

Gravitational wave data analysis



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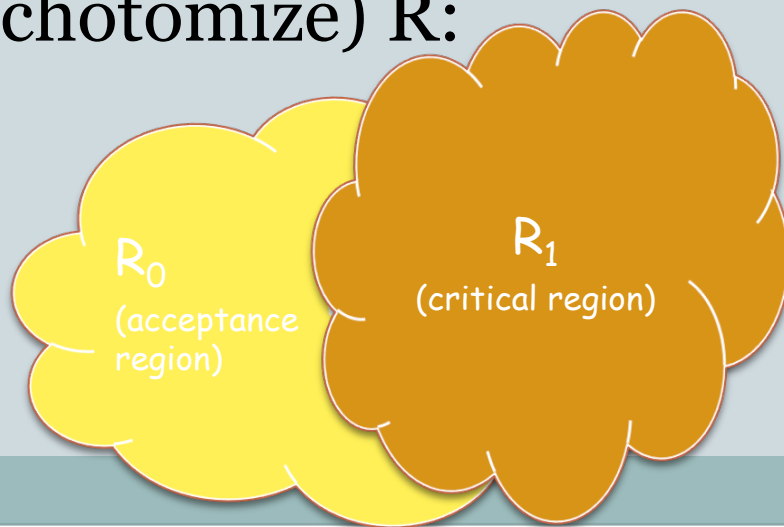
This presentation



- Aim: make it possible for you to read and understand the observational papers of LIGO and Virgo.
 - Very brief basics of signal detection
 - Searches for compact binary inspiral signals

Hypothesis testing

- Our data $\{y_i\} \in \mathcal{R}$
- Consider a detection problem: H_0 signal absent, H_1 signal present
- Question: has hypothesis H_0 or H_1 produced our data ?
- Deciding means finding a way to partition (dichotomize) \mathcal{R} :

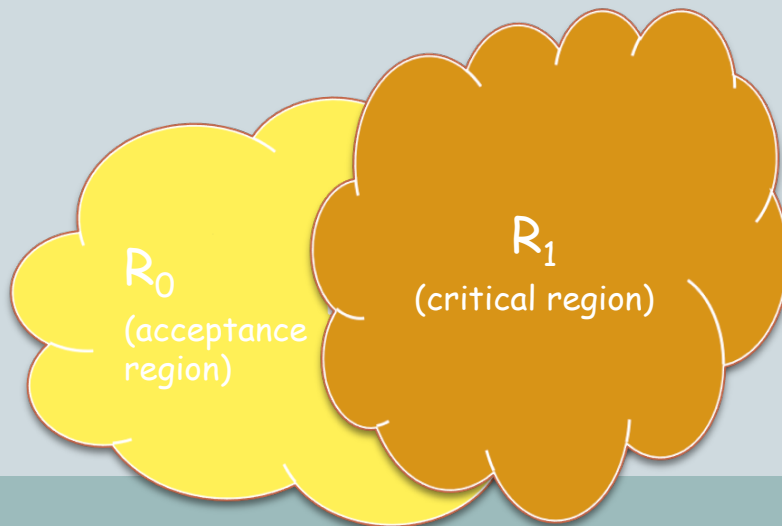


if $\{y_i\} \in R_0 \rightarrow D_0$
if $\{y_i\} \in R_1 \rightarrow D_1$

Hypothesis testing- types of errors



- Type I error: decide D_1 , when H_0 holds
 - False alarm probability, $P_{fa}=P(D_1|H_0)$, size of the test
- Type II errors:decide D_0 , when H_1 holds
 - False dismissal probability, $P_{fd}=P(D_0|H_1)$, $1-P_{fd}$ is the power of the test



if $\{y_i\} \in R_0 \rightarrow D_0$
if $\{y_i\} \in R_1 \rightarrow D_1$

Neymann-Pearson criterium



- The decision should be such that at fixed P_{fa} the P_{fd} is the smallest.
- It can be demonstrated that the corresponding partition is any level surface of a function of the data called the likelihood:

$$\Lambda(\mathbf{y}) = \frac{p_1(\mathbf{y})}{p_0(\mathbf{y})}$$

prob given H_1
prob given H_0

Neymann-Pearson criterium



- The specific level surface that one takes depends on convenience and defines the detection statistic.
- The partition is a threshold on the detection statistic that determines the P_{fa} and the P_{fd} .

$$R_1 \supset y \mid \Lambda(y) > \Lambda^*$$

$$P(\Lambda(y) > \Lambda^* \mid H_0) = P_{fa}$$

A very simple example

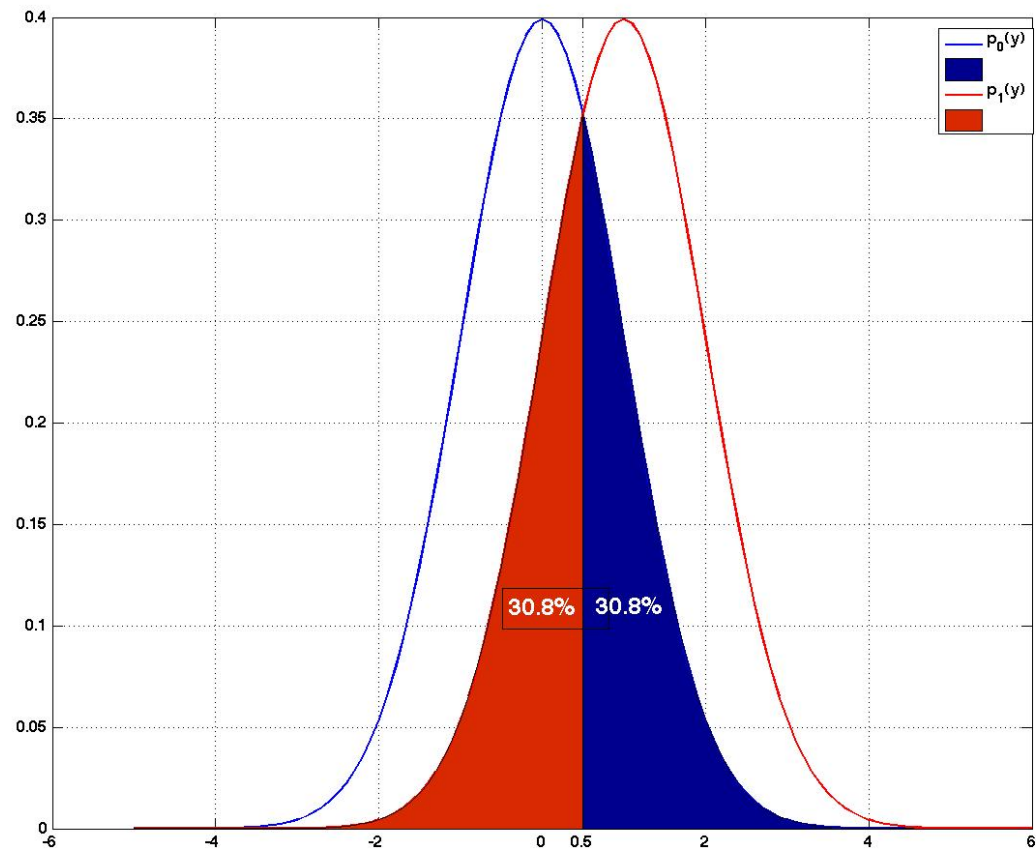


- Consider a single measurement
- $y = n + s$
 - n is Gaussian noise, zero mean and unit variance
 - $s = 1$ is a constant, our signal

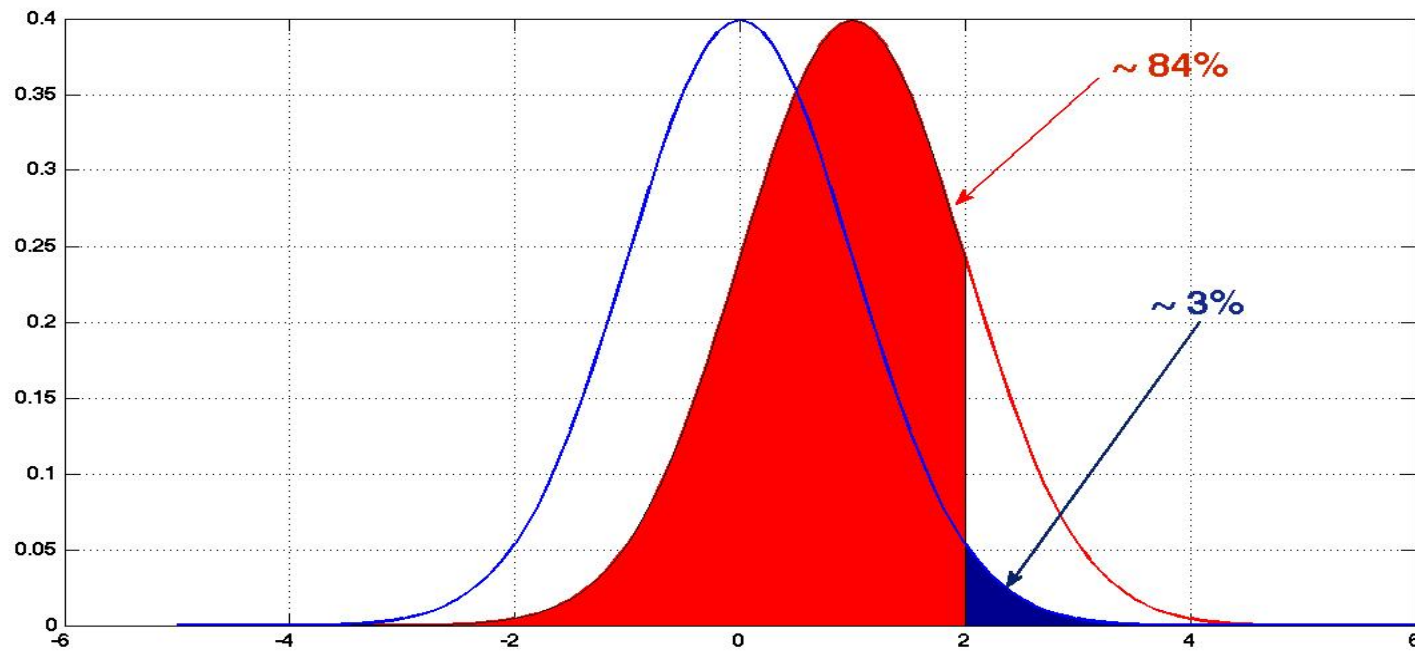
$$p_0 = \frac{1}{\sqrt{2\pi}} e^{-y^2/2} \quad p_1 = \frac{1}{\sqrt{2\pi}} e^{-(y-1)^2/2} \quad \text{and} \quad \Lambda(y) = \frac{p_1(y)}{p_0(y)}$$

- Neymann-Pearson Criterium:
 - $\Lambda(y) = e^{(y-1/2)}$ → Rule is: threshold on y

If we set $y^* = 0.5$

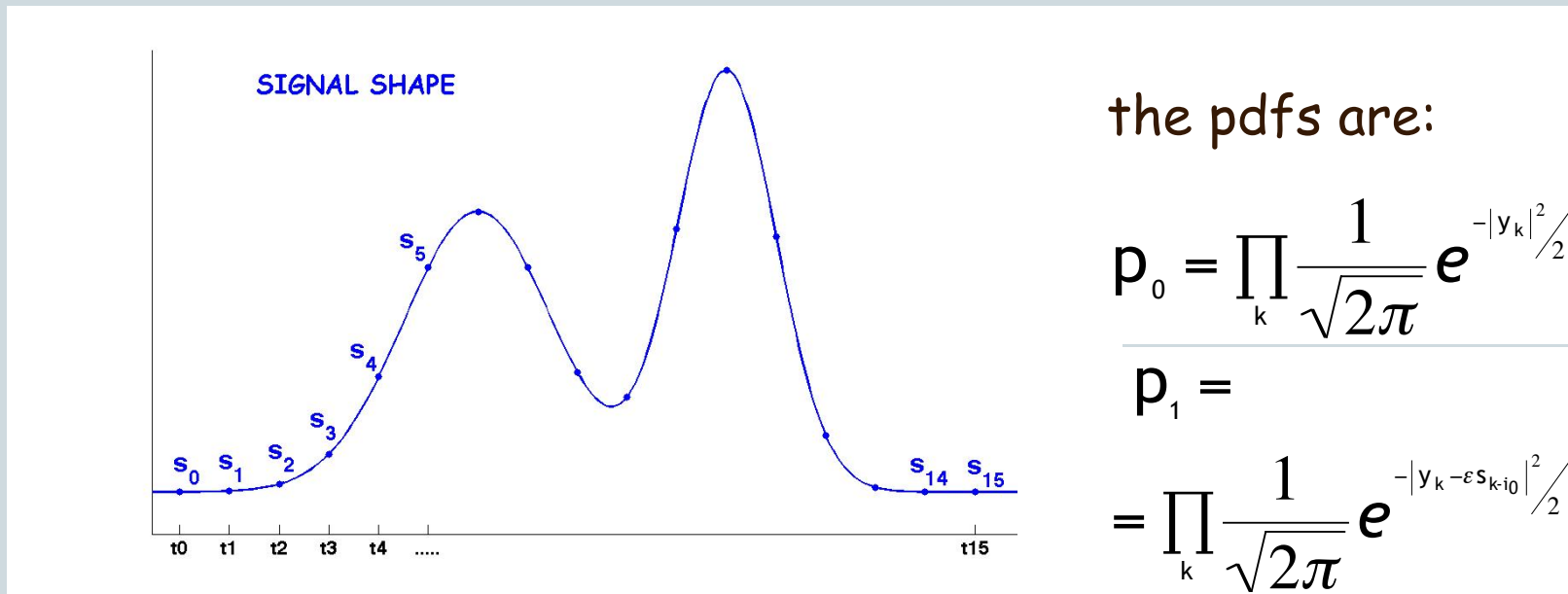


If we set $y^* = 2$



Another example, a set of measurements

- Consider some measurements
- $y_i(i_0) = n_i + \epsilon s_{i-i_0}$
 - n is Gaussian noise, zero mean and unit variance
 - s_{i-i_0} is a signal of known shape, arriving at time i_0



a convenient level surface of the likelihood is $\rho(\{y\}, i_0) = \sum_k \epsilon s_{k-i_0} y_k$

