Heavy Ion Collisions with General Purpose Event Generators

1607.04434 [hep-ph]

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Multiparton Interactions, from pp to pA

- Comparison of HI data to scaled *pp* is common but challenging.
- MPI models not generalized to heavy ion environments.
- Centrality observables are forward $\rightarrow$ diffraction.
- Goal: Tune all physics to e^+e^- and *pp*; introduce only nuclear geometry.
- *Necessary* starting point for microscopic models for collectivity.
- This talk:
  1. Beyond *pp*: fluctuations in Glauber model(s).
  2. Glauber-Gribov Colour Fluctuations and DIPSY.
  3. Particle production and FritiofP8.
  4. Outlook.

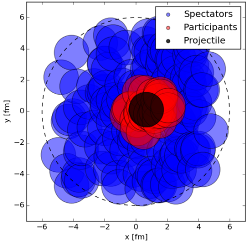
From pp to pA

- Standard Glauber approach: interaction through absorptive channels.
- What do we need to reproduce "centrality" $\propto$ forward particle production?
- Wounded nucleons updated to include fluctuations in target and projectile (SD + DD).
- Optical theorem in impact parameter space:

$$\Im(A_{el}) = \frac{1}{2}(|A_{el}|^2 + P_{abs}); \, T \equiv -iA_{el} \Rightarrow$$

$$\frac{d\sigma_{el}}{d^2b} = \langle T(b) \rangle^2, \, \frac{d\sigma_{tot}}{d^2b} = 2 \langle T(b) \rangle$$

$$\frac{d\sigma_{abs}}{d^2b} = 2 \langle T(b) \rangle - \langle T(b) \rangle^2$$



- No fluctuations! $T(b) = \Theta\left(\sqrt{\sigma_{abs}/\pi} - b\right)$

The wounded cross section

- Fluctuations related to diffractive excitations: Good-Walker.

$$\frac{d\sigma_{tot}}{d^2b} = 2\left\langle T \right\rangle_{t,p}, \; \frac{d\sigma_{el}}{d^2b} = \left\langle T \right\rangle_{t,p}^2, \; \frac{d\sigma_{SD,(p|t)}}{d^2b} = \left\langle \left\langle T \right\rangle_{(t|p)}^2 \right\rangle_{(p|t)} - \left\langle T \right\rangle_{p,t}^2$$

$$\frac{d\sigma_{DD}}{d^2b} = \left\langle T^2 \right\rangle_{p,t} - \left\langle \left\langle T \right\rangle_t^2 \right\rangle_p - \left\langle \left\langle T \right\rangle_p^2 \right\rangle_t + \left\langle T \right\rangle_{p,t}^2$$

- The *wounded* cross section is the sum of:

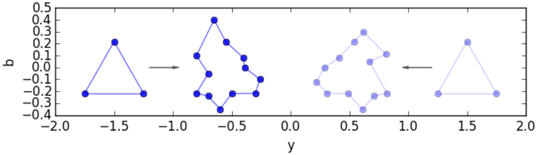
$$\frac{d\sigma_w}{d^2b} = \frac{d\sigma_{abs}}{d^2b} + \frac{d\sigma_{SD,t}}{d^2b} + \frac{d\sigma_{DD}}{d^2b} = 2\left\langle T \right\rangle_{p,t} - \left\langle \left\langle T \right\rangle_t^2 \right\rangle_p .$$

- Contributions to "centrality" observable: absorptively wounded, diffractively wounded, NOT elastically scattered.
- We need now to calculate $T(b)$.

