME in terms of form factors, decay constants, bag parameters, ...

# Lattice QCD Introduction

$$\mathcal{L}_{\text{QCD}} = \sum_f \bar{\psi}_f (\not{D} + m_f) \psi_f + \frac{1}{4} \text{tr} F_{\mu\nu} F^{\mu\nu}$$



- ◆ discrete Euclidean space-time (spacing  $a$ )  
derivatives  $\rightarrow$  difference operators, etc...
- ◆ finite spatial volume ( $L$ )
- ◆ finite time extent ( $T$ )

## adjustable parameters

❖ lattice spacing:

$$a \rightarrow 0$$



❖ finite volume, time:  $L \rightarrow \infty, T > L$



❖ quark masses ( $m_f$ ):

$$M_{H,\text{lat}} = M_{H,\text{exp}}$$



tune using hadron masses  
extrapolations/interpolations

$$m_f \rightarrow m_{f,\text{phys}}$$

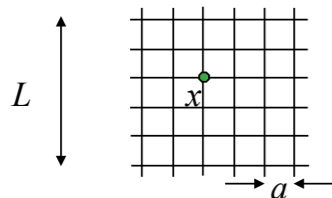
$m_{ud}$

$m_s$

$m_c$

$m_b$

❖ also:  $n_f =$  number of sea quarks: 3 (2+1), 4 (2+1+1)



# Lattice QCD Introduction

$$\langle \mathcal{O} \rangle \sim \int \mathcal{D}\psi \mathcal{D}\bar{\psi} \mathcal{D}A \mathcal{O}(\psi, \bar{\psi}, A) e^{-S}$$

$$S = \int d^4x \left[ \bar{\psi}(\not{D} + m)\psi + \frac{1}{4}(F_{\mu\nu}^a)^2 \right]$$

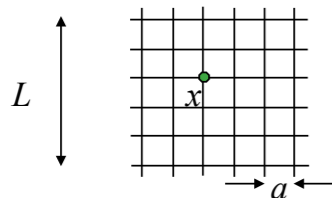
use monte carlo methods (importance sampling) to evaluate the integral.

**Note:** Integrating over the fermion fields leaves  $\det(\not{D} + m)$  in the integrand. The correlation functions,  $\mathcal{O}$ , are then written in terms of  $(\not{D} + m)^{-1}$  and gluon fields.

steps of a lattice QCD calculation:

1. generate gluon field configurations according to  $\det(\not{D} + m) e^{-S}$
2. calculate quark propagators,  $(\not{D} + m_q)^{-1}$ , for each valence quark flavor and source point
3. tie together quark propagators into hadronic correlation functions (usually 2 or 3-pt functions)
4. statistical analysis to extract hadron masses, energies, hadronic matrix elements, .... from correlation functions

**5. systematic error analysis**



# Lattice QCD Introduction

## systematic error analysis

...of lattice spacing, chiral, heavy quark, and finite volume effects is based on **EFT (Effective Field Theory)** descriptions of QCD

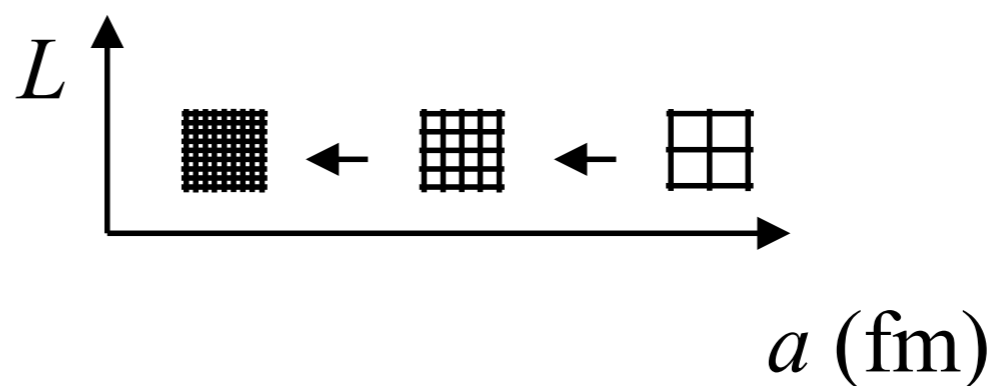
→ **ab initio**

The **EFT** description:

- provides functional form for extrapolation (or interpolation)
- can be used to build improved lattice actions/methods
- can be used to anticipate the size of systematic effects

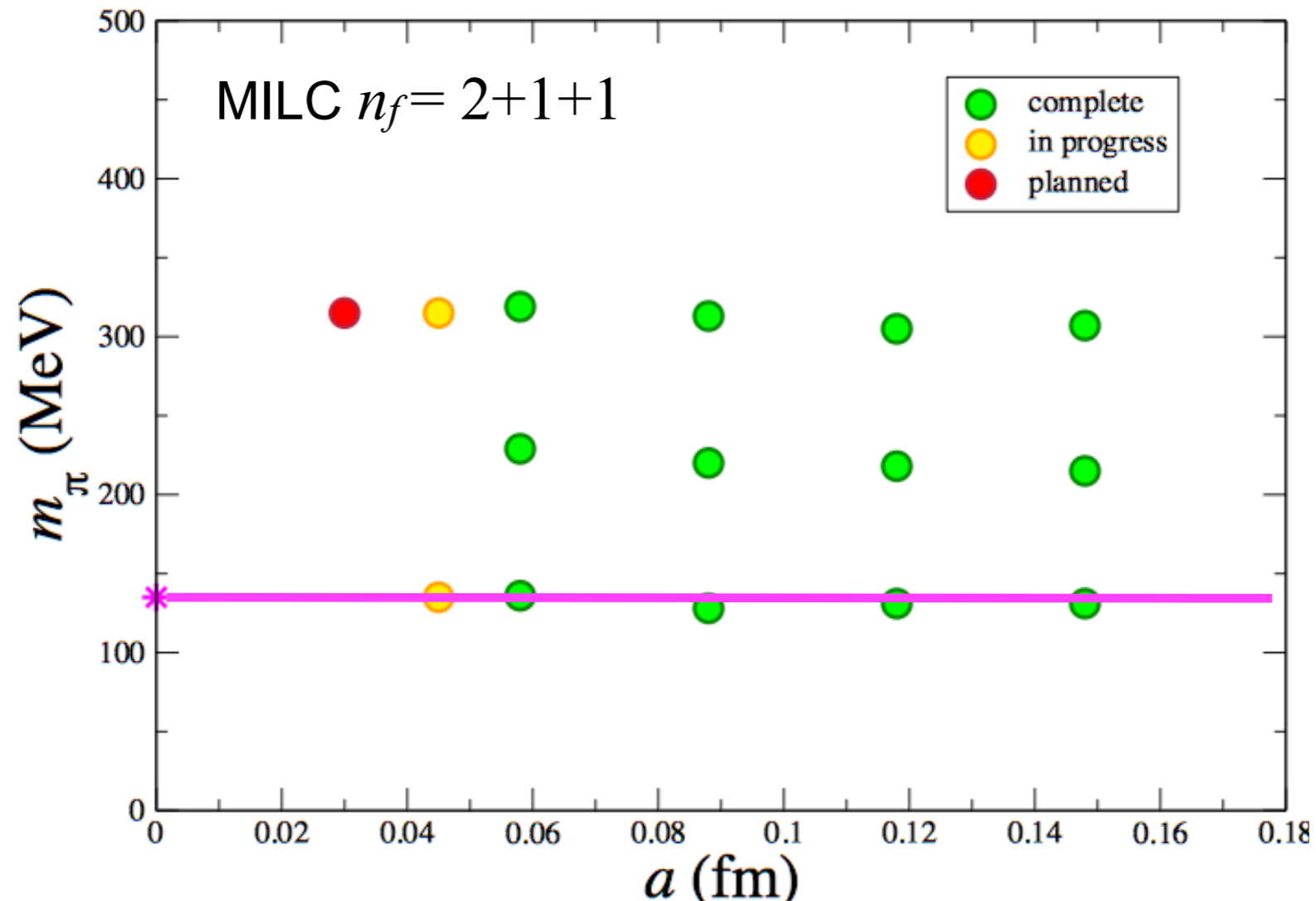
**To control and reliably estimate the systematic errors**

- repeat the calculation on several lattice spacings, light quark masses, spatial volumes, ...



# chiral-continuum extrapolation (interpolation)

Example: Set of ensembles by MILC collaboration



Five collaborations have now generated sets of ensembles that include sea quarks with physical light-quark masses:

PACS-CS, BMW, MILC, RBC/UKQCD, ETM

# Heavy Quark Treatment

---

- For light quarks (  $m_\ell < \Lambda_{\text{QCD}}$  ), discretization errors  $\sim \alpha_s^k (a\Lambda_{\text{QCD}})^n$
- For heavy quarks, discretization errors  $\sim \alpha_s^k (am_h)^n$   
with currently available lattice spacings  
for  $b$  quarks  $am_b > 1$   
for charm  $am_c \sim 0.15-0.6$

⇒ for charm can use light quark methods, if action is sufficiently improved (HISQ, tmWilson, NP imp. Wilson,...)

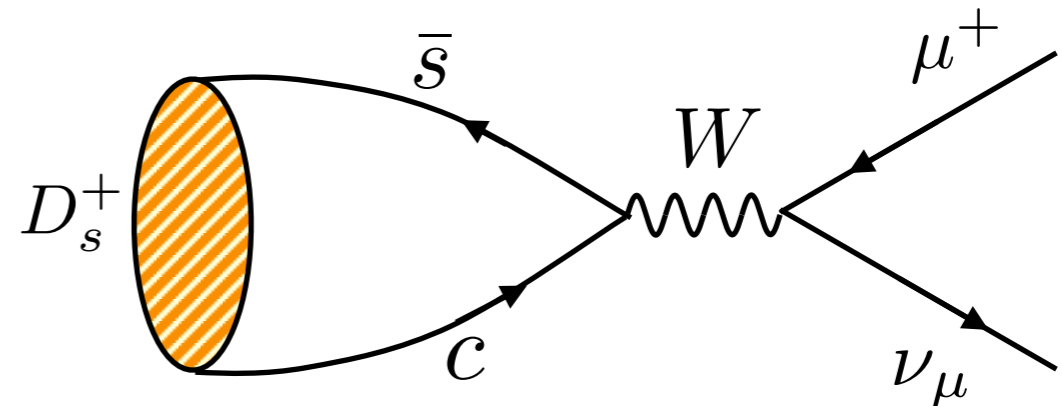
need effective field theory methods for  $b$  quarks

- avoid errors of  $(am_c)^n$  in the action by matching to continuum HQET:
  - ♦ relativistic HQ actions (Fermilab, Columbia, Tsukuba)  
can be used for charm and bottom



# Leptonic $D$ decay

example:  $D_s^+ \rightarrow \mu^+ \nu_\mu$

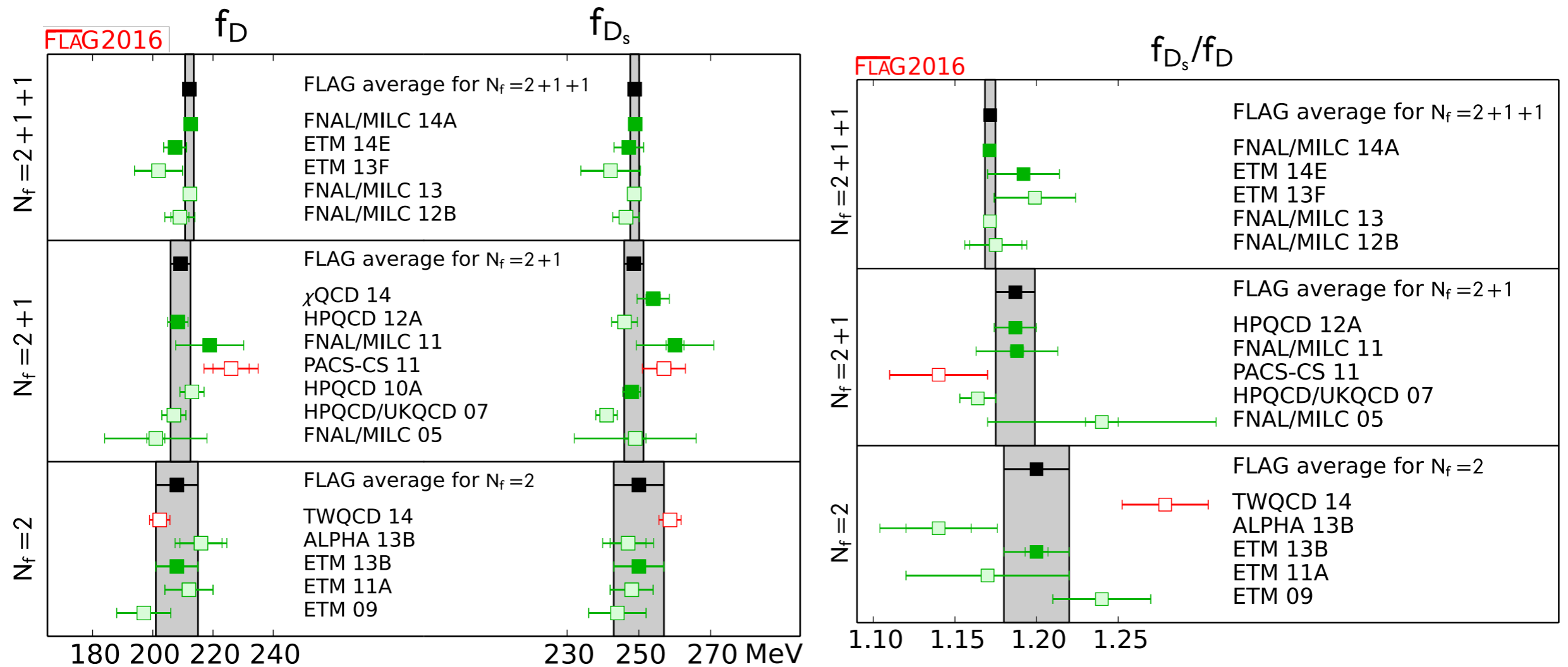


$$\Gamma(D_s^+ \rightarrow \ell^+ \nu_\ell(\gamma)) = (\text{known}) \times (1 + \delta_{\text{EM}}^\ell) \times |V_{cs}|^2 f_{D_s}^2$$

- use experiment + LQCD input for determination of CKM element
- similar for  $B$  ( $|V_{ub}|$ ) and  $K$  ( $|V_{us}|$ ) mesons
- **SU(3) ratio**  $f_{D_s}/f_{D^+}$ : statistical and systematic errors tend to cancel.
- $\delta_{\text{EM}}^\ell$  includes structure dependent EM corrections. It is needed to relate the “pure QCD” decay constant to experiment and is currently estimated phenomenologically.

