

QCD with functional methods



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Der Wissenschaftsfonds.



NAWI Graz
Natural Sciences

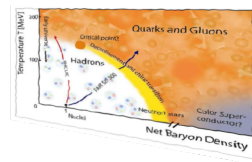
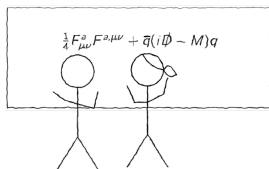


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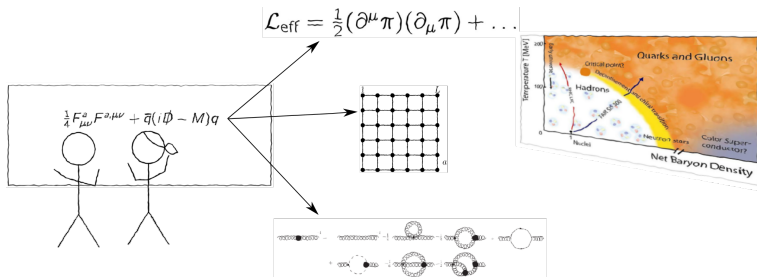
QCD: A simple theory?

$$\frac{1}{4} F_{\mu\nu}^a F^{a,\mu\nu} + \bar{q}(i\cancel{D} - M)q$$

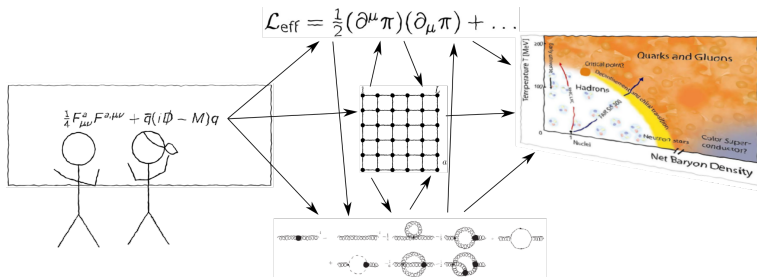
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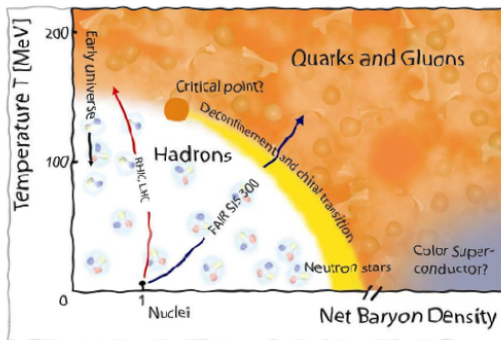
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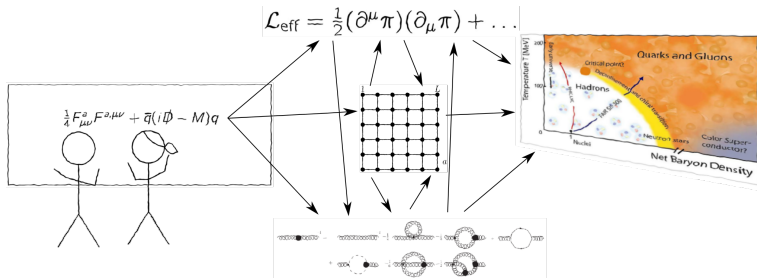


QCD: A simple theory?



- **Phases:** hadronic phase, quark-gluon plasma, color superconductor, quarkyonic?
- **Transitions:** first order line, crossover at $\mu = 0$
- **Critical point:** existence? position?

QCD: A simple theory?



- Challenges for all methods at $\mu > T$, e.g.
 - Lattice QCD: complex action problem
 - Models: parameters
 - Functional methods: reliability of truncations

Functional methods

Functional equations: Exact equations derived from QCD action.

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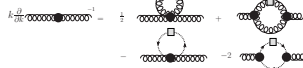
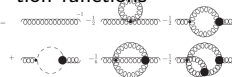
Functional methods

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Dyson-Schwinger eqs.:

eqs. of motion of correlation functions



funct. renorm. group eqs.

eqs. of motion from 3PI eff. action



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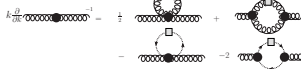
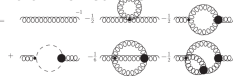
- Chiral limit accessible.
- No sign problem.
- Large scale separations easy.

eqs. of motion from 3PI eff. action



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Functional methods

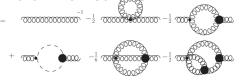
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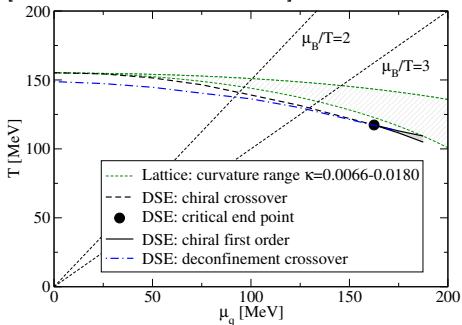
Difficulty

Infinitely large systems of equations without obvious ordering scheme.