Likelihood and matched filtering

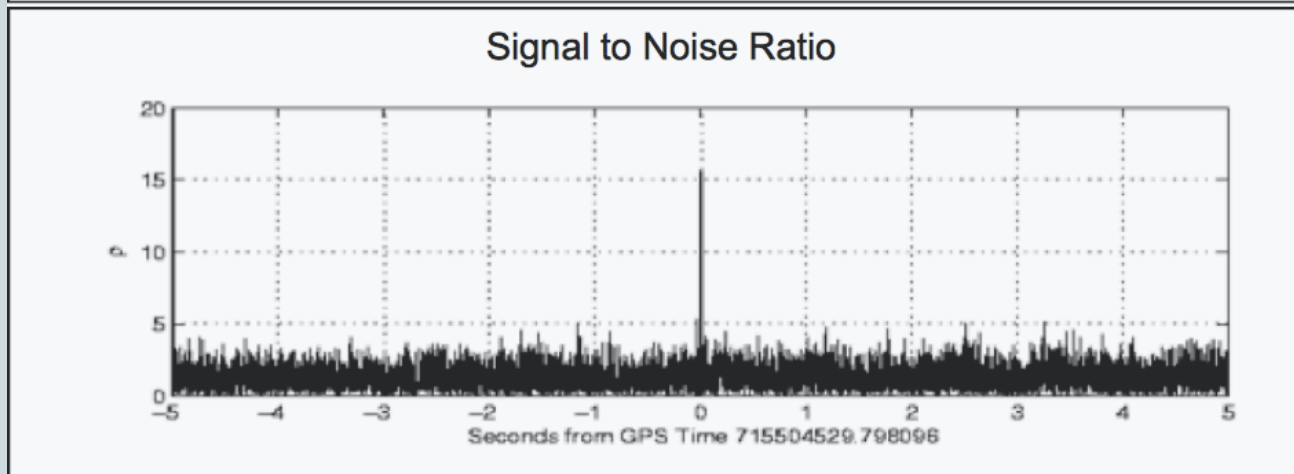
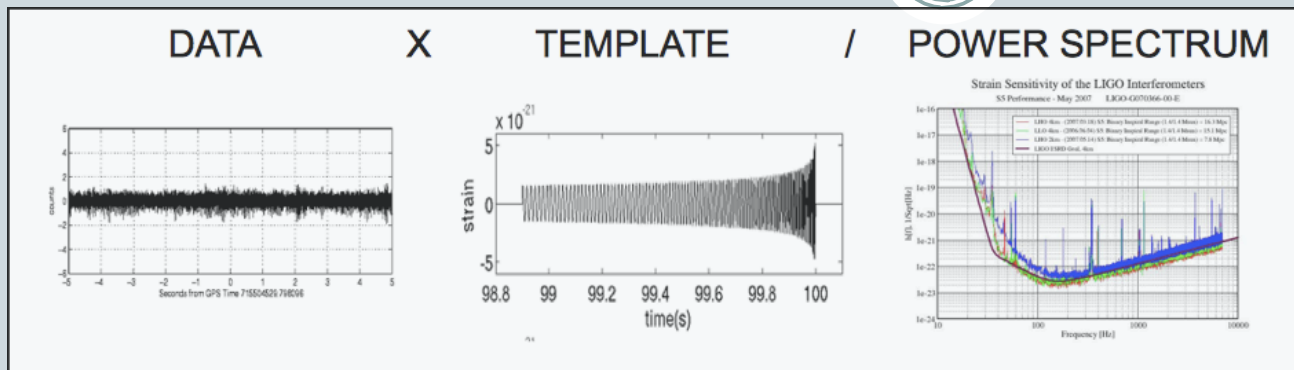


- so our detection statistic is $D(\{\mathbf{y}\}, \mathbf{i}_0) = \sum_k \varepsilon \mathbf{s}_{k-i_0} \mathbf{y}_k$
- in the continuum $\sum_k \varepsilon \mathbf{s}_{k-i_0} \mathbf{y}_k \Rightarrow \int_{-\infty}^{+\infty} dt s(t - \tau_0) y(t)$
- and in the Fourier domain:

$$\rho(\tau_0) = \int d\omega \frac{S^*(\omega, \tau_0) Y(\omega)}{N(\omega)}$$

the standard expression for matched filtering, N being the noise spectrum

Matched filter



FT of waveform with parameters a , at t_0

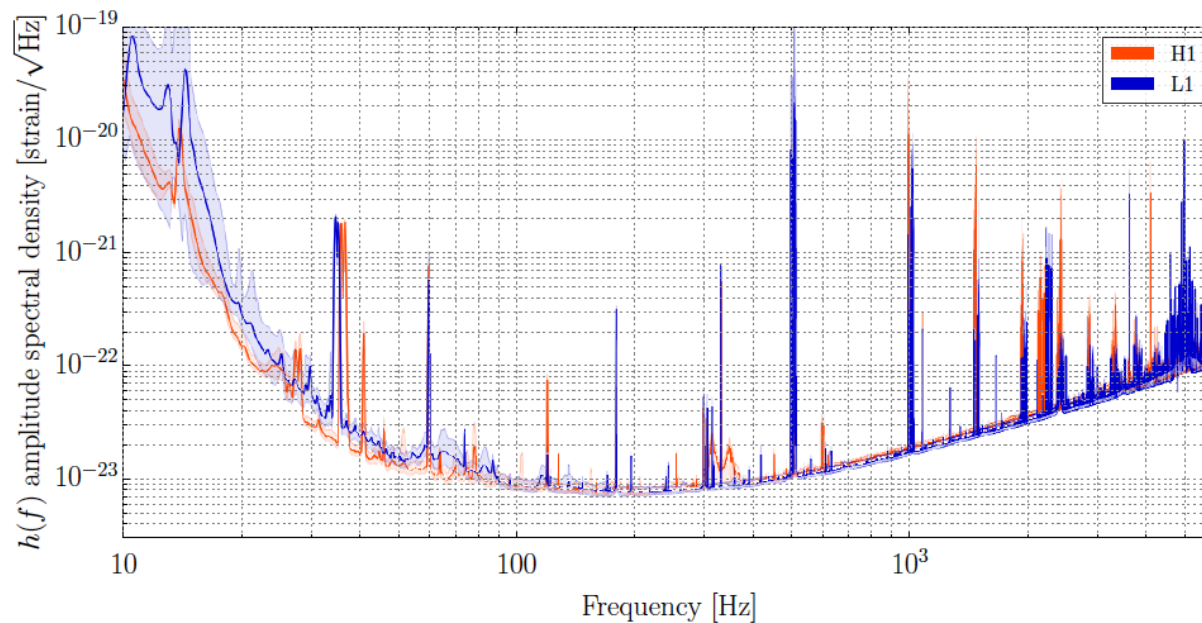
FT of data

$$\rho(t_0, a) = \int d\omega \frac{H^*(\omega, a, t_0) X(\omega)}{S_h(\omega)}$$

noise \rightarrow $S_h(\omega)$

What signals may be detectable ?

time scales of ms to s, compact objects, high accelerations:

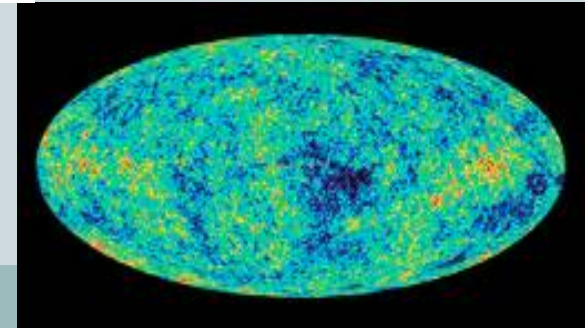
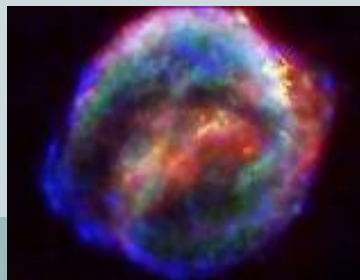


from inspiraling compact objects

bursts, typically arising from catastrophic events

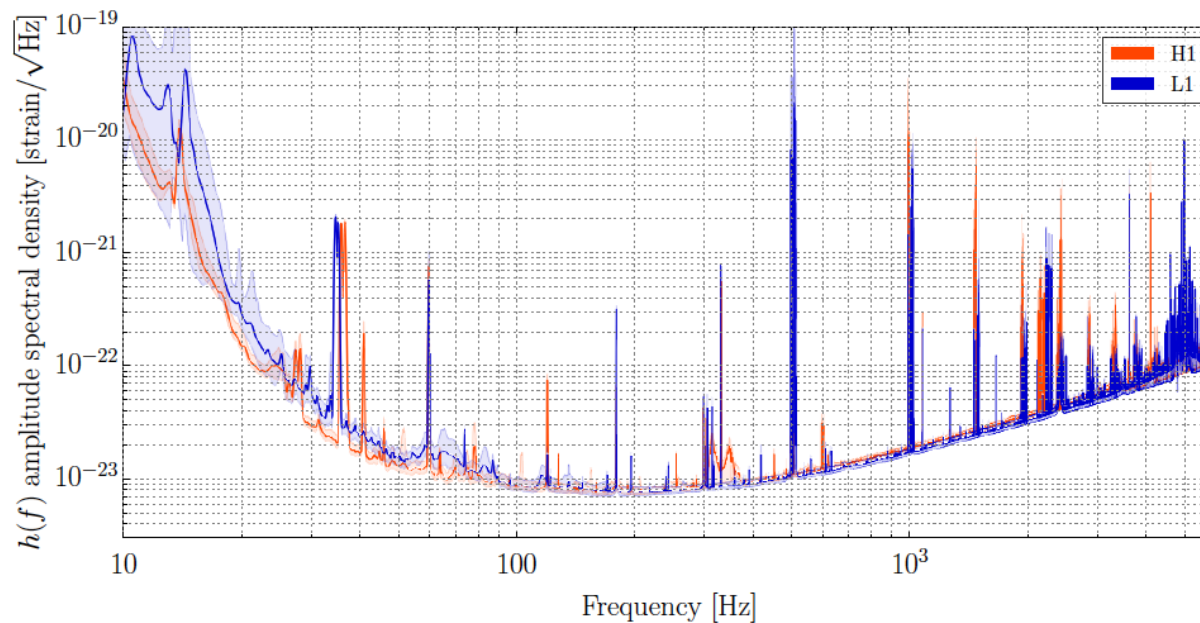
continuous quasi-periodic waves

stochastic background of gravitational radiation



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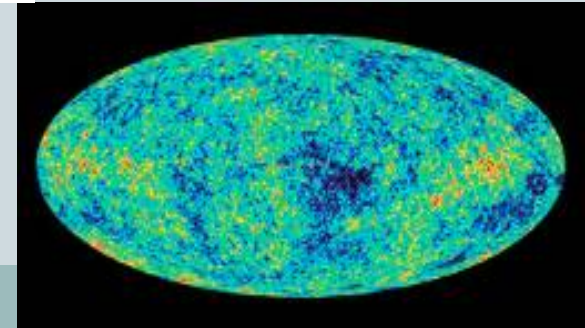
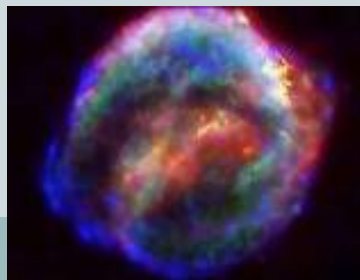


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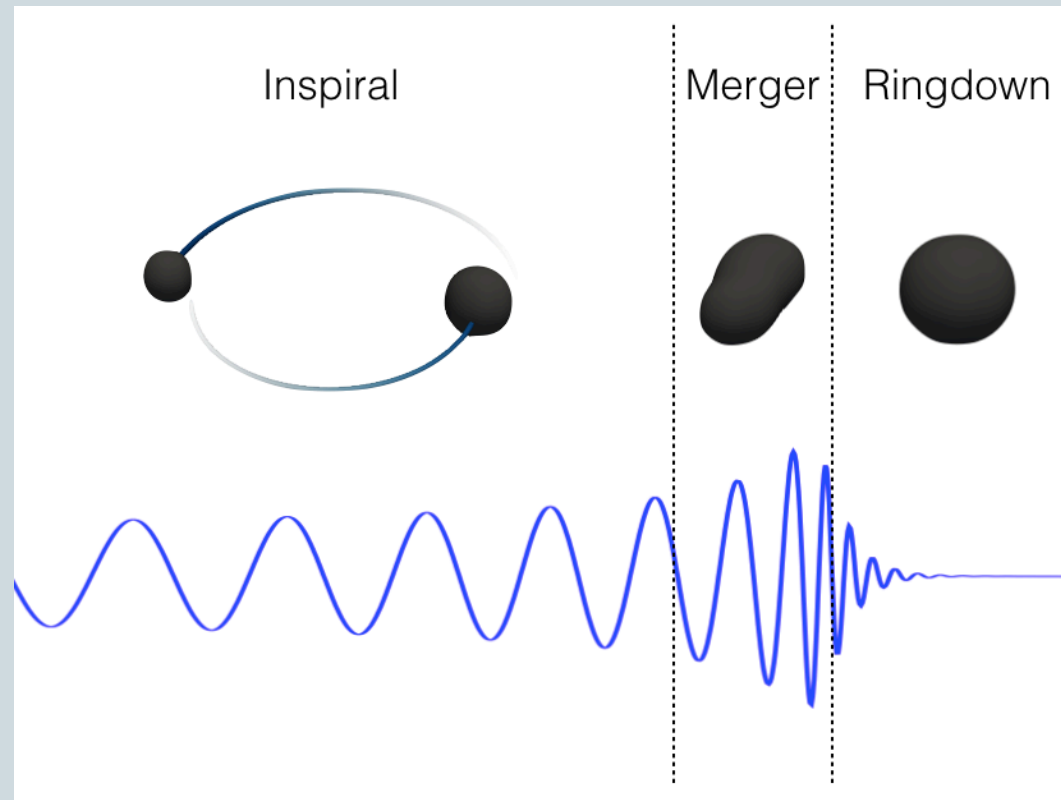
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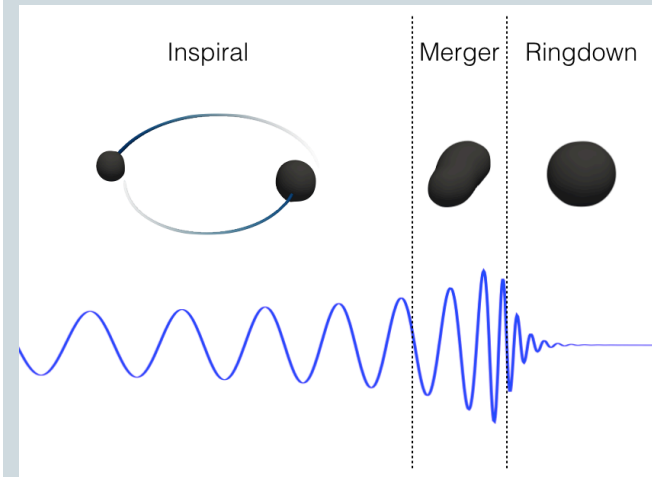
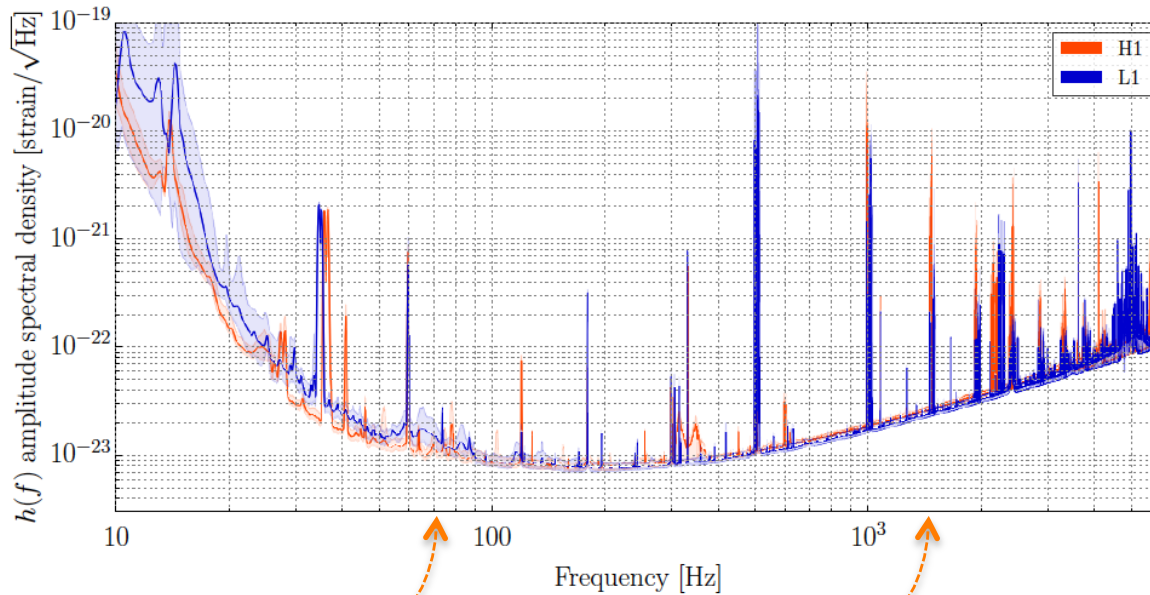
Compact binary coalescences as sources of GW