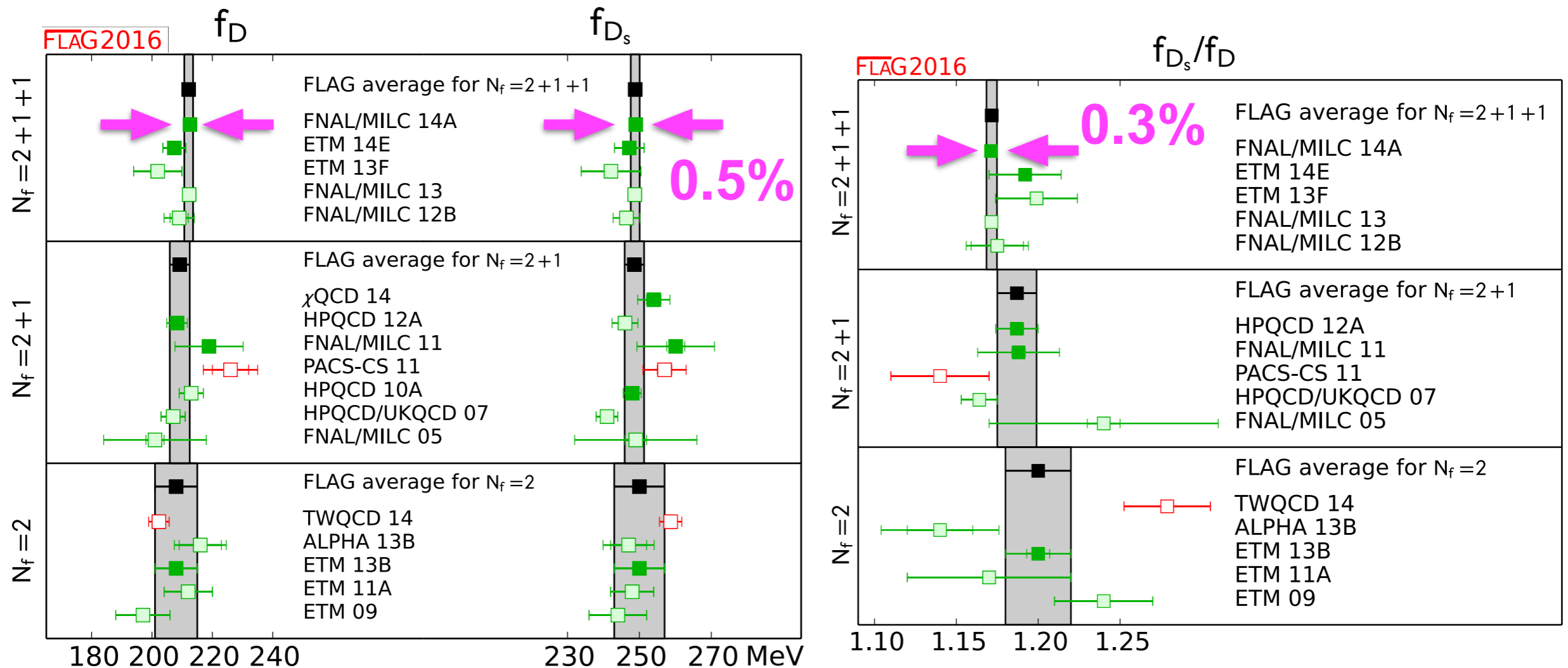
# $D_{(s)}$ decay constant results

S. Aoki et al (FLAG review, arXiv:1607.00299)



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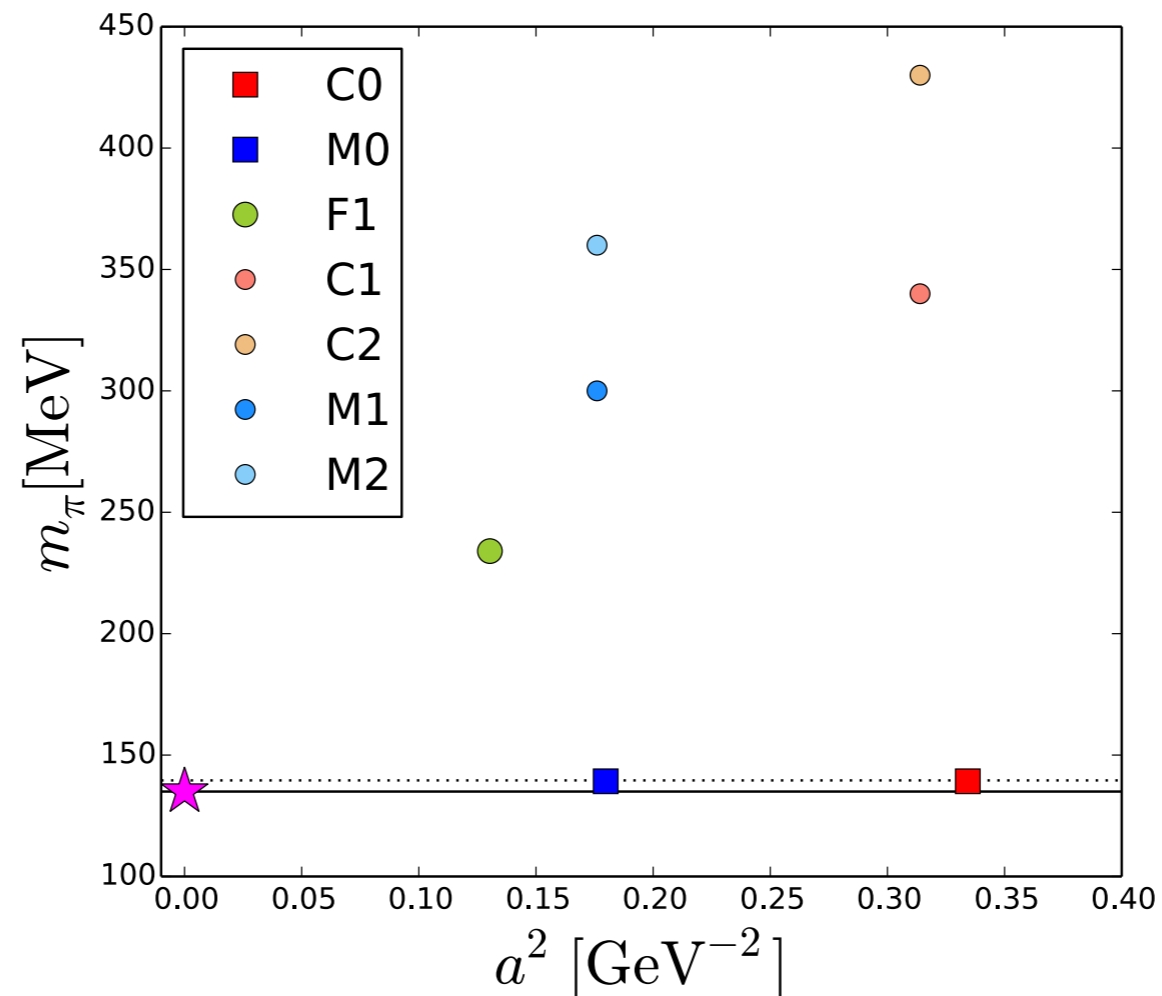
small errors in FNAL/MILC 14A (arXiv:1407.3772, 2014 PRD) due to

- physical mass ensembles
- improved action (small discretization errors)
- small lattice spacings
- PCAC (no renormalization)

# $D_{(s)}$ decay constant results

RBC/UKQCD (J.T. Tsang @ Lattice 2016):

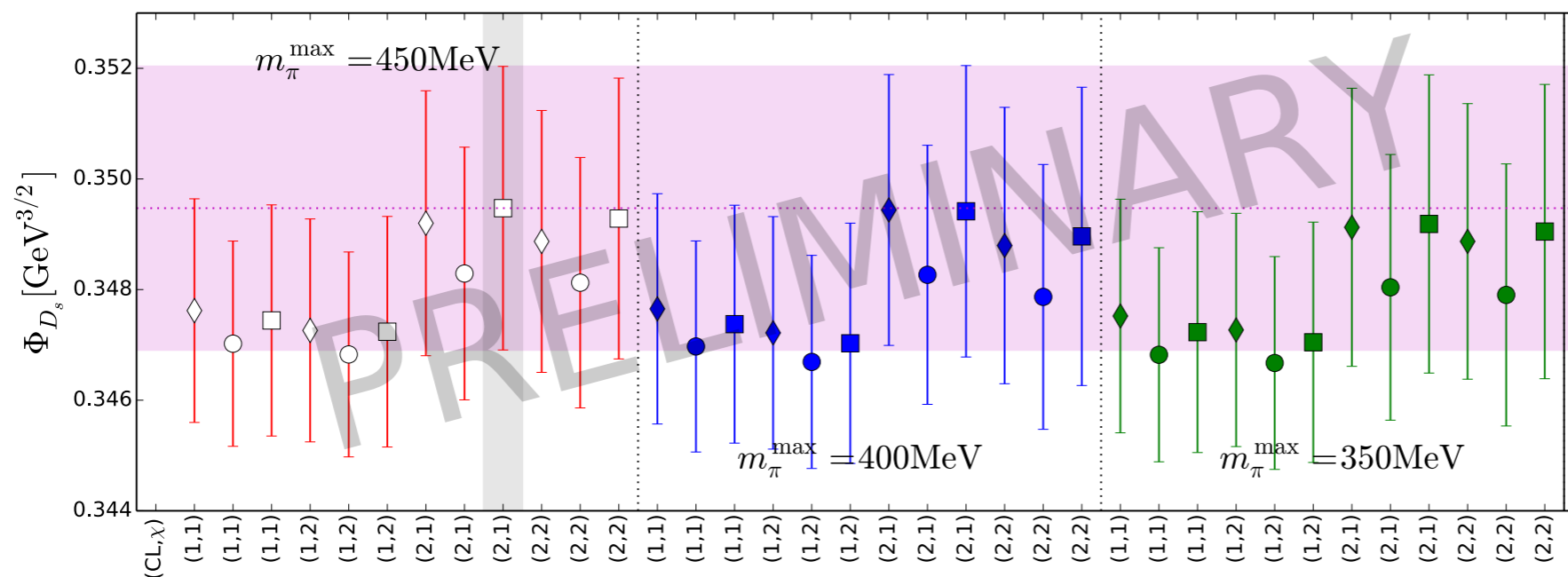
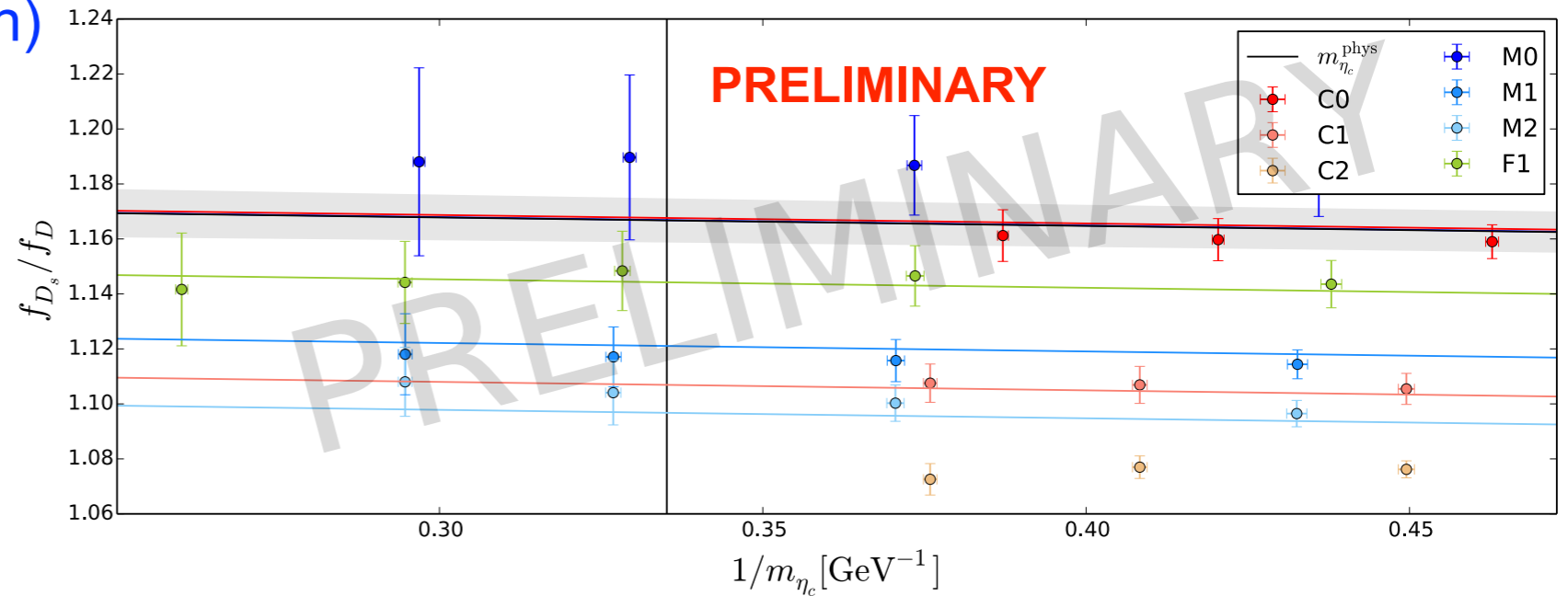
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- physical mass ensembles
- PCAC (no renormalization)



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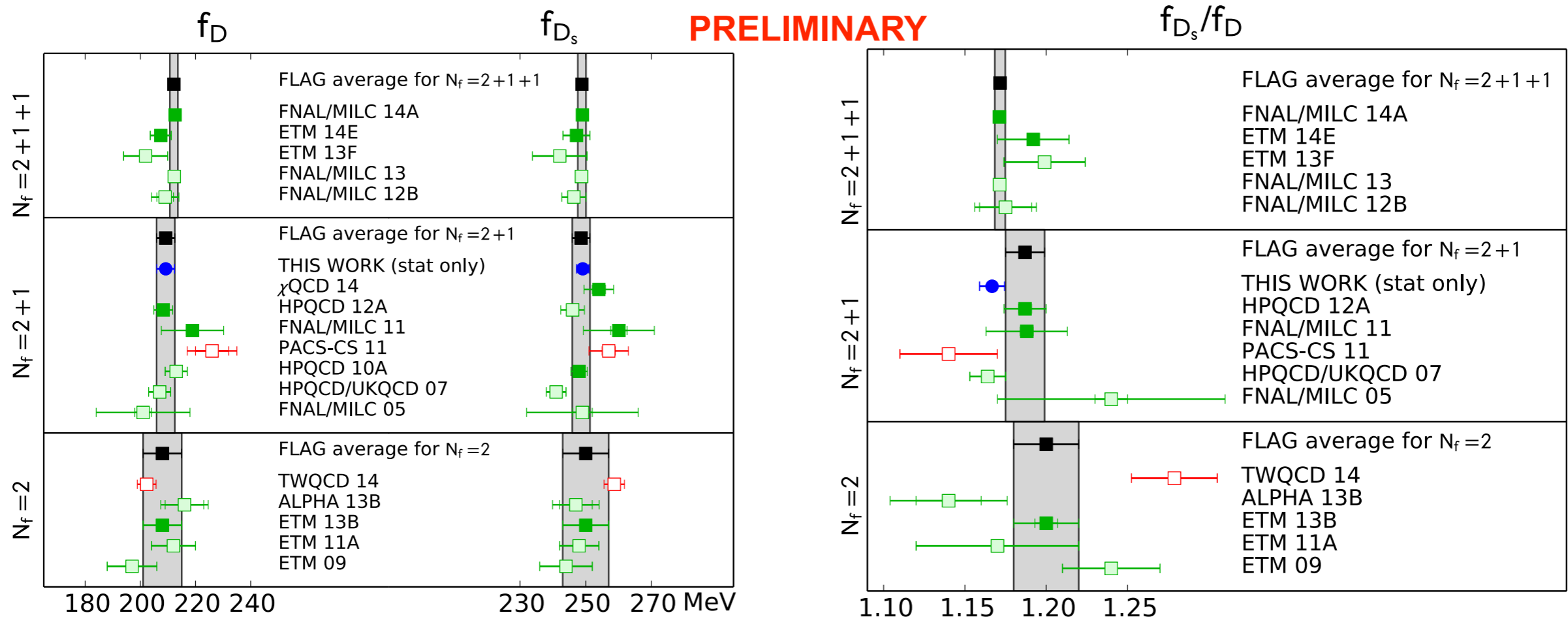
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# $D_{(s)}$ decay constant results

J. T. Tsang (RBC/UKQCD) @ Lattice 2016:

**PRELIMINARY**



RBC/UKQCD (J.T. Tsang @ Lattice 2016):

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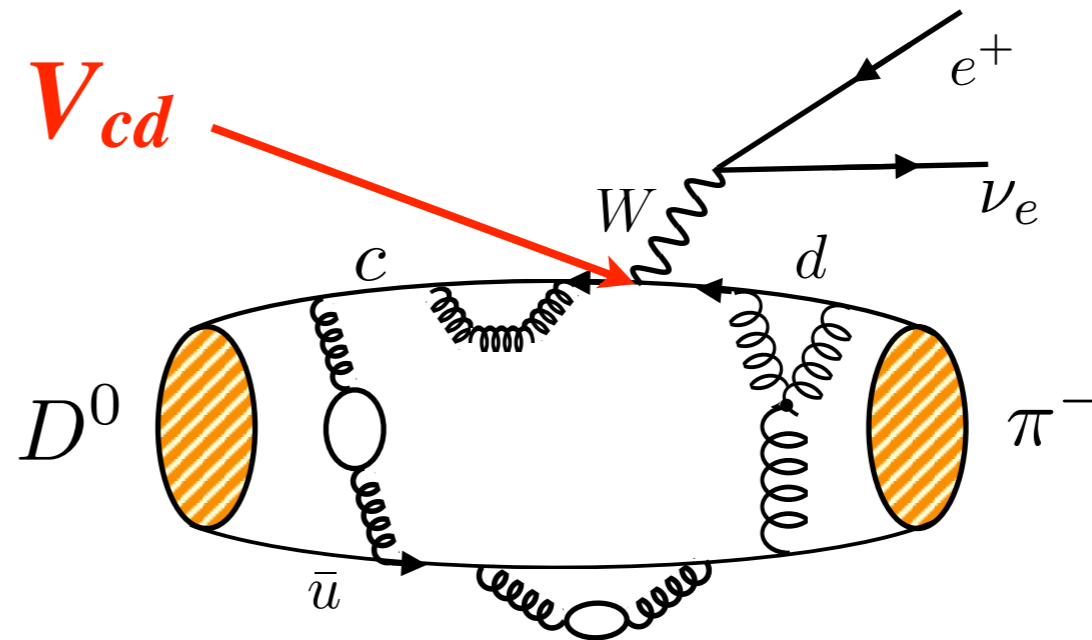
also ongoing work by:

- ALPHA/RQCD (imp. Wilson)
- FNAL/MILC (with Fermilab charm)

+ new results from ETM on  $f_{D^*(s)}$   
(Melis @ Lattice 2016)

# Semileptonic $D$ -meson decay

Example:  $D \rightarrow \pi \ell \nu$



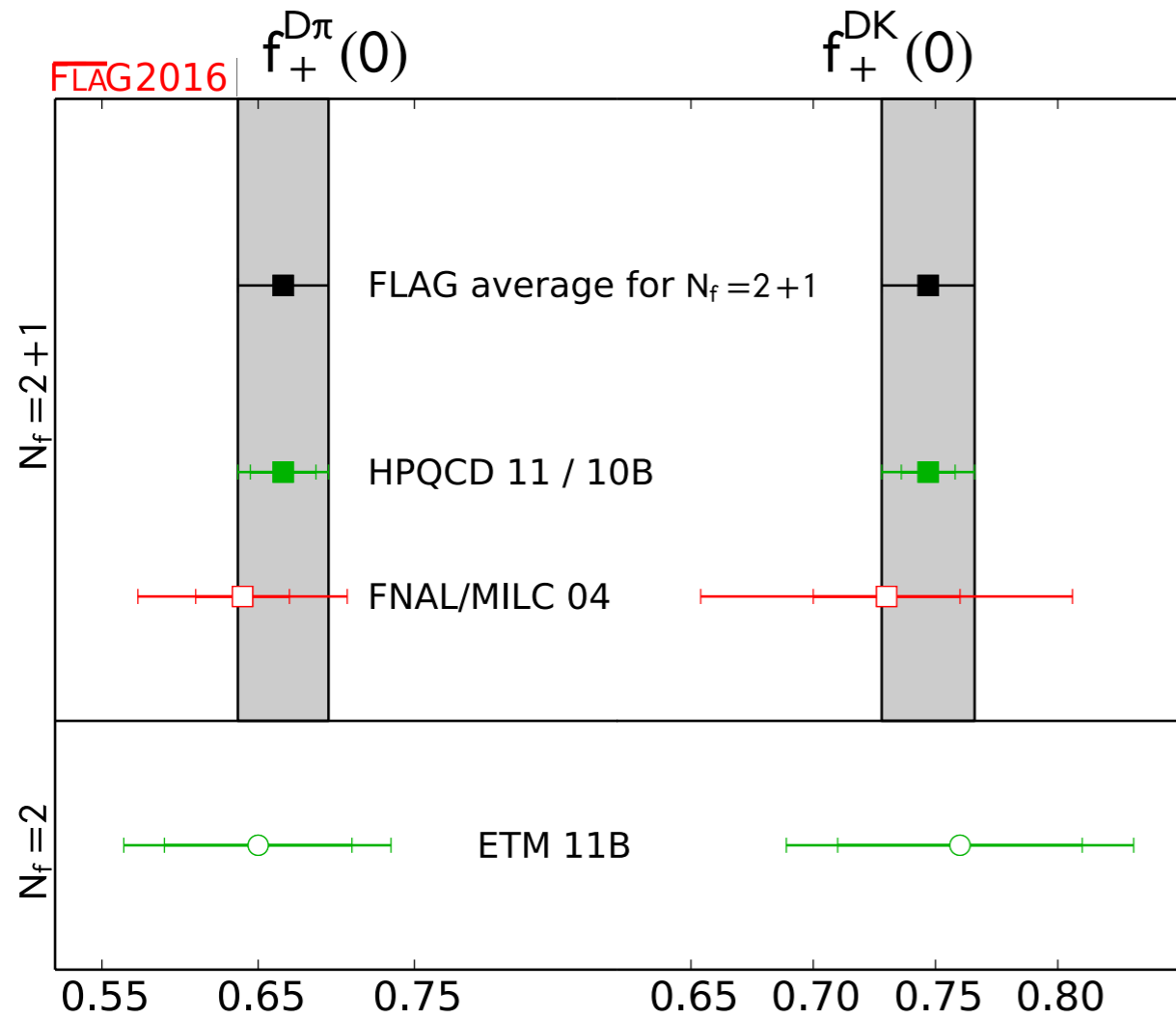
$$\frac{d\Gamma(D \rightarrow \pi \ell \nu)}{dq^2} = (\text{known}) \times |V_{cd}|^2 f_+^2(q^2) \quad \boxed{\ell = e, \mu}$$

