

Manifestations of Dark Matter and Variation of the Fundamental Constants in Atomic and Astrophysical Phenomena

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Physical Review Letters **113**, 081601 (2014)

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Physical Review D **90**, 096005 (2014)

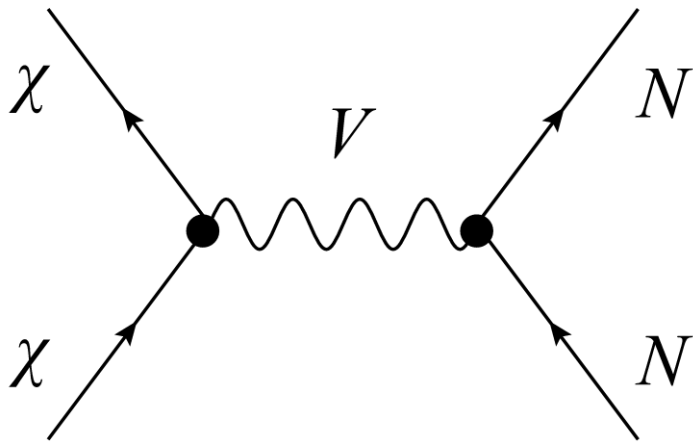
European Physical Journal C **75**, 110 (2015)

arXiv:1511.00447, 1604.04559, 1605.04028

Nature Physics **12**, 465 (2016)

Motivation

Traditional “scattering-off-nuclei” searches for heavy WIMP dark matter particles (χ) have not yet produced a strong positive result.



$$\mathcal{M}_{\text{scat}} \propto (e')^2$$
$$\Rightarrow \sigma_{\text{scat}} \propto (e')^4$$

Observable is **quartic** in the interaction constant e' , which is extremely small!

Motivation

We propose to search for other well-motivated forms of dark matter: *low-mass spin-0 particles*, which form a *coherently* oscillating classical†* field:

$\varphi(t) = \varphi_0 \cos(m_\varphi t)$, via effects that are **linear** in the interaction constant ($\Lambda_X =$ new-physics energy scale).

$$\mathcal{L}_{\text{eff}} = \frac{\phi}{\Lambda_X} X_{\text{SM}} X_{\text{SM}} \Rightarrow \mathcal{O} \propto \frac{1}{\Lambda_X}$$

Consideration of *linear effects* has already allowed us to improve on existing constraints on some interactions of dark matter by up to **15 orders of magnitude**, as well as derive the **first constraints** on some other interactions of dark matter.

* Coherently oscillating field \Rightarrow cold, $E_\varphi = m_\varphi c^2$

† $n_\varphi (\lambda_{\text{dB}}/2\pi)^3 \gg 1$

Low-mass Spin-0 Dark Matter

The mass range $10^{-22} \text{ eV} \leq m_\phi \leq 0.1 \text{ eV}$ is inaccessible to traditional “scattering-off-nuclei” and collider searches, but large regions are accessible to low-energy atomic and molecular experiments that search for **oscillating signals** produced by $\phi(t) = \phi_0 \cos(m_\phi t)$ [$10^{-8} \text{ Hz} \leq f \leq 10^{13} \text{ Hz}$].

In particular, ultra-low-mass spin-0 DM with mass $m_\phi \sim 10^{-22} \text{ eV}$ has been proposed to resolve several long-standing astrophysical puzzles (cusp-core, missing satellite and too-big-to-fail problems, etc.)

Coherence of Galactic DM

Gravitational interactions between DM and ordinary matter during galactic structure formation result in the virialisation of the DM particles ($v_{\text{vir}} \sim 10^{-3} c$), which gives the galactic DM field a finite coherence time and finite coherence length:

$$\tau_{\text{coh}} \sim \frac{2\pi}{m_{\phi} v_{\text{vir}}^2} \sim 10^6 \left(\frac{2\pi}{m_{\phi}} \right) \Rightarrow \frac{\Delta f}{f} \sim 10^{-6}$$

$$l_{\text{coh}} \sim \frac{1}{m_{\phi} v_{\text{vir}}} \sim \frac{10^3}{m_{\phi}} = \frac{10^3}{2\pi} \lambda_{\text{Compton}}$$

