

# Inelastic Thermal DM @ Accelerators

**Gordan Krnjaic**



w/ Eder Izaguirre & Brian Shuve



**1508.03050 PRD D93 (2016) no.6, 063523**

**Dark Sectors Workshop SLAC, April 29, 2016**

# How to make iDM?

Start with familiar model:  $A'$  + Dirac fermion  $\psi = (\eta, \xi^\dagger)$

$$A'_\mu \bar{\psi} \gamma^\mu \psi = A'_\mu (\eta^\dagger \bar{\sigma}^\mu \eta - \xi^\dagger \bar{\sigma}^\mu \xi)$$

$$m_D \bar{\psi} \psi = m_D \eta \xi + h.c.$$

# How to make iDM?

Start with familiar model:  $A'$  + Dirac fermion  $\psi = (\eta, \xi^\dagger)$

$$A'_\mu \bar{\psi} \gamma^\mu \psi = A'_\mu (\eta^\dagger \bar{\sigma}^\mu \eta - \xi^\dagger \bar{\sigma}^\mu \xi)$$

$$m_D \bar{\psi} \psi = m_D \eta \xi + h.c.$$

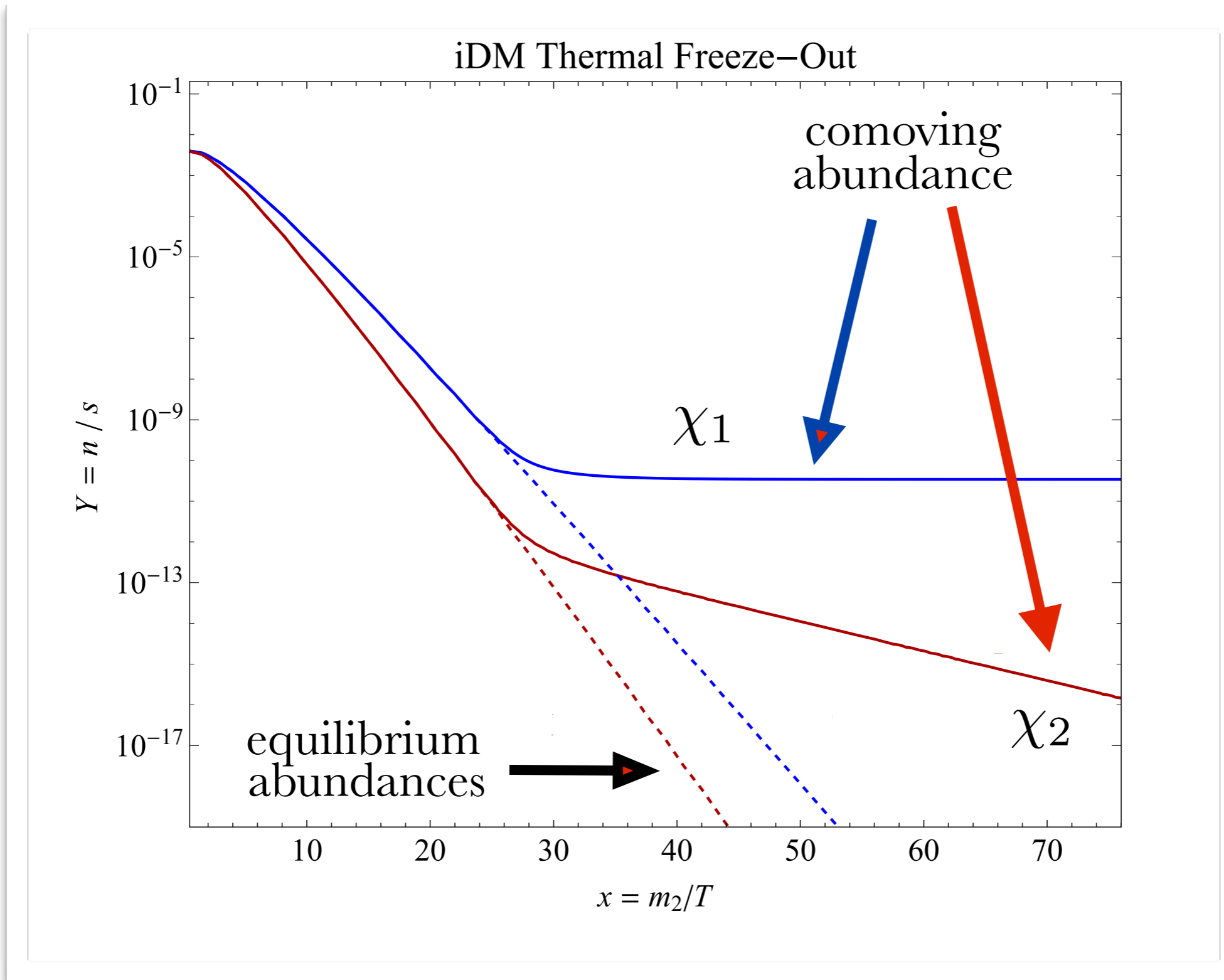
Dark Higgs can induce Majorana masses in IR

$$y H_D \eta \eta + y' H_D^\dagger \xi \xi \rightarrow m_\eta \eta \eta + m_\xi \xi \xi$$

Diagonalize to mass basis  $\chi_{1,2} = \xi \pm \eta$ ,  $m_{1,2} = m_D \mp \Delta$

$$\mathcal{L} \supset A'_\mu \chi_1^\dagger \bar{\sigma}^\mu \chi_2 + h.c.$$

# (Slightly) Different Cosmological History

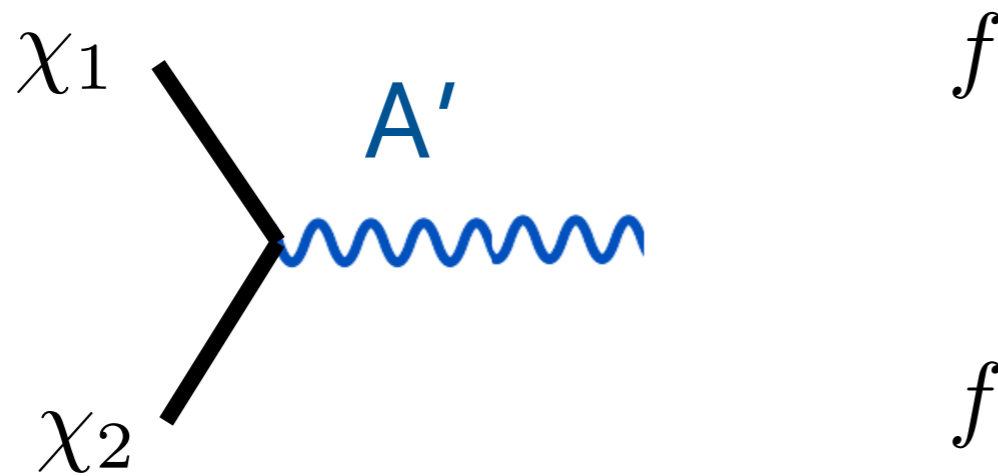


$$\alpha_D = 0.1, \quad m_1 = 10 \text{ GeV}, \quad m_1/m_{A'} = 1/3$$

# RECALL PHILIP SCHUSTER'S PLENARY

## Thermal Origin Target

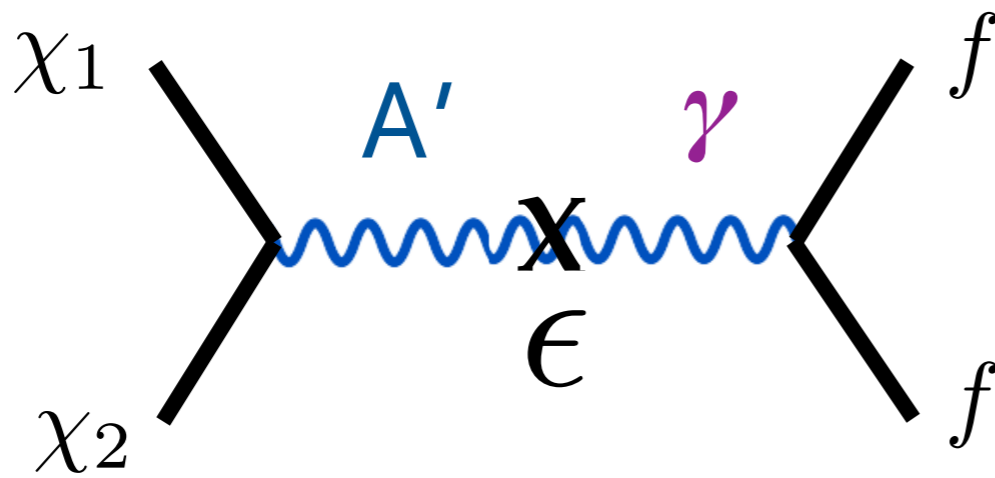
For invisibly decaying  $A'$        $m_{A'} > m_1 + m_2$