- ★ can calculate the form factors for the entire recoil energy range
- ★ can use  $z$ -expansion\* for model-independent parameterization of  $q^2$  dependence
- ★ calculate both form factors  $f_+(q^2), f_0(q^2)$
- ★ can compare shape between experiment and lattice
- ★ extension to rare SL decay form factors ( $f_T$ ) straightforward

\*see backup slides

# $D$ SL form factor results

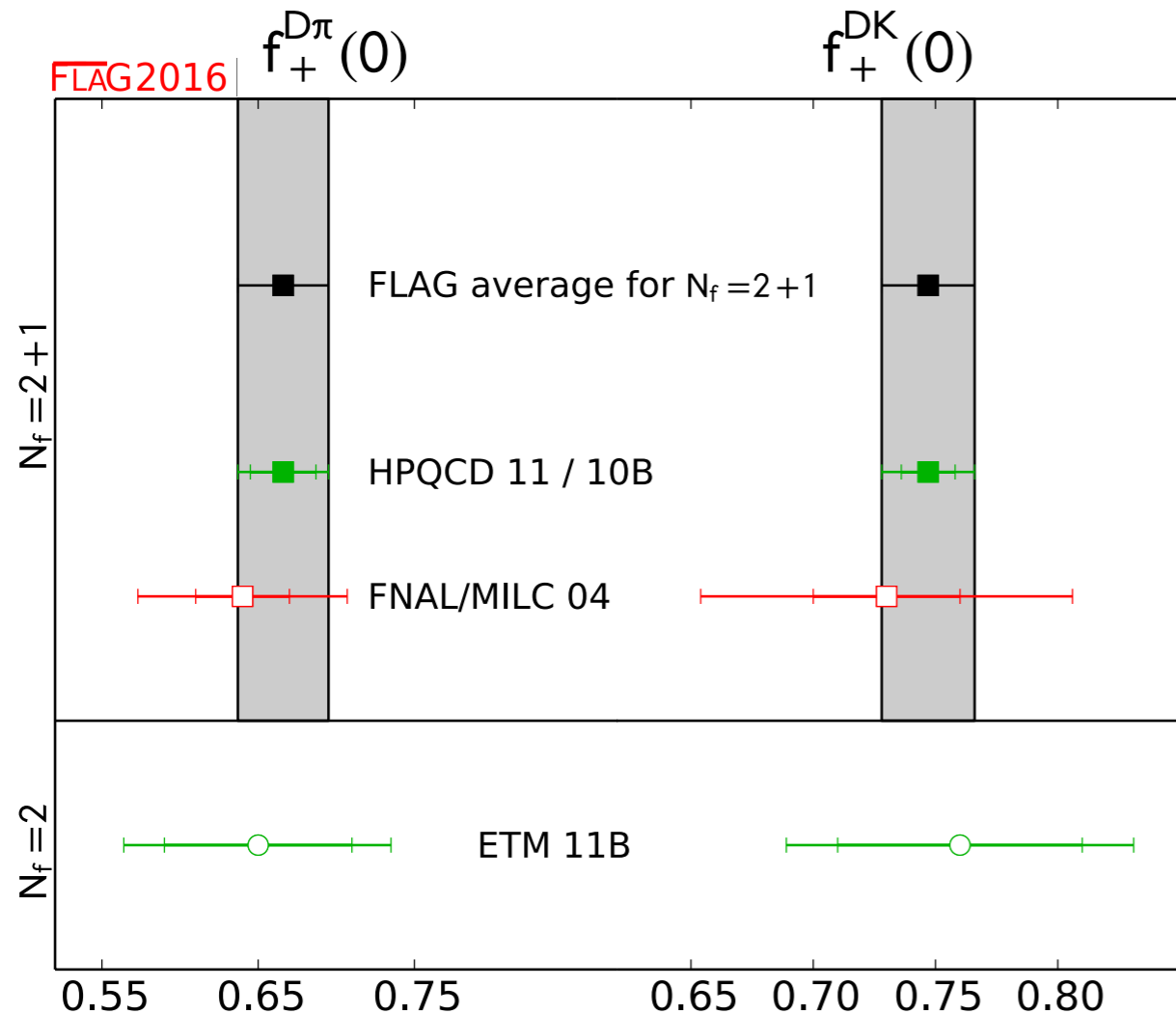
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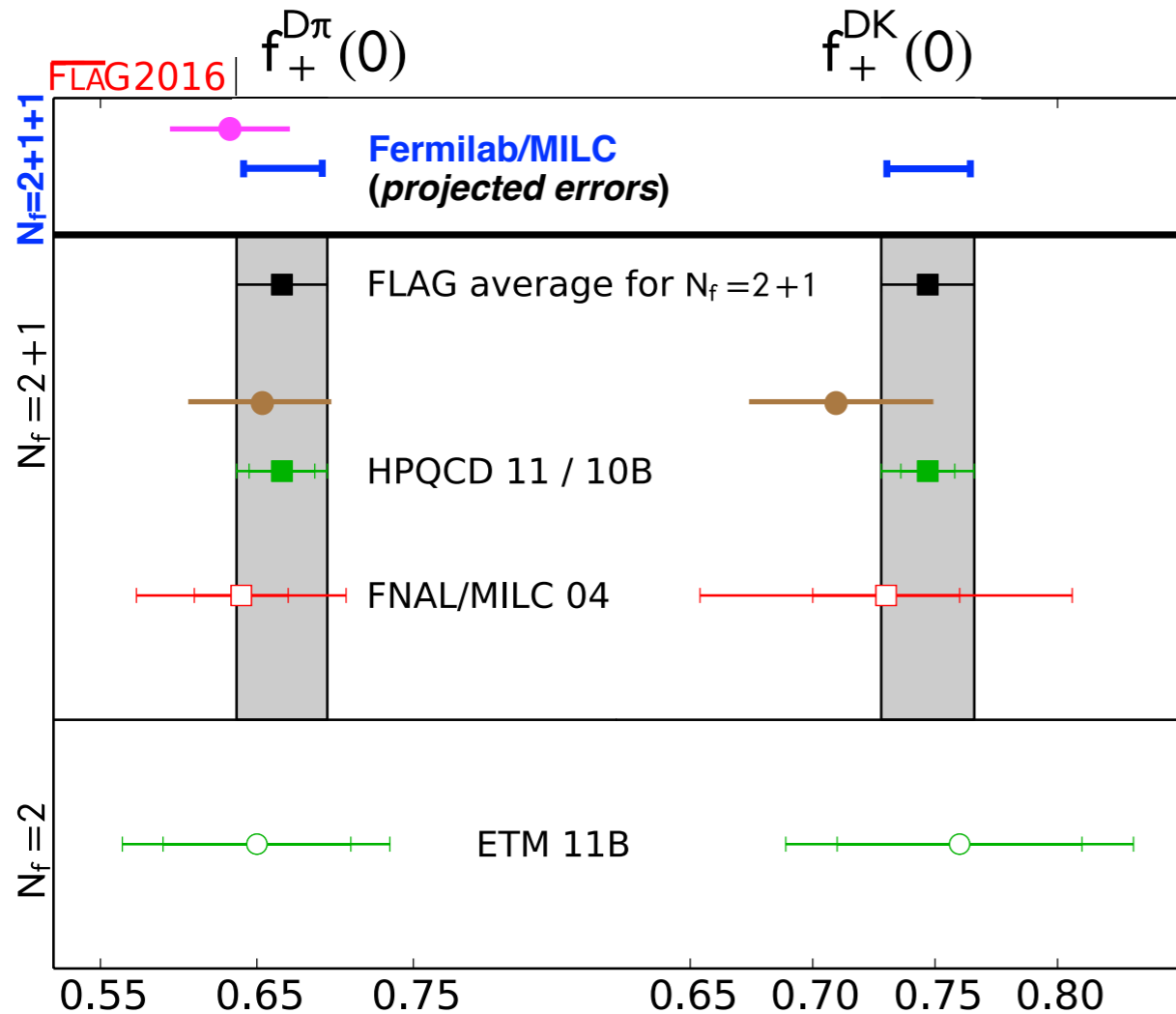


new preliminary results @ Lattice 2016:

- **ETM (G. Salerno)**  
2+1+1 flavors of tmWilson  
calculate all form factors over whole  $q^2$  range  
modified z-expansion  
preliminary sys. errors
- **FNAL/MILC (S. Gottlieb, T. Primer)**  
**no central values (yet)**  
2+1+1 flavors of HISQ  
physical mass ensembles  
calculate directly at zero  $q^2$
- **JLQCD (T. Kaneko)**  
2+1 flavors of DW fermions  
extrapolate to zero  $q^2$  with z-expansion  
chiral-continuum extrapolation  
still adding ensembles to analysis

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adapted from S. Aoki et al (arXiv:1607.00299)



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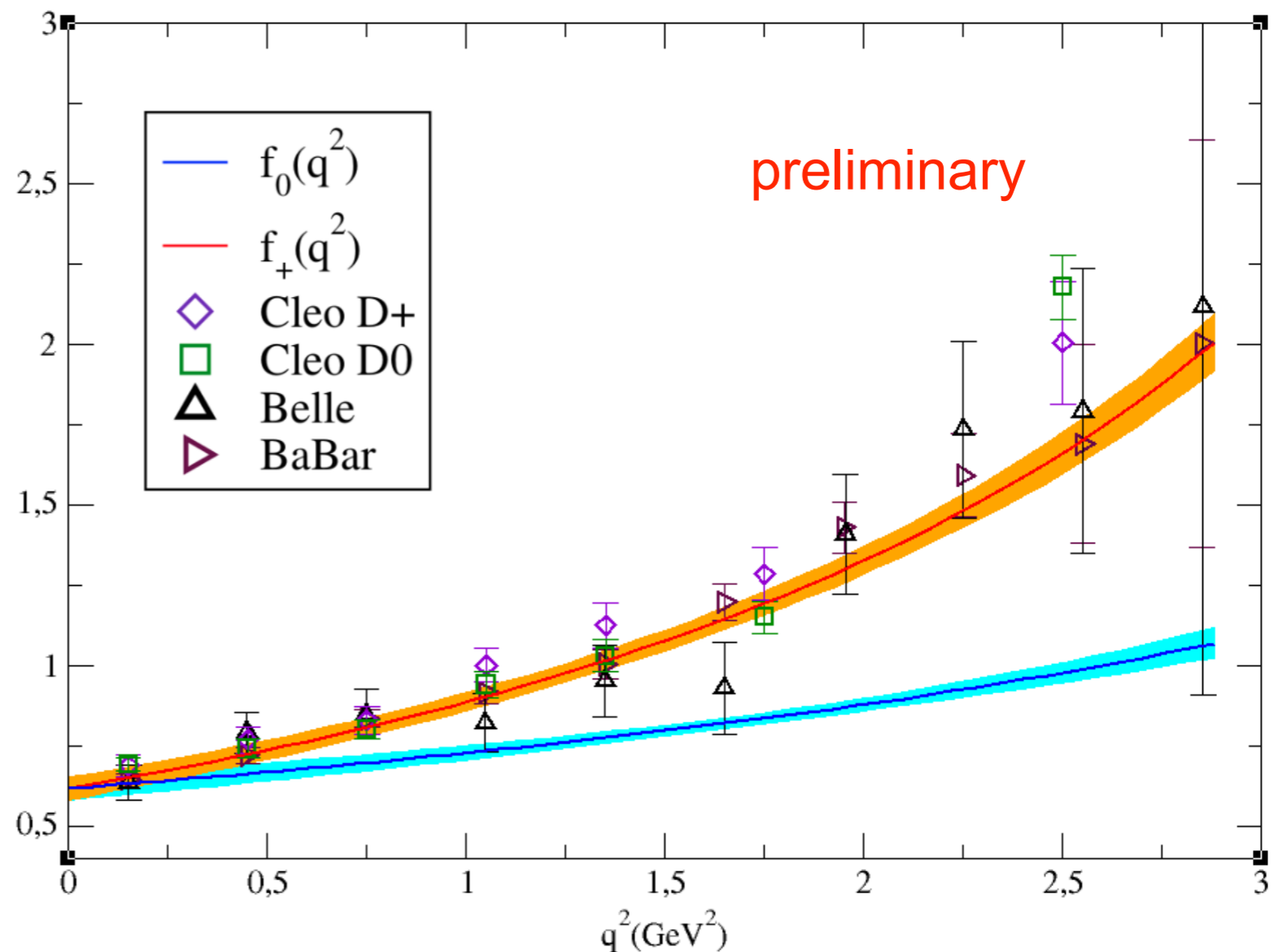
- 2+1+1 flavors of tmWilson
- calculate  $f_+, f_0$  over whole  $q^2$  range
- modified  $z$ -expansion
- correct for hypercubic discretization effects
- preliminary sys. errors

see talk by G. Salerno on Thursday

Three different values of the lattice spacing:  $0.06 \text{ fm} \div 0.09 \text{ fm}$

Different volumes:  $2 \text{ fm} \div 3 \text{ fm}$

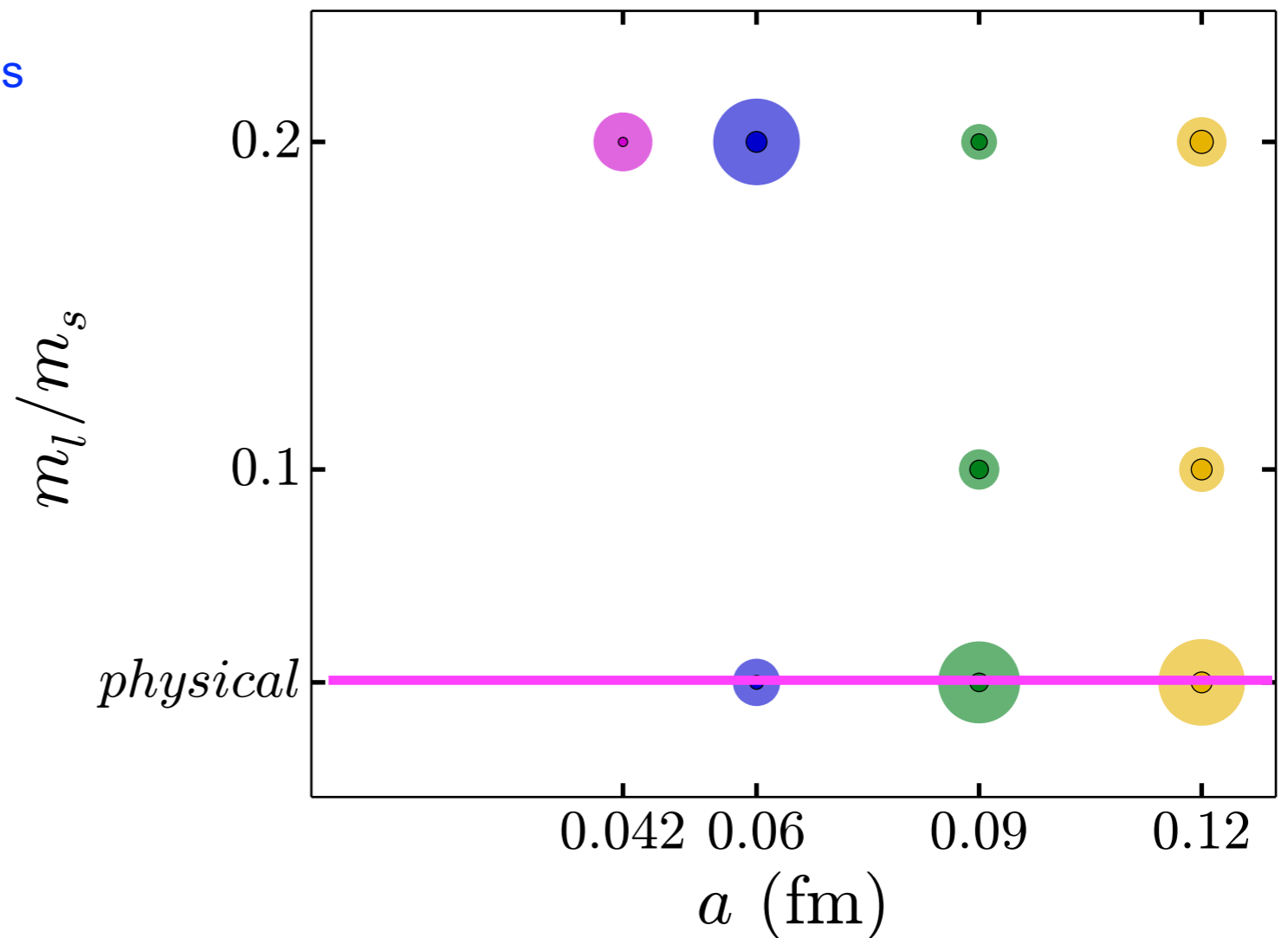
Pion masses in range  $210 \div 440 \text{ MeV}$