Low-mass Spin-0 Dark Matter

Dark Matter

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graph TD; DM[Dark Matter] --> S[Scalars: Even-parity]; DM --> PS[Pseudoscalars (Axions, ALPs): Odd-parity];
```

Scalars:
Even-parity

→ **‘Slow’ evolution and oscillating variation of fundamental constants**

- Atomic clocks
- Highly-charged ions
- Molecules
- Nuclear clocks
- Laser interferometers

**Pseudoscalars
(Axions, ALPs):**
Odd-parity

→ **Oscillating spin-dependent effects,
 P, T , Lorentz and Einstein
symmetry violation**

- Atomic magnetometry
- Ultracold neutrons
- Solid-state magnetometry

Variation of fundamental constants (fine structure constant α , α_s , masses) due to Dark matter

“**Fine tuning**” of fundamental constants is needed for life to exist. If fundamental constants would be even slightly different, life could not appear!

Variation of coupling constants in space provide natural explanation of the “fine tuning”: we appeared in area of the Universe where values of fundamental constants are suitable for our existence.

Source of the variation: Dark Matter?

We performed calculations to link change of atomic transition frequencies to change of fundamental constants:

optical transitions, atomic calculations for quasar absorption spectra and for atomic clocks transitions in Al II, Ca I, Sr I, Sr II, In II, Ba II, Dy I, Yb I, Yb II, Yb III, Hg I, Hg II, Tl II, Ra II, Th III ...

$$\omega = \omega_0 + q(\alpha^2/\alpha_0^2 - 1)$$

Molecular calculations

Microwave transitions: hyperfine frequency is sensitive to α , nuclear magnetic moments and nuclear radii.

We performed atomic, QCD and nuclear calculations.

Evidence for spatial variation of the fine structure constant

$$\alpha = e^2 / 2\epsilon_0 hc = 1/137.036$$

Quasar spectra

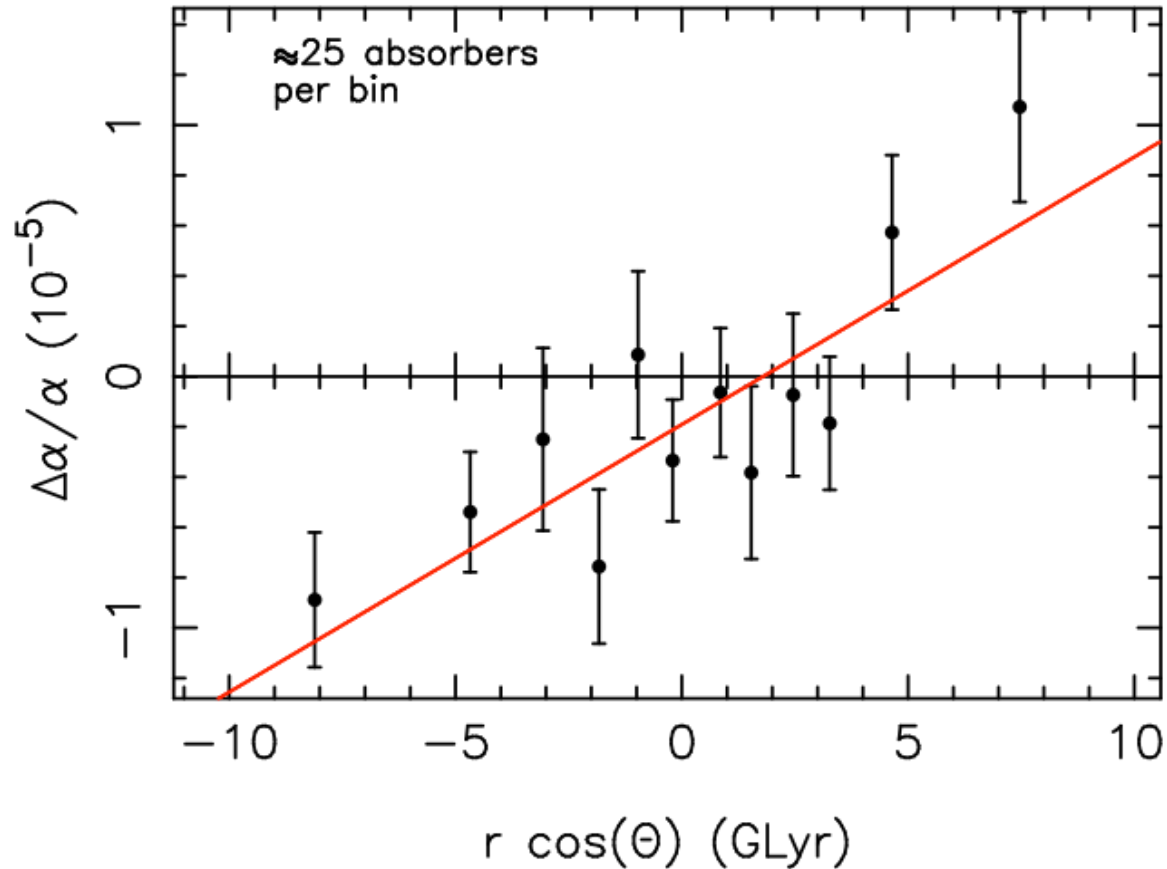
Webb, King, Murphy, Flambaum, Carswell,
Bainbridge, PRL2011, MNRAS2012

$$\alpha(x) = \alpha(0) + \alpha'(0)x + \dots$$

$x = r \cos(\phi)$, $r = ct$ – distance (t - light travel
time, c - speed of light)

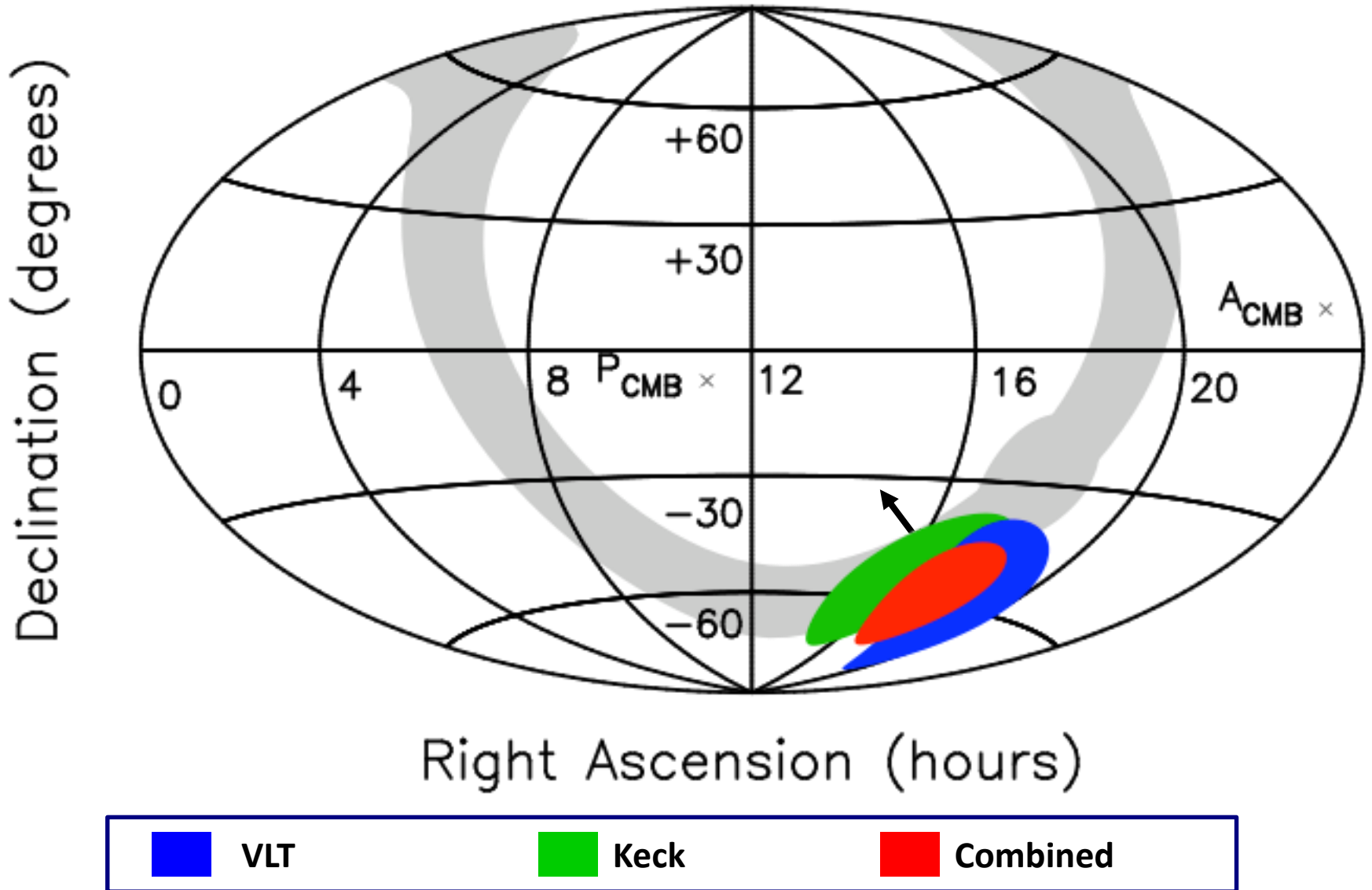
Reconciles all measurements of the variation

Distance dependence



$\Delta\alpha/\alpha$ vs $B r \cos\theta$ for the model $\Delta\alpha/\alpha = B r \cos\theta + m$ showing the gradient in α along the best-fit dipole. The best-fit direction is at right ascension 17.4 ± 0.6 hours, declination -62 ± 6 degrees, for which $B = (1.1 \pm 0.2) \times 10^{-6}$ GLyr^{-1} and $m = (-1.9 \pm 0.8) \times 10^{-6}$. This dipole+monopole model is statistically preferred over a monopole-only model also at the 4.1σ level. A cosmology with parameters $(H_0, \Omega_M, \Omega_\Lambda) = (70.5, 0.2736, 0.726)$.

Keck & VLT dipoles independently agree, $p=4\%$



Results for variation of fundamental constants: Clocks comparison

Source	Clock ₁ /Clock ₂	$d\alpha/dt/\alpha(10^{-16} \text{ yr}^{-1})$
Godun <i>et al</i> , 2014	Yb+opt/Yb+/Cs(hfs)	-0.07(0.21)
Leefer <i>et al</i> 2013	Dy/Cs(hfs)	-0.6(0.7)
Rosenband <i>et al</i> 08	Hg+(opt)/Al+(opt)	-0.16(0.23)
Huntemann <i>et al</i> 14	Yb+opt/Yb+/Cs(hfs)	-0.2(0.2)
Guena <i>et al</i> , 2012	Rb(hfs)/Cs(hfs)	3(2) ^a

^aassuming $m_{q,e}/\Lambda_{\text{QCD}} = \text{Const}$

**Combined results: $d/dt \ln \alpha = -1.5(1.0) \times 10^{-17} \text{ yr}^{-1}$
 $d/dt \ln(m_q/\Lambda_{\text{QCD}}) = 7(4) \times 10^{-15} \text{ yr}^{-1}$
 m_e/M_p or $m_e/\Lambda_{\text{QCD}} -0.1(1.0) \times 10^{-16} \text{ yr}^{-1}$**

Low-mass Spin-0 Dark Matter

Dark Matter



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graph TD; DM[Dark Matter] --> S[Scalars: Even-parity]
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Scalars:
Even-parity

→ **'Slow' evolution and oscillating variation of fundamental constants**

- Atomic clocks
- Highly-charged ions
- Molecules
- Nuclear clocks
- Laser interferometers

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

Consider an oscillating classical *scalar* field, $\phi(t) = \phi_0 \cos(m_\phi t)$, that interacts with SM fields (e.g. a fermion f) via quadratic couplings in ϕ .

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_f^{\text{SM}} = -m_f \bar{f} f \quad \Rightarrow \quad m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

$$\Rightarrow \frac{\delta m_f}{m_f} = \frac{\phi_0^2}{(\Lambda'_f)^2} \cos^2(m_\phi t) = \frac{\phi_0^2}{2(\Lambda'_f)^2} + \frac{\phi_0^2}{2(\Lambda'_f)^2} \cos(2m_\phi t)$$