The DIPSY model Flensburg et al. arXiv:1103.4321 [hep-ph]

- We can calculate $T(b)$ in DIPSY:
  Dipole evolution in **I**mpact **P**arameter **S**pace and rapidit**Y**.
- LL-BFKL with some corrections built on Mueller dipole model [Mueller and Patel arXiv:hep-ph/9403256].
- Proton/Nucleus structure built up dynamically from dipole splittings:

$$\frac{dP}{dY} = \frac{3\alpha_s}{2\pi^2} d^2\vec{z} \frac{(\vec{x}-\vec{y})^2}{(\vec{x}-\vec{z})^2(\vec{z}-\vec{y})^2}, \; f_{ij} = \frac{\alpha_s^2}{8}\left[\log\left(\frac{(\vec{x}_i-\vec{y}_j)^2(\vec{y}_i-\vec{x}_j)^2}{(\vec{x}_i-\vec{x}_j)^2(\vec{y}_i-\vec{y}_j)^2}\right)\right]^2$$



- Optical theorem gives: $T(b) = 1 - \exp\left(-\sum_{ij} f_{ij}\right)$
- Will serve as an initial state "truth".
- Also implemented as full event generator – out of scope for this talk.

Glauber-Gribov fluctuations (GG or GGCF)

- Parametrization of cross section fluctuations in Glauber-Gribov formalism [Alvioli and Strikman: arXiv:1301.0728 [hep-ph]]:
- Parameterization of total cross section distribution:

$$\sigma_{tot} = \int d\sigma \sigma P_{tot}(\sigma) = \int d\sigma \rho \frac{\sigma^2}{\sigma + \sigma_0} \exp\left[-\frac{(\sigma/\sigma_0 - 1)^2}{\Omega^2}\right]$$

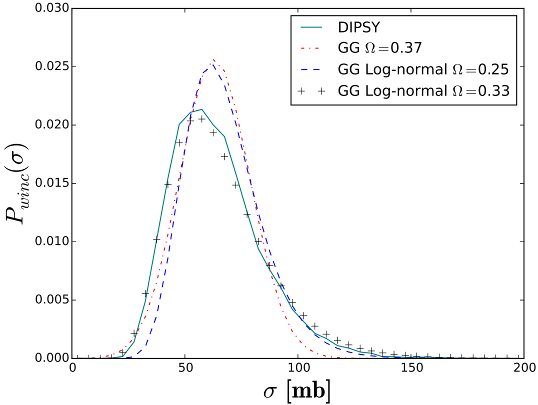
- Normal usage: With black disk, scale to total inelastic $\sigma_{in} = \lambda \sigma_{tot}$.
- From arguments above, should be σ_w
- BUT! $\sigma_{Glauber} = \sigma_w$ in GG/GGCF is not enough.
- Lack of information wrt. DIPSY, which calculates full $T(b)$.
- Assume semi-transparent disk:

$$T^{(pp)}(b, \sigma) = T_0 \Theta\left(\sqrt{\frac{\sigma}{2\pi T_0}} - b\right)$$

- Fit to semi-inclusive cross sections.
- Log-normal distribution fits DIPSY better.

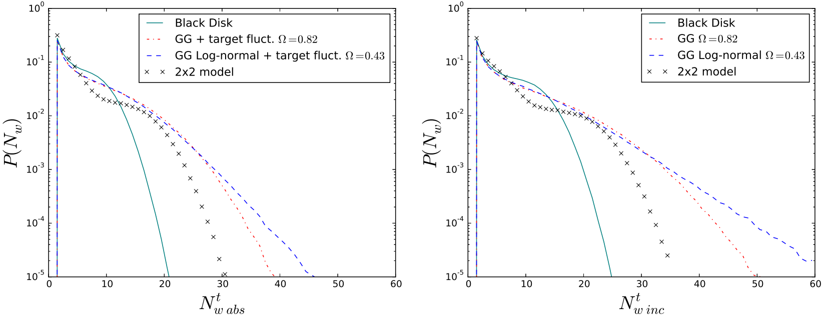
$$\sigma_{tot} = \int d^2b \int d\sigma P_{tot}(\sigma) 2T^{(pp)}(b,\sigma),\ \sigma_{el} = \int d^2b \left| \int d\sigma P_{tot} T^{(pp)}(b,\sigma) \right|^2,$$

$$\sigma_{w_{inc}} = \int d^2b \int d\sigma P_{tot}(\sigma) \left[2T^{(pp)}(b,\sigma) - T^{(pp)}(b,\sigma) \right],\ P_{tot}(\sigma, b) = \frac{1}{\Omega\sqrt{2\pi}} \exp\left(-\frac{\log^2(\sigma/\sigma_0)}{2\Omega^2} \right)$$



