

Soft lepton number violation in multi-Higgs doublet Seesaw models

Elke Aeikens

University Vienna
PhD-advisor Prof. Walter Grimus

28th June 2016

$\int dk$ Π Doktoratskolleg
Particles and Interactions



universität
wien

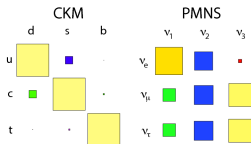
Faculty of Physics

Why we care about neutrinos

experimentally unsolved: anomalies...

theoretical unsolved: (all about mass)

- different mixing matrices than quarks
- normal or inverted mass hierarchy
- hierarchy problem: very light mass
- origin of mass: Dirac, Majorana



properties:

- just weak interacting
- no observed right handed partner

Desperately seeking sterile

The three known types of neutrino might be "balanced out" by a bashful fourth type

| ELECTRON NEUTRINO | MUON NEUTRINO | TAU NEUTRINO | STERILE NEUTRINO |
|------------------------|-------------------------|--------------|------------------|
| ν_e | ν_μ | ν_τ | ν_s |
| MASS | < 1 electronvolt | | > 1 electronvolt |
| FORCES THEY RESPOND TO | Weak force Gravity | | Gravity |
| DIRECTION OF SPIN | All three "left handed" | | "Right handed" |

Why we care about neutrinos

experimentally unsolved: anomalies...

theoretical unsolved: (all about mass)

- different mixing matrices than quarks
- normal or inverted mass hierarchy
- hierarchy problem: very light mass
- origin of mass: Dirac, Majorana

$$\mathcal{L}_M \stackrel{?}{=} \mathcal{L}_{Dirac} + \mathcal{L}_{Majorana}$$





$$\sim \bar{\nu}_R M_D \nu_L + \bar{\nu}_\alpha M_M \nu_\beta + h.c.$$

properties:

- just weak interacting
- no observed right handed partner

Desperately seeking sterile

The three known types of neutrino might be "balanced out" by a bashful fourth type

| ELECTRON NEUTRINO | MUON NEUTRINO | TAU NEUTRINO | STERILE NEUTRINO |
|---|--|---|---|
|  |  |  |  |
| MASS | < 1 electronvolt | | > 1 electronvolt |
| FORCES THEY RESPOND TO | Weak force Gravity | | Gravity |
| DIRECTION OF SPIN | All three "left handed" | | "Right handed" |

introduce:
right handed neutrinos ν_R

Advantage of right handed neutrinos

- **Explain mass hierarchy** in right handed neutrino mass models via the seesaw mechanism. [$m_{\nu_R} \gtrsim \text{TeV}$]
(with additional higgs doublets...)
- **Dark matter candidates** [$\text{keV} \lesssim m_{\nu_R} \lesssim \text{TeV}$]
- **Baryon asymmetry** via Leptogenesis in ν MSM models [$\text{keV} \lesssim m_{\nu_R} \lesssim \text{GeV}$]
- **Detected anomalies** at: LSND, MiniBooNE, gallium detectors: GALLEX, SAGE, reactor experiments... [$m_{\nu_R} \sim \text{eV}$]
(a.o. also IceCube)

tightest constrains from cosmology:

- **Boundaries from BBN**
- **CMB measurement** from PLANCK sets limits on N_ν and also the **Large Scale Structure**.

Advantage of right handed neutrinos

- **Explain mass hierarchy** in right handed neutrino mass models via the seesaw mechanism. [$m_{\nu_R} \gtrsim \text{TeV}$]
(with additional higgs doublets...)
- **Dark matter candidates** [$\text{keV} \lesssim m_{\nu_R} \lesssim \text{TeV}$]
- **Baryon asymmetry** via Leptogenesis in ν MSM models [$\text{keV} \lesssim m_{\nu_R} \lesssim \text{GeV}$]
- **Detected anomalies** at: LSND, MiniBooNE, gallium detectors: GALLEX, SAGE, reactor experiments... [$m_{\nu_R} \sim \text{eV}$]
(a.o. also IceCube)

tightest constrains from cosmology:

- **Boundaries from BBN** previous work
- **CMB measurement** from PLANCK sets limits on N_ν and also the **Large Scale Structure**.