QCD phase diagram from functional equations

2+1 flavor QCD from DSEs

[Fischer, Lücker, Welzbacher '14]:



Positions of critical endpoint:

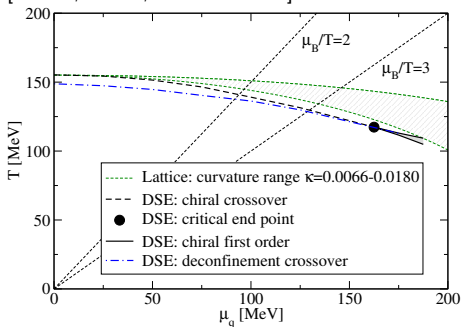
$$\sim (168 \text{ MeV}, 115 \text{ MeV})$$

lattice gluon from $T = 0$, vertex model

QCD phase diagram from functional equations

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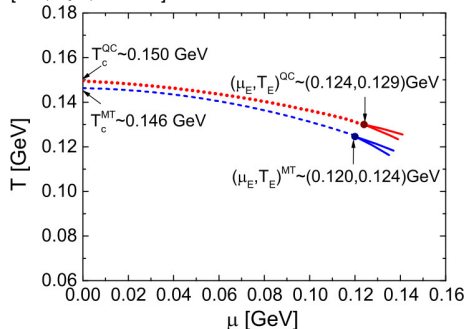
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lattice gluon from $T = 0$, vertex model

2 flavor QCD from DSEs

[Xin, Qin, Liu '15]:



$$\sim (122 \text{ MeV}, 126 \text{ MeV})$$

rainbow approximation

Landau gauge QCD

$$\mathcal{L} = \bar{\mathbf{q}}(-\not{D} + m)\mathbf{q} + \frac{1}{2}F^2 + \mathcal{L}_{gf} + \mathcal{L}_{gh}$$

$$F_{\mu\nu} = \partial_\mu \mathbf{A}_\nu - \partial_\nu \mathbf{A}_\mu + i g [\mathbf{A}_\mu, \mathbf{A}_\nu]$$



Landau gauge QCD

$$\mathcal{L} = \bar{\mathbf{q}}(-\not{D} + m)\mathbf{q} + \frac{1}{2}F^2 + \mathcal{L}_{gf} + \mathcal{L}_{gh}$$

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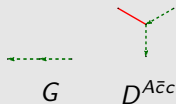


Landau gauge

- simplest one for functional equations

- $\partial_\mu \mathbf{A}_\mu = 0$: $\mathcal{L}_{gf} = \frac{1}{2\xi}(\partial_\mu \mathbf{A}_\mu)^2$, $\xi \rightarrow 0$

- requires ghost fields: $\mathcal{L}_{gh} = \bar{\mathbf{c}}(-\square + g \mathbf{A} \times) \mathbf{c}$



The tower of DSEs

$$\begin{aligned}
 \text{quark propagator} &= \text{tree} + \text{gluon loop} - \frac{1}{2} \text{ghost loop} - \frac{1}{2} \text{ghost-gluon loop} + \text{ghost-gluon loop} \\
 \text{ghost propagator} &= \text{tree} - \text{quark loop} \\
 \text{gluon propagator} &= \text{tree} - \text{quark loop} - \text{ghost loop} - \text{gluon loop}
 \end{aligned}$$

The diagrams represent the following terms:

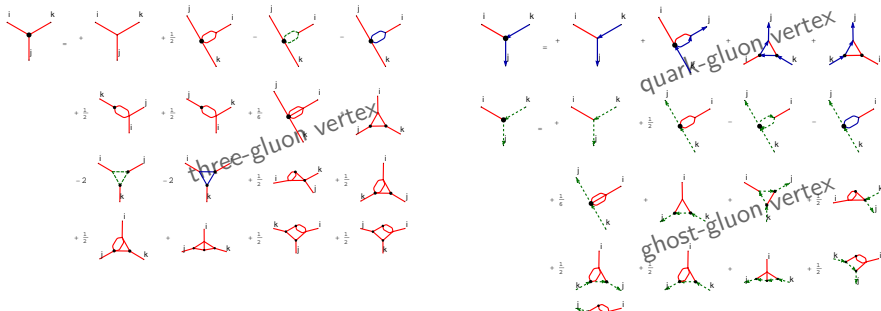
- quark propagator:** A red line with a black dot on the left, labeled i and j . It is equal to the tree-level propagator (red line) plus a gluon loop (red line with a red loop), minus half of a ghost loop (red line with a green loop), minus half of a ghost-gluon loop (red line with a red loop and a ghost loop), plus another ghost-gluon loop (red line with a red loop and a green loop).
- ghost propagator:** A green dashed line with a black dot on the left, labeled j and i . It is equal to the tree-level propagator (green dashed line) minus a quark loop (green dashed line with a red loop).
- gluon propagator:** A red line with a black dot on the left, labeled j and i . It is equal to the tree-level propagator (red line) minus a quark loop (red line with a red loop), minus a ghost loop (red line with a green loop), and minus a gluon loop (red line with a red loop).