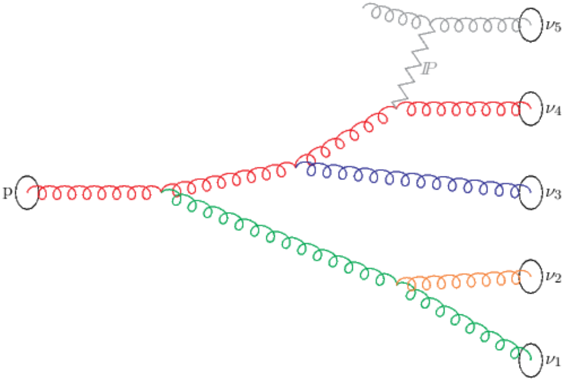
Types of wounded nucleons

- We can now fit to pp cross sections and obtain:
  1. The number of wounded nucleons inc. diffractive excitation.
  2. Given $T(b)$ assumption, which are which!
- We now have input for a model for particle production.



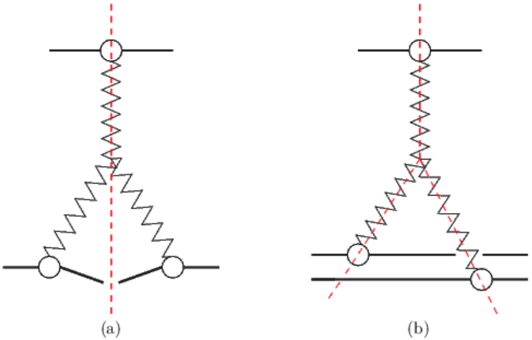
Full final states: Revival of Fritiof

- One absorptive collision contributes to full rapidity span.
- The rest contributes similarly to diffrative excitation (plus a colour exchange).
- Implementation in Pythia8 (FritiofP8), but idea is general.
- Baseline: Everything absorptive.



Similar to diffractive excitation? Really?

- Secondary absorptive interactions are similar to single diffractive ones.
- Consider cut Pomeron diagrams for:
  - (a) Single diffractive proton–proton.
  - (b) Double diffractive proton–deuteron.
- Not far fetched to assume that interactions are similarly distributed in rapidity.



Particle production: MPIs Sjöstrand and Skands: arXiv:hep-ph/0402078

- Several partons taken from the PDF.
- Hard sub-collisions with $2 \to 2$ ME:

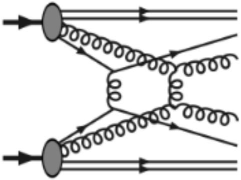


Figure T. Sjöstrand

$$\frac{d\sigma_{2\to 2}}{dp_\perp^2} \propto \frac{\alpha_s^2(p_\perp^2)}{p_\perp^4} \to \frac{\alpha_s^2(p_\perp^2 + p_{\perp 0}^2)}{(p_\perp^2 + p_{\perp 0}^2)^2}.$$

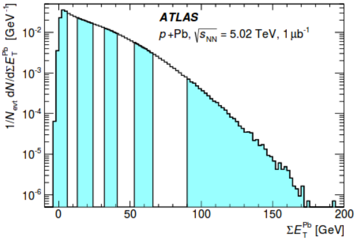
- Momentum conservation and PDF scaling.
- Ordered emissions: $p_{\perp 1} > p_{\perp 2} > p_{\perp 4} > ...$ from:

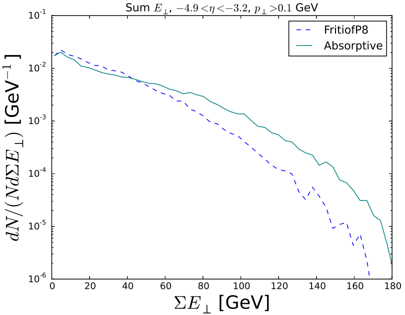
$$\mathcal{P}(p_\perp = p_{\perp i}) = \frac{1}{\sigma_{nd}} \frac{d\sigma_{2\to 2}}{dp_\perp} \exp\left[- \int_{p_\perp}^{p_{\perp i-1}} \frac{1}{\sigma_{nd}} \frac{d\sigma}{dp'_\perp} dp'_\perp \right]$$

- Number distribution narrower than Poissonian (momentum and flavour rescaling).