Advantage of right handed neutrinos

- **Explain mass hierarchy** in right handed neutrino mass models via the seesaw mechanism. [$m_{\nu_R} \gtrsim \text{TeV}$]
(with additional higgs doublets...)

actual work

- **Dark matter candidates** [$\text{keV} \lesssim m_{\nu_R} \lesssim \text{TeV}$]
- **Baryon asymmetry** via Leptogenesis in ν MSM models
[$\text{keV} \lesssim m_{\nu_R} \lesssim \text{GeV}$]

- **Detected anomalies** at: LSND, MiniBooNE, gallium detectors: GALLEX, SAGE, reactor experiments... [$m_{\nu_R} \sim \text{eV}$]
(a.o. also IceCube)

tightest constrains from cosmology:

- **Boundaries from BBN**
- **CMB measurement** from PLANCK sets limits on N_ν and also the **Large Scale Structure**.

previous work

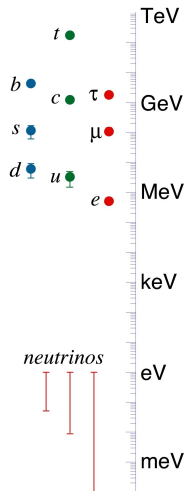
First ingredient for a good model

Hierarchy problem:

Neutrino mass is small $m_\nu < 0.1$ eV (exp. limits)

Masses are normally $m_e \simeq 0.5$ MeV to $m_t \simeq 173$ GeV

⇒ *small Yukawa masses seem to be unnatural*



First ingredient for a good model: ν_R

Hierarchy problem:

Neutrino mass is small $m_\nu < 0.1$ eV (exp. limits)

Masses are normally $m_e \approx 0.5$ MeV to $m_t \approx 173$ GeV

\Rightarrow *small Yukawa masses seem to be unnatural*

Solution: **Seesaw mechanism**

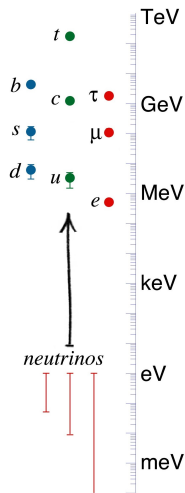
Majorana neutrinos with right handed partners ν_R

flavour scale: $m_R \gtrsim \text{TeV}$, $m_D \sim m_e$

$$M_{maj} = \begin{pmatrix} 0 & M_D^T \\ M_D & M_R \end{pmatrix}$$

diagonalisation:

mass scale: $m_\nu = -m_D^2/m_R \rightarrow m_\nu$ small



First ingredient for a good model: ν_R

Hierarchy problem:

Neutrino mass is small $m_\nu < 0.1$ eV (exp. limits)

Masses are normally $m_e \approx 0.5$ MeV to $m_t \approx 173$ GeV

\Rightarrow *small Yukawa masses seem to be unnatural*

Solution: **Seesaw mechanism**

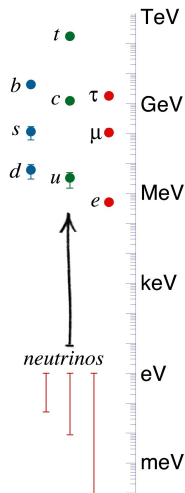
Majorana neutrinos with right handed partners ν_R

flavour scale: $m_R \gtrsim \text{TeV}$, $m_D \sim m_e$

$$M_{maj} = \begin{pmatrix} 0 & M_D^T \\ M_D & M_R \end{pmatrix}$$

diagonalisation:

mass scale: $m_\nu = -m_D^2/m_R \rightarrow m_\nu$ small

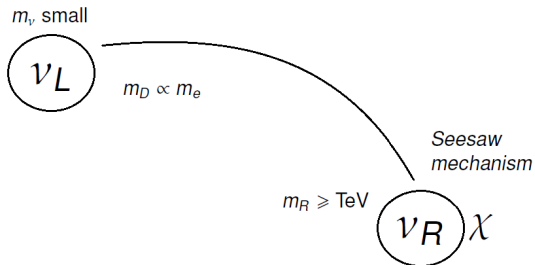


The model

m_ν small



The model



Second ingredient for a good model

Problem:

Yukawa couplings $Y_\nu \simeq \frac{m_\nu}{v}$ small, when m_ν small ($v \sim 246$ GeV, VEV)

Other gauge couplings large: e.g. Positron $e = \sqrt{4\pi\alpha} = 0.303$

\Rightarrow *small Y_ν seem to be unnatural*

Second ingredient for a good model: Φ_k

Problem:

Yukawa couplings $Y_\nu \simeq \frac{m_\nu}{v}$ small, when m_ν small ($v \sim 246$ GeV, VEV)

