

Including resummation in the NLO BK equation

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Outline

Outline of the talk:

- ▶ NLO BK equation
- ▶ Numerical result: $\ln r$ divergence
- ▶ Double and single log resummation
- ▶ Numerical results: resummation only and resummation + finite terms

Literature

- ▶ “Next-to-leading order evolution of color dipoles,” I. Balitsky and G. A. Chirilli, *Phys. Rev. D* **77** (2008) 014019, [[arXiv:0710.4330](#) [[hep-ph](#)]] .
- ▶ “Direct numerical solution of the coordinate space Balitsky-Kovchegov equation at next to leading order,” T. L., **H. Mäntysaari**,
Phys. Rev. D **91** (2015) 074016, [[arXiv:1502.02400](#) [[hep-ph](#)]]
- ▶ “Resumming double logarithms in the QCD evolution of color dipoles,”
E. Iancu et al *Phys. Lett. B* **744** (2015) 293 , [[arXiv:1502.05642](#) [[hep-ph](#)]] .
- ▶ “Collinearly-improved BK evolution meets the HERA data,”
E. Iancu et al *Phys. Lett. B* **750** (2015) 643 , [[arXiv:1507.03651](#) [[hep-ph](#)]] .
- ▶ “Next-to-leading order Balitsky-Kovchegov equation with resummation,”
T. L., **H. Mäntysaari**,
Phys. Rev. D **93** (2016) 094004, [[arXiv:1601.06598](#) [[hep-ph](#)]] .

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Motivation

Many ingredients available for NLO small- x calculations:

- ▶ NLO BK equation
- ▶ NLO JIMWLK equation
- ▶ NLO γ^* impact factor for DIS
- ▶ NLO single inclusive cross section for forward pA
- ▶ ...

Armed with these, want phenomenology @ NLO!

But first need to solve the evolution equation(s)!

The NLO BK equation

as derived by Balitsky and Chirilli, 2007

Equation

$$\partial_y S(r) = \frac{\alpha_s N_c}{2\pi^2} \mathbf{K}_1 \otimes [S(X)S(Y) - S(r)] + \frac{\alpha_s^2 N_F N_c}{8\pi^4} \mathbf{K}_r \otimes S(Y)[S(X') - S(X)] \\ + \frac{\alpha_s^2 N_c^2}{8\pi^4} \mathbf{K}_2 \otimes [S(X)S(z - z')S(Y') - S(X)S(Y)]$$

Notations & approximations

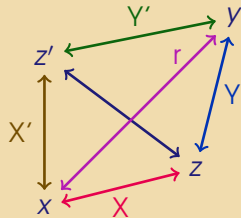
$$S(x - y) \equiv \langle \text{Tr } U^\dagger(x) U(y) \rangle$$

$$\otimes = \int d^2z \quad / \quad \int d^2z d^2z'$$

► Here large N_c & mean field:

$$\langle \text{Tr } U^\dagger U \text{Tr } U^\dagger U \rangle \rightarrow \langle \text{Tr } U^\dagger U \rangle \langle \text{Tr } U^\dagger U \rangle$$

Coordinates



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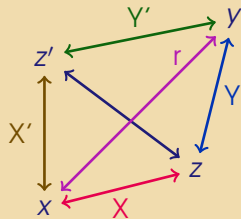
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Coordinates



Kernels

$$K_1 = \frac{r^2}{X^2 Y^2} \left[1 + \frac{\alpha_s N_C}{4\pi} \left(\frac{\beta}{N_C} \ln r^2 \mu^2 - \frac{\beta}{N_C} \frac{X^2 - Y^2}{r^2} \ln \frac{X^2}{Y^2} + \frac{67}{9} - \frac{\pi^2}{3} - \frac{10}{9} \frac{N_F}{N_C} - 2 \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2} \right) \right]$$

$$K_2 = -\frac{2}{(z-z')^4} + \left[\frac{X^2 Y'^2 + X'^2 Y^2 - 4r^2(z-z')^2}{(z-z')^4 (X^2 Y'^2 - X'^2 Y^2)} + \frac{r^4}{X^2 Y'^2 (X^2 Y'^2 - X'^2 Y^2)} + \frac{r^2}{X^2 Y'^2 (z-z')^2} \right] \ln \frac{X^2 Y'^2}{X'^2 Y^2}$$

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► Leading order

Kernels

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- ▶ Leading order
- ▶ Running coupling (Terms with β function coefficient)