Results [Data: ATLAS: 1508.00848 [hep-ex]]

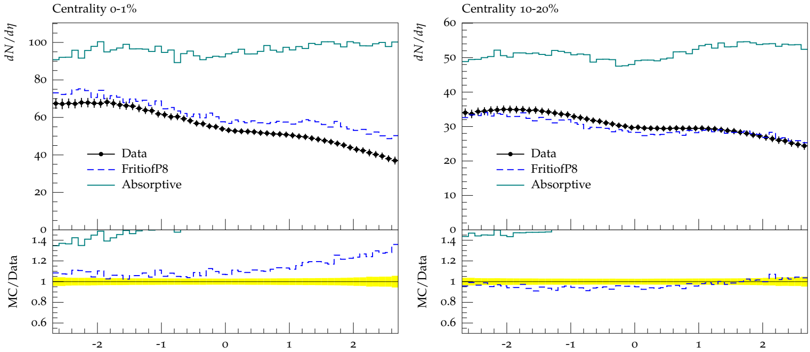
- Very good agreement with centrality observable.
- "Absorptive" overshoots.
- Measuring the exact region where diffractive excitation is important.





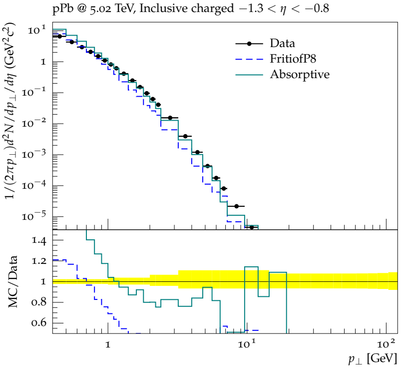
Multiplicity

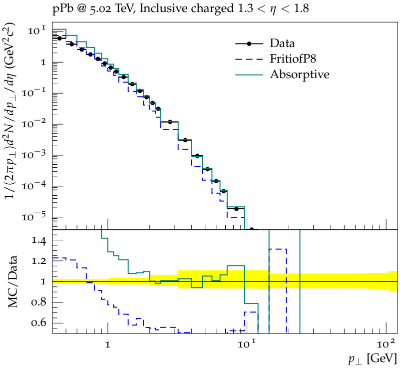
- Reproducing central collisions well.
- Comparison by own Rivet routine – implementation by exp. would be better.



Transverse momentum [Data: CMS: 1502.05287 [nucl-ex]]

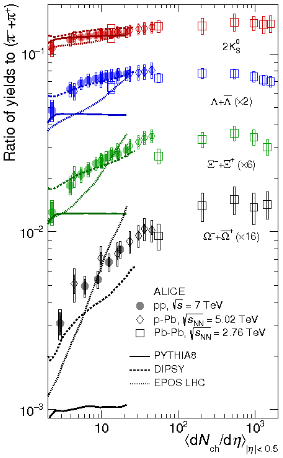
- Low-$p_\perp$ region improved from Absorptive model.
- Large uncertainties from pdf in this observable.





Outlook

- Currently: Merging with Rope Hadronization.
- Successful description of hadron yields in pp.
- Hard production in further in future.
- Framework for production already in place in Pythia.
- Important to include hard diffraction for forward production.