'Slow' drifts [Astrophysics
(high ρ_{DM}): BBN, CMB]

Oscillating variations
[Laboratory (high precision)]

Dark Matter-Induced Cosmological Evolution of the Fundamental Constants

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

We can consider a wide range of quadratic-in- ϕ interactions with the SM sector:

Photon:

$$\mathcal{L}_\gamma = \frac{\phi^2}{(\Lambda'_\gamma)^2} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \Rightarrow \alpha \rightarrow \frac{\alpha}{1 - \phi^2/(\Lambda'_\gamma)^2} \simeq \alpha \left[1 + \frac{\phi^2}{(\Lambda'_\gamma)^2} \right]$$

Fermions:

$$\mathcal{L}_f = -\frac{\phi^2}{(\Lambda'_f)^2} m_f \bar{f} f \Rightarrow m_f \rightarrow m_f \left[1 + \frac{\phi^2}{(\Lambda'_f)^2} \right]$$

Bosons W,Z (mediators of weak interactions):

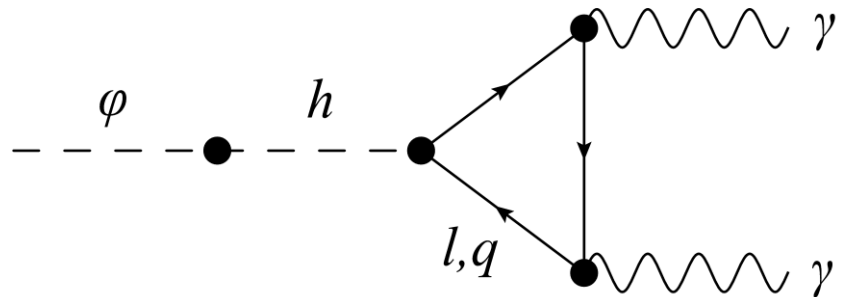
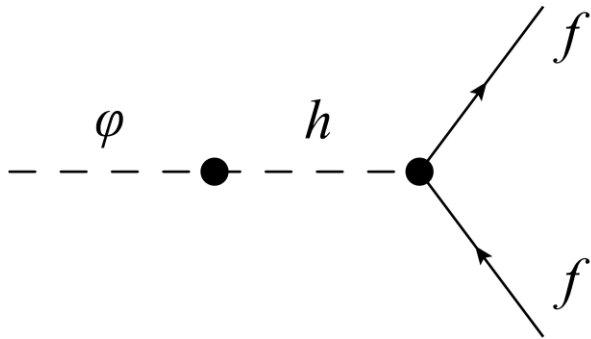
$$\mathcal{L}_V = \frac{\phi^2}{(\Lambda'_V)^2} \frac{M_V^2}{2} V_\nu V^\nu \Rightarrow M_V^2 \rightarrow M_V^2 \left[1 + \frac{\phi^2}{(\Lambda'_V)^2} \right]$$

Dark Matter-Induced Oscillating Variation of the Fundamental Constants

Also possible to have linear-in- ϕ interactions with the SM sector, which may be generated, e.g., through the super-renormalisable interaction of ϕ with the Higgs boson*

[Piazza, Pospelov, *PRD* **82**, 043533 (2010)]:

$$\mathcal{L}_H = -A\phi H^\dagger H$$



$$m_f \rightarrow m_f \left[1 - \frac{Ag_{hff} \langle h \rangle \phi}{m_f m_h^2} \right]$$

$$\alpha \rightarrow \alpha \left[1 + \frac{4Ag_{h\gamma\gamma} \phi}{m_h^2} \right]$$

* Produces logarithmically-divergent corrections to $(m_\phi)^2$, i.e., technically natural for $A < m_\phi$. Minimum of potential is stable (without adding extra ϕ^4 terms) for $(A/m_\phi)^2 < 2\lambda$.

Astrophysical Constraints on 'Slow' Drifts in Fundamental Constants Induced by Scalar Dark Matter (BBN)

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

- Largest effects of scalar dark matter are in the early Universe (highest $\rho_{\text{DM}} \Rightarrow$ highest φ_0^2).
- Earliest cosmological epoch that we can probe is Big Bang nucleosynthesis (from $t_{\text{weak}} = 1\text{s}$ until $t_{\text{BBN}} = 3\text{ min}$).
- Primordial ${}^4\text{He}$ abundance is sensitive to relative abundance of neutrons to protons (almost all neutrons are bound in ${}^4\text{He}$ by the end of BBN).

Weak interactions: freeze-out of weak interactions occurs at $t_{\text{weak}} = 1\text{s}$ ($T_{\text{weak}} = 0.75\text{ MeV}$).

$$\begin{aligned} p + e^- &\rightleftharpoons n + \nu \\ n + e^+ &\rightleftharpoons p + \bar{\nu} \end{aligned} \quad \left(\frac{n}{p}\right)_{\text{weak}} = e^{-(m_n - m_p)/T_{\text{weak}}}$$

Astrophysical Constraints on 'Slow' Drifts in Fundamental Constants Induced by Scalar Dark Matter (CMB)

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

- Weaker astrophysical constraints come from CMB measurements (lower ρ_{DM}).
- Variations in α and m_e at the time of electron-proton recombination affect the ionisation fraction and Thomson scattering cross section, $\sigma_{\text{Thomson}} = 8\pi\alpha^2/3m_e^2$, changing the mean-free-path length of photons at recombination and leaving distinct signatures in the CMB angular power spectrum.

$$\Lambda'_\gamma \gtrsim \frac{1 \text{ eV}^2}{m_\phi}, \quad \Lambda'_e \gtrsim \frac{0.6 \text{ eV}^2}{m_\phi}$$

Laboratory Searches for Oscillating Variations in Fundamental Constants Induced by Scalar Dark Matter

[Arvanitaki, Huang, Tilburg, PRD **91**, 015015 (2015); Stadnik, Flambaum, PRL **115**, 201301 (2015)]

- In the laboratory, we can search for **oscillating variations in the fundamental constants** induced by scalar DM, using **clock frequency comparison measurements**.

$$\frac{\delta(\omega_1/\omega_2)}{\omega_1/\omega_2} \propto \sum_X (K_{X,1} - K_{X,2}) \cos(\omega t)$$

- Exact frequency of oscillation is unknown: $\omega = m_\phi$ (linear) or $\omega = 2m_\phi$ (quadratic) [10^{-22} eV $\leq m_\phi \leq 0.1$ eV $\Rightarrow 10^{-8}$ Hz $\leq f \leq 10^{14}$ Hz], with $\Delta f/f \sim 10^{-6}$.
- Need to search over a broad range of frequencies.

Laboratory Searches for Oscillating Variations in Fundamental Constants Induced by Scalar Dark Matter

System	Laboratory	Constraints
