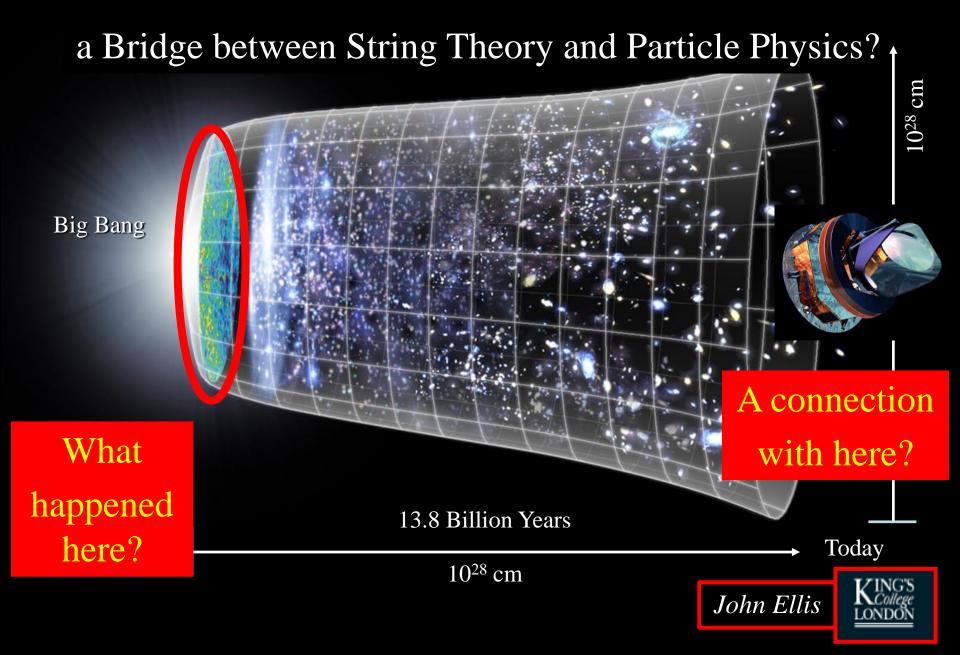
Models of Inflation:





Slow-Roll Inflation

Expansion driven by cosmological constant:

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{\Lambda}{3}$$

- Getting small density perturbations requires a "small" potential: $\left(\frac{V}{\epsilon}\right)^{\frac{1}{4}} = 0.0275 \times M_{Pl}$
- That is almost flat: $\epsilon = \frac{1}{2}M_{Pl}^2\left(\frac{V'}{V}\right)^2$, $\eta = M_{Pl}^2\left(\frac{V''}{V}\right)$ small so as to get sufficient e-folds of expansion:

$$N = \frac{v^2}{M_{Pl}^2} \int_{x_i}^{x_e} \left(\frac{V}{V'}\right) dx$$

Main CMB Observables

- Scalar and tensor perturbations
- Tilt in scalar spectrum (running down hill)

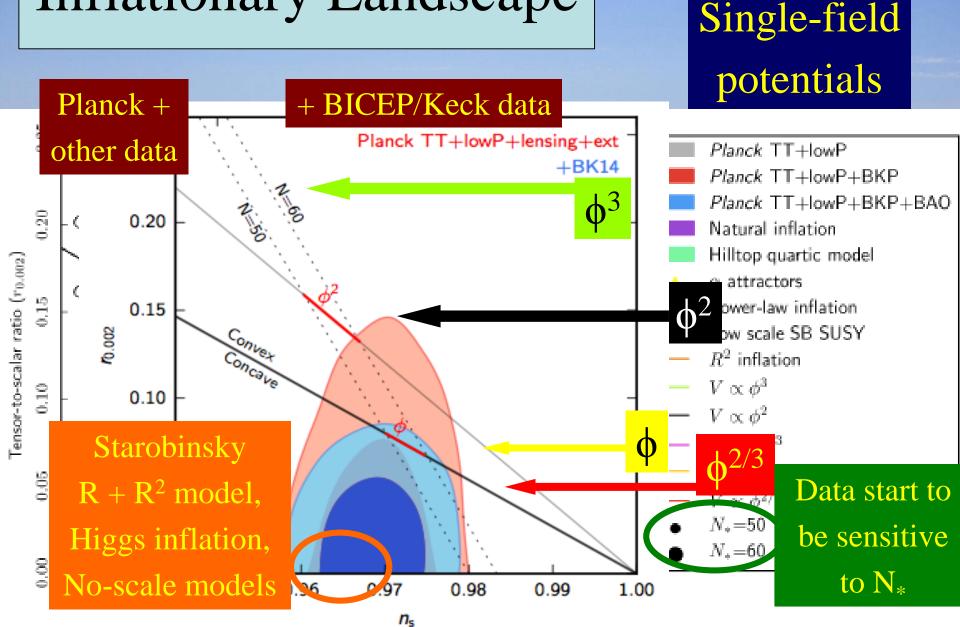
$$n_s = 1 - 6\epsilon + 2\eta$$

- Tensor perturbations = gravitational waves of quantum origin
- Tensor/scalar ratio:

$$r = 16\epsilon$$

- Are perturbations ~ Gaussian?
 - Look for deviations, e.g., f_{NL}
- Expected to be small in slow-roll models