[From ALICE: 1606.07424 [nucl-ex]]

Conclusions

- Including diffractive excitation is important for centrality observables.
- Using *wounded* cross section + semi-transparent, fluctuating disk.

$$\frac{d\sigma_w}{d^2b} = \frac{d\sigma_{abs}}{d^2b} + \frac{d\sigma_{SD,t}}{d^2b} + \frac{d\sigma_{DD}}{d^2b}$$

- Fritiof-like model for particle production.
- Reproducing charged particle spectra well.
- Prospects for introducing medium effects from microscopic interactions.

Thank you!

Bonus slides

Color reconnection

- Many partonic subcollisions $\Rightarrow$ Many hadronizing strings.
- But! $N_c = 3$, not $N_c = \infty$ gives interactions.
- Easy to merge low-$p_\perp$ systems, hard to merge two hard-$p_\perp$.

$$\mathcal{P}_{merge} = \frac{(\gamma p_{\perp 0})^2}{(\gamma p_{\perp 0})^2 + p_\perp^2}$$

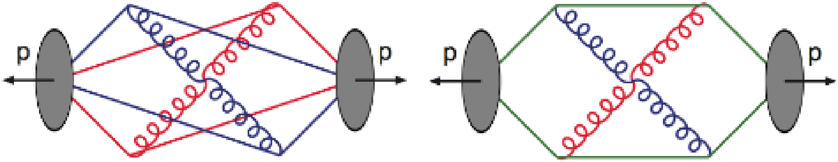


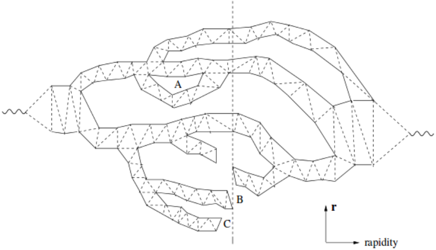
Figure T. Sjöstrand

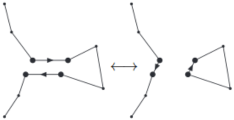
- Actual merging is decided by minimization of "potential energy":

$$\lambda = \sum_{dipoles} \log(1 + \sqrt{2}E/m_0)$$

Saturation and swings

- In DIPSY MPIs are fluctuations going on shell in interactions.
- Similar to saturation in another frame: Initial state swing.
- Multiple scatterings of a single dipole $\Leftrightarrow$ Several swings (Avsar, E.: arXiv:0709.1371 [hep-ph])
- Re-absorption of non-interacting branches.





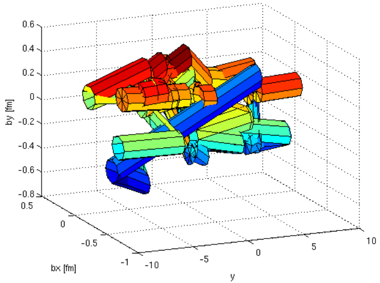
- Initial state swing competes with emission.
- All gluons get index from 1 to N_c^2, reconnect if compatible with:

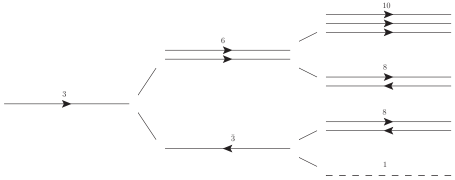
$$\frac{\mathcal{P}_{(12)(34)}}{\mathcal{P}_{(14)(32)}} = \frac{(\vec{x}_1 - \vec{x}_4)^2(\vec{x}_3 - \vec{x}_2)^2}{(\vec{x}_1 - \vec{x}_2)^2(\vec{x}_3 - \vec{x}_4)^2}.$$

Ropes, swings and junctions CB et al. arXiv:1412.6259 [hep-ph]

- Final state interactions: Many overlapping strings (like CR)

  Old in HI: Biro et al: Nucl.Phys. B245 (1984) 449-468.
- SU(3) multiplet structure decided by random walk.
- Effects implemented from perturbative (parton shower) to non-perturbative (hadronization) scales.