- Final evolution of compact binary systems involving neutron stars and/or black holes, driven by gravitational radiation



What mass ranges do we see ?

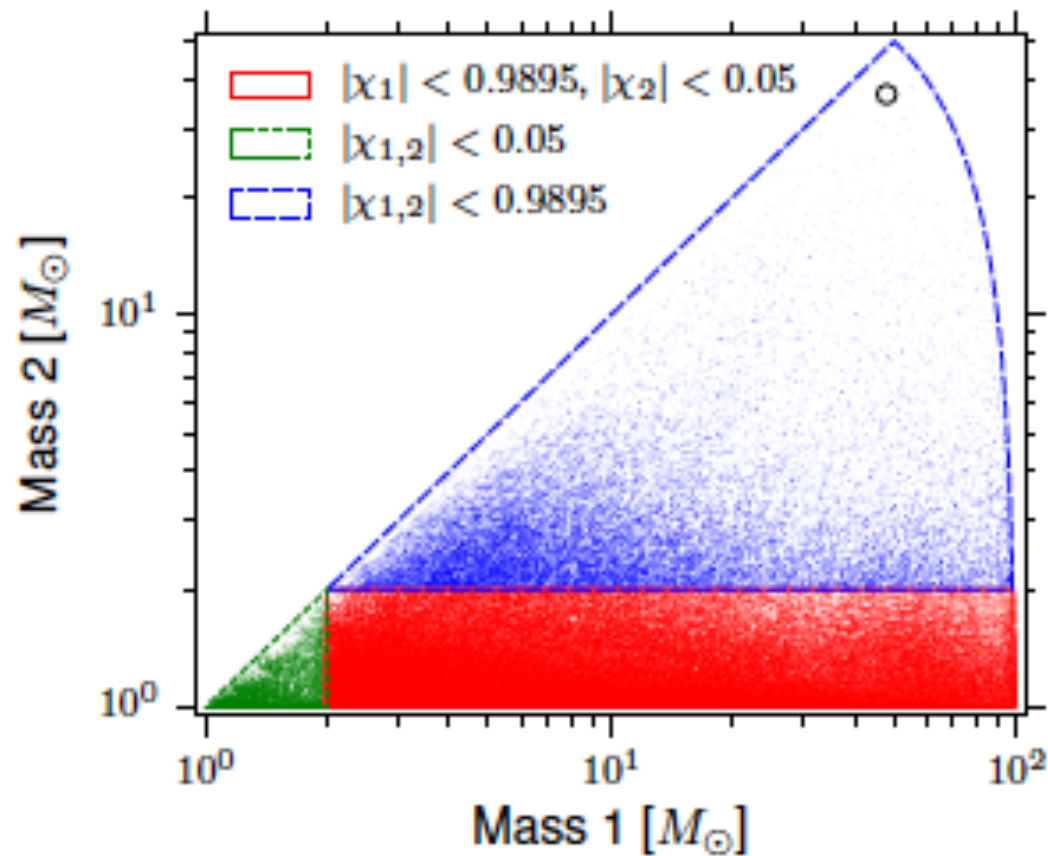


merger freq for
 $M_{\text{tot}} \approx 60 M_{\odot}$

merger freq for
 $M_{\text{tot}} \approx 2.4 M_{\odot}$

4300 Hz M_{\odot}/M_{tot}

Mass range (O1 search)



How do we search for signals ?

Matched filter

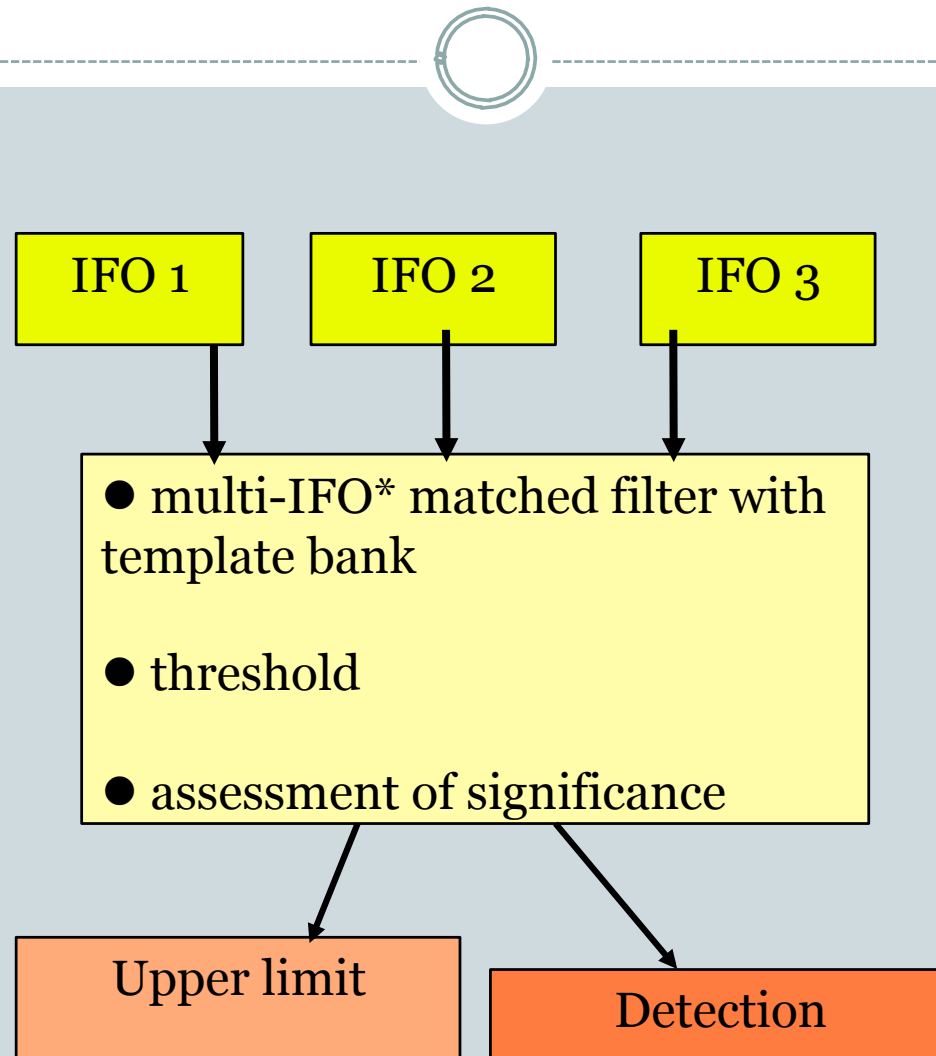
- At best you know what you're looking for; then you use a matched filter:

The diagram illustrates the matched filter equation with labels for its components:

- A box labeled "FT of waveform with parameters a, at t₀" has an arrow pointing to the $H^*(\omega, a, t_0)$ term in the numerator.
- A box labeled "FT of data" has an arrow pointing to the $X(\omega)$ term in the numerator.
- A box labeled "noise" has an arrow pointing to the $S_h(\omega)$ term in the denominator.

$$\rho(t_0, a) = \int d\omega \frac{H^*(\omega, a, t_0) X(\omega)}{S_h(\omega)}$$

Idealized pipeline schematics



*IFO : interferometer

but it's more complicated:



- the matched filter is optimal detection statistic for Gaussian stationary noise but our data are neither Gaussian nor stationary:
 - Weed-out spurious noise:
 - ✦ Data quality flags
 - ✦ Coincidence schemes
 - ✦ Signal-based noise rejection techniques
 - Ad-hoc inspection of interesting candidates:
 - ✦ Correlations with environmental channel
 - ✦ Examine overall status of detectors
 - But the problem remains of assessing the significance
 - ✦ Problem of background/noise estimation

The problem with large spurious noise events



Matched filter: is designed to give a large response when the signal waveform matches the template, but it also gives a large response when the instrumental noise has a large glitch. Even if the glitch shape looks nothing like a waveform, it can still drive the filter to give a large response.

Noisy data: the noise of *GW* detectors presents sporadic prominent non-Gaussian glitches. This is a problem.

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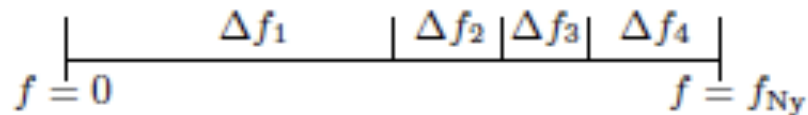
MITIGATION SCHEMES

A counter-measure: signal-based veto, the χ^2 test



If it looks like a duck, quacks like a duck, swims like a duck, then it is a duck.

Does it *really* look like a duck ? the χ^2 test



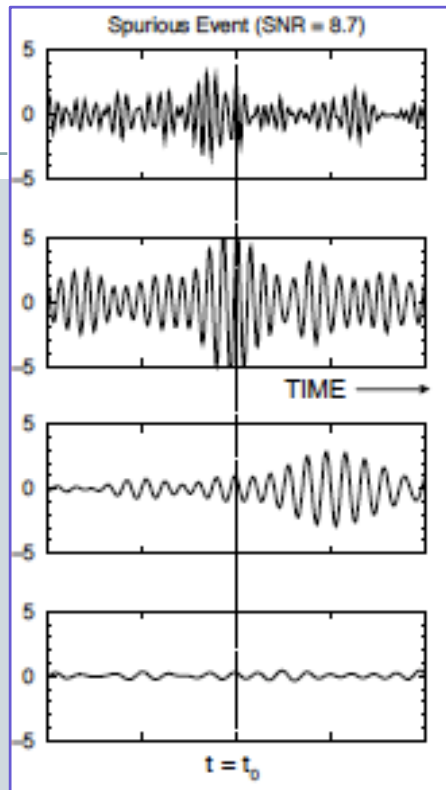
Main idea: consider p frequency sub-bands each contributing the same to the matched filter SNR, z , if a signal is present. Compute the matched filter detection statistic z_j for each of the sub-bands and verify that this is the case:

$$\chi^2 = \frac{p}{2p-2} \sum_{j=1}^p \left(z_j - \frac{z}{p} \right)^2 \quad E[\chi^2] = p - 1 \quad \text{var}[\chi^2] = 2(p - 1)$$

if the hypothesis is correct the residuals are random Gaussian variables and their square sum a chi square variable.

