

***Compactifying M-theory on a G2 manifold to describe/explain our world – Predictions for LHC (gluinos, winos, squarks), and dark matter***

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**CMS, Fermilab, April 2016**



# OUTLINE

- Testing theories in physics – some generalities - Testing 10/11 dimensional string/M-theories as underlying theories of *our world requires* compactification to four space-time dimensions!
- Compactifying M-theory on “G2 manifolds” to describe/ explain *our vacuum* – underlying theory - fluxless sector!
- Moduli – 4D manifestations of extra dimensions – stabilization - supersymmetry breaking – changes cosmology first 16 slides
- Technical stuff – 18-33 - quickly
- From the Planck scale to EW scale – 34-39
- LHC predictions – gluino about 1.5 TeV – also winos at LHC – but not squarks - 40-47
- Dark matter – in progress – surprising – 48
- (Little hierarchy problem – 49-51)
- Final remarks 1-5

String/M theory a powerful, very promising framework for constructing an **underlying theory** that **incorporates the Standard Models of particle physics and cosmology** and **probably addresses all the questions we hope to understand about the physical universe** – we hope for such a theory! – probably also a quantum theory of gravity

**Compactified M-theory** generically has gravity; **Yang-Mills forces like the SM**; **chiral fermions like quarks and leptons**; softly broken supersymmetry; **solutions to hierarchy problems**; **EWSB and Higgs physics**; unification; **small EDMs**; **no flavor changing problems**; **partially observable superpartner spectrum**; hidden sector DM; etc

**Simultaneously – generically**

**Argue compactified M-theory is *by far* the best motivated, and most comprehensive, extension of the SM – gets physics relevant to the LHC and Higgs and superpartners right – no ad hoc inputs or free parameters**

☺ Take it very seriously ☺

**So have to spend some time explaining derivations,  
testability of string/M theory**

**Don't have to be somewhere to test theory there**

**– E.g. no one at big bang, or dinosaur extinction, or not  
traveling faster than speed of light - but tests fully  
compelling**

**– Don't need experiments at Planck scale – always relics**

**-- If world supersymmetric, can connect EW scale data and  
Planck scale theory**

**String/M theory must be formulated in 10 (11) D to be a possible quantum theory of gravity, and obviously must be projected to 4D (“compactified”) for predictions, tests**

**Many string theorists do not know the techniques to study or evaluate compactified string/M-theories in 4 D**

**Most of what is written on this is very misleading, even by experts(!) – string theorists do not think much about it (“string theorists have temporarily given up trying to make contact with the real world” - 1999)**

**But string/M-theory's potential to provide a comprehensive underlying theory is too great to ignore it**

**String/M-theory is too important to be left to string theorists**

**Ideally theory would determine what corner of string/M theory to compactify to (heterotic? Type II? M-theory? Etc), and gauge group and matter content , and type of manifold etc – but not yet – small finite number – can try one at a time**

**Nevertheless, can address most issues – many major results do not depend on manifold, on details**

**COMPACTIFIED STRING THEORIES GIVE 4D TESTABLE RELATIVISTIC SUPERGRAVITY QUANTUM FIELD THEORIES – can calculate lots of predictions**



## There is a standard well-defined procedure to “compactify” (procedure for going to 4D)

- Choose Planck scale size manifold to compactify to
- Choose corner of string/M theory, e.g. heterotic, Type II, M-theory, etc, and gauge group, matter (e.g. SU(5)-MSSM)
- Write action, metric – project to 4D
  - Determine “superpotential”, essentially Lagrangian
  - Determine “gauge kinetic function”, metric for “gauge fields”
  - Determine “Kahler potential”, essentially metric for “scalar fields”
- Calculate potential energy, minimize it → 4D ground state

# ***Compactified string theory is analogous to Lagrangian of a system***

In all areas of physics one specifies the particular “theory” by giving the Lagrangian (Hamiltonian)

**Physical systems are described not by the Lagrangian but by *solutions* to the equations – look for set of solutions that might describe our world**

Normally find the ground state of a system, calculate energy levels and transitions

Analogous for string theory – our world corresponds to a metastable (or stable) ground state – called “vacuum”

***Curled up dimensions contain information on our world – particles and their masses, symmetries, forces, dark matter, superpartners, more – nature of compact dimensions observable indirectly via superpartner masses, etc***

# What would we need to understand and calculate to say we had an underlying theory (“final theory”) of our world?

- What are we made of? Why quarks and leptons?
- **What is light?**
- Why are there protons and nuclei and atoms? Why 3-2-1?
- **What is the origin of mass for fundamental particles (q, l, W and Z)?**
- Are the forces unified in form and strength?
- **Why are quark and charged lepton masses hierarchical?**
- Why are neutrino masses small and not hierarchical?
- **Is nature supersymmetric near the weak scale?**
- How is supersymmetry broken
- ***How is the hierarchy problem solved – stabilize hierarchy? – size of hierarchy? -  $\mu$ ?***
- Why matter asymmetry?
- **Quantum theory of gravity**
- What is an electron?
- ❑ What is dark matter? Ratio of DM to baryons?
- ❑ One and only one quark with Yukawa coupling  $\sim 1$
- ❑ Why families? Why 3?
- ❑ What is the inflaton? Why is the universe old and cold and dark?
- ❖ **Which corner of string/M-theory? Are several equivalent?**
- ❖ **Why three large dimensions?**
- ❖ Why is there a universe? More populated universes?
- ❖ Are the rules of quantum theory inevitable?
- ❖ Are the underlying laws of nature (forces, particles, etc) inevitable?
- ❖ CC problems?

➤ Answered (more or less) in compactified M-theory - **simultaneously**

❑ Addressable in compactified M-theory

❖ Can work on these

## Three new physics aspects:

- “Generic” – crucial to be predictive
- “Gravitino”- sets scale of superpartner masses
- “Moduli”
  - moduli 4D manifestation of existence of extra dimensions – generically present in all compactifications
  - New physics from compactifying
  - Describe sizes and shapes and metrics of small manifolds
  - Have definite values in vacuum – “stabilized” – if not, laws of nature time and space dependent
  - Supersymmetry breaking generates potential for all moduli, stabilizes
  - Dominate energy density of universe after inflation ends – oscillate, fall into minimum – we begin there
  - Can show lightest eigenvalue of moduli mass matrix about equal to gravitino mass
  - Decay of lightest moduli may determine matter asymmetry, and decay into DM

## GENERIC methods, results:

- Probably not a theorem (or at least not yet proved), might be avoided in special cases
- One has to work at constructing non-generic cases
- *No (or very few) adjustable parameters, no tuning*
- *Predictions NOT subject to qualitative changes from small input changes*

## GRAVITINO

- In theories with supersymmetry the graviton has a superpartner, gravitino – if supersymmetry broken, gravitino mass ( $M_{3/2}$ ) splitting from the massless graviton is determined by the form of supersymmetry breaking
- Gravitino mass sets the mass scale for all superpartners, for some dark matter

**“Naturalness” – superpartners should have masses like W,Z, top to solve hierarchy and other problems**

**“Naturalness” does suggest should have found superpartners at LHC Run 1, but naturalness is what you invoke if you *don't* have a theory – all superpartner predictions before about a decade ago were based on naturalness, not theory – some of our predictions were already made then, more recently**

**Theories need not be “natural” - Actual compactified string theories imply should *not* have found superpartners at LHC Run 1 (see below) – hierarchy problem etc still solved, in interesting ways**



# M-theory compactified on G2 manifold

## PAPERS ABOUT M-THEORY COMPACTIFICATIONS ON $G_2$ MANIFOLDS (11-7=4)

Earlier work 1995-2004 (stringy, mathematical) ; Witten 1995

- Papadopoulos, Townsend [hep-th/9506150](#), compactification on 7D manifold with  $G_2$  holonomy → resulting quantum field theory has **N=1 supersymmetry!!!**
- Acharya, [hep-th/9812205](#), **non-abelian gauge fields localized on singular 3 cycles**
- Atiyah and Witten, [hep-th/0107177](#), **analyze dynamics of M-theory on manifold of  $G_2$  holonomy with conical singularity and relations to 4D gauge theory**
- Acharya and Witten, [hep-th/0109152](#), **chiral fermions supported at points with conical singularities (quarks and leptons)**
- Witten, [hep-ph/0201018](#) – **M-theory embedding SU(5)-MSSM, solves doublet-triplet splitting in 4D supersymmetric GUT, GENERIC discrete symmetry sets  $\mu=0$**
- Beasley and Witten, [hep-th/0203061](#), **generic Kahler form**
- Friedmann and Witten, [hep-th/0211269](#), **SU(5) MSSM, scales – Newton's constant, GUT scale, proton decay – no susy breaking**
- Lukas, Morris [hep-th/0305078](#), **generic gauge kinetic function**
- Acharya and Gukov, [Physics Reports, 392\(2004\)2003](#)

Particles and forces!