# RECALL PHILIP SCHUSTER'S PLENARY

## Thermal Origin Target

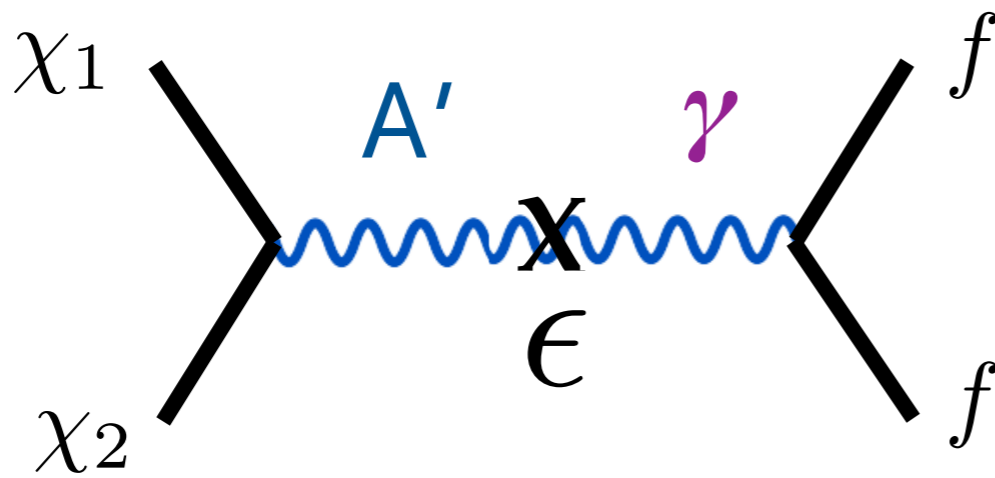
For invisibly decaying  $A'$        $m_{A'} > m_1 + m_2$



# RECALL PHILIP SCHUSTER'S PLENARY

## Thermal Origin Target

For invisibly decaying  $A'$        $m_{A'} > m_1 + m_2$

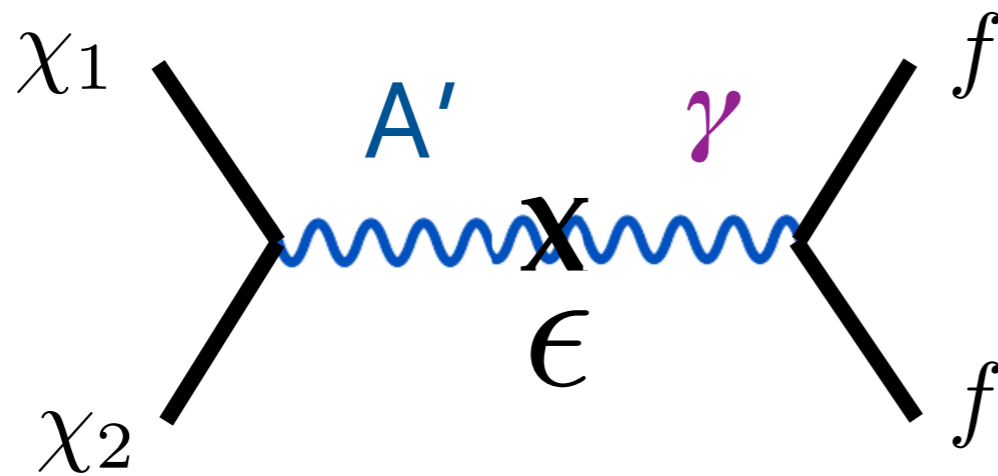


$$\times m_{\chi}^2 \times \frac{1}{m_{\chi}^2}$$

# RECALL PHILIP SCHUSTER'S PLENARY

## Thermal Origin Target

For invisibly decaying  $A'$        $m_{A'} > m_1 + m_2$

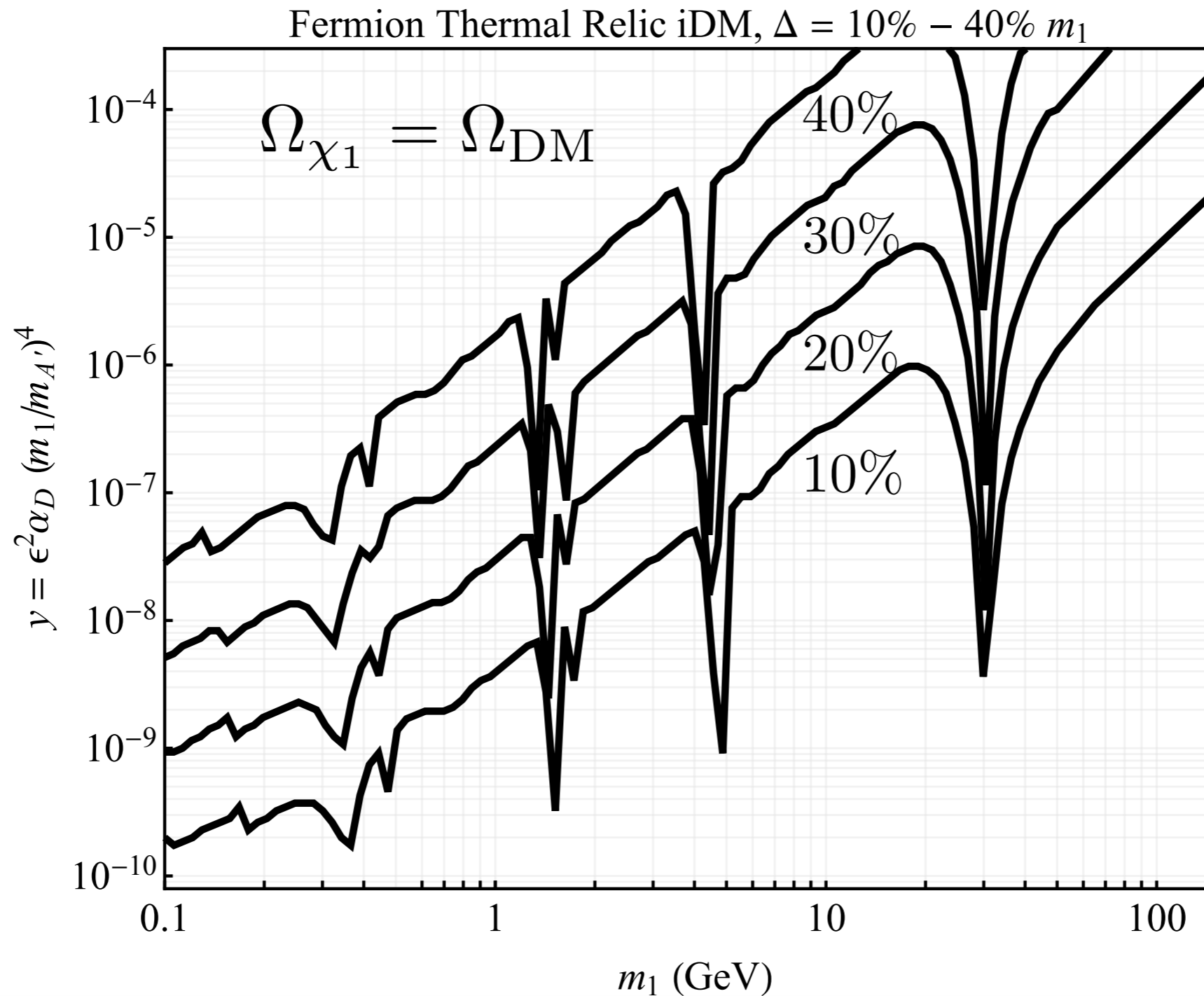


$$\sigma v \sim \alpha_D \epsilon^2 \alpha \times \frac{m_\chi^2}{m_{A'}^4} \times m_\chi^2 \times \frac{1}{m_\chi^2}$$

