

# Getting ready for LHC run 2: hunting dark matter at colliders

Adam Martin (amarti41@nd.edu)

University of Notre Dame



# Outline

- 1.) Introduction to Notre Dame High-Energy Physics
- 2.) Hunting for (supersymmetric) dark matter at colliders

# University of Notre Dame





# University of Notre Dame



physics





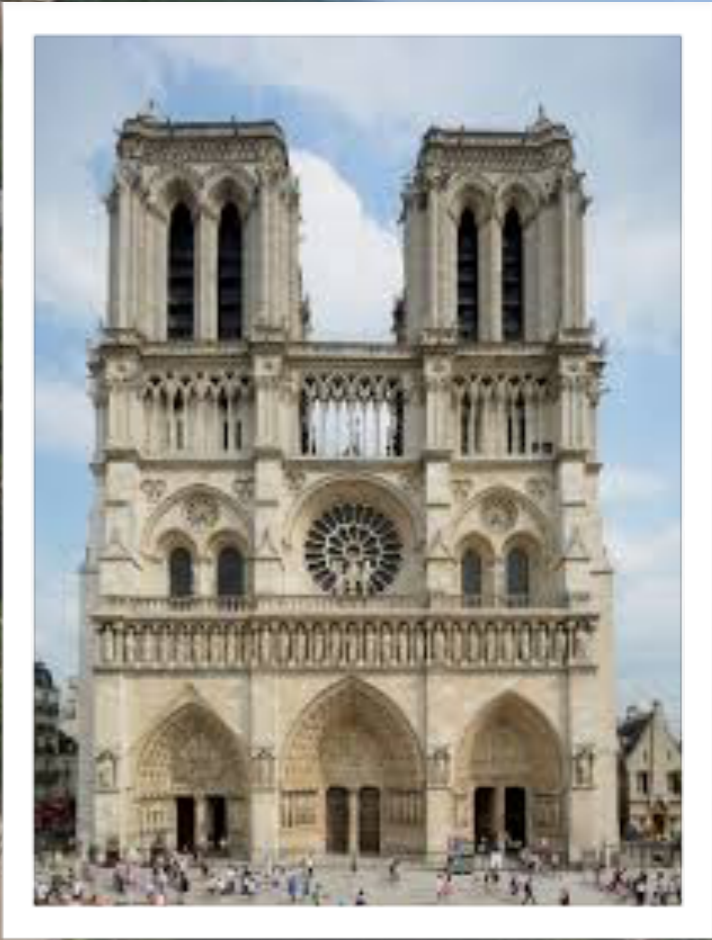










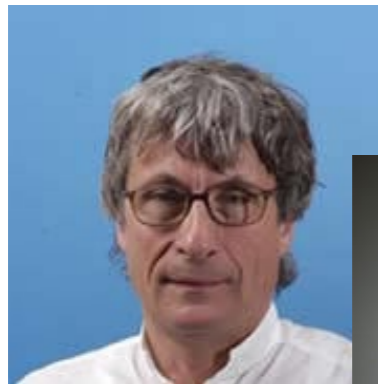




# Particle Physics at Notre Dame

theory

experiment



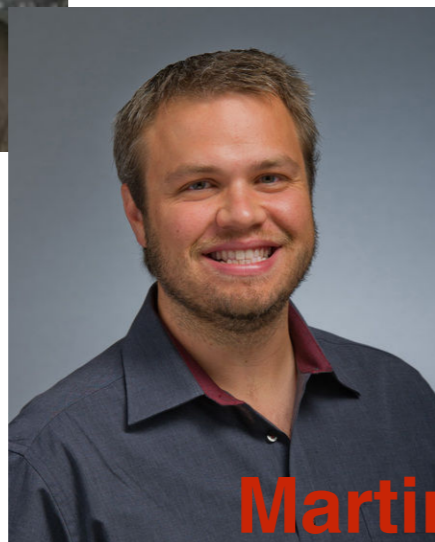
**Bigi**



**Delgado**



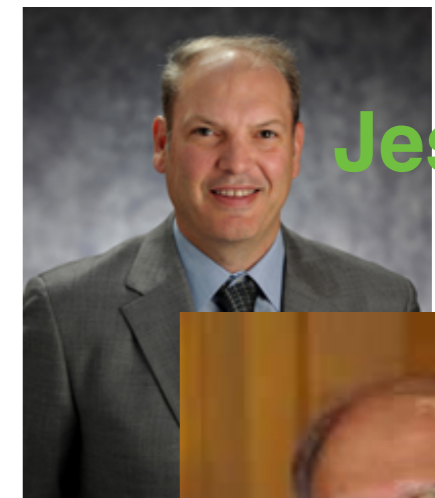
**Kolda**



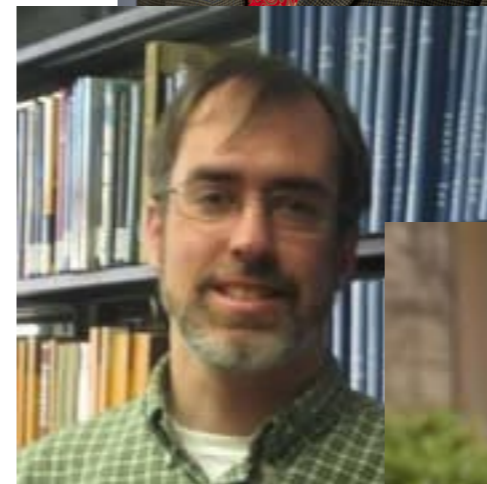
**Martin**



**Hildreth**



**Jessop**



**Lannon**



**Ruchti**



**Wayne**



**LoSecco**

# Particle Physics at Notre Dame

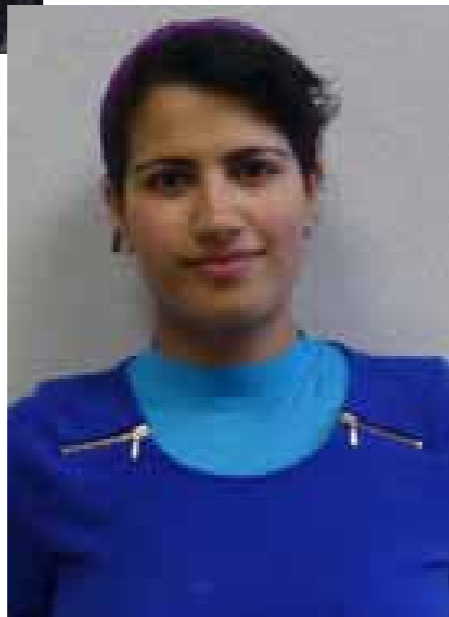
## students



*Alvarado, Carlos*



*Capdevilla, Rodolfo*



*Elahi, Fatemeh*



*Lehman, Landon*

## postdocs



Joe Bramante



Nirmal Raj



In addition to the topic I'll talk about  
there are several other ideas we're interested in

*inflation*

*compressed SUSY spectra*

*very light  $Z'$*

*simple MSSM extensions*

**Hilbert series**

*alternatives to freeze-out DM*

that I'd be happy to discuss

## Part II:

### Hunting Dark Matter at the LHC

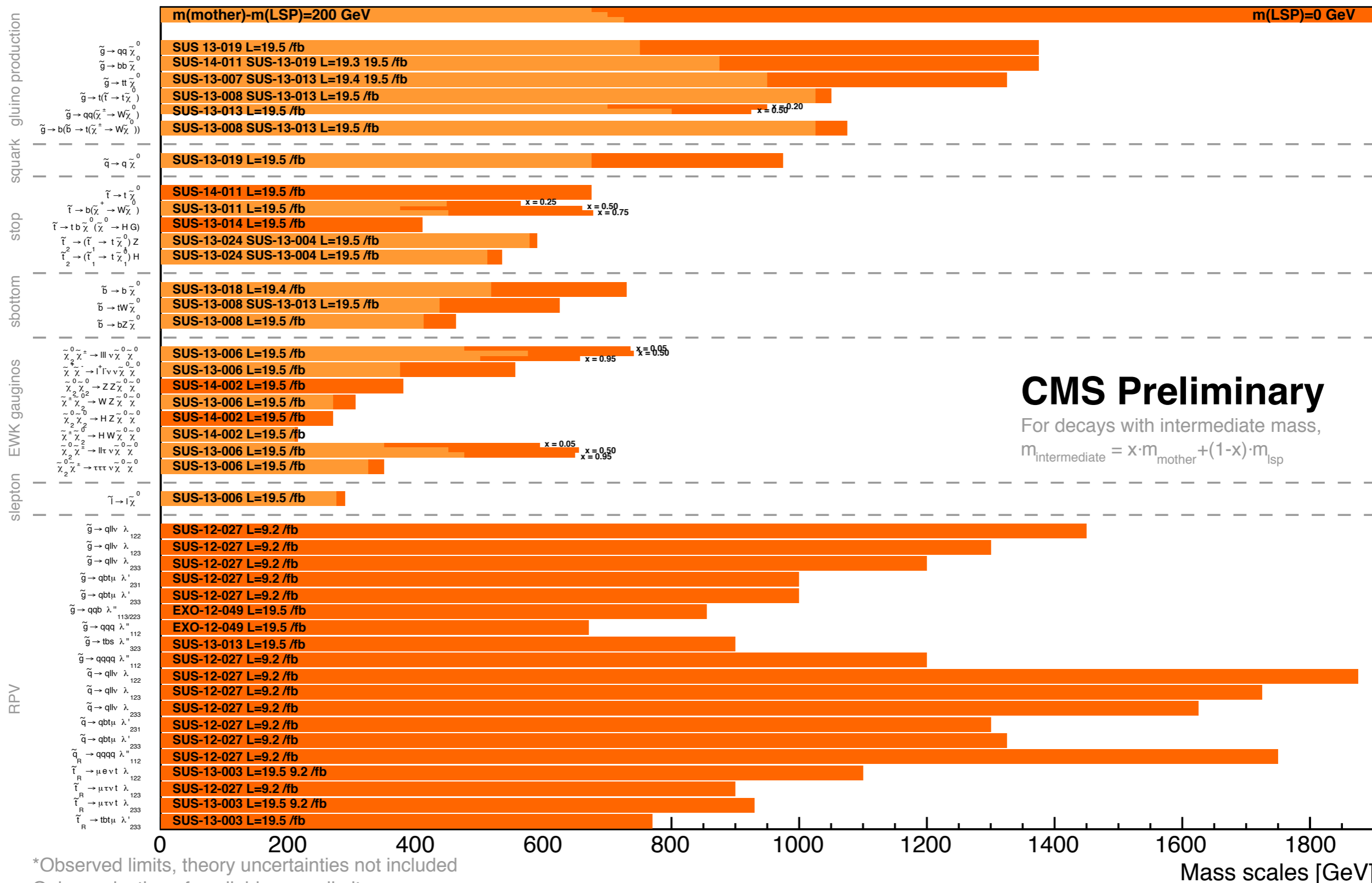
*specifically* = *superpartners of Higgs*  
*electroweakinos*      *boson, W, Z,  $\gamma$*



# Motivation

## Summary of CMS SUSY Results\* in SMS framework

ICHEP 2014

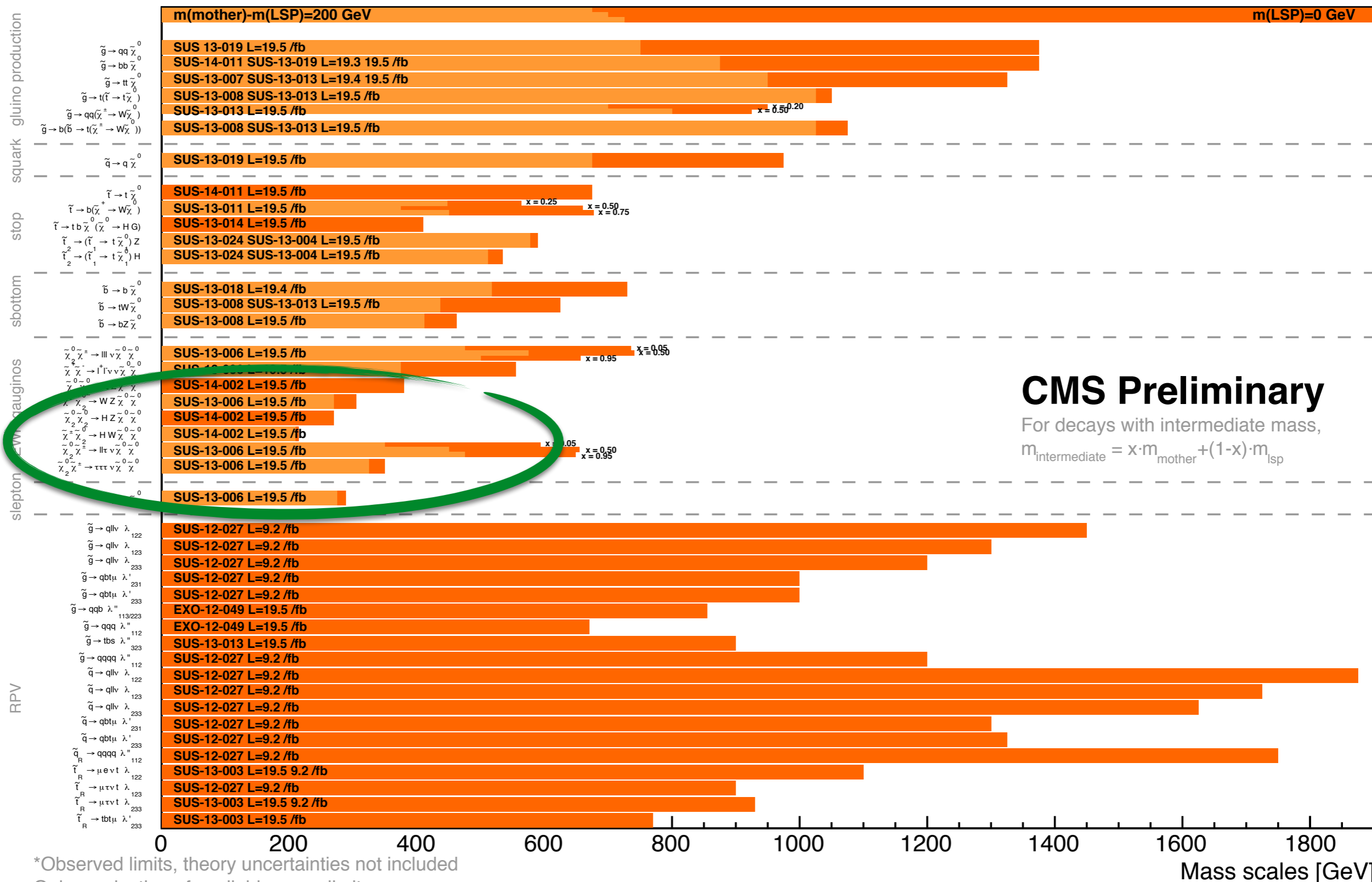




# Motivation

## Summary of CMS SUSY Results\* in SMS framework

ICHEP 2014



\*Observed limits, theory uncertainties not included

Only a selection of available mass limits

Probe \*up to\* the quoted mass limit

Mass scales [GeV]