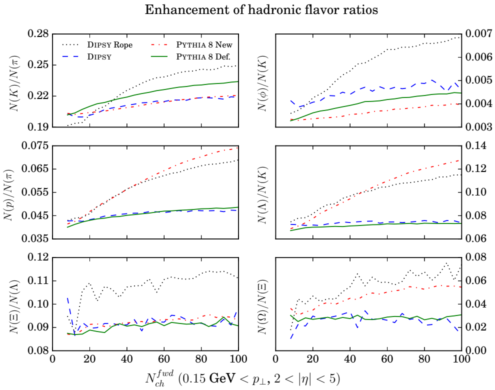


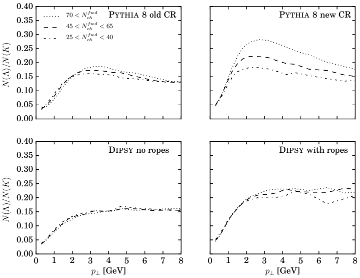


- Three options
  1. Highest multiplet (higher string tension).
  2. Lower multiplet (junction+higher st.).
  3. Singlet Final State swing (similar to CR).

Ropes, junctions and flavours CB and Christiansen: arXiv:1507.02091 [hep-ph]

- Strange enhancement: confirmed, baryons are not.
- Possible solution: Stepwise production mechanism for baryons.
- Flowlike behaviour from junction model.

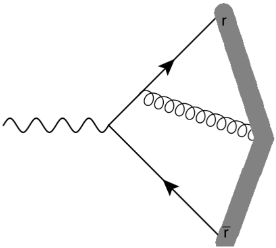




String Hadronization hep-ph/0603175

- Non-perturbative phase of final state.
- *Breaking*/tunneling with $\mathcal{P} \propto \exp\left(-\frac{\pi m_\perp^2}{\kappa}\right)$ gives hadrons.
- Left-right symmetry in the breaking gives

$$f(z) \propto z^{-1}(1-z)^a \exp\left(\frac{-bm_\perp}{z}\right).$$



- a and b related to total multiplicity.
- Flavours determined by relative probabilities:

$$\rho = \frac{\mathcal{P}_{\text{strange}}}{\mathcal{P}_{\text{u or d}}}, \xi = \frac{\mathcal{P}_{\text{diquark}}}{\mathcal{P}_{\text{quark}}}$$

- Probabilities are related to κ via tunneling equation.

Change of string tension

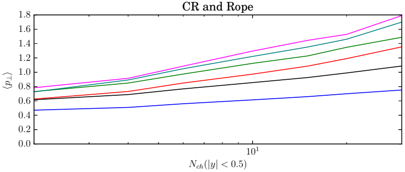
- Field changes when strings overlap - Simple Regge: $2\pi E/l = \kappa$.
- Effective string tension: $\kappa \mapsto \tilde{\kappa} = h\kappa$ from number of overlapping strings.
- Electrodynamics: Principle of superposition, simple.
- QCD: Not so simple. Secondary Casimir operator of multiplet.

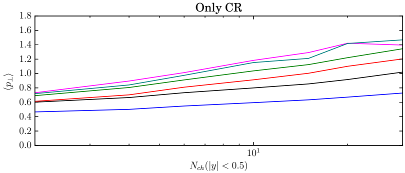
$$\kappa \propto C_2 \Rightarrow h = \tilde{\kappa}/\kappa = \frac{C_2(\text{multiplet})}{1\ \text{GeV/fm}}$$

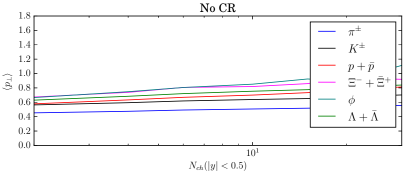
- Confirmed on the lattice, static case.

Ropes, CR and mass splitting

- Influenced heavily by FS effects.
- Tuning and quantitative comparison.
- Remember: Tuning $\neq$ fitting.