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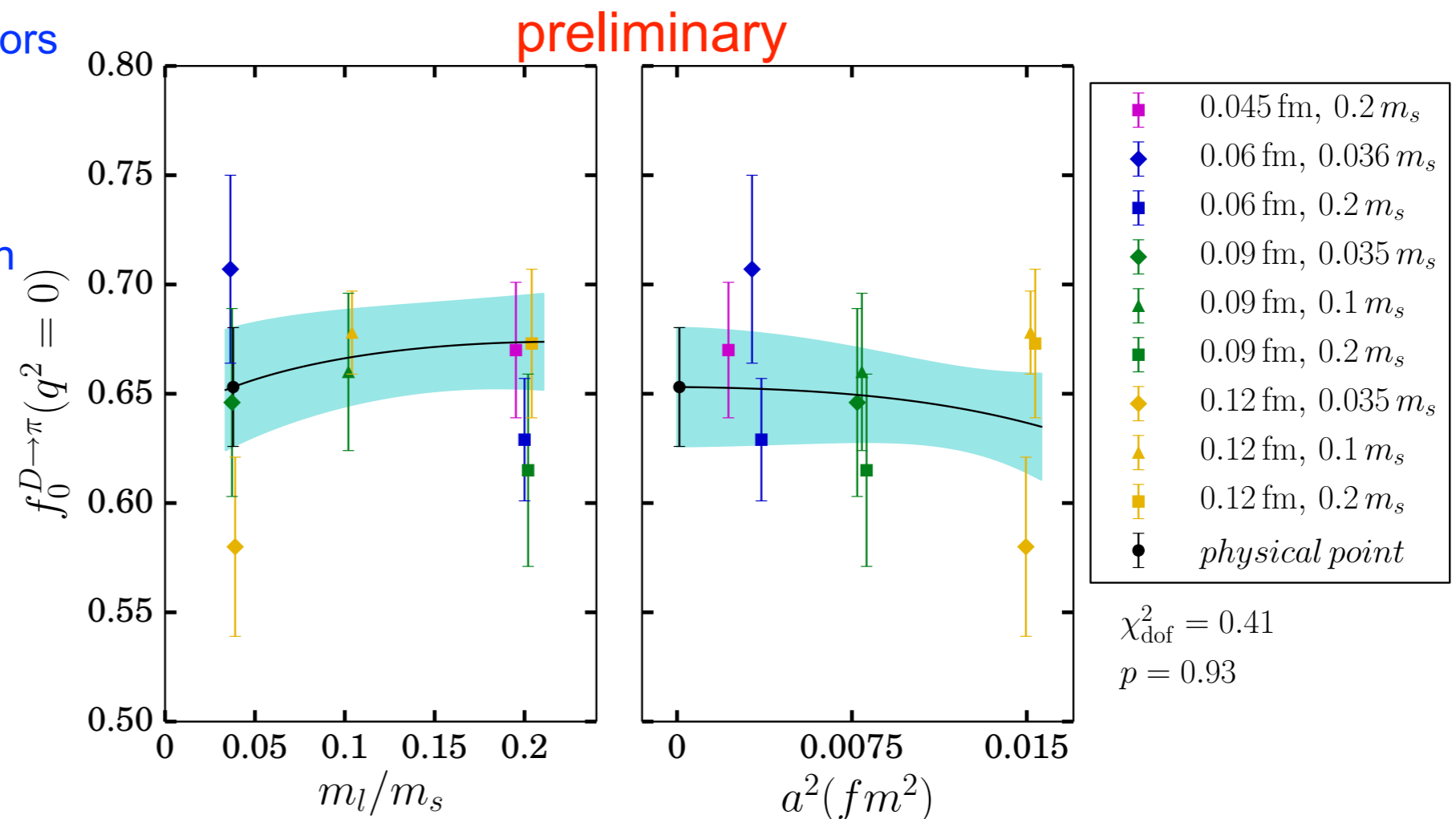
- 2+1+1 HISQ ensembles  
physical light quark masses
- HISQ valence charm, strange, light
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chiral-continuum extrapolation
- preliminary systematic error analysis
- next step:  
vector and scalar form factors  
+ range of recoil momenta  
 $\Rightarrow$  whole  $q^2$  range
- will yield better precision  
and shape comparison with  
experiment



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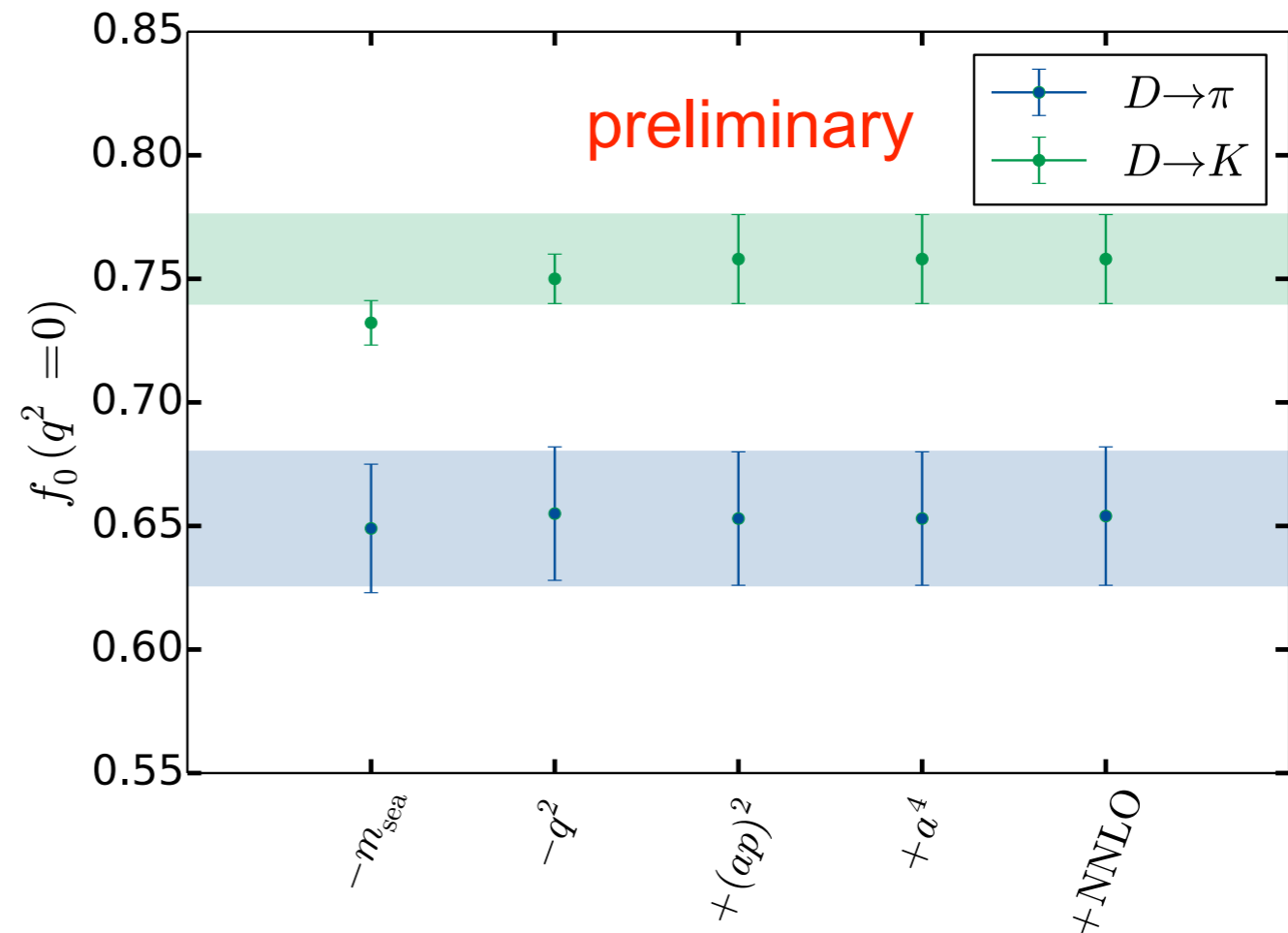
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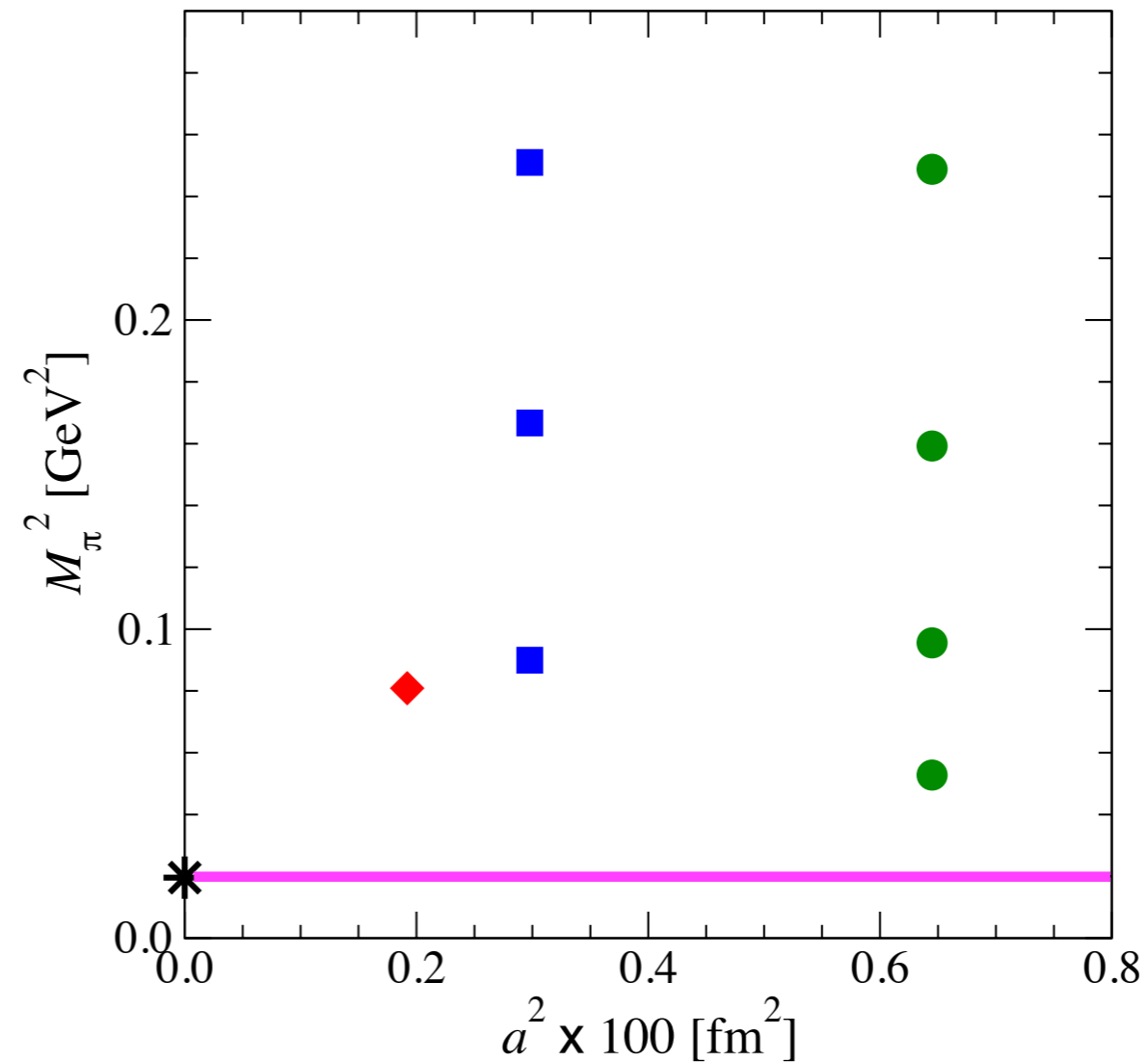


# $D$ SL form factor results

JLQCD (T. Kaneko) @ Lattice 2016

- 2+1 flavors of DW fermions
- extrapolate to zero  $q^2$  with  $z$ -expansion
- chiral-continuum extrapolation
- still adding ensembles to analysis

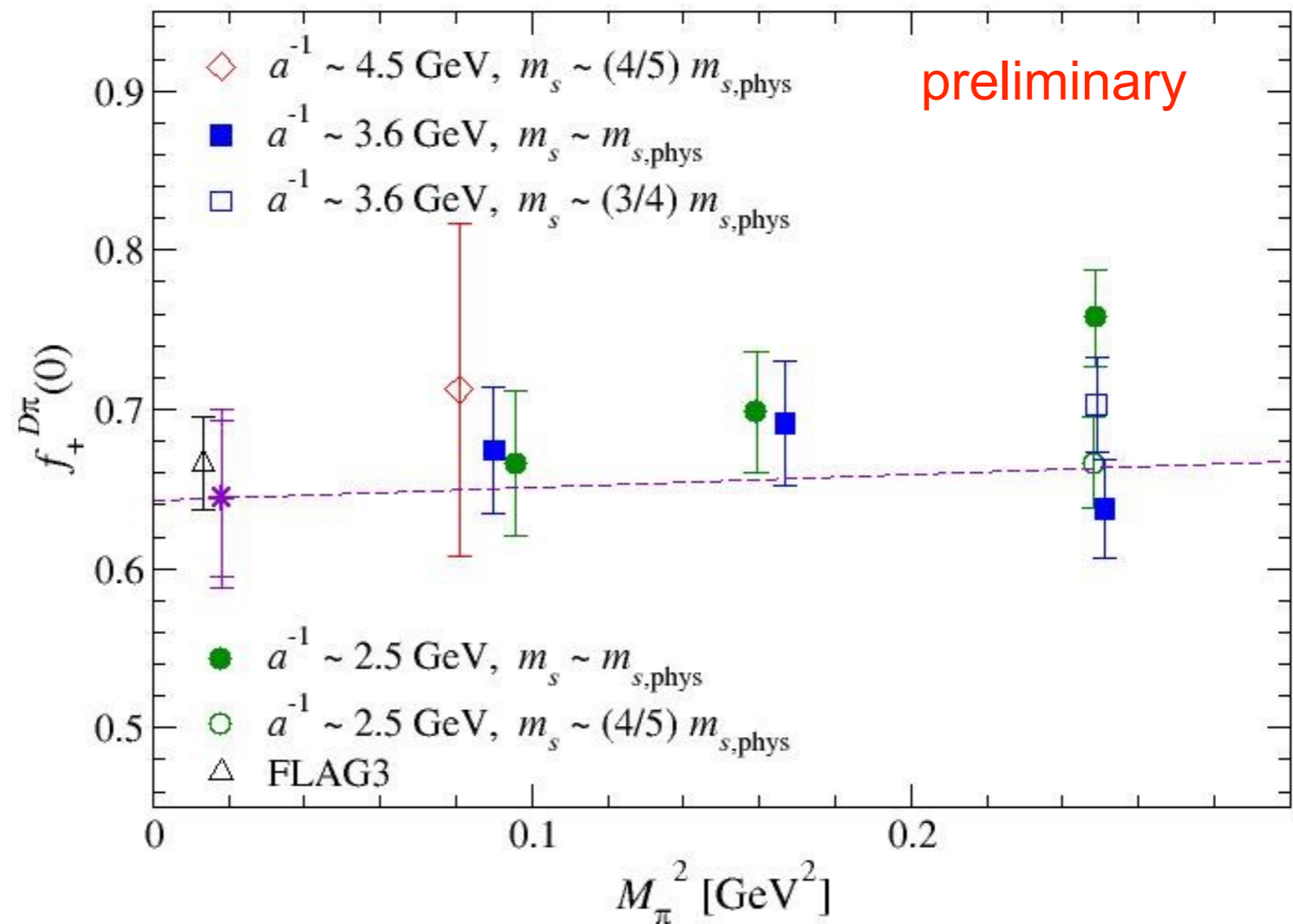
preliminary



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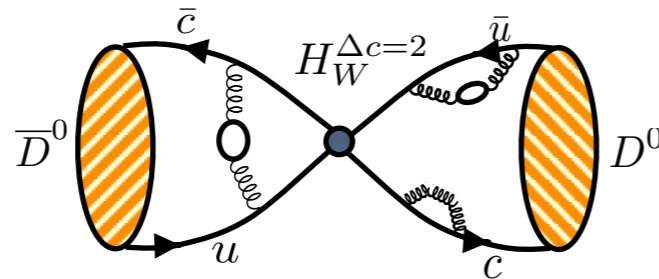
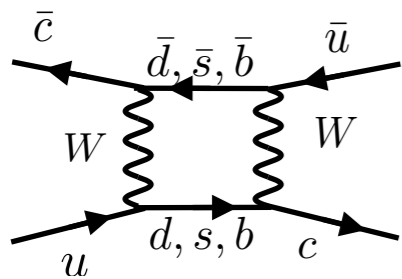
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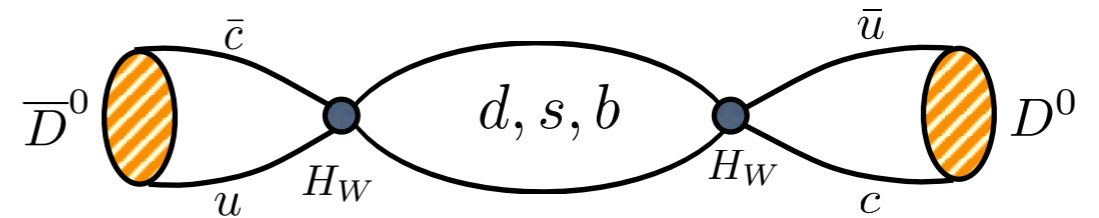


# Neutral $D$ meson mixing

$$M_{12} - \frac{i}{2}\Gamma_{12} \propto \langle D^0 | H_W^{\Delta c=2} | \bar{D}^0 \rangle + \sum_n \frac{\langle D^0 | H_W^{\Delta c=1} | n \rangle \langle n | H_W^{\Delta c=1} | \bar{D}^0 \rangle}{M_D - E_n + i\epsilon}$$



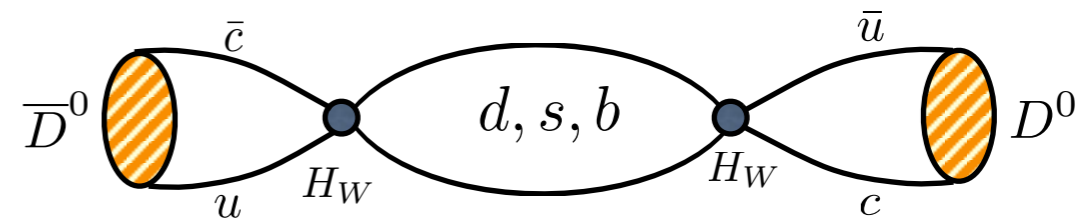
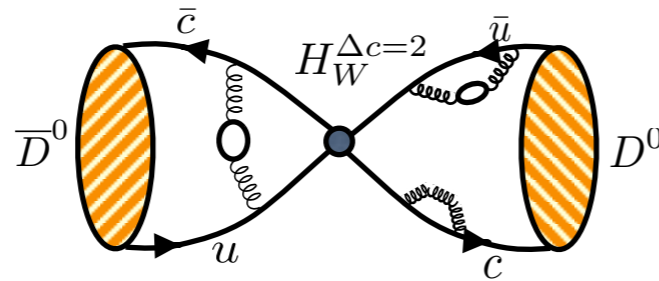
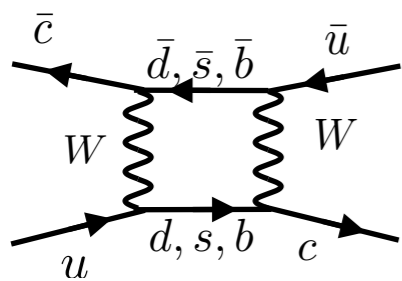
short distance



long distance

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“Simple”