Kernels

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- ▶ Conformal logs \implies vanish for $r = 0$ ($X = Y$ & $X' = Y'$)

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- ▶ Leading order
- ▶ Running coupling (Terms with β function coefficient)
- ▶ Conformal logs \implies vanish for $r = 0$ ($X = Y$ & $X' = Y'$)
- ▶ Nonconformal double log \implies blows up for $r = 0$

Running coupling

Absorb the β -terms into

- ▶ “Balitsky” running for LO term
- ▶ Parent dipole running for NLO terms

Now:

$$\frac{\alpha_s N_c}{2\pi^2} K_1 = \frac{\alpha_s(r) N_c}{2\pi^2} \left[\frac{r^2}{X^2 Y^2} + \frac{1}{X^2} \left(\frac{\alpha_s(X)}{\alpha_s(Y)} - 1 \right) + \frac{1}{Y^2} \left(\frac{\alpha_s(Y)}{\alpha_s(X)} - 1 \right) \right] \\ + \frac{\alpha_s(r)^2 N_c^2}{8\pi^3} \frac{r^2}{X^2 Y^2} \left[\frac{67}{9} - \frac{\pi^2}{3} - \frac{10 N_F}{9 N_c} - 2 \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2} \right]$$

Freeze coupling in the IR to $\alpha_s(r \rightarrow \infty) = 0.76$:

$$\alpha_s(r) = \frac{4\pi}{\beta \ln \left\{ \left[(2.5)^{10} + \left((4e^{-2\gamma_E}) / (r^2 \Lambda_{\text{QCD}}^2) \right)^5 \right]^{0.2} \right\}} \sim \frac{4\pi}{\beta \ln \left\{ \left(\frac{4e^{-2\gamma_E}}{r^2 \Lambda_{\text{QCD}}^2} \right) \right\}}$$

Initial condition

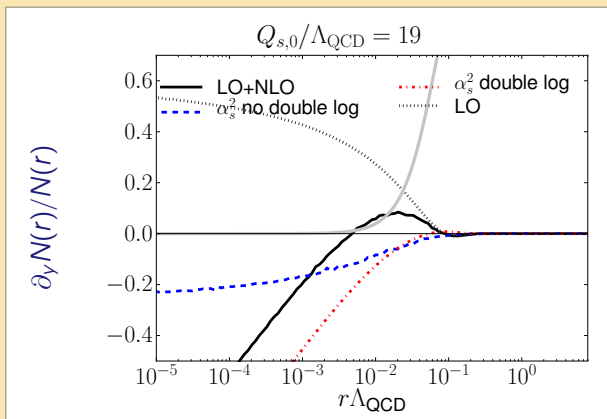
$$N(r) \equiv 1 - S(r) = 1 - \exp \left[-\frac{(r^2 Q_{s0}^2)^\gamma}{4} \ln \left(\frac{1}{r\Lambda_{\text{QCD}}} + e \right) \right],$$

2 tunable parameters

- ▶ $\frac{Q_{s0}}{\Lambda_{\text{QCD}}} \implies$ basically determines value of α_s at $y = 0$
- ▶ γ : anomalous dimension: shape $N(r) \sim r^{2\gamma}$
 - ▶ LO phenomenology prefers $\gamma \gtrsim 1$ at $y = 0$
 - ▶ At LO this eventually evolves into $\gamma \sim 0.8$ (running α_s)

Large double log

Relative change of $N(r)$ in one step dy diverges for small r



Behavior caused by the nonconformal double log term

T.L., H. Mäntysaari 2015

Double log resummation

Integral form of BK: double log \sim shift of rapidity variable

$$\partial_Y S_Y(r) = \alpha_s \int_z \frac{r^2}{X^2 Y^2} \left[1 + \frac{\alpha_s N_c}{2\pi} \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2} \right] [S_Y(X) S_Y(Y) - S_Y(r)]$$

\Leftrightarrow

$$S_Y(r) = \alpha_s \int_z^Y dy' \int_z \frac{r^2}{X^2 Y^2} \left[1 + \frac{\alpha_s N_c}{2\pi} \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2} \right] [S_{Y'}(X) S_{Y'}(Y) - S_{Y'}(r)]$$

\Leftrightarrow

$$S_Y(r) = \alpha_s \int d^2z \int^{y - \ln z^2 / r^2} dy' \frac{r^2}{X^2 Y^2} [S_{Y'}(X) S_{Y'}(Y) - S_{Y'}(r)]$$

