A Prototype Atomic Interferometer for Fundamental Physics

David Morris Liverpool Particle Physics Atom Interferometry Experiment

On behalf of the collaboration

**University of Sussex IOP** 

March 2016



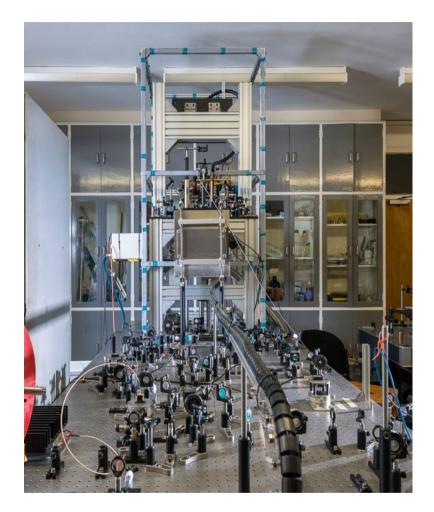






## Outline

- Atom interferometry concepts
- Motivation for building this detector
  - Developing the next generation of gravity sensors and gyroscopes for use in fundamental physics searches
  - Can dark energy be detection on the laboratory scale?<sup>[1]</sup>
- Status of the prototype
  - First step is to build a prototype atomic interferometer and use this to measure local gravity
  - Experiment upgrades
- Future development
  - Experiment goals



Experimental set-up of an Atomic interferometer

[Reference 1]

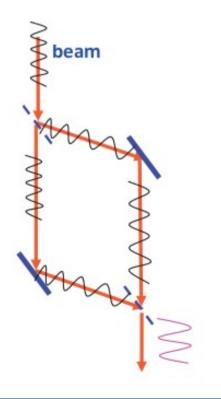
A terrestrial search for dark contents of the vacuum, such as dark energy, using atom interferometry Martin L. Perl arXiv:1101.5626

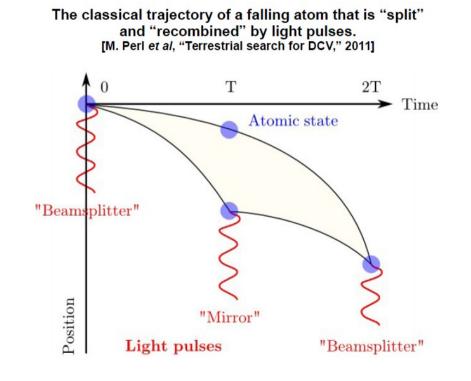
### **Atom Interferometry**

• Atomic interferometry - powerful tool for extremely precise measurements of fields.

- Laser Cooled atoms to nK temperatures
- Wave-like nature of atoms used to form a superposition which results in interference analogous to a laser interferometer

