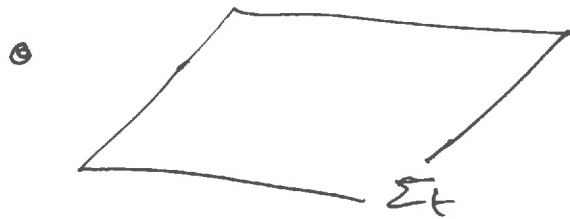




•  $M^{4D}$  manifold with metric  $g$ ; space and time are indistinguishable but — it is convenient to split time away from space



hypersurface at  $t = \text{const}$ : fully spatial

• First, fix unit normal  $\underline{n}$  such that  $\underline{n} \cdot \underline{n} = -1$  and

$$n_\mu \propto \nabla_\mu t = \Omega_\mu \quad \Rightarrow \quad \boxed{n_\mu = -\alpha \nabla_\mu t} \quad \underline{n} \cdot \underline{\tilde{\Sigma}} = \alpha^{-1} : \text{gradient of } t \text{ is a function.}$$

• Second, build projector  $\perp$  to  $\underline{n}$ :  $\underline{\gamma} = \underline{g} + \underline{n} \underline{n}$

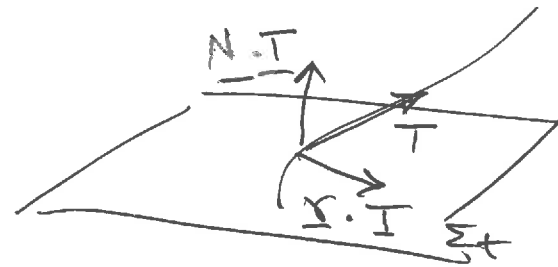
$$\gamma_{\mu\nu} = g_{\mu\nu} + n_\mu n_\nu \quad : \text{purely spatial} \quad \underline{\gamma} \cdot \underline{n} = 0$$

$\gamma$  can be used to raise/lower indices but only for purely spatial tensors.

- Thirdly, build projector along  $\underline{n}$ :  $\underline{N} = -\underline{n}\underline{n}$ ;  $N^M{}_N = -n^M n_N$

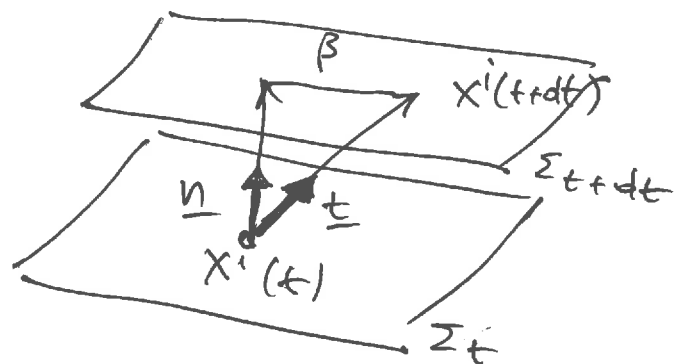
$$\underline{N} \cdot \underline{x} = 0$$

- In this way we can split any tensor in a purely spatial and in a purely time part



- $\underline{n}$  not parallel to  $\tilde{\Omega}$ ;  $\underline{n} \cdot \tilde{\Omega} = \alpha \neq \text{const} \Rightarrow$   
 need 4-vector that is timelike and parallel to  $\tilde{\Omega}$   
 $\underline{t} = \alpha \underline{n} + \underline{\beta} = (\text{time part}) + (\text{space part}) = \underline{e}_t$

$$\underline{t} \cdot \tilde{\Omega} = 1$$



$$x^i(t+dt) = x^i(t) - \beta^i(x^i, t) dt$$

In this way we can build all the metric functions

Using the lapse and shift we can write the  $tt$  and  $ti$  parts of the metric as

$$g_{tt} = \underline{t} \cdot \underline{t} = -\alpha^2 + \beta^i \beta_i \quad \left[ (\alpha n^\mu + \beta^\mu) (\alpha n_\mu + \beta_\mu) = \alpha^2 n^\mu n_\mu + 2\alpha n^\mu \beta_\mu + \beta^\mu \beta_\mu \right]$$

$$g_{ti} = \underline{t} \cdot \underline{x} = t^\mu \gamma_{\mu i} = t^\mu (g_{\mu i} + n_\mu \beta^i) = t_i =$$

$$\begin{aligned} & \stackrel{|\underline{t} \cdot \underline{n} = 0}{=} (\alpha n^\mu + \beta^\mu) g_{\mu i} = \alpha n_i + \beta^\mu g_{\mu i} \\ & = \alpha \cdot 0 + \beta^i \gamma_{ij} \\ & = + \beta_i \end{aligned}$$

so that the line element reads

$$ds^2 = -(\alpha^2 - \beta^i \beta_i) dt^2 + 2\beta_i dx^i dt + \gamma_{ij} dx^i dx^j$$

$$g_{\mu\nu} = \begin{pmatrix} -(\alpha^2 - \beta^i \beta_i) & \beta_i \\ \beta_i & \gamma_{ij} \end{pmatrix}; \quad g^{\mu\nu} = \begin{pmatrix} -1/\alpha^2 & \beta^i/\alpha^2 \\ \beta^i/\alpha^2 & \gamma^{ij} - \beta^i \beta^j/\alpha^2 \end{pmatrix}$$