notation note : $\rho \leftrightarrow z$ here

An example



$z_4(t)$

$z_3(t)$

$z_2(t)$

$z_1(t)$

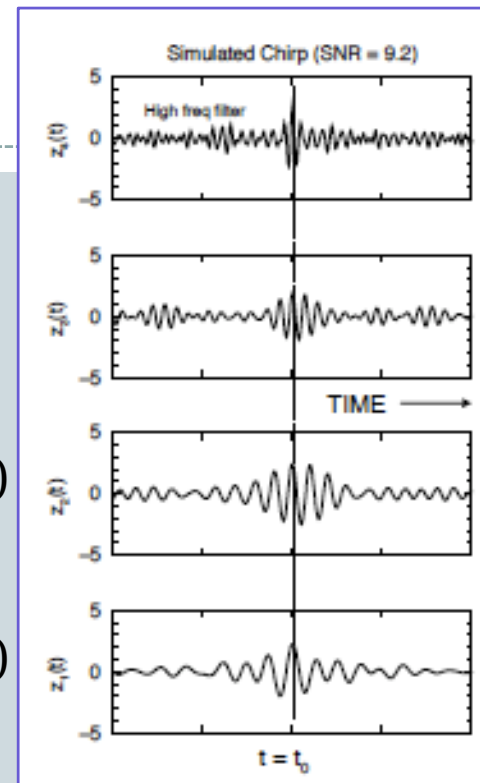


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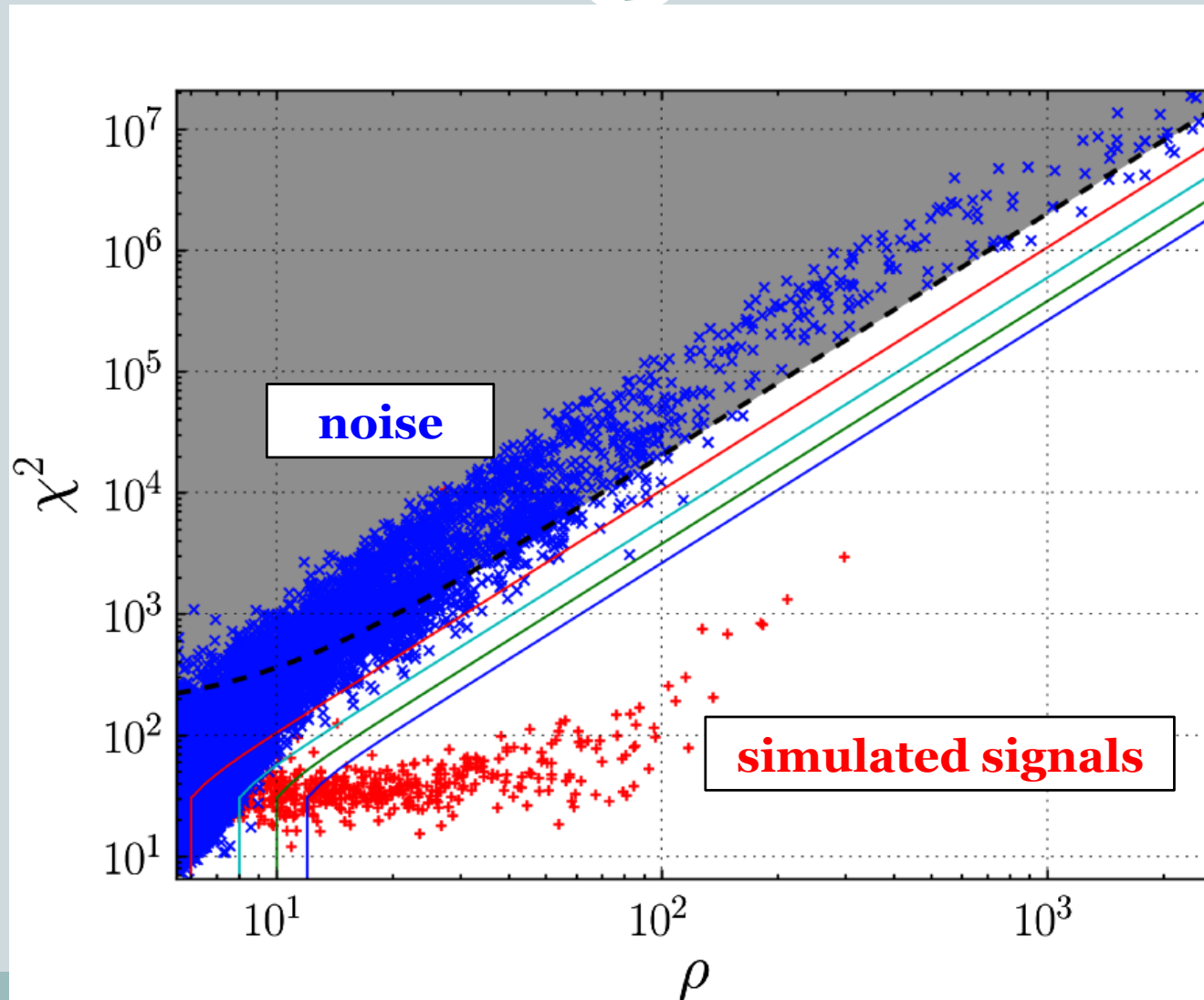


$$\begin{aligned}
 z_1 &= 0.23 \\
 z_2 &= 0.84 \\
 z_3 &= 5.57 \\
 z_4 &= 2.33 \\
 z &= z_1 + z_2 + z_3 + z_4 = 8.97 \\
 \chi^2 &= 4 \sum_{j=1}^4 (z_j - z/4)^2 = 68.4 \\
 P_{\chi^2 \geq 68.4} &= 1 - \frac{\gamma(3/2, 34.2)}{\Gamma(3/2)} = 9.4 \times 10^{-15}.
 \end{aligned}$$

$$\begin{aligned}
 z_1 &= 2.25 \\
 z_2 &= 2.44 \\
 z_3 &= 1.87 \\
 z_4 &= 2.64 \\
 z &= z_1 + z_2 + z_3 + z_4 = 9.2 \\
 \chi^2 &= 4 \sum_{j=1}^4 (z_j - z/4)^2 = 1.296 \\
 P_{\chi^2 \geq 1.296} &= 1 - \frac{\gamma(3/2, 0.648)}{\Gamma(3/2)} = 73\%.
 \end{aligned}$$

notation note : $\rho \leftrightarrow z$ here

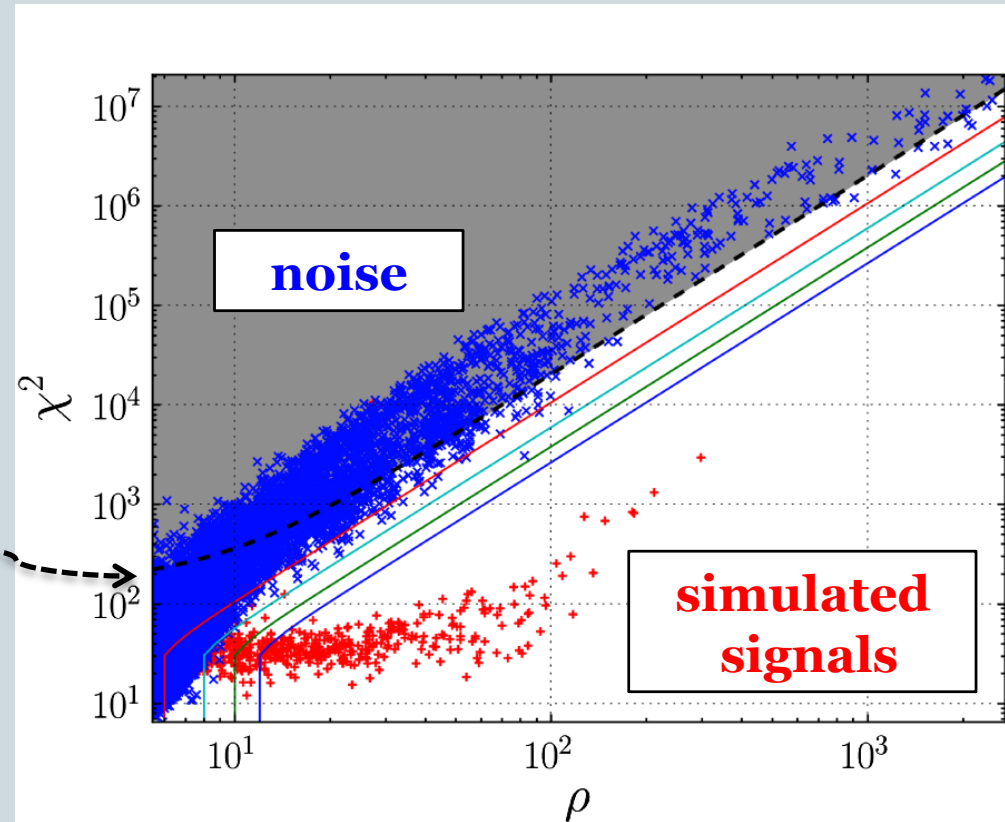
Cannot just threshold on χ^2



Use of χ^2 : a veto



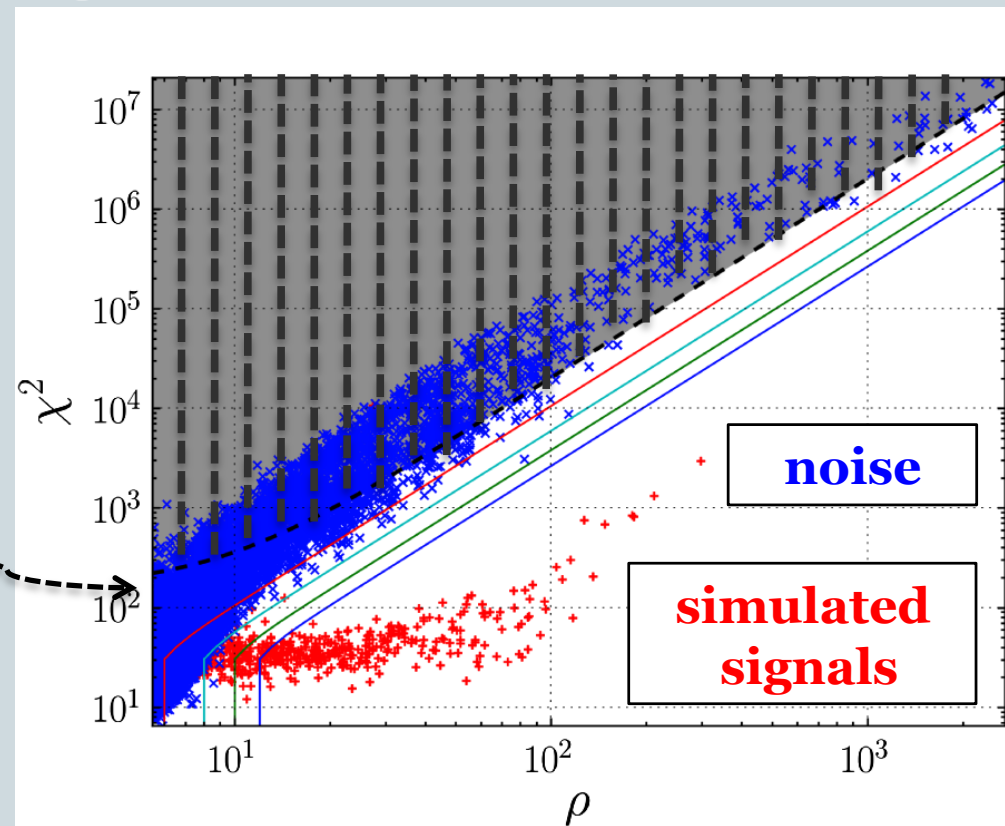
- in previous example number of freq bins $p=4$ usually $p=16$
- $n_{\text{dof}} = 2p-2$
- veto all triggers with $\chi^2 > 10 (p + 0.2 \rho^2)$



Use of χ^2 : a veto



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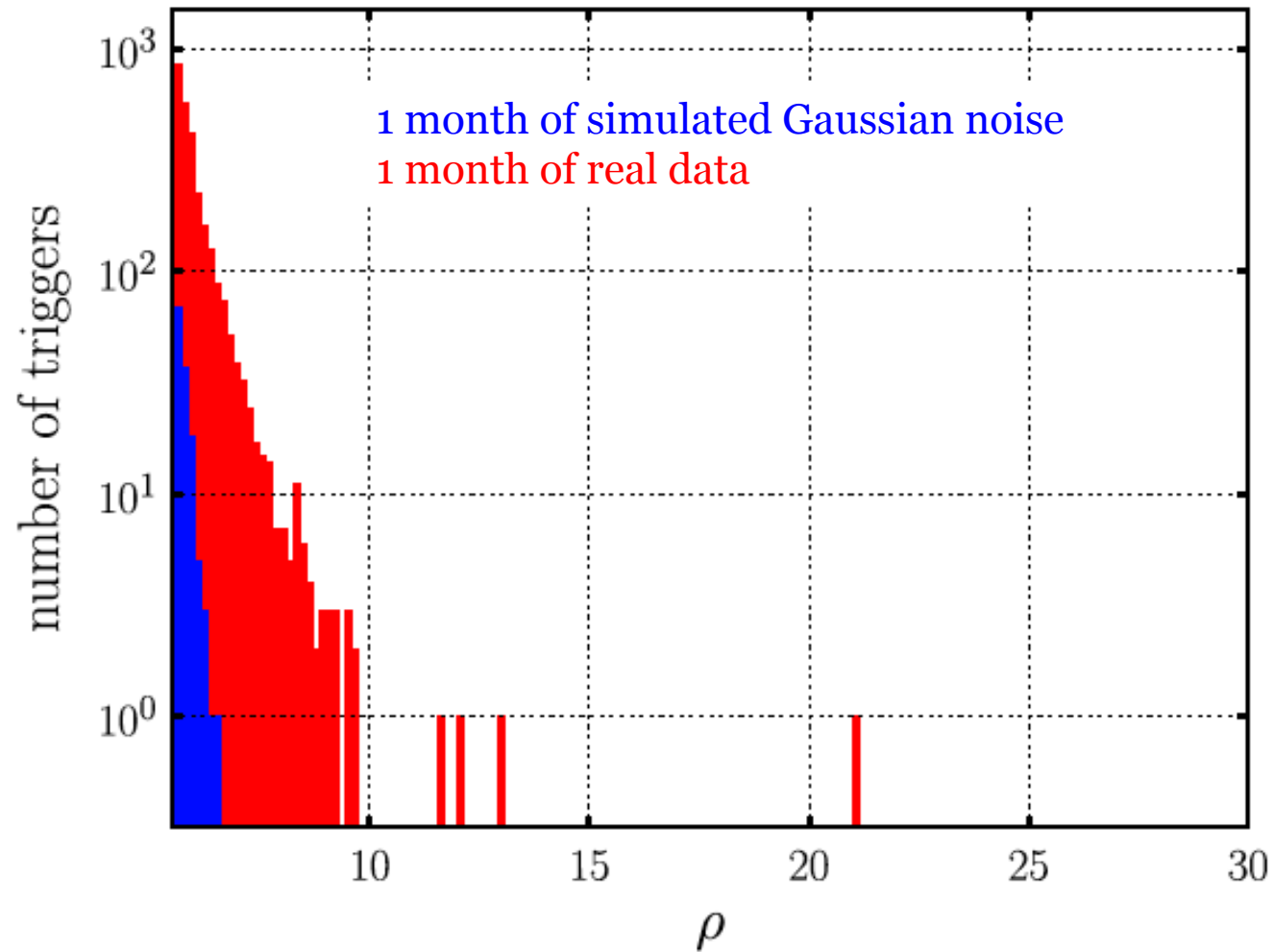
Coincidence requirement