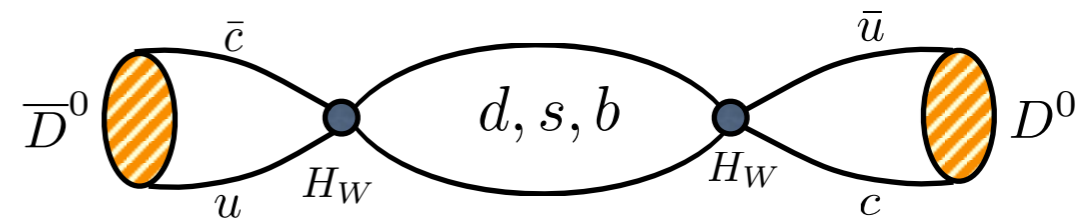
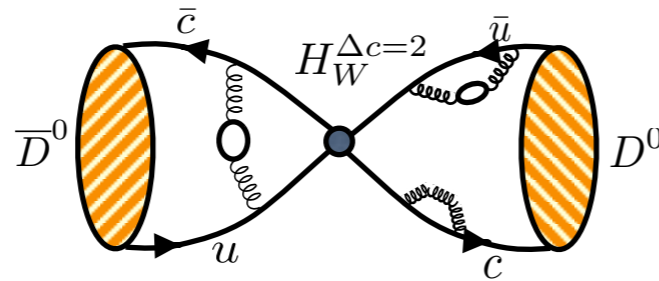
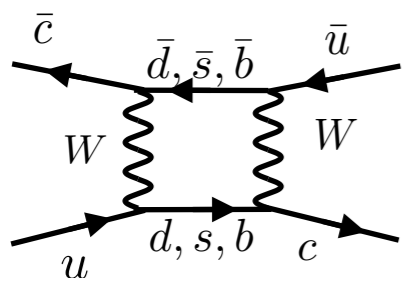
- can use the same methods as for  $B$  mixing (and decay constants, form factors)
- BSMs with heavy new particles can contribute here

“Hard”

- large contribution
- intermediate state can include multiple ( $>2$ ) hadrons:
  - ◆ formalism for multi-hadron states still under development (Hansen & Sharpe, arXiv:1602.00324, 2016 PRD)
  - ◆ not a problem for Kaon mixing
    - first calculation of long-distance contribution already exists (RBC/UKCD, arXiv:1406.0916, 2014 PRL)

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short distance

long distance



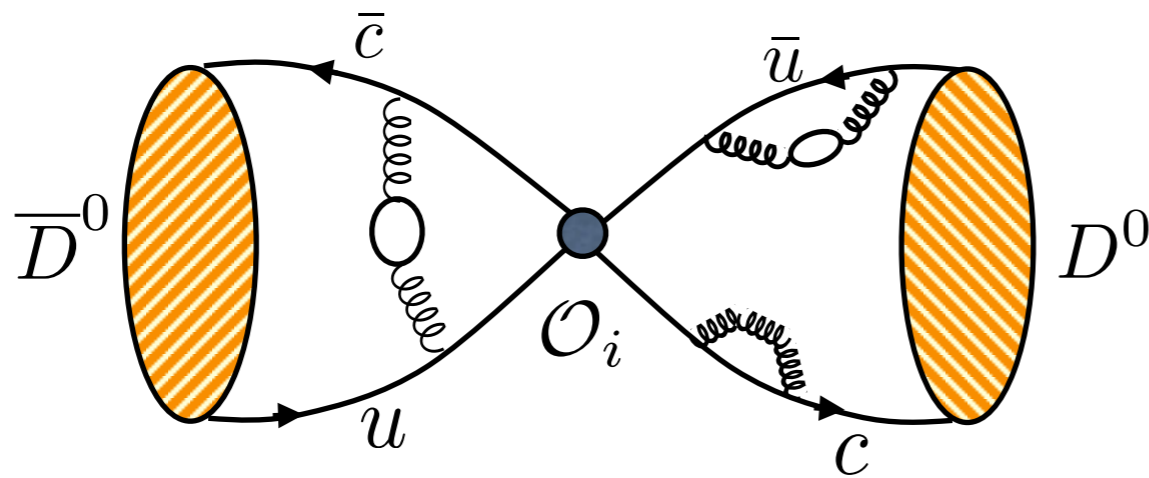
## “Simple”

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# Neutral $D$ meson mixing



In the SM and beyond:

$$\mathcal{H}_{\text{eff}} = \sum_{i=1}^5 c_i(\mu) \mathcal{O}_i(\mu)$$

$$\mathcal{O}_1 = \bar{c} \gamma^\mu L u \bar{c} \gamma^\mu L u$$

$$\mathcal{O}_2 = \bar{c} L u \bar{c} L u$$

$$\mathcal{O}_3 = \bar{c}^\alpha L u^\beta \bar{c}^\beta L u^\alpha$$

$$\mathcal{O}_4 = \bar{c} L u \bar{c} R u$$

$$\mathcal{O}_5 = \bar{c}^\alpha L u^\beta \bar{c}^\beta R u^\alpha c$$

$$\langle \mathcal{O}_i \rangle \equiv \langle D^0 | \mathcal{O}_i | \bar{D}^0 \rangle (\mu) = e_i M_D^2 f_D^2 B_D^{(i)}(\mu)$$

choose  $\mu = 3 \text{ GeV}$

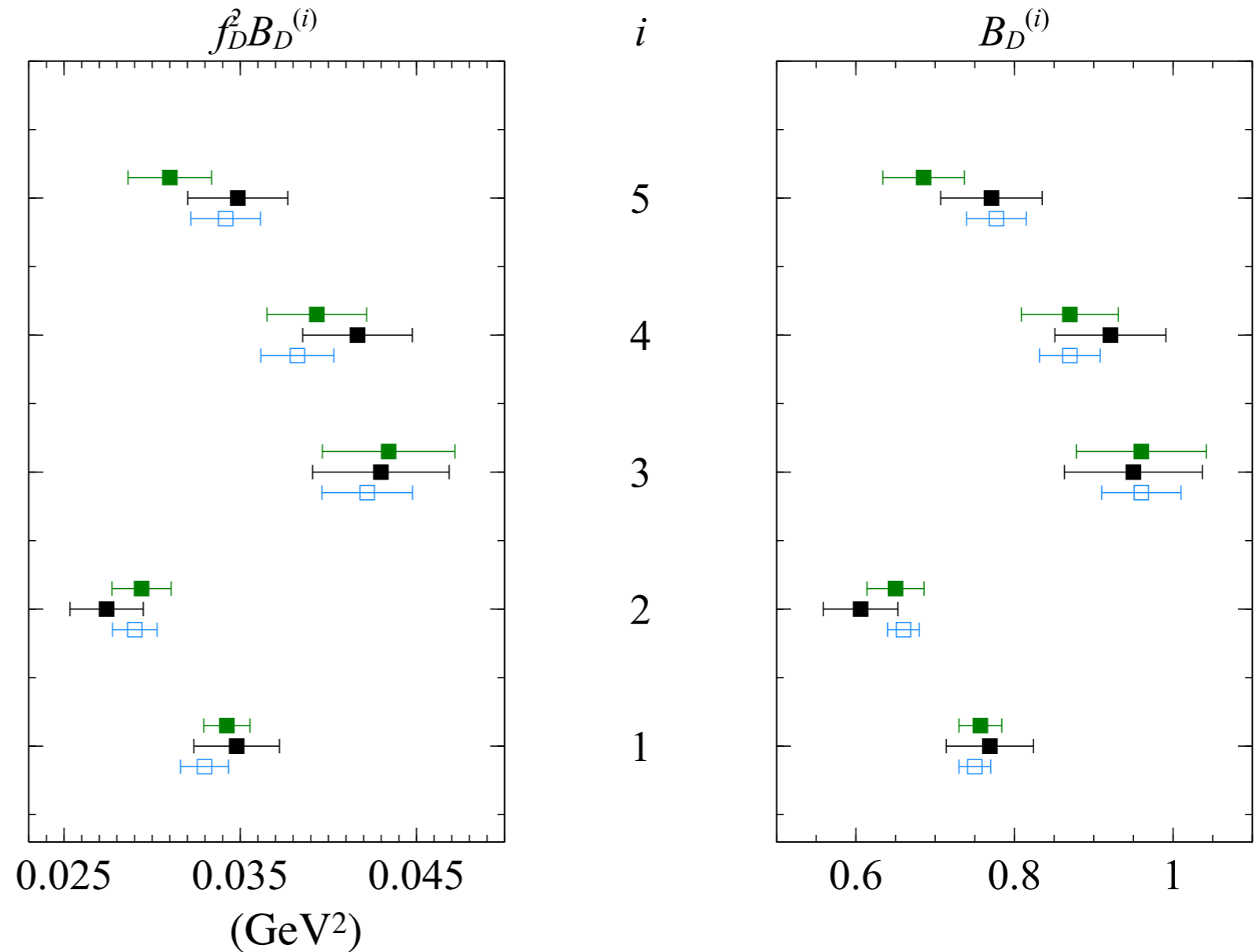
- calculate the matrix elements of all five local operators.

# $D$ mixing results in comparison

$\mu = 3 \text{ GeV}$

A. Kronfeld @ Lattice 2016 (plot by C.C. Chang)

- ETM:  
 $n_f = 2+1+1$   
arXiv:1505.06639
- Fermilab/MILC:  
 $n_f = 2+1$
- ETM:  
 $n_f = 2$   
arXiv:1403.7302

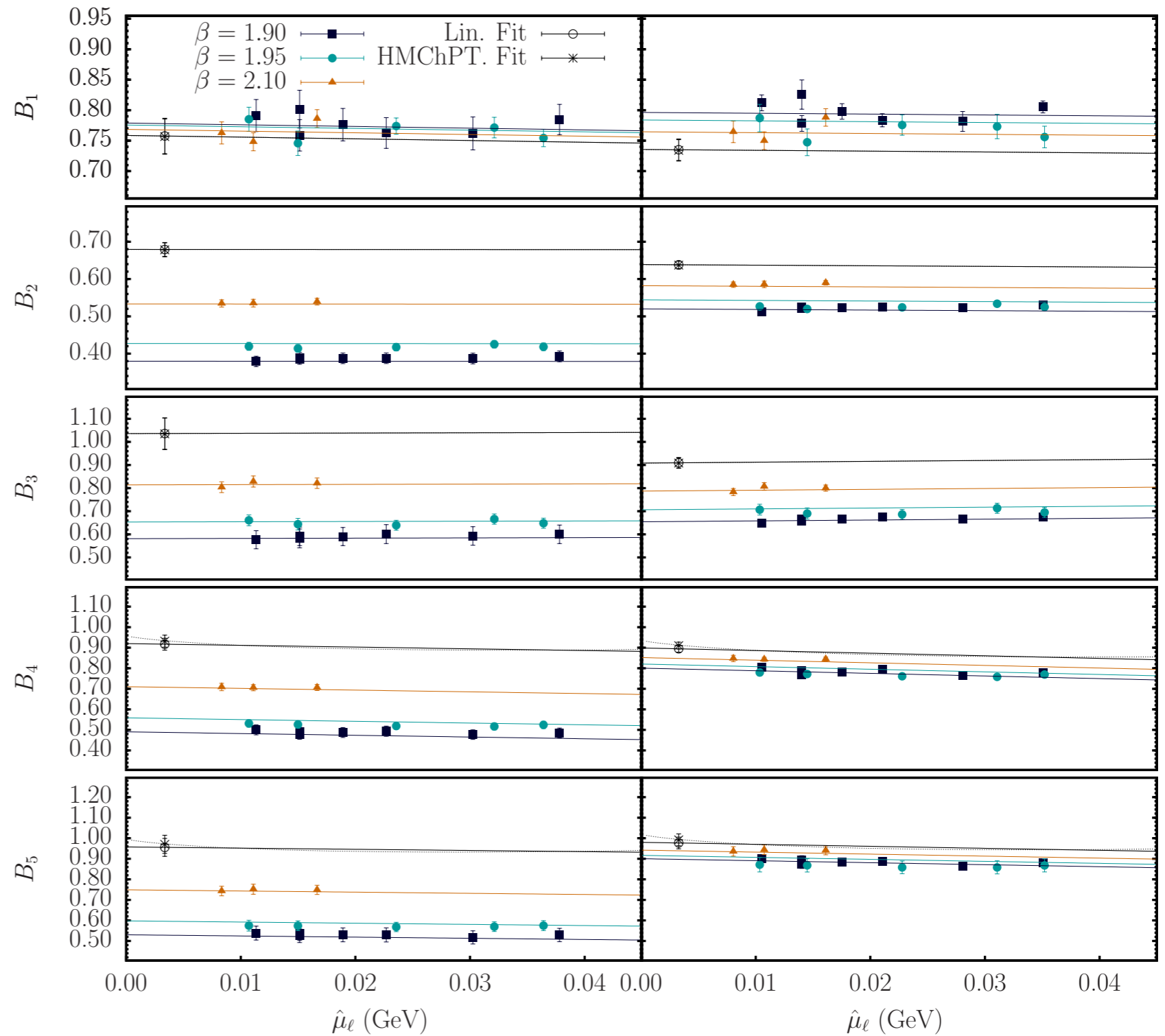


# D mixing results

ETM

$n_f = 2+1+1$  tmWilson

(arXiv:1505.06639, 2015 PRD)

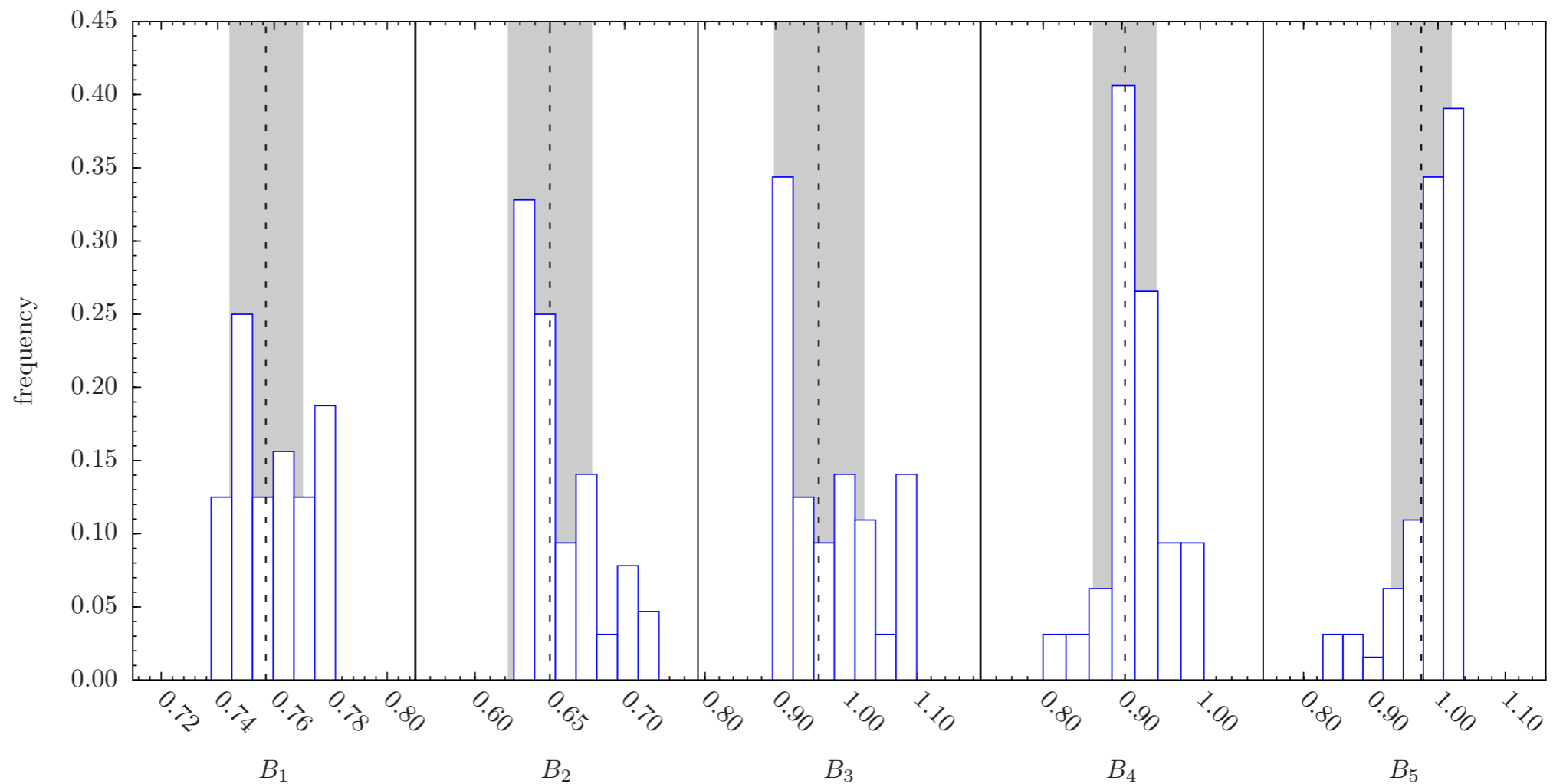
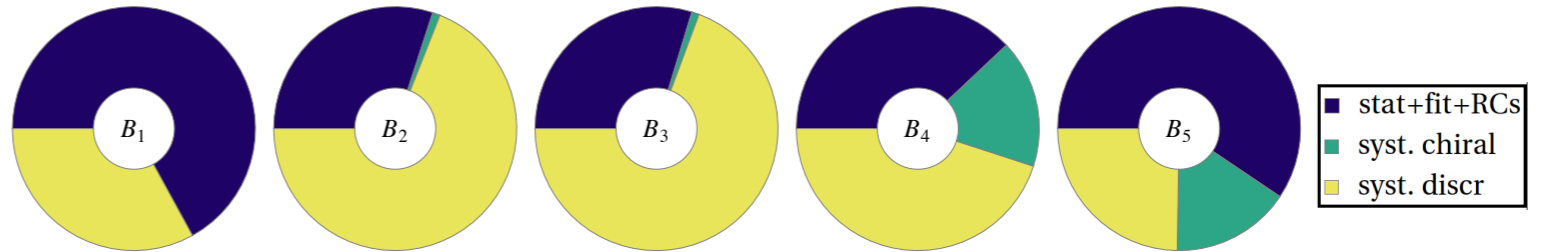