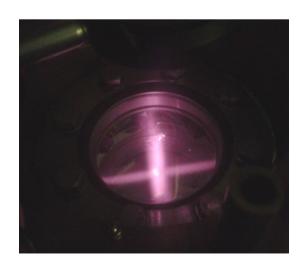
Interferometry

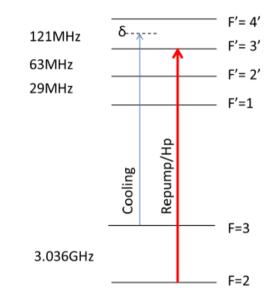


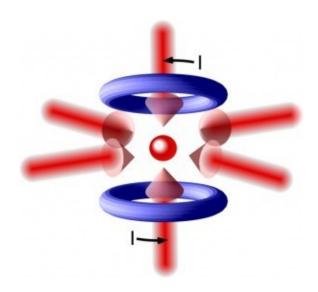


## **Magneto-Optical Trapping**

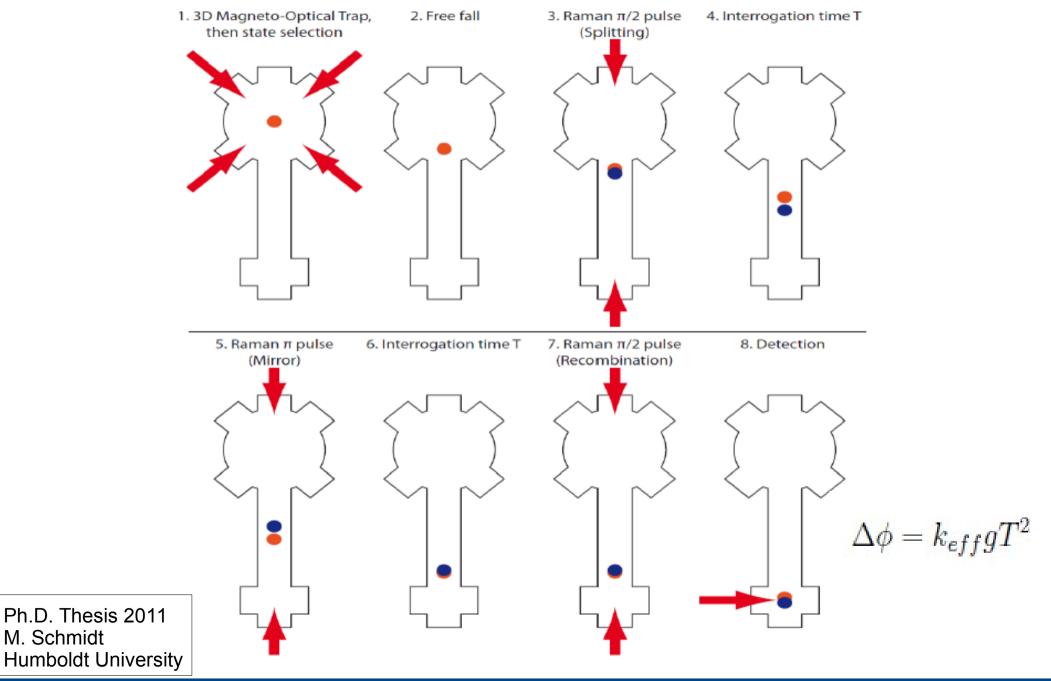
- Doppler cooling using 3 orthogonal pairs of counter propagating laser beams
- Magnetic field creates a position dependent trapping force
- Pumping laser keeps atoms in resonance with the cooling laser
- The atom trap is filled with <sup>85</sup>Rb from an atomic oven







## **Atom Interferometry**

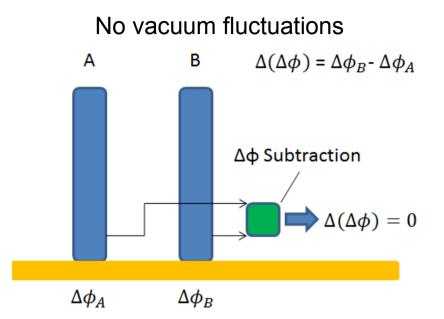


## **Experiment Aims**

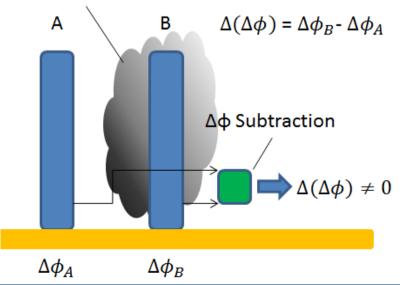
- Investigate the contents of the vacuum on the laboratory scale
- Two interferometers are set up side by side
- Measure a variation between the phase differences of the interferometers due to an external potential.

#### **Assumptions for detection**

- External potential has inhomogeneous spatial distribution.
- Exert a force on matter which is non-gravitational.



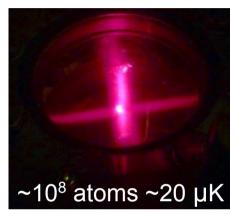
#### Phase difference from external potential



