

Quantum Simulation of High Energy Physics Models using Cold Atoms

Collaborators:

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A. Farace
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From the vacuum to the universe
Humboldt Kolleg for Particle Physics
Kitzbühl, Austria, July 1st, 2016



Simulating Physics with Computers

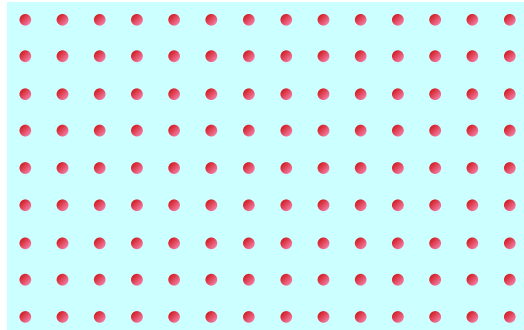
Richard P. Feynman

Department of Physics, California Institute of Technology, Pasadena, California 91107

Received May 7, 1981

1. INTRODUCTION

On the program it says this is a keynote speech—and I don't know what a keynote speech is. I do not intend in any way to suggest what should



$$c_1 |000\dots 0\rangle + c_2 |000\dots 1\rangle + \dots + c_{2^N} |111\dots 1\rangle$$

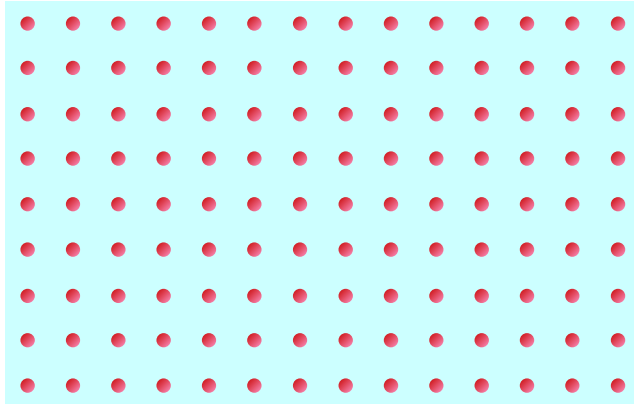


QUANTUM SIMULATORS

ANALOG



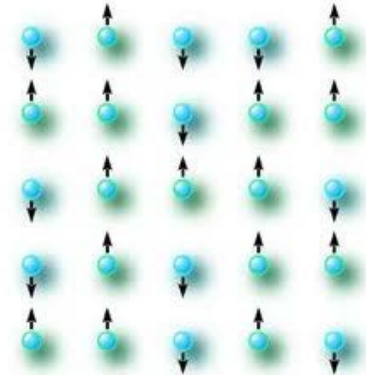
MODEL



Model Hamiltonian

$$H = \dots$$

QUANTUM SIMULATOR



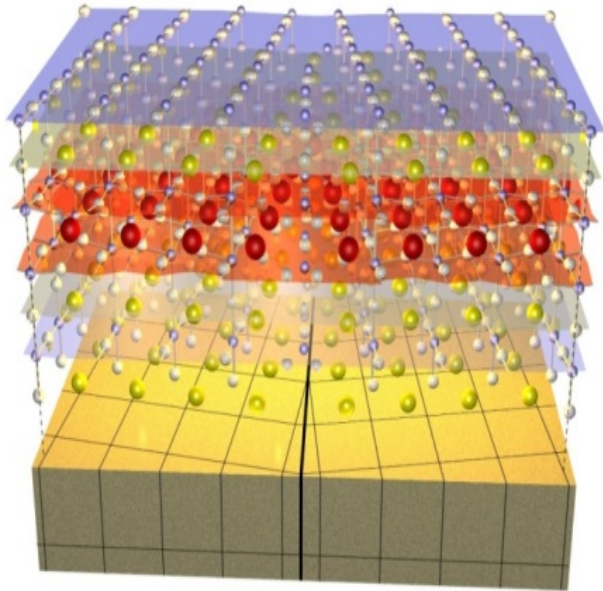
Model Hamiltonian

$$H = \dots$$

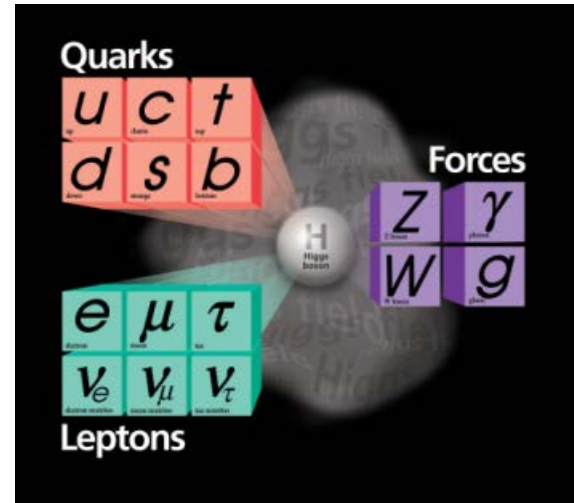


QUANTUM SIMULATORS

APPLICATIONS



Material Science



HEP?