**Basic framework established – powerful, rather complete**

➤ Acharya and I (and students, postdocs, collaborators) began there

# We make a few discrete assumptions, calculate

- Compactify **M-Theory** on manifold with  **$G_2$**  holonomy **in fluxless sector** – *well motivated and technically robust*
- Compactify to gauge matter group **SU(5)-MSSM** – can try others, one at a time
- Use generic Kahler potential and generic gauge kinetic function
- Assume needed singular mathematical manifolds exist – considerable progress recently – Simons Center workshops, Acharya, Simon Donaldson et al, etc
- **CC issues not relevant** - solving it doesn't help learn our vacuum, and not solving it doesn't stop learning our vacuum

**We started in 2005 – since LHC coming, focused on moduli stabilization, supersymmetry breaking, etc → LHC physics, Higgs physics, dark matter etc**

[**Acharya**, Bobkov, GK, Piyush Kumar, Kuflik, Shao, Watson, Lu, Zheng, S. Ellis – over 20 papers, over 500 arXiv pages]

- **Indeed we showed that in M theory supersymmetry automatically was spontaneously broken via gaugino and chiral fermion condensation**
- **Simultaneously moduli stabilized**, in unique de Sitter vacuum for given manifold
- **Calculated supersymmetry soft-breaking Lagrangian → radiative electroweak symmetry breaking, Higgs boson** – precise prediction of  $M_h/M_Z$  and  $h$  decays (in decoupling sector) – **gluino and wino masses, etc**

Get 4D effective supersymmetric field theory – in usual case coefficients of all operators are independent, so many coefficients – here all coefficients **DETERMINED**, calculable and connected

**NO adjustable parameters – sometimes coefficient of term hard to calculate, so constrained parameter, e.g. of order 1 but could be off  $\leq$  factor 2**

Generically two hidden sector 3D submanifolds do not intersect in a 7D space, so no light matter fields charged under both SM gauge group and hidden sector gauge groups  $\rightarrow$  supersymmetry breaking *generically gravity mediated* in these vacua

Technical aspects:

# MODULI STABILIZATION

- All moduli geometric, equivalent – All  $G_2$  moduli fields  $s_i$  have axionic partners  $t_i$  which have a shift symmetry in the absence of fluxes (different from heterotic or IIB) – such symmetries can only be broken by non-perturbative effects
- So in zero-flux sector only contributions to superpotential are non-perturbative, from “strong dynamics” (e.g. gaugino condensation or instantons) – focus on former
- In M theory superpotential and gauge kinetic function depend on all the moduli– all moduli on equal footing
  - so only need one term in  $W$  to stabilize all moduli
  - in practice use at least two to be sure supergravity approximation good numerically
  - **not racetrack**
- The hidden sector gaugino condensation produces an effective potential that stabilizes all moduli

A set of Kahler potentials, consistent with  $G_2$  holonomy and known to describe some explicit examples, was given by Beasley-Witten [th/0203061](#); Acharya, Denef, Valandro [th/0502060](#), with

$$K = -3 \ln(4\pi^{1/3} V_X)$$

$$V_X = \prod_{i=1}^N s_i^{a_i}, \quad \text{with} \quad \sum_{i=1}^N \underline{a_i} = 7/3$$

We assume we can use this.



The gauge kinetic functions here are integer linear combinations of all the moduli (Lukas, Morris th/0305078),

$$f_k = \sum_{i=1}^N \underline{N_i^k} z_i .$$

Focus on the (well-motivated) case where two hidden sector gauge kinetic functions are equal (the corresponding three-cycles are in the same homology class)]

Include massless hidden sector chiral fermion quark states  $Q$  with  $N_c$  colors,  $N_f$  flavors,  $N_f < N_c$  -- then (Affleck, Dine, Seiberg PRL 51(1983)1026, Seiberg hep-th/9402044, hep-th/9309335, Lebedev, Nilles, Ratz th/0603047),  $a = 2/(N_c - N_f)$

$$W = A_1 e^{i \frac{2\pi}{N_c - N_f} \sum_{i=1}^N N_i^{(1)} z_i} \det(Q\tilde{Q})^{-\frac{1}{N_c - N_f}} = A_1 \boxed{\phi^a} e^{i b_1 f_1}$$

and define an effective meson field

$$\phi \equiv \left( \det(Q\tilde{Q}) \right)^{1/2} = \phi_0 e^{i\theta}$$

- For pure  $SU(Q)$  super Yang-Mills hidden sector, non-perturbative dynamics generates an effective moduli superpotential of form  $W=AM_{pl}^3 e^{i2\pi bf}$  where  $f$  is the hidden sector gauge kinetic function  $f=\sum N_i z_i$  and  $b=1/Q$
- Integers  $N_i$  determined by homology class of the 3-cycle
- Hidden sectors with  $SU(P+1)$  gauge group with chiral charged matter, which arises from isolated conical singularities in the G2 manifold, also are included – superpotential from Seiberg et al
- Such a superpotential will stabilize all moduli, in de Sitter space
- Get unique de Sitter vacuum for a given manifold, and sector with  $Q-P=3$  has no high scale solutions, only  $M_{3/2} \approx 50$  TeV for number of moduli larger than about 60

# SUPERPOTENTIAL

$$W = \sum_{k=1}^M A_k e^{ib_k f_k} \sim \sum_k e^{ib_k \sum N_i^k z_i}$$

gauge kinetic function

Complex  
moduli

Keep two terms – enough to find solutions with good properties such as being in supergravity regime, simple enough to do most calculations semi-analytically (as well as numerically) – check some things with more terms numerically

Imagine expanding exponential – all terms get interactions

$b_k = 2\pi/c_k$  where  $c_k$  are dual Coxeter numbers of hidden sector gauge groups ---  
 $A_k$  are constants of order unity, and depend on threshold corrections to gauge couplings, some computed by Friedmann and Witten

The microscopic constants  $a_i, b_k, A_k, N_i^k$  are determined for a given  $G_2$  manifold (but not yet known for all relevant ones) --they completely characterize the vacua – not dependent on moduli

Finally mostly work with

$$W = A_1 \phi^a e^{ib_1 f} + A_2 e^{ib_2 f}$$

$$K = -3 \ln(4\pi^{1/3} V_X) + \phi \bar{\phi}$$

Can often get semi-analytic forms, and approximations good

We also looked at chiral families in both hidden sectors, more chiral families in each – no changes in qualitative results (in paper)

The N=1 SUGRA scalar potential is then given by:

$$V_X = V_7$$

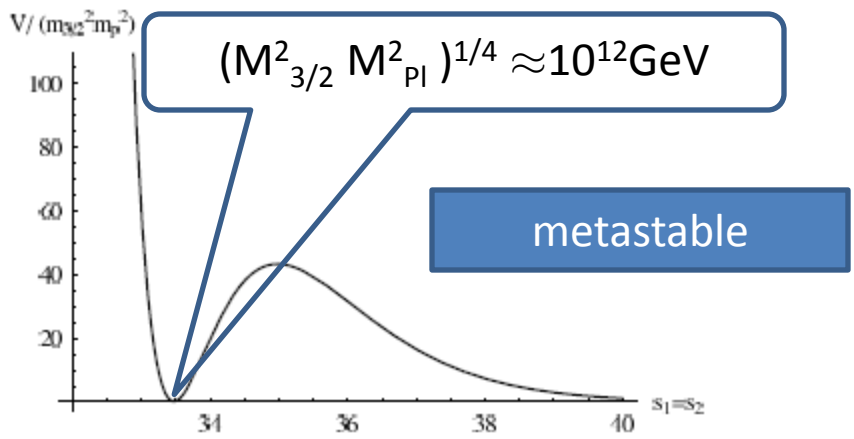
$$\begin{aligned}
 V = & \frac{e^{\phi_0^2}}{48\pi V_X^3} [(b_1^2 A_1^2 \phi_0^{2a} e^{-2b_1 \vec{\nu} \cdot \vec{a}} + b_2^2 A_2^2 e^{-2b_2 \vec{\nu} \cdot \vec{a}} + 2b_1 b_2 A_1 A_2 \phi_0^a e^{-(b_1+b_2) \vec{\nu} \cdot \vec{a}} \cos((b_1 - b_2) \vec{N} \cdot \vec{t} + a\theta)) \\
 & \times \sum_{i=1}^N a_i (\nu_i)^2 + 3(\vec{\nu} \cdot \vec{a})(b_1 A_1^2 \phi_0^{2a} e^{-2b_1 \vec{\nu} \cdot \vec{a}} + b_2 A_2^2 e^{-2b_2 \vec{\nu} \cdot \vec{a}} + (b_1 + b_2) A_1 A_2 \phi_0^a e^{-(b_1+b_2) \vec{\nu} \cdot \vec{a}} \\
 & \times \cos((b_1 - b_2) \vec{N} \cdot \vec{t} + a\theta)) + 3(A_1^2 \phi_0^{2a} e^{-2b_1 \vec{\nu} \cdot \vec{a}} + A_2^2 e^{-2b_2 \vec{\nu} \cdot \vec{a}} + 2A_1 A_2 \phi_0^a e^{-(b_1+b_2) \vec{\nu} \cdot \vec{a}} \\
 & \times \cos((b_1 - b_2) \vec{N} \cdot \vec{t} + a\theta)) + \frac{3}{4} \phi_0^2 (A_1^2 \phi_0^{2a} \left( \frac{a}{\phi_0^2} + 1 \right)^2 e^{-2b_1 \vec{\nu} \cdot \vec{a}} + A_2^2 e^{-2b_2 \vec{\nu} \cdot \vec{a}} \\
 & + 2A_1 A_2 \phi_0^a \left( \frac{a}{\phi_0^2} + 1 \right) e^{-(b_1+b_2) \vec{\nu} \cdot \vec{a}} \cos((b_1 - b_2) \vec{N} \cdot \vec{t} + a\theta)]. \tag{101}
 \end{aligned}$$