The tower of DSEs

$$\begin{aligned}
 \text{gluon propagator} &= \text{tree} + \text{self-energy} - \frac{1}{2} \text{ghost loop} - \frac{1}{2} \text{quark loop} + \text{ghost-gluon vertex} \\
 &+ \text{ghost-gluon vertex} - \frac{1}{6} \text{ghost loop} - \frac{1}{2} \text{quark loop}
 \end{aligned}$$

$$\text{ghost propagator} = \text{tree} - \text{ghost loop}$$

$$\text{quark propagator} = \text{tree} - \text{gluon loop}$$

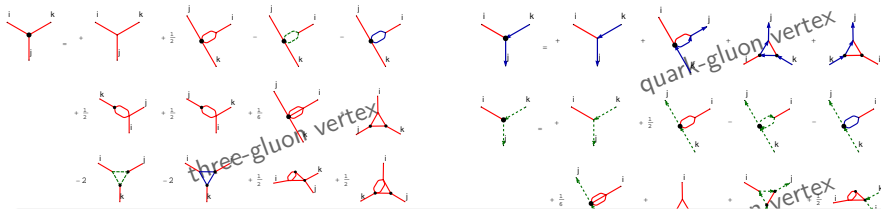


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 \text{gluon propagator} &= \text{tree} + \text{self-energy} - \frac{1}{2} \text{ghost loop} - \frac{1}{2} \text{quark loop} + \text{ghost-gluon loop} \\
 &+ \text{ghost-gluon-gluon loop} - \frac{1}{6} \text{ghost-gluon-gluon-gluon loop} - \frac{1}{2} \text{ghost-gluon-gluon-gluon-gluon loop}
 \end{aligned}$$

$$\text{ghost propagator} = \text{tree} - \text{ghost loop}$$

$$\text{quark propagator} = \text{tree} - \text{quark-gluon loop}$$



Infinitely many equations. In QCD, every n -point function depends on $(n + 1)$ - and possibly $(n + 2)$ -point functions.



Truncating the equations

Guides

- Perturbation theory
- Symmetries
- Lattice
- Analytic results

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Truncation

- Drop quantities (unimportant?)
- Use fits
- Model quantities (good models available? 'true' or 'effective'?)

Ideally: Find a truncation that has (I) **no parameters** and yields (II) **quantitative results**.

Example: Bottom-up approach

Quark gap equation:

$$\rightarrow \text{circle} \rightarrow S(p)^{-1} = \rightarrow S_0(p)^{-1} + \text{loop diagram}$$

Required input:

- Gluon propagator $D(p^2)$
- Quark-gluon interaction $\Gamma(p, q)$

Effective interaction $D(p^2)\Gamma(p, q) \rightarrow \mathcal{G}(p^2)$

Example: Maris-Tandy interaction with parameters ω , D

$$\frac{\mathcal{G}(s)}{s} = \frac{4\pi^2 D}{\omega^6} s e^{-s/\omega^2} + \frac{4\pi \gamma_m \pi \mathcal{F}(s)}{1/2 \ln[\tau + (1 + s/\Lambda_{\text{QCD}}^2)^2]}.$$

→ Use for **bound state** studies.

Example: Top-down for Yang-Mills theory

Neglect all non-primitively divergent Green functions. \rightarrow Self-contained.

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Full propagator equations (two-loop diagrams!):

$$\begin{aligned}
 & \text{Diagram 1: } i \text{---} \bullet \text{---} j \text{ }^{-1} = + \text{Diagram 2: } i \text{---} j \text{ }^{-1} - \frac{1}{2} \text{Diagram 3: } i \text{---} \text{loop} \text{---} j \text{ }^{-1} - \frac{1}{2} \text{Diagram 4: } i \text{---} \text{loop} \text{---} i \text{ }^{-1} + \text{Diagram 5: } j \text{---} \text{loop} \text{---} i \text{ }^{-1} \\
 & \quad - \frac{1}{6} \text{Diagram 6: } j \text{---} \text{loop} \text{---} i \text{ }^{-1} - \frac{1}{2} \text{Diagram 7: } j \text{---} \text{loop} \text{---} j \text{ }^{-1} \\
 & \text{Diagram 8: } j \text{---} \bullet \text{---} i \text{ }^{-1} = + \text{Diagram 9: } j \text{---} i \text{ }^{-1} - \text{Diagram 10: } j \text{---} \text{loop} \text{---} i \text{ }^{-1}
 \end{aligned}$$

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Full propagator equations (two-loop diagrams!):

$$\begin{aligned}
 & \text{Diagram 1} = \text{Diagram 2} + \text{Diagram 3} - \frac{1}{2} \text{Diagram 4} - \frac{1}{2} \text{Diagram 5} + \text{Diagram 6} + \text{Diagram 7} \\
 & \quad - \frac{1}{6} \text{Diagram 8} - \frac{1}{2} \text{Diagram 9}
 \end{aligned}$$

$$\text{Diagram 10} = \text{Diagram 11} - \text{Diagram 12}$$

Truncated three-point functions:

$$\begin{aligned}
 & \text{Diagram 13} = \text{Diagram 14} + \text{Diagram 15} + \text{Diagram 16} + \text{Diagram 17} \\
 & \text{Diagram 18} = \text{Diagram 19} + \text{Diagram 20} + \frac{1}{2} \text{Diagram 21} + \frac{1}{2} \text{Diagram 22} \\
 & \quad + \frac{1}{2} \text{Diagram 23} + \text{Diagram 24} - 2 \text{Diagram 25}
 \end{aligned}$$

Truncated four-gluon vertex:

$$\begin{aligned}
 & \text{Diagram 26} = \text{Diagram 27} + \frac{1}{2} \text{Diagram 28} + 3 \text{Diagram 29} \\
 & \quad + 3 \text{Diagram 30} + 3 \text{Diagram 31} - 6 \text{Diagram 32}
 \end{aligned}$$

Automated derivation

Derivation by hand becomes tedious:

- Large Lagrangians.
- Higher Green functions.
- Larger truncations.
- Error-prone.