## Prototype Interferometer

#### **Preparation**

• 10<sup>8</sup> ultracold atoms in free-fall under gravity

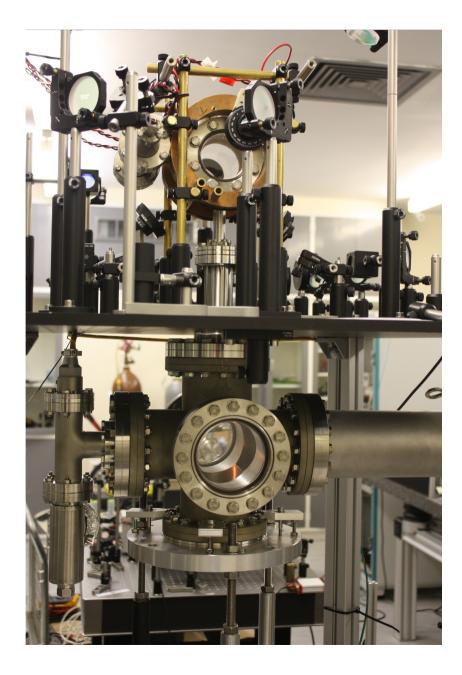


#### Interferometry Region

• Light-pulse interferometry sequence with Raman beams

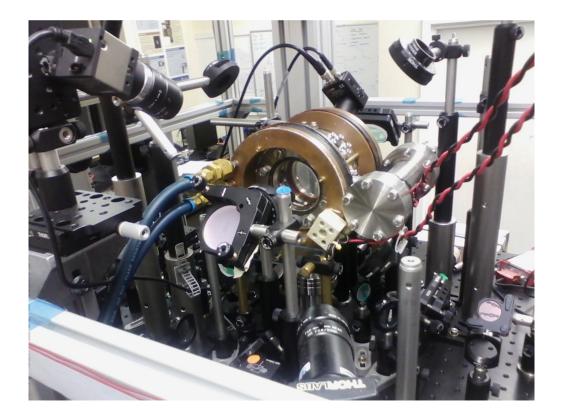
**Detection Region** 

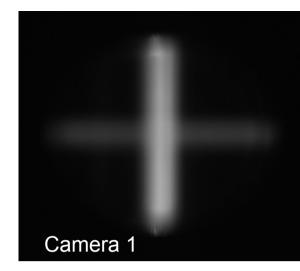
- Interference fringes
- Atom population ratio measurement are sensitive to phase difference

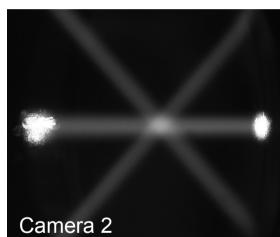


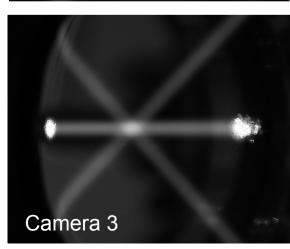
## **Beam Alignment**

- 3 quasi-orthogonal cameras used to align the MOT cooling beams
- Enables beams crossing to be placed within a pixel and set beams to be orthogonal.





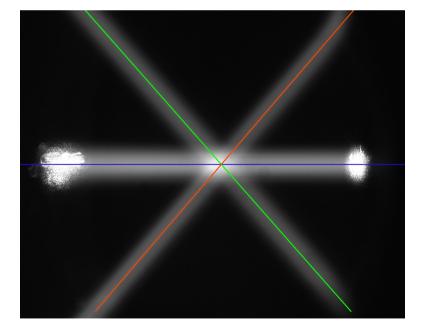


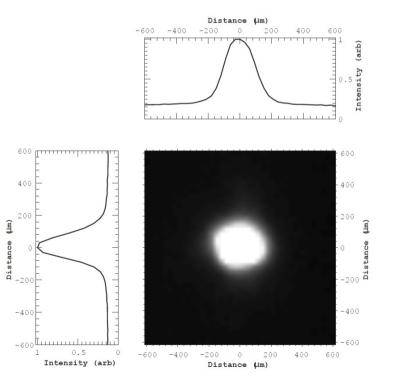


## **Beam Alignment**

• MOT cooling beams are set to all cross within 1 pixel on each camera and have a gradient of less then 1 pixel over 1000 pixels

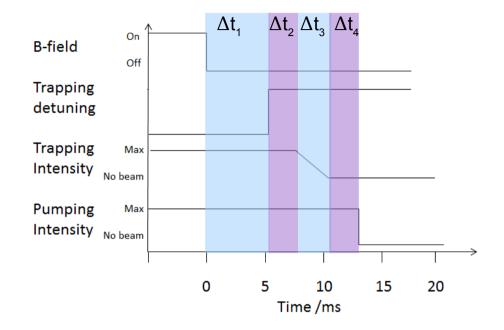
• After precision alignment a uniform MOT can be produced. With equal forces being applied by each of the 6 beams