Other gauge couplings large: e.g. Positron $e = \sqrt{4\pi\alpha} = 0.303$

\Rightarrow small Y_ν seem to be unnatural

Solution: **multi-Higgs doublet model (mHdm)**

include n_H Higgs doublets

$$\Phi_k = \begin{pmatrix} \Phi_k^+ \\ \Phi_k^0 \end{pmatrix}, \quad \langle 0 | \Phi_k^0 | 0 \rangle = \frac{v_k}{\sqrt{2}}, \quad \sum_k |v_k|^2 \sim (246 \text{ GeV})^2 \rightarrow \text{small } v_k \text{ so}$$

that $Y_\nu \sim \mathcal{O}(e)$.

lepton Yukawa couplings

$$\mathcal{L}_Y = - \sum_{k=1}^{n_H} \sum_{l,l' = e, \mu, \tau} [(\phi_k^-, \phi_k^{0*}) \mathbf{Y}_{lkl'} \bar{l}_R + (\phi_k^0, -\phi_k^+) \mathbf{Y}_{\nu kll'} \bar{\nu}_{lR}] \begin{pmatrix} \nu_{l'L} \\ l'_{L} \end{pmatrix} + \text{H.c.}$$

Second ingredient for a good model: Φ_k

Problem:

Yukawa couplings $Y_\nu \simeq \frac{m_\nu}{v}$ small, when m_ν small ($v \sim 246$ GeV, VEV)

Other gauge couplings large: e.g. Positron $e = \sqrt{4\pi\alpha} = 0.303$

\Rightarrow small Y_ν seem to be unnatural

Solution: **multi-Higgs doublet model (mHdm)**

include n_H Higgs doublets

$$\Phi_k = \begin{pmatrix} \Phi_k^+ \\ \Phi_k^0 \end{pmatrix}, \quad \langle 0 | \Phi_k^0 | 0 \rangle = \frac{v_k}{\sqrt{2}}, \quad \sum_k |v_k|^2 \sim (246 \text{ GeV})^2 \rightarrow \text{small } v_k \text{ so}$$

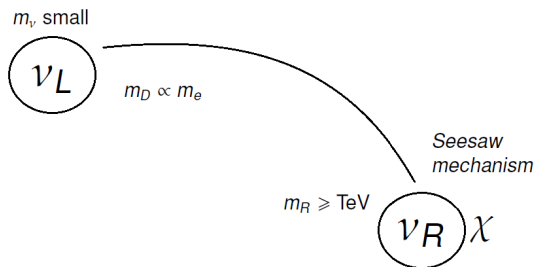
that $Y_\nu \sim \mathcal{O}(e)$.

lepton Yukawa

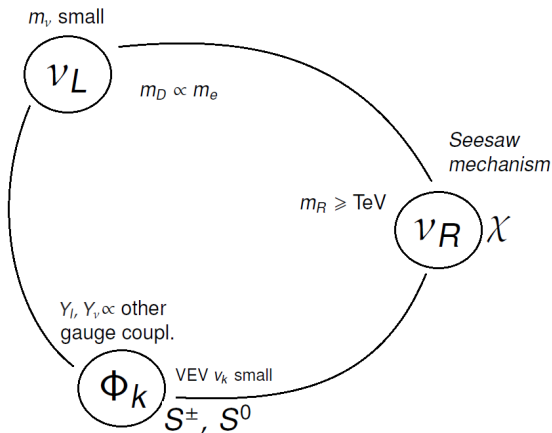
$$\mathcal{L}_Y = - \sum_{k=1}^{n_H} \sum_{l, l' = e, \mu, \tau} [(\phi_{k^-}, \phi_{k^0}^*) Y_{lkl'} \bar{l}_R + (\phi_{k^0}^+, -\phi_{k^+}) Y_{\nu kll'} \bar{\nu}_{lR}] \begin{pmatrix} \nu_{l'L} \\ l'_L \end{pmatrix} + \text{H.c.}$$

interesting effect:
observable processes!

The model



The model



experimentally testable processes

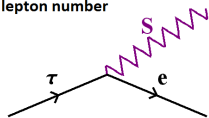
Third ingredient for a good model

mHdm Problem:

Flavour-changing neutral scalar interactions (FCNIs) at tree level appear.