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$$\left[2 g^2 N_c Z_1 \text{DAAA} \left[y, q_s + y + 2 \text{sp}[q, q_1], \frac{-y - \text{sp}[q, q_1]}{\sqrt{y (q_s + y + 2 \text{sp}[q, q_1])}} \right] \right]$$

$$\text{DAAA} \left[x_2 + y + 2 \text{sp}[p, q], q_s + x_2 - 2 \text{sp}[p, q_1], \frac{-x_2 - \text{sp}[p, q] + \text{sp}[p, q_1] + \text{sp}[q, q_1]}{\sqrt{(x_2 + y + 2 \text{sp}[p, q]) (q_s + x_2 - 2 \text{sp}[p, q_1])}} \right] \text{Dg1}[q_s] \text{Dg1}[q_s + x_2 - 2 \text{sp}[p, q_1]] \text{Dg1}[$$

$$\text{sp}[p, q]^4 (\text{sp}[p, q_1]^2 \text{sp}[q, q_1] (y + \text{sp}[q, q_1]) + q_s x_2 (y (9 q_s + 6 (x_2 + y)) + (5 q_s + 6 x_2 + 10 y) \text{sp}[q, q_1]) - \text{sp}[p, q_1] (q_s y (5 q_s$$

$$\text{sp}[p, q]^3 (2 \text{sp}[p, q_1]^3 (q_s y - \text{sp}[q, q_1]^2) + \text{sp}[p, q_1] (q_s y (10 q_s^2 + (-5 x_2 - 3 y) y + q_s (19 x_2 + 3 y)) + (3 q_s^3 + 8 q_s x_2 y + 21 q_s^2$$

$$q_s x_2 (y (-9 q_s^2 + 3 x_2^2 + 7 x_2 y + 3 y^2 + 2 q_s (x_2 + y)) + (-10 q_s^2 + q_s (-3 x_2 - 19 y) + x_2 (3 x_2 + 5 y)) \text{sp}[q, q_1] + (-16 q_s - 7 x_2 - 11 y$$

$$\text{sp}[p, q_1]^2 (q_s (-16 q_s - 11 x_2 - 7 y) y + (-5 q_s^2 + q_s (-9 x_2 - 19 y) + 2 y (5 x_2 + 3 y)) \text{sp}[q, q_1] + (-5 q_s + 12 (x_2 + y)) \text{sp}[q, q_1]^2 +$$

$$\text{sp}[p, q]^2 (\text{sp}[p, q_1]^4 \text{sp}[q, q_1] (q_s + \text{sp}[q, q_1]) + \text{sp}[p, q_1]^3 (q_s y (7 q_s + 11 x_2 + 16 y) + (-6 q_s^2 + y (9 x_2 + 5 y) + q_s (-10 x_2 + 19 y)$$

$$q_s x_2 (y (-3 q_s^3 - 10 q_s^2 (x_2 + y) - 6 x_2 y (x_2 + y) + q_s (-3 x_2^2 - 19 x_2 y - 3 y^2)) + (-6 q_s^3 + q_s^2 (-21 x_2 - 32 y) + q_s (-9 x_2^2 - 60 x_2 y$$

$$(-15 q_s^2 - 15 x_2^2 + q_s (-46 x_2 - 41 y) - 41 x_2 y - 12 y^2) \text{sp}[q, q_1]^2 + (-7 q_s - 16 x_2 - 11 y) \text{sp}[q, q_1]^3) + \text{sp}[p, q_1]^2 (q_s y (-15 q_s$$

$$(3 q_s^3 + q_s^2 (5 x_2 - 39 y) + q_s (-81 x_2 - 39 y) y + y^2 (5 x_2 + 3 y)) \text{sp}[q, q_1] + (12 q_s^2 + 12 x_2^2 + 3 x_2 y + 12 y^2 + 3 q_s (x_2 + y)) \text{sp}[q$$

$$\text{sp}[p, q_1] (q_s y (6 q_s^3 + q_s^2 (32 x_2 + 21 y) + q_s (25 x_2^2 + 60 x_2 y + 9 y^2) + x_2 (3 x_2^2 + 25 x_2 y + 15 y^2)) + (15 q_s^3 (x_2 + y) + x_2 y (-3 x_2$$

$$(-3 q_s^3 + x_2^2 (-3 x_2 - 5 y) + q_s^2 (39 x_2 - 5 y) + q_s x_2 (39 x_2 + 81 y)) \text{sp}[q, q_1]^2 + (-6 q_s^2 + q_s (19 x_2 - 10 y) + x_2 (5 x_2 + 9 y)) \text{sp}[q$$