COLD ATOMS IN OPTICAL LATTICES

COLD ATOMS

OPTICAL LATTICES



- Laser standing waves: dipole-trapping

$$H = \int \Psi_{\sigma}^{\dagger} \left(-\nabla^2 + V(\mathbf{r}) \right) \Psi_{\sigma} + u_{\sigma_i} \int \Psi_{\sigma_1}^{\dagger} \Psi_{\sigma_2}^{\dagger} \Psi_{\sigma_3} \Psi_{\sigma_4}$$

Lattice theory: Bose/Fermi-Hubbard model

$$H = -t \sum_n \left(a_n^{\dagger} a_{n+1} + h.c \right) + U \sum_n a_n^{\dagger 2} a_n^2$$

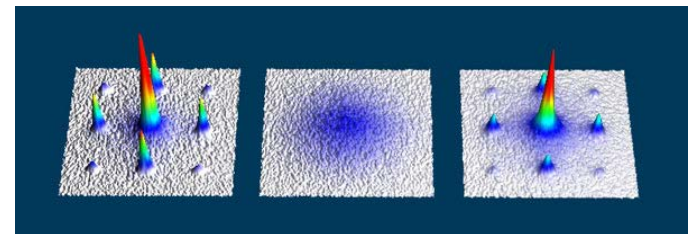
articles

Quantum phase transition from a superfluid to a Mott insulator in a gas of ultracold atoms

Markus Greiner*, Olaf Mandel*, Tilman Esslinger†, Theodor W. Hänsch* & Immanuel Bloch*

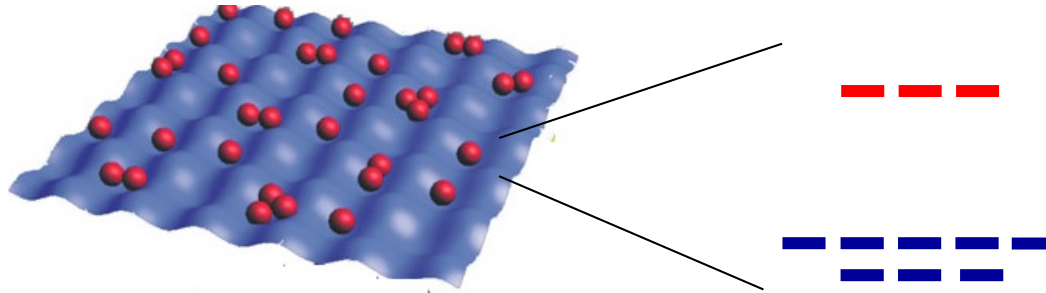
* Sektion Physik, Ludwig-Maximilians-Universität, Schellingstrasse 4/III, D-80799 Munich, Germany, and Max-Planck-Institut für Quantenoptik, D-85748 Garching, Germany

† Quantenelektronik, ETH Zürich, 8093 Zurich, Switzerland



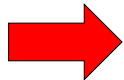
COLD ATOMS

QUANTUM SIMULATION



▪ Bosons/Fermions:
$$H = - \sum_{\substack{\langle n,m \rangle \\ \sigma, \sigma'}} (t_{\sigma, \sigma'} a_{n, \sigma}^\dagger a_{m, \sigma'} + h.c.) + \sum_{\substack{n \\ \sigma, \sigma'}} U_{\sigma, \sigma'} a_{n, \sigma}^\dagger a_{n, \sigma'}^\dagger a_{n, \sigma} a_{n, \sigma}$$

▪ Spins:
$$H = - \sum_{\substack{\langle n,m \rangle \\ \sigma, \sigma'}} (J_x S_n^x S_m^x + J_y S_n^y S_m^y + J_z S_n^z S_m^z) + \sum_{\substack{n \\ \sigma, \sigma'}} B_n S_n^z$$