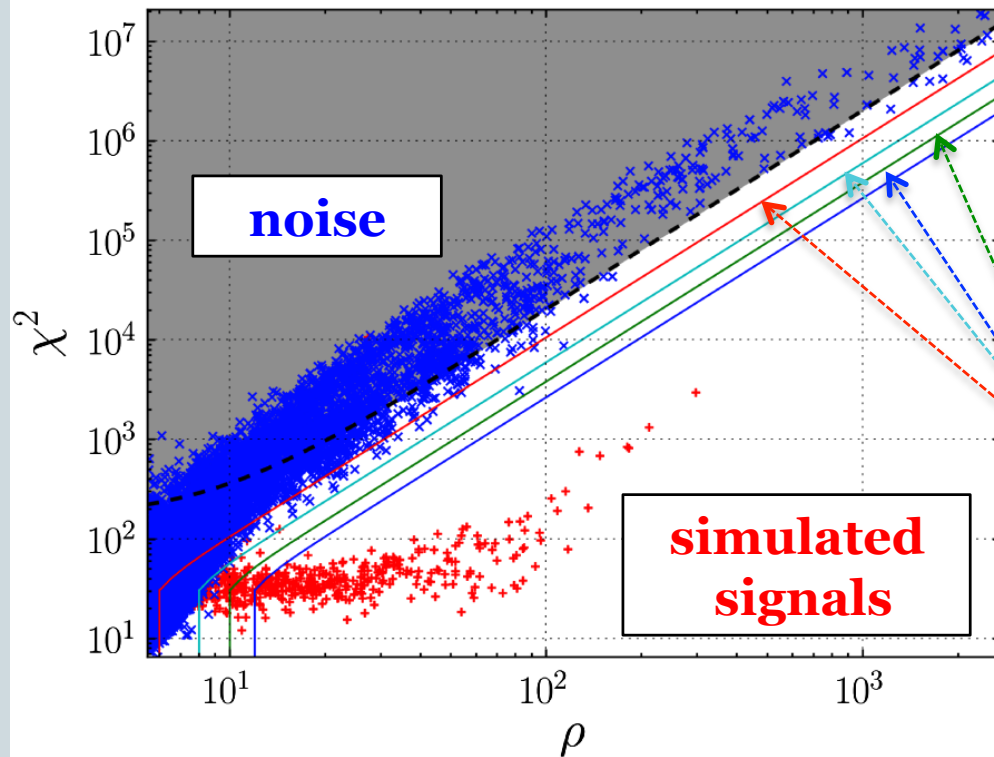
- After the χ^2 veto we also veto triggers in one detector that do not have a consistent counterpart trigger in the other detector
 - Consistency in waveform parameters
 - Close enough in time



But still many spurious triggers



Go back to single-detector triggers and construct a new detection statistic

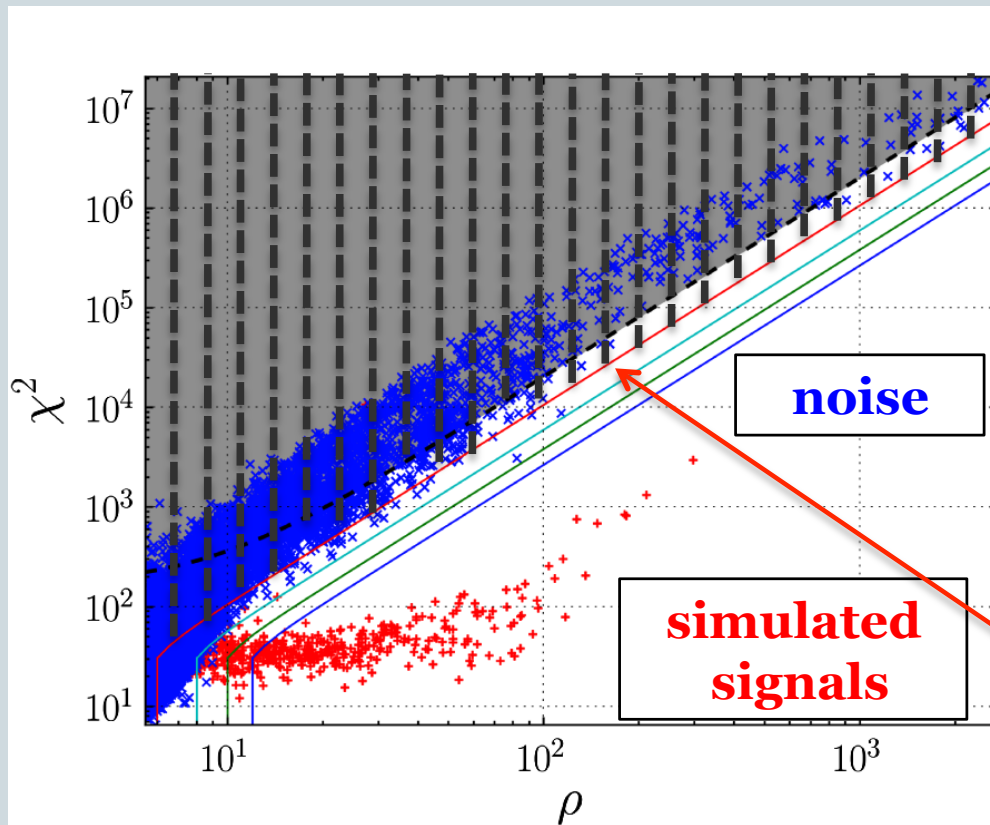


$$\rho_{\text{new}}^2 = \begin{cases} \rho^2 & \text{if } \chi_r^2 \leq 1 \\ \frac{\rho^2}{\left\{ \frac{1}{2} [1 + (\chi_r^2)^3] \right\}^{1/3}} & \text{if } \chi_r^2 > 1 \end{cases}$$

with $\chi_r^2 = \chi^2 / n_{\text{dof}}$

This statistic matches the constant false alarm contours and discriminates well between signals and noise.

A new detection statistic



$$\rho_{\text{new}}^2 = \begin{cases} \rho^2 & \text{if } \chi_r^2 \leq 1 \\ \frac{\rho^2}{\left\{ \frac{1}{2} [1 + (\chi_r^2)^3] \right\}^{1/3}} & \text{if } \chi_r^2 > 1 \end{cases}$$

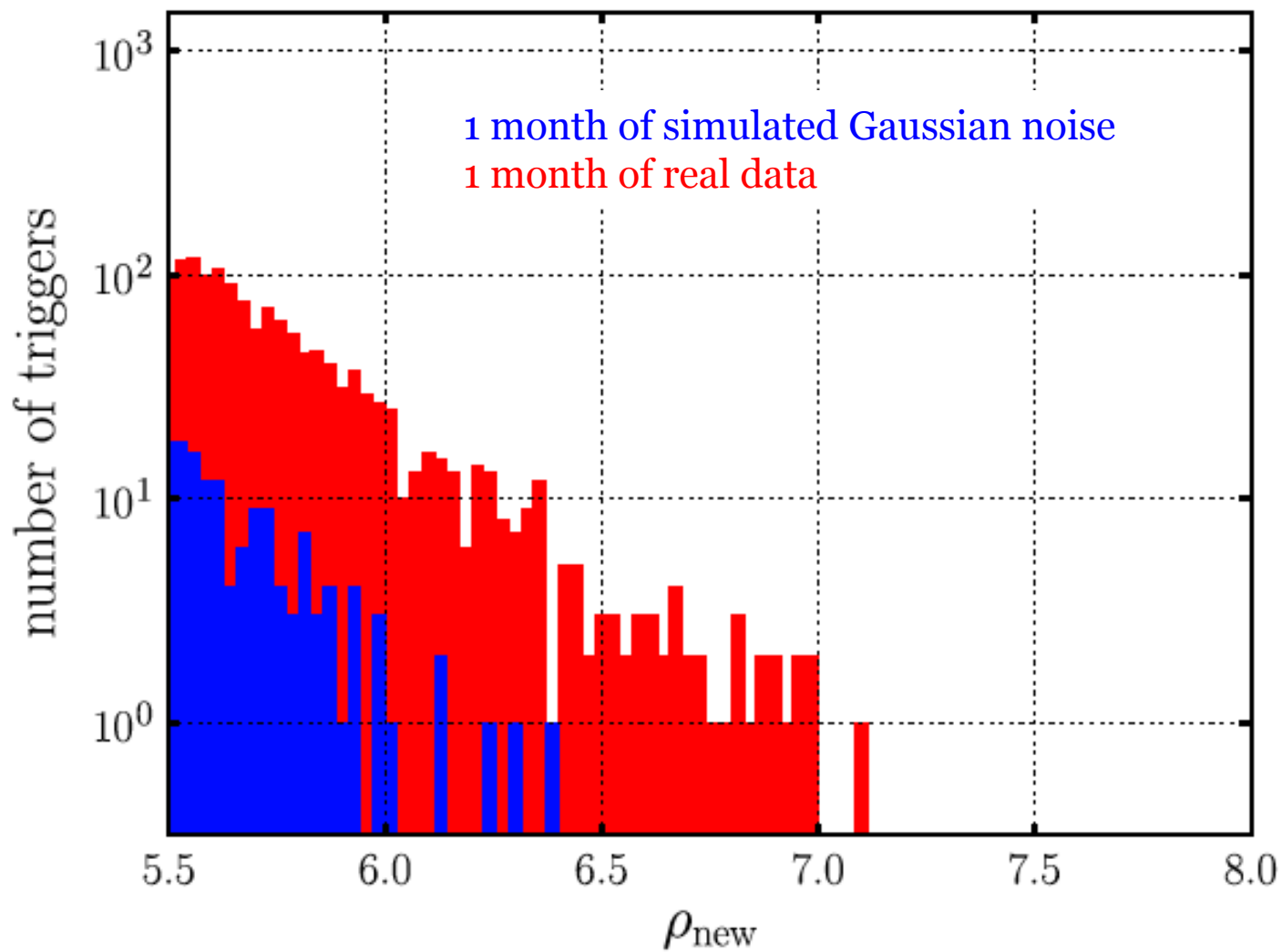
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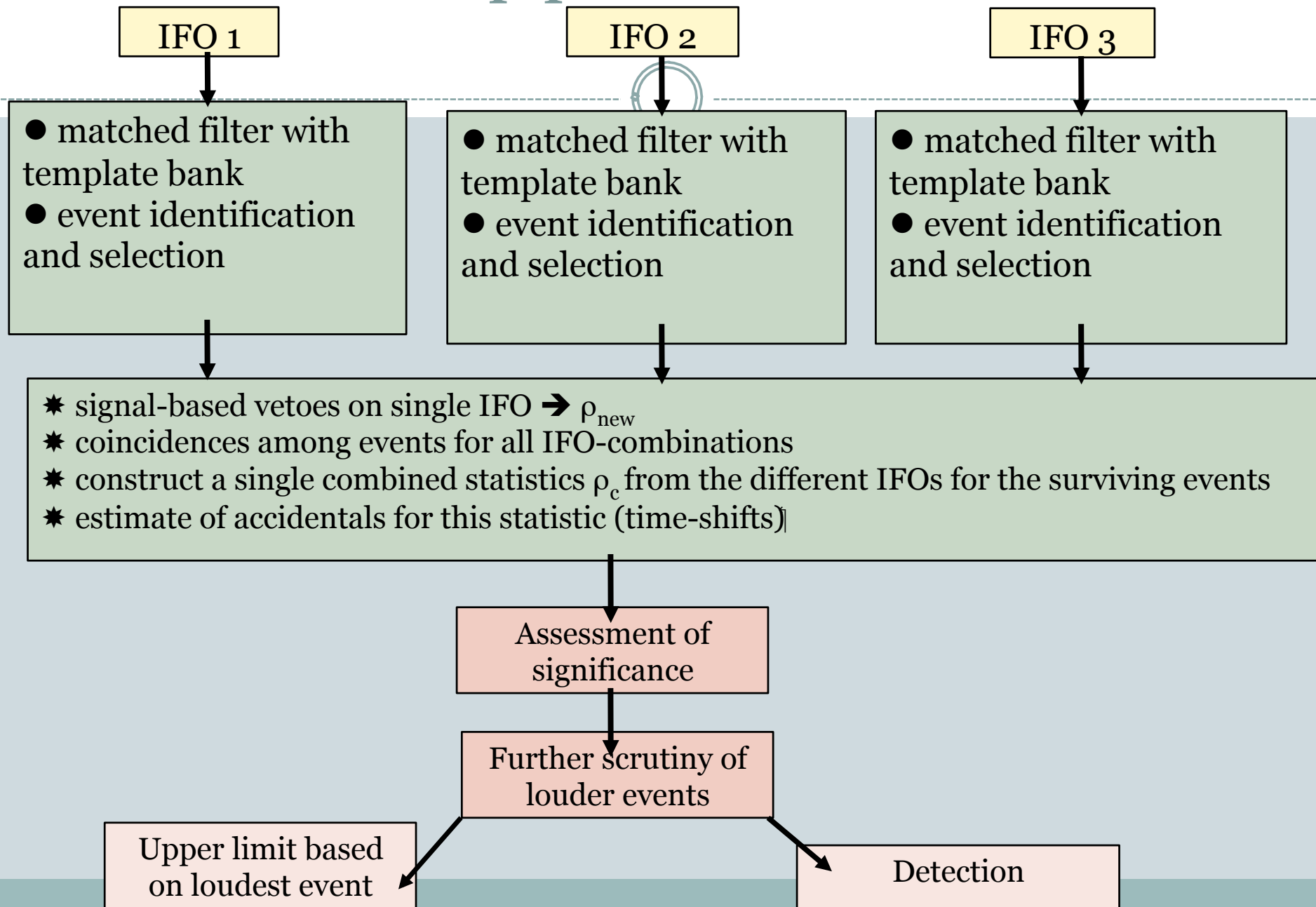
Triggers with $\rho_{\text{new}}^2 < 5.5$ are vetoed
Then again coincidences

The final combined detection statistic: $\rho_c^2 = \sum \rho_{\text{new},i}^2$

Remarkable improvement



Actual pipeline schematics



Assessment of significance

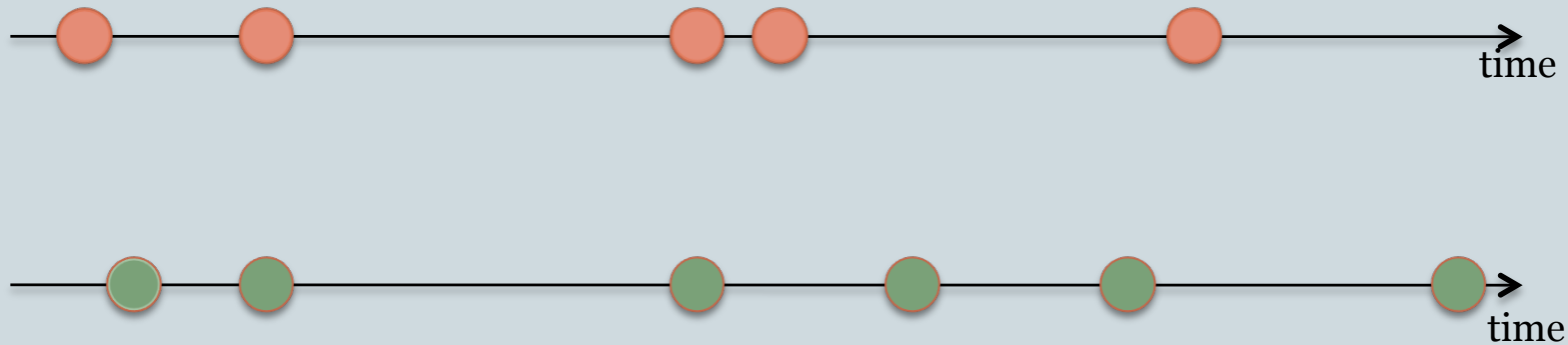


- The analysis produces a list of coincident triggers, each with an associated combined SNR, ρ_c .
- These triggers need to be compared with those that one would obtain by chance, i.e. the accidentals, the background. We do this by comparing the distributions.
- How do we estimate the background? We repeat the analysis on off-source data (by time-shifting the data streams).
- If an on-source coincidence trigger is significantly above the estimated background, then it is a candidate event that warrants further inspection.