# Why electroweakinos?

electroweakinos = admixtures of Higgsino, Bino, Wino

charginos (Higgsinos, Winos):  $(\tilde{W}^+, \tilde{H}_u^+, \tilde{W}^-, \tilde{H}_d^-)$

$$\mathbf{X} = \begin{pmatrix} M_2 & gv_u \\ gv_d & \mu \end{pmatrix} = \begin{pmatrix} M_2 & \sqrt{2}s_\beta m_W \\ \sqrt{2}c_\beta m_W & \mu \end{pmatrix}$$

neutralinos (all 4):  $\psi^0 = (\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0)$

$$\mathbf{M}_{\tilde{N}} = \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix}$$



# Why electroweakinos?

electroweakinos = admixtures of Higgsino, Bino, Wino

charginos (Higgsinos, Winos):  $(\tilde{W}^+, \tilde{H}_u^+, \tilde{W}^-, \tilde{H}_d^-)$

**X in the limit  $m_{\tilde{Q}}, m_{\tilde{L}} \gg m_\chi$ :**  
**phenomenology set by 4 parameters:**  
 **$M_1, M_2, \mu, \tan\beta$**

neutralinos (all 4):  $\psi^0 = (\tilde{B}, \tilde{W}^0, \tilde{H}_d^0, \tilde{H}_u^0)$

$$\mathbf{M}_{\tilde{N}} = \begin{pmatrix} M_1 & 0 & -c_\beta s_W m_Z & s_\beta s_W m_Z \\ 0 & M_2 & c_\beta c_W m_Z & -s_\beta c_W m_Z \\ -c_\beta s_W m_Z & c_\beta c_W m_Z & 0 & -\mu \\ s_\beta s_W m_Z & -s_\beta c_W m_Z & -\mu & 0 \end{pmatrix}$$

# Why electroweakinos?

Naturalness: Higgsino mass is related to the Z mass

$$\frac{1}{2}M_Z^2 = -m_{H_u}^2 - |\mu|^2$$

*weak scale*

*to avoid big cancellations,  
both of these should be weak  
scale (~100 GeV) too!*



# Why electroweakinos?

Naturalness: Higgsino mass is related to the Z mass

$$\frac{1}{2}M_Z^2 = -m_{H_u}^2 - |\mu|^2$$

*weak scale*

*to avoid big cancellations,  
both of these should be weak  
scale (~100 GeV) too!*

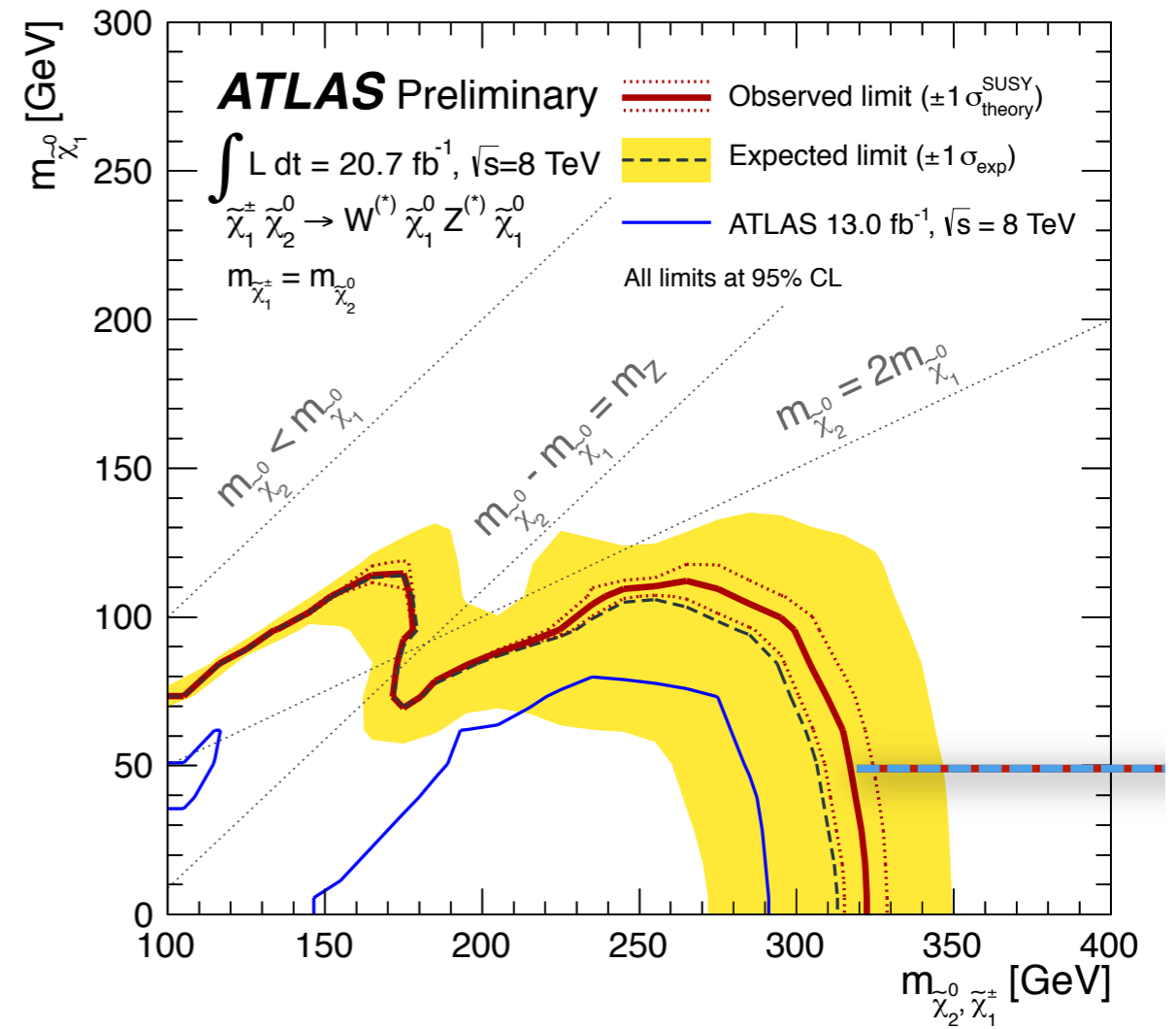
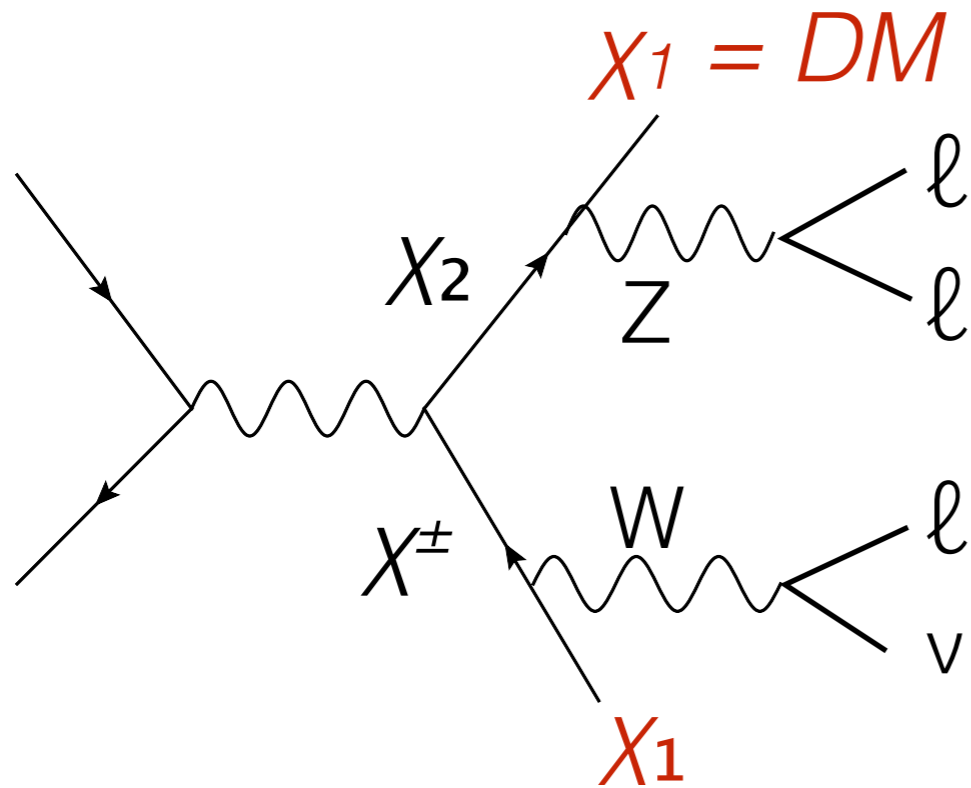
Dark Matter: lightest superparticle is stable, can be DM

right amount of DM achieved from admixtures  
of electroweakinos, mass ~ 100 GeV - TeV

“Well-tempered electroweakino” scenario

# Current electroweakino searches

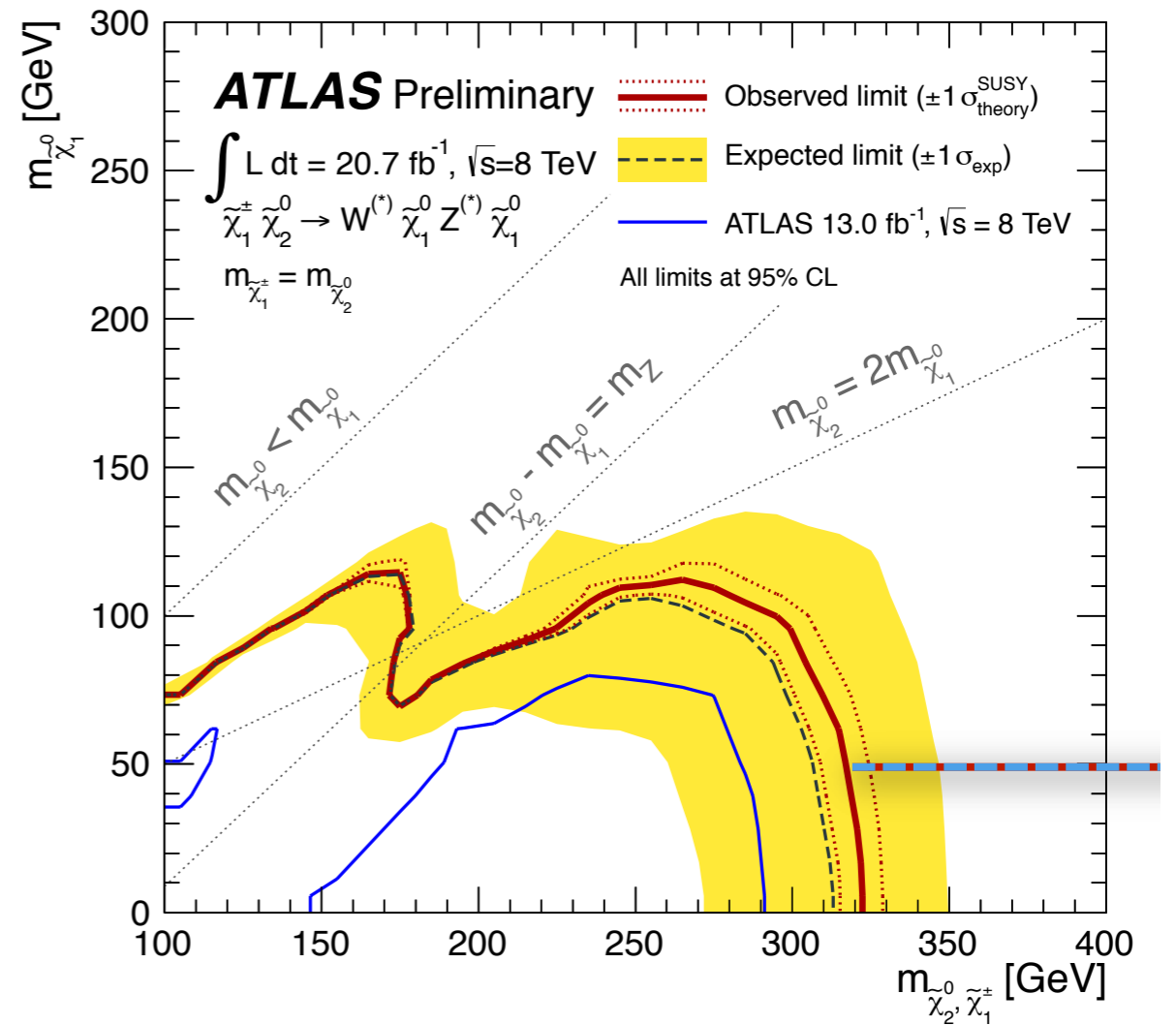
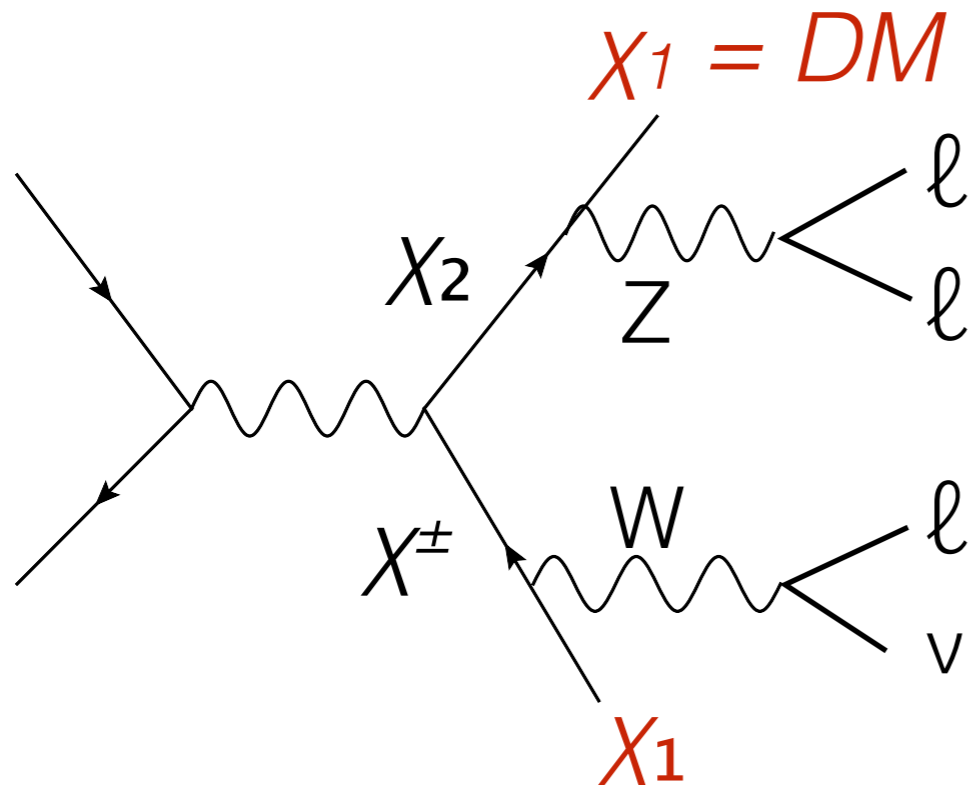
“tri-lepton” searches:





# Current electroweakino searches

“tri-lepton” searches:



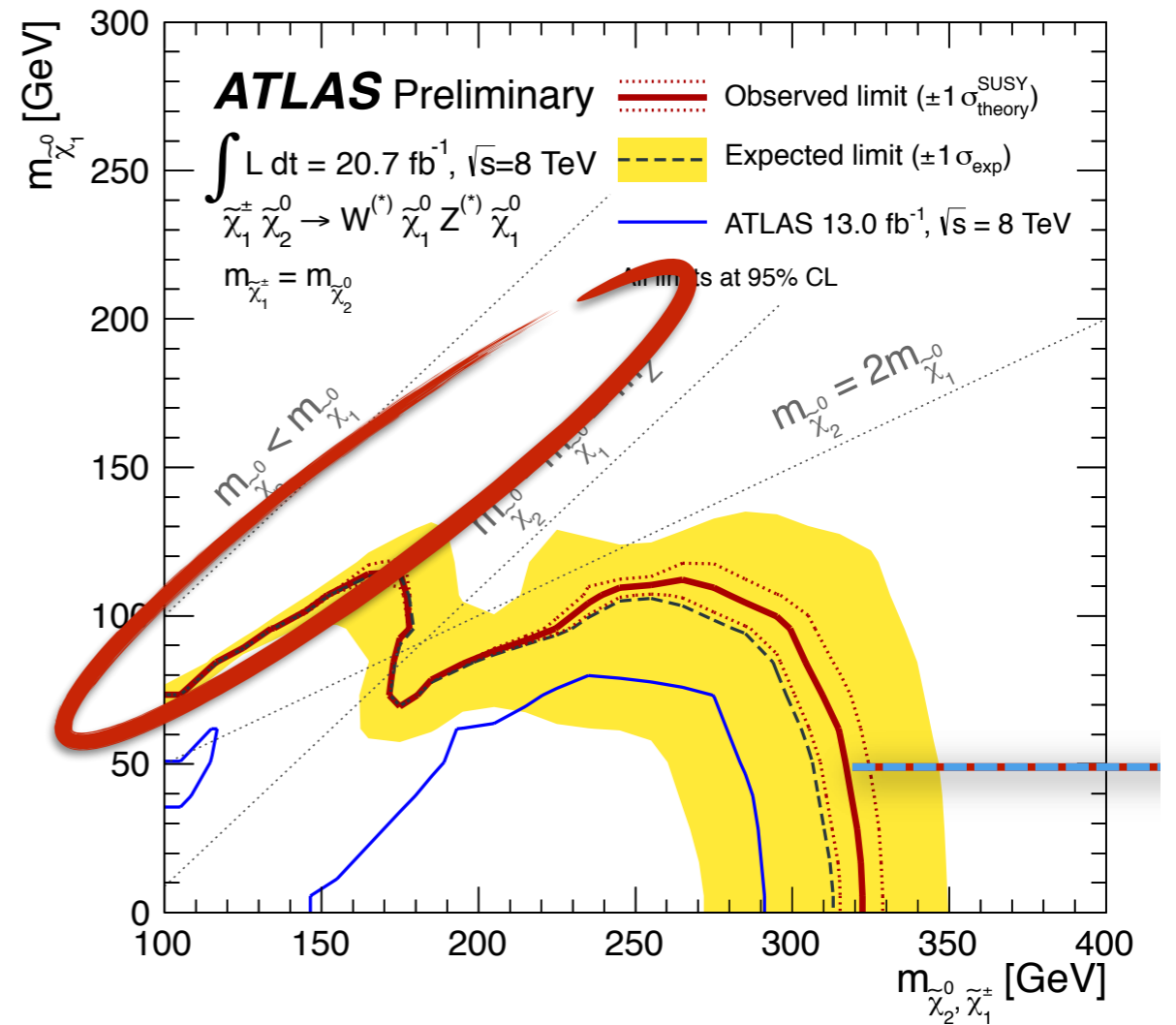
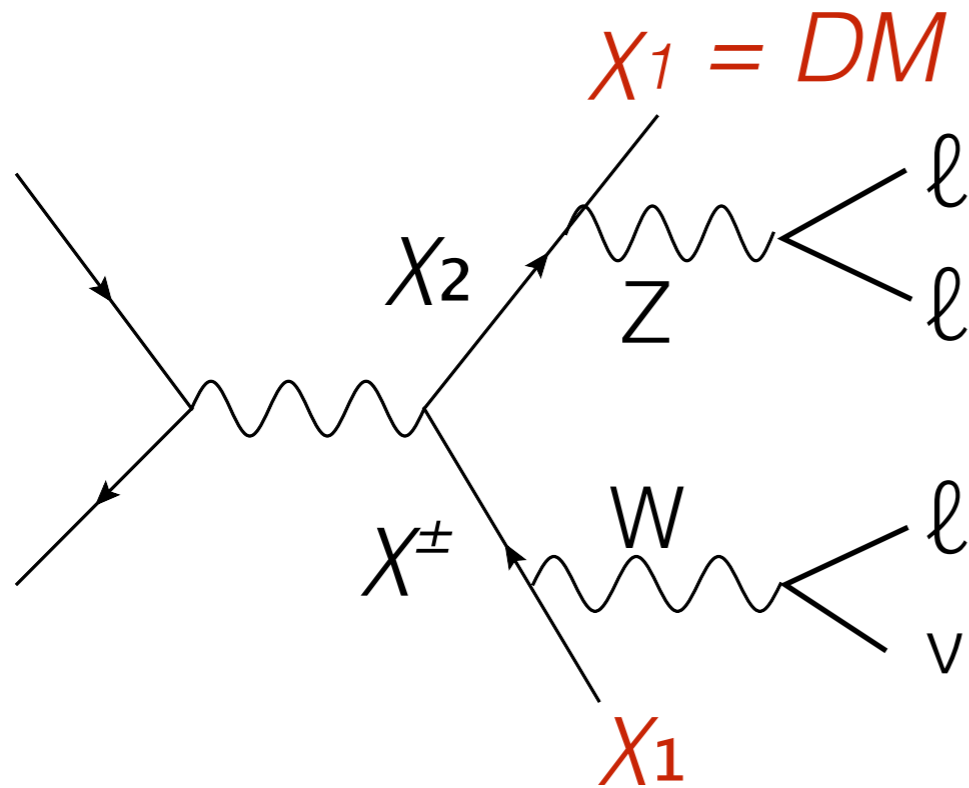
struggle when the states are nearly degenerate ( $m_{\tilde{\chi}_2} \sim m_{\tilde{\chi}_\pm} \sim m_{\tilde{\chi}_1}$ )

leptons become inefficient to trigger upon (25 GeV for single lepton trigger, 17 GeV/8 GeV for 2-lepton)...

Doesn't improve at 14 TeV

# Current electroweakino searches

“tri-lepton” searches:



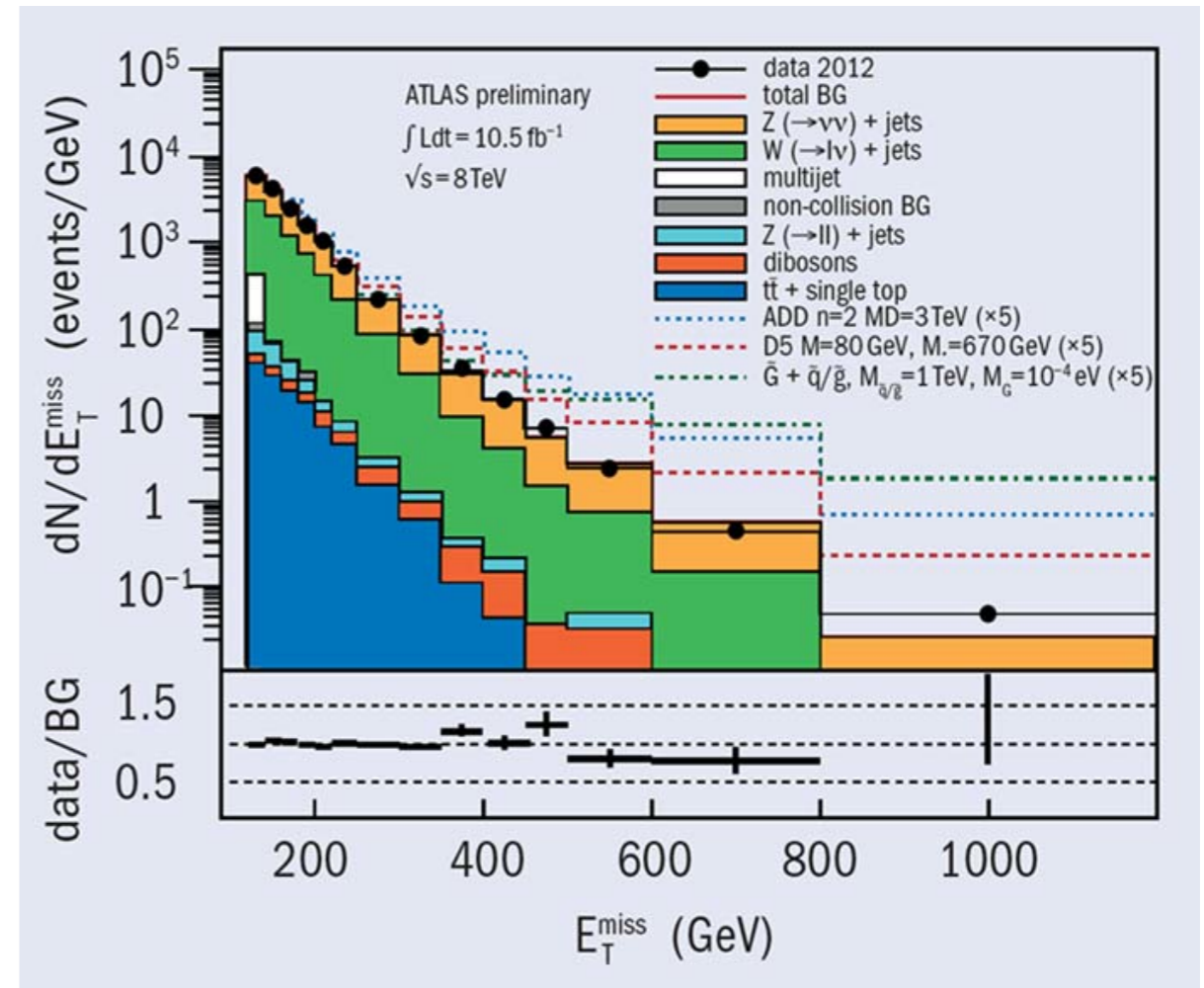
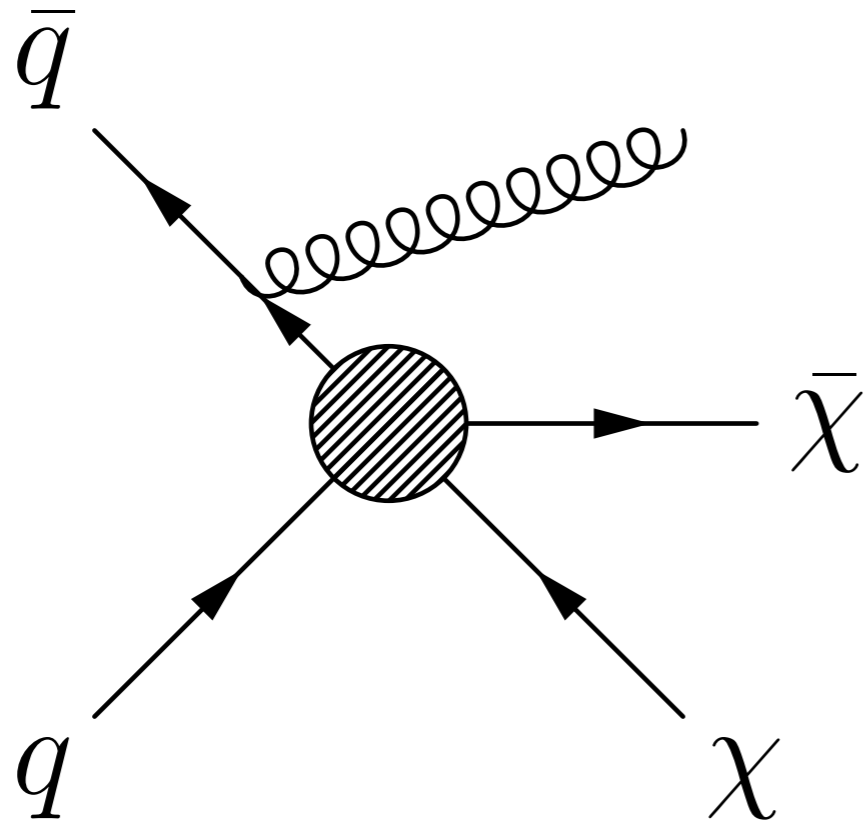
struggle when the states are nearly degenerate ( $m_{\chi_2} \sim m_{\chi_{\pm}} \sim m_{\chi_1}$ )

leptons become inefficient to trigger upon (25 GeV for single lepton trigger, 17 GeV/8 GeV for 2-lepton)...