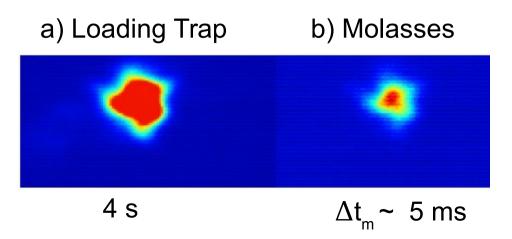




### Molasses Sequence

- Doppler limit is ~ 200 µk
- Additional cooling by using polarisation gradient cooling ~2  $\mu K$
- FGPA timing sequences varies magnetic field, laser power and detuning
- Parameters to vary
- $\Delta t_1$  eddies currents
- $\Delta t_{_2} \text{cooling}$
- $\Delta t_3^{}$  ramping
- $\Delta t_{_{4}}$  pump beam

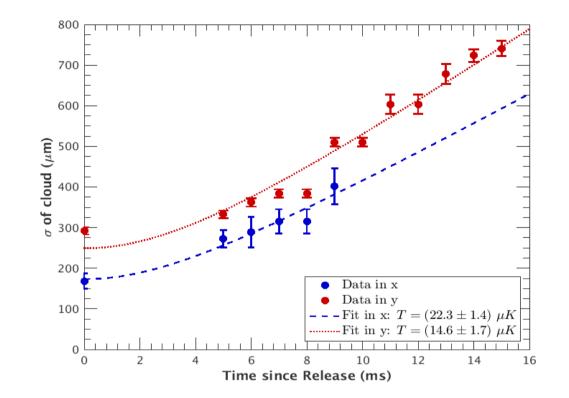


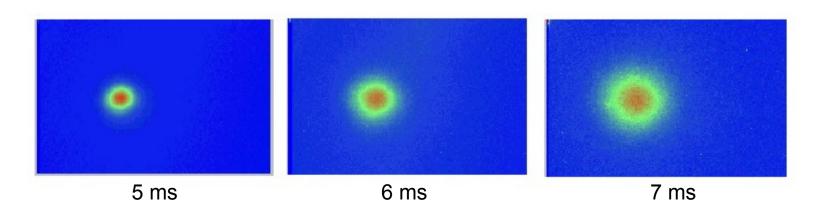


### **Temperature Measurement**

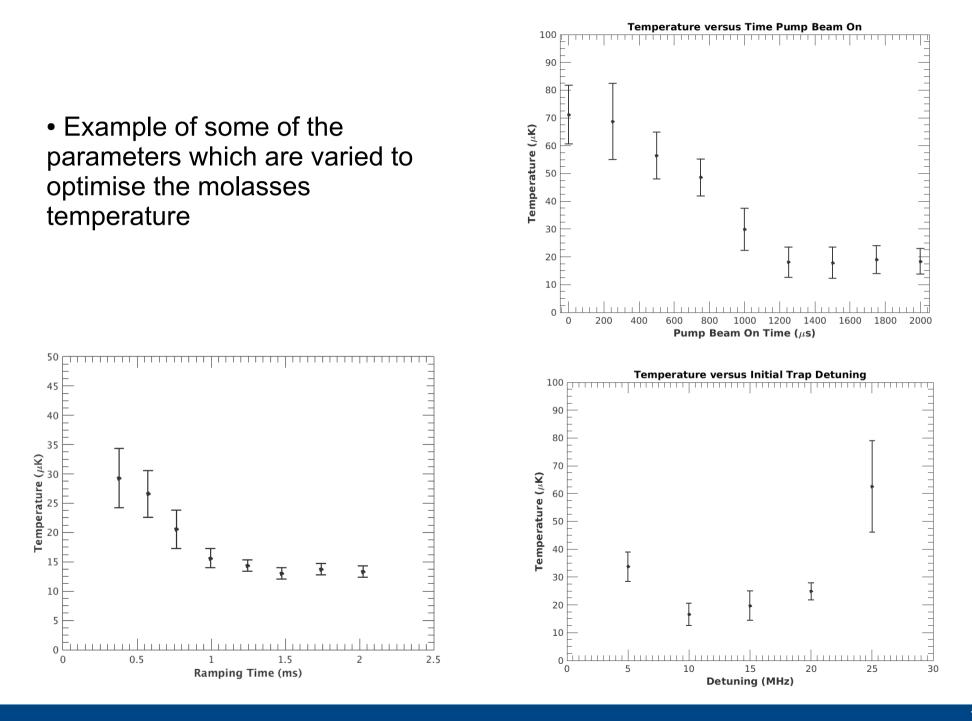
- Imaging expanding atomic cloud at varying times after release from molasses
- Radius of cloud is related to velocity and therefore temperature