$y \equiv$  Dimensionless parameter  
controlling cross-section



# Thermal Relic Targets

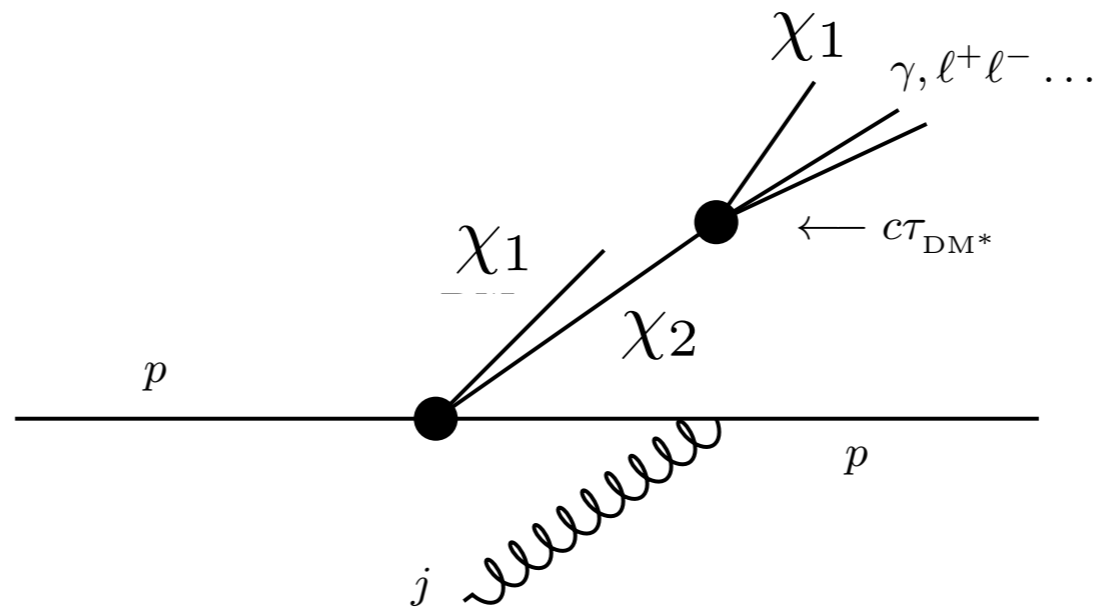


**For given splitting, required “y” fixed by DM mass**

# Signal Topology

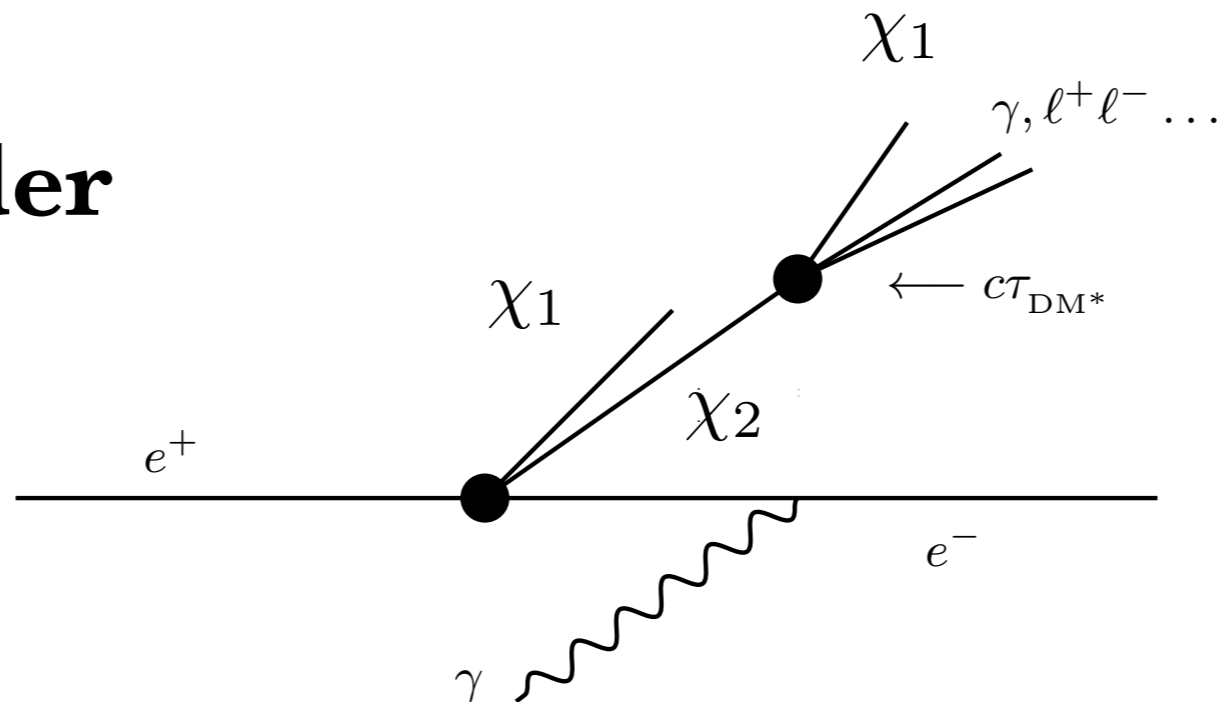
## Hadron Collider

$$J + \cancel{E}_T + l^+ l^-$$

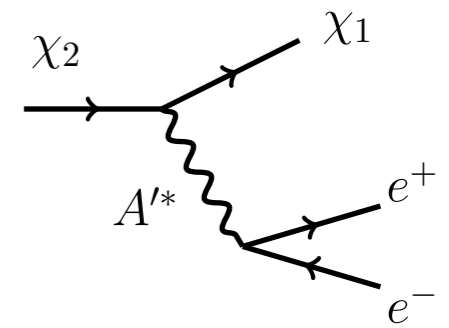
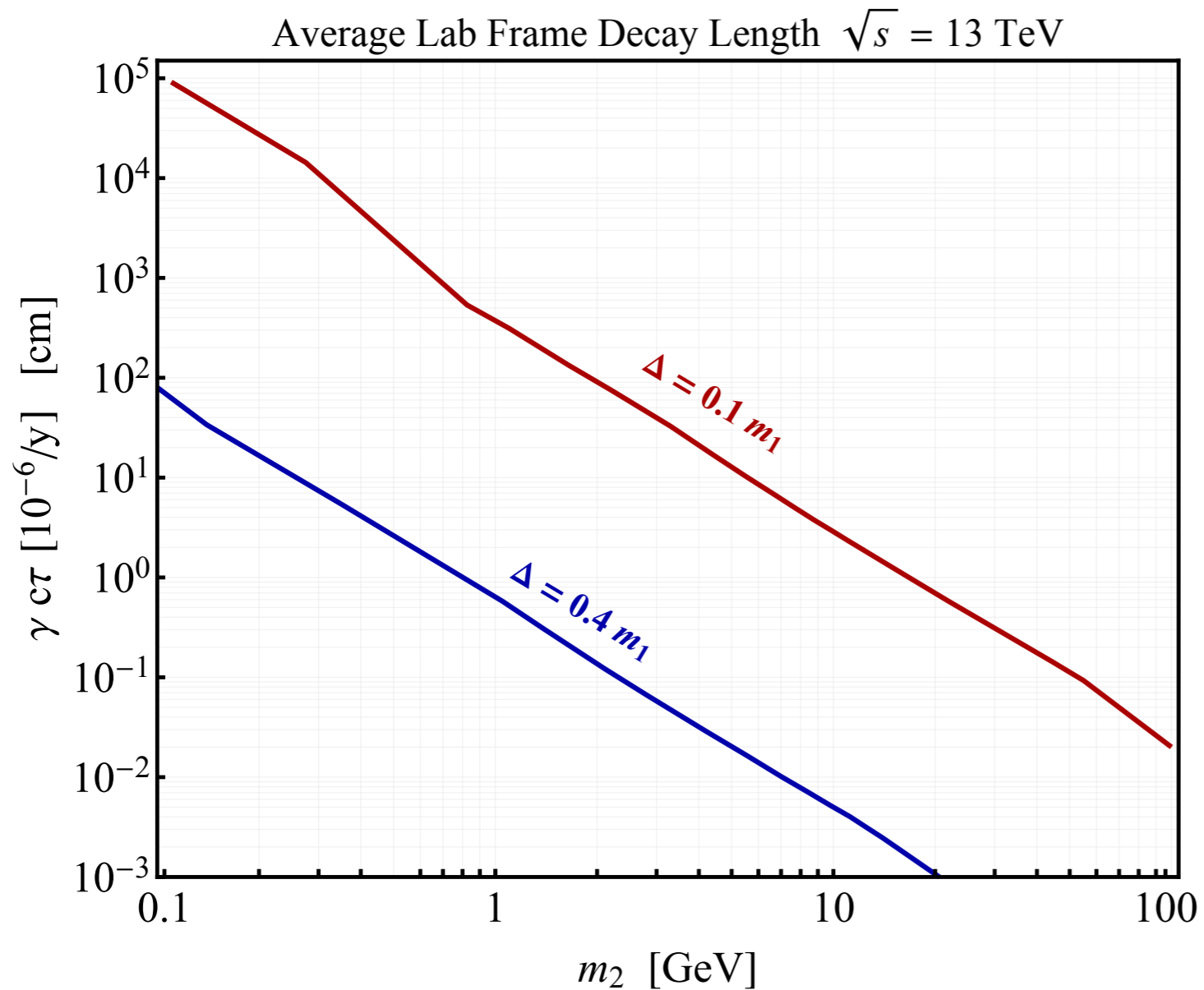