Inflationary Landscape



Monomial

Challenges for Inflationary Models

- Links to low-energy physics?
 - Only SM candidate for inflaton is Higgs
 - •BUT negative potential....
- Link to other physics?
 - -SUSY partner of RH (singlet) neutrino?
 - -Some sort of axion?
- Links to Planck-scale physics?
 - Inflaton candidates in string theory?
 - Inflaton candidates in compactified string models

Inflationary Models in Light of Planck

- Planck CMB observations consistent with inflation
- Tilted scalar perturbation spectrum (rolling down):

$$n_s = 0.968 \pm 0.006$$

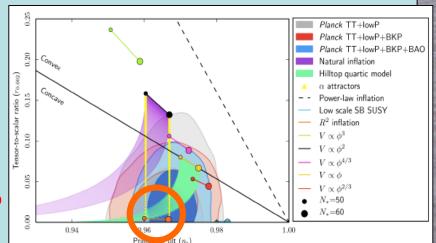
• BUT strengthen upper limit on tensor

perturbations: r < 0.1

• Challenge for many simple inflationary models







Starobinsky Model

- Non-minimal general relativity (singularity-free cosmology): $S = \frac{1}{2} \int d^4x \sqrt{-g} (R + R^2/6M^2)$ • No scalar!?

• Inflationary interpretation, calculation of perturbations:

 $\delta S_b = \frac{1}{2} \int d^4x \left[\phi'^2 - \nabla_\alpha \phi \nabla^\alpha \phi + \left(\frac{a''}{a} + M^2 a^2 \right) \phi^2 \right]$

Conformally equivalent to scalar field model:

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[\tilde{R} + (\partial_{\mu}\varphi')^2 - \frac{3}{2} M^2 (1 - e^{-\sqrt{2/3}\varphi'})^2 \right]$$

Higgs Inflation: a Single Scalar?

Bezrukov & Shaposhnikov, arXiv:0710.3755

• Standard Model with non-minimal coupling to gravity: $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty$

gravity:
$$S_{J} = \int d^{4}x \sqrt{-g} \left\{ -\frac{M^{2} + \xi h^{2}}{2} R \right\}$$
$$+ \frac{\partial_{\mu} h \partial^{\mu} h}{2} - \frac{\lambda}{4} \left(h^{2} - v^{2} \right)^{2} \right\}$$

• Consider case $1 \ll \sqrt{\xi} \ll 10^{17}$: in Einstein frame

$$S_E = \int d^4x \sqrt{-\hat{g}} \left\{ -\frac{M_P^2}{2} \hat{R} + \frac{\partial_\mu \chi \partial^\mu \chi}{2} - U(\chi) \right\}$$

- With potential: $U(\chi) = \frac{\lambda M_P^4}{4\xi^2} \left(1 + \exp\left(-\frac{2\chi}{\sqrt{6}M_P}\right)\right)^{-2}$ Similar to Starobinsky, but not identical
- Successful inflationary potential at $\chi \gg M_P$

Problem for Higgs Inflation

Degrassi, Di Vita, Elias-Miro, Giudice, Isodori & Strumia, arXiv:1205.6497

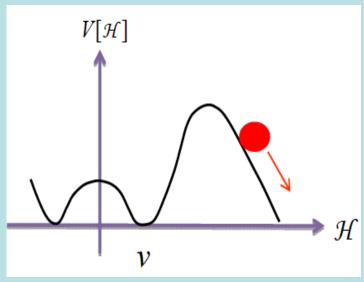
• Large $M_h \rightarrow large self-coupling \rightarrow blow up at$

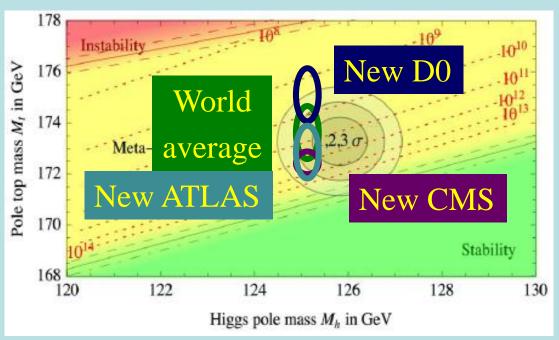
$$\lambda(Q) = \lambda(v) - \frac{3m_t^4}{2\pi^2 v^4} \log \frac{Q}{v}$$