As a result

$$n_\mu = (-\alpha, 0, 0, 0);$$

$$n^\mu = \frac{1}{\alpha} (1, -\beta^i)$$

$$n^i = g^{\mu i} n_\mu = g^{0i} n_0 = -\alpha \left( \frac{\beta^i}{\alpha^2} \right) = -\beta^i / \alpha$$

$$n^0 = g^{00} n_0 = -\frac{1}{\alpha^2} \cdot (-\alpha) = \frac{1}{\alpha}$$

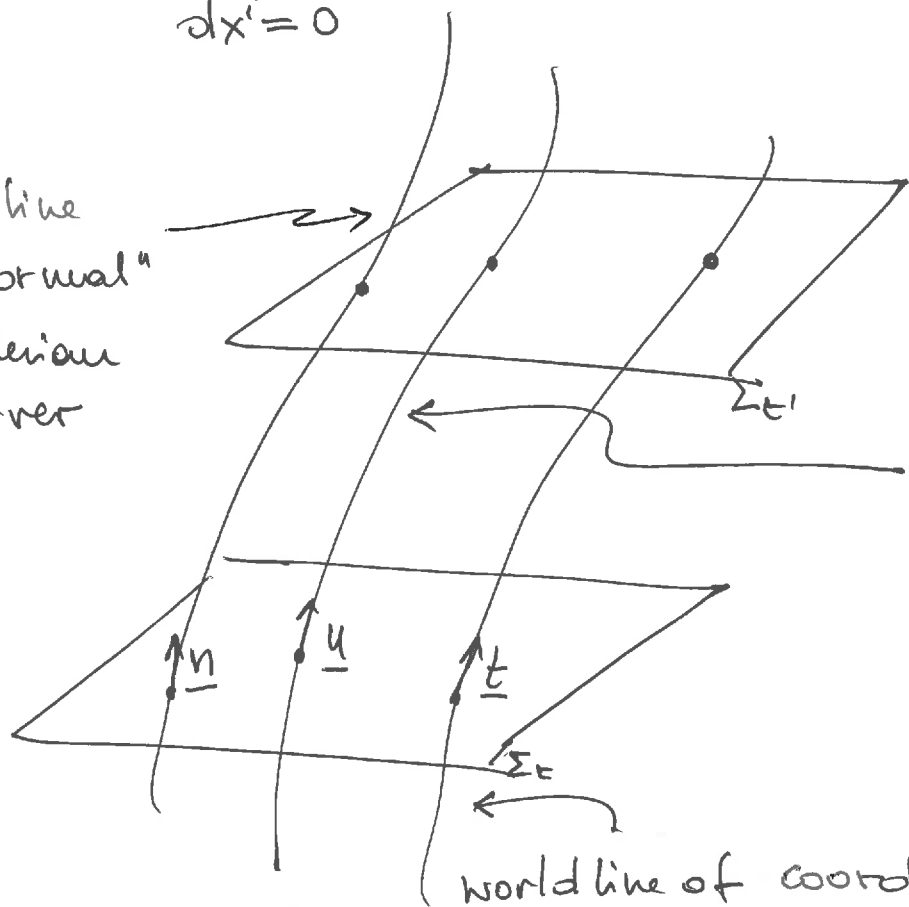
$$d\tau^2 = -ds^2 = +\alpha^2 dt^2$$

$$dx^i = 0$$

$$d\tau = \pm \alpha dt$$

The lapse function expresses the rate of change of proper time relative to the change of coordinate time.

World line of "normal" or Eulerian observer



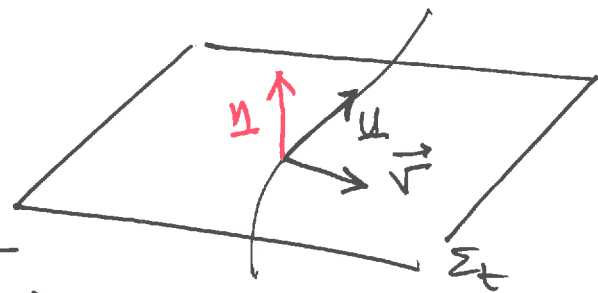
World line of fluid element with 4-velocity  $\underline{u}$