$$\sigma_x(t) = \sqrt{\sigma_x^2(0) + \frac{k_{\rm B}T_x}{m}t^2}.$$



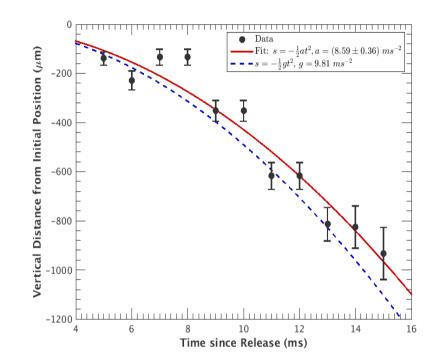


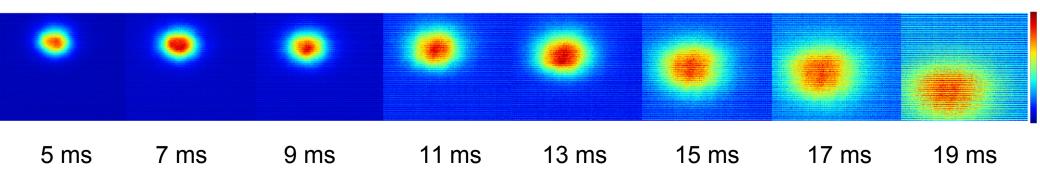
#### **Temperature Optimisation**



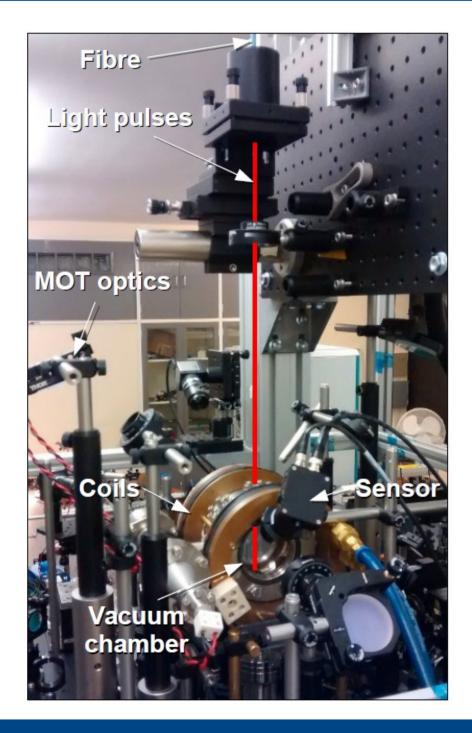
#### **Atom Drop**

- Sufficient thermal energy removed from atoms
- See the atom cloud influenced by the Earth's gravitational field
- Another step closer to being able to do atomic interferometry

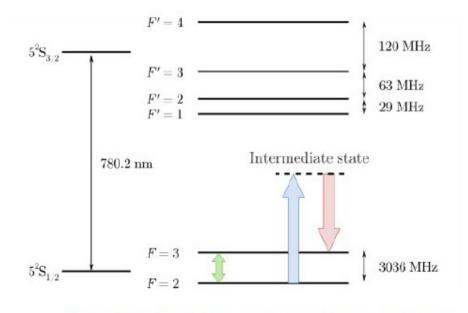




#### **Raman Beams**



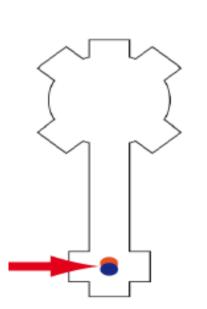
- Delivered by a fibre
- Two photon Raman transition, separated by 3 GHz
- "π-pulse" = "mirror" pulse
- "π/2-pulse" = "beamsplitter" pulses (half duration)