⇒ *strong experimental bounds on FCNIs*

FCNI from hart broken
lepton number

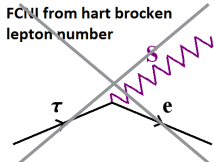


Third ingredient for a good model: L_α

mHdm Problem:

Flavour-changing neutral scalar interactions (FCNIs) at tree level appear.

⇒ *strong experimental bounds on FCNIs*



Solution: **soft lepton number L_α violation** $\{\alpha = e, \mu, \tau\}$

L_α conservation:

in Yukawa interactions

⇒ diag. Y_l, Y_ν

$$M_l \sim \sum_k v_k^* Y_{lk}, \quad M_D \sim \sum_k v_k Y_{\nu k}$$

L_α explicit soft breaking:

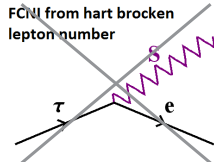
in Majorana term ⇒ **non-diag. M_R**

Third ingredient for a good model: L_α

mHdm Problem:

Flavour-changing neutral scalar interactions (FCNIs) at tree level appear.

⇒ *strong experimental bounds on FCNIs*



Solution: **soft lepton number L_α violation** $\{\alpha = e, \mu, \tau\}$

L_α conservation:

in Yukawa interactions

⇒ $\text{diag. } Y_l, Y_\nu \Rightarrow \text{diag. } \mathbf{M}_l, \mathbf{M}_D = \text{diag}(m_e, m_\mu, m_\tau)$

L_α explicit soft breaking:

in Majorana term ⇒ **non-diag. \mathbf{M}_R**

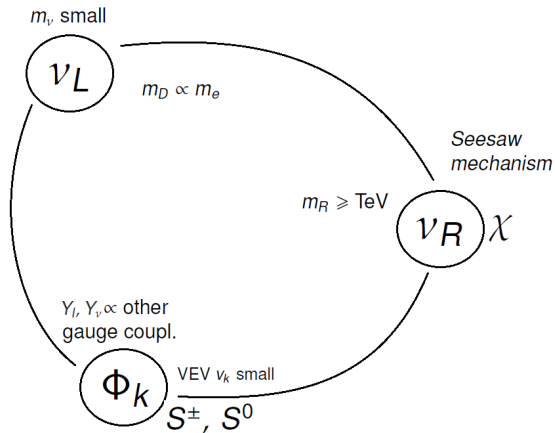
additional advantages:

- explain atm. & sol. maximal mixing [Grimus, 01]
- ampl. of FC processes are finite at one-loop
- ampl. are stable under radiative corrections

PMNS

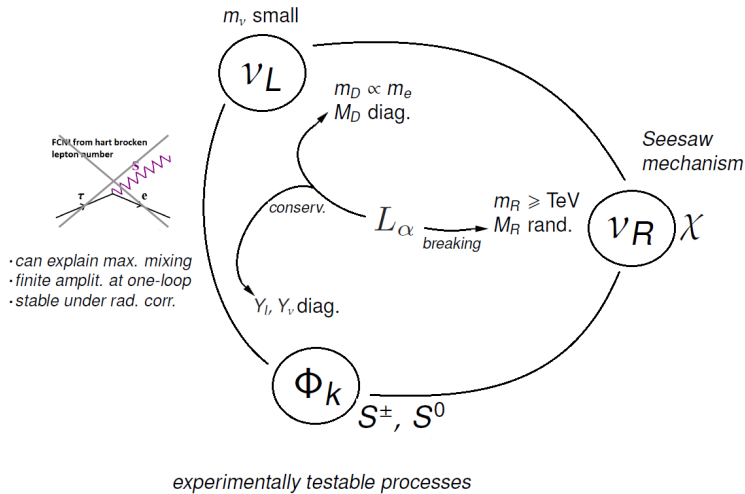
| | ν_1 | ν_2 | ν_3 |
|------------|---------|---------|---------|
| ν_e | Yellow | Blue | Red |
| ν_μ | Green | Blue | Yellow |
| ν_τ | Green | Blue | Yellow |