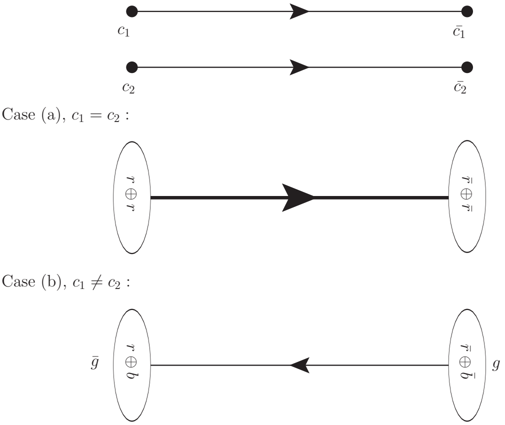






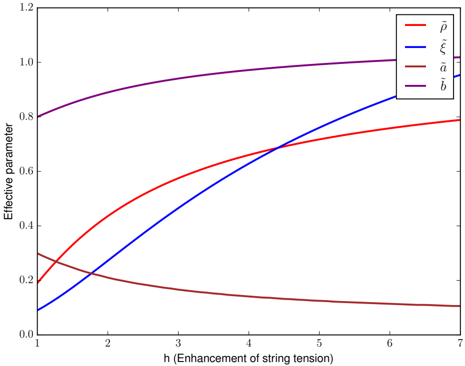
Example arXiv:1412.6259 [hep-ph]

- The simplest example: Two $q\bar{q}$ pairs act coherently.
- Two distinct possibilities:



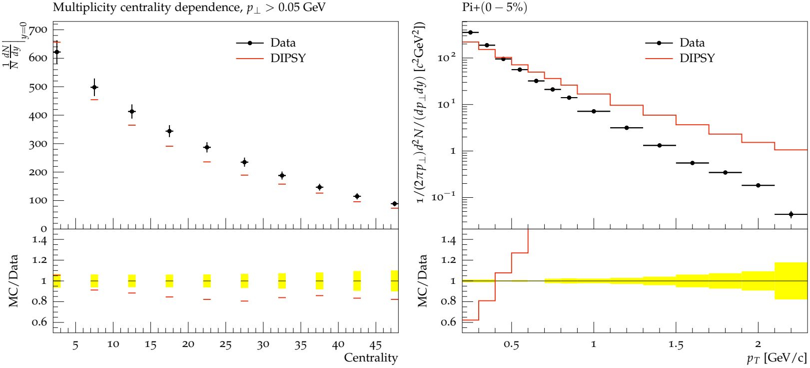
Effect on hadronization parameters

- All parameters related through string tension.
- ρ (strange) and ξ (baryon) are very sensitive.



- Large effect on hadronic flavours.
- Smaller effect on hadron $p_{\perp}$ and multiplicity (tunable).

DIPSY and HI



Multiplicity centrality dependence, $p_\perp > 0.05$ GeV
$\frac{1}{N}\frac{dN}{dy}|_{y=0}$
Data
DIPSY
MC/Data
Centrality
Pi+(0 − 5%)
$1/(2\pi p_\perp) d^2N/(dp_\perp dy)$ [c^2GeV^2])
p_T [GeV/c]

Highest multiplet

- All higher multiplets represents a coherent interaction.
- Fundamental quantum numbers p and q from recursion relations.
- Number of random (anti)-triplets added decided by overlaps.

$$\{p, q\} \otimes \vec{3} = \{p+1, q\} \oplus \{p, q+1\} \oplus \{p, q-1\}$$

$$\underbrace{\begin{smallmatrix}\square\\\square\end{smallmatrix} \otimes \begin{smallmatrix}\square\\\square\end{smallmatrix} \otimes \ldots \otimes \begin{smallmatrix}\square\\\square\end{smallmatrix}}_{\text{All anti-triplets}} \underbrace{\otimes \square \otimes \square \otimes \ldots \otimes \square}_{\text{All triplets}}$$

- Transform to $\tilde{\kappa} = \frac{2p+q+2}{4}\kappa$ and $2N = (p+1)(q+1)(p+q+2)$.
- N (multiplicity of the multiplet) serves as a state's weight.
- String hadronized with $\tilde{\kappa}$.