Doesn't improve at 14 TeV



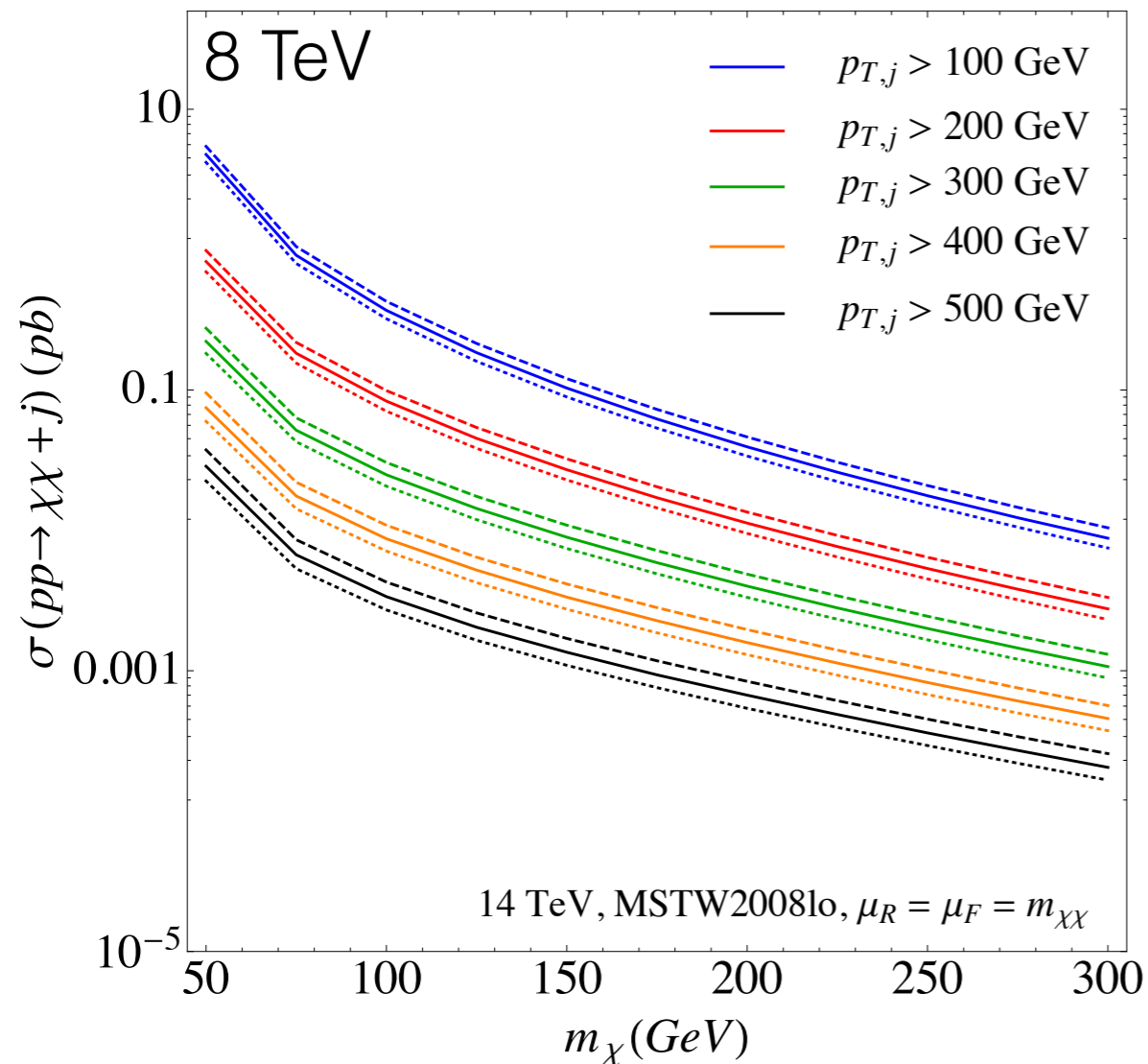
# Monojet searches



model independent, ideal for degenerate spectra

less to distinguish signal from background; limits degrade as the invisible-SM interaction becomes less contact-like

# Monojet searches



[Han, Kribs, AM, Menon 1401.1235]

for pure Higgsinos

recasting mono-jet searches on  
8 TeV data, limits are

$$m_\chi \gtrsim 80 \text{ GeV}$$

**no better than LEP II bound!**

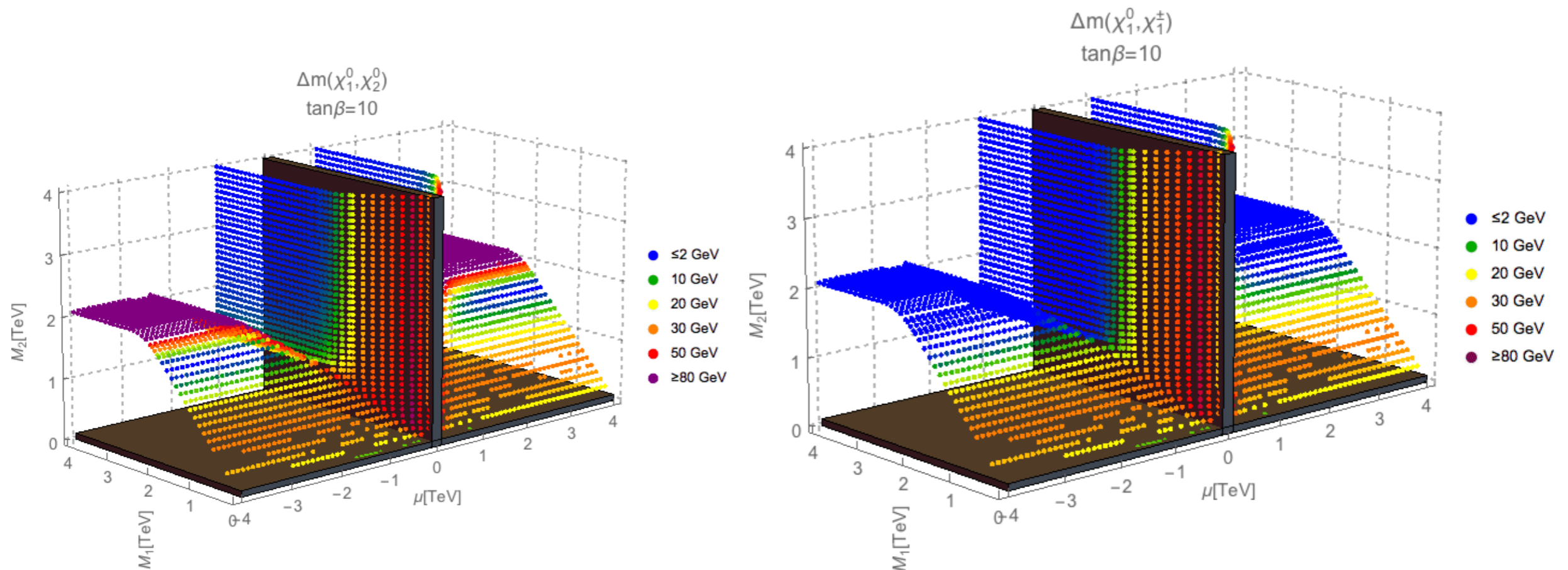
projected to 14 TeV, high L, limits still feeble



# Why electroweakinos?

LHC limits are weakest when states are nearly degenerate

but approximately-degenerate spectra are exactly what is required for well-tempered scenarios to fit DM abundance

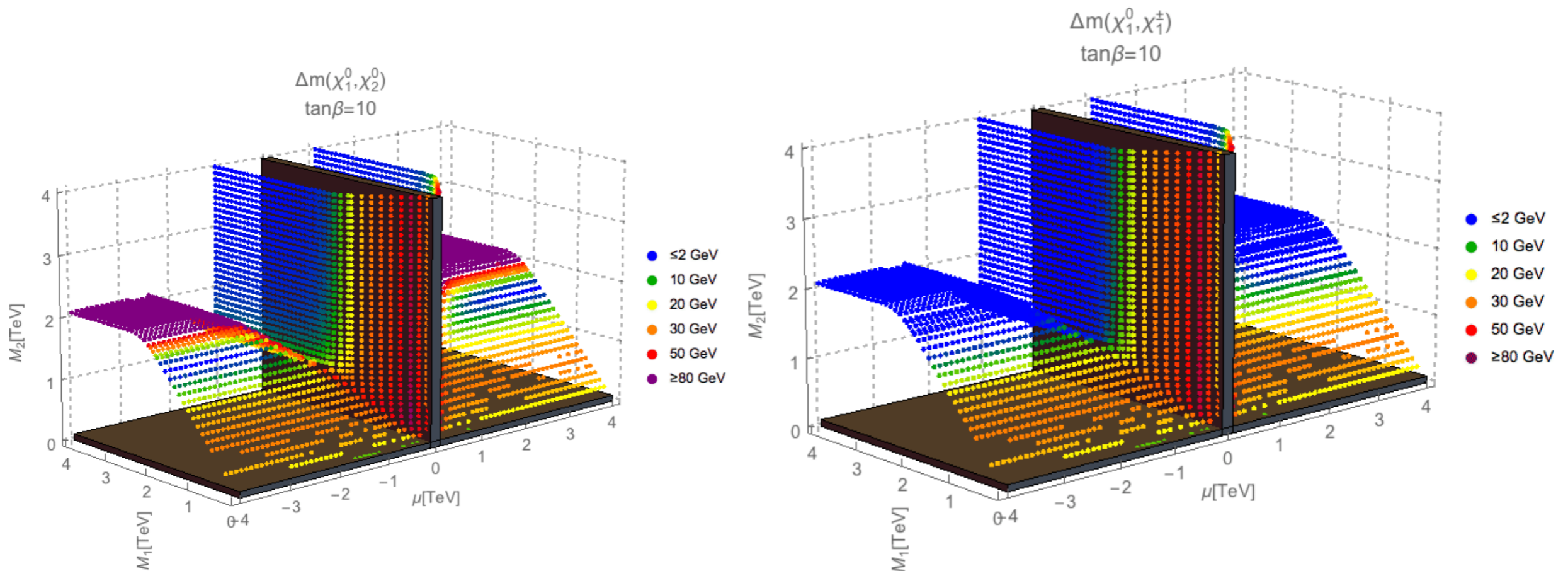


improved searches needed

# Why electroweakinos?

LHC limits are weakest when states are nearly degenerate

but approximately-degenerate spectra are exactly what is required for well-tempered scenarios to fit DM abundance