## Lepton Collider

$$\gamma + \cancel{E} + l^+ l^-$$



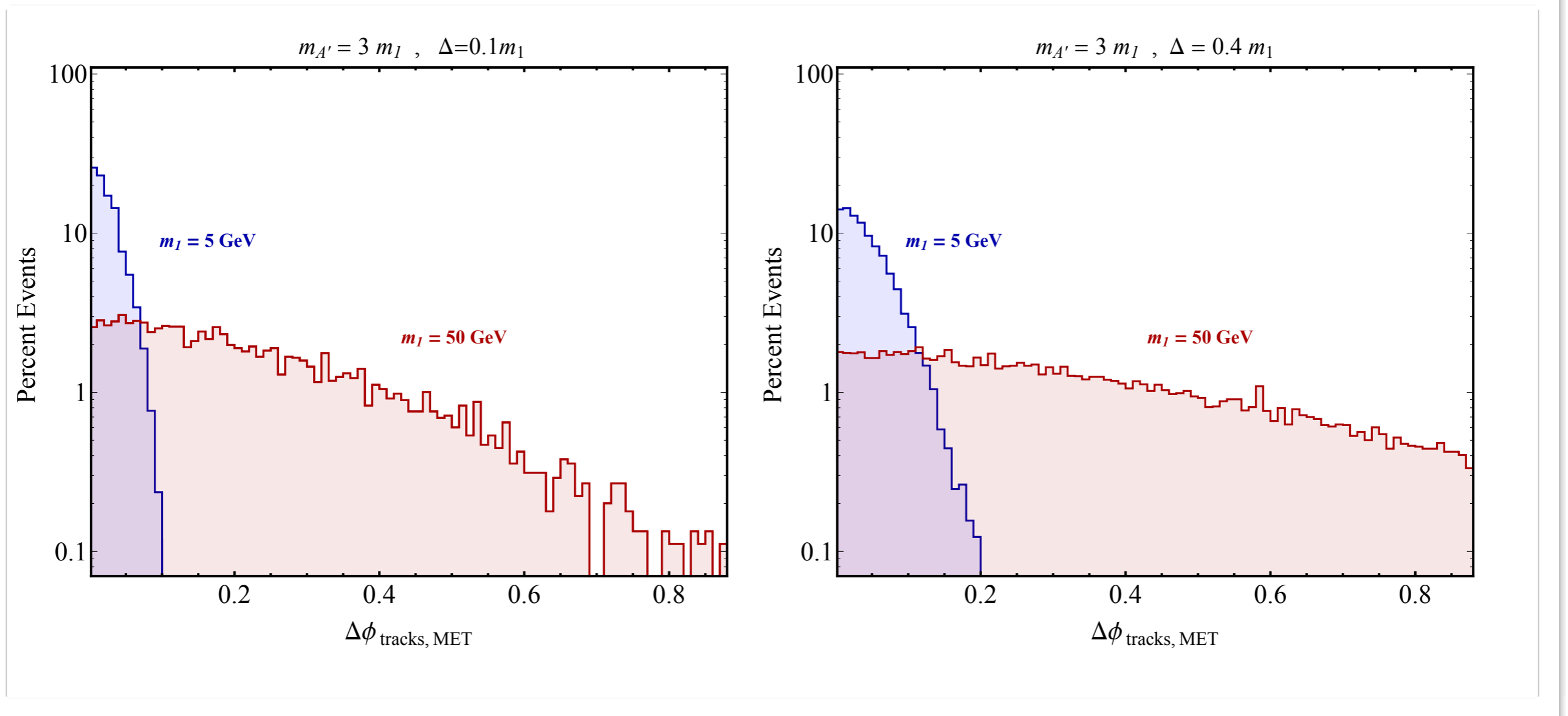
# Signal Feature: Displaced Vertex



$$\Gamma_{\chi_2} = \frac{4\epsilon^2 \alpha \alpha_D \Delta^5}{15\pi m_{A'}^4}$$

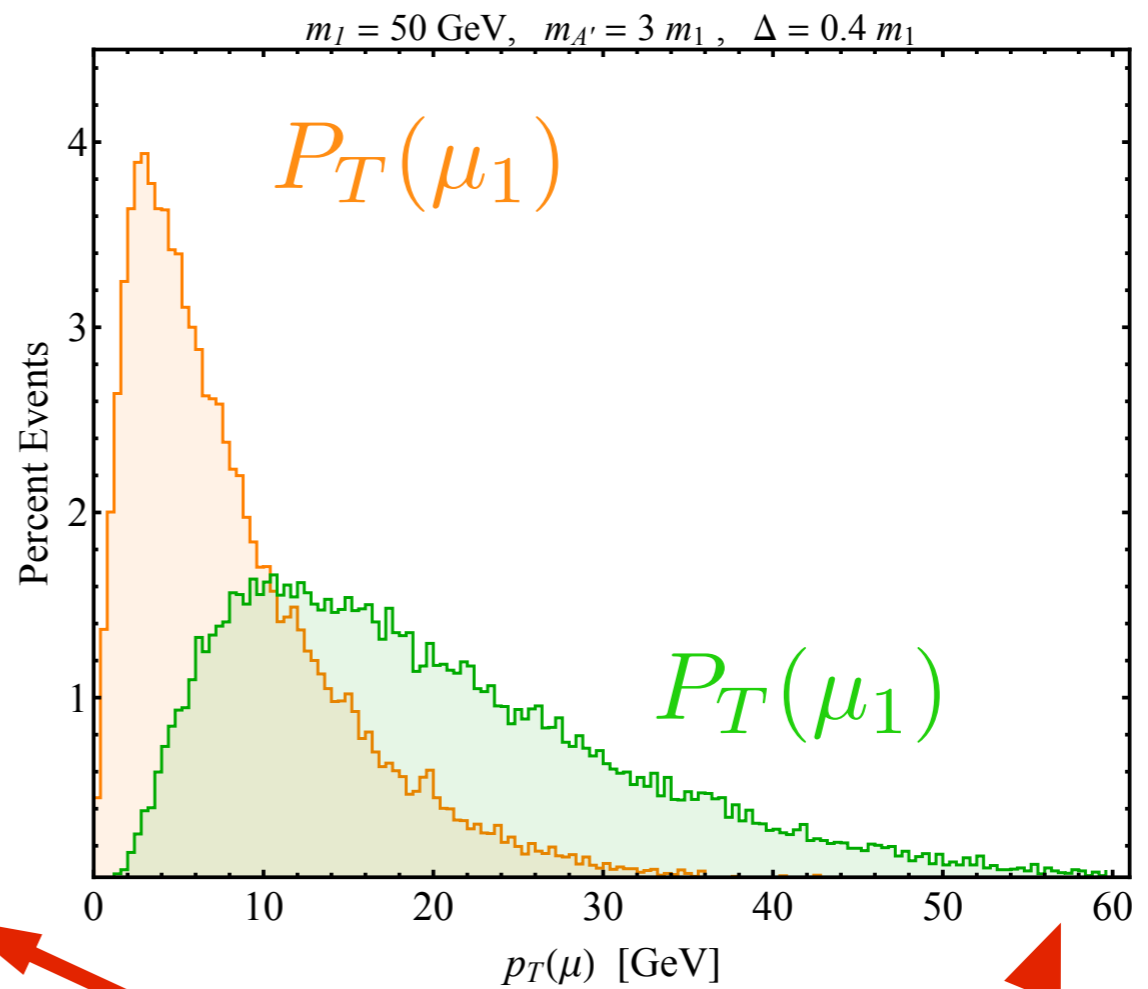
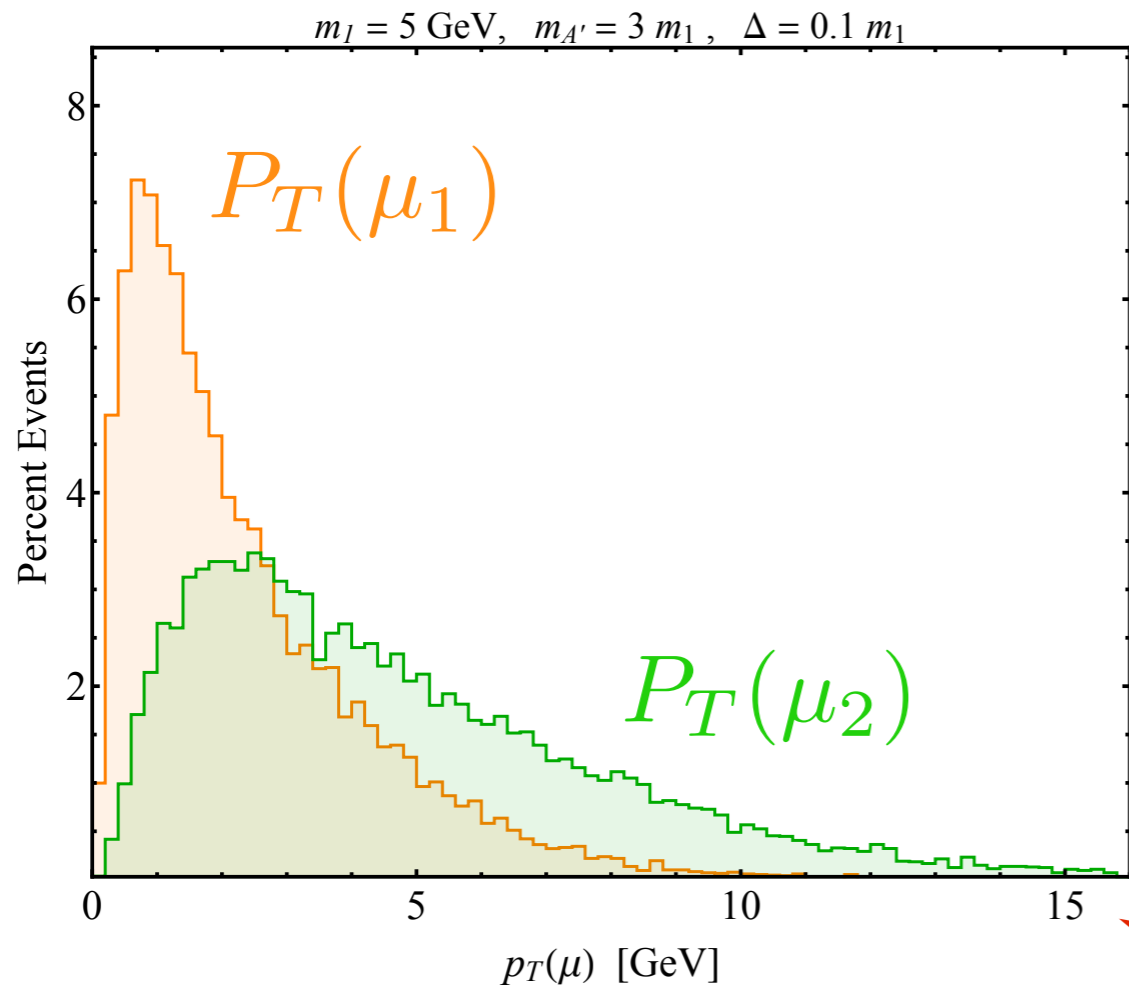
LHC 13 TeV  $\alpha_D = 0.1$ ,  $m_1/m_{A'} = 1/3$