$^{162,164}\text{Dy}/^{133}\text{Cs}$	UC Berkeley	Van Tilburg, LEEFER, Bougas, Budker, <i>PRL</i> 115 , 011802 (2015); Stadnik, Flambaum, <i>PRL</i> 115 , 201301 (2015) + arXiv:1605.04028
$^{87}\text{Rb}/^{133}\text{Cs}$	LNE-SYRTE Paris	Hees, Guena, Abgrall, Bize, Wolf, arXiv:1604.08514; Stadnik, Flambaum, arXiv:1605.04028

Laser Interferometry (LIGO, Virgo, GEO600, TAMA300, smaller-scale)

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015)]

Extremely sensitive laser interferometers can be used to search for oscillating effects produced by scalar field.



Laser Interferometry (LIGO, Virgo, GEO600, TAMA300, smaller-scale)

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015)]

Laser interferometers can be used to search for oscillating effects produced by scalar field.

Accumulated phase in an arm, $\Phi = \omega L/c$, changes if fundamental constants change ($L = Na_B$ and ω_{atomic} depend on the fundamental constants).

$$\Phi = \frac{\omega_{\text{electronic}} L}{c} \approx \left(\frac{e^2}{a_B \hbar} \right) \left(\frac{Na_B}{c} \right) = N\alpha$$
$$\Rightarrow \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

$\Phi = 2\pi L/\lambda$, $\delta\Phi = \Phi \delta\alpha/\alpha = 10^{11} \delta\alpha/\alpha$ single passage,
up to $10^{14} \delta\alpha/\alpha$ for maximal number of reflections

Laboratory Searches for Oscillating Variations in Fundamental Constants Induced by Scalar Dark Matter

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); arXiv:1511.00447]

- We can compare a *photon wavelength* with an *interferometer arm length*.
- Accumulated phase in an arm, $\Phi = \omega L/c$, changes if the fundamental constants change ($L \sim Na_B$ and ω_{atomic} depend on the fundamental constants).

$$\Phi = \frac{\omega L}{c} \propto \left(\frac{e^2}{a_B \hbar} \right) \left(\frac{Na_B}{c} \right) = N\alpha \Rightarrow \frac{\delta\Phi}{\Phi} \approx \frac{\delta\alpha}{\alpha}$$

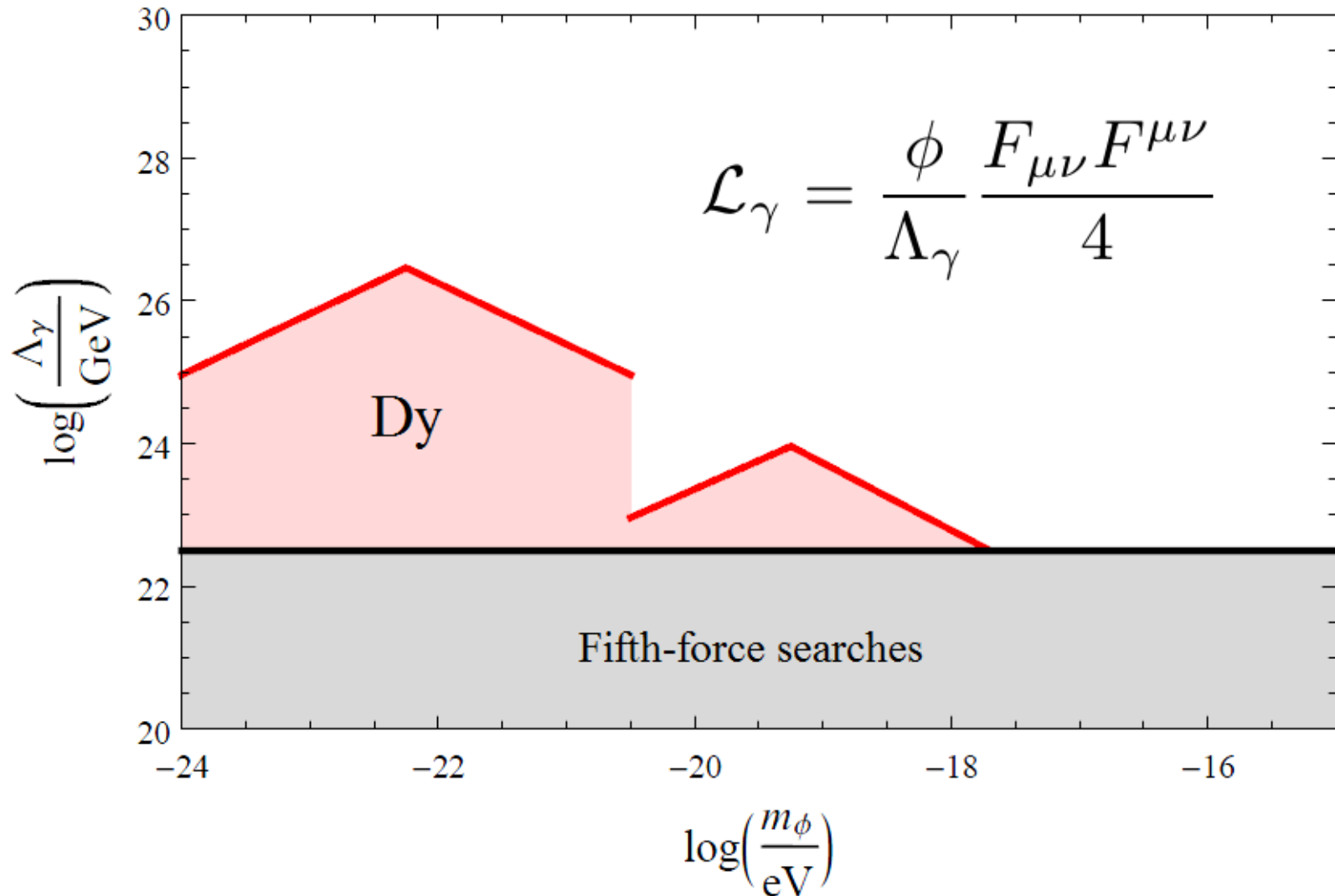
- Multiple reflections enhance observable effects due to variation of the fundamental constants by the effective mean number of passages N_{eff} (e.g. $N_{\text{eff}} \sim 10^5$ in a strontium clock – silicon cavity interferometer).

Laboratory Searches for Oscillating Variations in Fundamental Constants Induced by Scalar Dark Matter

System	Λ'_y	Λ'_e	Λ'_p	Λ'_q
Atomic (Dy, optical clock)	+	-	-	-
Atomic (hyperfine)	+	+	+	+
Highly charged ionic	+	-	-	-
Molecular (hyperfine/rotational)	+	+	+	+
Molecular (fine-structure/vibrational)	+	+	+	+
Molecular (Ω -doubling/hyperfine)	+	+	+	+
Nuclear (e.g. ^{229}Th)	+	-	+	+
Laser interferometer, Bar	+	+	+	+

Laboratory Search for Oscillating Variations in Fundamental Constants using Atomic Dysprosium

[Van Tilburg, Leefer, Bougas, Budker, *PRL* **115**, 011802 (2015)]

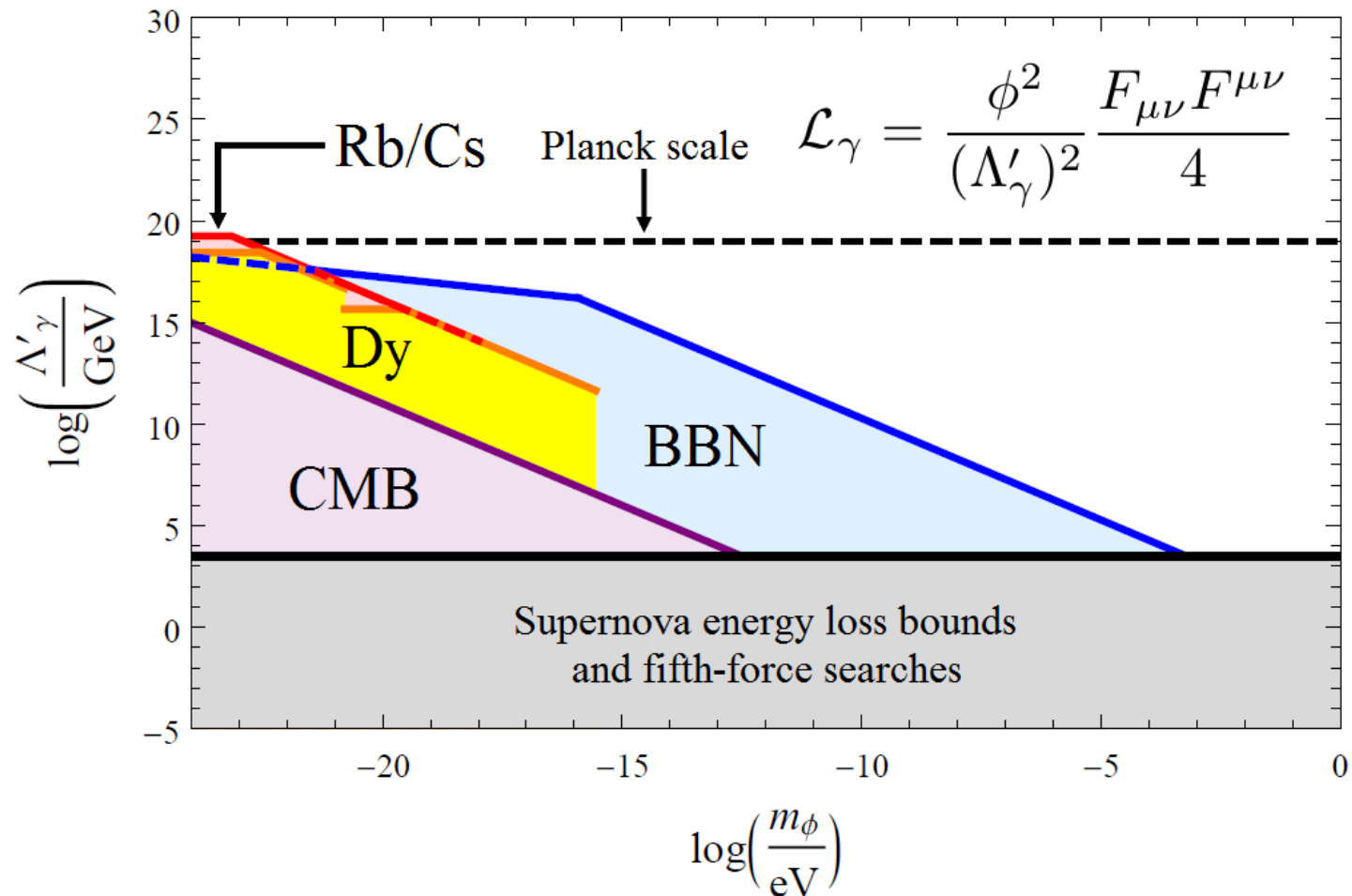


Constraints on Quadratic Interaction of Scalar Dark Matter with the Photon

BBN, CMB, Dy and Rb/Cs constraints:

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015) + arXiv:1605.04028]

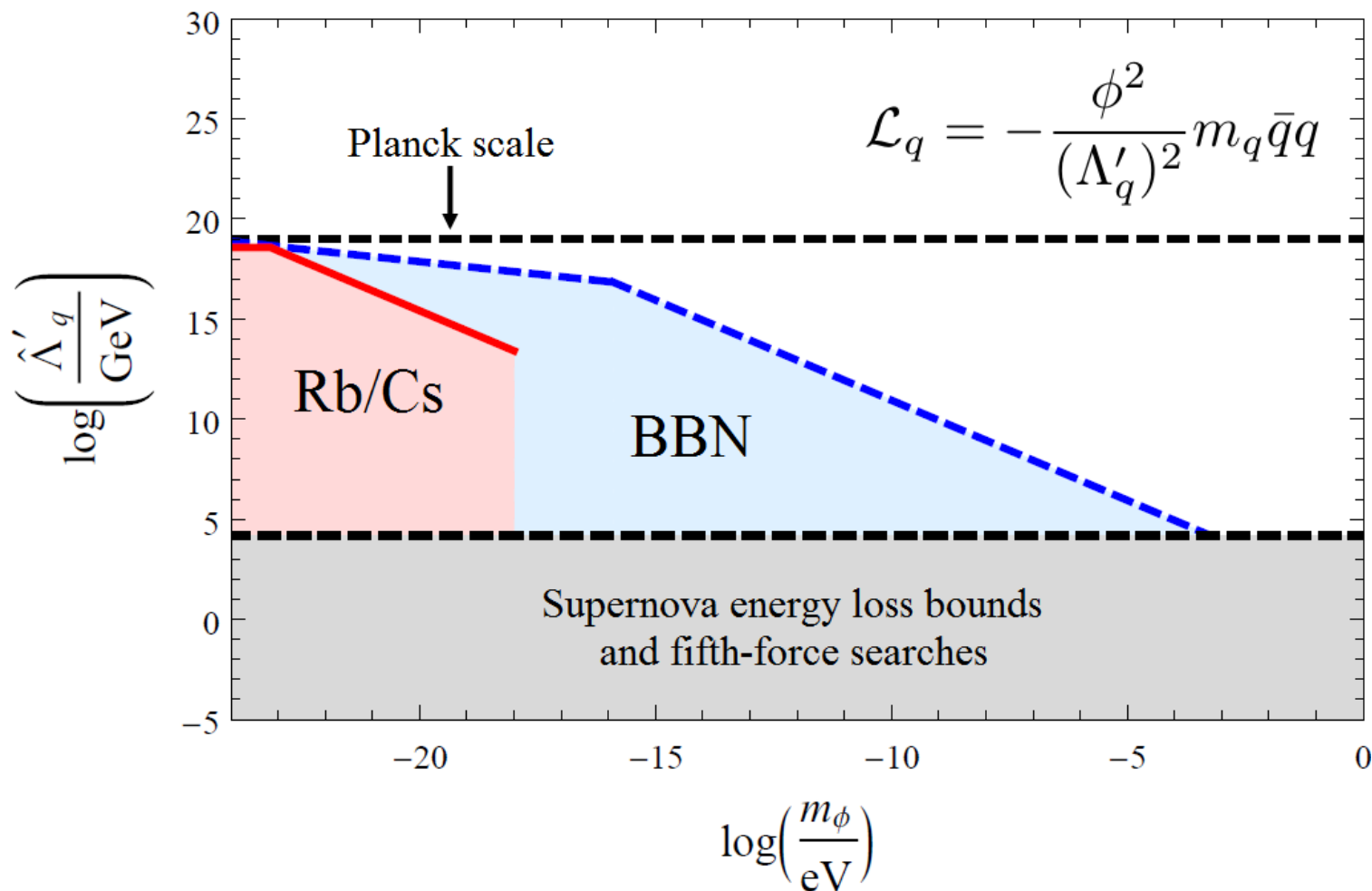
15 orders of magnitude improvement!



Constraints on Quadratic Interactions of Scalar Dark Matter with Light Quarks

BBN and Rb/Cs constraints:

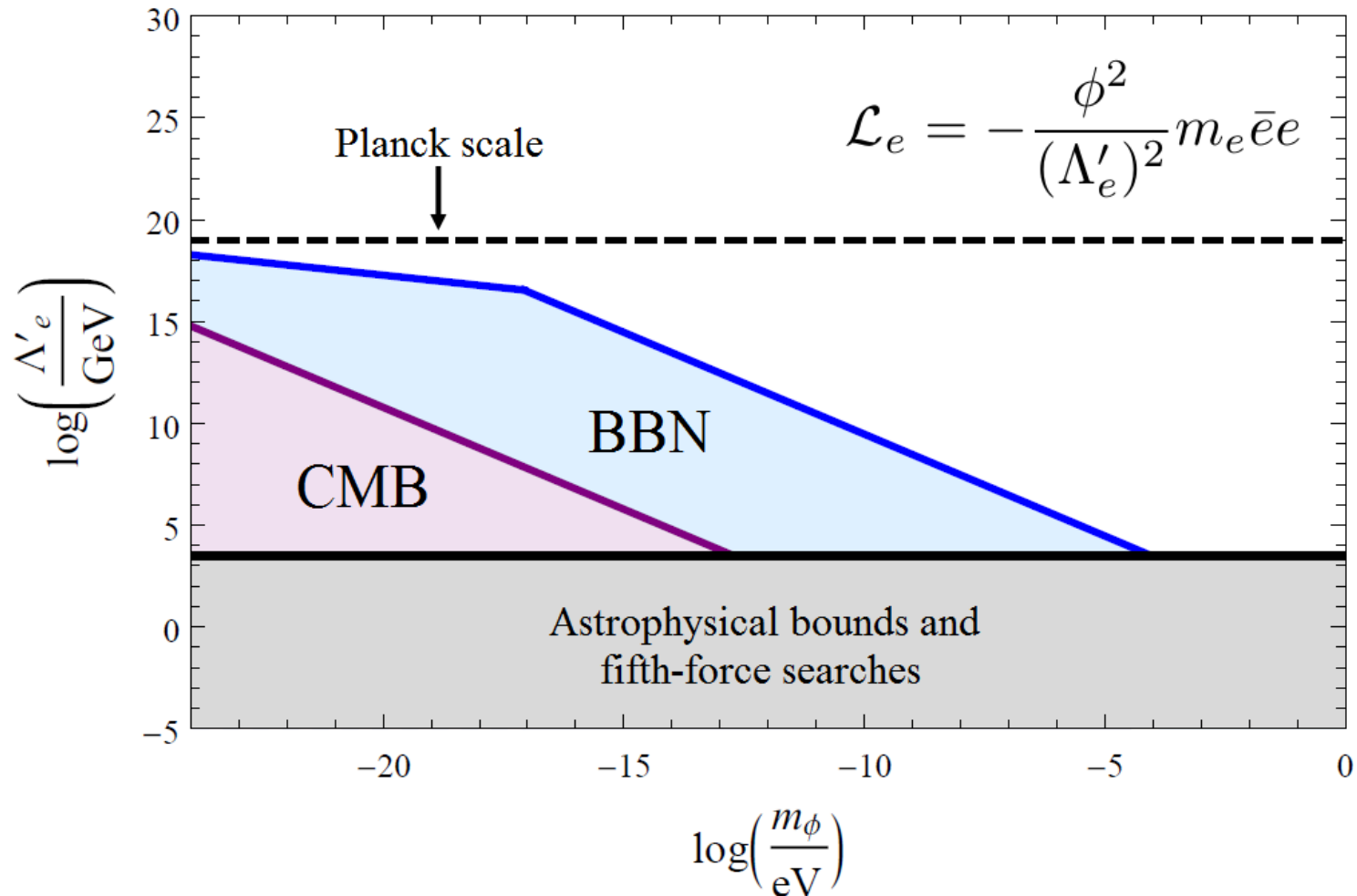
[Stadnik, Flambaum, *PRL* 115, 201301 (2015) + arXiv:1605.04028]



Constraints on Quadratic Interaction of Scalar Dark Matter with the Electron

BBN and CMB constraints:

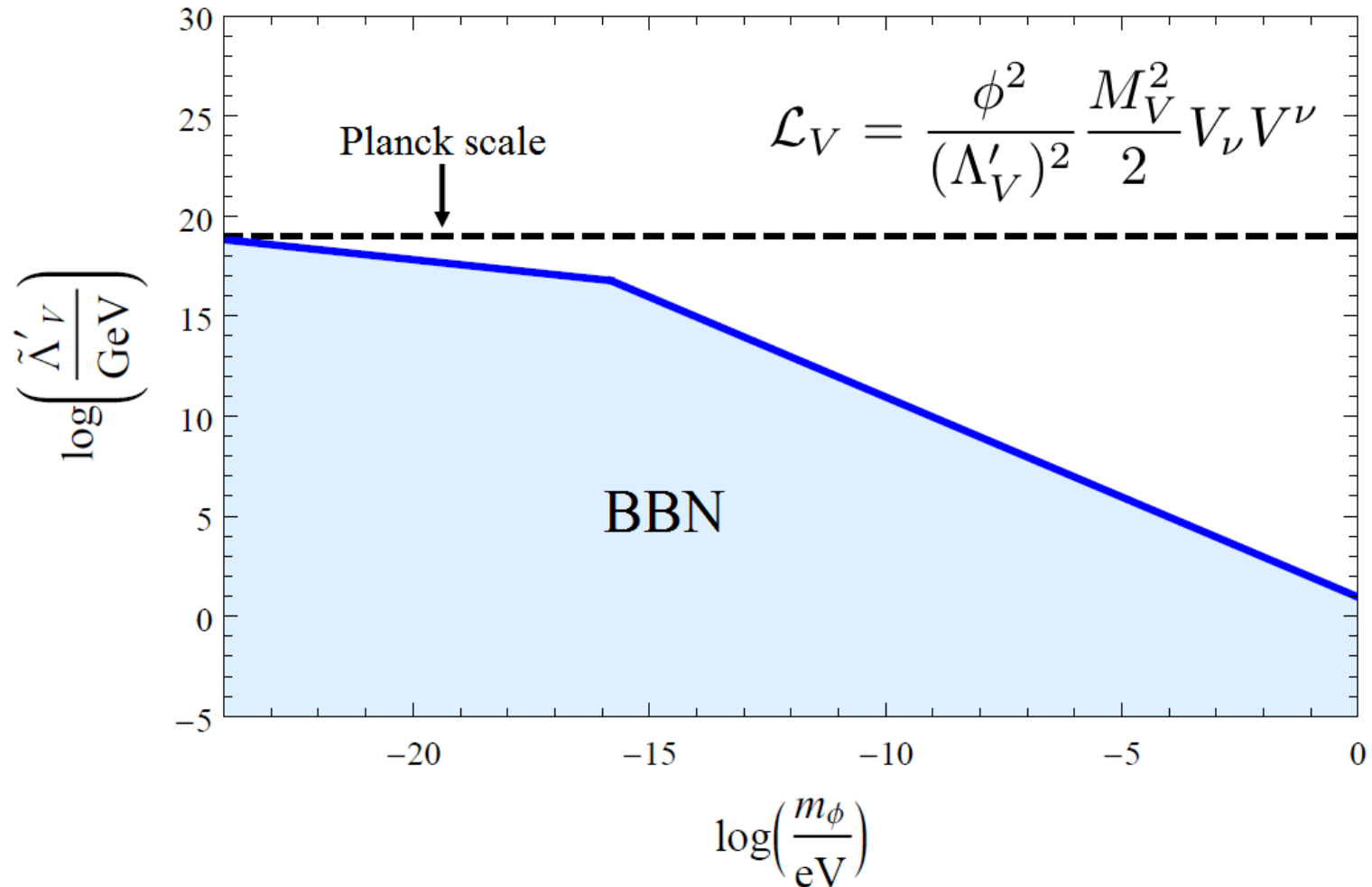
[Stadnik, Flambaum, *PRL* 115, 201301 (2015)]



Constraints on Quadratic Interactions of Scalar Dark Matter with W and Z Bosons

BBN constraints:

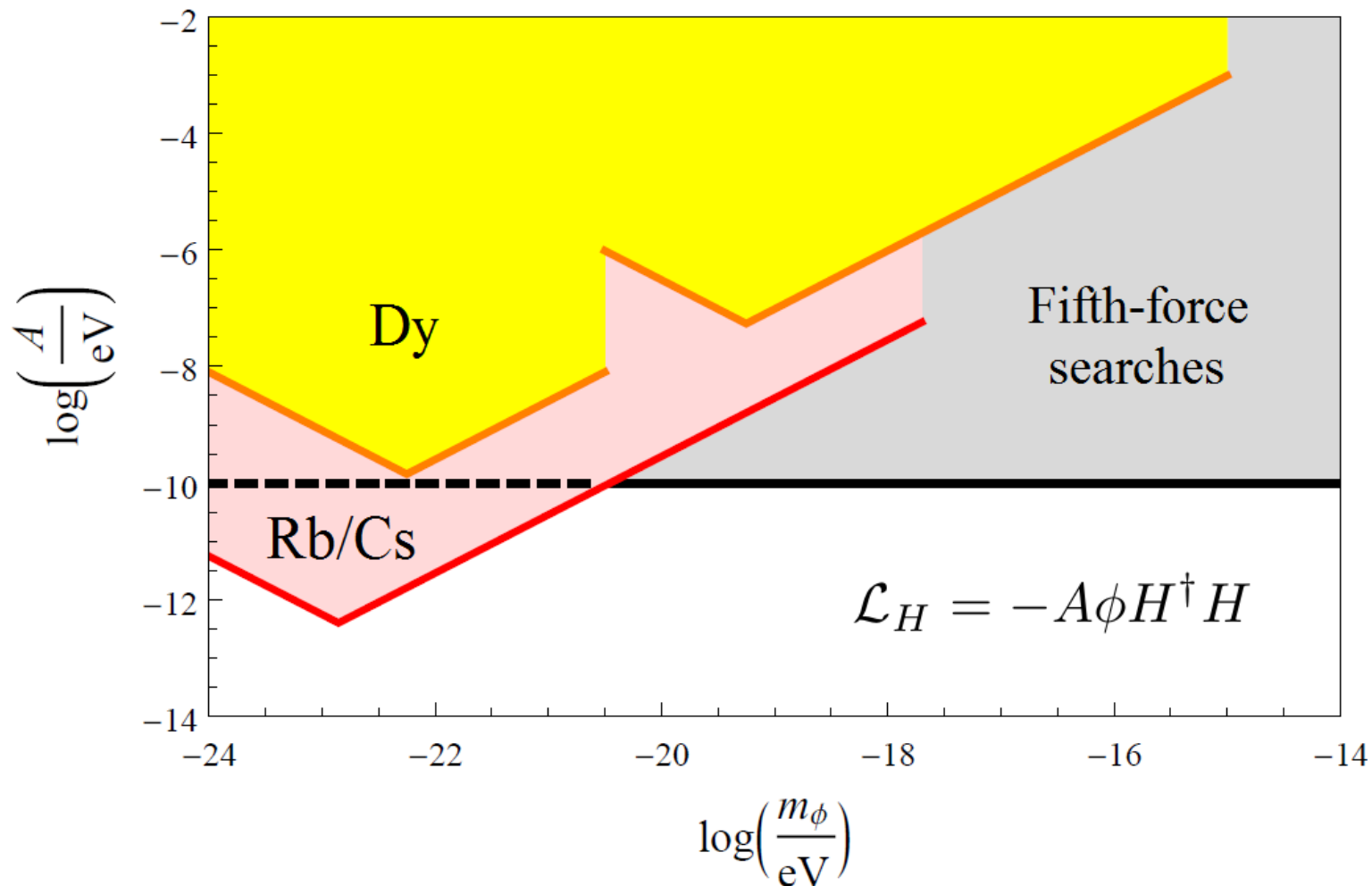
[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]



Constraints on Linear Interaction of Scalar Dark Matter with the Higgs Boson

Dy and Rb/Cs constraints:

[Stadnik, Flambaum, arXiv:1605.04028]



Topological Defect Dark Matter

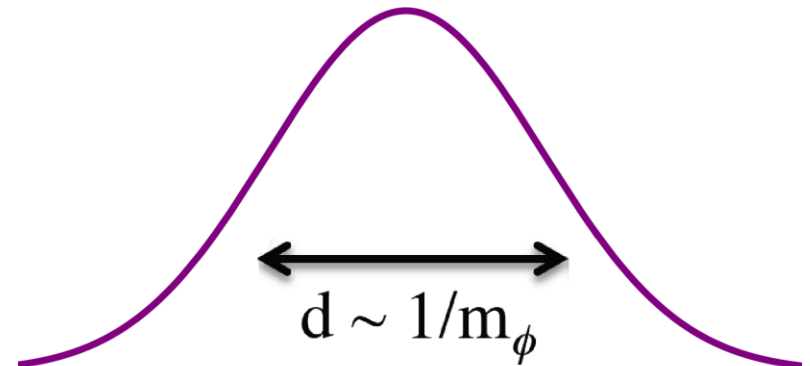
Take a simple scalar field and give it a self-potential, e.g. $V(\phi) = \lambda(\phi^2 - v^2)^2$. If $\phi = -v$ at $x = -\infty$ and $\phi = +v$ at $x = +\infty$, then a stable domain wall will form in between, e.g. $\phi = v \tanh(xm_\phi)$ with $m_\phi = \lambda^{1/2} v$.

The characteristic “span” of this object is $d \sim 1/m_\phi$, and it is carrying energy per area $\sim v^2/d \sim v^2 m_\phi$. Networks of such topological defects can give contributions to dark matter/dark energy and act as seeds for structure formation.