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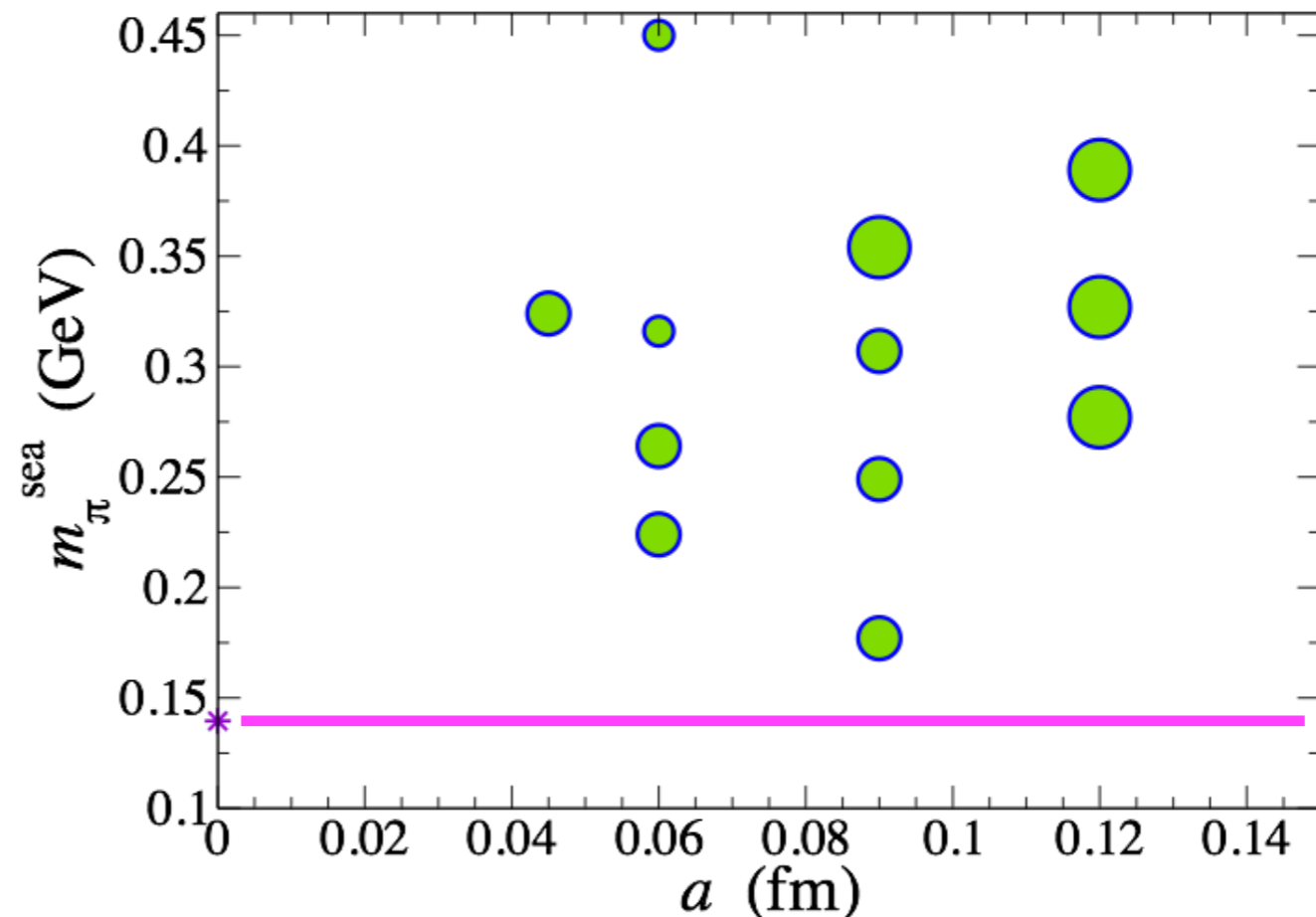
(arXiv:1505.06639, 2015 PRD)



# $D$ mixing results

FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

- 14 MILC asqtad ensembles  
4 lattice spacings  
~ 4 sea quark masses per lattice spacing  
~ 600 - 2000 configurations  
× 4 time-sources per configuration
- Fermilab  $c$  quarks
- mNPR renormalization

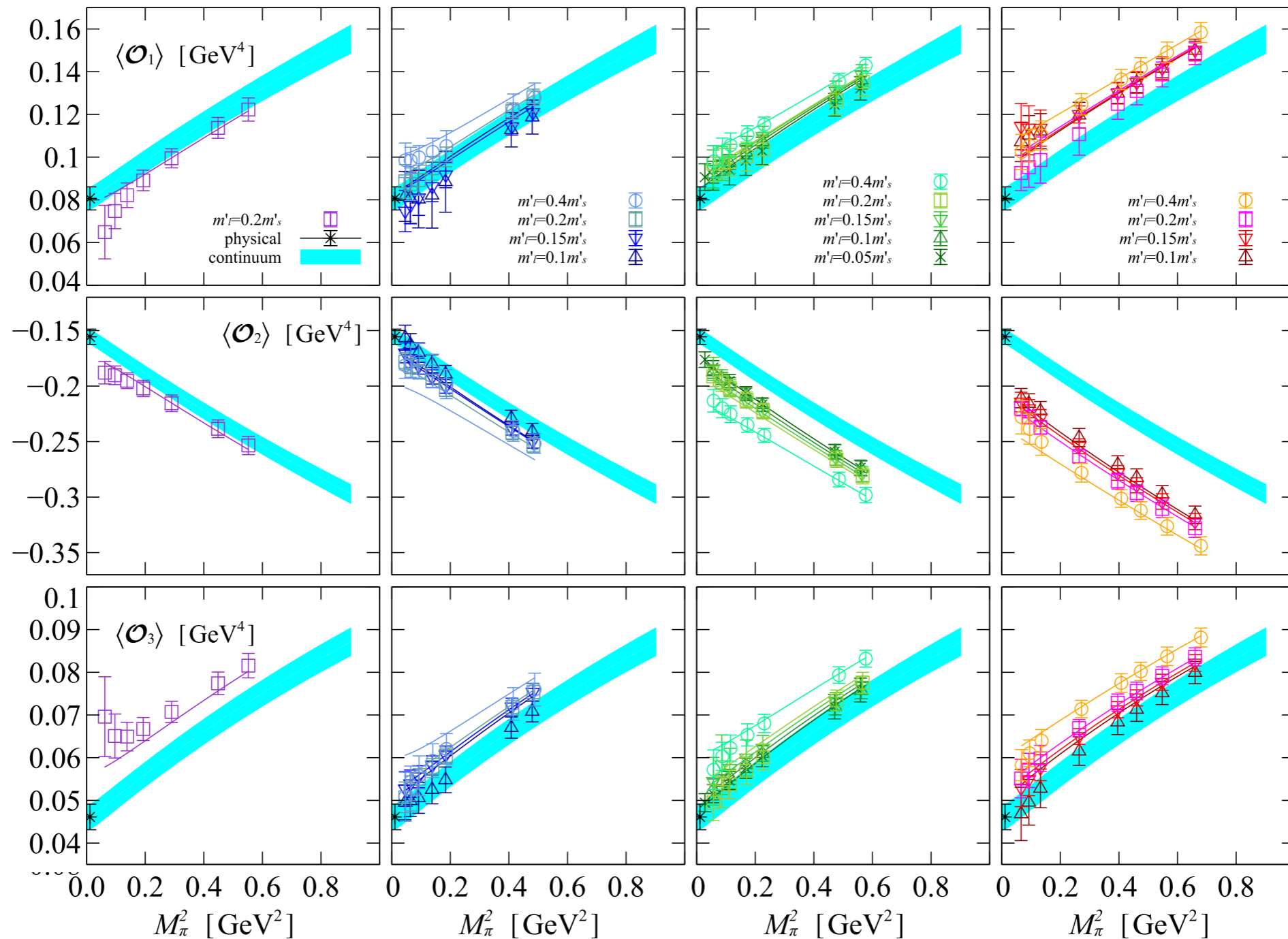




# $D$ mixing results

FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

Operators 1,2,3

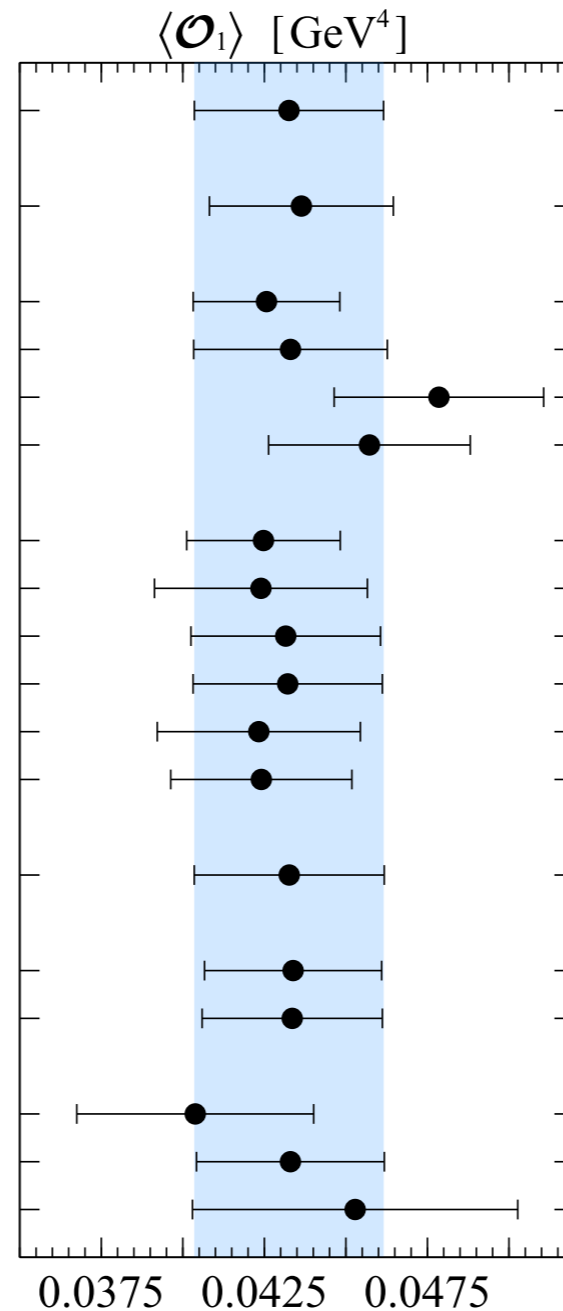


# D mixing results

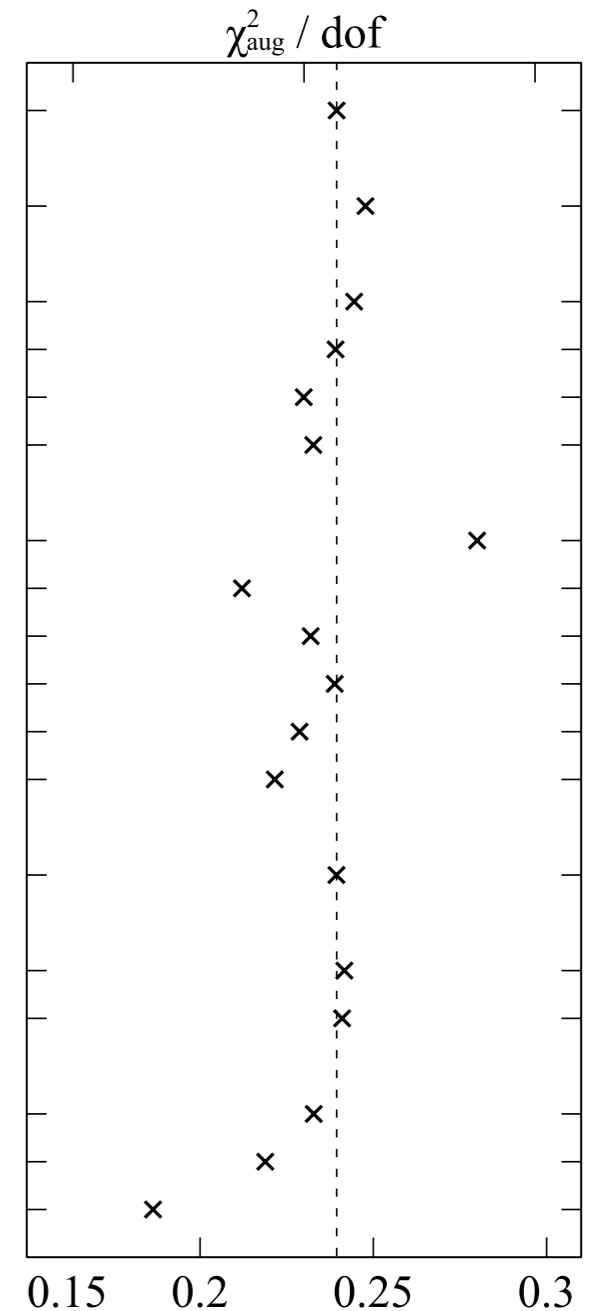
FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

systematic error study

- remove or add higher order terms in fit function:
  - chiral expansion
  - heavy meson expansion
  - light quark discretization effects
  - HQ discretization effects
  - renormalization (perturbative expansion)
- change data included
- change inputs



base  
 $f_k$  vs.  $f_\pi$   
 mNPR  
 mNPR+ $\alpha_s^3$   
 PT<sub>P</sub>+ $\alpha_s^2$   
 PT<sub>L</sub>+ $\alpha_s^2$   
 NLO ( $m_q < 0.65 m_s$ )  
 N<sup>3</sup>LO  
 LO x 2  
 NLO x 2  
 NNLO x 2  
 no splitting  
 generic  $O(\alpha_s a^2)$   
 HQ  $O(\alpha_s a)$  only  
 HQ  $O(\alpha_s a, a^2)$  only  
 no  $a \approx 0.12$  fm  
 no  $a \approx 0.045$  fm  
 individual