**Minimize**

# Results:

$$m^2_{\text{scalar}} \approx M^2_{3/2} + V_0 + \text{small corrections calculable from } W, K, f$$

$V_0$  is value of potential at minimum  $\approx$  cosmological constant, set it to be small for any particular vacuum



*So scalar masses essentially equal to gravitino mass*

## DE SITTER VACUUM, GAUGINO MASSES SUPRESSED

$$M_{1/2} \sim K_{mn} F_m \partial_n f_{SM}$$

Standard Model  
gauge kinetic  
function

- $f_{SM}$  doesn't depend on chiral fermions, whose F-term gives the largest contribution to supersymmetry breaking
- $F_{\text{chiral fermion}} \sim V_7$  but  $F_{\text{moduli}} \sim V_3$ ,  $V_7 \gg V_3$
- matter Kahler potential does not enter, so results more reliable
- moduli dependence is entirely in Volume factors, so same for all G2 manifolds for tree level gaugino masses



## □ Including $\mu$ parameter in string theory ( $W = \mu H_u H_d + \dots$ so $\mu \sim 10^{16}$ GeV ?)

- Normally  $\mu$  and  $\tan\beta$  treated as parameters, constrained to get EWSB
- **Ultimately want to derive them from first principles**
- If  $\mu$  in  $W$  then it should be of order string scale
- **Need symmetry to set  $\mu=0$**
- **Witten, hep-ph/0201018 – found generic discrete symmetry for  $G_2$  compactifications, closely connected to doublet-triplet splitting problem proton lifetime**
- **Unbroken discrete symmetry so  $\mu \equiv 0$**  – but when moduli are stabilized the effects generally not invariant so in M-theory with moduli stabilized the symmetry is broken
- **$\mu$  proportional to  $M_{3/2}$  since  $\mu \rightarrow 0$  if susy unbroken**
- **$\mu$  proportional to moduli vev since  $\mu \rightarrow 0$  if moduli not stabilized**
- Stabilization led to moduli vevs/ $M_{pl} \lesssim 0.1$
- **So finally expect  $\mu \leq 0.1 M_{3/2}$**
- But answer for residual symmetry not known – interesting mathematics – value of  $\mu$  depends on manifold – maybe  $0.04 M_{3/2}$

arXiv:1102.0556, Acharya, Kane,  
Kuflik, Lu

# MAIN RESULTS, PREDICTIONS FOR M-THEORY SO FAR, and in progress – ONE THEORY

- **Moduli stabilized** – vevs calculable and  $\lesssim 1/10 M_{pl}$ , masses multi TeV ✓
- Calculate gravitino mass approximately, from Planck scale  $\sim 50$  TeV
- **Scalars heavy** (squarks, higgs sector, sleptons)  $\sim$  gravitino mass (2006) **PREDICTION, LHC**
- **Gaugino masses suppressed** (by volume ratios),  $\sim$  factor 40 **PREDICTION, LHC**
- **Hierarchy problem solved** ✓
- **Non-thermal cosmological history** via late time moduli decay (before BBN) **PREDICTION**
- **Moduli decay can provide ratio of baryogenesis and DM** **PREDICTION**
- **Axions stabilized, give solution to strong CP problem, spectrum of axion masses** ✓
- **Anticipated Higgs boson mass and BR (SM-like) before data** **PREDICTION** ✓
- **SM quark and lepton charges, Yang-Mills 3-2-1 forces, parity violation, generic**
- **Gauge coupling unification, proton decay all right**
- **No flavor problem, weak CPV ok**
- **EDMs calculable, smallness explained (could have been wrong)** **PREDICTION** ✓
- $\mu \approx 2-3$  TeV – included in theory, approximately calculable
- $\tan\beta \approx 5-7$  **PREDICTION**
- **LHC predictions** – gluinos  $\sim 1.5$  TeV, 3<sup>rd</sup> family decays enhanced  
-- wino, bino  $\sim \frac{1}{2}$  TeV
- **Need future collider for higgsinos, scalars** – not at LHC **PREDICTION**
- **Hidden sector DM, under study** – **LSP decays, LSP generically never dark matter**

ALL FOLLOW FROM few DISCRETE ASSUMPTIONS – no free parameters – all **SIMULTANEOUS**

# Scales

Planck scale

GUT  $\sim 2 \times 10^{16}$

String, KK, etc

$\Lambda \approx 10^{14}$  GeV

gaugino, chiral fermion condensation, F-terms  $\neq 0$  (susy broken)

$$\Lambda \approx \exp\{-2\pi V_3/3Q\} M_{Pl}/V_7^{1/2}$$

( $V_3 \sim Q$  so not sensitive)

Supersymmetry breaking dynamical, automatic!

Hierarchy problem solved

$M_{3/2} \sim 50$  TeV

Gravitino mass (so squarks heavy)

$M_{3/2} = e^{K/2} W/M_{Pl}^2, W \sim \Lambda^3$

$\mu \sim 2-3$  TeV

Gaugino mass suppression

$$M_{1/2} \sim F_{mod} \frac{\partial f_{vis}}{\partial F_{mod}} + F_{ChiFerm} \frac{\partial f_{vis}}{\partial F_{ChiFerm}}$$

and  $F_{mod}/F_{ChiFerm} \sim V_3/V_7 \ll 1$

gluino  $\sim 1.5$  TeV, wino, bino 0.5 TeV

TeV scale

REWSB

$\mu \approx \langle \text{mod} \rangle M_{3/2}$  (Witten+mod stabilization)  $\sim$  few TeV

$M_{Hu} \sim f_{M0}(t) M_0^2 - f_{A0}(t) A_0^2 \ll M_{3/2}$  ( $f_{M0} \approx f_{A0}; A_0 \gtrsim M_0$ )

# Qualitative gravitino and gluino masses

$$W \sim \Lambda^3$$

$$M_{3/2} = \frac{e^{K/2} W}{M_{Pl}^2} \approx \left( \frac{\Lambda}{M_{Pl}} \right)^3 \frac{1}{V_3} M_{Pl}$$

$$M_{gluino} \approx \left( \frac{\Lambda}{M_{Pl}} \right)^3 \frac{1}{V_3} \frac{V_3}{V_7} M_{Pl}$$

+ anomaly  
mediation terms

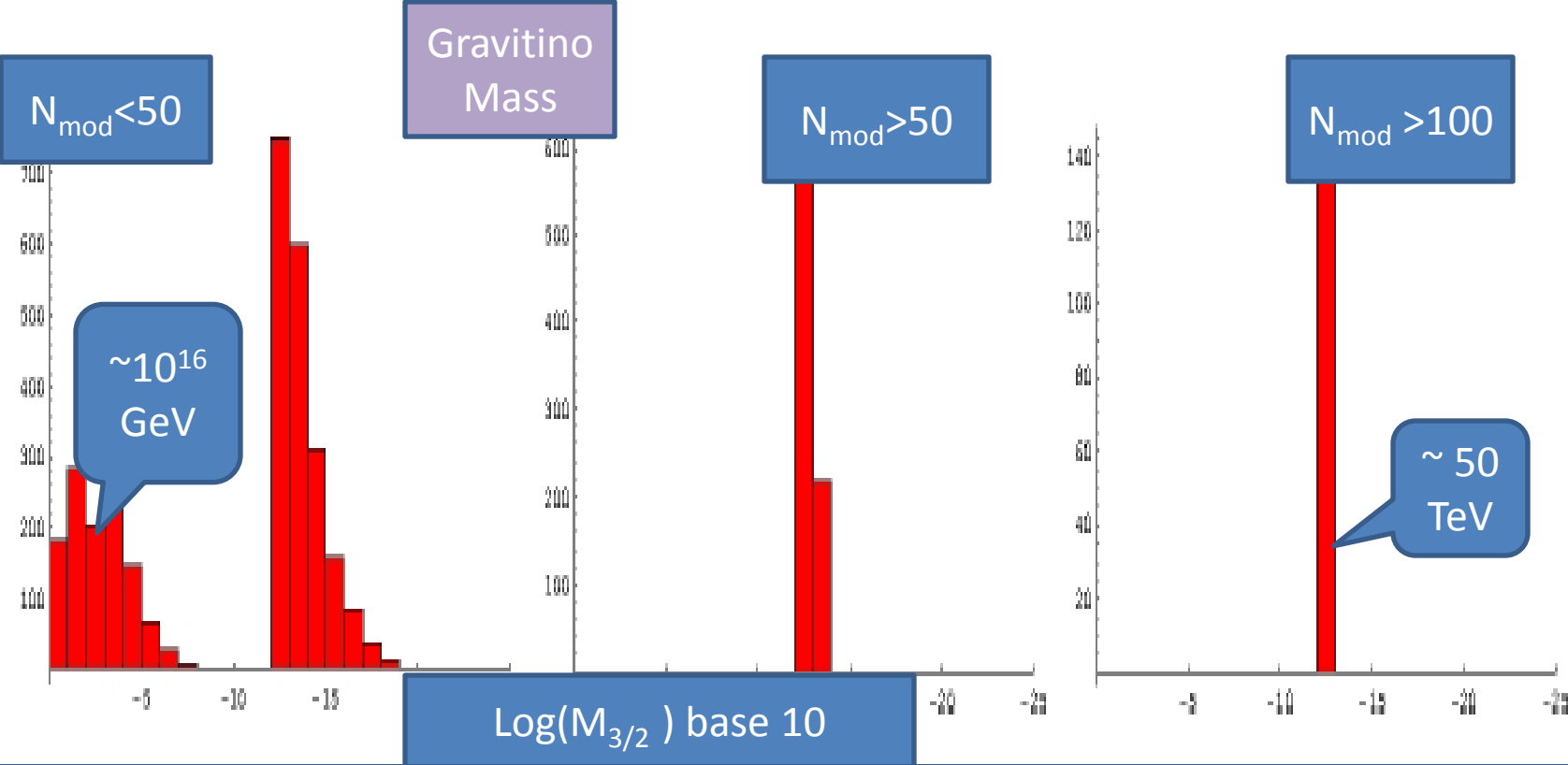
[Acharya, Bobkov calculated cross term between matter and moduli Kahler potential – coefficient C of order 1 but C hard to calculate – we include that term in careful calculation of gaugino masses – use Higgs mass to help fix  $C \approx 1/2$