- Small M_h: renormalization due to t quark drives quartic coupling < 0 at some scale Λ
 → vacuum unstable
- Instability @ 80.0 0.06 0.04 0.02 0.00 -0.02 $10^{10} \ 10^{12} \ 10^{14} \ 10^{16} \ 10^{18} \ 10^{20}$ RGE scale μ in GeV
- Negative potential not suitable for inflation
- Problem avoided with supersymmetry

Vacuum Instability in the Standard Model

• Very sensitive to m_t as well as M_H





• Instability scale: Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio & Strumia, arXiv:1307.3536

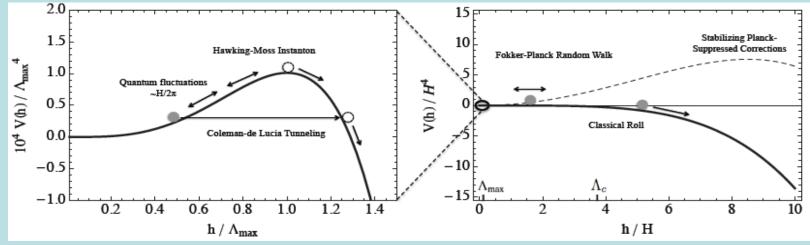
$$\log_{10} \frac{\Lambda_I}{\text{GeV}} = 11.3 + 1.0 \left(\frac{M_h}{\text{GeV}} - 125.66 \right) - 1.2 \left(\frac{M_t}{\text{GeV}} - 173.10 \right) + 0.4 \frac{\alpha_3(M_Z) - 0.1184}{0.0007}$$

 $m_t = 173.3 \pm 1.0 \text{ GeV} \rightarrow \log_{10}(\Lambda/\text{GeV}) = 11.1 \pm$

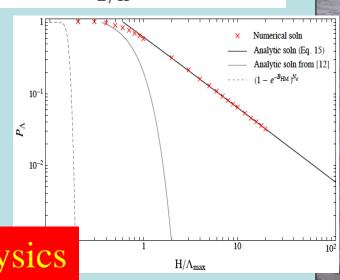
Instability during Inflation?

Hook, Kearns, Shakya & Zurek: arXiv:1404.5953

Do inflation fluctuations drive us over the hill?



- Then Fokker-Planck evolution
- Do AdS regions eat us?
 - Disaster if so
 - If not, OK if more inflation



Avoid with new physics

What lies beyond the Standard Model?

Supersymmetry

Stabilize electroweak vacuum

New motivations From LHC Run 1

- Successful prediction for Higgs mass
 - Should be < 130 GeV in simple models
- Successful predictions for couplings
 - Should be within few % of SM values
- Naturalness, GUTs, string, ..., dark matter

Inflation Cries out for Supersymmetry

- Want "elementary" scalar field $(at least looks elementary at energies << M_P)$
- To get right magnitude of perturbations prefer mass << M_P ($\sim 10^{13}$ GeV in simple ϕ^2 models)
- And/or prefer small self-coupling $\lambda << 1$
- Both technically natural with supersymmetry

No-Scale Supergravity Inflation

- Supersymmetry + gravity = Supergravity
- Include conventional matter?
- Potentials in generic supergravity models have 'holes' with depths \sim M_P^4
- Exception: no-scale supergravity
- Appears in compactifications of string
 Witten, 1985
- Flat directions, scalar potential ~ global model + controlled corrections