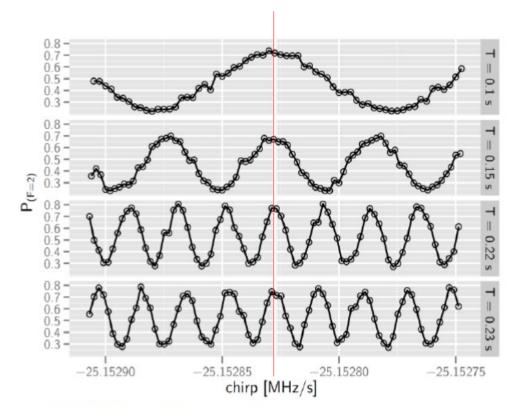
Rubidium-85 atomic energy level fine structure

### Measurement of Gravity

- Varying the chirp rate of the Raman beams traces out the interferometer fringes
- Changing the time T between Raman pluses traces out different fringe patterns
- Where the fringes meet corresponds to the chirp rate required to cancel out the Doppler shift due to gravity



Detection



A Mobile Atom Interferometer for High-Precision Measurements of Local Gravity, Alexander Senger

## **Future Plans**

- Finishing commissioning in the coming months
  - increase the power in the raman beams
- Measurement of local gravity as a benchmark
- Optimisation of device
- Implement double interferometer



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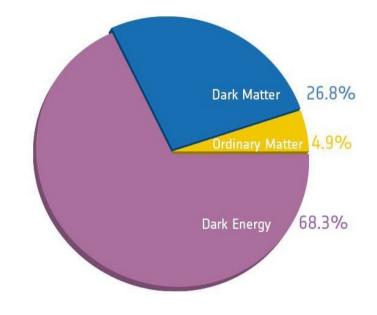




## **Additional Slides**

## Nature of Dark Energy

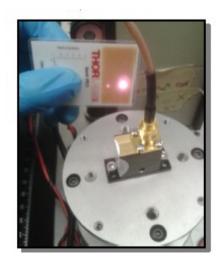
- Cosmological observations indicate 68% of the universe is dark energy.
- Present theory offers no fundamental understanding of the nature of dark energy
- Dark energy has a small but non-zero density 1.67x10<sup>-27</sup> kg m<sup>-3</sup>. Is this measurable on a terrestrial scale?

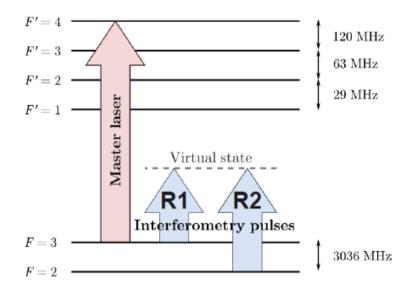


Planck Results arXiv:1303.5062

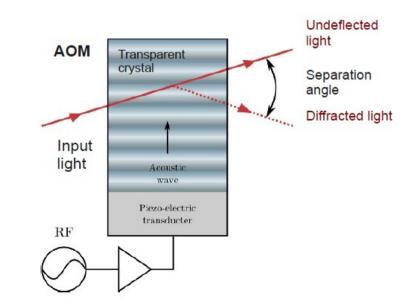
# **Frequency Control System**

- Many frequencies required for the atom interferometer – from only two lasers
- Optical circuit generates all required frequencies from and acousto-optical modulators (AOM)
- Interferometry requires a two-photon transition
- AOM's allow the lasers to be vary in frequency and intensity





Rubidium-85 hyperfine structure



## **Raman Transitions**

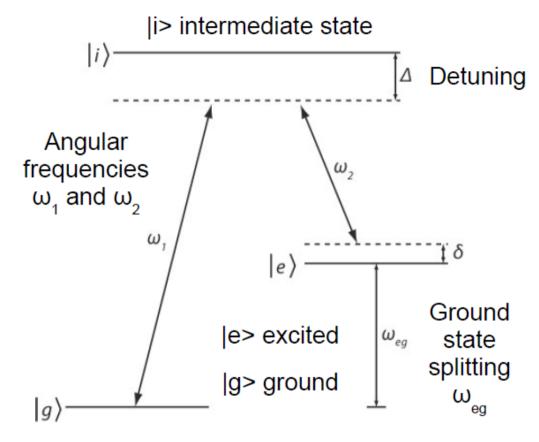
$$P_e(\tau) = \frac{\Omega_{\text{eff}}^2}{\Omega_{\text{eff}}^2 + (\delta_{12} - \delta^{AC})^2} \sin^2 \left( \sqrt{\Omega_{\text{eff}}^2 + (\delta_{12} - \delta^{AC})^2} \cdot \frac{\tau}{2} \right)$$