Time-slides



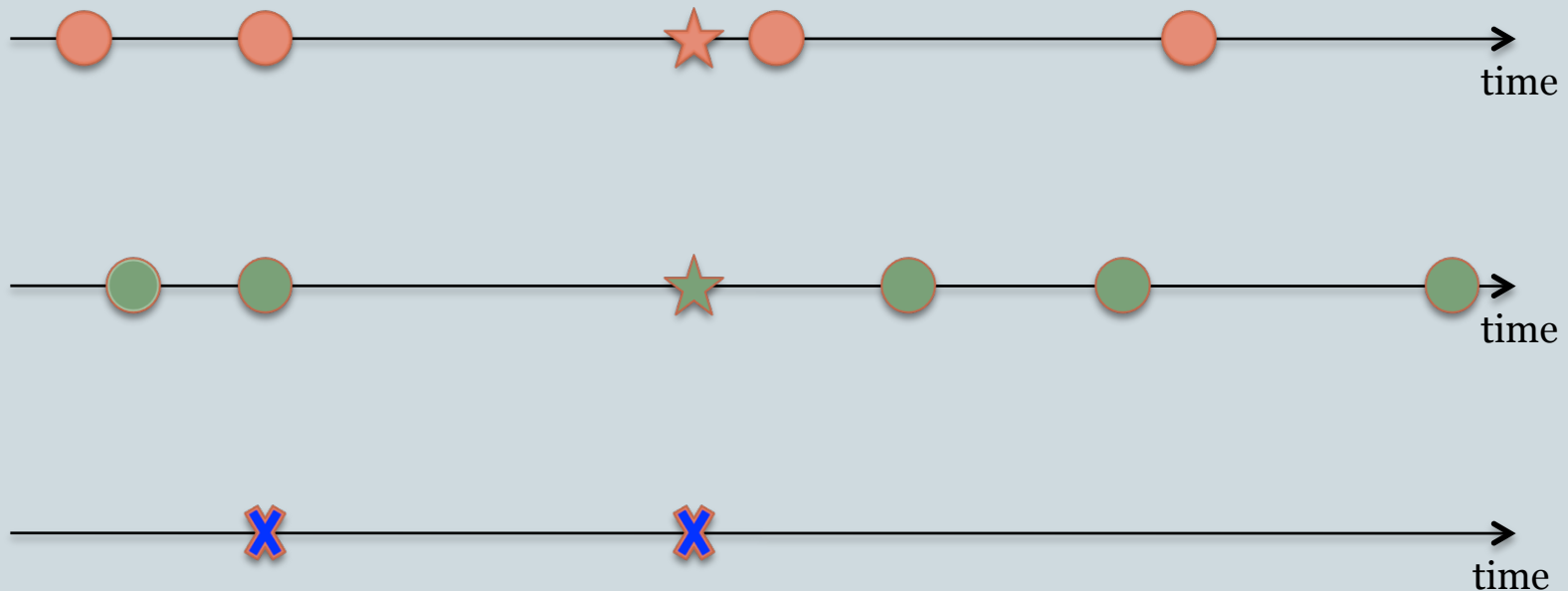
- triggers from detector 1
- triggers from detector 2



Time-slides



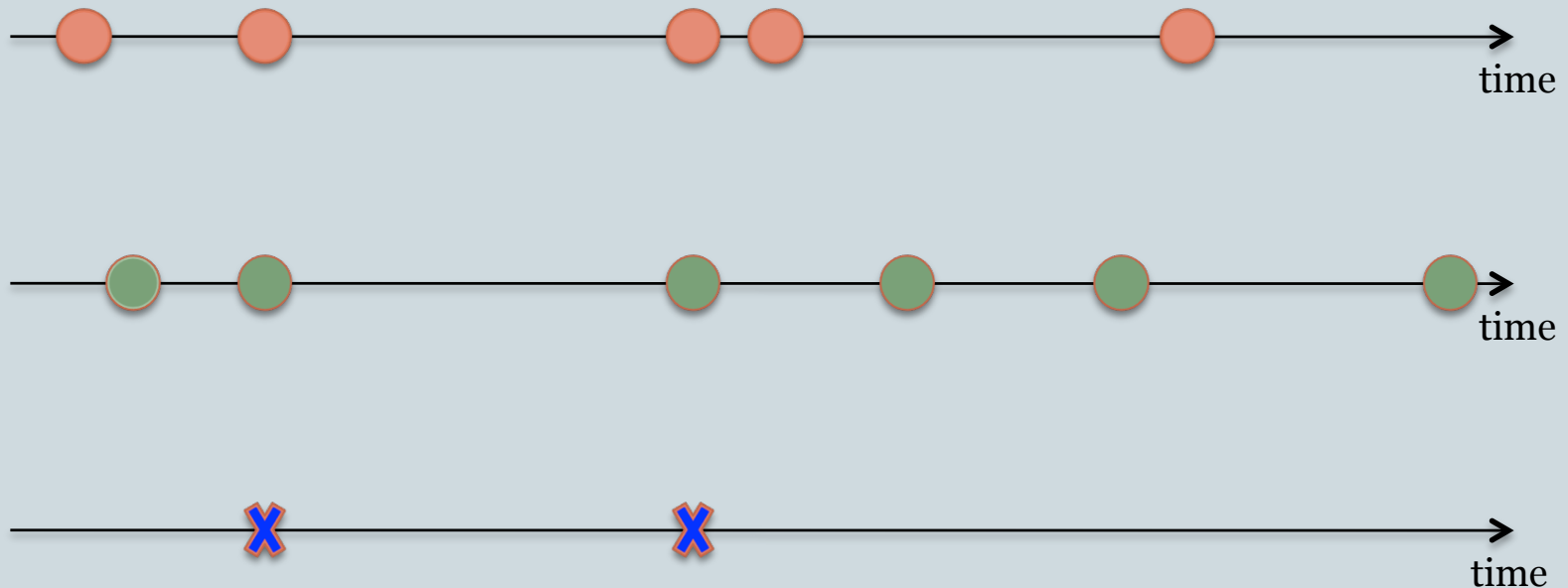
- triggers from detector 1
- triggers from detector 2
- ✕ coincident triggers
- ★ trigger due to GWs
- ★ trigger due to GWs



Time-slides



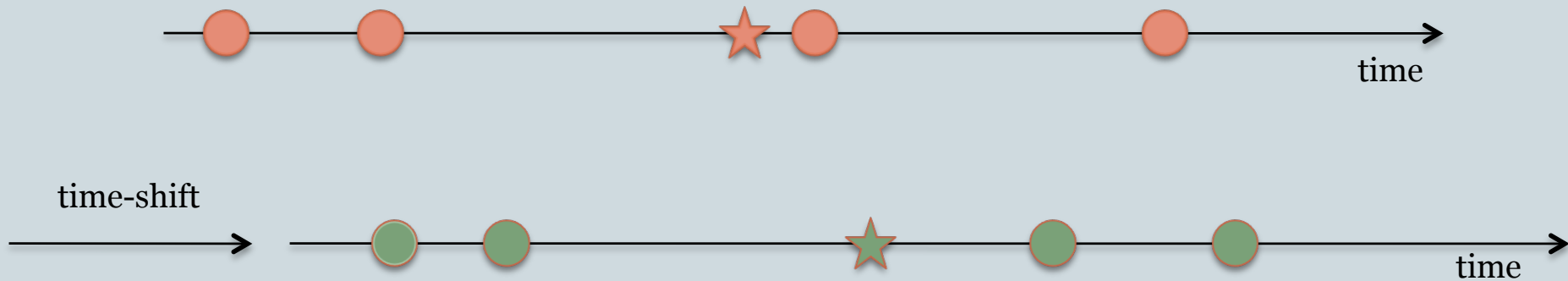
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Time-slides



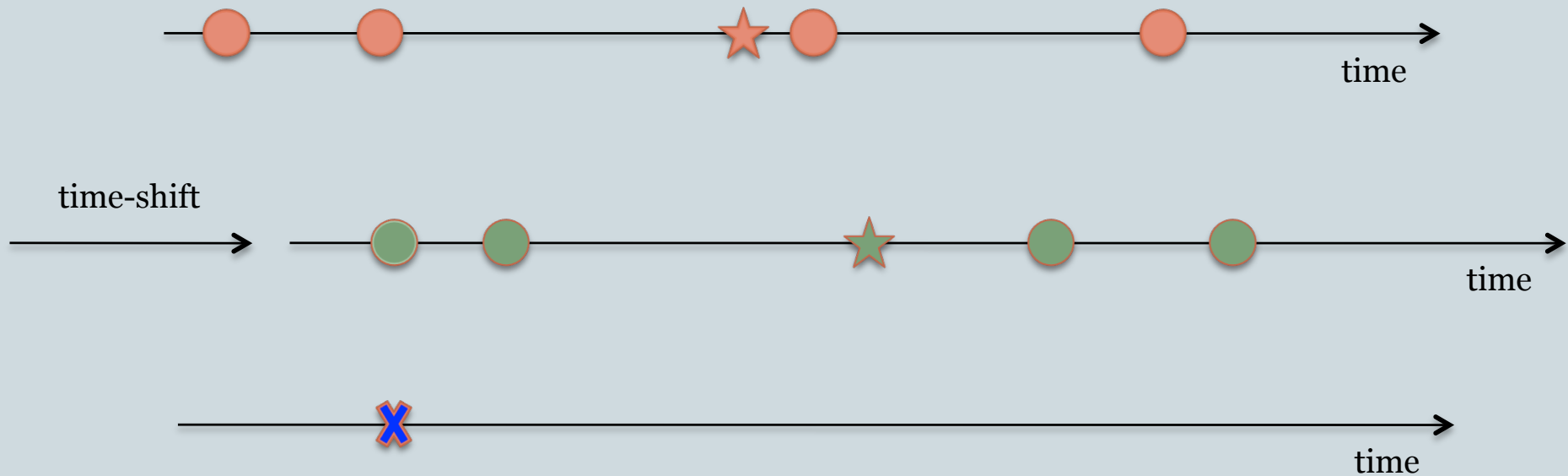
- ★ trigger from detector 1 due to a GW
- ★ trigger from detector 2 due to a GW



Time-slides



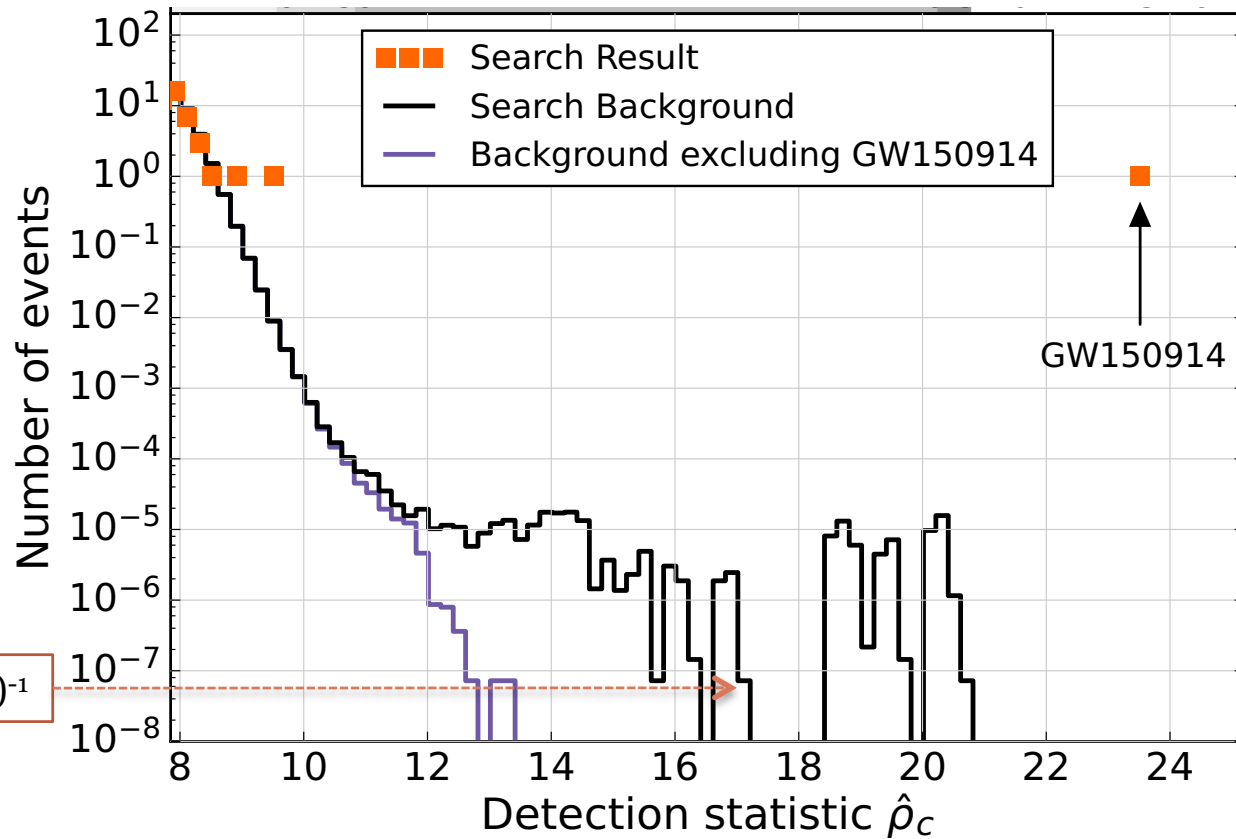
- ★ trigger from detector 1 due to a GW
- ★ trigger from detector 2 due to a GW
- ✕ coincident triggers



The first GW signal (GW150914)



1.4×10^7 time slides corresponding to 608 000 yrs of simulated background.