$$x_2 y (-\text{sp}[p, q_1]^5 (q_s + \text{sp}[q, q_1]) + \text{sp}[p, q_1]^4 (q_s (6 q_s + 6 x_2 + 9 y) + (10 q_s + 6 x_2 + 5 y) \text{sp}[q, q_1]) - q_s (q_s y - \text{sp}[q, q_1]^2) (x_2 (-$$

$$(6 q_s + 9 x_2 + 6 y) \text{sp}[q, q_1]^2 + \text{sp}[q, q_1]^3) + \text{sp}[p, q_1]^3 (q_s (-3 q_s^2 - 3 x_2^2 + q_s (-7 x_2 - 2 y) - 2 x_2 y + 9 y^2) + (-3 x_2^2 + 3 x_2 y +$$

$$\text{sp}[p, q_1]^2 (q_s (-3 q_s^2 (2 x_2 + y) + q_s (-6 x_2^2 - 19 x_2 y - 10 y^2) + y (-3 x_2^2 - 10 x_2 y - 3 y^2)) + (-3 q_s^3 - 25 q_s^2 (x_2 + y) + q_s (-15 x_2$$

$$(-12 q_s^2 - 15 x_2^2 - 46 x_2 y - 15 y^2 - 41 q_s (x_2 + y)) \text{sp}[q, q_1]^2 + (-11 q_s - 16 x_2 - 7 y) \text{sp}[q, q_1]^3) +$$

Automated derivation

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- Automated derivation of DSEs and **flow equations**:

Mathematica package *DoFun* [Alkofer, MQH, Schwenzer '08; MQH, Braun '11]

<http://tinyurl.com/dofun2>

- Framework for numeric handling:

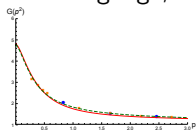
C++ program *CrasyDSE* [MQH, Mitter '11]

<http://tinyurl.com/crasydse>

Go ahead and calculate . . .

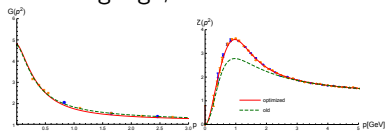
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Landau gauge, vacuum:



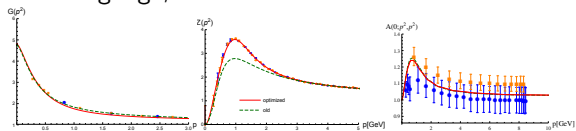
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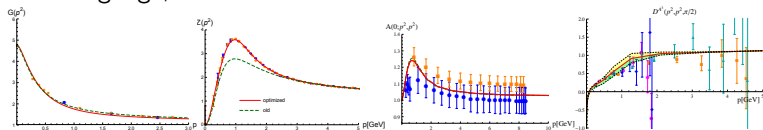
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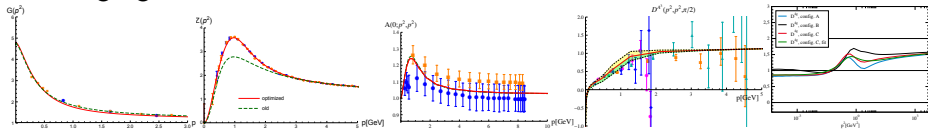
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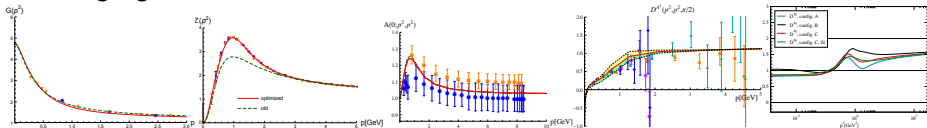
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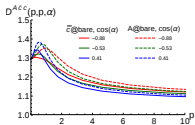


Go ahead and calculate ...

Landau gauge, vacuum:

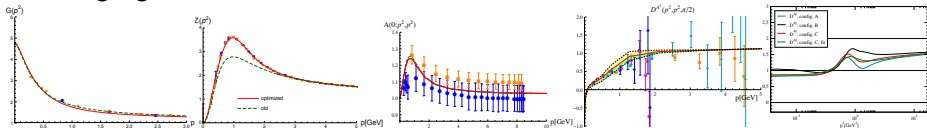


Coulomb and linear covariant gauges, vacuum:

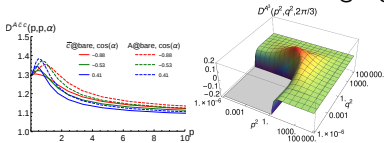


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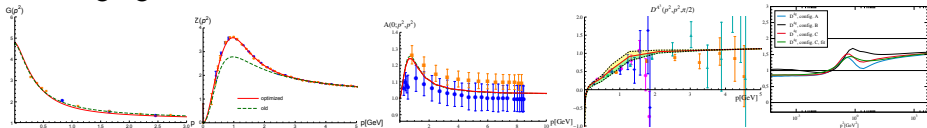


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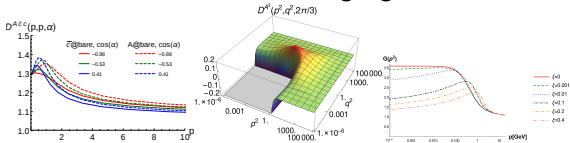