Excited-state population  $P_e$  as a function of Rabi frequency  $\Omega$  and detuning  $\delta$ . [Schmidt thesis, Eq. 2.55]

"For a resonant Raman process the frequency detuning  $\delta \approx 0$  and the detuning  $\Delta$  from the intermediate state remains large, so that excitation by the single-photon absorption is negligible in comparison to the coherent transfer from |g> to |e>" [C. J. Foot Atomic Physics p209]

$$\omega_1, k_1$$
 Atom  $\omega_2, k_2$ 

Recoil velocity =  $2hk_{eff}/M$ where  $k_{eff} = k_1 - k_2$ 



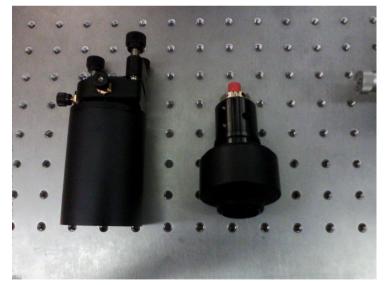
Level scheme of three-level Raman transitions, simplified to exclude AC Stark shifts

$$\omega_{\text{eff}} = \omega_1 - \omega_2 \approx \omega_{\text{eg}}$$

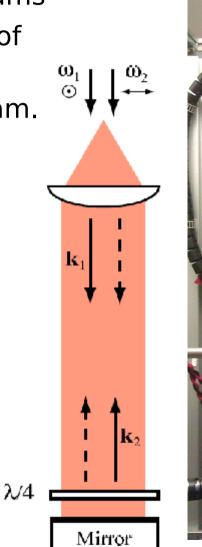
[Schmidt thesis, Fig. 2.1]

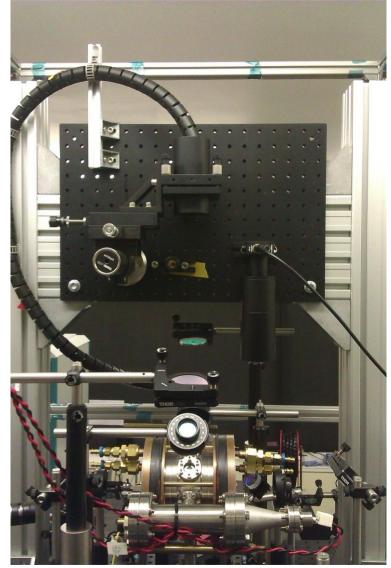
## Raman Beam Delivery

- Fibre delivers the Raman beams
- 10 m fibre cleans up modes of beam.
- Collimated into 10.8 mm beam.