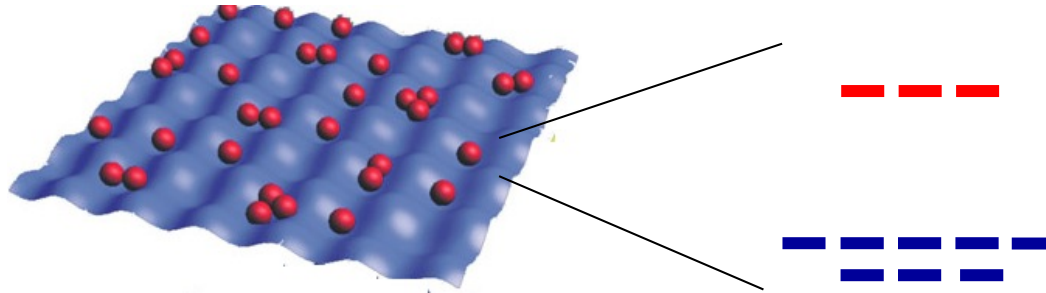


CONDENSED MATTER PHYSICS

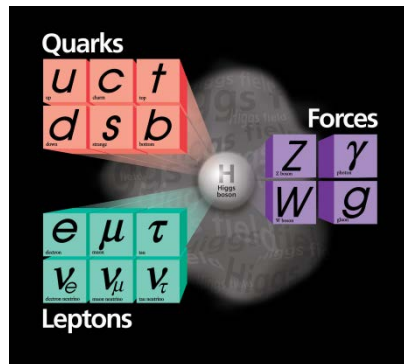


COLD ATOMS

QUANTUM SIMULATION



HIGH ENERGY PHYSICS?



QUANTUM SIMULATIONS OF HEP MODELS

QUANTUM SIMULATION HEP MODELS

INGREDIENTS



$$S = \int \bar{\Psi}(i\gamma^\mu \partial_\mu - m)\Psi - Q \int A_\mu \bar{\Psi}\gamma^\mu\Psi - \frac{1}{4} \int F_{\mu\nu}F^{\mu\nu} + \dots$$

- Matter + Gauge Fields
- Relativistic theory
- Gauge invariant

▪ Hamiltonian formulation: $i\partial_t |\Psi\rangle = H |\Psi\rangle$

- Gauss law

$$G(x) |\Psi\rangle = 0$$

$$[H, G(x)] = 0$$



QUANTUM SIMULATION HEP MODELS

INGREDIENTS



Lattice



J. Kogut and L. Susskind, Phys. Rev. D **11**, 395 (1975).
J. B. Kogut, Rev. Mod. Phys. **51**, 659 (1979).
J. B. Kogut, Rev. Mod. Phys. **55**, 775 (1983).

$$H = \int \Psi_{\sigma}^{\dagger} (-\nabla^2 + V(r)) \Psi_{\sigma} + u \int \Phi_{\mu}^{\dagger} \Phi_{\sigma} \Psi_{\sigma}^{\dagger} \Psi_{\sigma} + v \int \Phi_{\sigma}^{\dagger} \Phi_{\sigma}^{\dagger} \Phi_{\sigma} \Phi_{\sigma} + \dots$$