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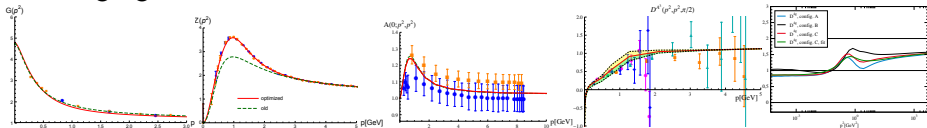


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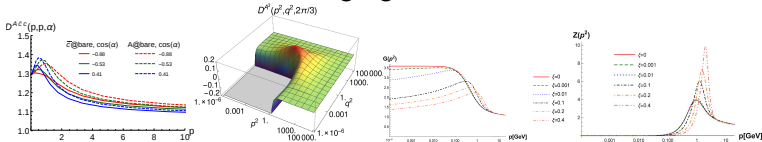


Go ahead and calculate ...

Landau gauge, vacuum:

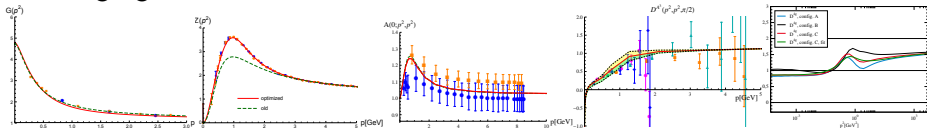


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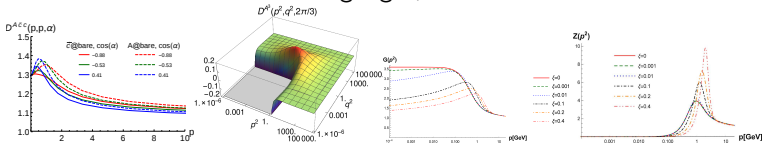


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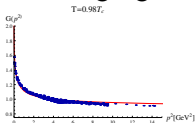
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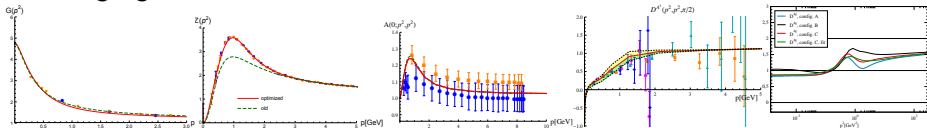


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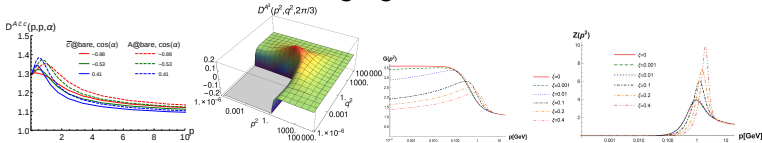


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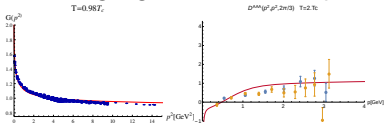
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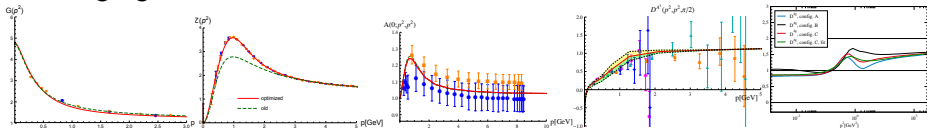


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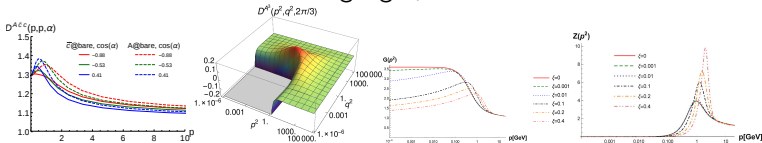


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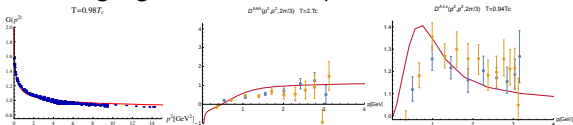
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- Vacuum
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→ Comparison with lattice is helpful, but finally self-consistent checks are required.

Two words of caution:

- One cannot assume naturally that the hierarchy is the same for all T and μ .
- Even the effect of a single correlation function is difficult to estimate.

Example: What do we need to go beyond modern QCD phase diagram calculations?

Beyond effective interaction approximation: ✓ [Fischer, Lücker, Welzbacher '14]

Input for DSEs:

- model for quark-gluon vertex (parameters fixed at $\mu = 0$)
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Ultimately, **full control** over Yang-Mills part required!

Yang-Mills theory

Up to now

(separately and partly combined):

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- Three-point functions
- Four-gluon vertex (subset of dressings)

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Challenges:

- Find **truncation**.
- **Solve** large system of equations.
- Spurious (quadratic) divergences in gluon propagator: Consistent subtraction [MQH, von Smekal, '14]?

Testing truncations

Switch to three dimensions:

- UV finite: no renormalization, no anomalous running
- Spurious divergences easier to handle

⇒ Many complications from $d = 4$ absent!