0D object – a *Monopole*

1D object – a *String*

2D object – a *Domain wall*



Searching for Topological Defects

Detection of topological defects via transient-in-time effects requires searching for **correlated signals** using a terrestrial or space-based **network of detectors**.

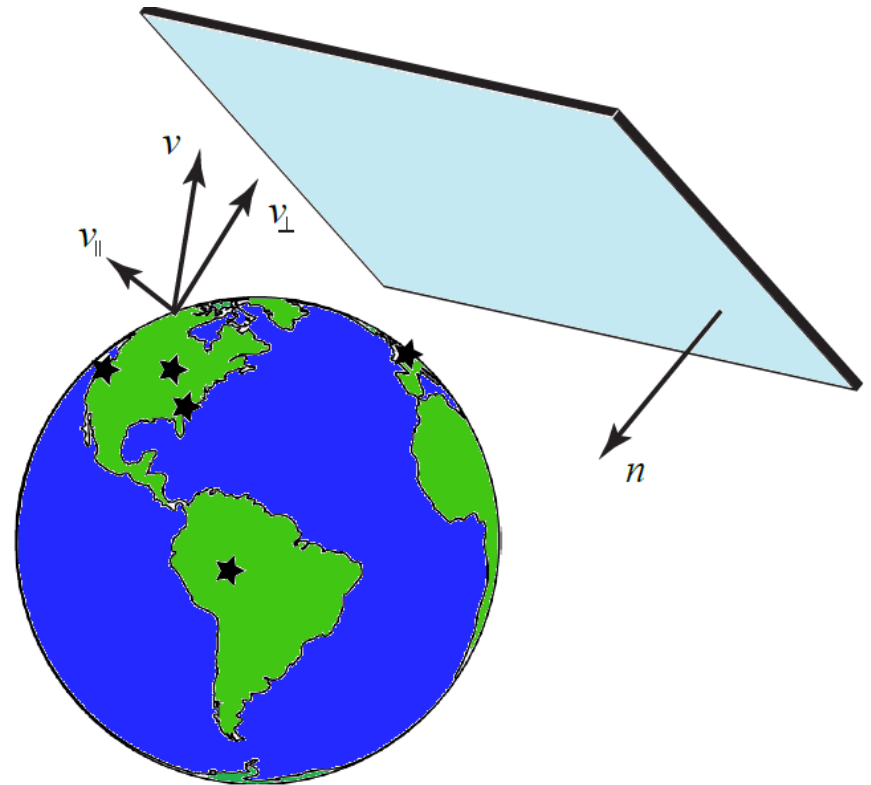
Recent proposals include:

Magnetometers [Pospelov *et al.*, *PRL* **110**, 021803 (2013)]

Pulsar Timing [Stadnik, Flambaum, *PRL* **113**, 151301 (2014)]

Atomic Clocks [Derevianko, Pospelov, *Nature Physics* **10**, 933 (2014)]

Laser Interferometers
[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); arXiv:1511.00447]



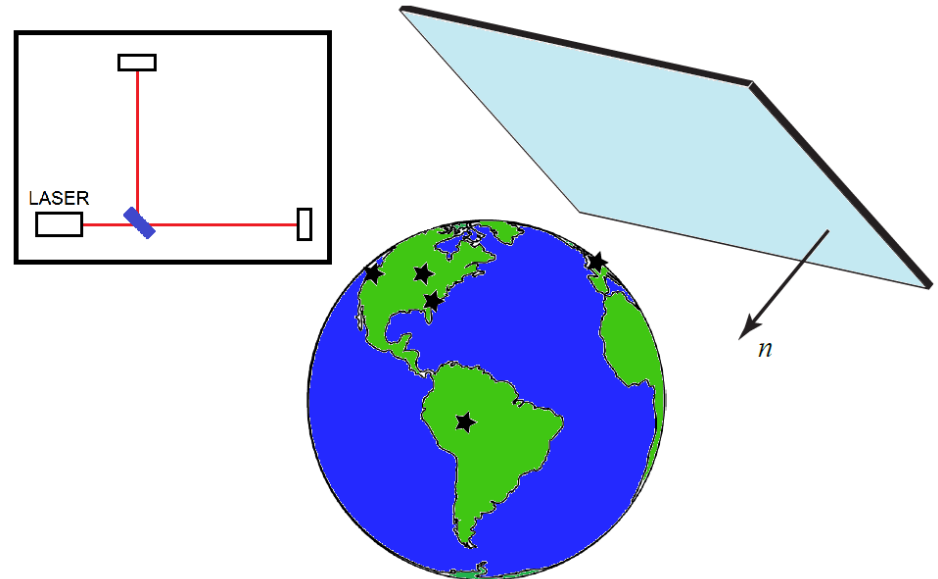
Global Network of Laser/Maser Interferometers (LIGO, Virgo, GEO600, TAMA300)

Stadnik, Flambaum, Phys. Rev.Lett. 2015 + Ongoing collaboration with LIGO and VIRGO (Klimenko, Mitselmakher)

$$\mathcal{L}_{\text{int}}^f = - \sum_{f=e,p,n} m_f \left(\frac{\phi c}{\Lambda'_f} \right)^2 \bar{f} f$$

$$\mathcal{L}_{\text{int}}^\gamma = \left(\frac{\phi}{\Lambda'_\gamma} \right)^2 \frac{F_{\mu\nu} F^{\mu\nu}}{4}$$

Topological defects, which consist of scalar particles, temporarily alter the masses of the electron, proton, neutron and photon, as well as the fine-structure constant α . This may produce a difference in the phases of light propagating in the two arms ($\Phi = kL$). One can search for defects through correlated signals in a global network of interferometers ($v_{\text{TD}} \sim 10^{-3} c$).

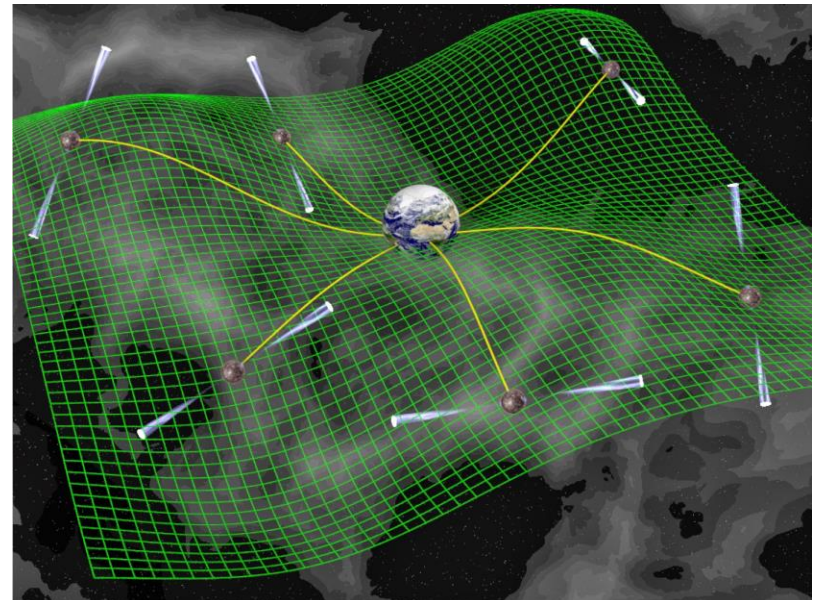
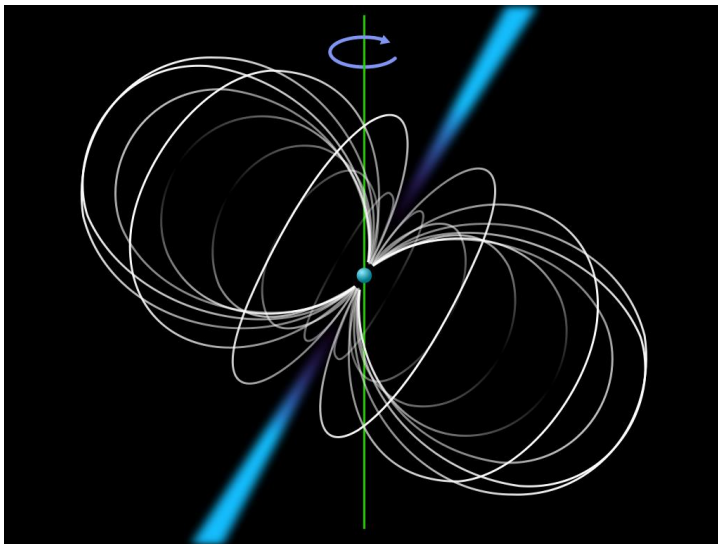


Pulsar Timing

[Stadnik, Flambaum, *PRL* **113**, 151301 (2014)]

Pulsars are highly-magnetised, rapidly rotating neutron stars ($T_{\text{rot}} \sim 1 \text{ ms} - 10 \text{ s}$), with very high long-term period stability ($\sim 10^{-15}$).

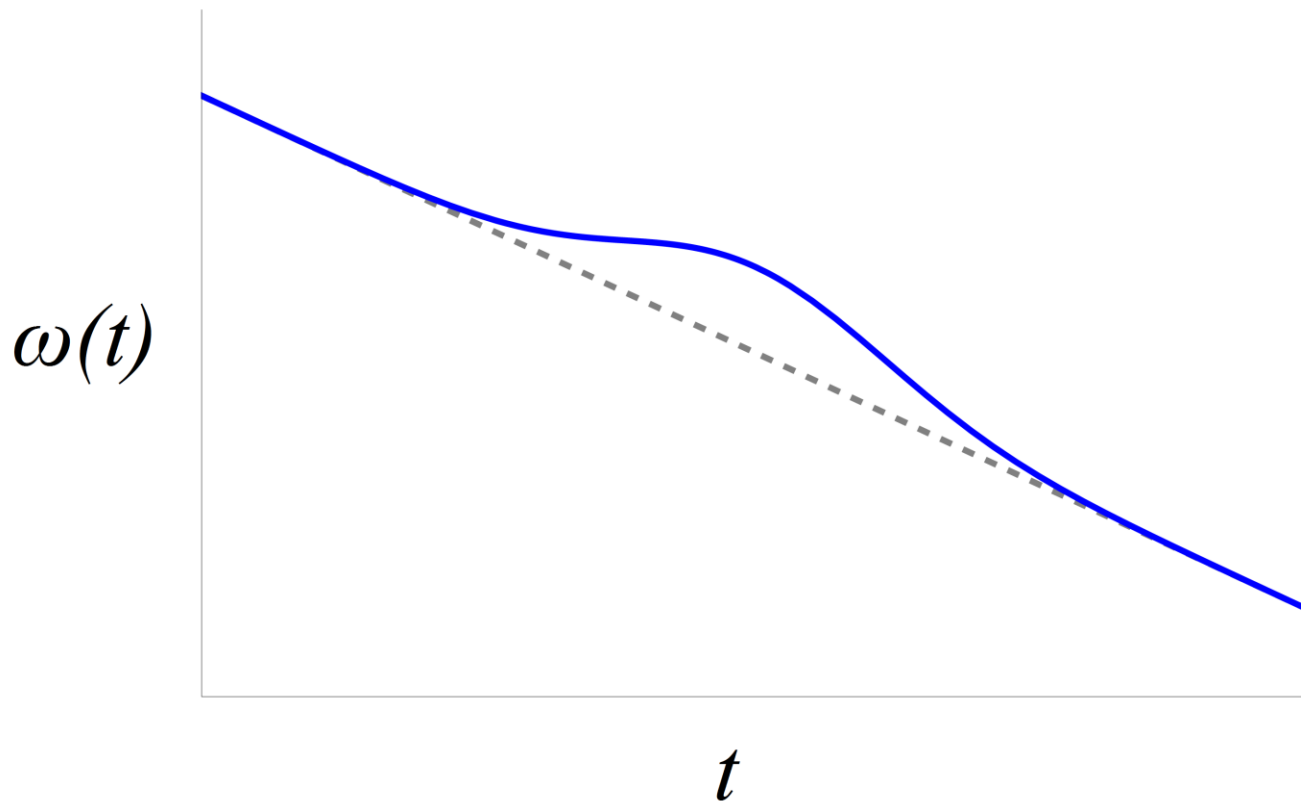
A network of pulsars can be used to search for correlated effects ($v_{\text{TD}} \sim 10^{-3}c$) produced by dark matter topological defects.



Pulsar Timing

[Stadnik, Flambaum, *PRL* **113**, 151301 (2014)]

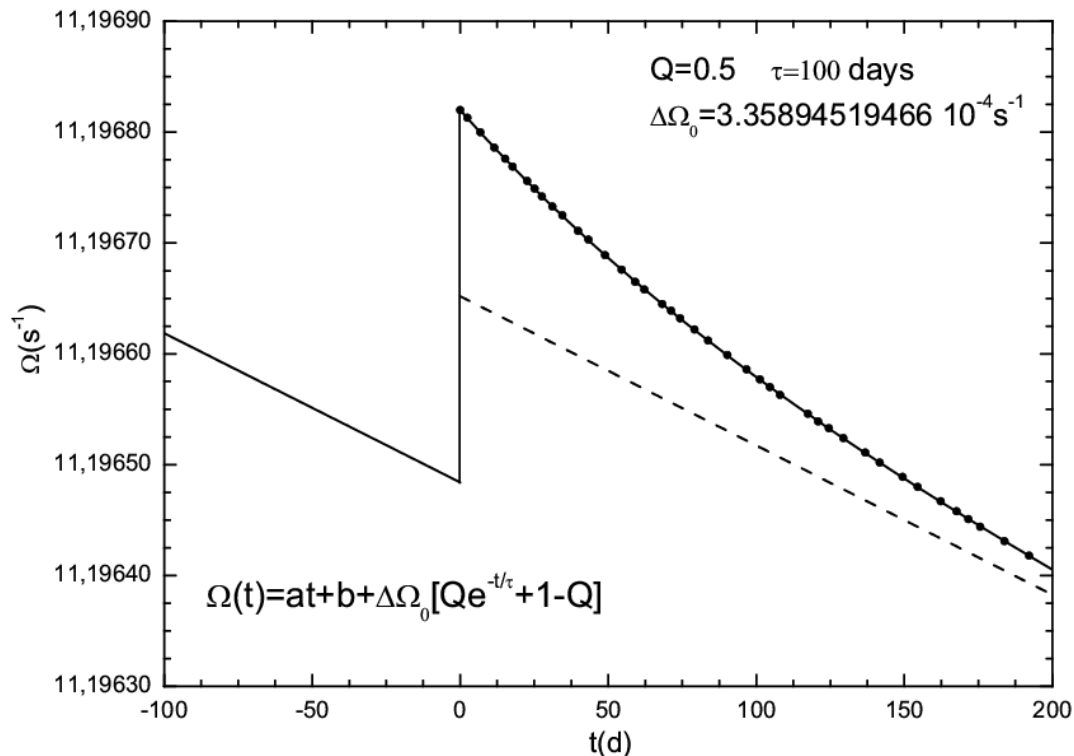
Adiabatic passage of a topological defect through a pulsar produces a Gaussian-shaped modulation in the pulsar rotational frequency profile



Pulsar Timing

[Stadnik, Flambaum, *PRL* **113**, 151301 (2014)]

Non-adiabatic passage of a topological defect through a pulsar may trigger a pulsar 'glitch' event (which have already been observed, but their underlying cause is still disputed).



Glitch Theory

- Model pulsar as 2-component system: neutron superfluid core, surrounded by neutron crust
- 2 components can rotate independently of one another
- Rotation of neutron superfluid core quantified by area density of quantised vortices (which carry angular momentum)
- Strong vortex 'pinning' to neutron crust
- Can vortices be unpinned by topological defect?
- Vortices avalanche = pulsar glitch

Low-mass Spin-0 Dark Matter

Dark Matter



**Pseudoscalars
(Axions, ALPs):**
Odd-parity

Axions explain the absence of CP violation in the strong interaction and are a leading dark matter candidate

**→ Oscillating spin-dependent effects,
 P, T , Lorentz and Einstein symmetry violation**

- Atomic magnetometry
- Ultracold neutrons
- Solid-state magnetometry

“Axion Wind” Spin-Precession Effect

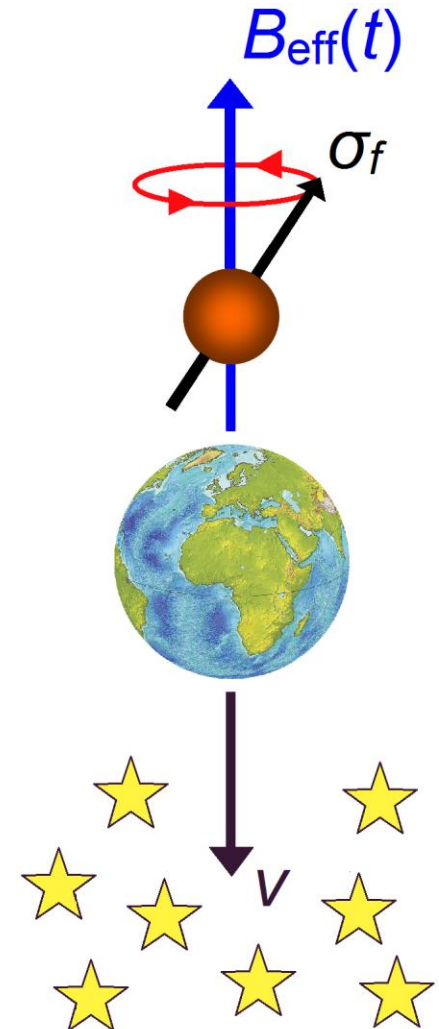
[Flambaum, *Patras Workshop*, 2013], [Graham, Rajendran, *PRD* **88**, 035023 (2013)],
[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Motion of Earth through galactic axions gives rise to the interaction of fermion spins with a time-dependent *pseudo*-magnetic field $B_{\text{eff}}(t)$, producing **spin-precession effects**.

$$\mathcal{L}_{aff} = -\frac{C_f}{2f_a} \partial_i [a_0 \cos(\varepsilon_a t - p_a \cdot r)] \bar{f} \gamma^i \gamma^5 f$$

$$\Rightarrow H_{\text{eff}}(t) \simeq \underbrace{\frac{C_f a_0}{2f_a} \sin(m_a t) p_a \cdot \sigma_f}_{B_{\text{eff}}(t)}$$

$B_{\text{eff}}(t)$



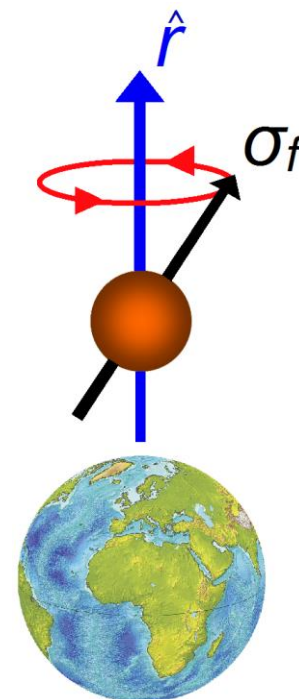
Axion-Induced Oscillating Spin-Gravity Coupling

[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Distortion of axion field by gravitational field of Sun or Earth induces **oscillating spin-gravity couplings**.