The model



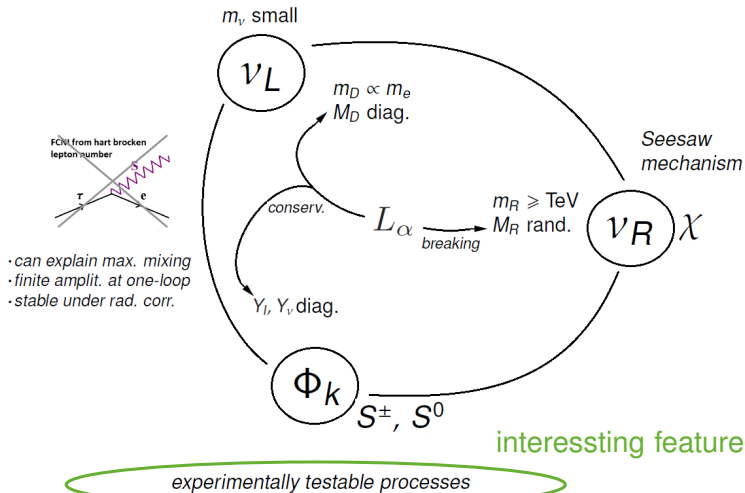
experimentally testable processes

The model



The model

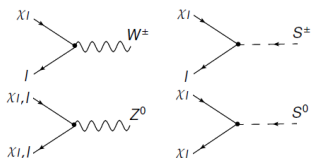
Nice model but: **Can it be tested? Does it bring Limits?**



The model

Evtl. experimentally testable processes:

Additional fermion interactions:



$$\mathcal{A}(\tau^- \rightarrow \mu^- \gamma) \propto 1/m_R^2,$$

$$\mathcal{A}(Z \rightarrow \tau^+ \mu^-) \propto 1/m_R^2,$$

$$\mathcal{A}(\tau^- \rightarrow \mu^- e^+ e^-) \propto \begin{cases} 1/m_R^2 & n_H = 1 \\ \text{const.} & n_H > 1 \end{cases}$$

Processes including the sub-process $l^- \rightarrow l^- S^{0*}$, ($S^{0*} \rightarrow e^+ e^-$) have ($n_H \geq 2$) non- m_R -suppressed contributions from graphs with charged-scalar exchange S^\pm (plot) in their Amplitudes \mathcal{A} , [Grimus, Lavoura, 02].

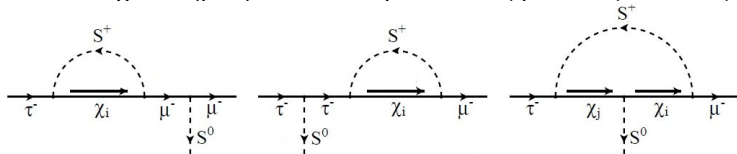
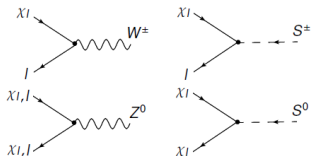


Figure: The tree diagrams for $\tau^- \rightarrow \mu^- S^{0*}$

The model

Evtl. experimentally testable processes:

Additional fermion interactions:



$$\mathcal{A}(\tau^- \rightarrow \mu^- \gamma) \propto 1/m_R^2,$$

$$\mathcal{A}(Z \rightarrow \tau^+ \mu^-) \propto 1/m_R^2,$$

$$\mathcal{A}(\tau^- \rightarrow \mu^- e^+ e^-) \propto \begin{cases} 1/m_R^2 & n_H = 1 \\ \text{const.} & n_H > 1 \end{cases}$$

Processes including the sub-process $l^- \rightarrow l^- S^{0*}$, ($S^{0*} \rightarrow e^+ e^-$) have ($n_H \geq 2$) non- m_R -suppressed contributions from graphs with charged-scalar exchange S^\pm (plot) in their Amplitudes \mathcal{A} , [Grimus, Lavoura, 02].

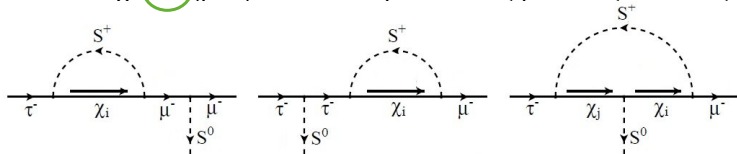


Figure: The tree diagrams for $\tau^- \rightarrow \mu^- S^{0*}$

Expected outcome and goals

Nice model but: **Can it be tested? Does it bring Limits?**

Expectations:

- Finding upper bounds on flavour diagonal Yukawa couplings (Y_l, Y_ν) at one loop (with $m_R \rightarrow \infty$)
- Finding lower benchmarks on seesaw scale m_R
 \Rightarrow with comparing them to the experimental upper bounds on branching ratios.
- Pointing out **experimental signatures**.

Thank you!