Lattice

Fermion-gauge field
coupling

Gauge field
dynamics

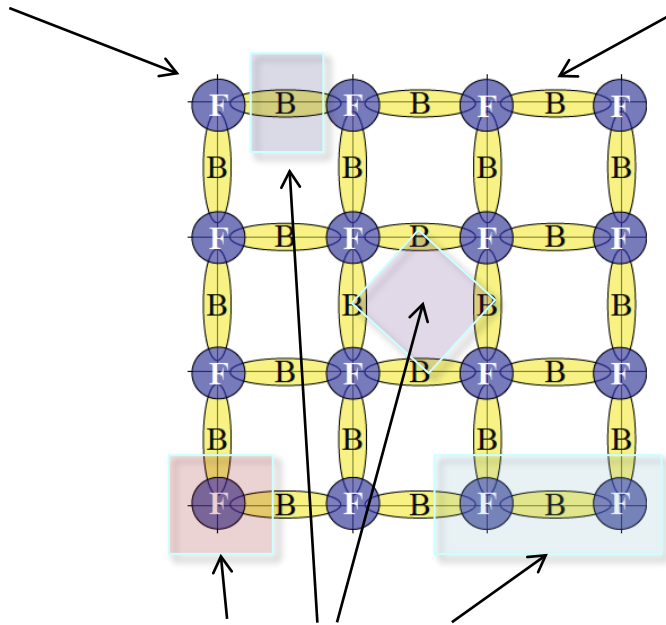
HEP LATTICE MODELS

HAMILTONIAN FORMULATION



Matter (Fermions): can move

Gauge fields (Bosons): Static



- Hamiltonian: $H = H_M + H_{KS} + H_{\text{int}}$

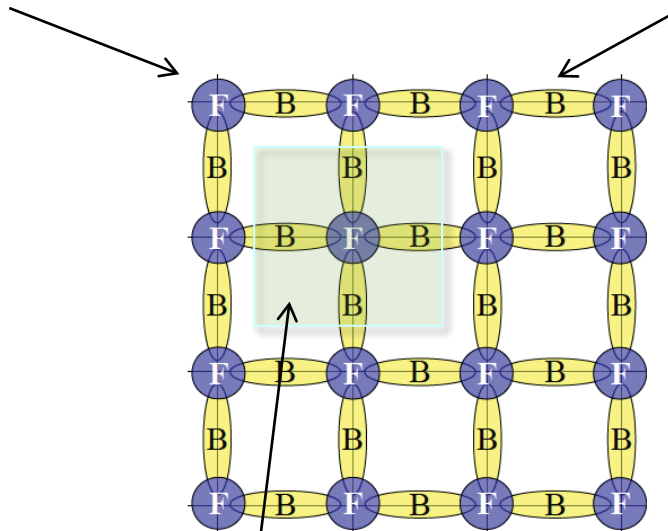
HEP LATTICE MODELS

HAMILTONIAN FORMULATION



Matter (Fermions): can move

Gauge fields (Bosons): Static



- **Hamiltonian:** $H = H_M + H_{KS} + H_{\text{int}}$
- **Gauge invariance:** Gauge group: $U(1)$, Z_N , $SU(N)$, etc

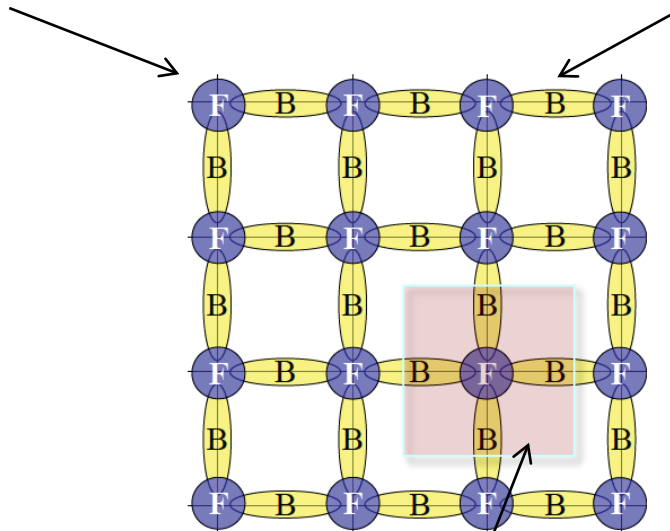
HEP LATTICE MODELS

HAMILTONIAN FORMULATION



Matter (Fermions): can move

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- **Hamiltonian:** $H = H_M + H_{KS} + H_{\text{int}}$
- **Gauge invariance:** Gauge group: $U(1)$, Z_N , $SU(N)$, etc
- **Gauss law:** $G_{\text{plaquette}} | \text{phys} \rangle = 0$

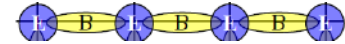
HEP LATTICE MODELS

HAMILTONIAN FORMULATION



Example: compact-QED in 1D

- Hamiltonians:



$$H_M = \sum_{\mathbf{n}} M_{\mathbf{n}} \psi_{\mathbf{n}}^{\dagger} \psi_{\mathbf{n}}$$

$$H_{int} = \epsilon \sum_{\mathbf{n}, \mathbf{k}} \left(\psi_{\mathbf{n}}^{\dagger} e^{i\phi_{\mathbf{n}, \mathbf{k}}} \psi_{\mathbf{n}+\mathbf{k}} + \psi_{\mathbf{n}+\mathbf{k}}^{\dagger} e^{-i\phi_{\mathbf{n}, \mathbf{k}}} \psi_{\mathbf{n}} \right)$$

$$H_{KS} = H_E = \frac{g^2}{2} \sum_{\mathbf{n}, \mathbf{k}} E_{\mathbf{n}}^2$$

$$[E_{\mathbf{n}, \mathbf{k}}, \phi_{\mathbf{m}, \mathbf{l}}] = -i\delta_{\mathbf{n}\mathbf{m}}\delta_{\mathbf{k}\mathbf{l}} \quad (\text{ie, compact})$$