[ $\Phi$  visible sector matter,  $\varphi$  moduli,  $\kappa_{\alpha\beta}$  Kahler metric]

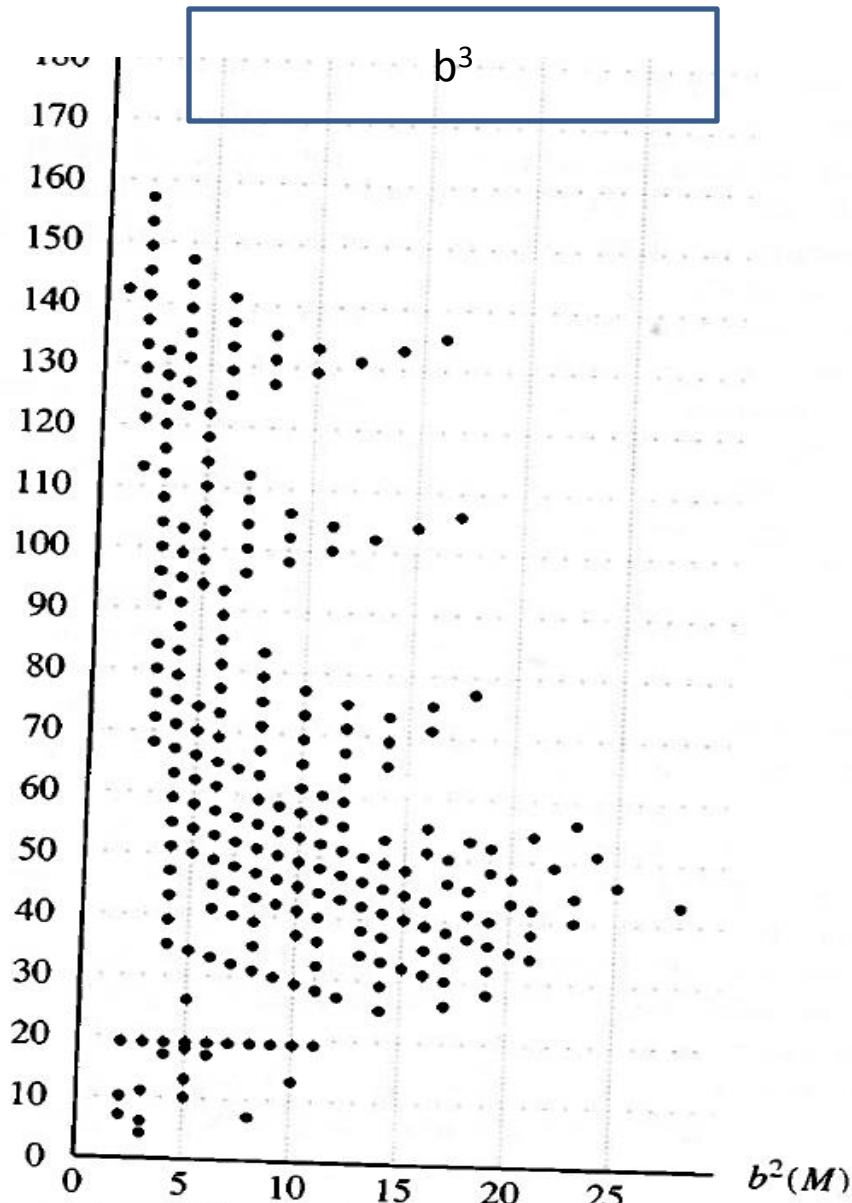
$$K = -3\ln(4\pi^{1/3}V_7) + \kappa_{\alpha\beta} \frac{\Phi_\alpha^\dagger \Phi_\beta}{V_7} + \frac{\bar{\varphi}\varphi}{V_7} + \frac{C}{3} \kappa_{\alpha\beta} \frac{\bar{\varphi}\varphi}{V_7} \frac{\Phi_\alpha^\dagger \Phi_\beta}{V_7}$$

Use  $M_h$  value to pin down  $M_{3/2}$  rather precisely,  $M_{3/2} = 35$  TeV]

# Hierarchy problem solved IF number of moduli ( $b_3$ ) large enough!



This is after setting potential to zero at minimum – do not have to separately set  $V_0$  to zero and also  $M_{3/2}$  to TeVs – does not happen in other corners !



Dominic Joyce,  
 “Compact Manifolds  
 with Special  
 Holonomy” – graph  
 for non-singular  
 manifolds

**Figure 12.3.** Betti numbers  $(b^2(M), b^3(M))$  of compact, simply-connected 7-manifolds with holonomy  $G_2$

## HIGGS MASS, DECAYS

Two Higgs doublets in supersymmetry – large scalar terms in soft-breaking Lagrangian ( $M_{Hu}, M_{Hd}$ ) plus radiative electroweak symmetry breaking imply one light Higgs boson and four heavy ones, “decoupling sector”

Calculate ratio  $M_{\text{higgs}}/M_Z$  – determined by “ $\lambda$ ” of Higgs potential – write theory at string scale – do “renormalization group running” down to electroweak scale, known through three loops with heavy scalars – use “match and run”

Compactified M-theory (with generic gauge kinetic function and kahler potential) anticipated  $M_{\text{higgs}}=126.4 \text{ GeV}$  summer 2011, before data – predicted all decay branching ratios would be within few per cent of Standard Model ones (as observed) – BR not a mystery

Electroweak scale spread of about  $\pm 1.2 \text{ GeV}$  purely because top quark yukawa and  $\alpha_s$  enter RGE running from high scale

Higgs data exactly as expected from compactified M-theory  
MSSM decoupling sector and electroweak symmetry breaking



Here is where supersymmetry is “hiding” at LHC

# LHC

Squark masses  $\sim$  gravitino mass  $\sim$  few tens of TeV

**GAUGINO MASSES  $\sim$  TeV**

arXiv:1408.1961 [Sebastian Ellis, GK, Bob Zheng]

$M_{\text{gluino}} \approx 1.5 \text{ TeV},$

$M_{\text{bino}} \approx 450 \text{ GeV},$

$M_{\text{wino}} \approx 614 \text{ GeV}$

all consistent with current data

Lesson from (compactified M-)theory: should not have expected superpartners at LHC Run 1

$\sigma_{\text{gluino}} \approx 12 \text{ fb}$  (smaller because squarks heavy),

$\sigma_{\text{wino pairs}} \approx 15 \text{ fb}$

For 1.5 TeV,  $3\sigma$  gluino signal probably needs  $\sim 45 \text{ fb}^{-1}$   
because of backgrounds (top pairs about 300 times gluino pairs)

Any bets?