# Signal Feature: MET/Lepton Correlated



LHC 13 TeV  $\alpha_D = 0.1, m_1/m_{A'} = 1/3$

# Signal Feature(Bug): Soft Leptons



**note range**

LHC 13 TeV  $\alpha_D = 0.1, m_1/m_{A'} = 1/3$

# LHC 13 Signal Region

- Trigger on monojet +  $\cancel{E}_T > 120$  GeV
- Leading jet  $P_T(j) > 120$  GeV
- Leading jet &  $\cancel{E}_T$  back-to-back
- Displaced muon jet  $\sim 1\text{mm} - 30\text{cm}$
- Muon  $P_T(\mu) > 5$  GeV
- Muons not isolated  $|\Delta\phi(\cancel{E}_T, \mu J)| < 0.4$ .

# LHC Backgrounds

**Dileptons from photon conversion in detector**

**Dileptons from displaced QCD**

**Pile Up**

**Jets +  $\tau\tau$**

**Verdict: small, reducible (see backup slides)**

# BaBar/Belle Search

$$e^+e^- \rightarrow \gamma A' \rightarrow \gamma \chi_1 \chi_2 \rightarrow \gamma \cancel{E} + \ell^+ \ell^-$$

## Signal Region

- Trigger on lepton  $p > 100 \text{ MeV}$
- Transverse impact param.  $\sim 1\text{mm} - 30\text{cm}$

## Potential BGs low:

Hadronic resonances (can reconstruct)

Conversion from  $e^+e^- \rightarrow \gamma \pi^+ \pi^-$   $e^+e^- \rightarrow \gamma \gamma$

reducible w/ missing mass and displacement



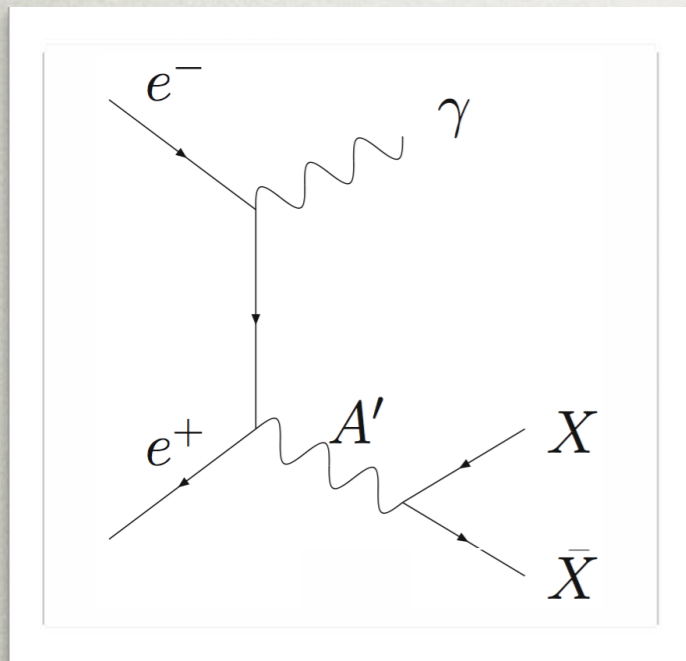
# RECALL FROM PHILIP'S PRIMER

For direct DM production  
(on-shell mediator):

$$y \equiv \alpha_D \epsilon^2 \left( \frac{m_\chi}{m_{A'}} \right)^4$$

$$\sigma \sim \epsilon^2 \frac{1}{E_{CM}^2}$$

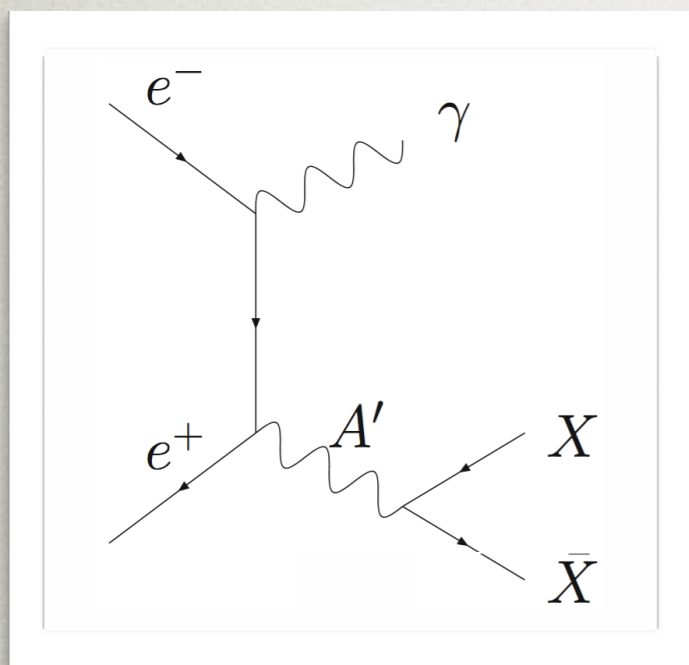
$$\sim y \times \underbrace{\frac{1}{\alpha_D} \left( \frac{m_{A'}}{m_\chi} \right)^4}_{}$$



# RECALL FROM PHILIP'S PRIMER

For direct DM production  
(on-shell mediator):

$$y \equiv \alpha_D \epsilon^2 \left( \frac{m_\chi}{m_{A'}} \right)^4$$



$$\sigma \sim \epsilon^2 \frac{1}{E_{CM}^2}$$

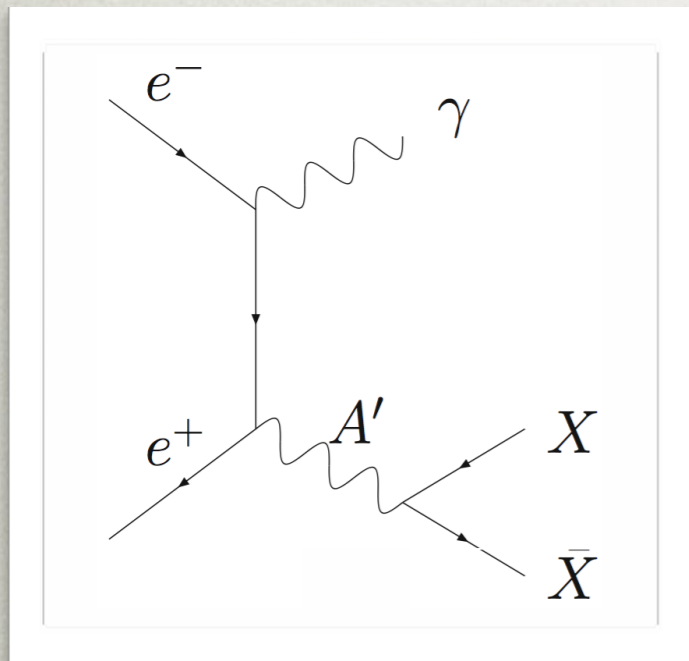
$$\sim y \times \underbrace{\frac{1}{\alpha_D} \left( \frac{m_{A'}}{m_\chi} \right)^4}_{> 100}$$

In this scenario  $\left[ \begin{array}{l} \alpha_D \leq 0.1 - 1 \\ m_{A'} \geq (2 - 3)m_\chi \end{array} \right] \Rightarrow \text{combination} > 100$

# RECALL FROM PHILIP'S PRIMER

For direct DM production  
(on-shell mediator):

$$y \equiv \alpha_D \epsilon^2 \left( \frac{m_\chi}{m_{A'}} \right)^4$$



$$\sigma \sim \epsilon^2 \frac{1}{E_{CM}^2}$$

$$\sim y \times \underbrace{\frac{1}{\alpha_D} \left( \frac{m_{A'}}{m_\chi} \right)^4}_{}$$

In this scenario  $\left[ \begin{array}{l} \alpha_D \leq 0.1 - \textcircled{1} \\ m_{A'} \geq (2 - \textcircled{3}) m_\chi \end{array} \right] \Rightarrow \text{combination} > 100$