 JE, Engvist, Nanopoulos, Olive & Srednicki, 1984

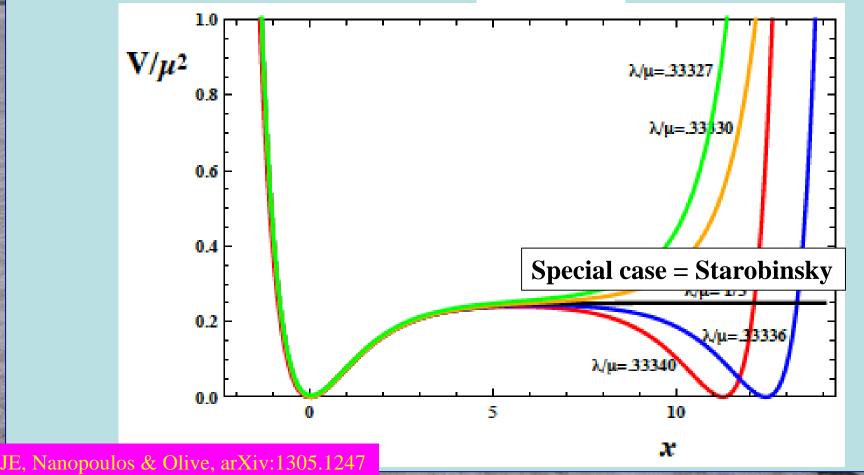
No-Scale Supergravity Inflation Revived

JE, Nanopoulos & Olive, arXiv:1305.1247

- Simplest SU(2,1)/U(1) example:
- Kähler potential: $K = -3\ln(T + T^* |\phi|^2/3)$
- Wess-Zumino superpotential: $W = \frac{\mu}{2}\Phi^2 \frac{\lambda}{3}\Phi^3$
- Assume modulus T = c/2 fixed by 'string dynamics'
- Ef $\mathcal{L}_{eff} = \frac{c}{(c |\phi|^2/3)^2} |\partial_{\mu}\phi|^2 \frac{\hat{V}}{(c |\phi|^2/3)^2}$ $\hat{V} \equiv \left|\frac{\partial W}{\partial \phi}\right|^2$
- Modifications to globally supersymmetric case
- Good inflation possible ...

No-Scale Supergravity Inflation

• Inflationary potential for $\lambda \simeq \mu/3$



Is there more profound connection?

• Starobinsky model:

$$S = \frac{1}{2} \int d^4x \sqrt{-g} (R + R^2/6M^2)$$

• After conformal transformation:

$$S = \frac{1}{2} \int d^4x \sqrt{-\tilde{g}} \left[\tilde{R} + (\partial_{\mu} \varphi')^2 - \frac{3}{2} M^2 (1 - e^{-\sqrt{2/3}\varphi'})^2 \right]$$

- Effective potential: $V = \frac{3}{4}M^2(1 e^{-\sqrt{2/3}\varphi'})^2$
- Identical with the no-scale Wess-Zumino model for the case $\lambda = \mu/3$

... it actually IS Starobinsky

Cecotti, 1987

Beyond Starobinsky

• Exponential approach to constant potential:

$$V = A \left(1 - \delta e^{-Bx} + \mathcal{O}(e^{-2Bx}) \right)$$

• Relations between observables:

$$n_s = 1 - 2B^2 \delta e^{-Bx},$$

 $r = 8B^2 \delta^2 e^{-2Bx},$
 $N_* = \frac{1}{B^2 \delta} e^{+Bx}.$

$$n_s = 1 - \frac{2}{N_*}, r = \frac{8}{B^2 N_*^2}$$

• E.g., multiple no-scale moduli:

$$K \ni -\Sigma_i N_i \ln(T_i + T_i^*): N_i > 0, \Sigma_i N_i = 3$$

- Characteristic of generic string compactifications
- Tensor/scalar ratio may be < prediction of Starobinsky $B = \sqrt{\left(\frac{2}{N_i}\right)}$

$$r = \frac{4N_i}{N_*^2}.$$

• String phenomenology via the

CMB?

model:

JE, Nanopoulos & Olive, arXiv:1307.3537

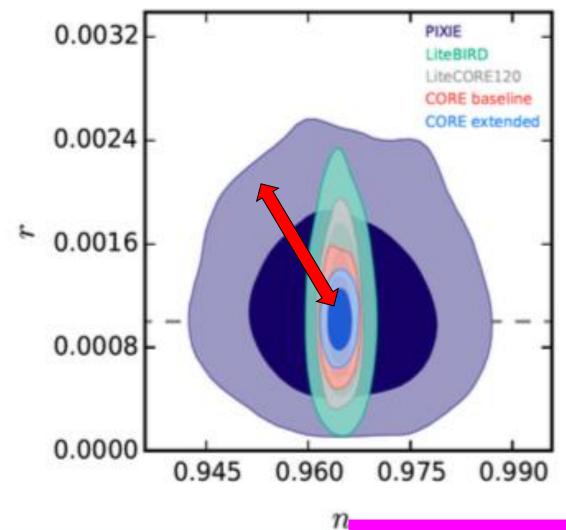
Inflationary Dream

 String-inspired inflationary model with inflation by a Kähler modulus:

$$K \ni -\Sigma_i N_i \ln(T_i + T_i^*)$$

$$n_s = 1 - \frac{2}{N_*}, r = \frac{8}{B^2 N_*^2}$$

- $N_i = 1$
- N_{*} in [44, 59]



See also ...