Junction handling

- Extra junctions handled through simplistic, popcorn-based approach.

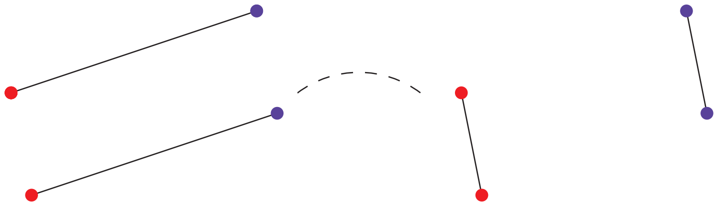
- Extra parameter for colour fluctuations (no data handle).
- Better: Dynamical handling in a "swing".

- Related: recent Pythia 8 model arXiv:1505.01681 [hep-ph]

The singlet swing

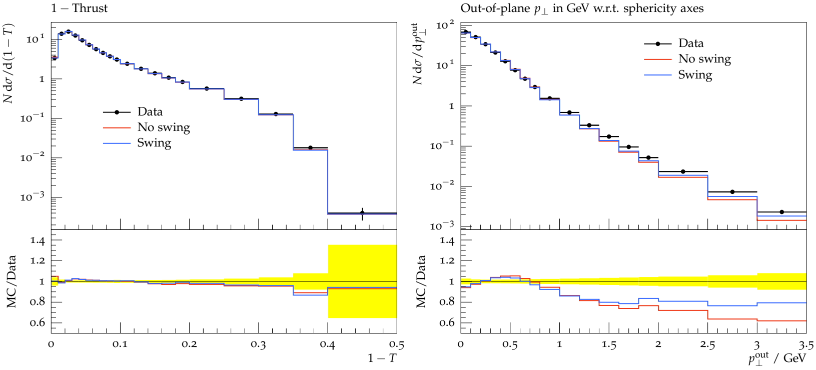
- Singlets are handled already in the FS shower (Ariadne).
- Matching colours *swing* with each other, competing w. emission.

$$\frac{dP_e}{d\ln(p_\perp^2)} \approx dy\frac{C_F\alpha_s}{2\pi} \text{ and } \frac{dP_r}{d\ln(p_\perp^2)} = \lambda\frac{(\vec{p_1}+\vec{p_2})^2(\vec{p_3}+\vec{p_4})^2}{(\vec{p_1}+\vec{p_4})^2(\vec{p_2}+\vec{p_3})^2}$$



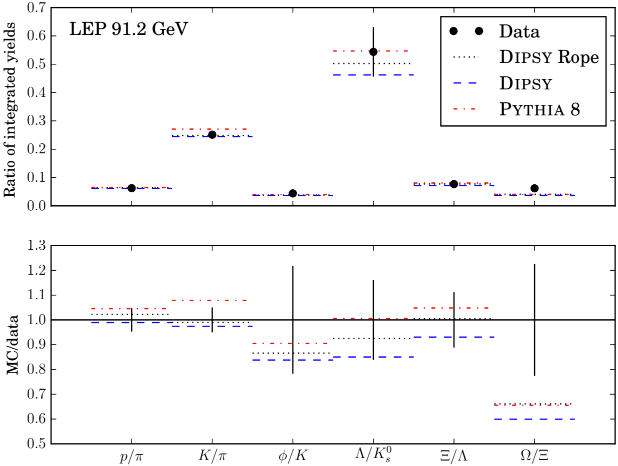
Singlet swing and LEP Data: DELPHI

- Comes in already at perturbative level.
- Retuning of shower is neccesary.
- No large difference, $p_\perp^{out}$ somewhat improved.



Flavour ratios - LEP Data: SLD, LEP and PDG Avg.

- String at LEPs. Agreement with data.
- Jet universality: Gain predictive power in *pp* by fixing parameters here.



Flavour ratios - LHC Data: CMS and ALICE

- Ropes at LHC. Overall better agreement, problem with p/π.
- Integrated quantities, need per event quantities as function of activity.