- Gauss law: $G_n |phys\rangle = 0$

- Gauge invariance: $e^{-i\theta G_n} H e^{i\theta G_n} = H$

$$G_n = E_{n+1} - E_n - \psi_n^{\dagger} \psi_n$$

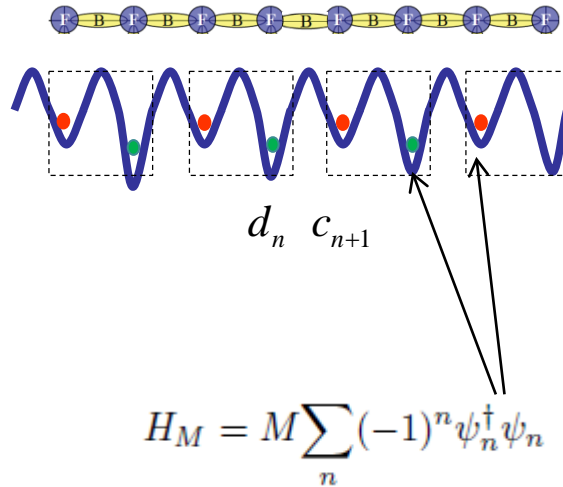
SCHWINGER MODEL

QUANTUM SIMULATION

SCHWINGER MODEL 1+1



Fermions:



internal states



$$\{c_n, c_n^\dagger\} = \{d_n, d_n^\dagger\} = 1$$

$$H = \int \Psi_\sigma^\dagger (-\nabla^2 + V(r)) \Psi_\sigma + u \int \Phi_\mu^\dagger \Phi_\sigma \Psi_\sigma^\dagger \Psi_\sigma + v \int \Phi_\sigma^\dagger \Phi_\sigma^\dagger \Phi_\sigma \Phi_\sigma + \dots$$

- Even sites: hole = particle
- Odd sites: fermion = antiparticle

Staggered Fermions:

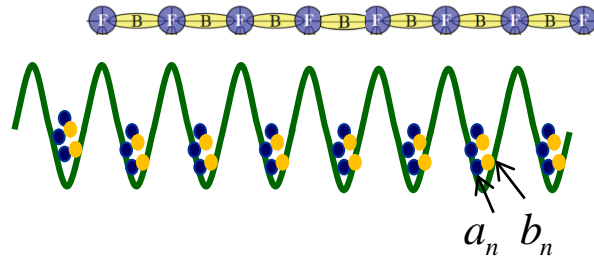
- L. Susskind, Phys. Rev. D **16**, 3031 (1977).
- G. 't Hooft, Nucl. Phys. B **75**, 461 (1974)

QUANTUM SIMULATION

SCHWINGER MODEL 1+1



□ Bosons:



internal states

$$[a_n, a_n^\dagger] = [b_n, b_n^\dagger] = 1$$

● Schwinger rep:

$$\begin{aligned}
 L_+ &= a^\dagger b \\
 L_z &= \frac{1}{2} (N_a - N_b) \\
 \ell &= \frac{1}{2} (N_a + N_b)
 \end{aligned}
 \quad \xrightarrow{\ell \gg 1} \quad
 \begin{aligned}
 L_+ &\approx \ell e^{i\phi} \\
 L_z &\approx i\partial_\phi
 \end{aligned}$$

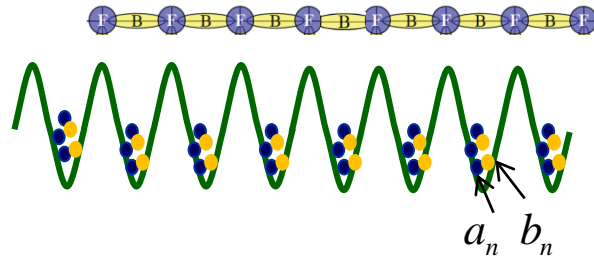
- If ℓ is small (eg 2 atoms), we obtain a truncated version
- One can also use a single atom with few internal levels (Z_M is the gauge group)

QUANTUM SIMULATION

SCHWINGER MODEL 1+1



□ Bosons:



internal states



$$[a_n, a_n^\dagger] = [b_n, b_n^\dagger] = 1$$

Schwinger rep:

$$L_+ = a^\dagger b$$

$$L_z = \frac{1}{2} (N_a - N_b)$$

$$\ell = \frac{1}{2} (N_a + N_b)$$

$$H_E = \frac{g^2}{2} \sum_n L_{z,n}^2$$

$$= \frac{g^2}{8} \sum_n (N_{a,n}^2 + N_{b,n}^2 - 2N_{a,n}N_{b,n})$$

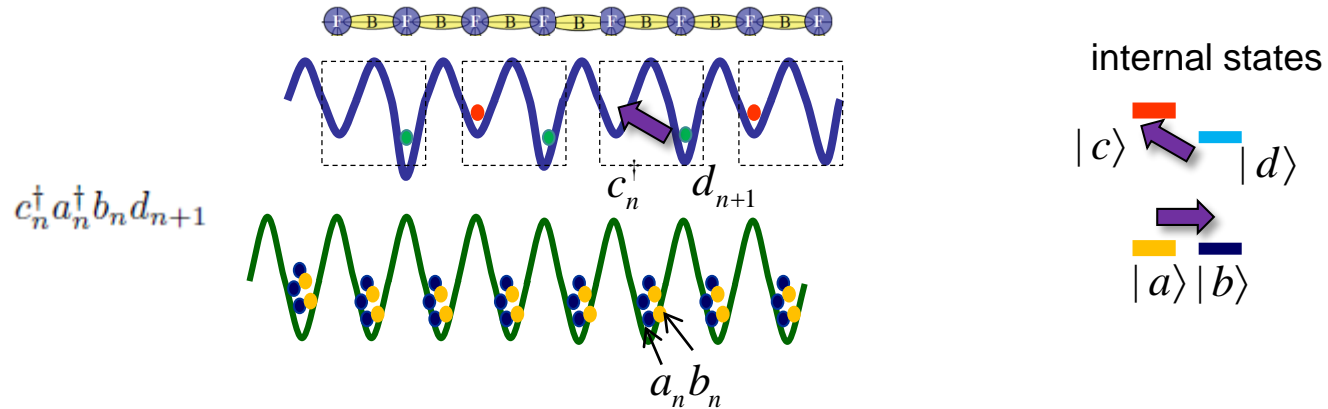
$$H = \int \Psi_\sigma^\dagger (-\nabla^2 + V(r)) \Psi_\sigma + u \int \Phi_\mu^\dagger \Phi_\sigma \Psi_\sigma^\dagger \Psi_\sigma + v \int \Phi_\sigma^\dagger \Phi_\sigma^\dagger \Phi_\sigma \Phi_\sigma + \dots$$

QUANTUM SIMULATION

SCHWINGER MODEL 1+1



Interactions:



$$H = \int \Psi_\sigma^\dagger (-\nabla^2 + V(r)) \Psi_\sigma + u \int \Phi_\mu^\dagger \Phi_{\sigma'} \Psi_\sigma^\dagger \Psi_{\sigma'} + v \int \Phi_\sigma^\dagger \Phi_{\sigma'} \Phi_{\sigma'} \Phi_\sigma + \dots$$

conserves angular momentum locally

→ Gauge invariance

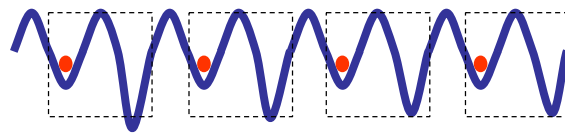
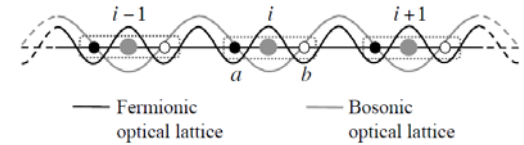
$$H_{int} \frac{1}{\sqrt{\ell(\ell+1)}} \psi_n^\dagger a_n^\dagger b_n \psi_{n+1} \approx \psi_n^\dagger e^{i\phi_n} \psi_{n+1}$$

QUANTUM SIMULATION

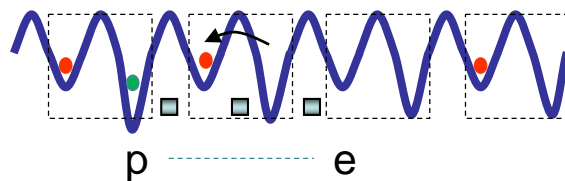
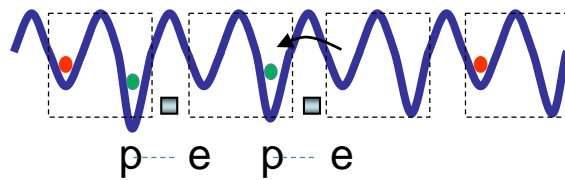
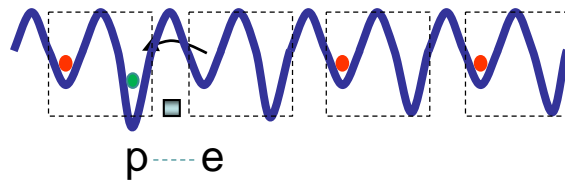
SCHWINGER MODEL 1+1



Physical processes:

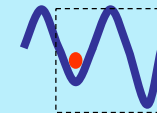


non-interacting
vacuum

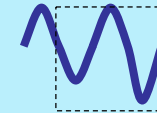


TABLE

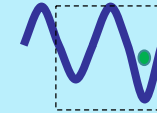
$$|0\rangle_e |0\rangle_p$$



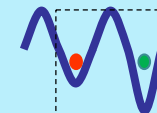
$$|1\rangle_e |0\rangle_p$$



$$|1\rangle_e |1\rangle_p$$



$$|0\rangle_e |1\rangle_p$$

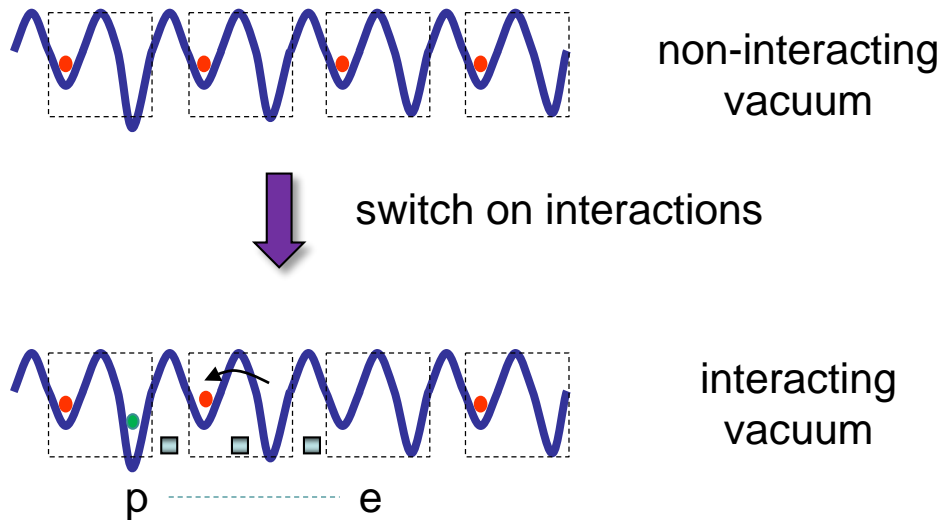
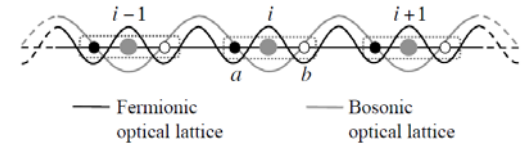


QUANTUM SIMULATION

SCHWINGER MODEL 1+1



Preparation:



- Confinement
- Excitations: vector + scalar
- Time-dependent phenomena
- First experiments: few bosonic atoms

TABLE

$ 0\rangle_e 0\rangle_p$	
$ 1\rangle_e 0\rangle_p$	
$ 1\rangle_e 1\rangle_p$	
$ 0\rangle_e 1\rangle_p$	

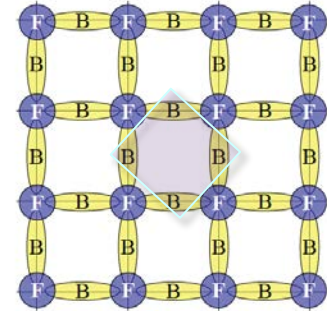
QUANTUM SIMULATION

HIGHER DIMENSIONS, NON-ABELIAN



- Plaque interaction:

$$H_B = -\frac{2\epsilon^4}{\lambda^3} \sum_{\mathbf{n}} \left(U_{\mathbf{n},1} U_{\mathbf{n}+\hat{1},2} U_{\mathbf{n}+\hat{2},1}^\dagger U_{\mathbf{n},2}^\dagger + h.c. \right) = -\frac{4\epsilon^4}{\lambda^3} \sum_{\mathbf{n}} \cos \left(\phi_{\mathbf{n},1} + \phi_{\mathbf{n}+\hat{1},2} - \phi_{\mathbf{n}+\hat{2},1} - \phi_{\mathbf{n},2} \right)$$

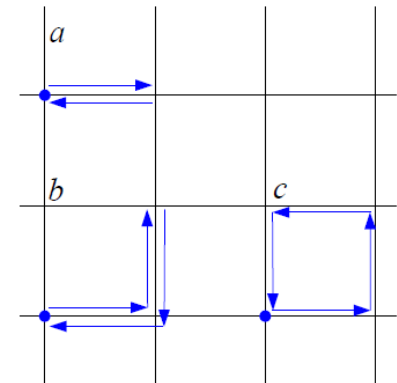


- Non-abelian gauge theories:

$$H_E = \frac{g^2}{2} \sum_{\mathbf{n},k,a} (E_{\mathbf{n},k})_a (E_{\mathbf{n},k})_a$$

$$H_B = -\frac{1}{g^2} \sum_{\text{plaquettes}} \left(\text{Tr} \left(U_1 U_2 U_3^\dagger U_4^\dagger \right) + h.c. \right)$$

$$H_{int} = \epsilon \sum_{\mathbf{n},k} \left(\psi_{\mathbf{n}}^\dagger U_{\mathbf{n},k}^r \psi_{\mathbf{n}+\hat{k}} + h.c. \right)$$