## 3 and only 3 channels at LHC:

$$pp \rightarrow gg$$

$$pp \rightarrow W^+ W^-$$

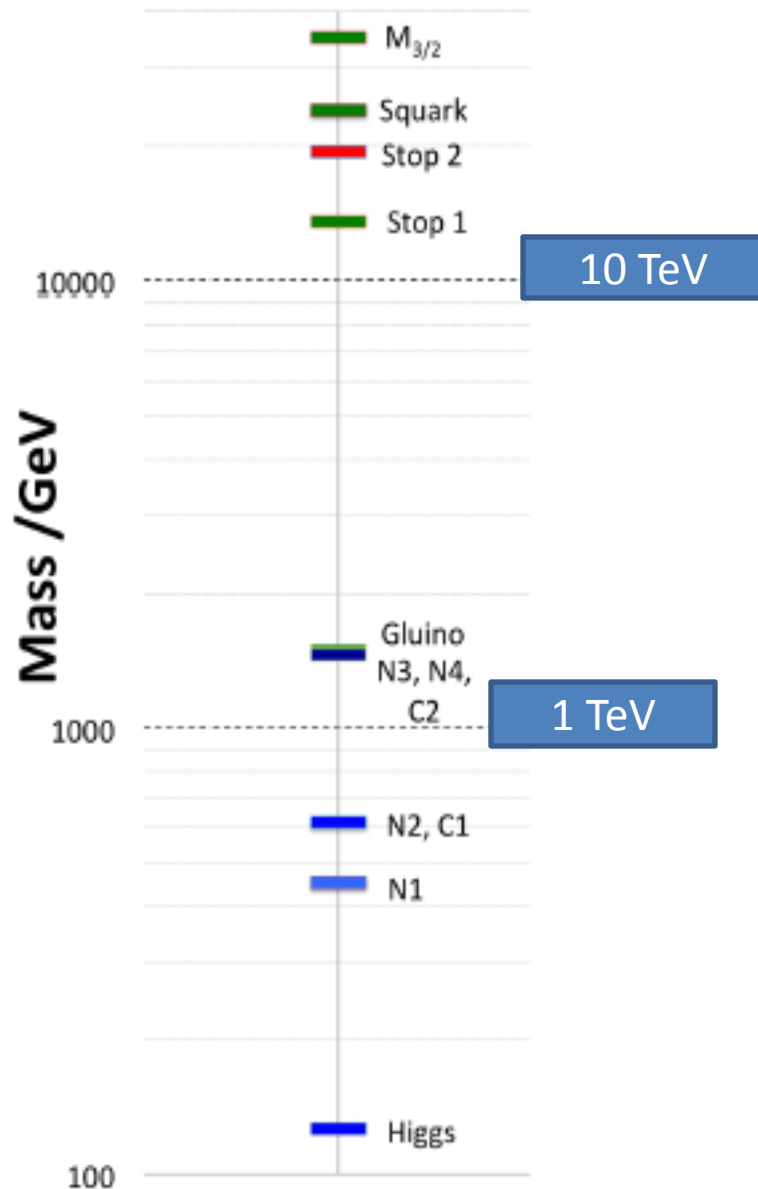
$$pp \rightarrow W^\pm W^0$$

$$M(W^+) = M(W^0) = 614 \text{ GeV}$$

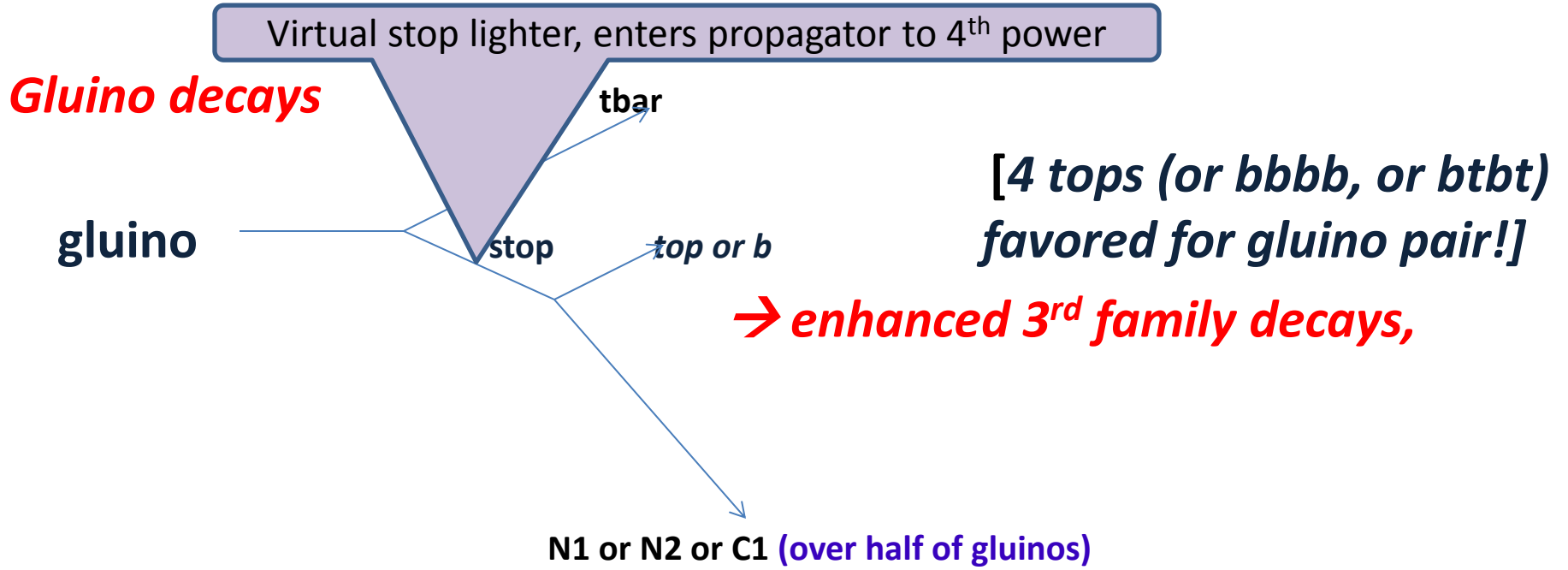
$$M(LSP) = 450 \text{ GeV}$$

$$W^\pm \rightarrow W^\pm + LSP \sim 100\%$$

$$W^0 \rightarrow h + LSP \sim 98\%$$



| Particle          | Mass (GeV)         |
|-------------------|--------------------|
| $m_0$             | 24200              |
| $M_{3/2}$         | 35000              |
| $\tilde{q}_{L,R}$ | 24000              |
| $\tilde{t}_2$     | 19300              |
| $\tilde{t}_1$     | 13500              |
| $b_2$             | 23900              |
| $\tilde{b}_1$     | 19300              |
| $\tilde{g}$       | 1500               |
| $\chi_1^0$        | 450                |
| $\chi_2^0$        | 614                |
| $\chi_3^0$        | 1460               |
| $\chi_4^0$        | 1460               |
| $\chi_{1\pm}$     | 614                |
| $\chi_{2\pm}$     | 1460               |
| $h$               | 125.2 <sup>2</sup> |



**Glino lifetime  $\sim 10^{-19}$  sec, decays in beam pipe**

**Glino decays flavor-violating: 3<sup>rd</sup> family / (1<sup>st</sup> + 2<sup>nd</sup>)  $\approx 1.2$  (naively 0.5)**

For heavy squarks,  $\sigma(\text{gluinos, 13 TeV}) / \sigma(\text{gluinos, 8 TeV}) \approx 30-45$  for 1.5 TeV gluino

### Glauino BR

| Decay   | BR (%) |
|---|--------|
| $\tilde{g} \rightarrow \chi_1^+ q_{1,2} \bar{q}_{1,2}$  | 25     |
| $\tilde{g} \rightarrow \chi_1^\pm b \bar{t}, t \bar{b}$ | 23     |
| $\tilde{g} \rightarrow \chi_1^0 t \bar{t}$              | 20     |
| $\tilde{g} \rightarrow \chi_2^0 q_{1,2} \bar{q}_{1,2}$  | 12     |
| $\tilde{g} \rightarrow \chi_1^0 q_{1,2} \bar{q}_{1,2}$  | 8      |
| $\tilde{g} \rightarrow \chi_2^0 b \bar{b}$              | 7      |
| $\tilde{g} \rightarrow \chi_2^0 t \bar{t}$              | 4      |
| $\tilde{g} \rightarrow \chi_1^0 b \bar{b}$              | 1      |

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### Neutralino BR

| Decay                                   | BR (%) |
|---|--------|
| $\chi_4^0 \rightarrow \chi_1^\pm W^\mp$ | 60     |
| $\chi_4^0 \rightarrow \chi_2^0 h$       | 27     |
| $\chi_4^0 \rightarrow \chi_1^0 h$       | 8      |
| $\chi_4^0 \rightarrow \chi_2^0 Z$       | 4      |
| $\chi_4^0 \rightarrow \chi_1^0 Z$       | 2      |
| $\chi_3^0 \rightarrow \chi_1^\pm W^\mp$ | 60     |
| $\chi_3^0 \rightarrow \chi_2^0 Z$       | 26     |
| $\chi_3^0 \rightarrow \chi_1^0 Z$       | 8      |
| $\chi_3^0 \rightarrow \chi_2^0 h$       | 4      |
| $\chi_3^0 \rightarrow \chi_1^0 h$       | 2      |
| $\chi_2^0 \rightarrow \chi_1^0 h$       | 98     |
| $\chi_2^0 \rightarrow \chi_1^0 Z$       | 2      |

### Chargino BR

| Decay                                   | BR (%) |
|---|--------|
| $\chi_2^\pm \rightarrow \chi_1^\pm h$   | 31     |
| $\chi_2^\pm \rightarrow \chi_1^\pm Z$   | 30     |
| $\chi_2^\pm \rightarrow \chi_2^0 W^\pm$ | 30     |
| $\chi_2^\pm \rightarrow \chi_1^0 W^\pm$ | 9      |
| $\chi_1^\pm \rightarrow \chi_1^0 W^\pm$ | 100    |

Future colliders – 100 TeV--gluino + squark associated production

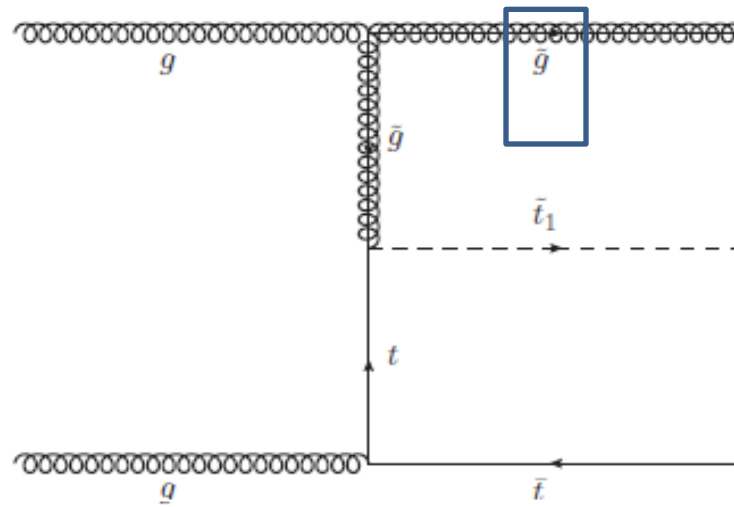


Figure 4: Dominant Feynman graph for stop associated production by gluon splitting.