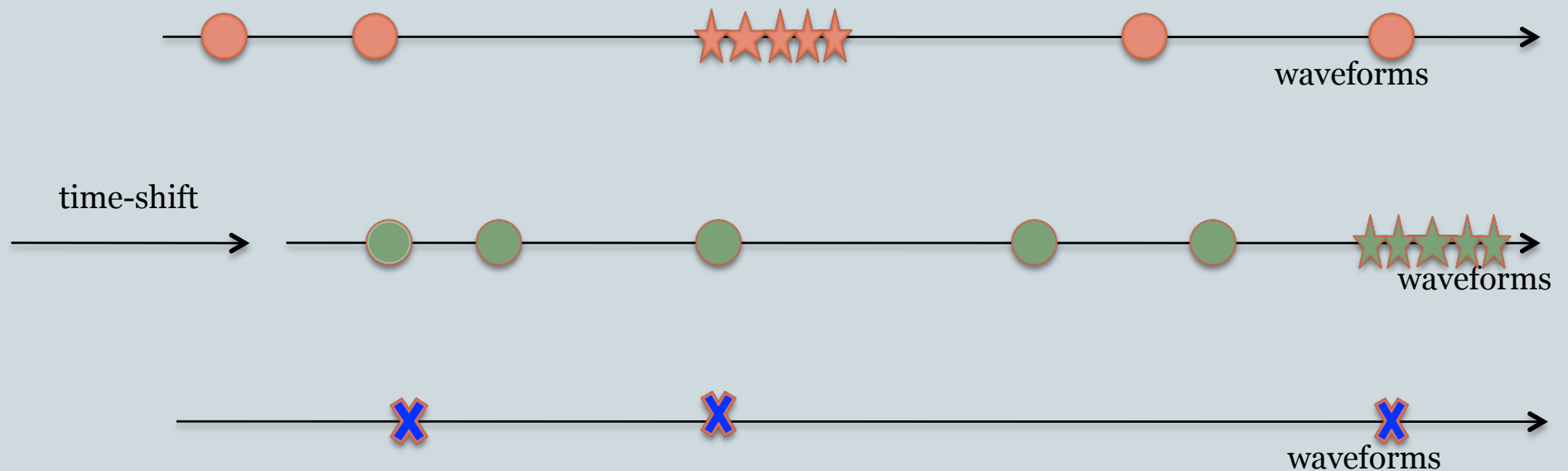


$7 \times 10^{-8} \approx (1.4 \times 10^7)^{-1}$

Time-slides: a conservative estimate of the background



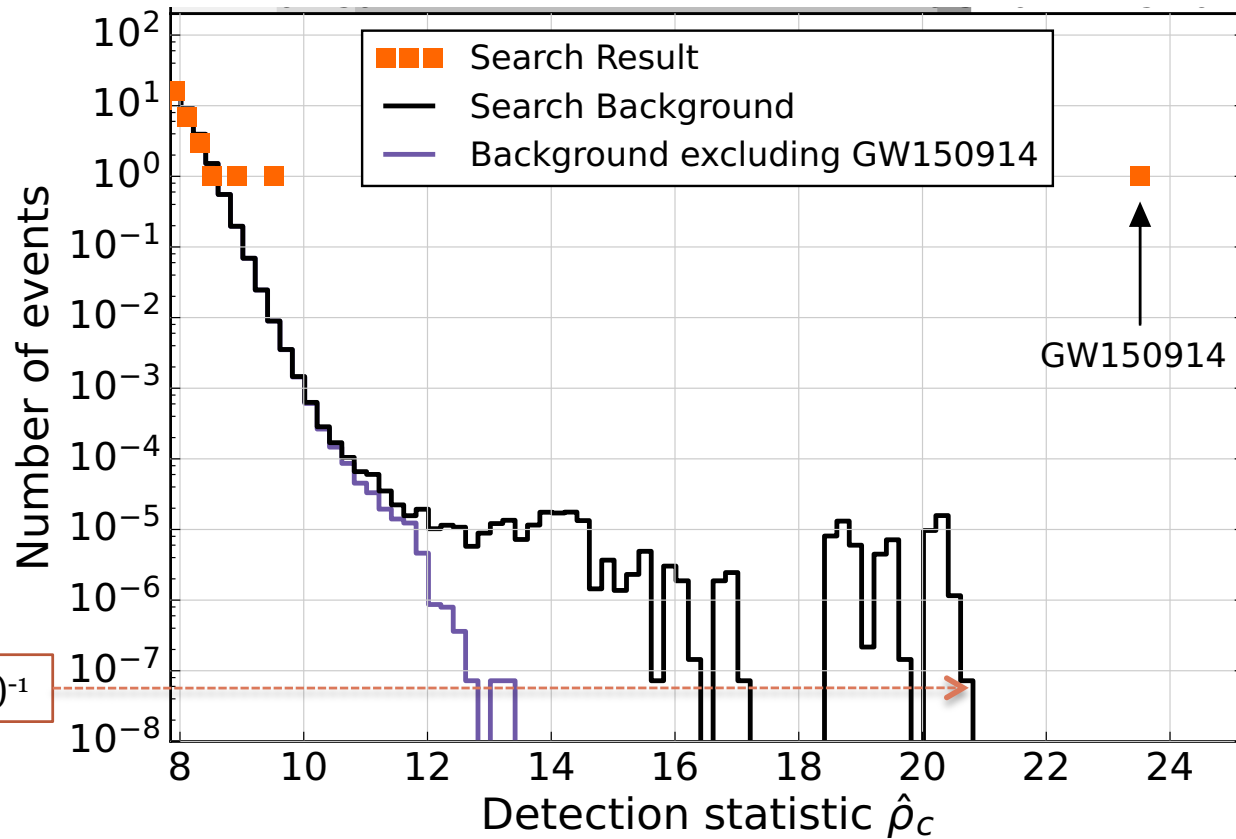
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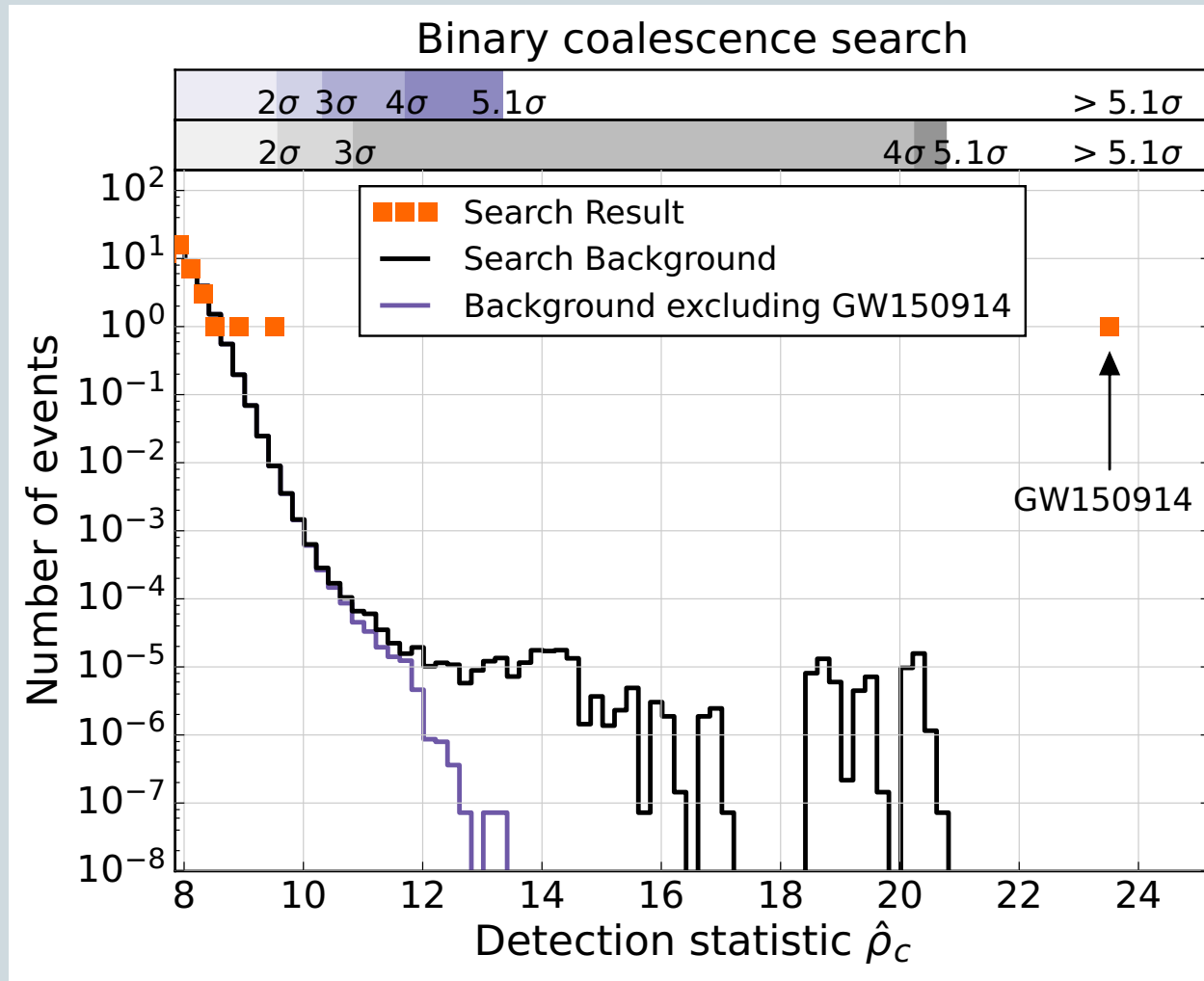
$7 \times 10^{-8} \approx (1.4 \times 10^7)^{-1}$

Significance numbers for GW150914



- Analysis time (livedtime): 14.14 days
- Could not get as significant detection statistic value in 1.375×10^7 realizations of the experiment.
- This corresponds to an analysis background time of $1.375 \times 10^7 \times 14.14 \text{ days} / 365 \text{ days} = 608\,000 \text{ yrs}$
- Let's take the most significant event at $\rho_c \sim 21$
- False alarm rate (FAR) = 1 event / background-time = $1.6 \times 10^{-6} \text{ yr}^{-1}$
- FAR \rightarrow 3 x FAR X because 3 independent searches were performed (trials factor). FAR = $4.9 \times 10^{-6} \text{ yr}^{-1}$
- Poisson process with $\lambda = \text{FAR} \times \text{livedtime} = 2 \times 10^{-7}$
- The probability to measure one event or more in a Poisson process with that average rate λ is FAP = 2×10^{-7} (it's the same as λ because $\lambda \ll 1$)
- The Gaussian sigma level corresponding to such FAP is 5.1

The first GW signal (GW150914)