In particular: Spurious divergences of 'simple' enough form ($\propto a\Lambda + b \ln \Lambda$).

Quantitative study of truncation effects possible:

Vary equations and truncations.

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Truncation

Complete system of propagators, three-point function and four-gluon vertex.

Varying the four-gluon vertex

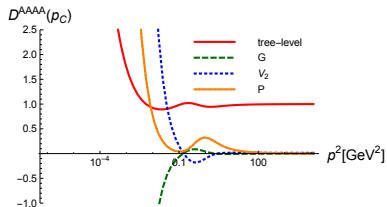
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Compare:

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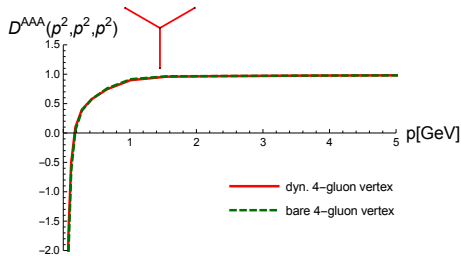
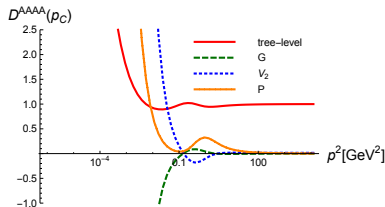
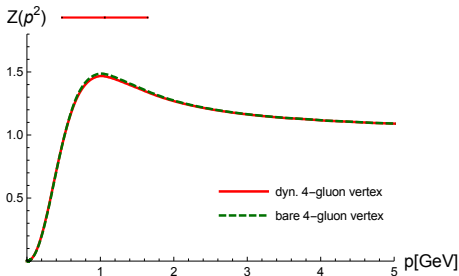


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→ Very similar results.

[MQH '16]

Solution from the 3PI effective action

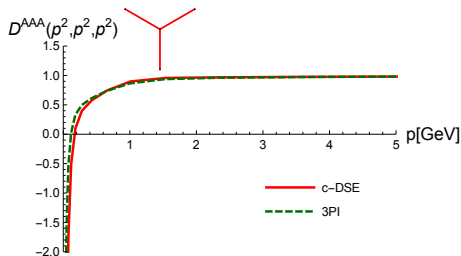
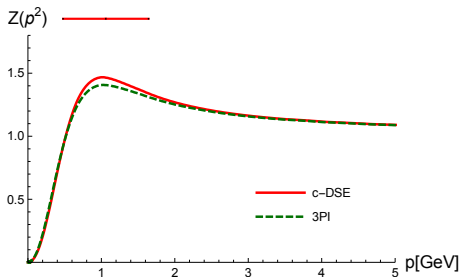
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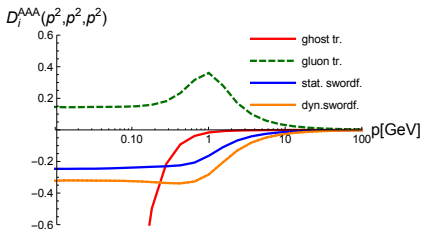
[MQH '16]

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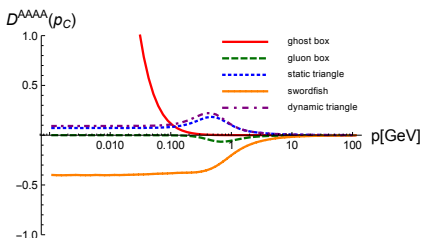
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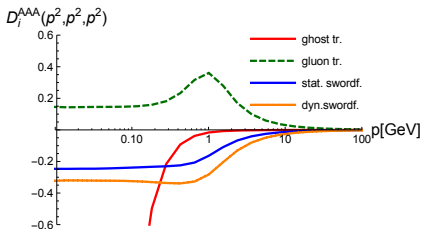
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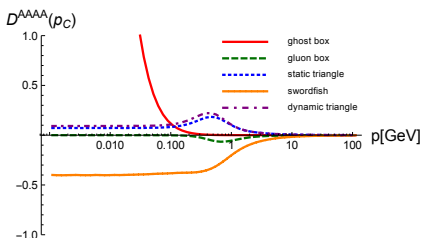
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Four-gluon vertex:



Higher contributions:

- Higher vertices close to 'tree-level'?
→ Small.
- If pattern changes (higher vertices large): cancellations required.