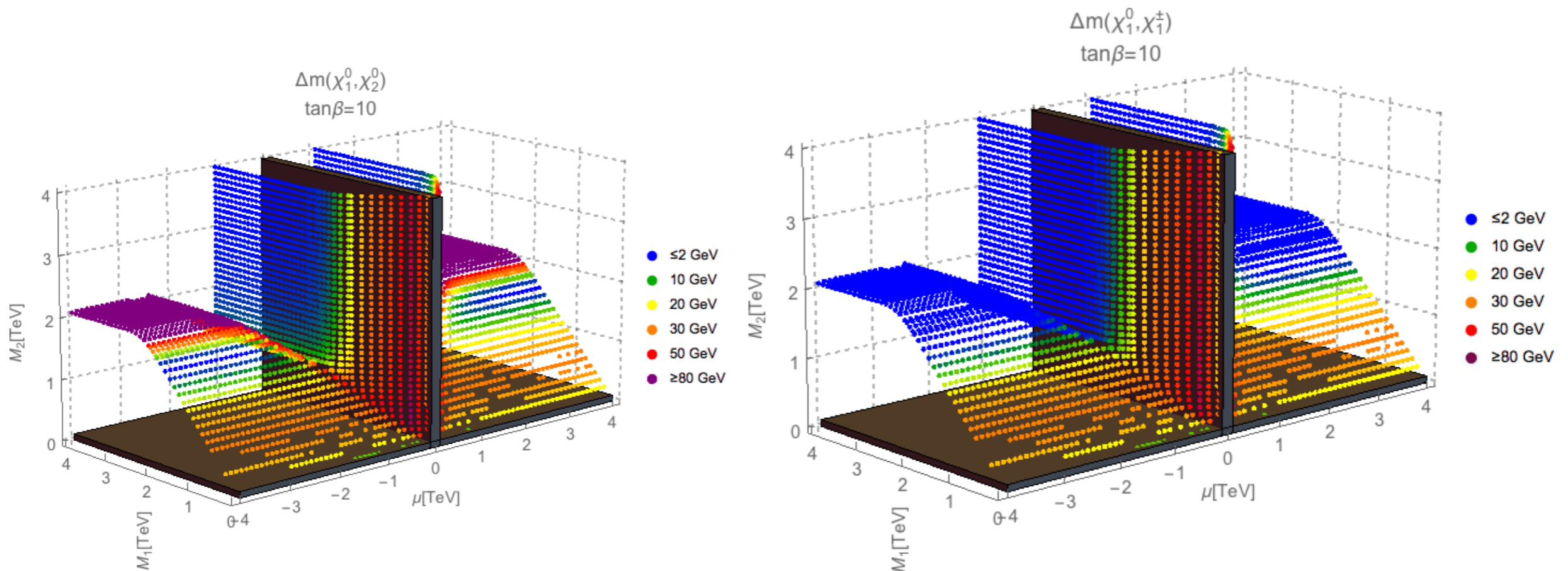
improved searches needed



# Why electroweakinos?

LHC limits are weakest when states are nearly degenerate

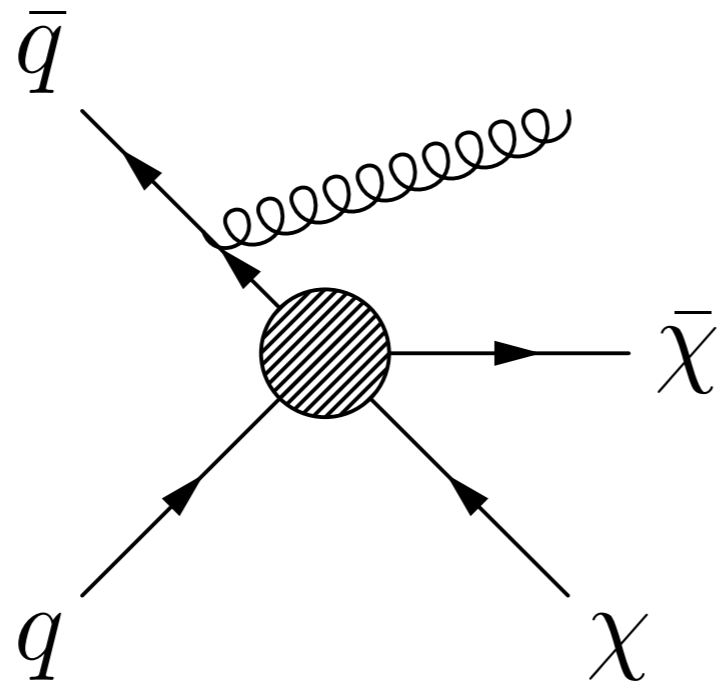
but approximately-degenerate spectra are exactly what is required for well-tempered scenarios to fit DM abundance



improved searches needed

# One possibility: Improve mono-jet searches

the problem with tri-lepton searches is triggering; solve this by requiring hard initial-state jet, as in mono-jet

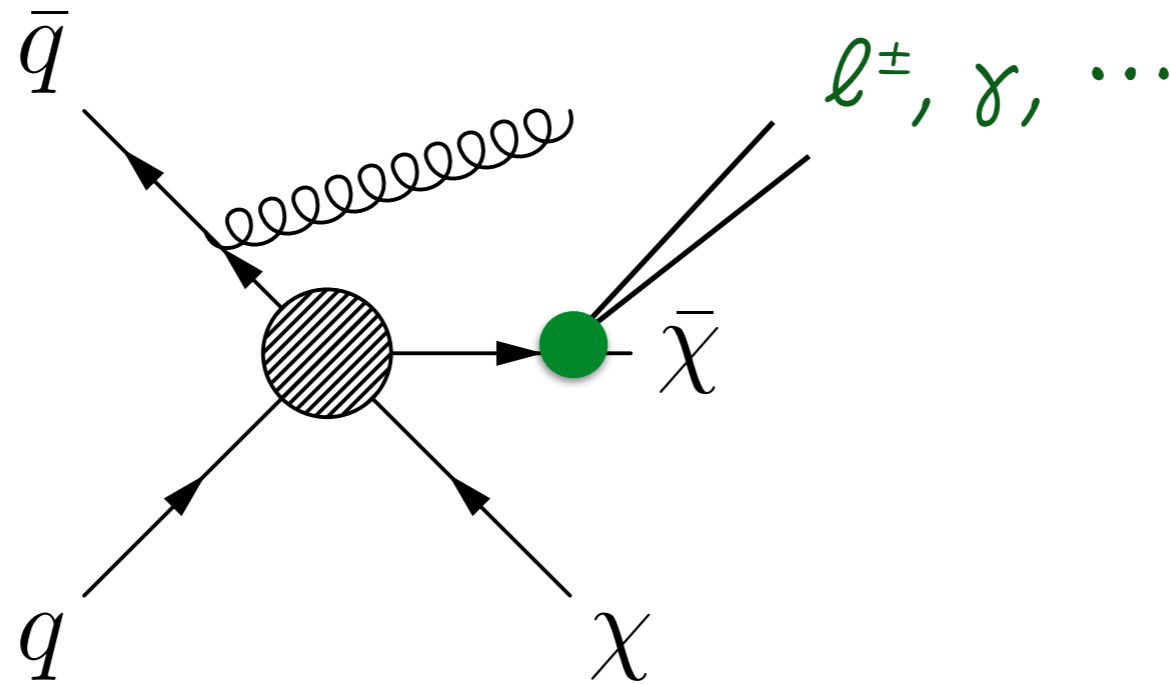


once triggered upon, event can be analyzed using lower, off-line thresholds ( $\gtrsim 7$  GeV for  $e^\pm$ ,  $\gtrsim 3$  GeV for  $\mu^\pm$ )



# One possibility: Improve mono-jet searches

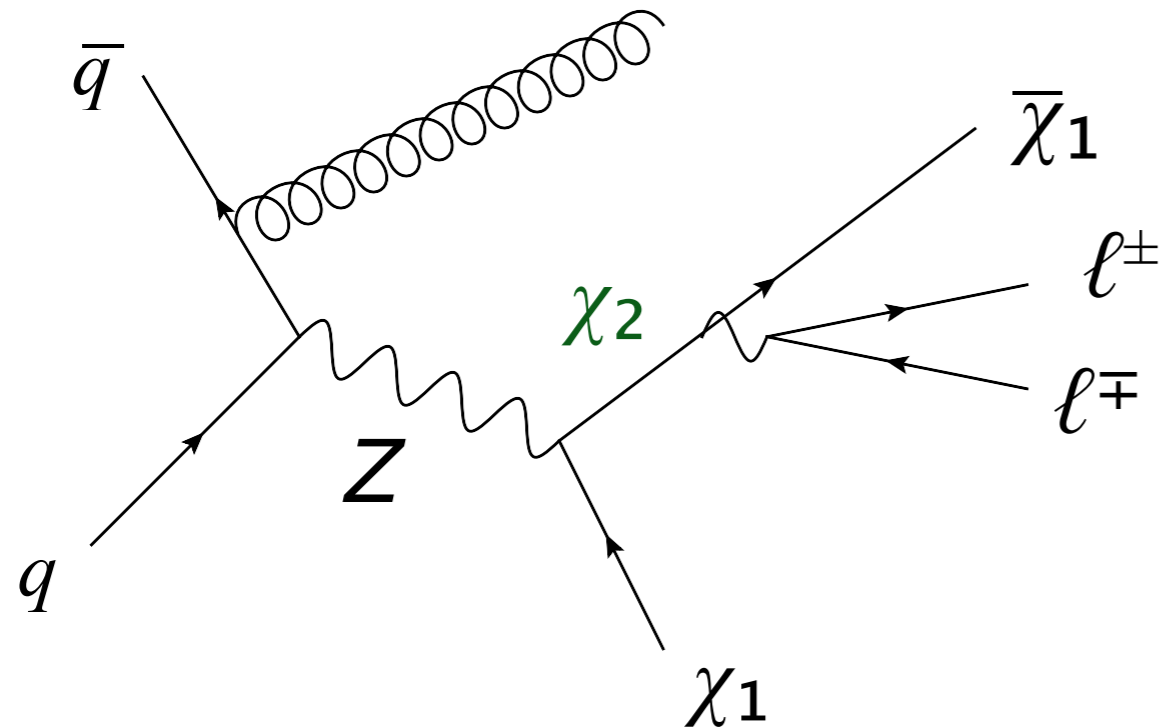
the problem with tri-lepton searches is triggering; solve this by requiring hard initial-state jet, as in mono-jet



once triggered upon, event can be analyzed using lower, off-line thresholds ( $\gtrsim 7$  GeV for  $e^\pm$ ,  $\gtrsim 3$  GeV for  $\mu^\pm$ )

# Example: jet + $\ell^+\ell^-$ + MET

[Han, Kribs, AM, Menon 1401.1235]



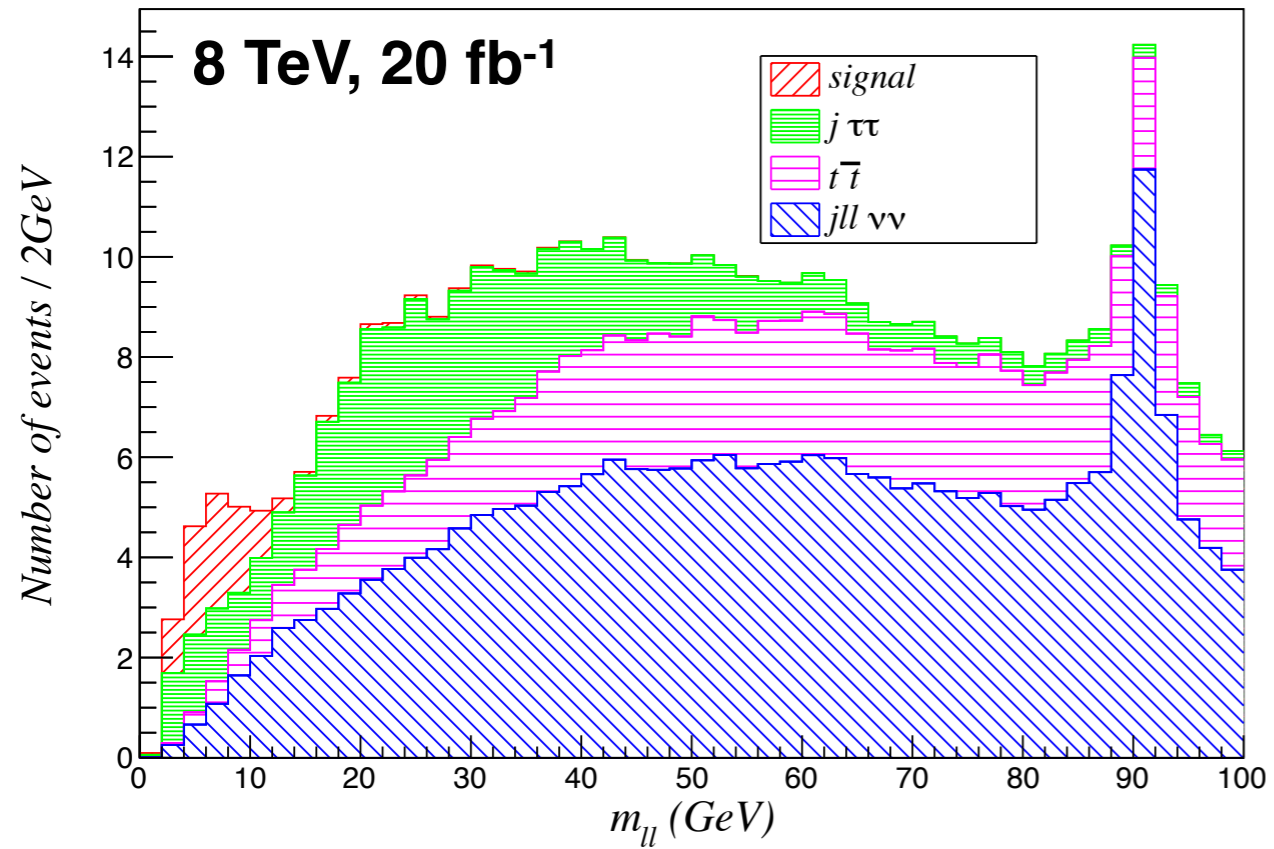
basic cuts: jet  $p_T > 100$  GeV,  
 $|\eta_j| < 2.5$ ,  
MET  $> 100$  GeV  
for triggering

then require 2 leptons,  
 $p_T > 7$  GeV,  $|\eta_\ell| < 2.5$

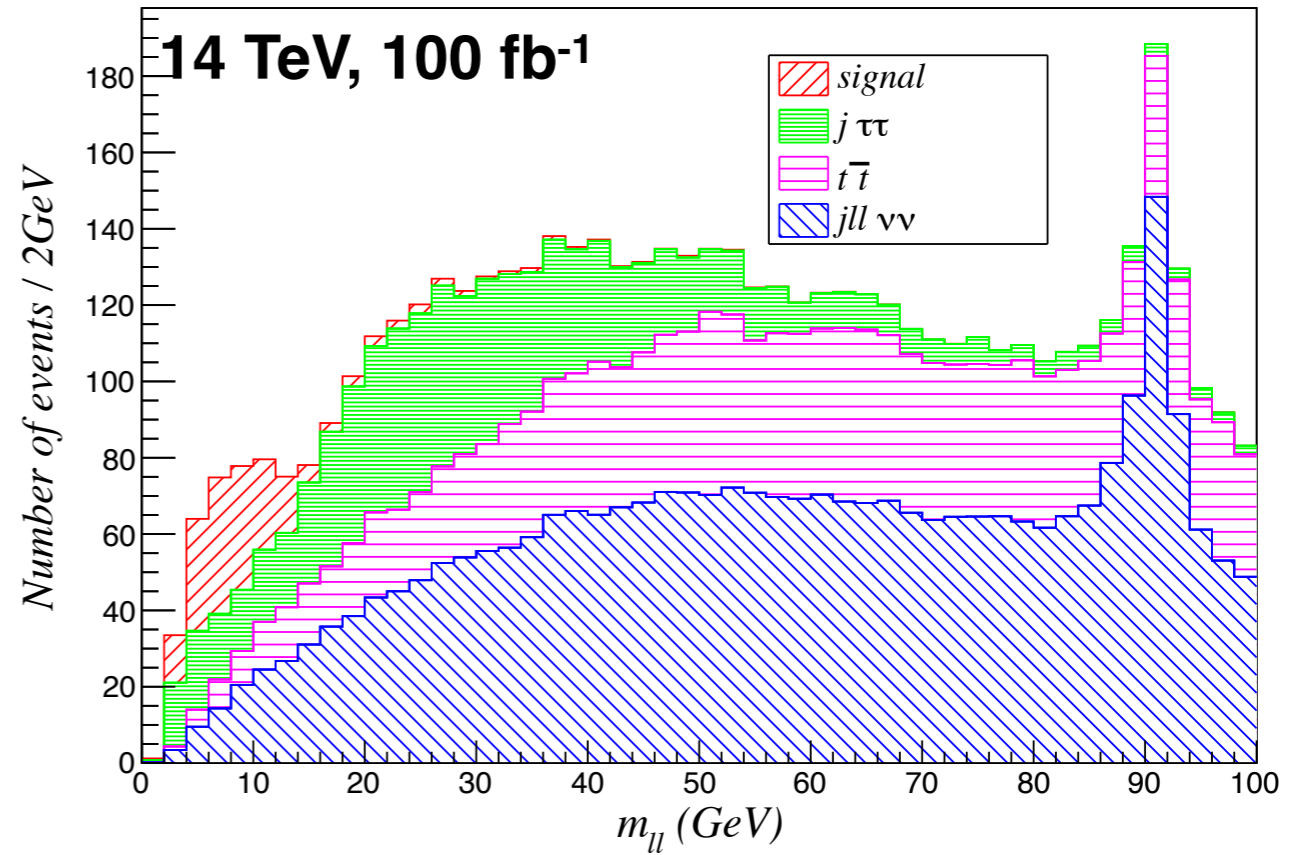
there are SM backgrounds ( $WW+j$ ,  $t\bar{t}$ ,  $\tau^+\tau^-+j$ ), but they  
can be controlled with cuts