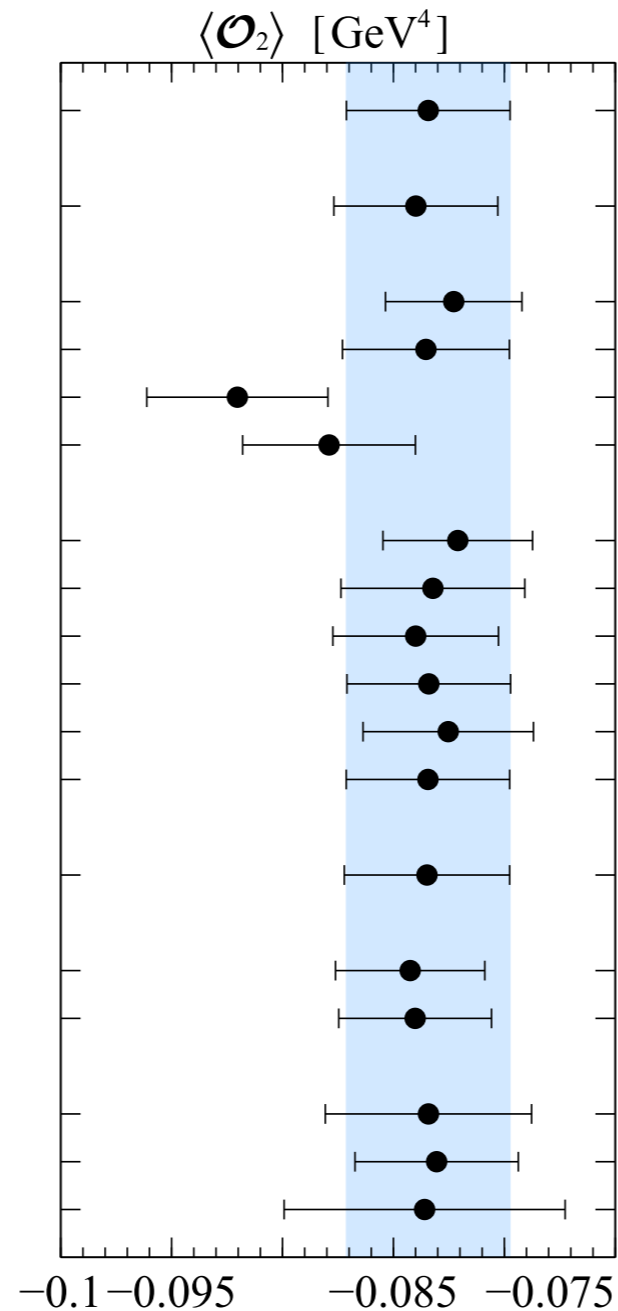


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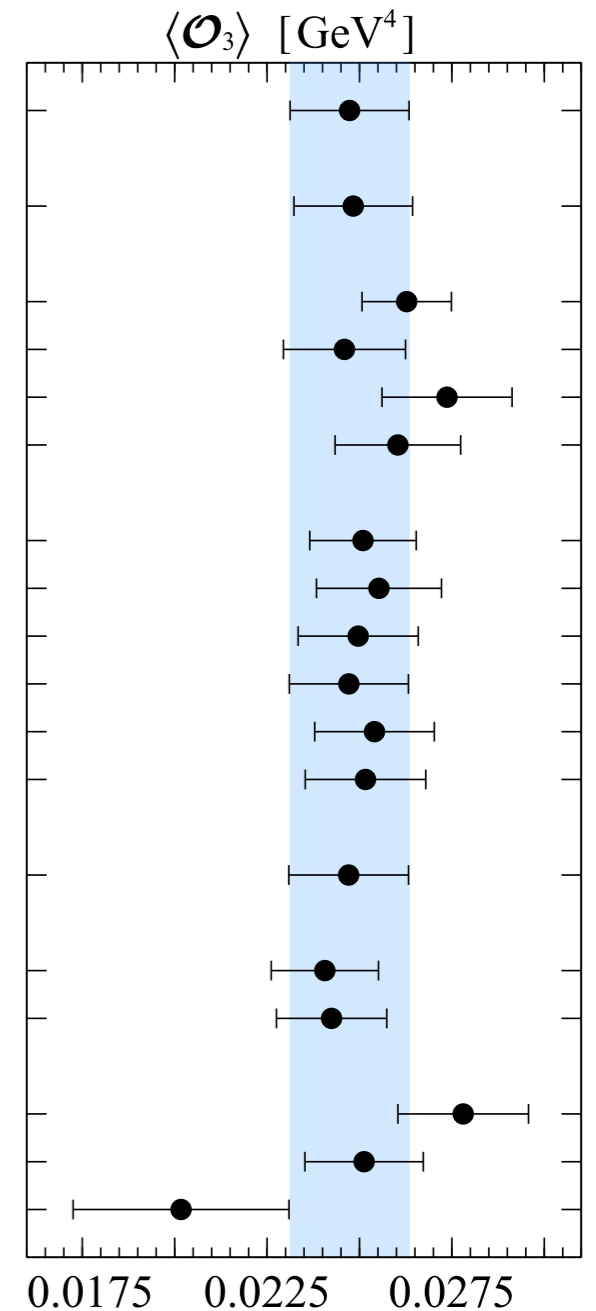
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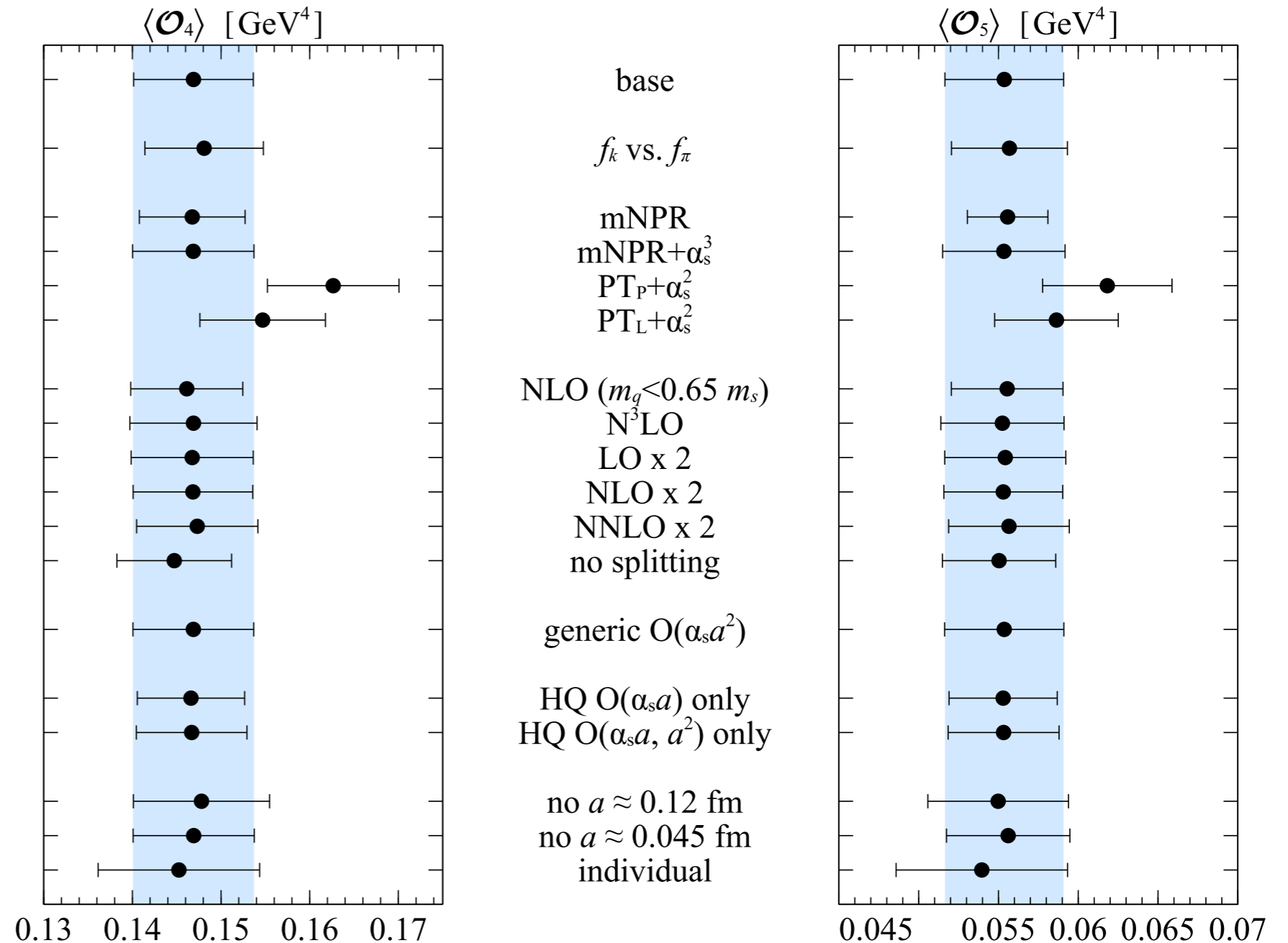


# D mixing results

FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

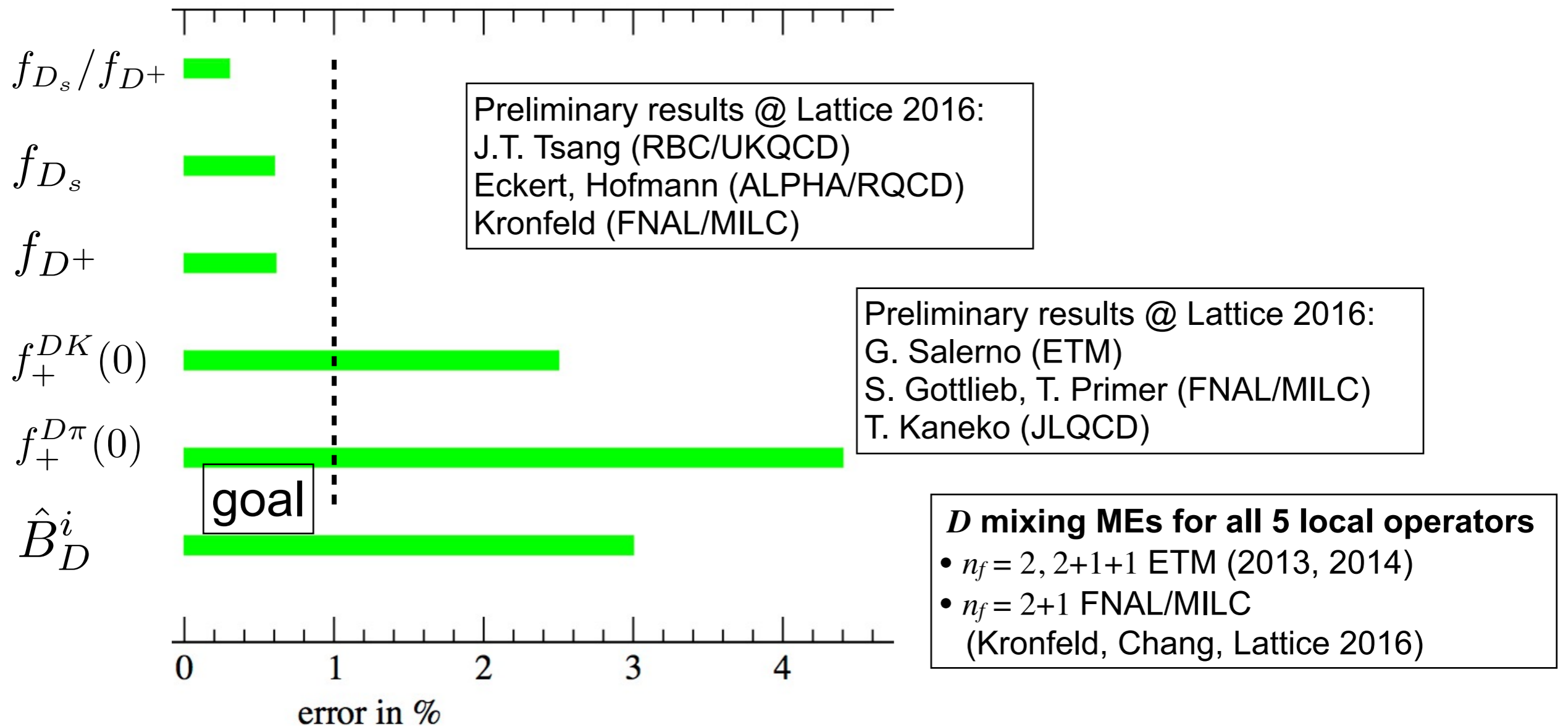
systematic error study

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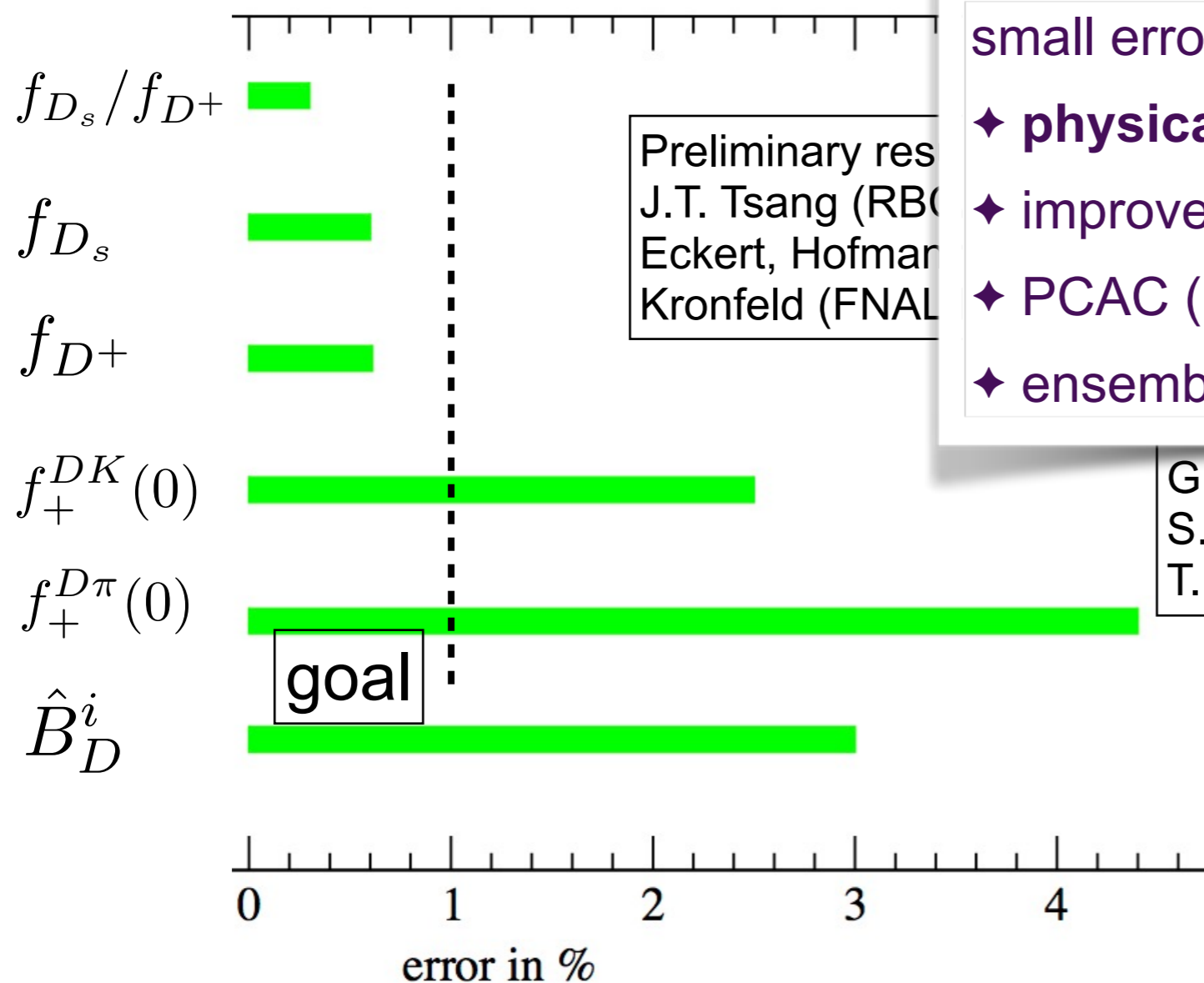
# D meson summary

errors (in %) comparison:



# D meson summary

errors (in %) comparison:



small errors due to

- ◆ physical light quark masses
- ◆ improved charm-quark action (HISQ)
- ◆ PCAC (no renormalization)
- ◆ ensembles with small lattice spacings

G. Salerno (ETM)  
S. Gottlieb, T. Primer (FNAL/MILC)  
T. Kaneko (JLQCD)

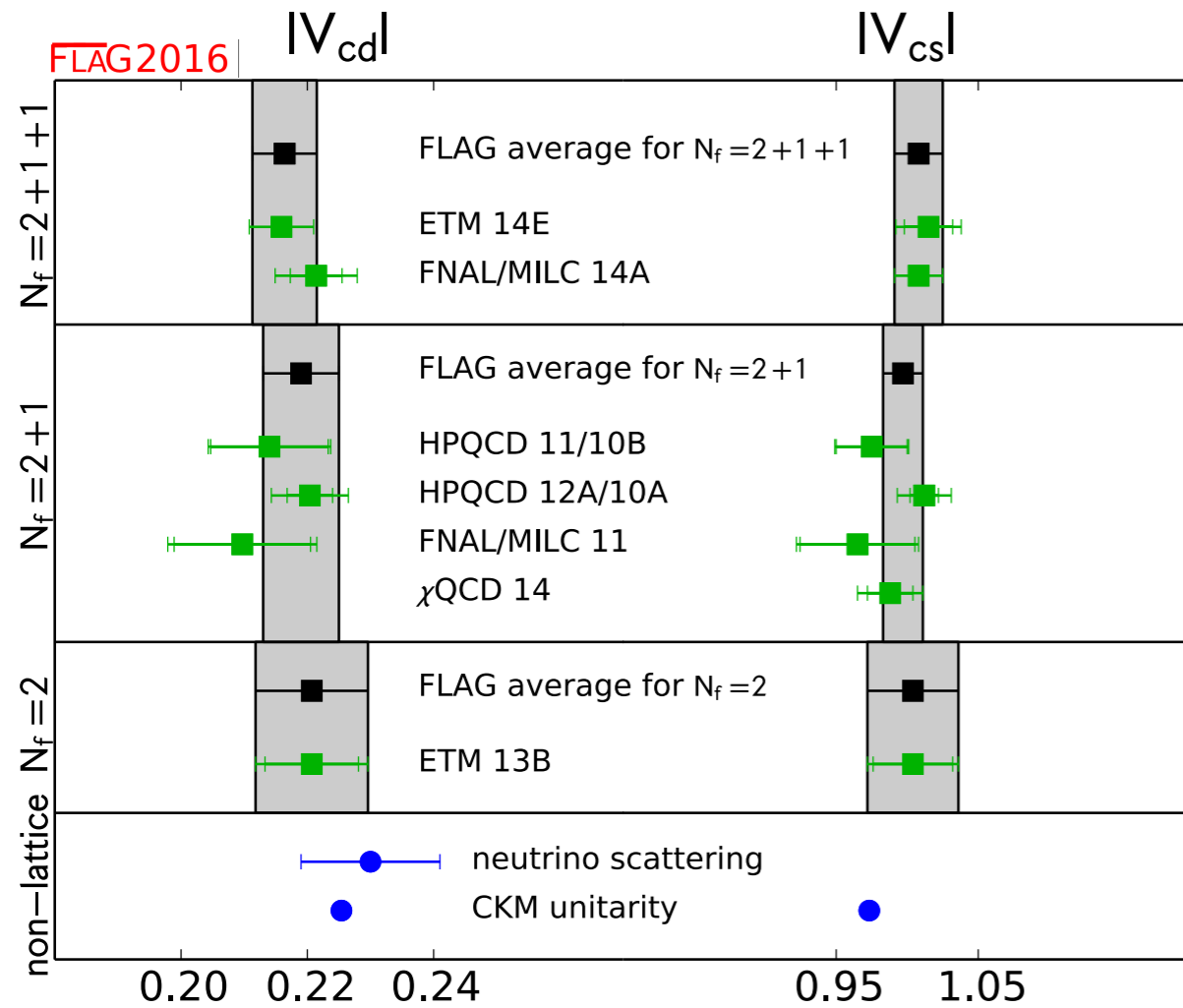
**D mixing MEs for all 5 local operators**

- $n_f = 2, 2+1+1$  ETM (2013, 2014)
- $n_f = 2+1$  FNAL/MILC  
(Kronfeld, Chang, Lattice 2016)