Saturate inequalities to show worst-case sensitivity

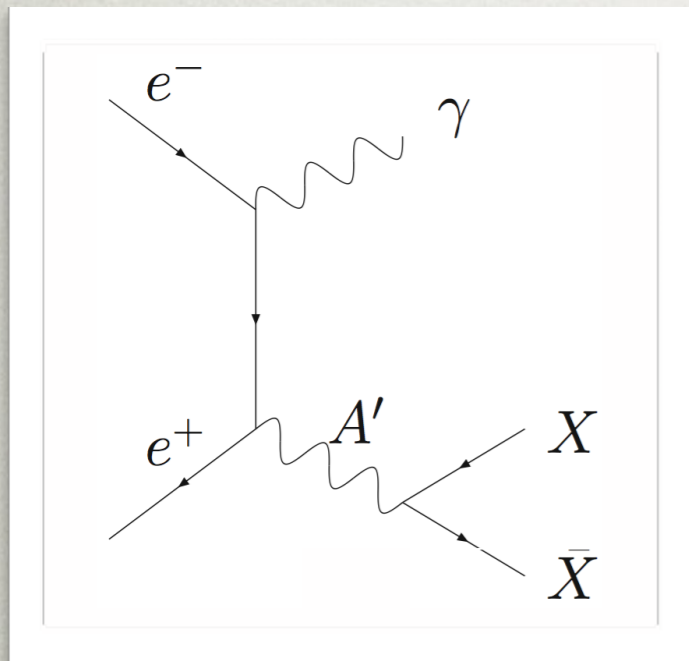
# RECALL FROM PHILIP'S PRIMER

For direct DM production  
(on-shell mediator):

$$y \equiv \alpha_D \epsilon^2 \left( \frac{m_\chi}{m_{A'}} \right)^4$$

$$\sigma \sim \epsilon^2 \frac{1}{E_{CM}^2}$$

$$\sim y \times \underbrace{\frac{1}{\alpha_D} \left( \frac{m_{A'}}{m_\chi} \right)^4}_{> 100}$$

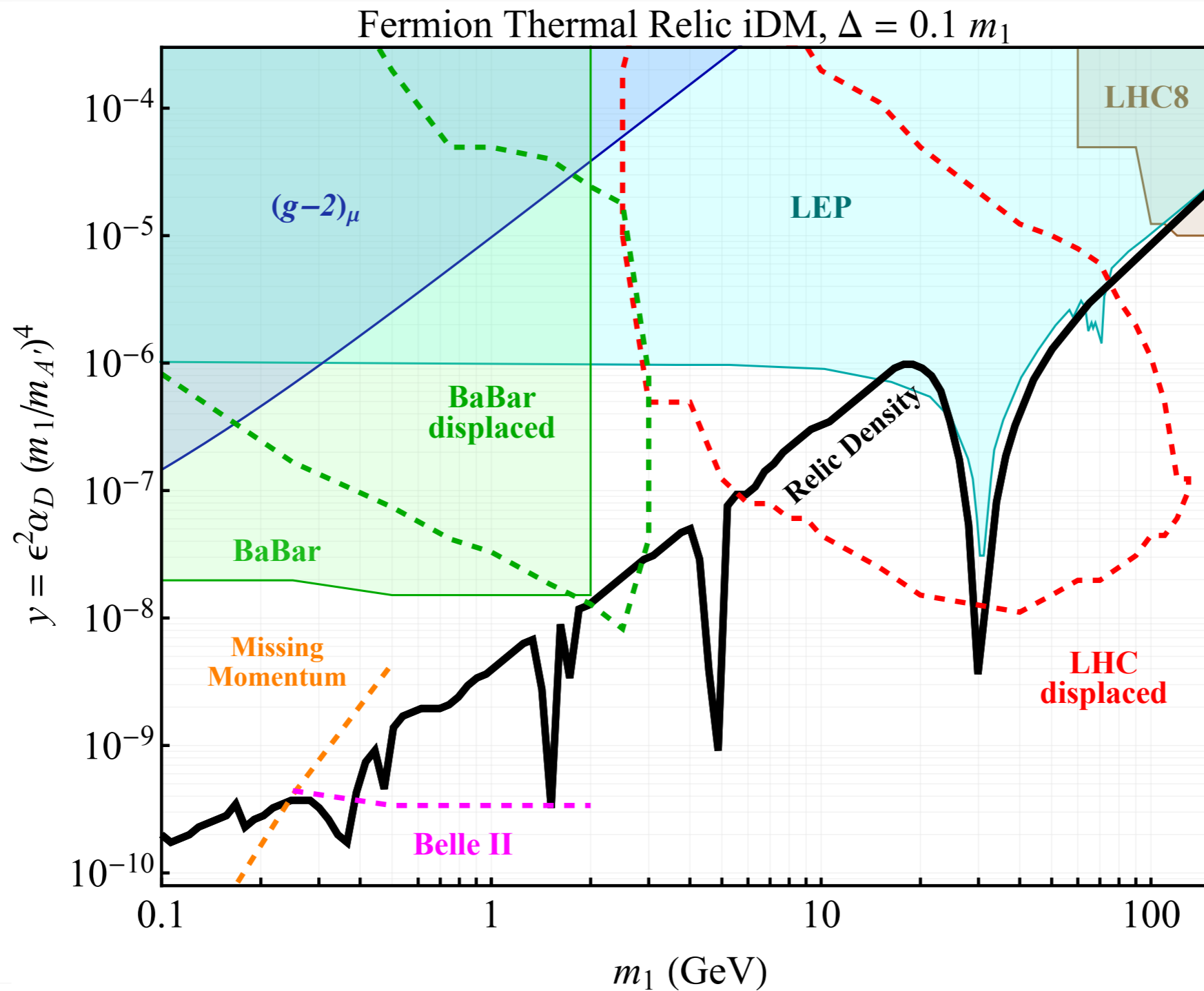


In this scenario  $\left[ \begin{array}{l} \alpha_D \leq 0.1 - \textcircled{1} \\ m_{A'} \geq (2 - \textcircled{3}) m_\chi \end{array} \right] \Rightarrow \text{combination} > 100$