# Example: jet + $\ell^+\ell^-$ + MET

M1=1000GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$



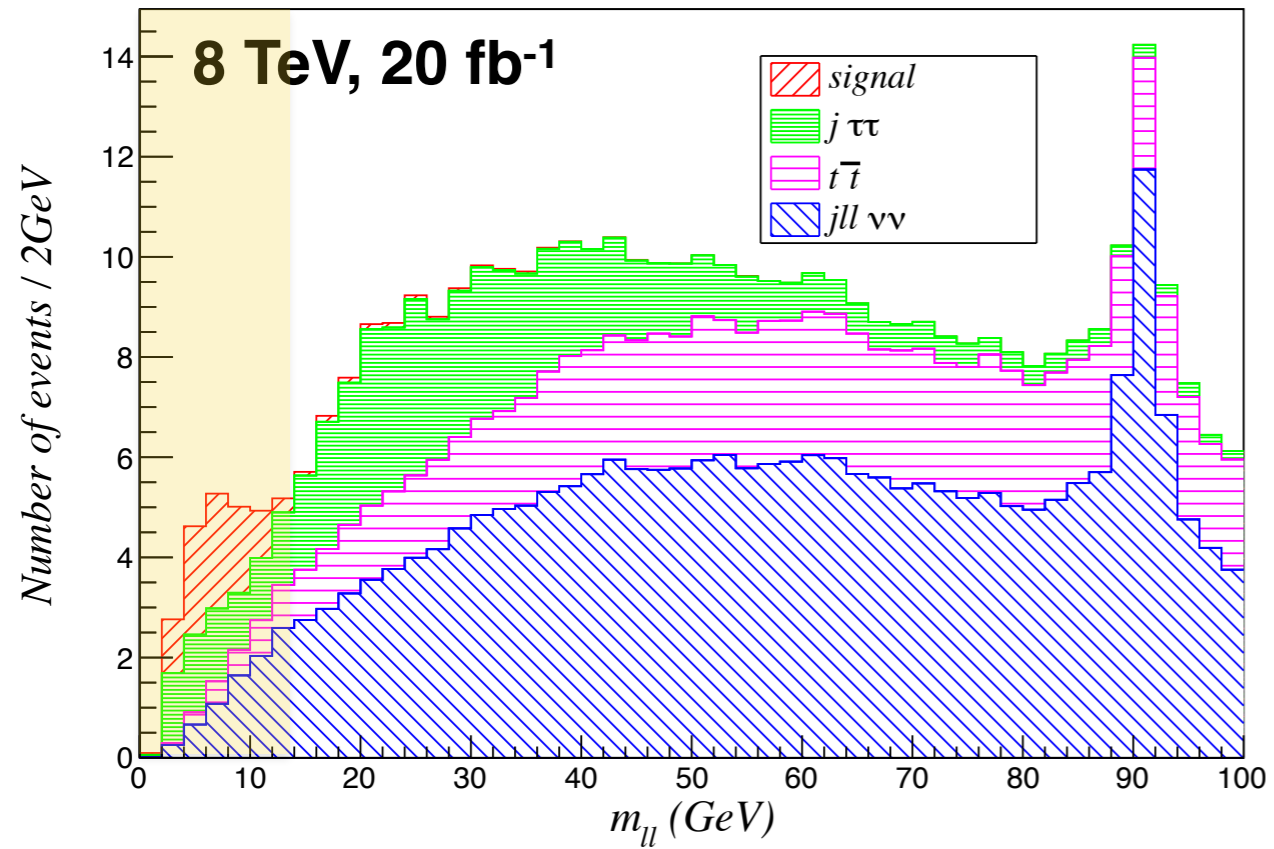
M1=500GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$