$$\mathcal{L}_{aff} = -\frac{C_f}{2f_a} \partial_i [a_0(r) \cos(\varepsilon_a t - p_a \cdot r)] \bar{f} \gamma^i \gamma^5 f$$

$$\Rightarrow H'_{\text{eff}}(t) \propto \frac{C_f a_0}{f_a} \sin(m_a t) \sigma_f \cdot \hat{r}$$



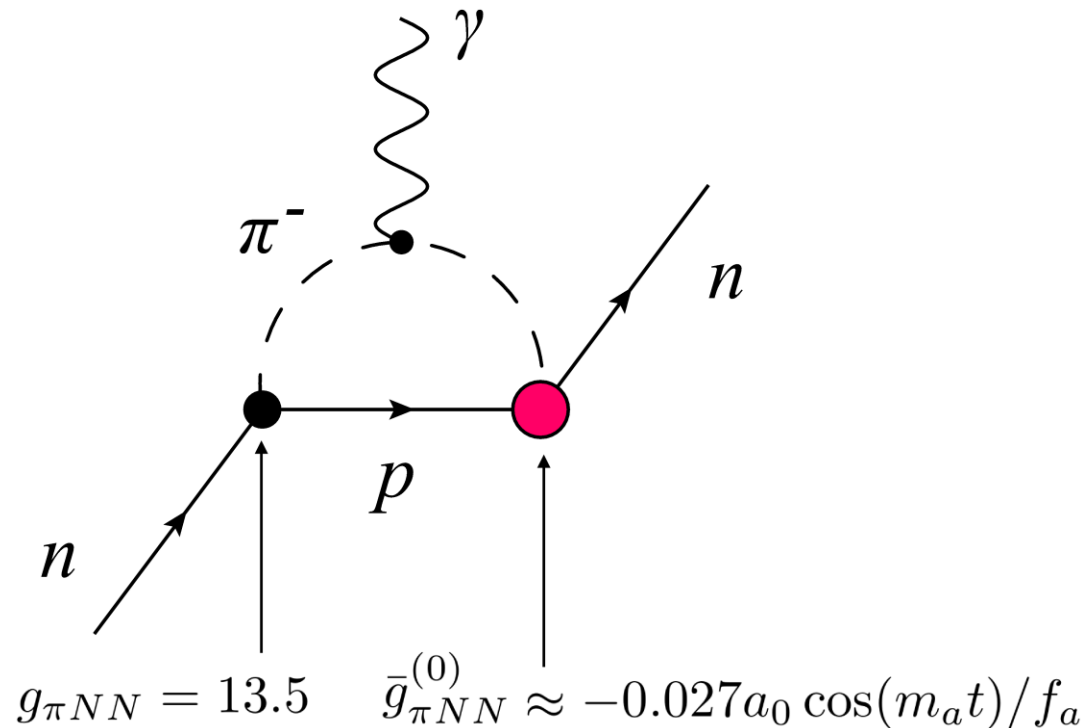
Spin-axion momentum and axion-induced oscillating spin-gravity couplings to nucleons may have isotopic dependence ($C_p \neq C_n$) – calculations of proton and neutron spin contents for nuclei of experimental interest have been performed, see, e.g., [Stadnik, Flambaum, *EPJC* **75**, 110 (2015)].

Axion-Induced Oscillating Neutron EDM

[Graham, Rajendran, *PRD* **84**, 055013 (2011)]

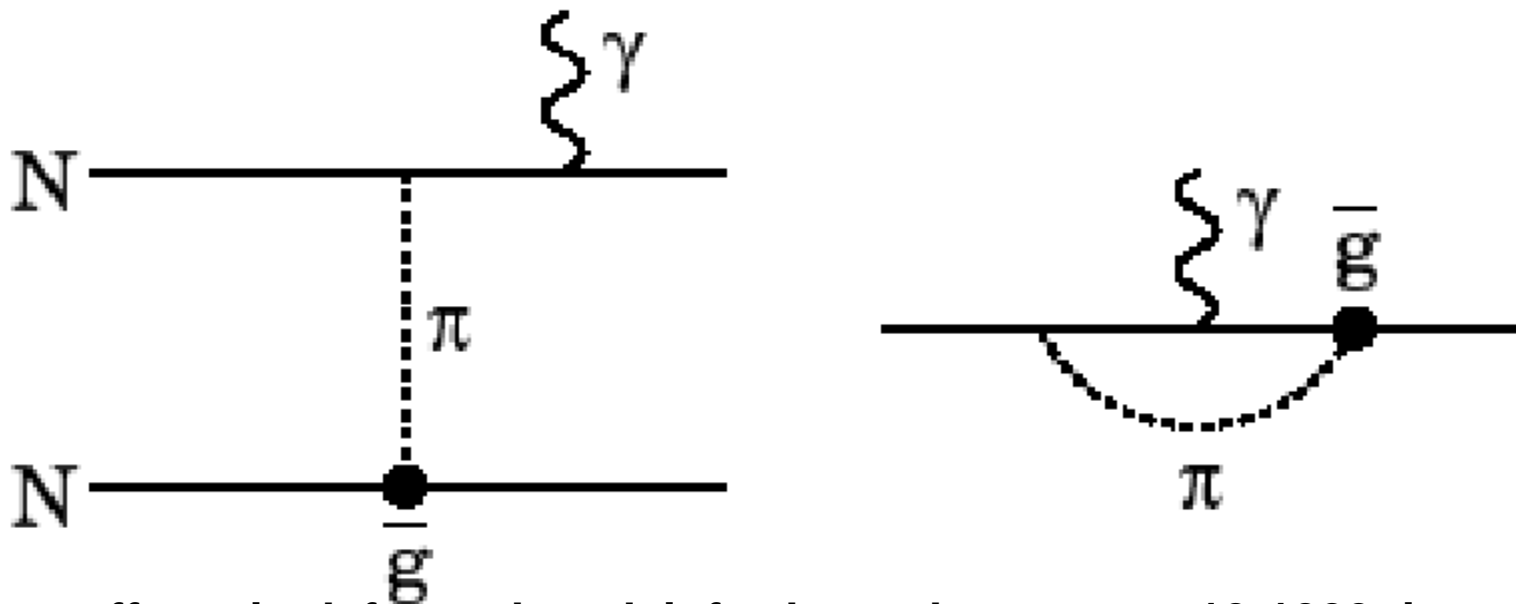
An oscillating axion field induces an **oscillating neutron electric dipole moment** via its coupling to gluons.

$$\mathcal{L}_{agg} = \frac{a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G\tilde{G} \quad d_n(t) \approx 2.4 \times 10^{-16} \frac{a_0}{f_a} \cos(m_a t) \text{ e} \cdot \text{cm}$$



Nuclear EDM: P, T -odd NN interaction gives 40 times larger contribution than intrinsic nucleon EDM

[Sushkov, Flambaum, Khriplovich, *JETP* **60**, 873 (1984)]



Collective effects in deformed nuclei: further enhancement 10-1000 times, magnetic quadrupole F. 1994, Schiff moment A. F. S. 1996

Atomic EDM produced by nuclear magnetic quadrupole moment

Magnetic interaction is not screened!

MQM produced by nuclear T,P-odd forces
(Khriplovich, Sushkov, Flambaum)

Collective enhancement in deformed nuclei
(Flambaum). T,P-odd nuclear interaction
produces spin hedgehog- correlation (s r)

Spherical – magnetic monopole forbidden

Deformed- collective magnetic quadrupole

Paramagnetic molecules ThO, TaN, YbF, ...
(Flambaum, DeMille, Kozlov)

Axion-Induced Oscillating Atomic and Molecular EDMs

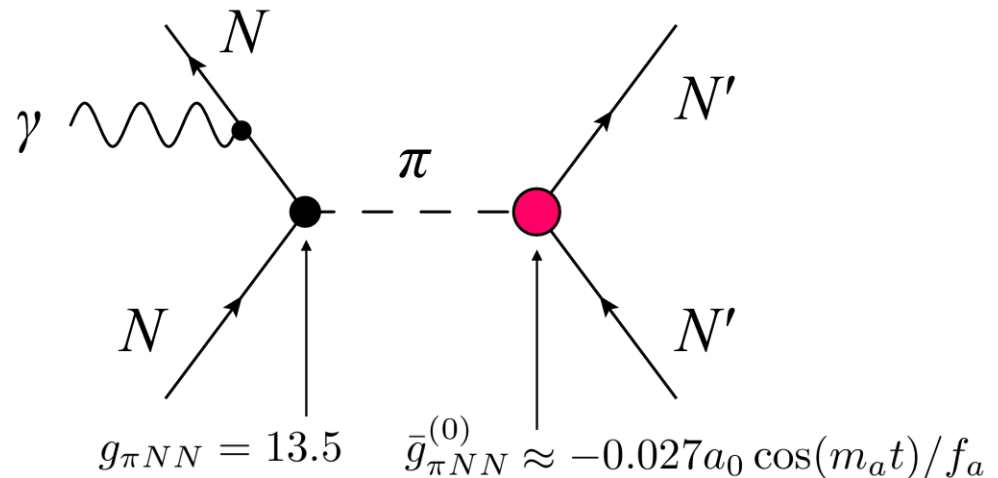
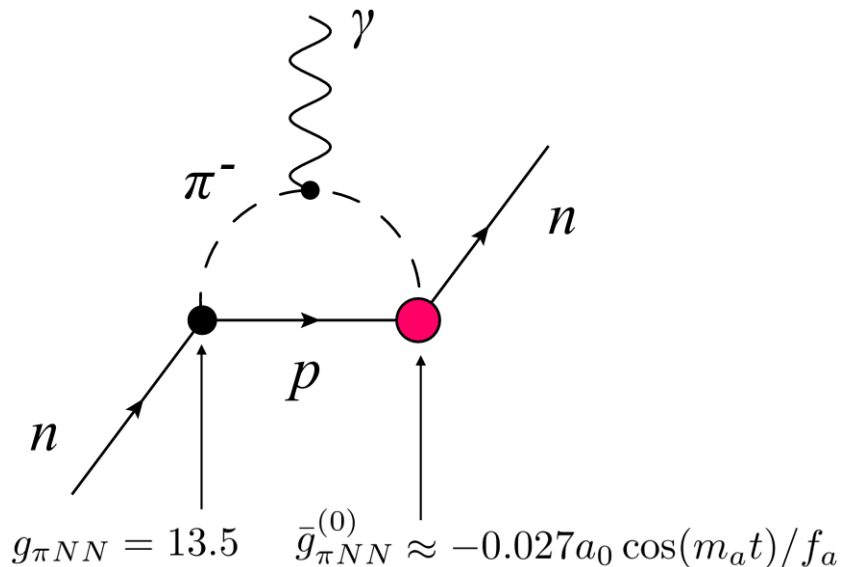
[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Oscillating atomic and molecular EDMs are induced through oscillating Schiff ($J \geq 0$) and oscillating magnetic quadrupole ($J \geq 1/2$, **no Schiff screening**) moments of nuclei, which arise from *intrinsic oscillating nucleon EDMs* and *oscillating P, T -violating intranuclear forces* (**larger by factor of several – 1000**).

$$\mathcal{L}_{agg} = \frac{a_0 \cos(m_a t)}{f_a} \frac{g^2}{32\pi^2} G\tilde{G}$$

$$d(^{199}\text{Hg})(t) \approx -1.8 \times 10^{-19} \frac{a_0}{f_a} \cos(m_a t) e \cdot \text{cm}$$

$$d(^{225}\text{Ra})(t) \approx 9.3 \times 10^{-17} \frac{a_0}{f_a} \cos(m_a t) e \cdot \text{cm}$$

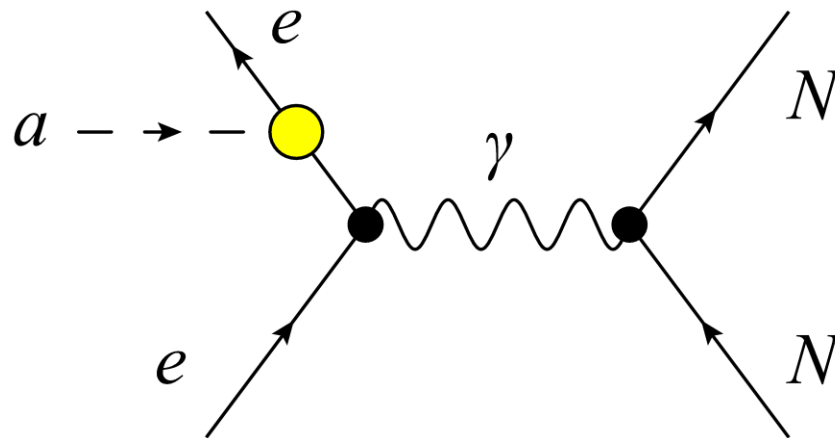


Axion-Induced Oscillating EDMs of Paramagnetic Atoms and Molecules

[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)], [Roberts, Stadnik, Dzuba, Flambaum, Leeper, Budker, *PRL* **113**, 081601 (2014) + *PRD* **90**, 096005 (2014)]

In *paramagnetic* atoms and molecules, **oscillating EDMs** are also induced through *mixing of opposite-parity states* via the interaction of the oscillating axion field with atomic/molecular electrons.

$$\mathcal{L}_{aee} = -\frac{C_e}{2f_a} \partial_0 [a_0 \cos(m_a t)] \bar{e} \gamma^0 \gamma^5 e \quad d_{\text{atomic}}(t) \sim -\frac{C_e a_0 m_a^2 \alpha_s}{f_a e} \cos(m_a t)$$



Relativistic effects increase ionisation by dark matter scattering on electrons by up to 3 orders of magnitude!

[Roberts, Flambaum, Gribakin, PRL **116**, 023201 (2016)]

- Important for numerous existing and future dark matter detectors.
- Detailed relativistic many-body calculations in [Roberts, Dzuba, Flambaum, Pospelov, Stadnik, arXiv:1604.04559]
- DAMA collaboration claims detection of dark matter, others – no detection. Possible explanation: scattering of dark matter on electrons (instead of scattering on nuclei).
- Our calculations show tension between DAMA and XENON results.

Conclusions

- New classes of dark matter effects that are **linear** in the underlying interaction constant (traditionally-sought effects of dark matter scale as second or fourth power)
- **15 orders of magnitude improvement** on quadratic interactions of scalar dark matter with the photon, electron, and light quarks (u, d).
- Improved limits on linear interaction with the Higgs boson.
- **First limits** on linear and quadratic interactions of scalar dark matter with vector bosons (W^+, W^-, Z^0)