We computed using MadGraph5 [39] the production cross-sections to leading order for  $t\bar{t}$  channels for both  $\sqrt{s} = 50$  TeV and 100 TeV. The results are tabulated below, including number of events  $N$  expected given  $3000 \text{ fb}^{-1}$  of data.

| Channel                                      | $\sigma_{50 \text{ TeV}}$ (fb) | $N_{50 \text{ TeV}}$ | $\sigma_{100 \text{ TeV}}$ (fb) | $N_{100 \text{ TeV}}$ |
|--|--------------------------------|----------------------|---------------------------------|-----------------------|
| $pp \rightarrow t\bar{t}_1\tilde{g}$         | $7.1 \times 10^{-5}$           | 0                    | $1.6 \times 10^{-2}$            | 47                    |
| $pp \rightarrow b\bar{b}_1\tilde{g}$         | $2.6 \times 10^{-6}$           | 0                    | $3.0 \times 10^{-3}$            | 9                     |
| $pp \rightarrow \tilde{q}_{1(L,R)}\tilde{g}$ | $3.2 \times 10^{-4}$           | 1                    | $3.0 \times 10^{-1}$            | 900                   |
| $pp \rightarrow \chi_3^0\chi_4^0$            | $9.2 \times 10^{-1}$           | 2800                 | 3.4                             | 10200                 |
| $pp \rightarrow \chi_3^0\chi_2^\pm$          | 1.8                            | 5400                 | 6.4                             | 19200                 |
| $pp \rightarrow \chi_4^0\chi_2^\pm$          | 1.8                            | 5400                 | 6.4                             | 19200                 |
| $pp \rightarrow \chi_2^\pm\chi_2^\mp$        | 1.0                            | 3000                 | 3.7                             | 11100                 |

**Glينو, wino, bino mass predictions are generic and robust – not just “a little above current limits” – clear to any knowledgeable person who goes through derivation**

**Qualitatively:**

- **Compactification, RGE running down**
- **F-terms  $\neq 0$  from hidden sector gaugino and chiral fermion condensation, so supersymmetry broken – largest gauge groups on 3-cycles run fastest  $\rightarrow$  scale  $\approx 10^{14}$  GeV**  
[ $\Lambda \approx (M_{\text{pl}}/V_7) \exp(-2\pi V_3/3Q) \approx 10^{14} \text{ GeV}$ ]
- **Then calculate gravitino mass  $\approx 40$  TeV [ $W \sim \Lambda^3/M_{\text{pl}}^3$ ,  $M_{3/2} \approx e^{K/2} W/M_{\text{pl}}^2$ ]**
- **Gaugino masses automatically suppressed to  $\sim$  TeV since largest susy-breaking source of mass absent,  $V_3/V_7 \sim 1/40$**   
 $\rightarrow$  gluino mass  $\sim 1.5$  TeV ( $\pm 10\text{-}15\%$ )
- **Gluino cross section  $\approx 12$  fb - top pair background large – note limits weaker for heavy squarks and for realistic decays**

## HIDDEN SECTOR DARK MATTER – in progress – predictions and tests

[Acharya, Sebastian Ellis, GK, Brent Nelson, Malcolm Perry, Bob Zheng]

- In M-theory, curled up 7D space has 3D submanifolds (“3-cycles”) that generically have (orbifold) singularities and therefore have particles in gauge groups  $\sim 100$  submanifolds (3<sup>rd</sup> Betti number) – **we live on one, “visible sector”**
- **Supersymmetry breaking due to ones with large gauge groups**
- **Gravitational interactions, same gravitino and moduli for all**
- **Other hidden sectors have their own matter, some stable and DM candidates – can calculate spectra, relic densities**
- **Calculations underway: already published general relic density calculations with a non-thermal cosmological history, arXiv:1502.05406 (Acharya, GK, Nelson, Zheng)**
- **Now analyzing actual hidden sectors systematically for M-Theory**
- **Examples of stable relics exist, with relic density of order what is observed – e.g. M-theory case  $U(1)^3$ , DM mass  $\sim 10$  MeV**
- **Generically, LSP decays to lighter hidden sector states in some hidden sector – LSP “never” dark matter**

**U(1)’s generic explicitly and via larger gauge groups breaking - kinetic mixing portals generic (other portals too) – light gauginos generic – light chiral fermions generic via hierarchical couplings**

**IT IS NOT GENERIC TO NOT HAVE SIZABLE KINETIC MIXING AND LIGHT HIDDEN SECTOR STATES**



# LITTLE HIERARCHY ~ 2 TEV, NOT 40 TEV – MAYBE EVEN SOLVED

-- derive  $\tan\beta$  too

Usual EWSB conditions [so higgs potential minimum away from origin]:

$$M_Z^2 = -2\mu^2 + 2(M_{Hd}^2 - M_{Hu}^2 \tan^2\beta)/\tan^2\beta = -2\mu^2 + 2M_{Hd}^2/\tan^2\beta - 2M_{Hu}^2$$
$$2B\mu = \sin 2\beta (M_{Hu}^2 + M_{Hd}^2 + 2\mu^2)$$

$M_{Hu}^2$  runs to be negative,  $M_{Hd}^2$  and B don't run much,  $\mu$  suppressed,  
 $\sin 2\beta \approx 2/\tan\beta$

If no  $\mu$  from superpotential, and visible sector Kahler metric and Higgs bilinear coefficient independent of meson field, and if  $F_{\text{mod}} \ll F_\phi$  then B (high scale)  $\approx 2M_{3/2}$  – recall  $\mu < 0.1M_{3/2}$

$$\rightarrow \tan\beta \approx M_{Hd}^2/B\mu \approx M_{3/2}^2/B\mu \rightarrow \tan\beta \approx M_{3/2}/2\mu (\sim 6)$$

➤  $\mu, |M_{Hu}| \sim 2 \text{ TeV}$ , so little hierarchy  $\sim 10\text{-}20$ , not  $\sim M_{3/2}/M_Z$

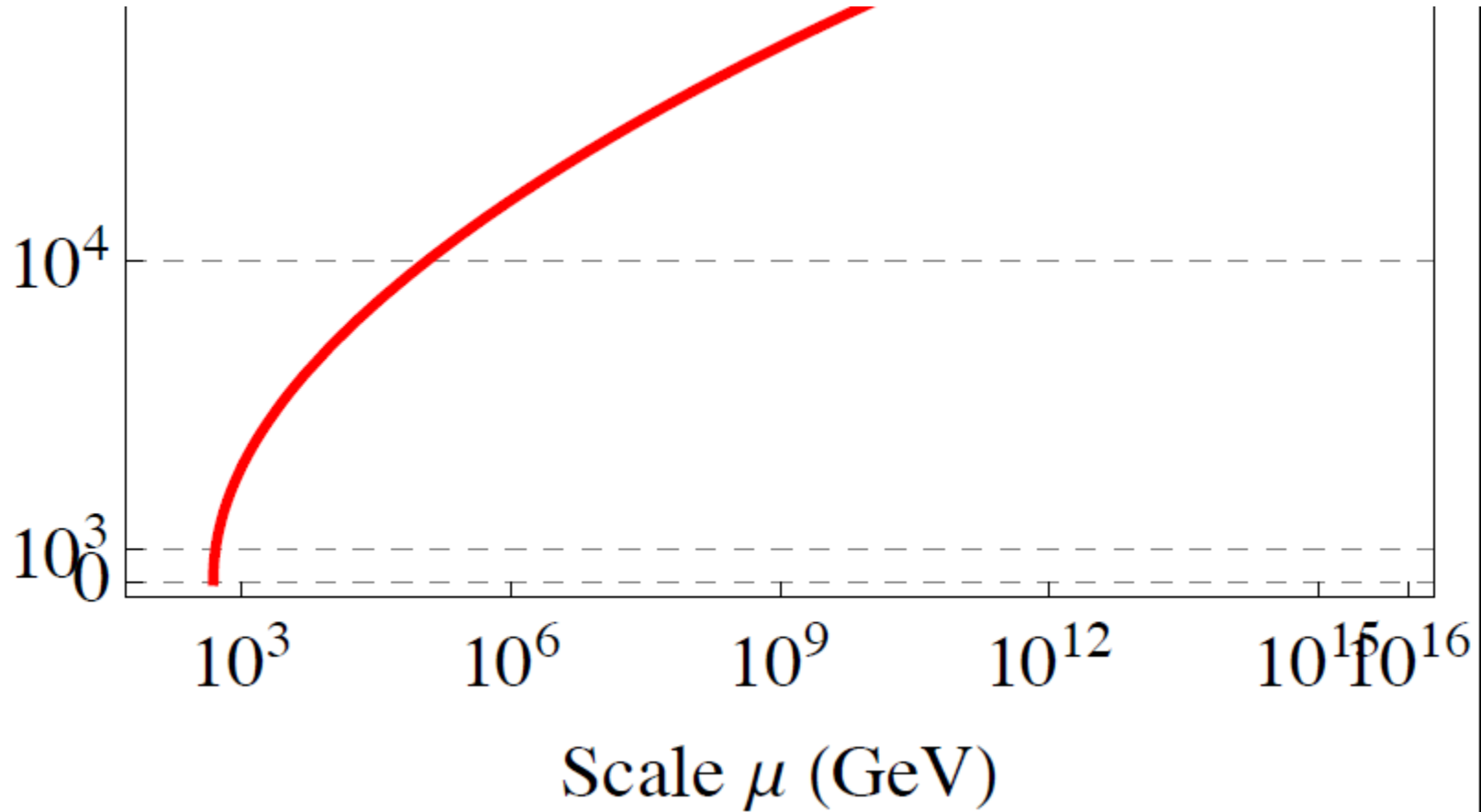
➤ Maybe cancellations – have a theory, so meaningful