For $r \ll z \sim 1/Q_s$. \Rightarrow "Kinematical constraint" Beuf

Iancu et al 2015 rewrite as practical rapidity-local form

$$\left[1 + \frac{\alpha_s N_c}{2\pi} \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2} \right] \Rightarrow \frac{J_1 \left(2\sqrt{\frac{\alpha_s N_c}{\pi} \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2}} \right)}{\sqrt{\frac{\alpha_s N_c}{\pi} \ln \frac{X^2}{r^2} \ln \frac{Y^2}{r^2}}} \quad \text{with} \quad \int^y dy'$$

Single log resummation

lanou et al 2015 : there is also a single log,
with DGLAP anomalous dimension $A_1 = 11/12$:

$$\begin{aligned}\partial_Y S(r) &= \alpha_s^2 \int_{z'} \left[-\frac{2}{(z-z')^4} + \frac{X^2 Y'^2 + X'^2 Y^2 - 4r^2(z-z')^2}{(z-z')^4(X^2 Y'^2 - X'^2 Y^2)} \ln \frac{X^2 Y'^2}{X'^2 Y^2} \right] \\ &\quad \times [S(X)S(z-z')S(Y') - S(X)S(Y)] + \dots \\ &\sim \alpha_s^2 A_1 \int_z \frac{r^2}{X^2 Y^2} \ln \frac{\min\{X^2, Y^2\}}{r^2} + \dots\end{aligned}$$

Should be resummed into

$$K_{STL} = \exp \left\{ -\frac{\alpha_s N_C A_1}{\pi} \left| \ln \frac{C_{\text{sub}} r^2}{\min\{X^2, Y^2\}} r^2 \right| \right\}$$

- ▶ lanou et al take $C_{\text{sub}} = 1$ and throw other NLO terms away
- ▶ This captures leading logs, but leaves other α_s^2 -terms unknown, and result dependent on arbitrary C_{sub}

Full equation

$$\begin{aligned}\partial_Y S(r) = & \frac{\alpha_s N_c}{2\pi^2} K_1 \otimes [S(X)S(Y) - S(r)] \\ & + \frac{\alpha_s^2 N_c^2}{8\pi^4} K_2 \otimes [S(X)S(z-z')S(Y') - S(X)S(Y)] + N_F\text{-part}\end{aligned}$$

Now solve resummed equation

$$\begin{aligned}\partial_Y S(r) = & \frac{\alpha_s N_c}{2\pi^2} \left[K_{DLA} K_{STL} K_{Bal} - K_{sub} + K_1^{\text{fin}} \right] \otimes [S(X)S(Y) - S(r)] \\ & + \frac{\alpha_s^2 N_c^2}{8\pi^4} K_2 \otimes [S(X)S(z-z')S(Y') - S(X)S(Y)] + N_F\text{-part}\end{aligned}$$

- ▶ $K_{DLA} \sim J_1(x)/x$ would give double log in K_1 if expanded in α_s
- ▶ K_{Bal} resums β -function terms in K_1
- ▶ $K_{STL} = \exp \{ -\alpha_s N_c A_1 \# \ln r^2 \}$ resums single log
- ▶ K_{sub} subtracts α_s -part of K_{STL} , which is already in K_2
- ▶ K_1^{fin} is rest (nonlog) part in K_1

Resummation and rest

$$\begin{aligned}\partial_Y S(r) = & \frac{\alpha_s N_c}{2\pi^2} \left[K_{DLA} K_{STL} K_{Bal} - K_{sub} + K_1^{fin} \right] \otimes [S(X)S(Y) - S(r)] \\ & + \frac{\alpha_s^2 N_c^2}{8\pi^4} K_2 \otimes [S(X)S(z-z')S(Y') - S(X)S(Y)] + N_F\text{-part}\end{aligned}$$

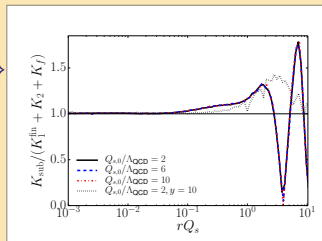
Split this into three parts

LO (running coupling) $\frac{\alpha_s N_c}{2\pi^2} K_{Bal}$: used so far

Resummation $\frac{\alpha_s N_c}{2\pi^2} K_{Bal} [K_{DLA} K_{STL} - 1]$: (\sim used by Iancu et al)