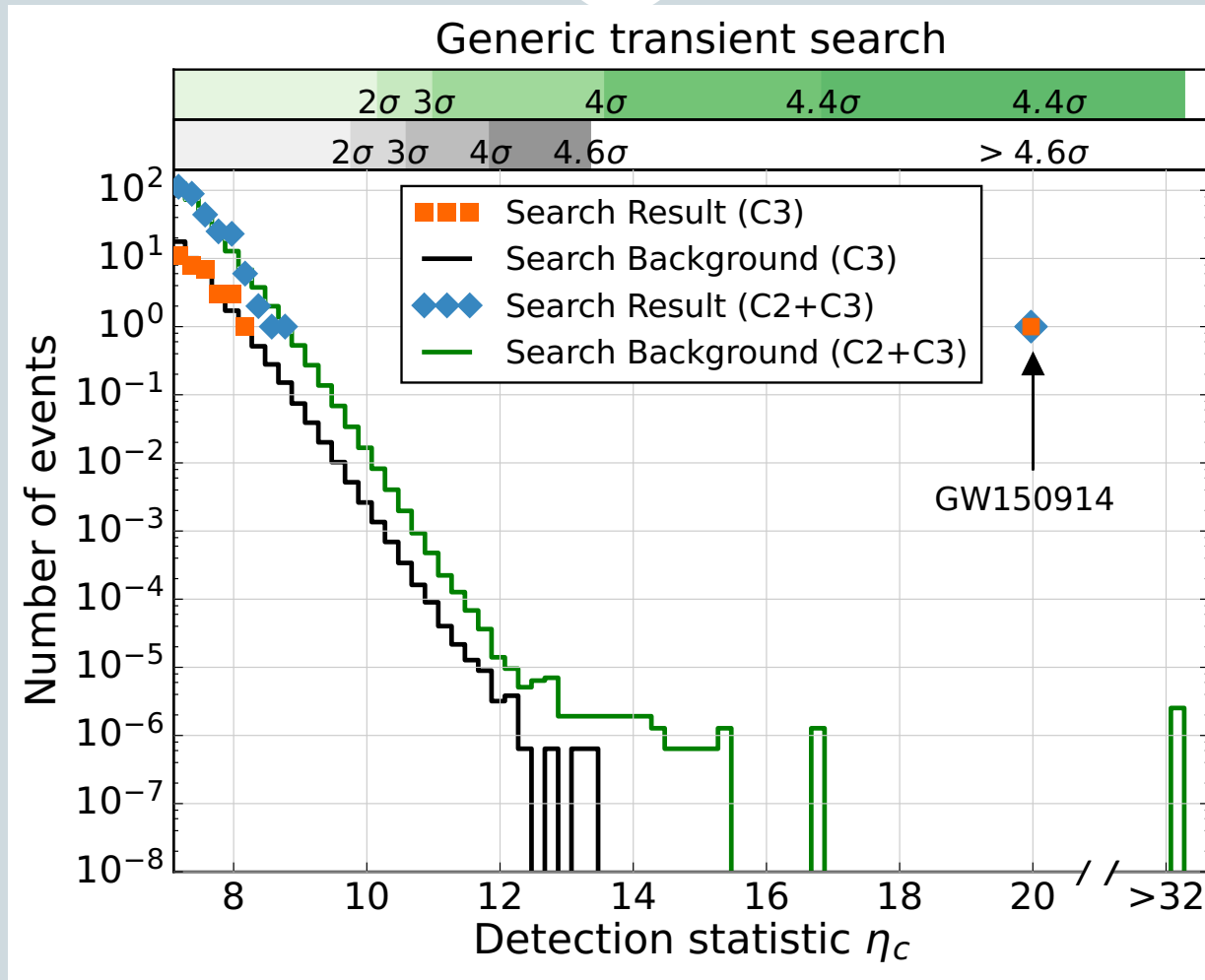


Further inspection means

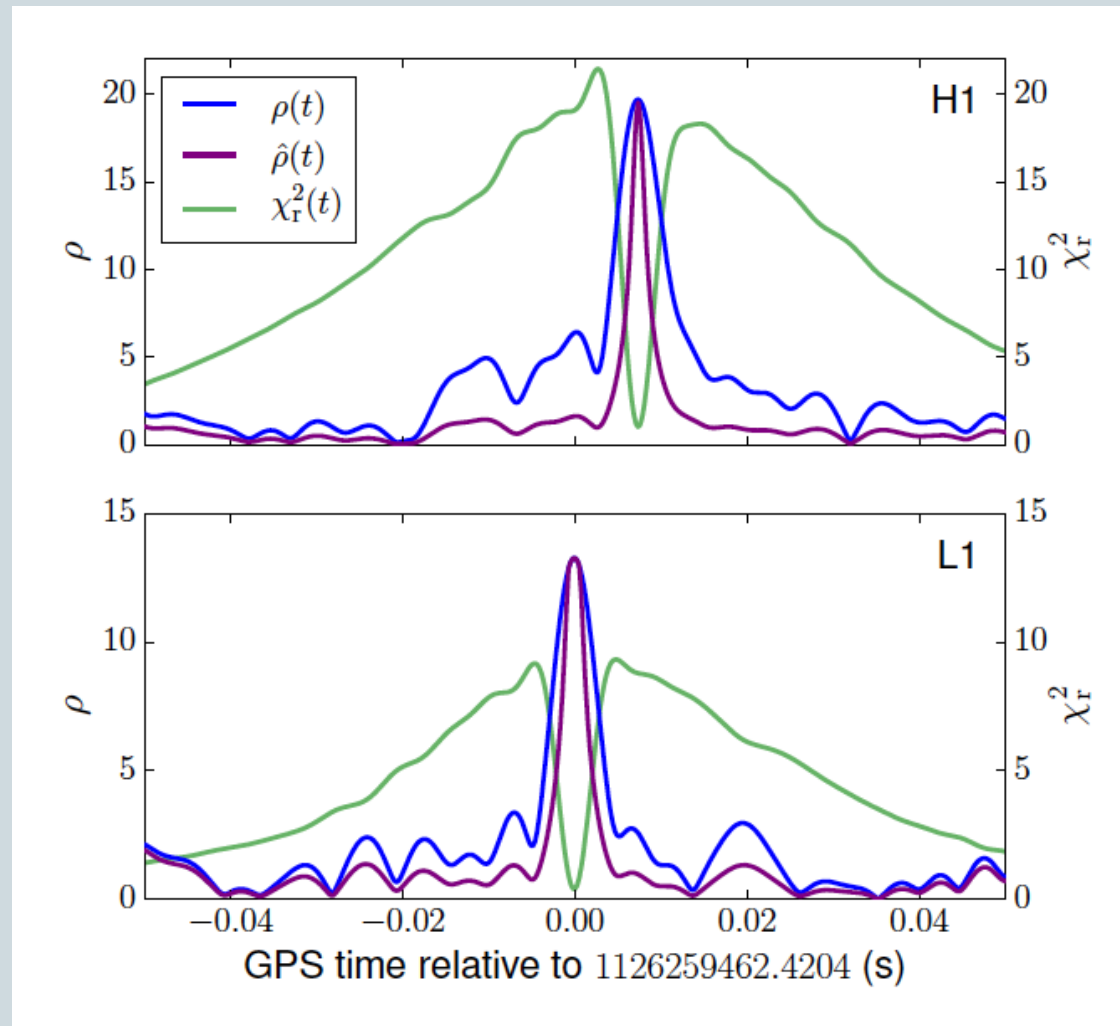


- Statistical significance of the candidate (cross check with other pipelines)
- Status of the interferometers
- Check for environmental or instrumental causes
- Check intermediate stages of the analysis
- Check for coincidences with non-GW searches: other E/M or particle detectors when relevant

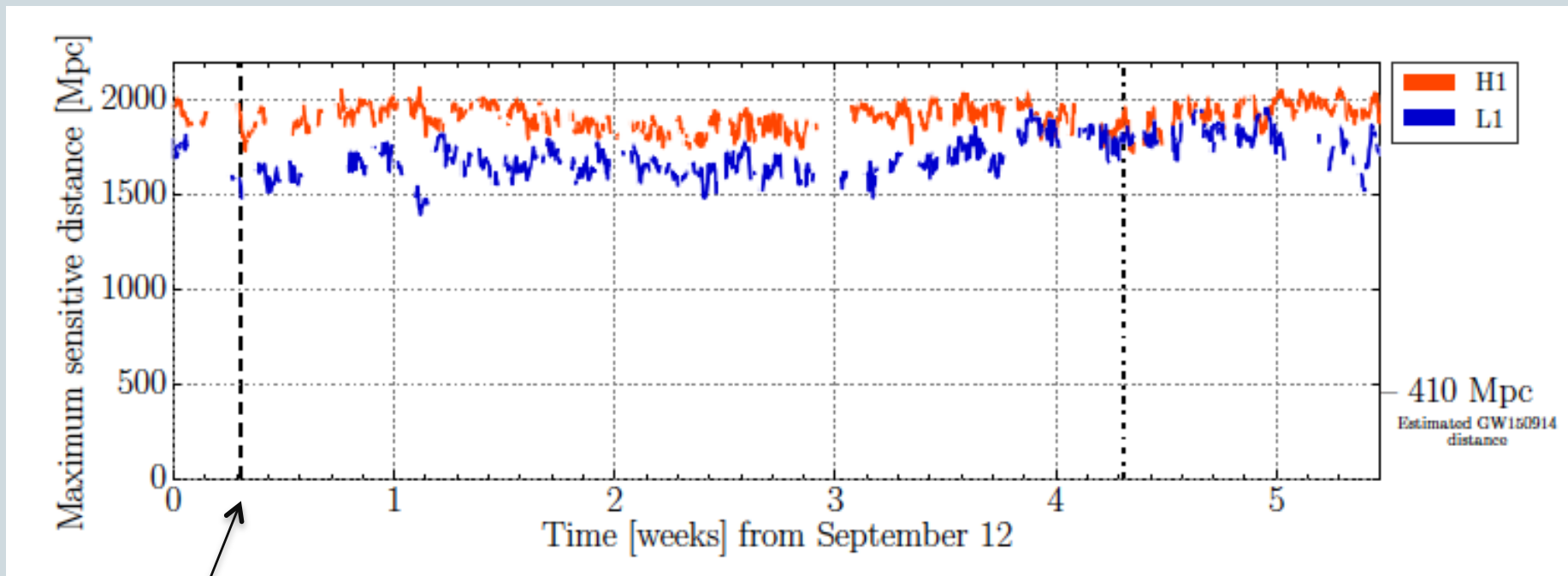
GW150914 as seen by the generic short-transient signals search



Intermediate analysis products



State of the interferometers: stability



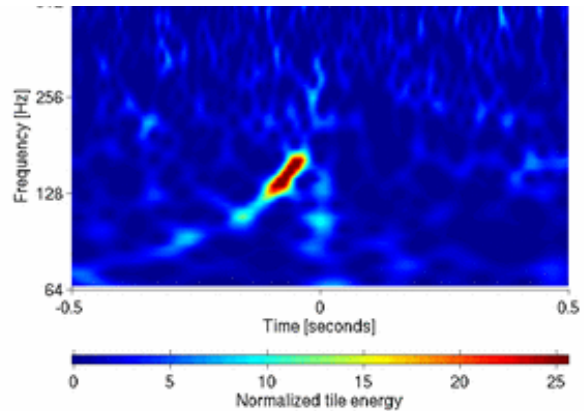
GW150914

Candidate appearance, examples from LIGO science runs

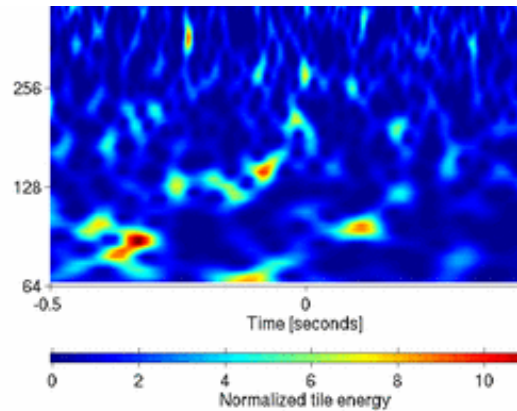


Trigger from simulated signal

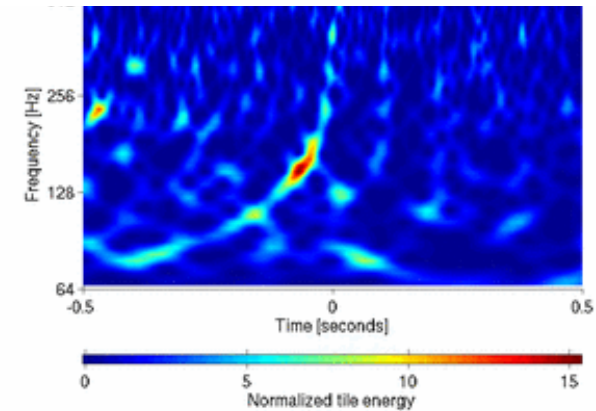
GW channel: H1



GW channel: H2

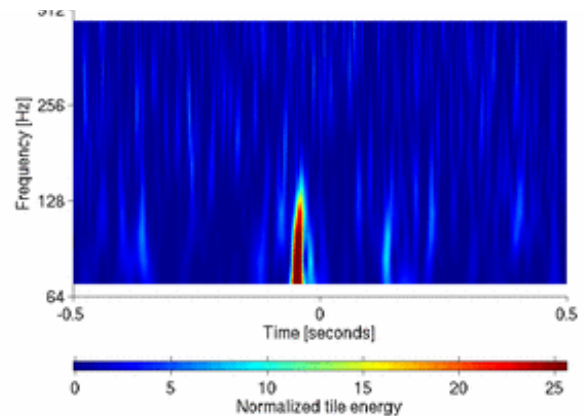


GW channel: L1

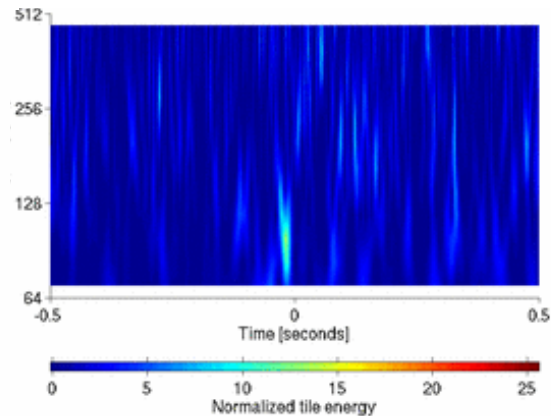


Background trigger

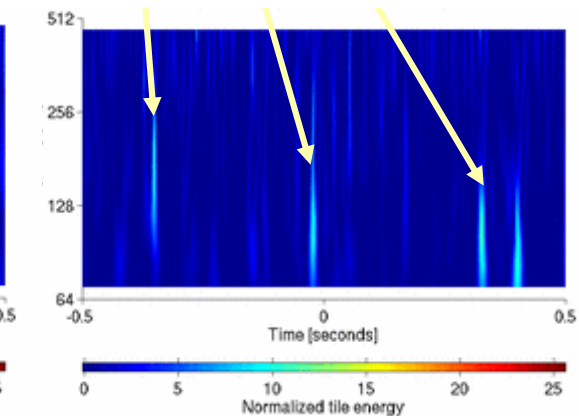
GW channel: H1



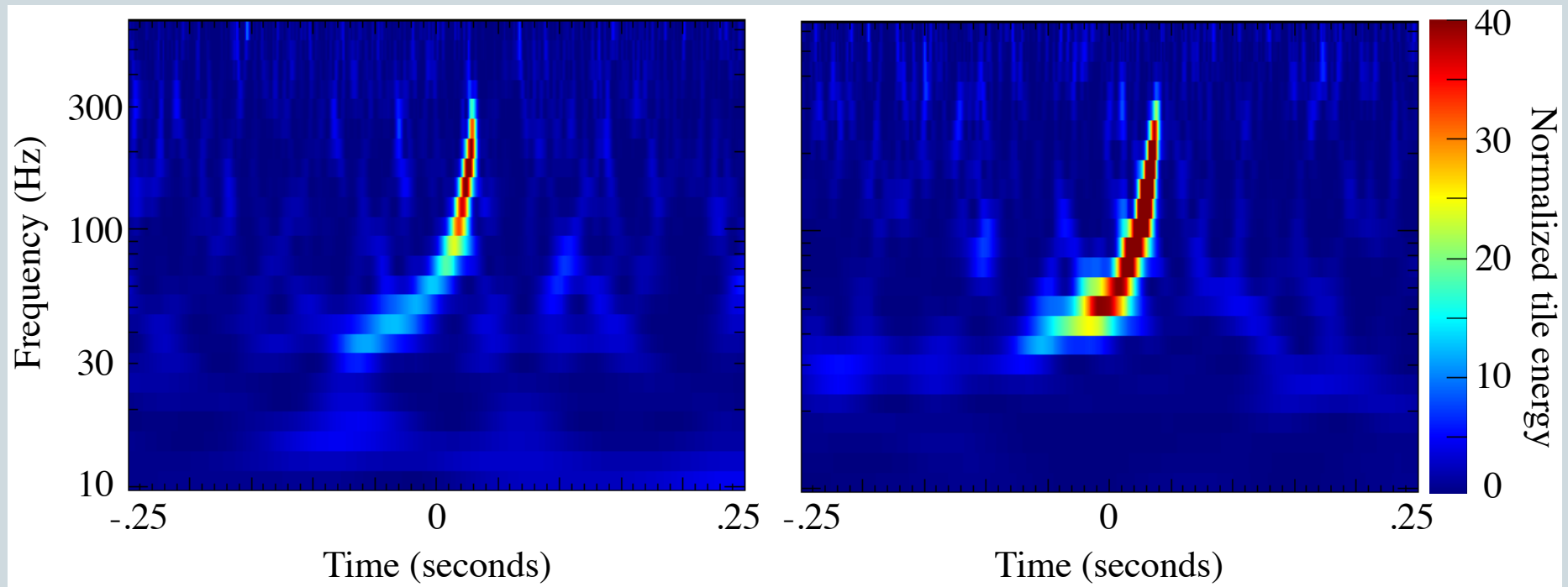
GW channel: H2



GW channel: L1



GW150914 candidate appearance



GW150914, the first GW detection



Washington, 12 February 2016