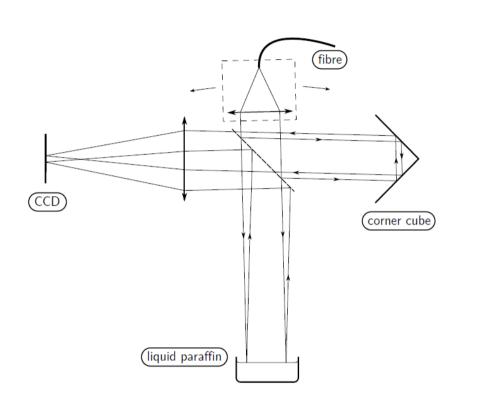


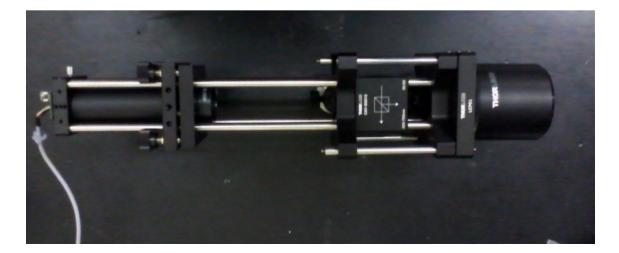
Collimator mount





# **Precision Alignment**



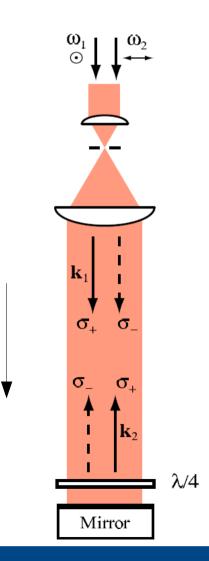


Using precision optics with

- Extremely flat surfaces,  $\lambda/20$ 

g

- Small beam deviations, 10 urads



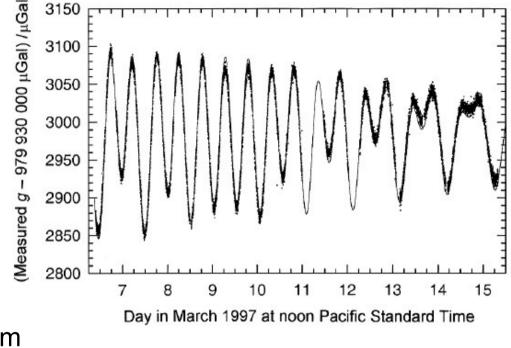
# Improving Sensitivity

Using atom interferometers to measure tidal effects:

 $\sigma_{
m g}$  need to be 10'  $\Delta \phi = k_{eff} g T^2$  =

If long data taking takes several days to reach 10<sup>7</sup> than daily tidal variations would not be observable.

- Require having large phase shift  $\Delta \Phi$ :
- T is time of flight for the atom cloud
- Prototype under construction height  $\sim 1 \text{ m}$
- T<sup>2</sup> is proportional to h,
- h = 10 m, approx 10 x improvement in  $\Delta \Phi$ ,



1 μGal = 10-8ms-2

Increased atom number gives increased statistical power Increased drop length gives greater sensitivity of each

measurement

### **Polarisation Gradient Cooling**

[Himsworth thesis, "Coherent Manipulation ... " 2009]

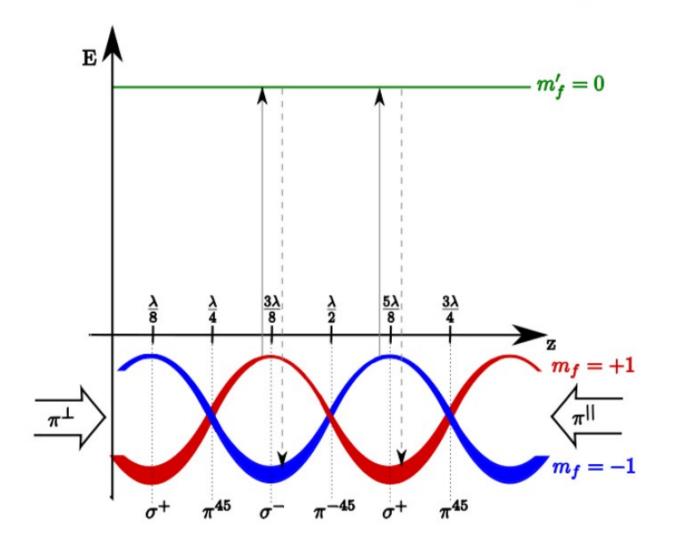
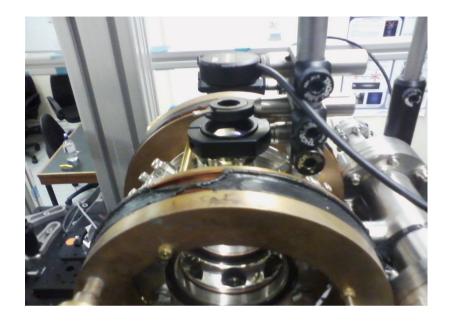


FIGURE 5.4: The Sysiphus effect of Sub-Doppler cooling. See text for details. The width of the lower  $m_F$  states (red and blue) representative of the state populations.

## **Atom Number**



Power: Atom No. Solid angle ,Scatting rate Photon energy

#### Calculation

- •Measure flux for some small area
- •Extrapolate  $4\pi$  flux
- •Subtract laser fluorescence background
- •Take ratio: divide by single atom flux (scattering rate \* photon energy)
  - i.e. N = total flux / single atom flux

#### Time of flight measurement

$$\sigma_x(t) = \sqrt{\sigma_x^2(0) + \frac{k_{\rm B}T_x}{m}t^2}.$$

Thesis Kevin J. Weatherill Durham Atomic physics A CO2 Laser Lattice Experiment for Cold Atoms

2007