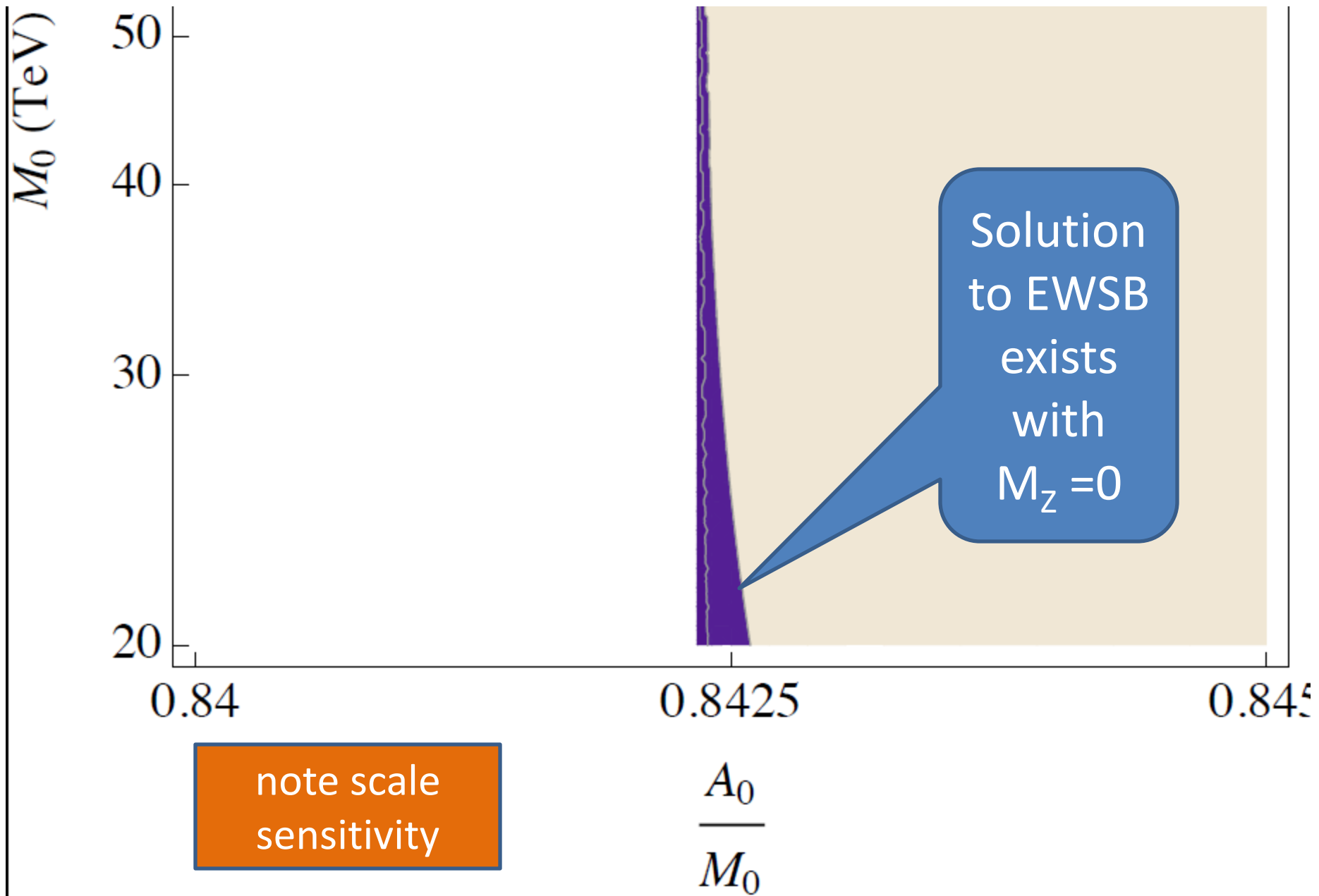
- BUT Calculations of kahler potential, trilinears have corrections – not yet calculable – so can't calculate running well enough

• There are  $M_0$  and  $A_0$  and  $\mu$  in the range  $M_Z \approx 0$

GeV

$M_{H_u}$





## FINAL REMARKS (1)

- **String/M-theory too important to be left to string theorists**
- **10/11 D String/M-theory with curled up small dimensions may seem complicated – but probably it is the SIMPLEST FRAMEWORK THAT COULD SIMULTANEOUSLY INCORPORATE AND *EXPLAIN* ALL THE PHENOMENA WE WANT TO UNDERSTAND – 10/11D needed → meaningful predictions**
- **Compactified M-theory promising candidate for our vacuum – at least shows not premature to study such compactifications**

## FINAL REMARKS (2)

- **Moduli generically present – inevitable in M Theory – implies non-thermal cosmological history – maybe ratio baryons/DM**
- **$M_h/M_Z$  and Higgs decay branching ratios anticipated**
- **LHC: gluino  $\sim 1.5$  TeV, wino, bino  $\sim 0.5$  TeV ( $\pm \sim 10\%$ ) – good signatures – need  $\sim 40$  fb $^{-1}$  because of backgrounds**
- **Hidden sector dark matter candidates generic, probably inevitable – LSP generically always decays**

## FINAL REMARKS (3)

➤ **Many results generic, don't depend on manifold**

- **gravity mediation;**
- **moduli stabilized;**
- **gravitino mass;**
- **scalars heavy;**
- **gauginos light (gluino, LSP etc);**
- **small EDMs**
- **matter dominated cosmological history**
- **EWSB,  $M_h/M_Z$ , h BR** (2 doublets, heavy scalars, EWSB solutions)
- **LSP decays to hidden sector matter**

# FINAL REMARKS (4)

## Possible issues:

- $g_{\mu-2}$ ;
- $N_{\text{eff}}$  (Acharya, Chakrit Pong....1512.07907);
- No clear  $X(760) \rightarrow \gamma\gamma$  candidate

## FINAL REMARKS (5)

- **Landscape? – Obviously many solutions**
- **Examples already show not an obstacle to finding candidate descriptions of our world – then study properties of compactifications to see implications for multiverse populations**
- **Use phenomenology and theory constraints to find regions of landscape like our world**
- **Maybe in each vacuum can calculate all major results (?)**
- **Crucial question - are the many solutions populated? – maybe not [Perry et al; Greene et al; Shiu et al]**



“if people don’t want to come to the ballpark nobody’s  
going to stop them”

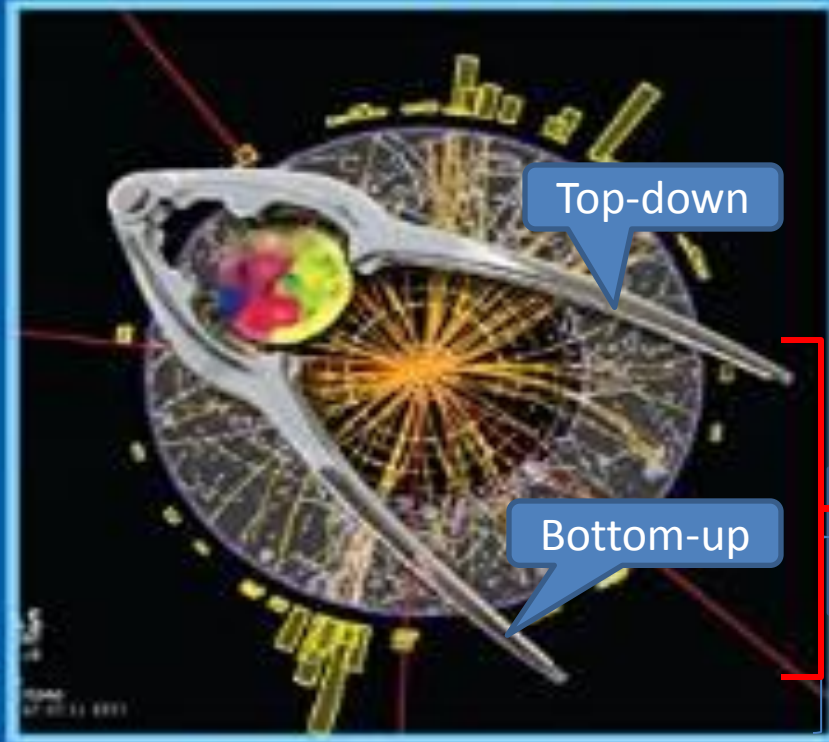
Yogi Berra

ISBN 1782-1328

Advanced Series on  
Directions in High Energy Physics — Vol. 22

# PERSPECTIVES ON STRING PHENOMENOLOGY

Editors  
Bobby Acharya, Gordon I. Kane and Piyush Kumar



Top-down

Bottom-up

String  
phenomenology

Nutcracker!