- Nakayama, Takahashi & Yag
- Kallosh & Linde arXiv:13
- Buchmuller, Domcke & Kar
- Kallosh & Linde arXiv:13
- Farakos, Kehagias and Riot
- Roest, Scalisi & Zavala ar
- Kiritsis arXiv:1307.5873
- Ferrara, Kallosh, Linde & P
- ... over 100 papers

No-Scale Inflation

Recent review

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July 2015

Abstract. Supersymmetry is the most natural framework for physics above the TeV scale, and the corresponding framework for early-Universe cosmology, including inflation, is supergravity. No-scale supergravity emerges from generic string compactifications and yields a non-negative potential, and is therefore a plausible framework for constructing models of inflation. No-scale inflation yields naturally predictions similar to those of the Starobinsky model based on $R+R^2$ gravity, with a tilted spectrum of scalar perturbations: $n_s\sim 0.96$, and small values of the tensor-to-scalar perturbation ratio r<0.1, as favoured by Planck and other data on the cosmic microwave background (CMB). Detailed measurements of the CMB may provide insights into the embedding of inflation within string theory as well as its links to collider physics.

KCL-PH-TH/2015-28, LCTS/2015-20, CERN-PH-TH/2015-144 ACT-05-15, MI-TH-1521, UMN-TH-3442/15, FTPI-MINN-15/32

A No-Scale Inflationary Model to Fit Them All

JE, García, Nanopoulos & Olive, arXiv:1405.0271

• A no-scale supergravity model of inflation that allows larger tensor/scalar values:

$$K = -3\ln(T+\bar{T}) + \frac{|\phi|^2}{(T+\bar{T})^3} \quad W = \sqrt{\frac{3}{4}} \frac{m}{a} \phi(T-a) \quad V = \frac{3m^2}{4a^2} |T-a|^2$$

Motivated by orbifold compactification

• Identify inflaton with components of modulus T:

$$T = a \left(e^{-\sqrt{\frac{2}{3}}\rho} + i\sqrt{\frac{2}{3}} \sigma \right)$$

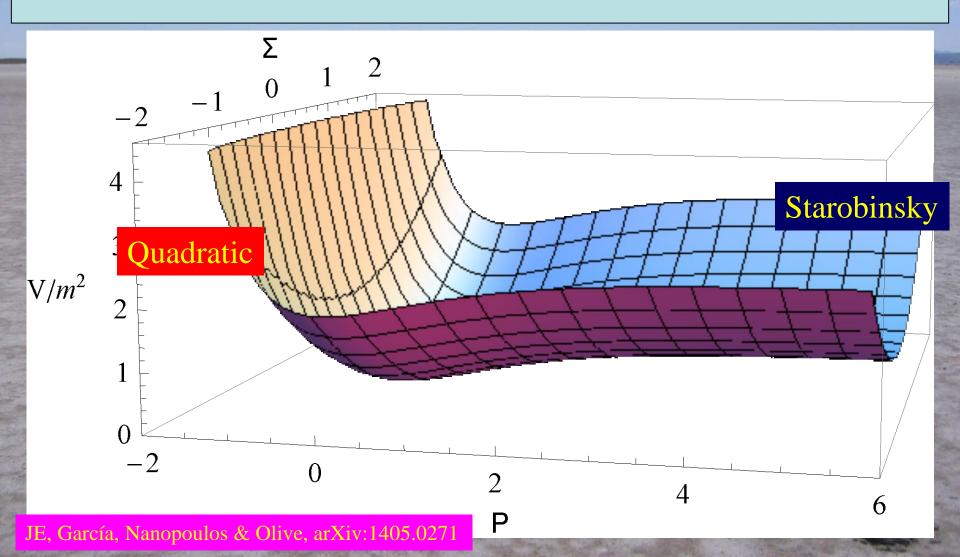
• Effective Lagrangian:

Starobinsky

Quadratic

$$\mathcal{L} \ = \ \frac{1}{2} \partial_\mu \rho \partial^\mu \rho + \frac{1}{2} e^{2\sqrt{\frac{2}{3}}\rho} \partial_\mu \sigma \partial^\mu \sigma - \frac{3}{4} m^2 \left(1 - e^{-\sqrt{\frac{2}{3}}\rho}\right)^2 - \frac{1}{2} m^2 \sigma^2$$

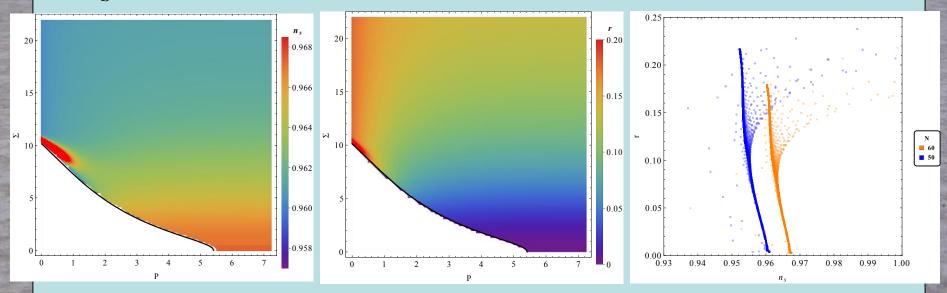
A No-Scale Inflationary Model to Fit Them All