Saturate inequalities to show worst-case sensitivity

Logic applies to both LHC & LEP

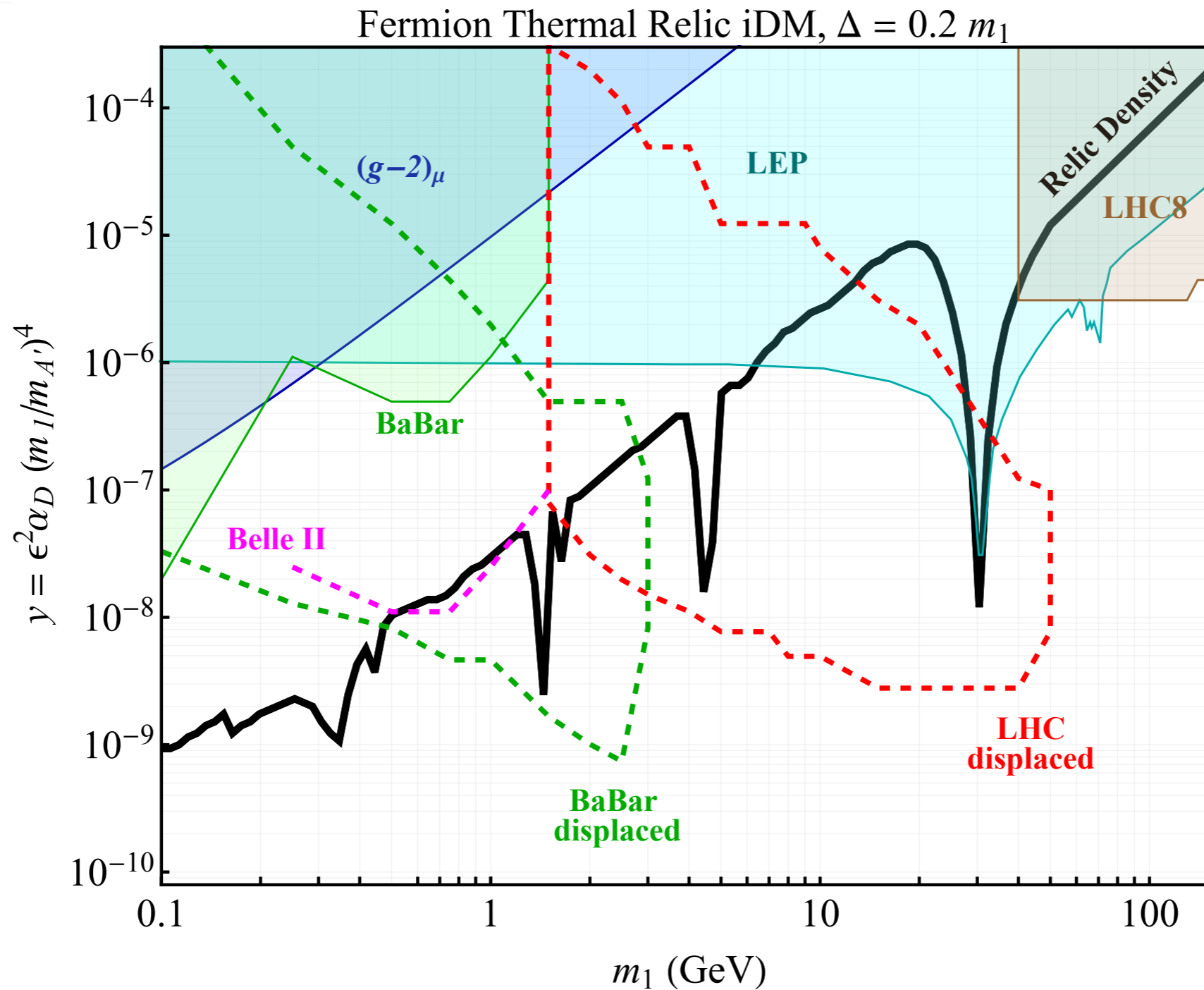
# 10% Splitting



Conservative

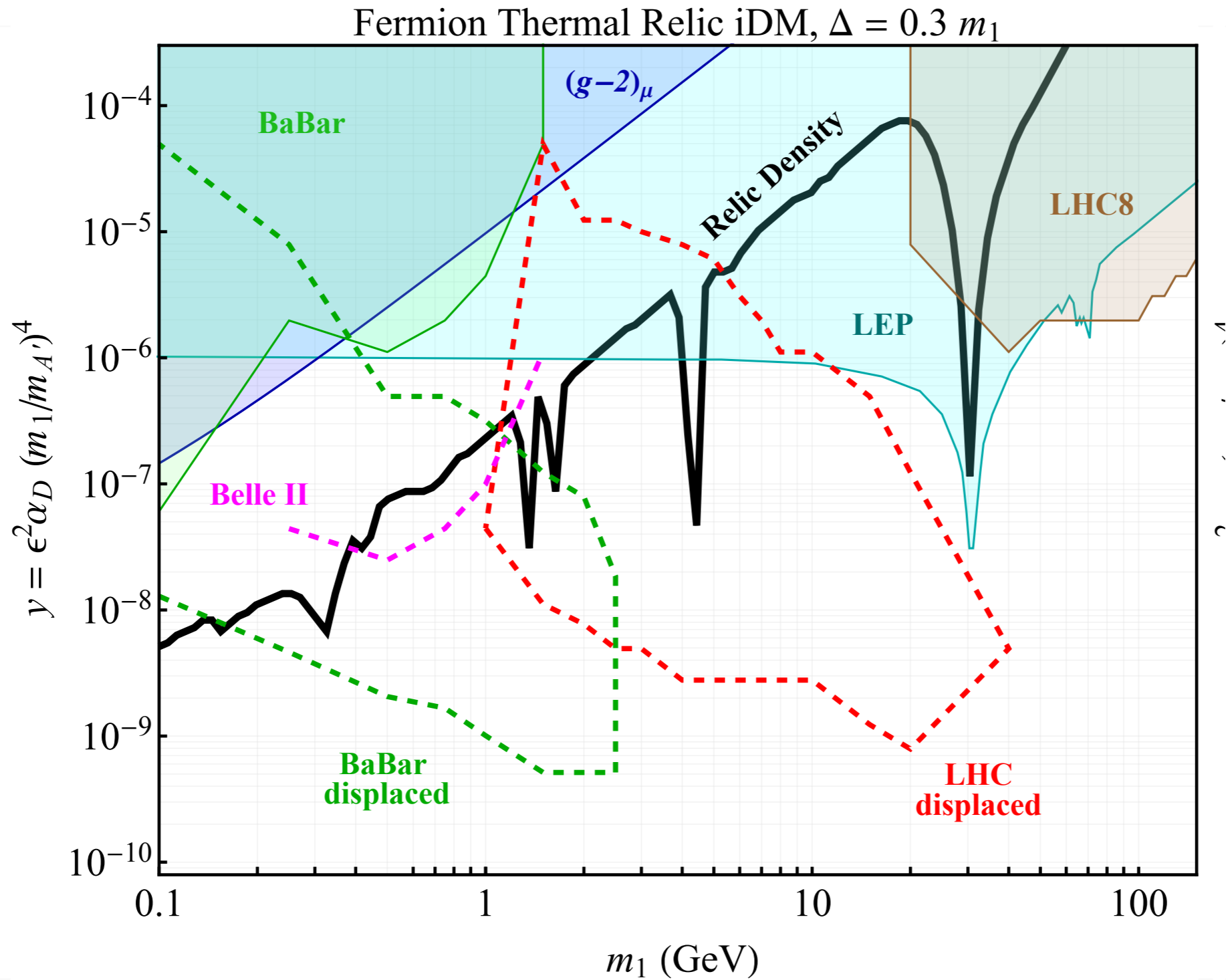
$$\alpha_D = 0.1, m_{A'}/m_1 = 3$$

# 20% Splitting



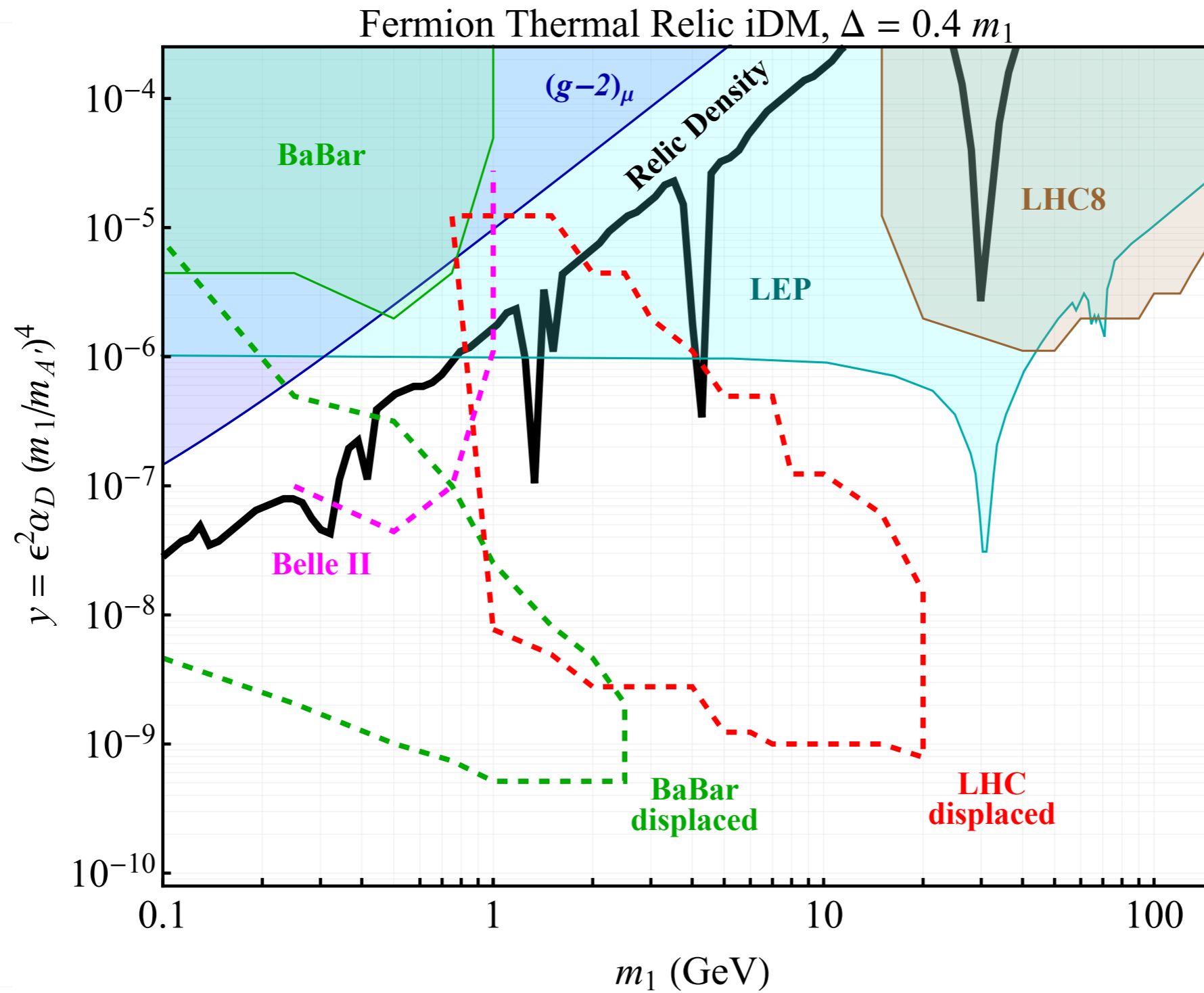
Conservative  $\alpha_D = 0.1$ ,  $m_{A'}/m_1 = 3$