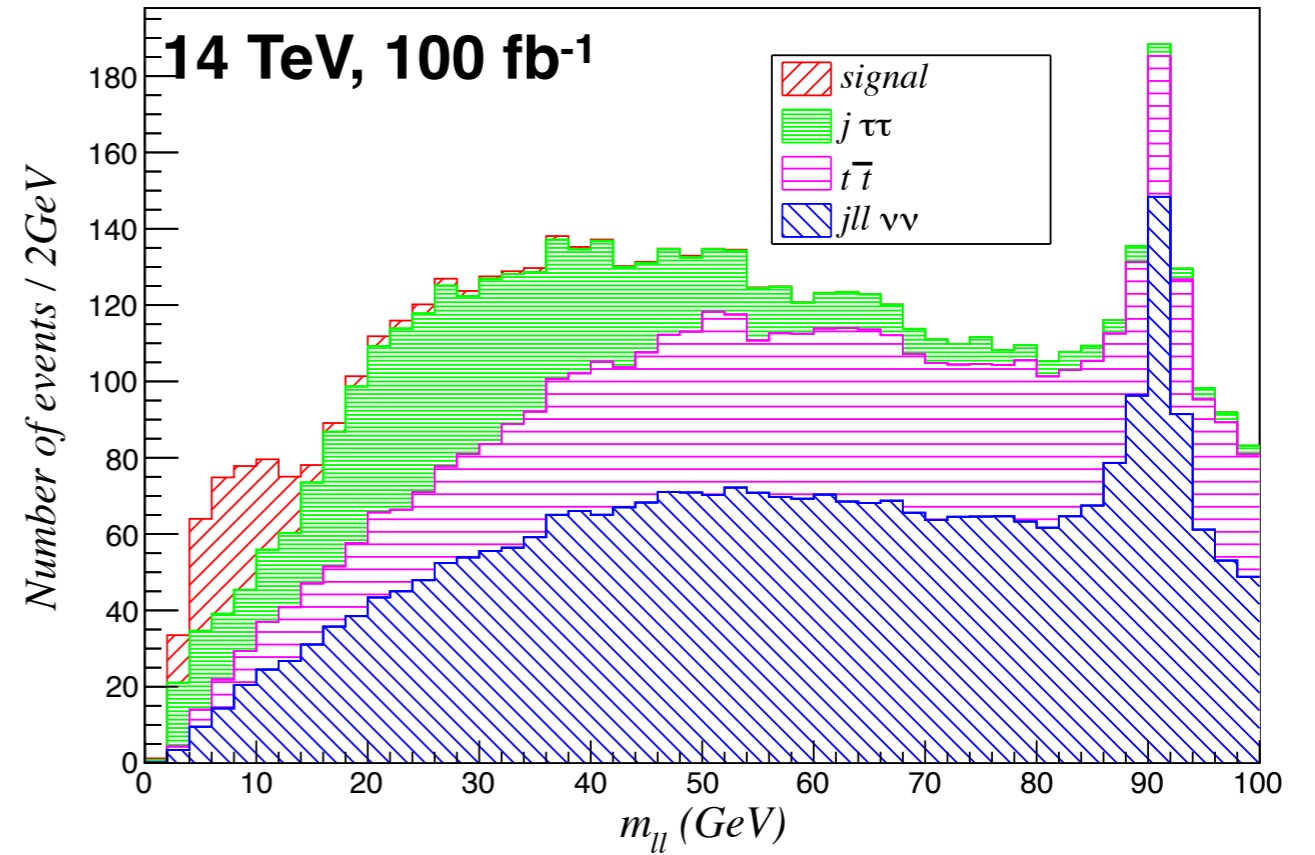


# Example: jet + $\ell^+\ell^-$ + MET

M1=1000GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$

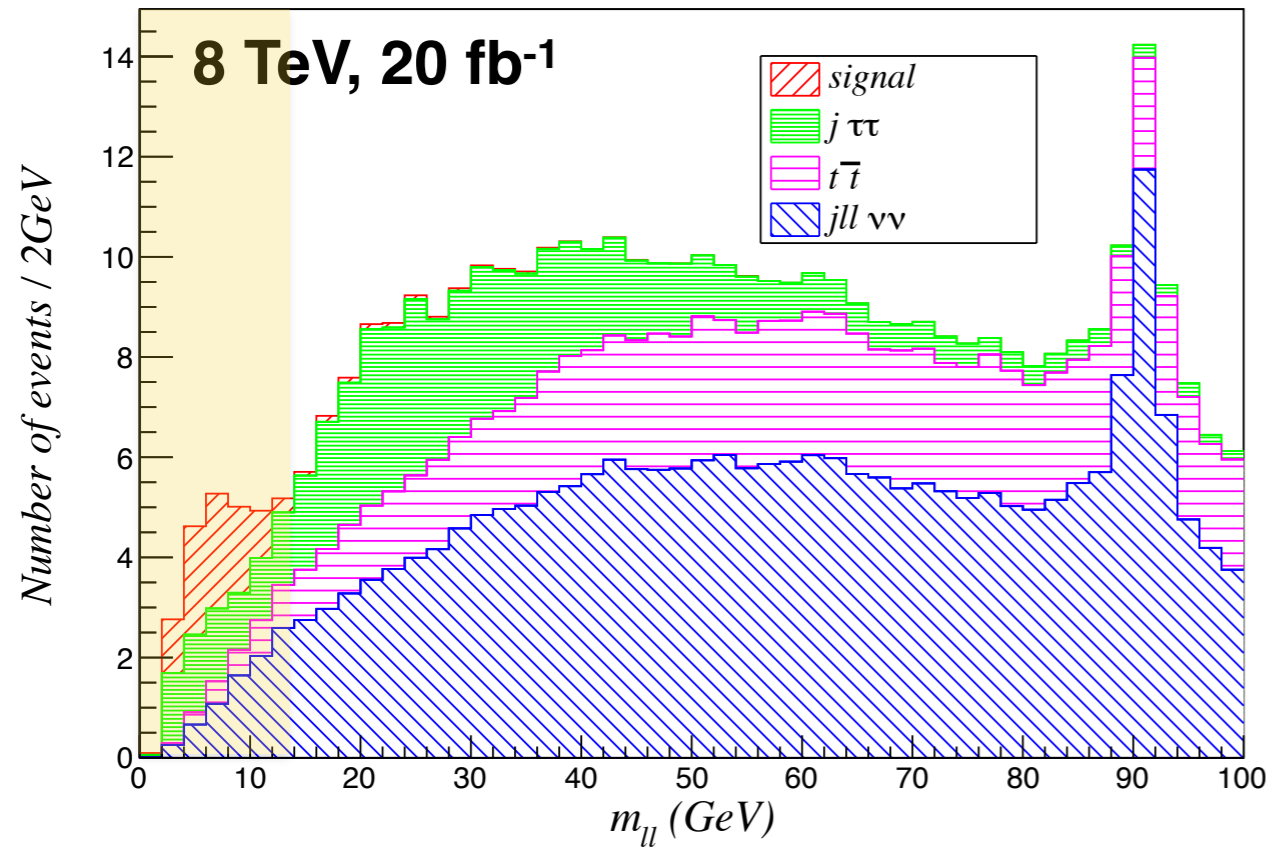


M1=500GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$

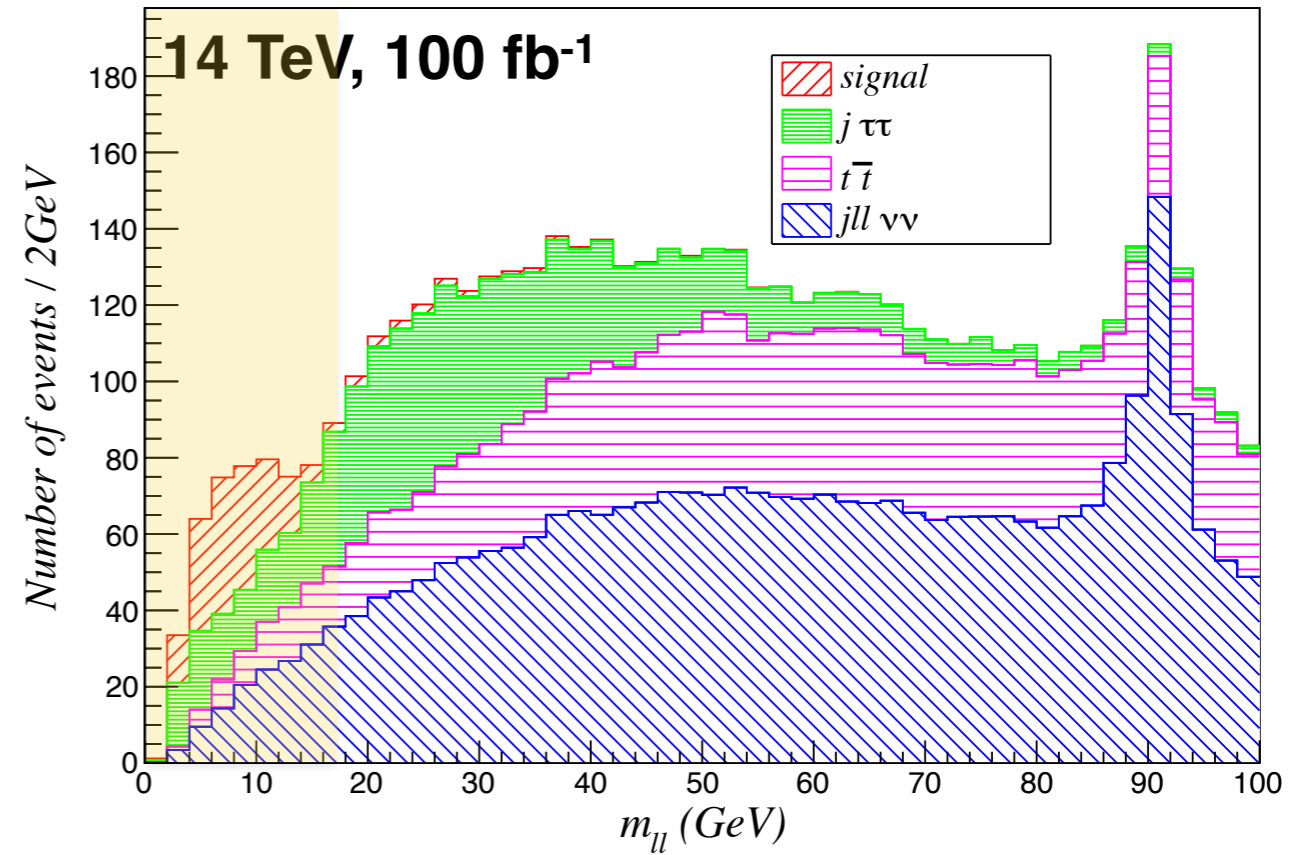


# Example: jet + $\ell^+\ell^-$ + MET

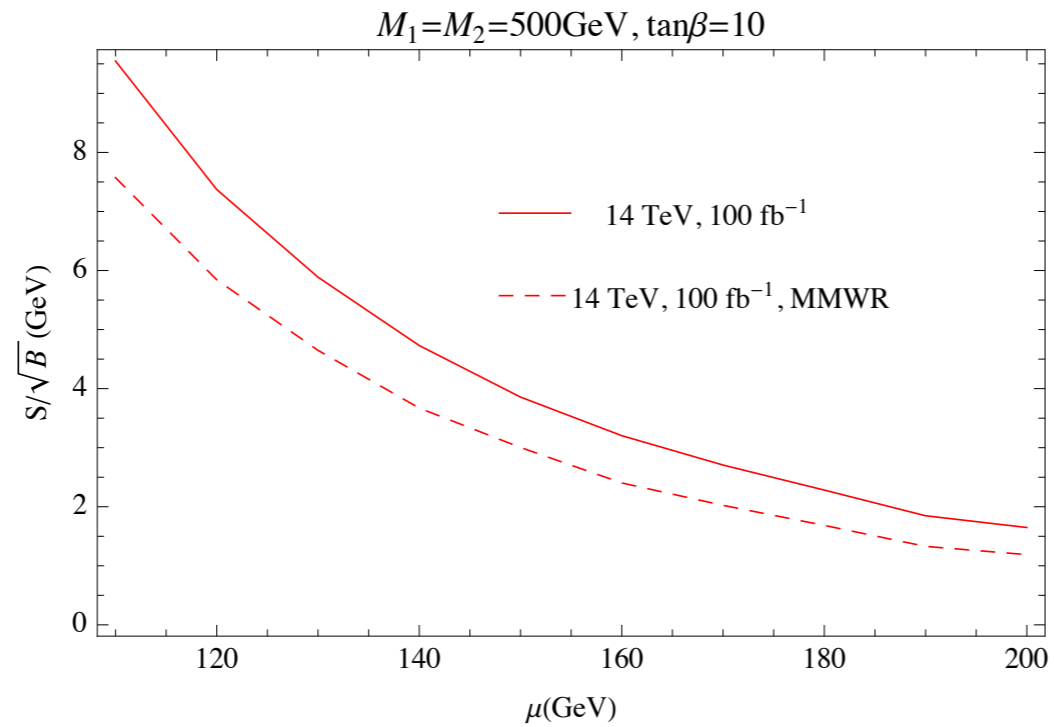
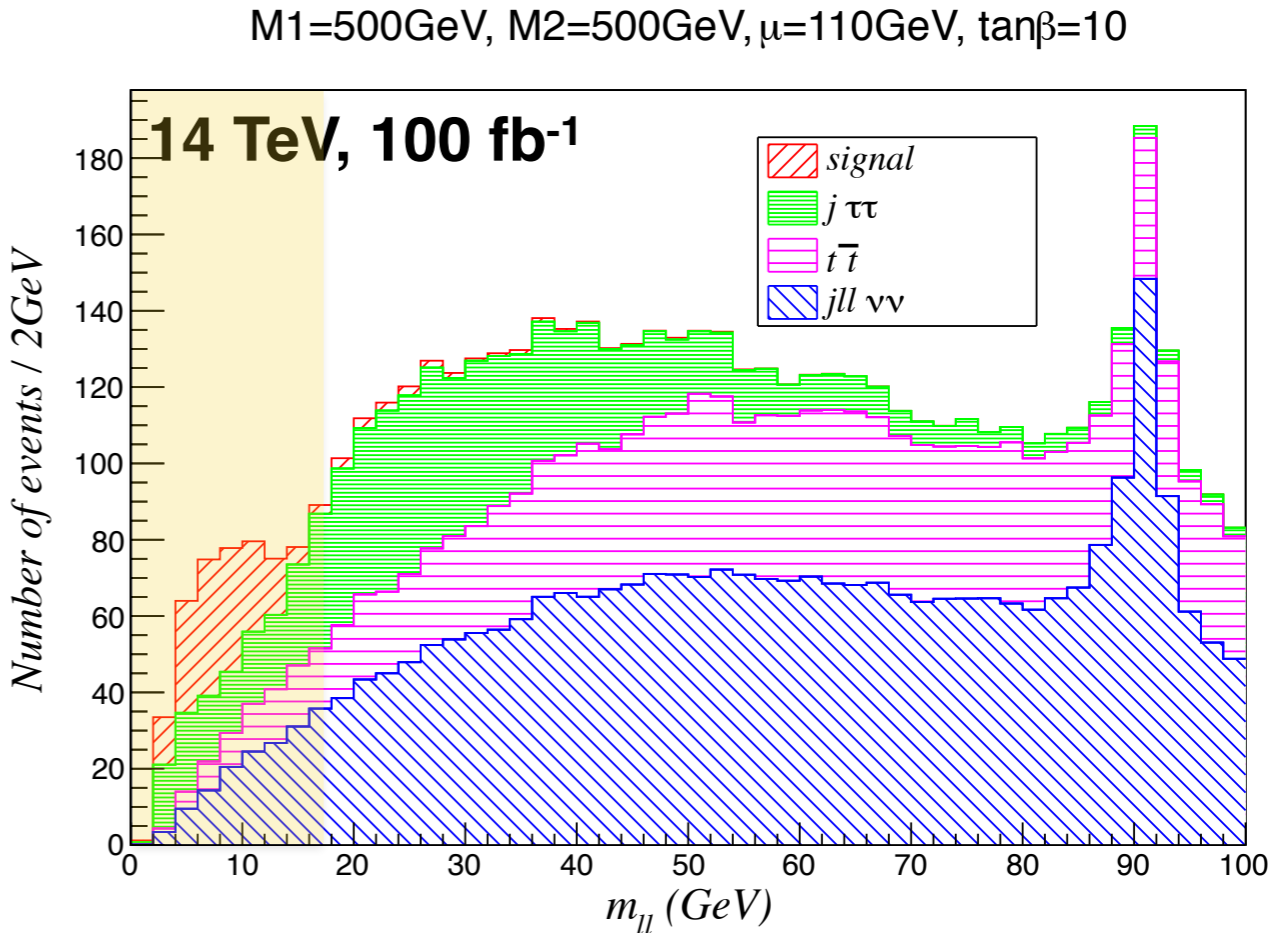
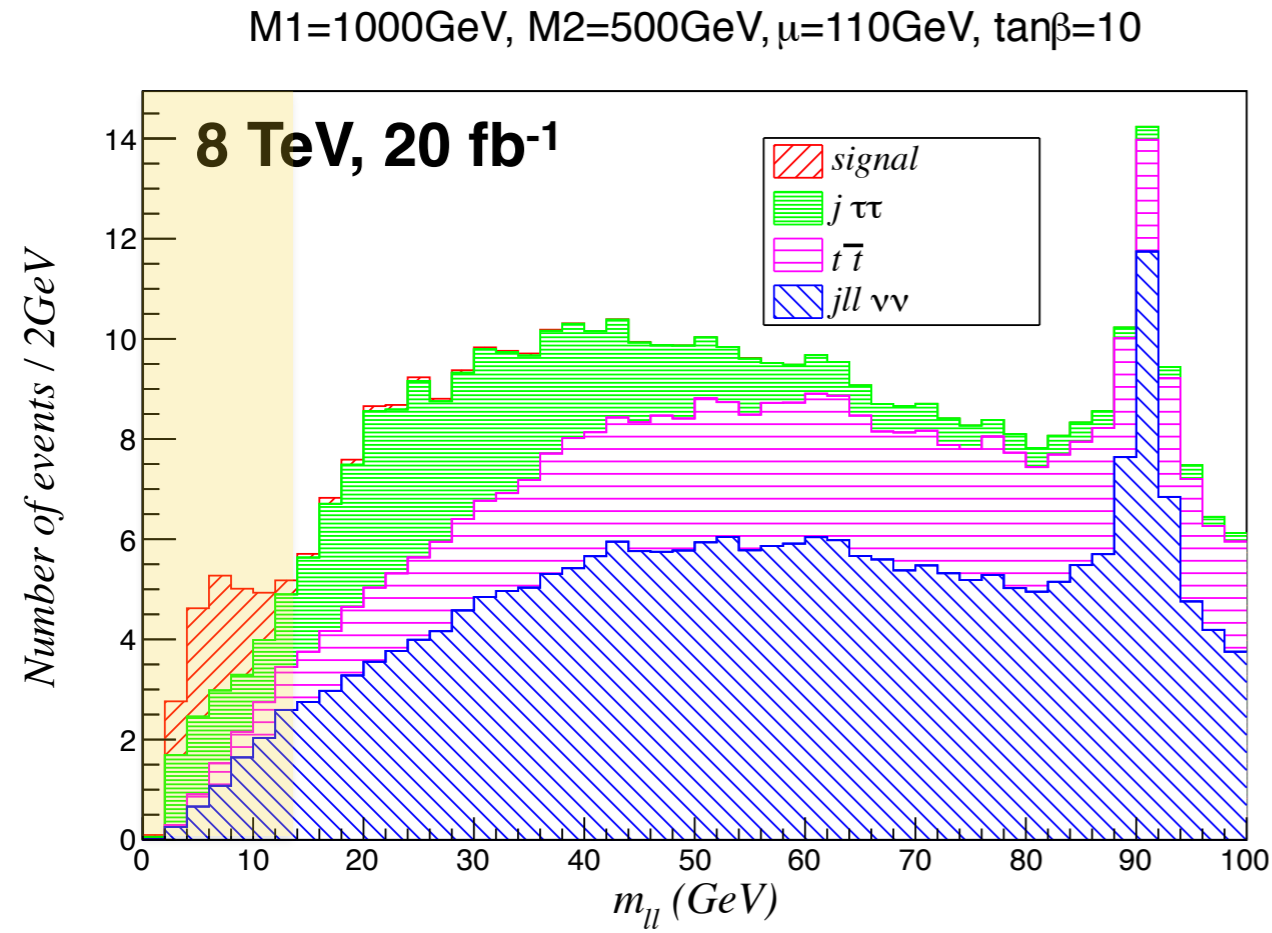
M1=1000GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$



M1=500GeV, M2=500GeV,  $\mu=110$ GeV,  $\tan\beta=10$



# Example: jet + $\ell^+\ell^-$ + MET



sensitive to spectra  
other searches  
cannot see!



# Improving mono-jet searches: jet + $\ell^+\ell^-$ + MET

why are these working?

# Improving mono-jet searches: jet + $\ell^+\ell^-$ + MET

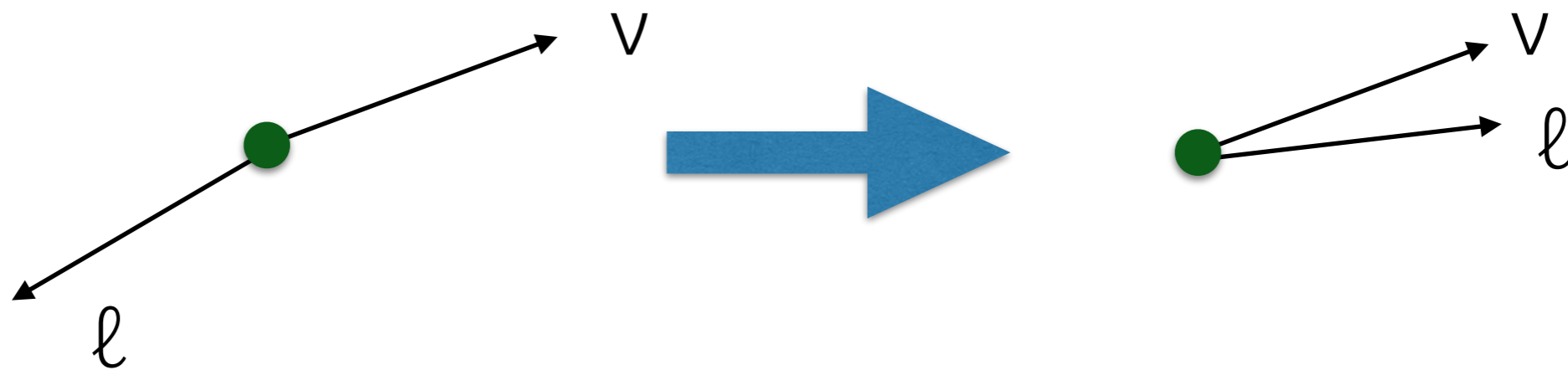
why are these working?