Link



L R

{a₁, a₂} {b₁, b₂}

bosonic modes

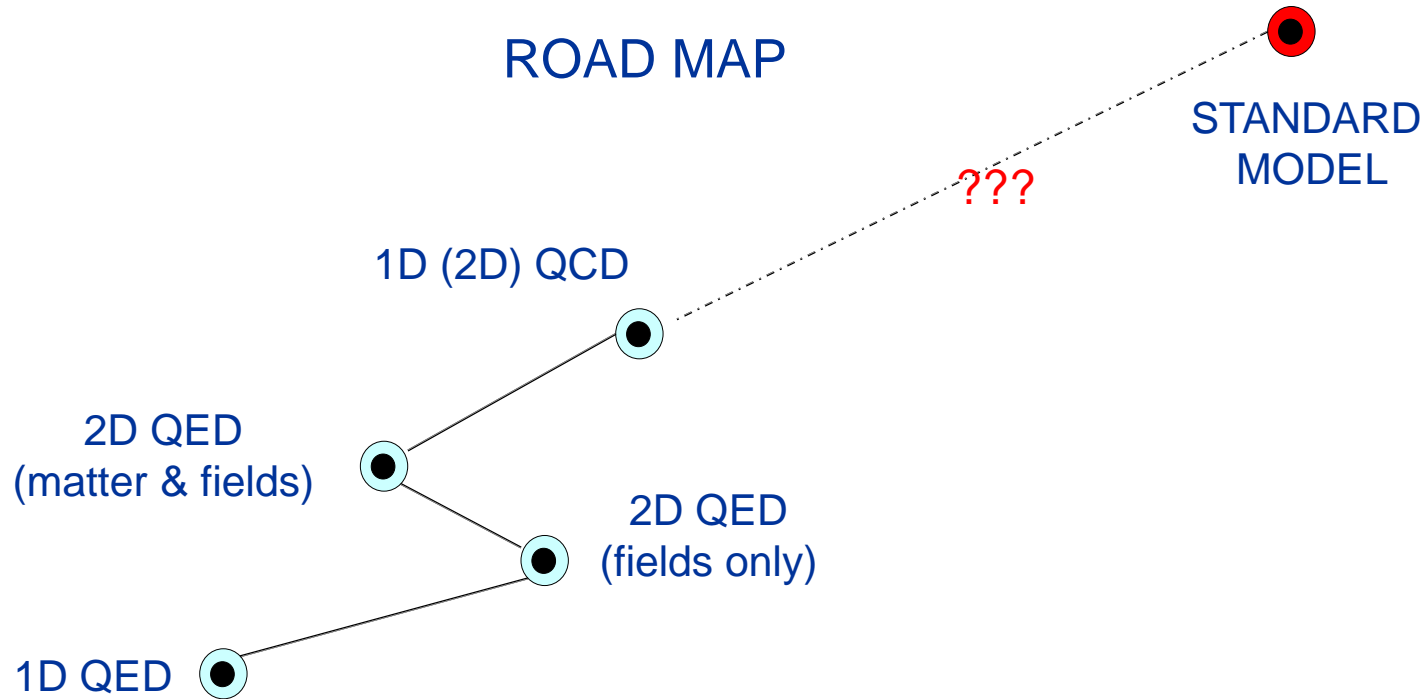


QUANTUM SIMULATION HIGH ENERGY MODELS





QUANTUM SIMULATION HIGH ENERGY MODELS



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Zohar, IC, Reznik, PRL **107**, 275301 (2011)
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Zohar, IC, Reznik, PRL **110**, 125304 (2013)
[Zohar, IC, Reznik, PRA **88**, 023617 \(2013\)](#)
Zohar, Farace, Reznik, JIC, in preparation

See also:

Kapit, Mueller, PRA **83**, 033625 (2011)
Banerjee, ..., Wiese, Zoller, PRL **109**, 175302 (2013)
Banerjee, ..., Wiese, Zoller, PRL **110**, 125303 (2013)
Gauge fields: Lewenstein et al