A No-Scale Inflationary Model to Fit Them All

JE, García, Nanopoulos & Olive, arXiv:1405.0271

- Predictions for general initial conditions
- n_s, r as functions of initial values

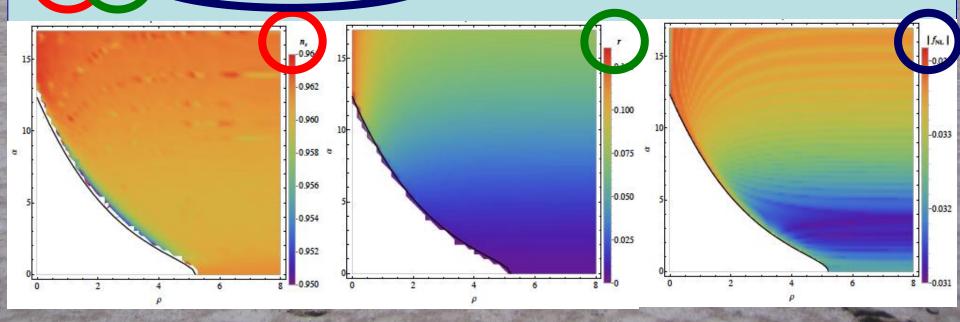


Time will tell what Nature has chosen!

Two-Field Analysis of No-Scale Inflationary Model

JE, García, Nanopoulos & Olive, arXiv:1409.819

- Isocurvature effects on curvature perturbations may suppress tensor/scalar ratio r
- (n_s)r, ion-Gaussianity dependences on initial values



How many e-Folds of Inflation?

• General expression:

JE, García, Nanopoulos & Olive, arXiv:1505.06986

$$N_* = 67 - \ln\left(\frac{k_*}{a_0 H_0}\right) + \frac{1}{4} \ln\left(\frac{V_*^2}{M_P^4 \rho_{\text{end}}}\right) + \frac{1 - 3w_{\text{int}}}{12(1 + w_{\text{int}})} \ln\left(\frac{\rho_{\text{reh}}}{\rho_{\text{end}}}\right) - \frac{1}{12} \ln g_{\text{th}}$$

• In no-scale supergravity models:

Amplitude of perturbations

$$N_{*} = 68.659 - \ln\left(\frac{k_{*}}{a_{0}H_{0}}\right) + \frac{1}{4}\ln\left(A_{S*}\right) - \frac{1}{4}\ln\left(N_{*} - \sqrt{\frac{3}{8}}\frac{\phi_{\text{end}}}{M_{P}}\right) + \frac{3}{4}e^{\sqrt{\frac{2}{3}}\frac{\gamma_{\text{end}}}{M_{P}}}\right) + \frac{1}{12(1+w_{\text{int}})}\left(2.030 + 2\ln\left(\Gamma_{\phi}/m\right) - 2\ln(1+w_{\text{eff}}) - 2\ln(0.81 - 1.10\ln\delta)\right) - \frac{1}{12}\ln g_{\text{th}},$$