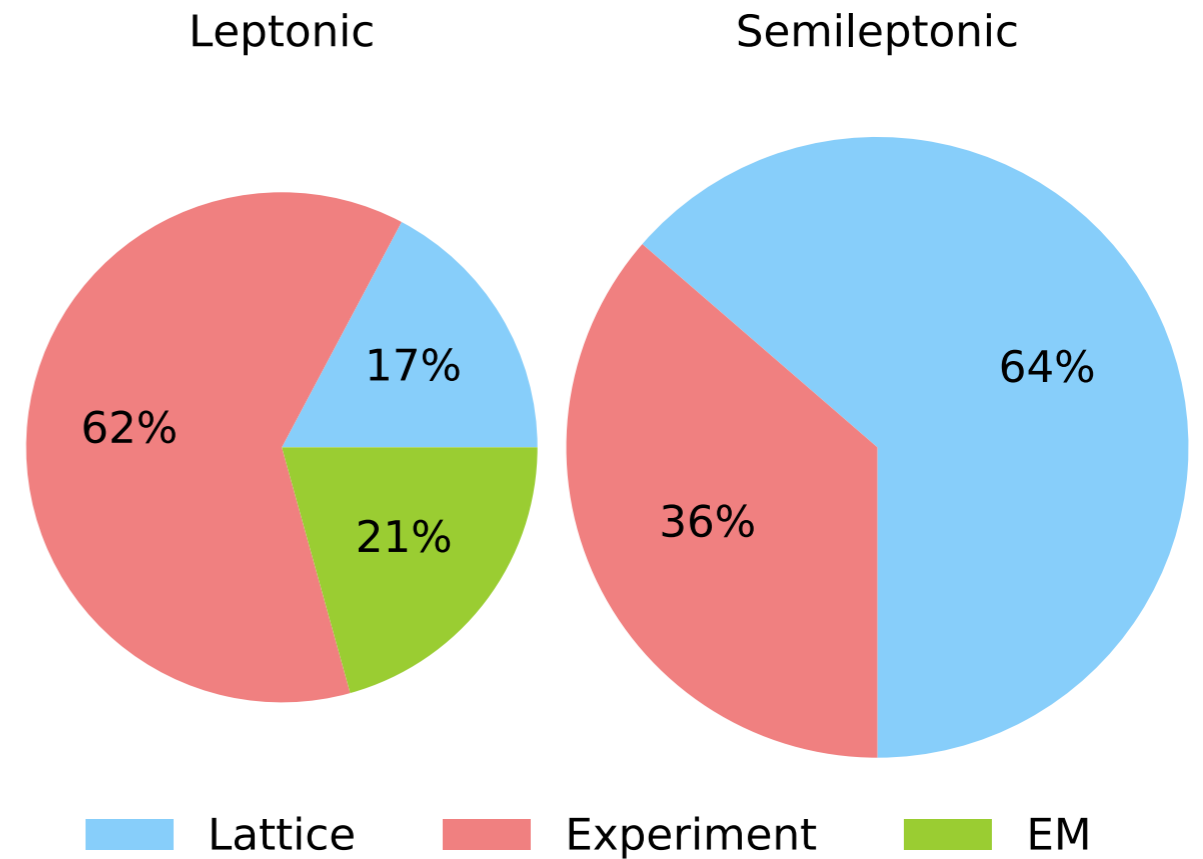
# Implications for $|V_{cs}|, |V_{cd}|$

S. Aoki et al (FLAG review, arXiv:1607.00299)

S. Gottlieb, T. Primer (FNAL/MILC) @ Lattice 2016

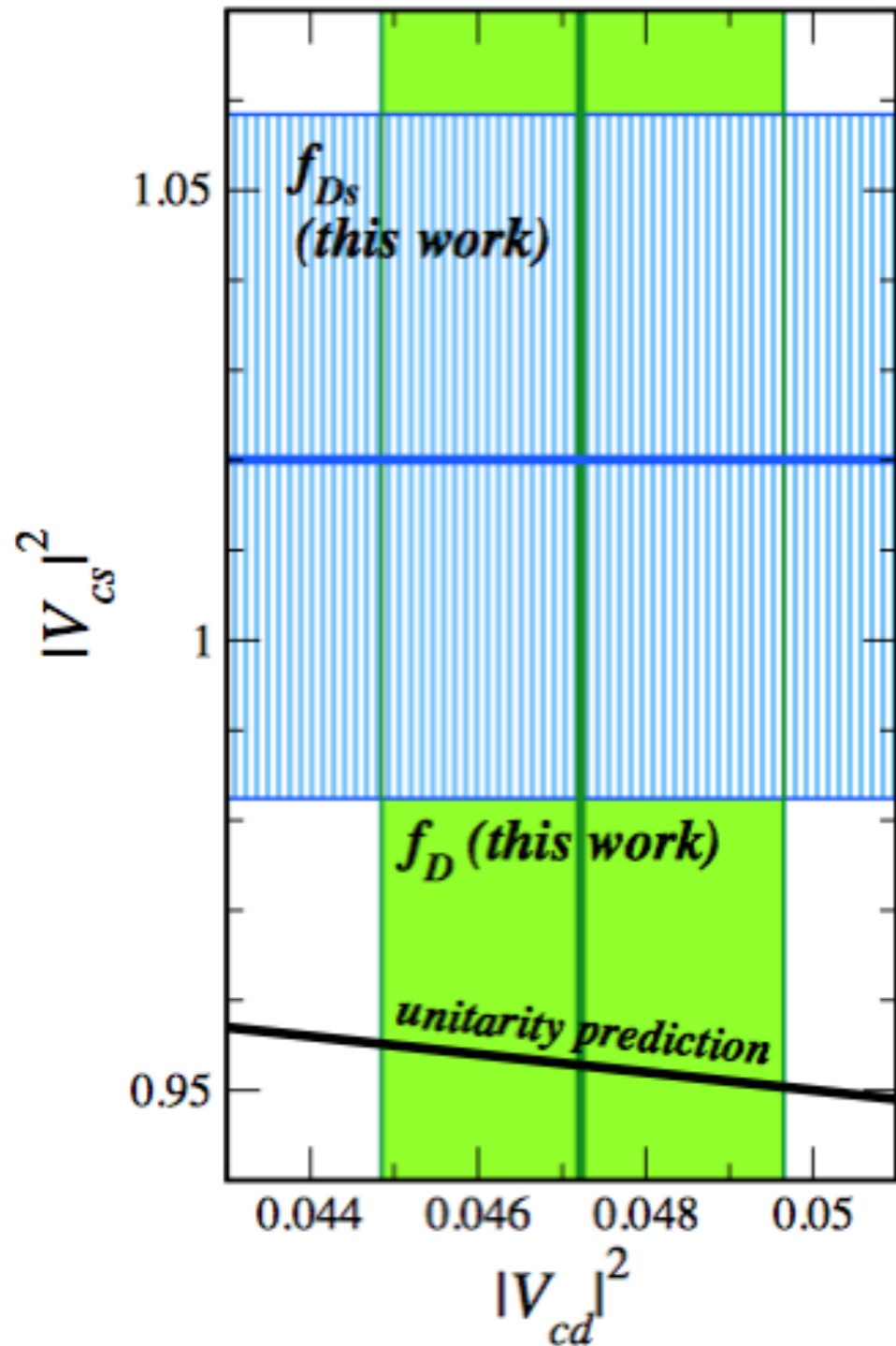


## $|V_{cs}|$ comparison



# Implications for the 2<sup>nd</sup> row of the CKM Matrix

FNAL/MILC (arXiv:1407.3772, 2014 PRD)



errors on  $|V_{cs}|$  and  $|V_{cd}|$  are dominated by experiment (PDG 2015, arXiv:509.02220):

$$\begin{aligned} |V_{cd}| &= 0.217 (1)_{\text{LQCD}} (5)_{\text{exp}} \\ |V_{cs}| &= 1.007 (4)_{\text{LQCD}} (16)_{\text{exp}} \end{aligned}$$

(based on the PDG average of 2+1 & 2+1+1 flavor LQCD results; average is dominated by FNAL/MILC)

2 $\sigma$  tension with unitarity:

$$|V_{cs}|^2 + |V_{cd}|^2 + |V_{cb}|^2 - 1 = 0.064(32)$$

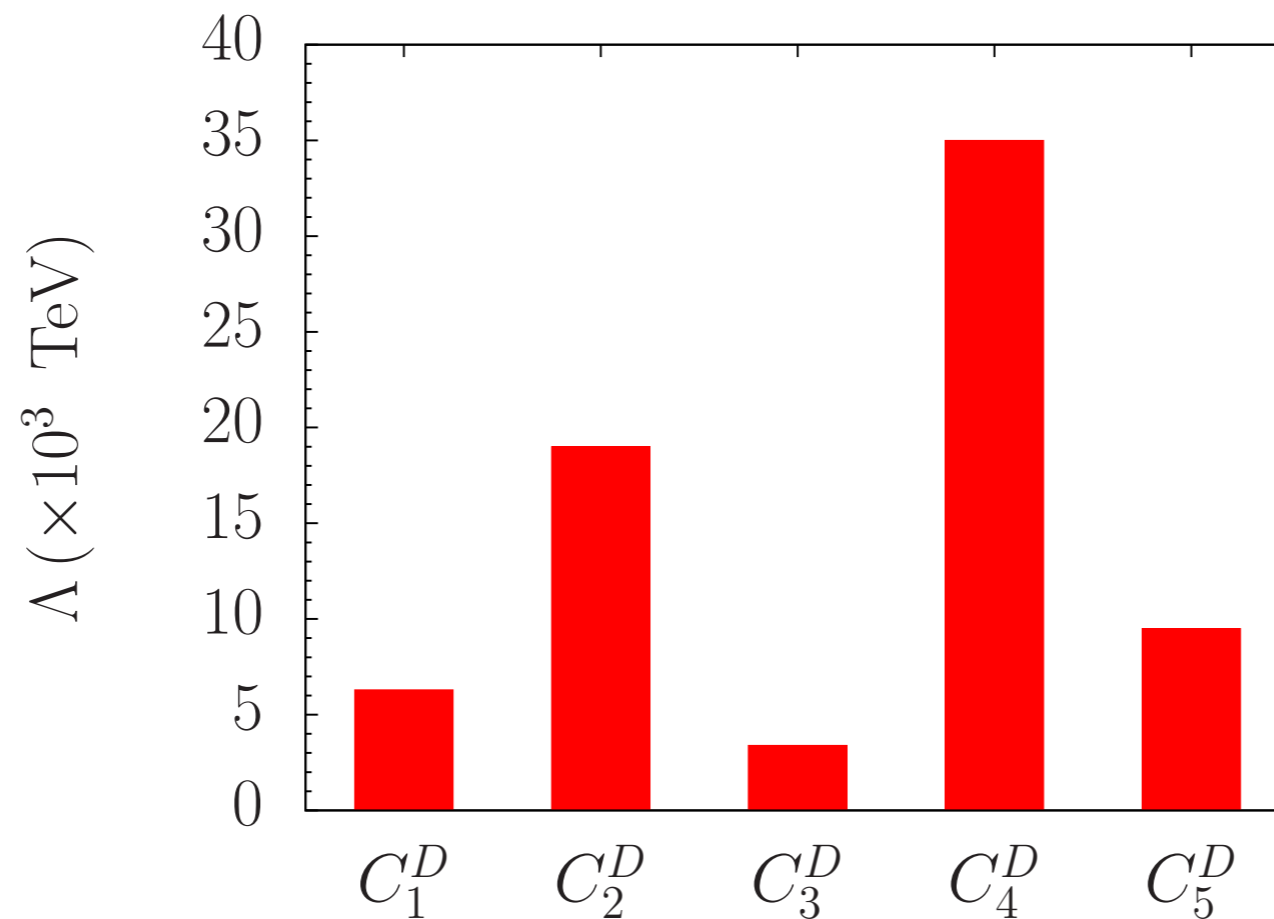


# Lower bounds on New Physics

ETM (arXiv:1403.7302, 2014 PRD)

Wilson coefficients:  $C_i(\Lambda) \sim \frac{F_i L_i}{\Lambda^2}$

Generic tree-level, strongly interacting:  $L_i \sim F_i \sim 1$



# Summary

---

- ★ Gauge field ensembles with light sea quarks at their **physical masses** are being used in a growing number of LQCD calculations of  $D$  meson quantities
  - ▣▣▣▣ removes chiral extrapolation errors ▣▣▣▣ better precision
- ★ LQCD results for  $D$ ,  $D_s$  decay constants are already very precise ( $\sim 0.5\%$  errors)
  - ▣▣▣▣ uncertainties in CKM determinations are dominated by experimental contributions
  - slight ( $2\sigma$ ) tension with 2nd row unitarity
- ★ For  $D$  semileptonic form factors, LQCD calculations still need better precision
  - goal:  $\sim 1\%$  errors
  - ▣▣▣▣ need to calculate the form factors over the entire recoil range
  - work in progress by several lattice groups
  - extension to FCNC form factors ( $f_T$ ) straightforward
- ★ For neutral  $D$  meson mixing there are now two independent LQCD calculations of the matrix elements of the full set of five local operators.
  - Further improvements are not needed until there are reliable predictions of the long-distance contributions.
- ★ For semileptonic decays into vector meson final states, the finite volume formalism has recently been developed (Briceño et al, arXiv:1406.5965, 2015 PRD)
  - Pilot studies for  $B \rightarrow K^*$  are underway

# Outlook



Amala Willenbrock

# Outlook

How do/did we get to 1% total errors (or below)?

- ★ physical mass ensembles are essential
  - ★ small lattice spacings
  - ★ calculate renormalizations nonperturbatively
  - ★ small statistical errors (straightforward, but expensive)
  
  - ★ will eventually need to include
    - ◆ strong isospin breaking ( $m_u \neq m_d$ ) effects ✓
    - ◆ QED effects
- program being developed for kaon quantities, muon  $g-2$

Extend the reach of LQCD to include

- SL decay to vector meson final states (in progress)
  - hadronic  $D$  decays
  - long-distance contributions to neutral  $D$  mixing
- ⇒ formalism is being developed for multi-hadron states in finite volume

Already done for kaons

- ★ excited state spectra, resonances, scattering states  
(see the talks by Sinead Ryan, Graham Moir, Gavin Cheung)



**Thank you!**

# Backup slides

# Heavy Quark Treatment

---

## Relativistic Heavy Quarks - Fermilab formulation

- start with the relativistic Wilson action +  $O(a)$  improvement
- with mass-dependent matching conditions, cut-off effects are

$$\alpha_s^k f(m_h a) (a\Lambda)^n \text{ with}$$

$$am_h \sim 1 : f(m_h a) \sim O(1)$$

## FNAL/MILC implementation for action and currents:

tree-level tadpole  $O(a)$  improved  
mostly nonperturbative renormalization (mNPR)

# Heavy Quark Treatment

---

## HISQ action for charm:

- like asqtad, the HISQ action is a tree-level tadpole improved staggered action, with discretization errors for light quarks:

$$\alpha_s (a\Lambda)^2, (a\Lambda)^4$$

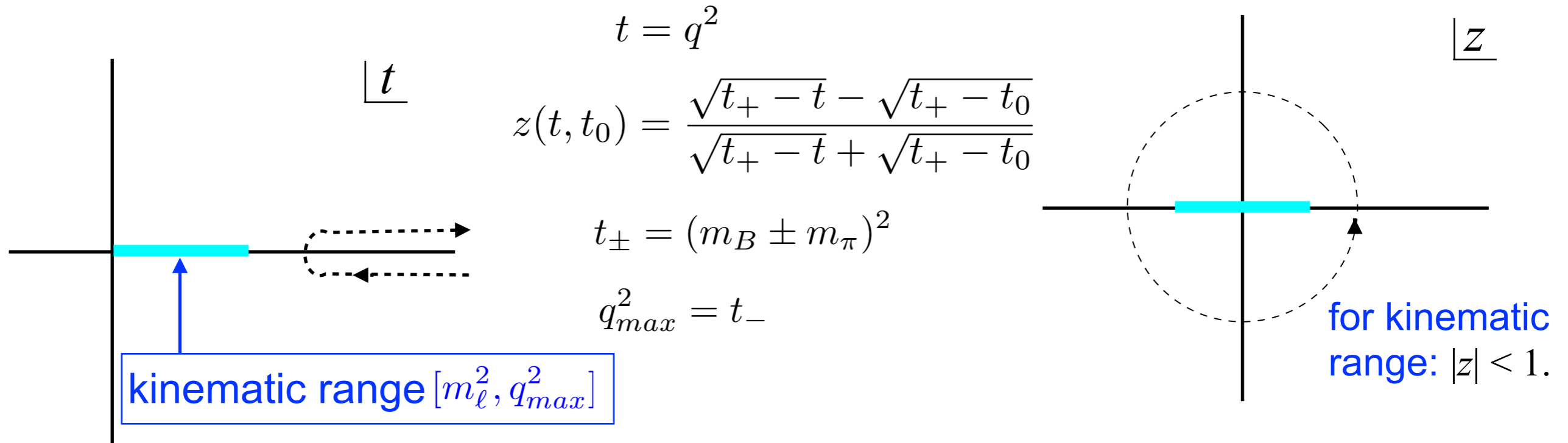
- HISQ action is highly improved for charm quarks:

$$\sim \alpha_s \Lambda/m_h (am_h)^2, (\Lambda/m_h)^2 (am_h)^4$$

- can also be used for heavier than charm quarks



# The $z$ -expansion



The form factor can be expanded as:

$$f(t) = \frac{1}{P(t)\phi(t, t_0)} \sum_{k=0} a_k(t_0) z(t, t_0)^k$$

Bourelly et al (Nucl.Phys. B189 (1981) 157)  
 Boyd et al (hep-ph/9412324, PRL 95)  
 Lellouch (arXiv:hep-ph/9509358, NPB 96)  
 Boyd & Savage (hep-ph/9702300, PRD 97)  
 Bourelly et al ( arXiv:0807.2722, PRD 09)

- $P(t)$  removes poles in  $[t_-, t_+]$
- The choice of outer function  $\phi$  affects the unitarity bound on the  $a_k$ .
- In practice, only first few terms in expansion are needed.

# chiral-continuum extrapolation

---

Some ensembles still have

$$m_{\text{light}} > 1/2 (m_u + m_d)_{\text{phys}}$$

$\chi^{\text{PT}}$  guides the extrapolation/interpolation to the physical point.

- include (light quark) discretization effects (for example, staggered  $\chi^{\text{PT}}$ )
- combined continuum-chiral extrapolation
- Heavy meson  $\chi^{\text{PT}}$ :  $\chi^{\text{PT}} + 1/M$  expansion
- can also add HQ discretization terms to chiral-continuum fits

# $D$ mixing results

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FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

chiral-continuum extrapolation

SU(3) heavy-meson partially-quenched rooted staggered  $\chi$ PT

- NLO chiral logs + staggered discretization corrections
- + analytic terms (up to N<sup>3</sup>LO)
- + leading  $1/M$  terms in HM expansion
- + HQ discretization terms
- + higher order PT terms (up to  $O(\alpha_s)^3$ )

# $D$ mixing results

FNAL/MILC (Kronfeld, Chang @ Lattice 2016)

SU(3) heavy-meson partially-quenched rooted staggered  $\chi$ PT

- NLO chiral logs + taste-splittings + “wrong-spin” corrections
- + analytic terms (up to N<sup>3</sup>LO)
- +  $B$ -meson hyperfine and flavor splittings
- + HQ discretization terms
- + higher order PT terms (up to  $O(\alpha_s)^3$ )

Schematically

$$\langle O_1^q \rangle = \beta_1 \left( 1 + \text{NLO chiral logs + taste-splittings} + \text{wrong spin terms} \right) + (2\beta_2 + 2\beta_3) \text{w.s.} + (2\beta'_2 + 2\beta'_3) \text{w.s.}$$

+ analytic terms

LECs for  $\langle O_1 \rangle, \langle O_2 \rangle, \langle O_3 \rangle$

C. Bernard (Phys.Rev. D87 (2013) 114503, arXiv: 1303.0435)

- no new LECs with simultaneous fits to the operators that mix at NLO  
 $[\langle O_1 \rangle, \langle O_2 \rangle, \langle O_3 \rangle]$  and  $[\langle O_4 \rangle, \langle O_5 \rangle]$