World Scientific

- **Derive** solution to large hierarchy problem
- Generic solutions with **EWSB derived**
- main F term drops out of **gaugino masses** so **dynamically suppressed**
- **Trilinears**  $> M_{3/2}$  necessarily
- **$\mu$  incorporated in theory**
- Little hierarchy significantly reduced
- **Scalars** =  $M_{3/2} \sim 50 \text{ TeV}$  necessarily, scalars not very heavy
- **Glino lifetime**  $\lesssim 10^{-19}$  sec, decay in beam pipe
- **$M_h \approx 126 \text{ GeV}$  unavoidable** from ratio to Z

## SPLIT SUSY (ETC) MODELS

- Assumes **no solution (possible) for large hierarchy problem**
- **EWSB assumed**, not derived
- **Gauginos suppressed by assumed R-symmetry**, suppression arbitrary
- Trilinears small, suppressed compared to scalars
- **$\mu$  not in theory** at all; guessed to be  $\mu \sim M_{3/2}$
- **No solution to little hierarchy**
- Scalars **assumed** very heavy, whatever you want, e.g.  $10^{10} \text{ GeV}$
- **Long lived gluino**, perhaps meters or more
- **Any  $M_h$  allowed**

# M-THEORY – 11D [M-theory, string theory not yet fully defined – standard in physics ]

- Must “**compactify**” to 4D for our world – geometry is  $X \times \mathbb{R}^{(3,1)}$  ,  $\mathbb{R}$  Minkowski,  $X$  compact manifold [expected to be near Planck scale size (want natural size, time, energy scale set by  $G_N$  ,  $h$ ,  $c$ )]
- $X$  are compact manifolds with  $G_2$  holonomy – admit one covariantly constant spinor  $\rightarrow$  N=1 supersymmetry, a symmetry of the 4D massless modes and interactions and Lagrangian under bosons(integer spin fields)  $\leftrightarrow$  fermions (spin  $\frac{1}{2}$  fields)
- Metrics with  $G_2$  holonomy are Ricci flat, metric is solution of Einstein’s equations in 11D, has finite 4D Newton’s constant, spin 2 massless graviton
- If  $X$  smooth no interesting physics – want solutions with singularities

Why  $N_2 \rightarrow N_1 + h$  dominates:

- $N_2$ - $N_1$ - $h$  coupling from wino-higgsino- $h$  and bino-higgsino- $h$  couplings in gauge eigenstates
- $N_1 \sim$  bino
- $N_2 \sim$  wino
- So  $N_2 \rightarrow N_1 + h$  suppressed by one power of gaugino-higgsino mixing, which is  $\sim M_Z/\mu \leq 1/10$
- Only higgsinos couple directly to  $Z$ , via  $Z$ -higgsino-higgsino vertex, so  $Z$ - $N_1$ - $N_2$  vertex suppressed by two powers, so  $N_2 \rightarrow N_1 + Z$  suppressed by  $\sim (M_Z/\mu)^2$

In Witten's  $M$  theory approach to  $SU(5)$ , the combination of the discrete symmetry, the Wilson lines and the fact that GUT multiplets are localised at points, allows one to prevent the MSSM Higgs doublets,  $H_u$  and  $H_d$ , from having a mass (the  $\mu$ -term) whilst the colour triplets  $D$  and  $\overline{D}$  could have large masses. For simplicity we assume that the symmetry is  $Z_N$ . We use the following notation:  $\overline{\mathbf{5}}^w$  is the multiplet containing  $H_d$  and  $\overline{D}$  and is localised along the Wilson line (which is a circle in the extra dimensions);  $\mathbf{5}^h$  is the multiplet containing  $H_u$ ;  $\overline{\mathbf{5}}^m$  and  $\mathbf{10}^m$  are the matter multiplets. Then the transformation rules for these multiplets under  $Z_N$  are:

$$\begin{aligned}\overline{\mathbf{5}}^w &\rightarrow \eta^\omega \left( \eta^\delta H_d^w \oplus \eta^\gamma \overline{D}^w \right), \\ \mathbf{5}^h &\rightarrow \eta^\chi \mathbf{5}^h, \\ \overline{\mathbf{5}}^m &\rightarrow \eta^\tau \overline{\mathbf{5}}^m, \\ \mathbf{10}^m &\rightarrow \eta^\sigma \mathbf{10}^m,\end{aligned}\tag{1}$$

where  $\eta \equiv e^{2\pi i/N}$ ,  $2\delta + 3\gamma = 0 \pmod N$ . By requiring that Yukawa couplings, Majorana neutrino masses, and colour-triplet masses must be present, we obtain constraints on the charges as can be seen in Table I where we chose  $\omega = 0$ .

TABLE I. Couplings and charges for  $SU(5)$  operators.

| Coupling  | Constraint                           |
|---|--------------------------------------|
| $H_u^h \mathbf{10}^m \mathbf{10}^m$                           | $2\sigma + \chi = 0 \pmod N$         |
| $H_d^w \mathbf{10}^m \overline{\mathbf{5}}^m$                 | $\sigma + \tau + \delta = 0 \pmod N$ |
| $H_u^w H_u^w \overline{\mathbf{5}}^m \overline{\mathbf{5}}^m$ | $2\chi + 2\tau = 0 \pmod N$          |
| $\overline{D}^w D^h$  | $\chi + \gamma = 0 \pmod N$          |

One can solve these by writing all angles in terms of, say,  $\sigma$

$$\begin{aligned}\chi &= -\gamma = -2\sigma \pmod N, \\ \delta &= -3\sigma + N/2 \pmod N, \\ \tau &= 2\sigma + N/2 \pmod N,\end{aligned}\tag{2}$$

which automatically forbids the  $\mu$ -term and dimension four and five proton decay operators.

Generically the vacua of the potential will spontaneously break the  $Z_N$  symmetry. This then generates an effective  $\mu$  term from, e.g. Kähler potential operators of the form

$$K \supset \frac{s}{m_{pl}} H_u H_d + h.c.,\tag{3}$$

à la Giudice-Masiero [20], where  $s$  generically denotes a modulus field of the appropriate charge and  $m_{pl}$  is the Planck scale. Note that such terms are forbidden in the superpotential due to holomorphy and the axion shift symmetries. From [7–9, 19] we know that the moduli vevs are approximately  $\langle s \rangle \sim 0.1 m_{pl}$ ,  $\langle F_s \rangle \sim m_{1/2} m_{pl}$  and from the standard supergravity Lagrangian [21] we get an effective  $\mu$ -term:

$$\mu = \langle m_{3/2} K_{H_u H_d} - F^k K_{H_u H_d k} \rangle,\tag{4}$$

which leads to

$$\mu \sim \frac{\langle s \rangle}{m_{pl}} m_{3/2} + \frac{\langle F_s \rangle}{m_{pl}}.\tag{5}$$

Since gaugino masses are suppressed [7–9, 19], the  $F$ -term vev is subleading and we get

$$\mu \sim 0.1 m_{3/2} \sim \mathcal{O}(TeV).\tag{6}$$

## □ From Planck scale to 50 TeV “dimensional transmutation”

Scale of gaugino condensation  $\Lambda \approx M_{pl} \exp(-8\pi^2 / 3Qg^2) \approx \exp(2\pi \text{Im}f / 3Q)$

where  $\text{Im}f = \sum N_i s_i$

Q is rank of condensing gauge group

With  $Q-P=3$ ,  $\text{Im}f=14Q/\pi \rightarrow \Lambda \approx M_{pl} e^{-28/3} \approx 2 \times 10^{14}$  GeV, so

$\Lambda \approx 10^{-4} M_{pl} \approx$  *scale at which supersymmetry broken*

Then  $W \sim \Lambda^3 \sim 10^{-12} M_{pl} \sim 2 \times 10^6$  GeV =  $2 \times 10^3$  TeV. Also expect inverse volume factor  $1/V_7$  from  $e^{K/2}$  so

$$M_{3/2} \approx e^{K/2} W \sim 50 \text{ TeV}$$

Note  $\text{Im}f/Q$  not explicitly dependent on  $Q$  – still dependent because of  $V_7$  and  $P_{\text{eff}}$ , but weakly – so  $\Lambda$  rather well determined