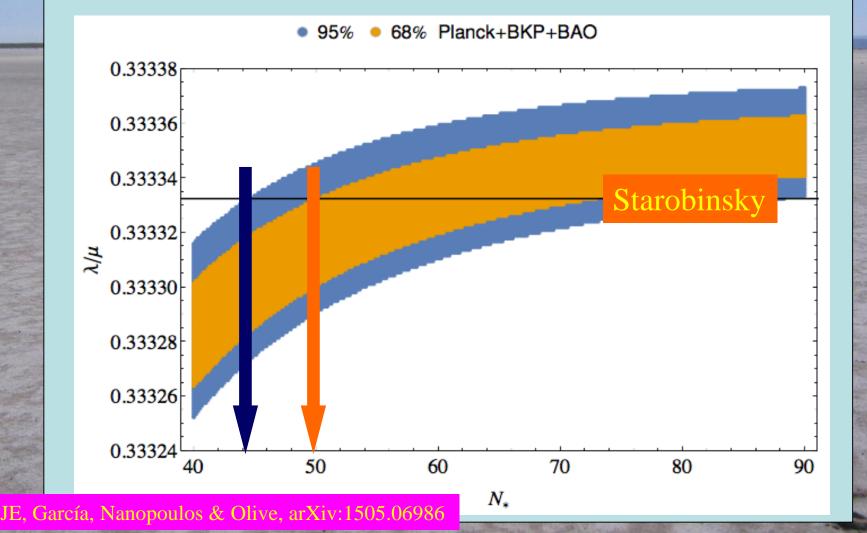
Equation of state during inflaton decay

Inflaton decay rate

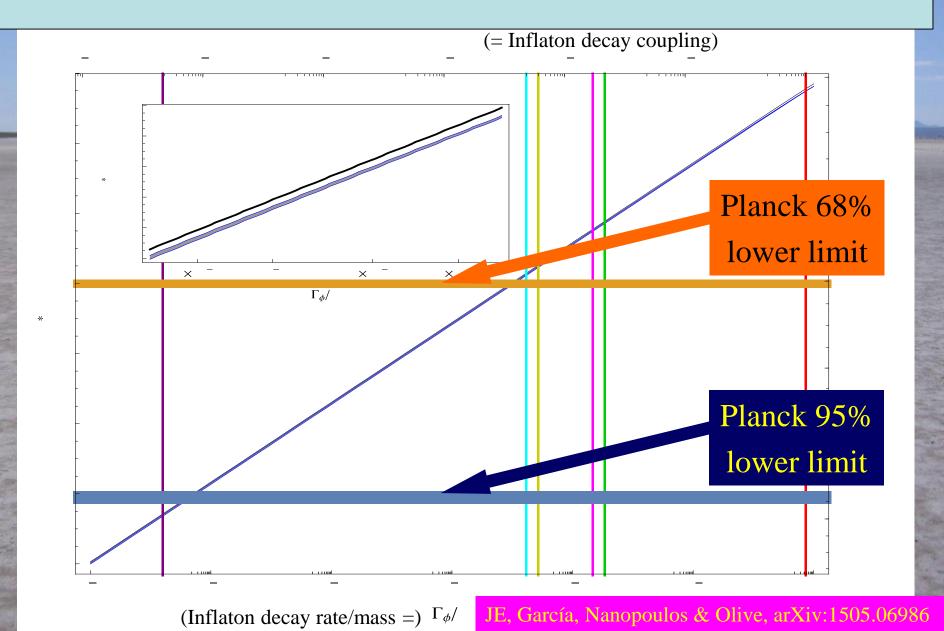
Prospective constraint on inflaton models?

Planck Constraints on # of e-Folds

Starobinsky-like no-scale models



Planck Constraints on # of e-Folds



General Analysis of Reheating

Selected models:

1): R², Higgs

2): Φ^2

3): Φ^4

4): $\Phi^{2/3}$

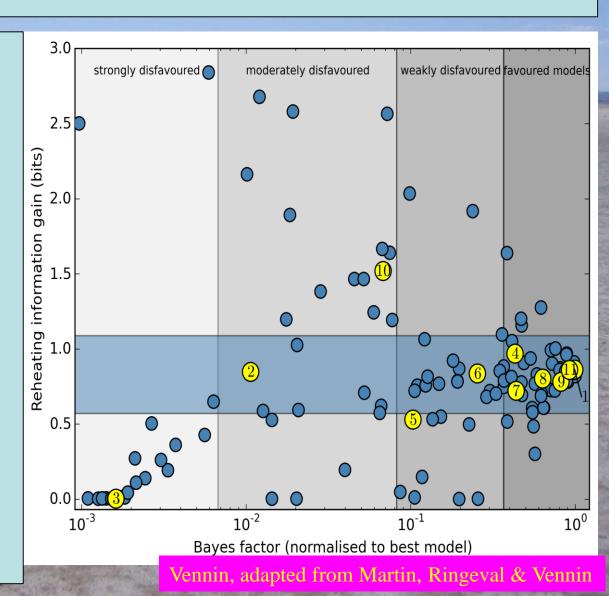
5): Φ^p , p in [0.2, 6]