[MQH '16]

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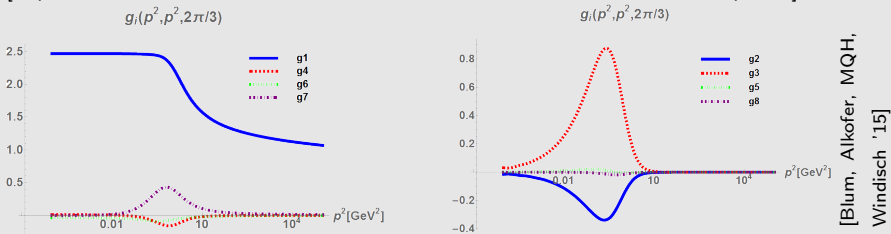
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Explicit quark-gluon vertex solution

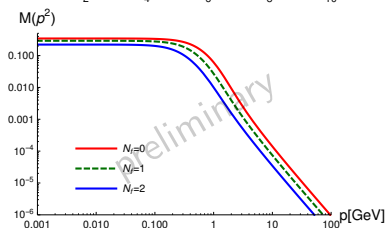
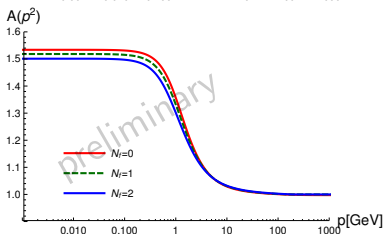
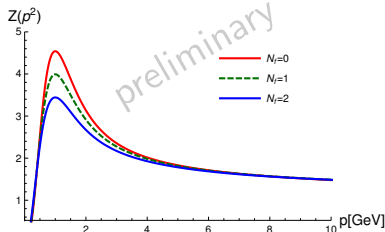
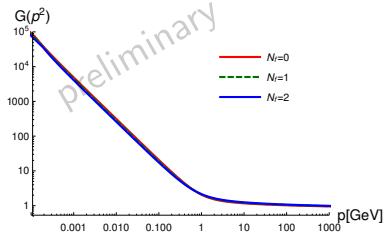
[Hopfer '14; Windisch '14; Mitter, Pawlowski, Strodthoff '14; Williams, Fischer, Heupel '15]



\rightarrow Also non-tree-level dressings become important.

First step: Unquenching of propagators

Subtraction of spurious divergences as in Yang-Mills part [MQH, von Smekal '14]!
Models for vertices.



→ Poster by Contant.

[Contant, MQH, unpublished]

Summary and conclusions

- Functional equations: Non-perturbative approach to QCD.
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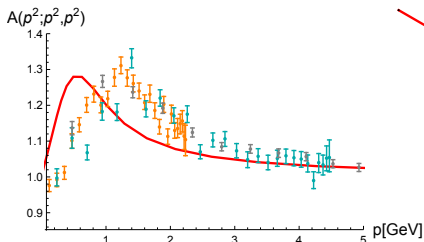
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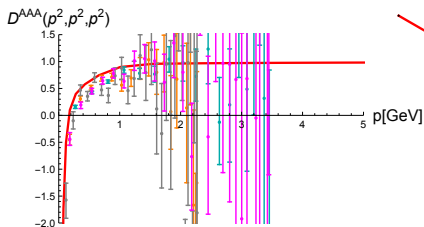
Results: Three-point functions

Dressings:



[MQH '16; lattice: Maas, unpublished]

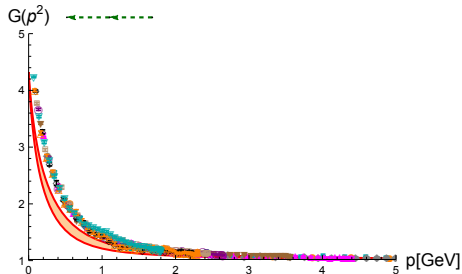
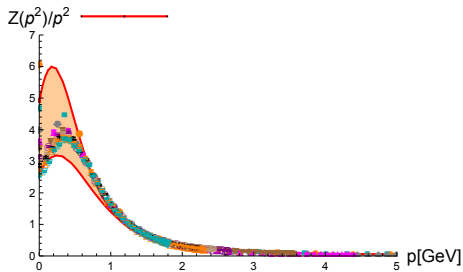
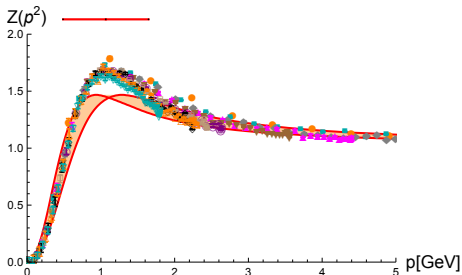
- Maximum position shifted.
- Bump height ok.



[MQH '16; lattice: Cucchieri, Maas, Mendes]

- Good agreement with lattice data.
- Linear IR divergence.

Results: Propagators



Bands from uncertainty in setting the physical scale.

[MQH '16; lattice: Maas '14]

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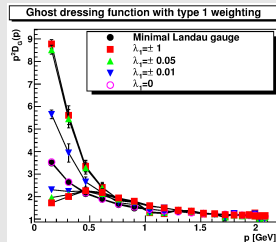
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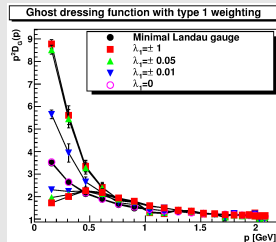
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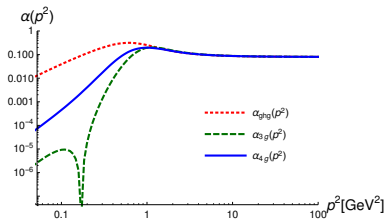
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$$\alpha_{4g}(p^2) = \frac{g^2}{4\pi} D^{AAAA}(p^2, p^2, p^2) Z(p^2)^2$$



⇒ Good agreement down to a few GeV!

[MQH '16]