- Oscillating effects of variation of fundamental constants and violation of the fundamental symmetries: Parity, Time reversal, Lorentz, Einstein equivalence principle
- Enormous potential for low-energy atomic experiments to search for dark matter with unprecedented sensitivity

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References (Scalar Dark Matter)

Y. V. Stadnik and V. V. Flambaum. *Can Dark Matter Induce Cosmological Evolution of the Fundamental Constants of Nature?* Physical Review Letters **115**, 201301 (2015). arXiv:1503.08540.

Y. V. Stadnik and V. V. Flambaum. *Searching for Dark Matter and Variation of Fundamental Constants with Laser and Maser Interferometry.* Physical Review Letters **114**, 161301 (2015). arXiv:1412.7801.

Y. V. Stadnik and V. V. Flambaum. *Improved limits on interactions of low-mass spin-0 dark matter from atomic clock spectroscopy.* arXiv:1605.04028.

Y. V. Stadnik and V. V. Flambaum. *Enhanced effects of variation of the fundamental constants in laser interferometers and application to dark matter detection.* arXiv:1511.00447.

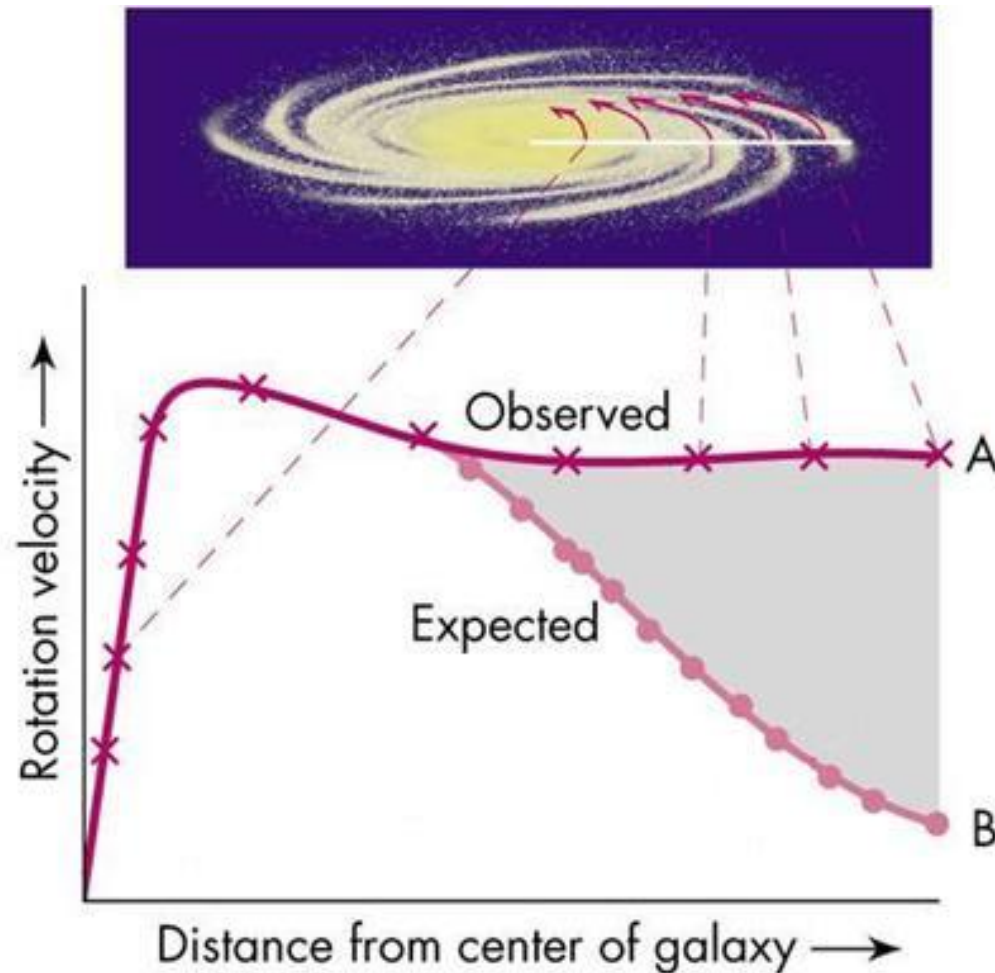
Y. V. Stadnik and V. V. Flambaum. *Searching for Topological Defect Dark Matter via Nongravitational Signatures.* Physical Review Letters **113**, 151301 (2014). arXiv:1405.5337.

References (Axion Dark Matter)

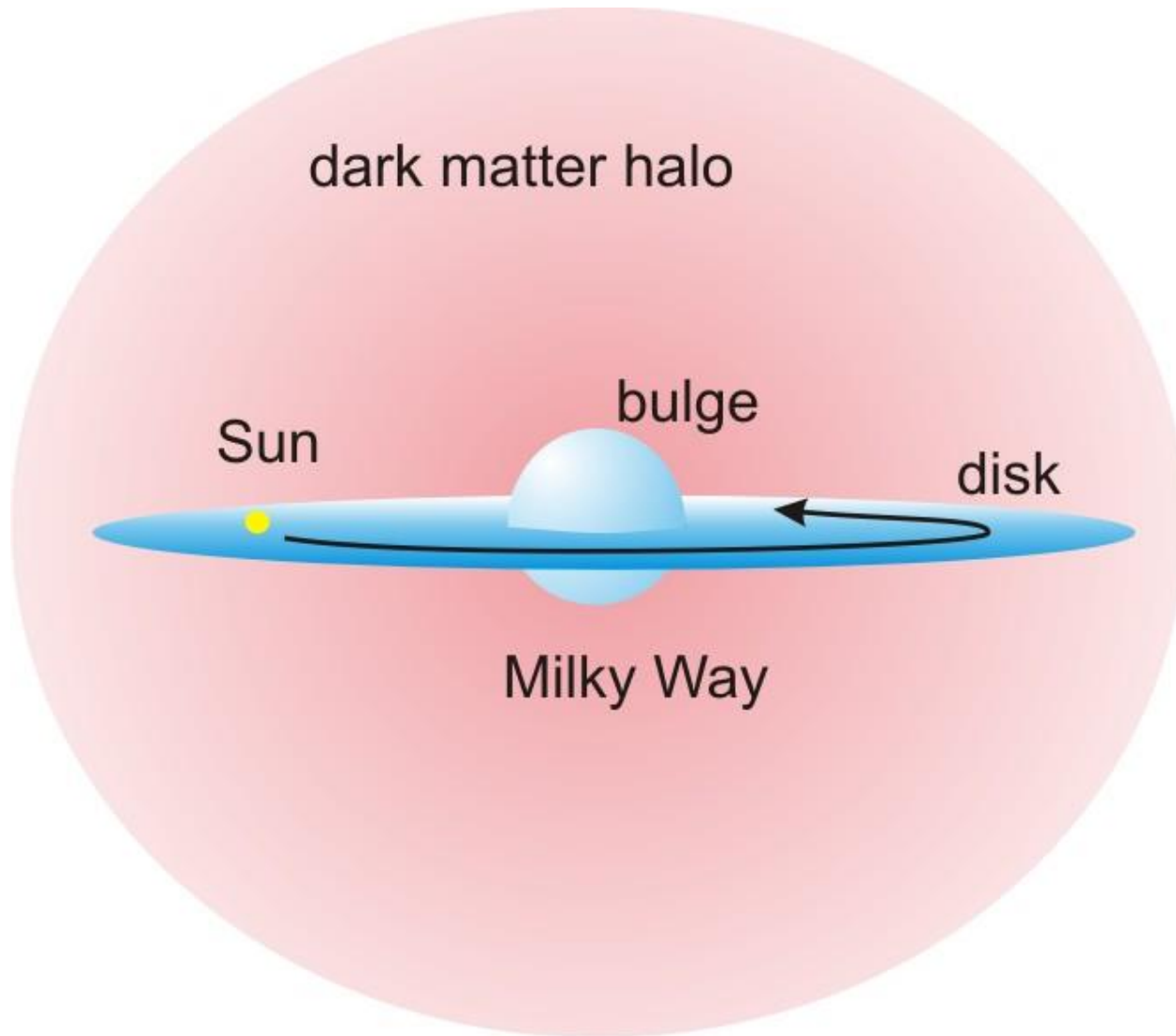
- Y. V. Stadnik and V. V. Flambaum. *Axion-induced effects in atoms, molecules and nuclei: Parity nonconservation, anapole moments, electric dipole moments, and spin-gravity and spin-axion momentum couplings*. Physical Review D **89**, 043522 (2014). arXiv:1312.6667.
- B. M. Roberts, Y. V. Stadnik, V. A. Dzuba, V. V. Flambaum, N. Leeper and D. Budker. *Limiting P-odd interactions of Cosmic Fields with Electrons, Protons and Neutrons*. Physical Review Letters **113**, 081601 (2014). arXiv:1404.2723.
- B. M. Roberts, Y. V. Stadnik, V. A. Dzuba, V. V. Flambaum, N. Leeper and D. Budker. *Parity-violating interactions of cosmic fields with atoms, molecules and nuclei: Concepts and calculations for laboratory searches and extracting limits*. Physical Review D **90**, 096005 (2014). arXiv:1409.2564.
- Y. V. Stadnik and V. V. Flambaum. *Nuclear spin-dependent interactions: searches for WIMP, axion and topological defect dark matter, and tests of fundamental symmetries*. European Physical Journal C **75**, 110 (2015). arXiv:1408.2184.

Motivation

Studies of galactic rotation curves (Zwicky 1930s; Rubin *et al.* 1970s)



Motivation



Astrophysical Constraints on 'Slow' Drifts in Fundamental Constants Induced by Scalar Dark Matter (BBN)

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

- Largest effects of scalar dark matter are in the early Universe (highest $\rho_{\text{DM}} \Rightarrow$ highest φ_0^2).
- Earliest cosmological epoch that we can probe is Big Bang nucleosynthesis (from $t_{\text{weak}} \approx 1\text{s}$ until $t_{\text{BBN}} \approx 3\text{ min}$).
- Primordial ${}^4\text{He}$ abundance is sensitive to relative abundance of neutrons to protons (almost all neutrons are bound in ${}^4\text{He}$ by the end of BBN).

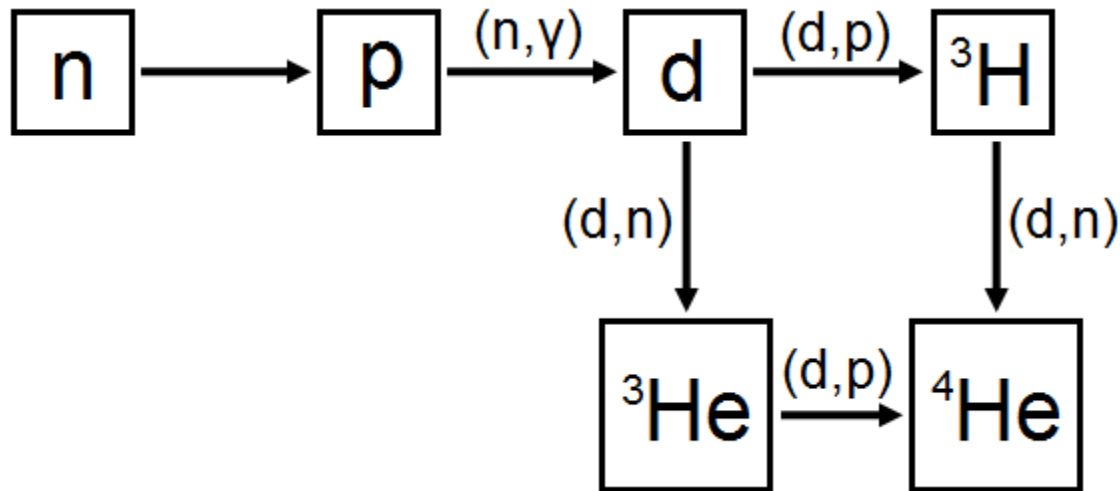
Weak interactions: freeze-out of weak interactions occurs at $t_{\text{weak}} \approx 1\text{s}$ ($T_{\text{weak}} \approx 0.75\text{ MeV}$).

$$\begin{aligned} p + e^- &\rightleftharpoons n + \nu \\ n + e^+ &\rightleftharpoons p + \bar{\nu} \end{aligned} \quad \left(\frac{n}{p}\right)_{\text{weak}} = e^{-(m_n - m_p)/T_{\text{weak}}}$$

Astrophysical Constraints on 'Slow' Drifts in Fundamental Constants Induced by Scalar Dark Matter (BBN)

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

BBN reactions: reaction channels that produce ${}^4\text{He}$ last until $t_{\text{BBN}} \approx 3 \text{ min}$ ($T_{\text{BBN}} \approx 60 \text{ keV}$).



$$\frac{\Delta Y_p({}^4\text{He})}{Y_p({}^4\text{He})} \approx \frac{\Delta(n/p)_{\text{weak}}}{(n/p)_{\text{weak}}} - \Delta \left[\int_{t_{\text{weak}}}^{t_{\text{BBN}}} \Gamma_n(t) dt \right] \Rightarrow \text{Limits on } \Lambda'_x$$

Astrophysical Constraints on 'Slow' Drifts in Fundamental Constants Induced by Scalar Dark Matter (CMB)

[Stadnik, Flambaum, *PRL* **115**, 201301 (2015)]

- Weaker astrophysical constraints come from CMB measurements (lower ρ_{DM}).