6), 7), 8): hilltop

9): brane

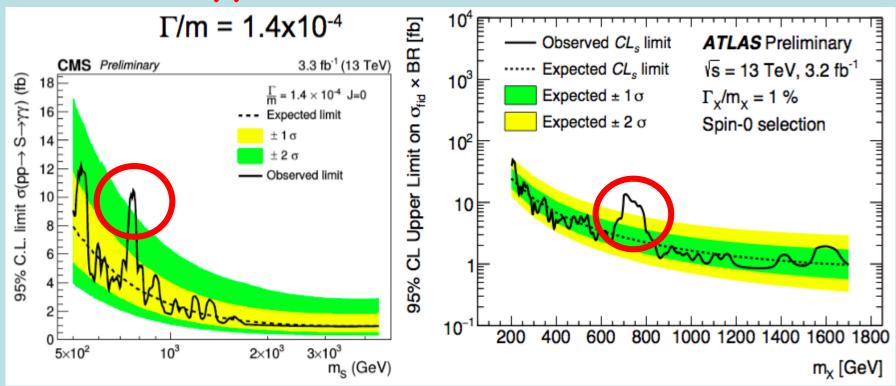
10): natural

11): α attractors

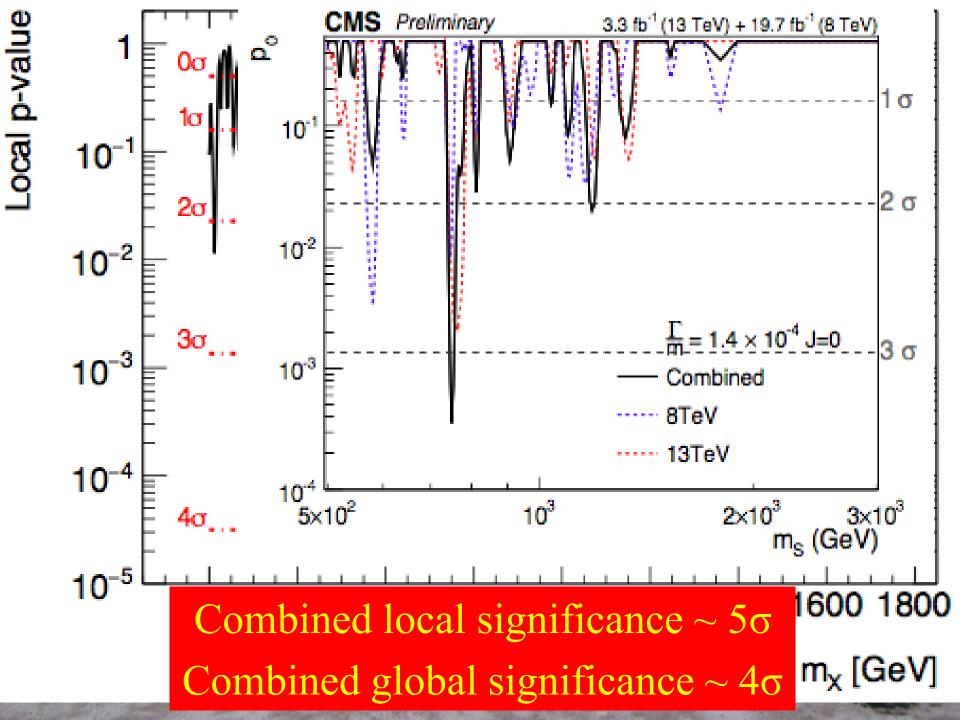


New particle @ 750 GeV?

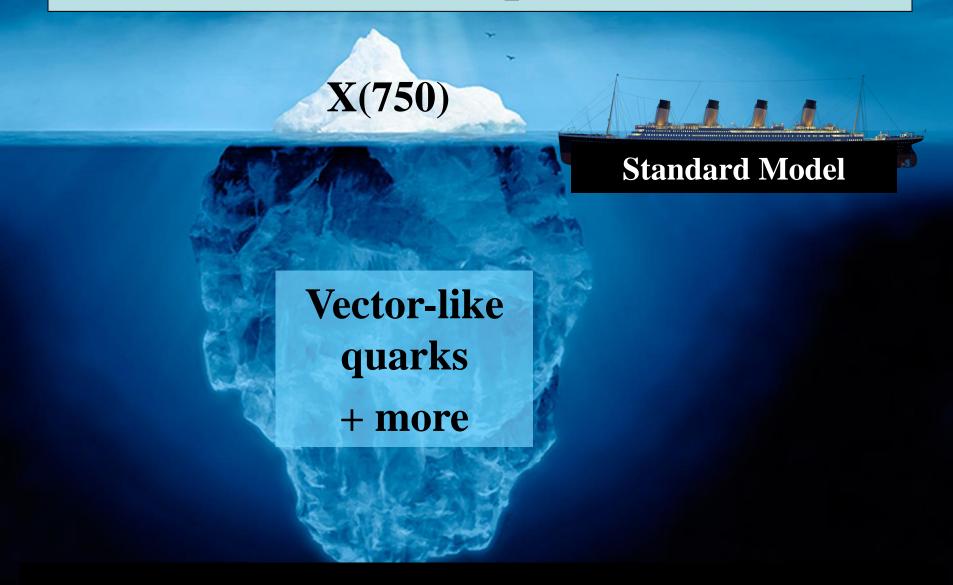
Peaks in γγ invariant mass distributions



- Hint of new X(750) particle has not gone away
- Wait and see!



The Potential Impact of X(750)



X(750) and Inflation

- Another (pseudo)scalar particle?
- Extrapolation to inflation energies non-trivial
- Modifies possibility of Higgs inflation?
 - Need new study of renormalization effects
- Another candidate for the inflaton?
 - Possible with non-minimal X²R coupling
 - Predictions similar to Higgs inflation
- Two-field inflation?
- Wait and see! Should know in a few months