Finite NLO rest: $-K_{sub} + K_1^{fin}$ and K_2, K_f

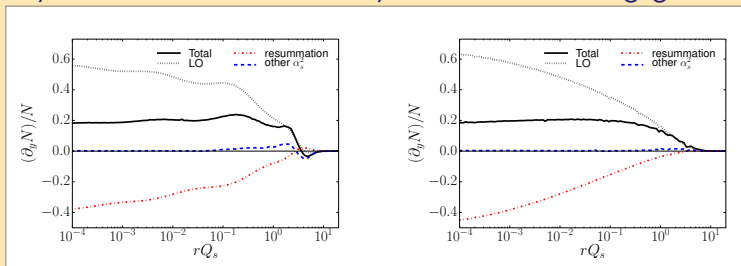
- ▶ Separation depends on C_{sub} in $K_{STL} = \exp \left\{ -\frac{\alpha_s N_c A_1}{\pi} \left| \ln \frac{C_{sub} r^2}{\min\{X^2, Y^2\}} r^2 \right| \right\}$
- ▶ Numerically **choose** C_{sub} to **minimize** the finite NLO part:
result $C_{sub} = 0.65$



Contribution to evolution speed

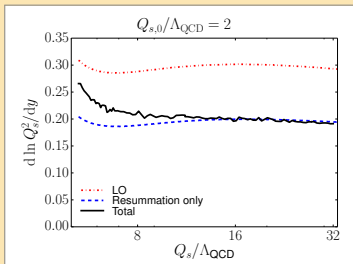
$y=0$: finite NLO small

$y=10$: finite NLO negligible



Finite NLO terms

- ▶ Significant for small y
 \Rightarrow phenomenology
- ▶ Negligible at $y \rightarrow \infty$



Anomalous dimension

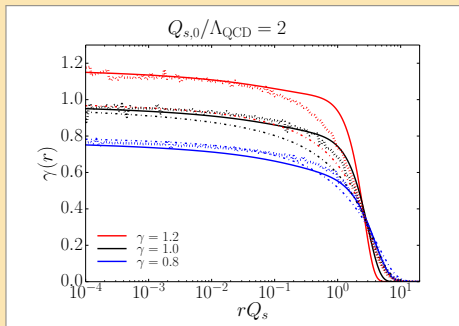
Recall initial condition: $N(r) = 1 - e^{-\frac{(r^2 Q_s^2)^\gamma}{4}} \ln\left(\frac{1}{r\Lambda_{\text{QCD}}} + e\right)$,

Define

$$\gamma(r) \equiv -\frac{d \ln N(r)}{d \ln r^2}$$

Geometric scaling?

- ▶ LO: fast to $\gamma \sim 0.8$
- ▶ NLO: stay at initial γ



- ▶ Solid: initial condition
- ▶ Dotted: $y = 5$ NLO
- ▶ Dot-dashed: $y = 5$ LO (rc)

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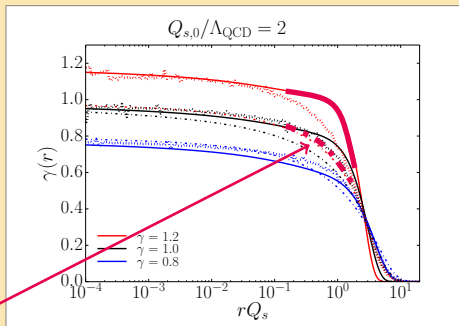
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LO $y = 0$ to $y = 5$

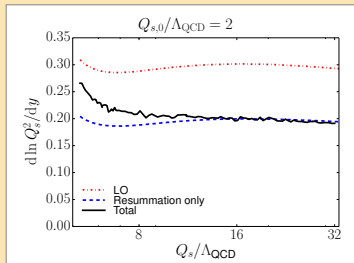


- ▶ Solid: initial condition
- ▶ Dotted: $y = 5$ NLO
- ▶ Dot-dashed: $y = 5$ LO (rc)

Conclusions

Numerical solution of NLO BK equation

- ▶ Resummation
 - ▶ Resummation of double \perp logs stabilizes equation
 - ▶ Can also resum single \perp logs, ambiguity in constant C_{sub}
 - ▶ Numerically: $C_{\text{sub}} = 0.65$ minimizes finite NLO
- ▶ Properties of solution
 - ▶ Finite NLO terms important for realistic Q_s/Λ_{QCD}
 - ▶ Geometric scaling: sets in slowly if at all



Outlook:

- ▶ Fit DIS
- ▶ Apply to pA