World line of coordinates

An observer with tangent vector  $\underline{n}$  is said to be a "normal" observer or "Eulerian" observer. This is the standard observer in a 3+1 split and it is relative to this observer that 3+1 quantities are measured. This is the case also for the fluid four-velocity  $\underline{u}$

$$\Downarrow: \left( \begin{array}{c} \text{spatial part} \\ \text{of } \underline{u} \end{array} \right) = \frac{(\text{projection of } \underline{u} \text{ on } \Sigma_t)}{(\text{projection of } \underline{u} \text{ along } \underline{n})}$$

$$= \frac{(\text{space})}{(\text{time})} = \frac{\delta^i_{\mu} u^{\mu}}{-u_{\mu} n^{\mu}}$$



$$W := -n_{\mu} u^{\mu} = \alpha u^t : \text{Lorentz factor}$$

$$W = (1 - v^i v_i)^{-1/2} \text{ as in special relativity} \quad (\text{Exercise})$$

$$W \rightarrow \infty \text{ for } v \rightarrow 1$$

Prove  $W = (1 - v^i v_i)^{-1/2}$

Proof

$$v^i = \frac{1}{\alpha} \left( \frac{u^i}{u^0} + \beta^i \right) ; \quad v_i = \delta_{ij} v^j = \frac{\delta_{ij}}{\alpha} \left( \frac{u^j}{u^0} + \beta^j \right)$$

$$v^i v_i = \frac{1}{\alpha^2} \left[ \left( \frac{u^i}{u^0} + \beta^i \right) \delta_{ij} \left( \frac{u^j}{u^0} + \beta^j \right) \right] = \frac{1}{\alpha^2} \left[ \delta_{ij} \frac{u^i u^j}{(u^0)^2} + 2 \beta^i \frac{u_i}{u^0} + \beta^i \beta_i \right]$$

$$-1 = u^\mu u_\mu =$$

$$= g_{00} (u^0)^2 + 2g_{0i} u^0 u^i + u^i u_i$$

$$= -(\alpha^2 - \beta^i \beta_i) (u^0)^2 + 2\beta_i u^i u^0 + u^i u_i \Rightarrow$$

$$= -1 + (\alpha u^0)^2$$

$$= \frac{1}{(\alpha u^0)^2} \left[ u^i u_i + 2\beta^i u_i u^0 + (u^0)^2 \right]$$

$$= \frac{1}{(\alpha u^0)^2} (-1 + (\alpha u^0)^2) = \frac{-1 + W^2}{W^2} \Rightarrow W^2 (1 - v^i v_i) = 1 \Rightarrow$$

$$W = (1 - v^i v_i)^{-1/2} \quad \checkmark$$

In component form:

$$v^t = 0 \quad ; \quad v^i = \frac{\delta_{\mu}^i u^{\mu}}{\alpha u^t} = \frac{(\delta_{\mu}^i + n^i n_{\mu}) u^{\mu}}{\alpha u^t} = \frac{u^i}{\alpha u^t} + \left(-\frac{\beta^i}{\alpha}\right) \frac{n_{\mu} u^{\mu}}{\alpha u^t} \quad \text{--- } \alpha u^t$$

$$= \frac{u^i}{\alpha u^t} + \frac{\beta^i}{\alpha} = \frac{1}{\alpha} \left( \frac{u^i}{u^t} + \beta^i \right)$$

$$v_t = 0 \quad ; \quad v_i = \frac{\delta_{\mu}^i u^{\mu}}{\alpha u^t} = \frac{\delta_{ij}}{\alpha} \left( \frac{u^i}{u^t} + \frac{\delta_{i0} u^0}{u^0} \right) = \frac{\delta_{ij}}{\alpha} \left( \frac{u^i}{u^t} + \beta^j \right) ;$$

In other words, using  $W = \alpha u^t$

$$v^i = \frac{u^i}{W} + \frac{\beta^i}{\alpha}$$

$$v_i = \frac{u_i}{W} = \delta_{ij} \left( \frac{u^j}{W} + \frac{\beta^j}{\alpha} \right)$$

Note that  $\underline{\delta}$  does not lower indices of  $\underline{u}$  (4D object)

I recall that in special relativity

$$v^i = \frac{u^i}{u^t} = \frac{dx^i/dt}{dt/d\tau} = \frac{dx^i}{d\tau}$$

Recap  $\underline{v}^i = \frac{u^i}{W} + \frac{\beta^i}{\alpha} = \frac{u^i}{\alpha u^t} + \frac{\beta^i}{\alpha}$

Compare with special-relativistic expression

$$\underline{v}_{SR}^i = \frac{u^i}{u^t} = \frac{u^i}{\gamma}$$

Hence, the three-velocity gains a dependence on  $\alpha$  and  $\underline{\beta}$

$$\underline{v}^i = \underline{v}_{SR}^i \text{ if } \alpha=1; \underline{\beta}=0$$

It's easy to show that

$$\underline{u} = \underbrace{W \underline{n}}_{\text{purely time part}} + \underbrace{W \underline{v}}_{\text{purely spatial part}}$$

$$\left( \begin{aligned} W &= -\underline{n} \cdot \underline{u} \Rightarrow \underline{n} W = \underline{u} \\ u^i &= W v^i - \frac{\beta^i}{\alpha} W = W n^i + W v^i \\ u^0 &= W n^0 = +\frac{1}{\alpha} W \end{aligned} \right)$$