# 30% Splitting



Conservative  $\alpha_D = 0.1$ ,  $m_{A'}/m_1 = 3$

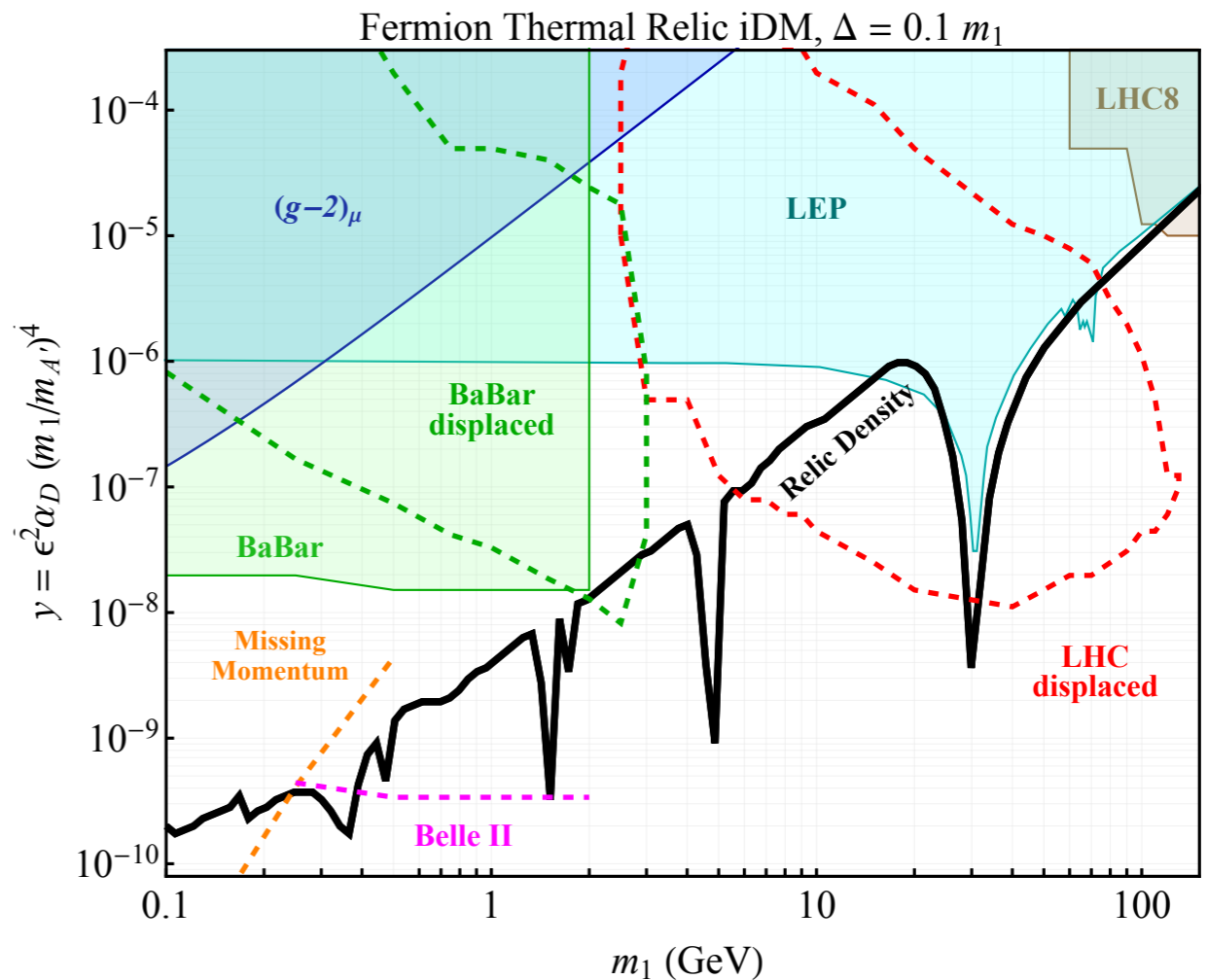
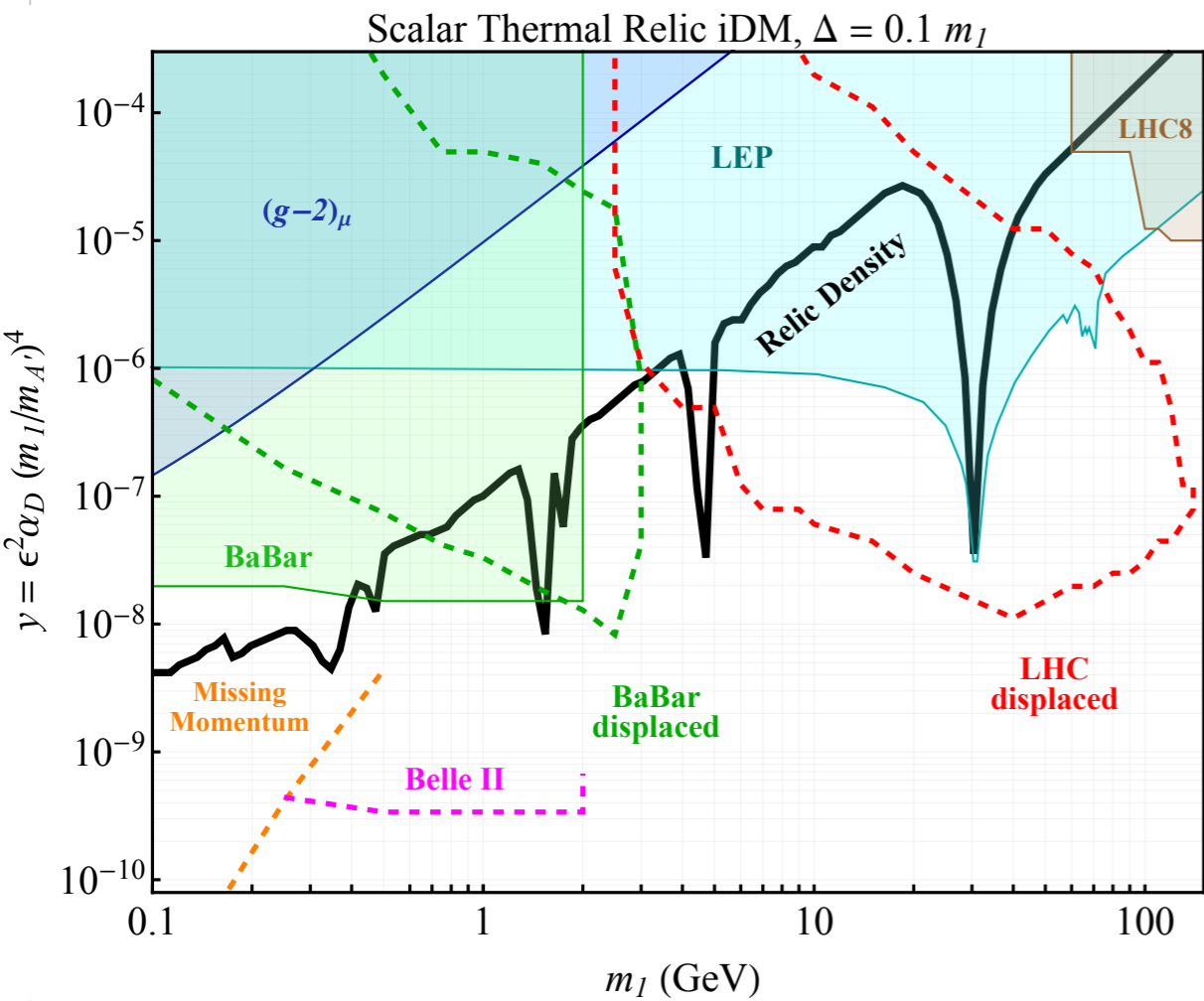
# 40% Splitting



Conservative  $\alpha_D = 0.1$ ,  $m_{A'}/m_1 = 3$



# Scalar vs. Fermion



$$\alpha_D = 0.1, \quad m_{A'}/m_1 = 3$$

$$\sigma v_{\text{scalar}} \sim \sigma v_{\text{fermion}} v^2$$

# Concluding Remarks

## **Inelastic DM: simple, viable, & needs colliders**

- **Clear, bounded parameter space:**  
conservatively test thermal iDM for  $O(1)$  splittings, comparable masses, small kinetic mixing
- **Cosmology @ Colliders** (jet/photon+MET +displaced dimuons) can close nearly all remaining thermal territory w/ LHC run 2 and BaBar
- **Future Challenge:** Exploit direct and/or resonant production mode (no jet/photon) LHCb Triggerless readout?...

# LHC Backgrounds

## Leptons from photon conversion in detector

$$pp \rightarrow j\gamma Z \rightarrow j\gamma(Z \rightarrow \nu\nu), \quad \sigma \approx 100 \text{ fb}$$

### Reduction Strategy

- Veto (leptons point to detector region)
- Veto (strict lepton isolation)
- Veto (dilepton invariant mass near  $\sim 0$ )
- Demand muons, reduce conversion prob.  
 $(m_e/m_\mu)^2 \approx 10^{-5}$

**Verdict: Very Small**

# LHC Backgrounds

## dileptons from displaced QCD

Difficult to calculate, but (theorist) estimate w/

- QCD event w/ hard jet + 2 muons
- Muon displacement 1 cm - 30 cm
- Point of closest approach < 1 mm

Total prob.  $\sim 10^{-7} \implies \sigma_{\text{QCD,BG}} < 100 \text{ fb}$

All this is before demanding large MET

**Verdict: Probably Very Small**

**Similar argument for j + W/Z BG**

# LHC Backgrounds

## Pile Up

### High Impact-parameter muons from other vertex

- Signal muons highly collimated from decay of boosted particle
- Dimuon momentum points back to primary vertex
- Same primary vertex as leading jet

**Verdict: Probably Very Small, Very Reducible**

# LHC Backgrounds

Jets +  $\tau\tau$

## Boosted taus decay to yield displaced muons

- Total cross section  $\sim 10$  fb
- Add muon decay penalty  $\sim 0.1$  fb
- Also need both to decay within  $\sim \mu\text{m}$
- Dimuon distribution will be different (single parent)

**Verdict: Very Small, Very Reducible**