- Variations in α and m_e at the time of electron-proton recombination affect the ionisation fraction and Thomson scattering cross section, $\sigma_{\text{Thomson}} = 8\pi\alpha^2/3m_e^2$, changing the mean-free-path length of photons at recombination and leaving distinct signatures in the CMB angular power spectrum.

$$\Lambda'_\gamma \gtrsim \frac{1 \text{ eV}^2}{m_\phi}, \quad \Lambda'_e \gtrsim \frac{0.6 \text{ eV}^2}{m_\phi}$$

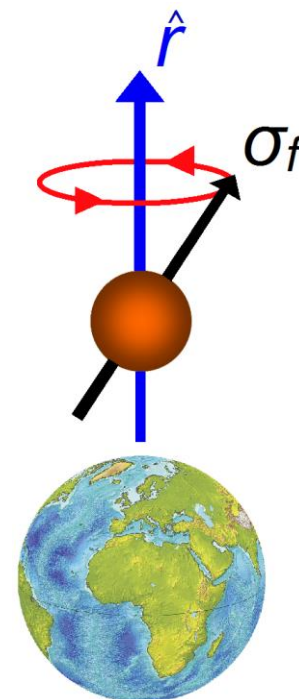
Axion-Induced Oscillating Spin-Gravity Coupling

[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)]

Distortion of axion field by gravitational field of Sun or Earth induces **oscillating spin-gravity couplings**.

$$\mathcal{L}_{aff} = -\frac{C_f}{2f_a} \partial_i [a_0(r) \cos(\varepsilon_a t - p_a \cdot r)] \bar{f} \gamma^i \gamma^5 f$$

$$\Rightarrow H'_{\text{eff}}(t) \propto \frac{C_f a_0}{f_a} \sin(m_a t) \sigma_f \cdot \hat{r}$$



Spin-axion momentum and axion-induced oscillating spin-gravity couplings to nucleons may have isotopic dependence ($C_p \neq C_n$) – calculations of proton and neutron spin contents for nuclei of experimental interest have been performed, see, e.g., [Stadnik, Flambaum, *EPJC* **75**, 110 (2015)].

Axion-Induced Oscillating Parity Non-conservation in Atoms and Molecules

[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)], [Roberts, Stadnik, Dzuba, Flambaum, Leefer, Budker, *PRL* **113**, 081601 (2014) + *PRD* **90**, 096005 (2014)]

Interaction of the oscillating axion field with atomic/molecular electrons mixes opposite-parity states, producing **oscillating PNC effects in atoms and molecules**.

$$\mathcal{L}_{aee} = -\frac{C_e}{2f_a} \partial_0 [a_0 \cos(m_a t)] \bar{e} \gamma^0 \gamma^5 e \quad E_{\text{PNC}}(t) = -\frac{C_e a_0 m_a}{2f_a} \sin(m_a t) K_{\text{PNC}}$$

Axion-induced oscillating atomic PNC effects are determined entirely by relativistic corrections (in the non-relativistic approximation, $K_{\text{PNC}} = 0$)*.

* Compare with the Standard Model *static* atomic PNC effects in atoms, which are dominated by Z^0 -boson exchange between atomic electrons and nucleons in the nucleus, where the effects arise already in the non-relativistic approximation.

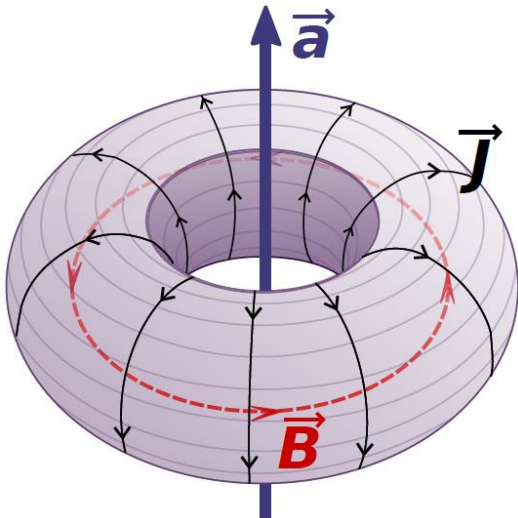
Axion-Induced Oscillating Nuclear Anapole Moments

[Stadnik, Flambaum, *PRD* **89**, 043522 (2014)], [Roberts, Stadnik, Dzuba, Flambaum, Leefler, Budker, *PRL* **113**, 081601 (2014) + *PRD* **90**, 096005 (2014)]

Interaction of the oscillating axion field with nucleons in nuclei induces **oscillating nuclear anapole moments**.

$$\mathcal{L}_{aNN} = -\frac{C_N}{2f_a} \partial_0 [a_0 \cos(m_a t)] \bar{N} \gamma^0 \gamma^5 N$$

$$\mathbf{a}(t) = -\frac{C_N a_0 m_a}{f_a} \frac{\pi e \mu}{m} \frac{KI}{I(I+1)} \langle r^2 \rangle \sin(m_a t)$$



Topological Defect Dark Matter

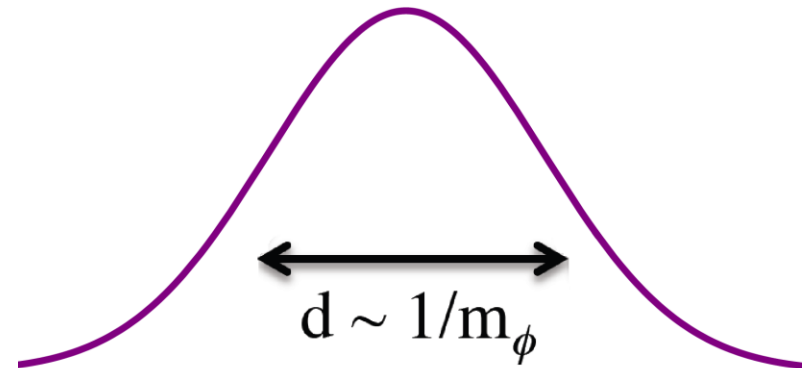
Take a simple scalar field and give it a *self-potential*, e.g. $V(\phi) = \lambda(\phi^2 - v^2)^2$. If $\phi = -v$ at $x = -\infty$ and $\phi = +v$ at $x = +\infty$, then a stable **domain wall** will form in between, e.g. $\phi(x) = v \tanh(xm_\phi)$ with $m_\phi = \lambda^{1/2} v$.

The characteristic “span” of this object is $d \sim 1/m_\phi$, and it is carrying energy per area $\sim v^2/d \sim v^2 m_\phi$. *Networks of such topological defects can give contributions to dark matter/dark energy and act as seeds for structure formation.*

0D object – a *Monopole*

1D object – a *String*

2D object – a *Domain wall*



Topological Defect Dark Matter

Topological defects may have *large amplitude*, *large transverse size* (possibly macroscopic) and *large distances* (possibly astronomical) between them.



=> *Signatures of topological defects are very different* from other forms of dark matter!

*Topological defects produce **transient-in-time effects.***

Searching for Topological Defects

Detection of topological defects via transient-in-time effects requires searching for **correlated signals** using a terrestrial or space-based **network of detectors**.

Recent proposals include:

Magnetometers [Pospelov *et al.*, *PRL* **110**, 021803 (2013)]

Pulsar Timing [Stadnik, Flambaum, *PRL* **113**, 151301 (2014)]

Atomic Clocks [Derevianko, Pospelov, *Nature Physics* **10**, 933 (2014)]

Laser Interferometers
[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); arXiv:1511.00447]