- signal is large MET + **soft** leptons (low  $m_{\ell\ell}$ )

# Improving mono-jet searches: jet + $\ell^+\ell^-$ + MET

why are these working?

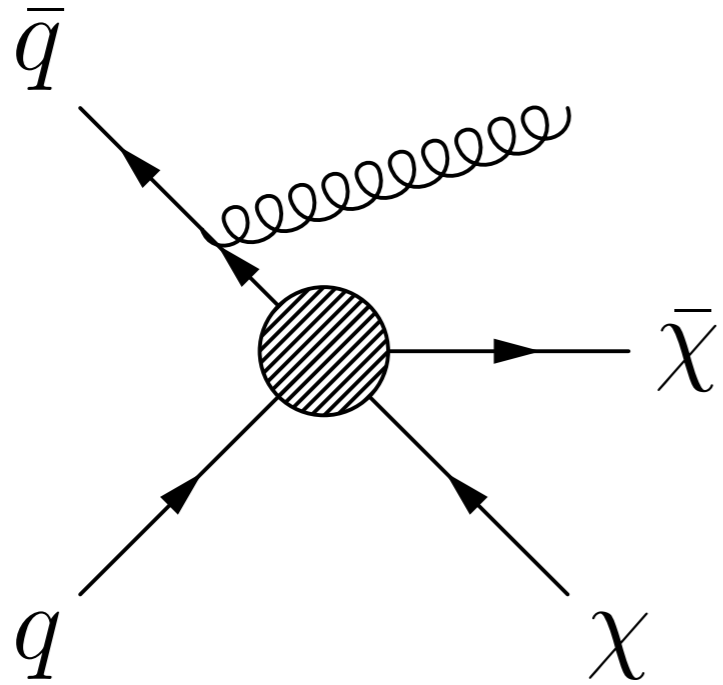
- signal is large MET + **soft** leptons (low  $m_{\ell\ell}$ )
- in background (di-boson), MET and lepton come from the **same** particle (W)



large MET  $\leftrightarrow$  **high  $p_T$**  leptons



Now that we know the trick, can be applied to lots of other SUSY DM signals/final states



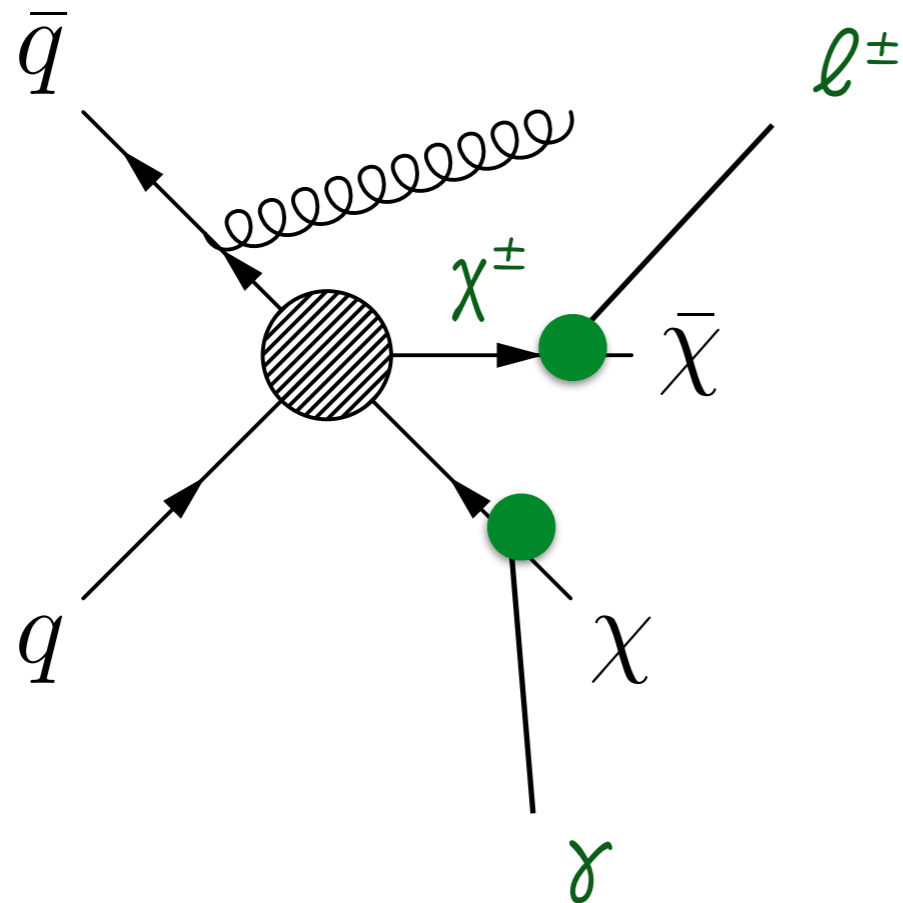
Ex.

$l + \gamma + \text{MET}$  final state  
occurs when the DM is a  
wino-bino mixture

important, as wino-bino DM often  
can't be found by  
direct/indirect detection

[See *Han 1409.7000, Bramante et al 1412.4789, 1510.03460*]

Now that we know the trick, can be applied to lots of other SUSY DM signals/final states



Ex.

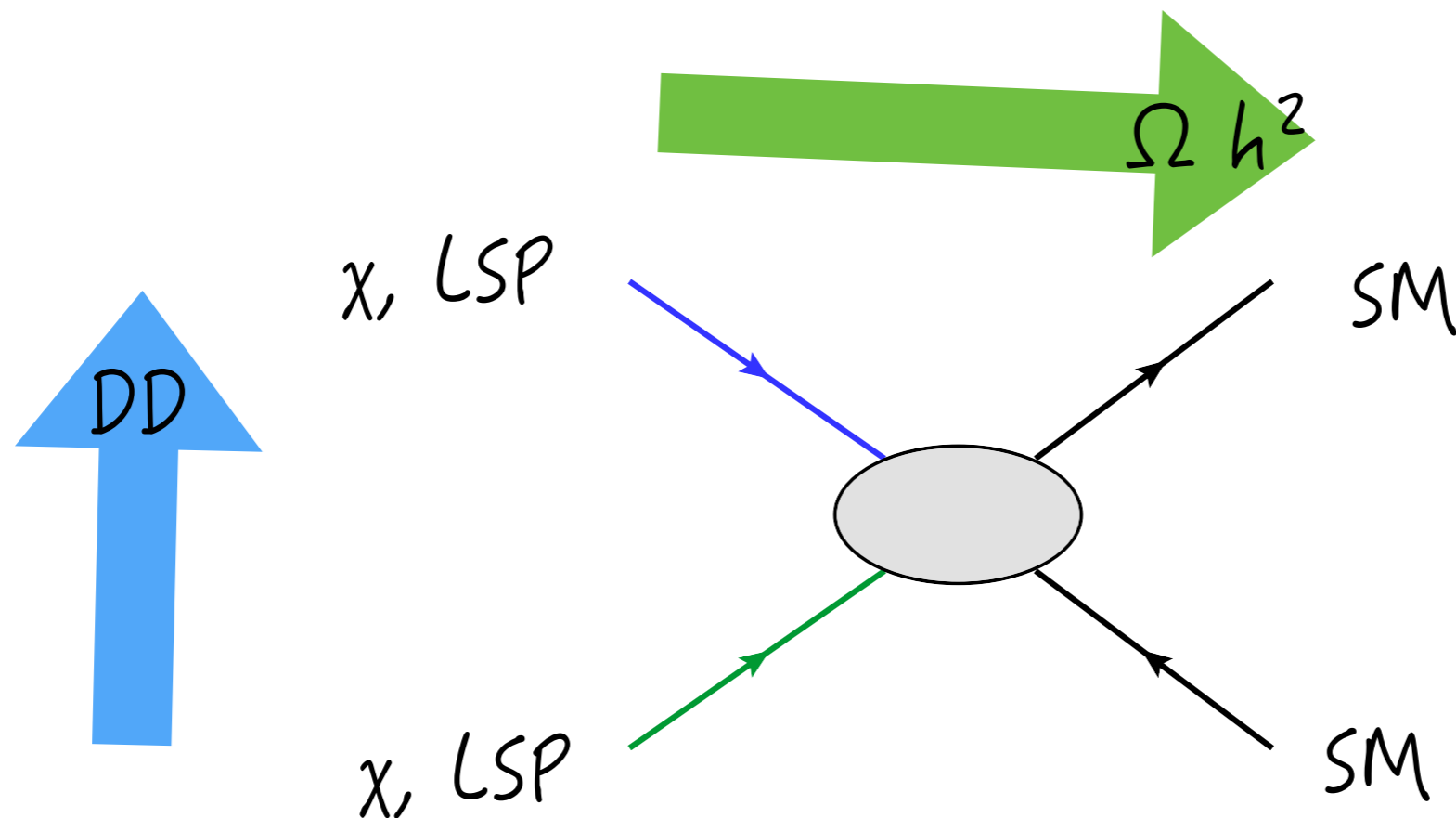
$\ell + \gamma + \text{MET}$  final state  
occurs when the DM is a  
wino-bino mixture

important, as wino-bino DM often  
can't be found by  
direct/indirect detection

[See Han 1409.7000, Bramante et al 1412.4789, 1510.03460]

Not restricted to SUSY dark matter:

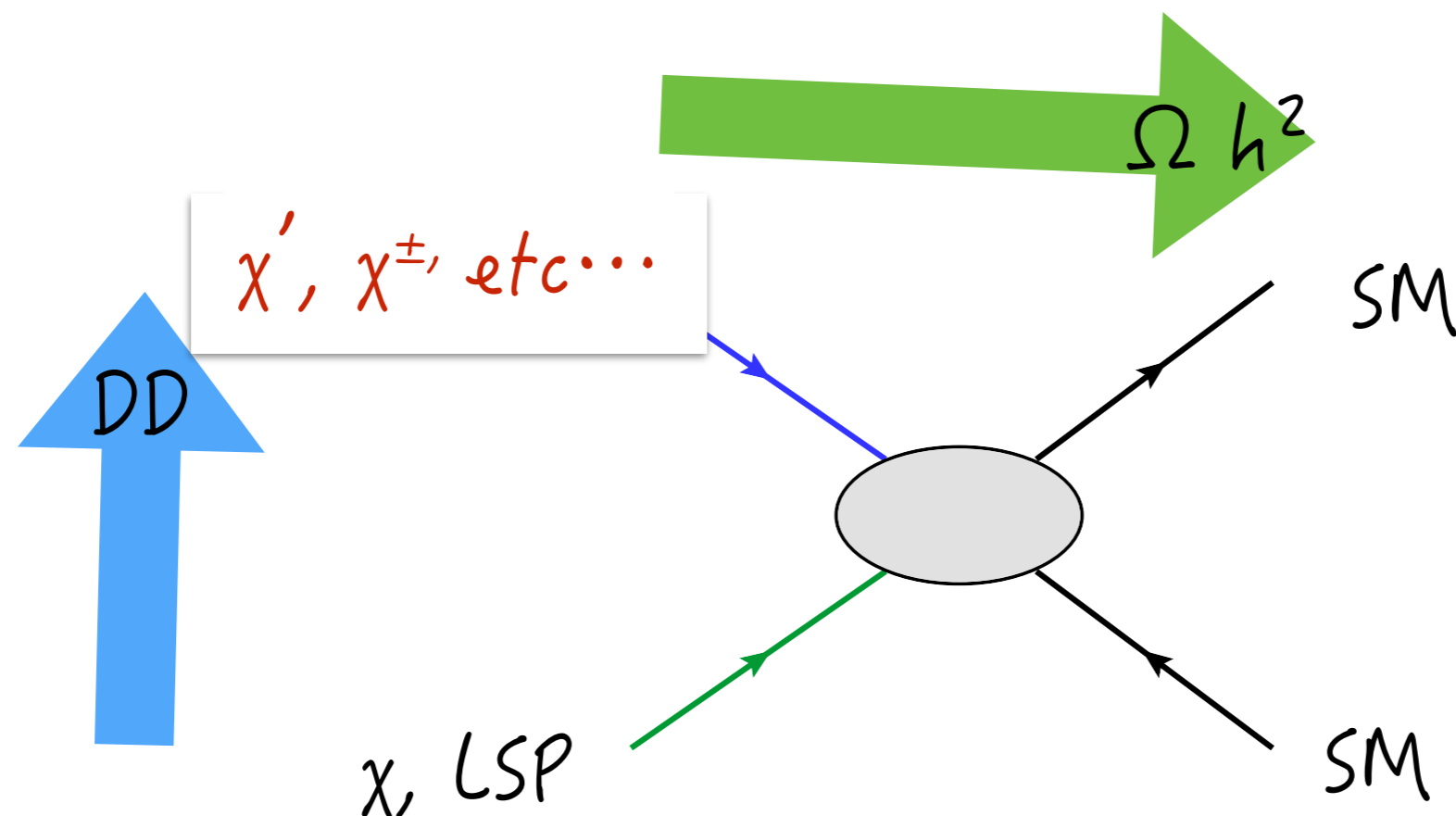
in fact, most current, viable models of EW DM will have a set of states with similar mass





Not restricted to SUSY dark matter:

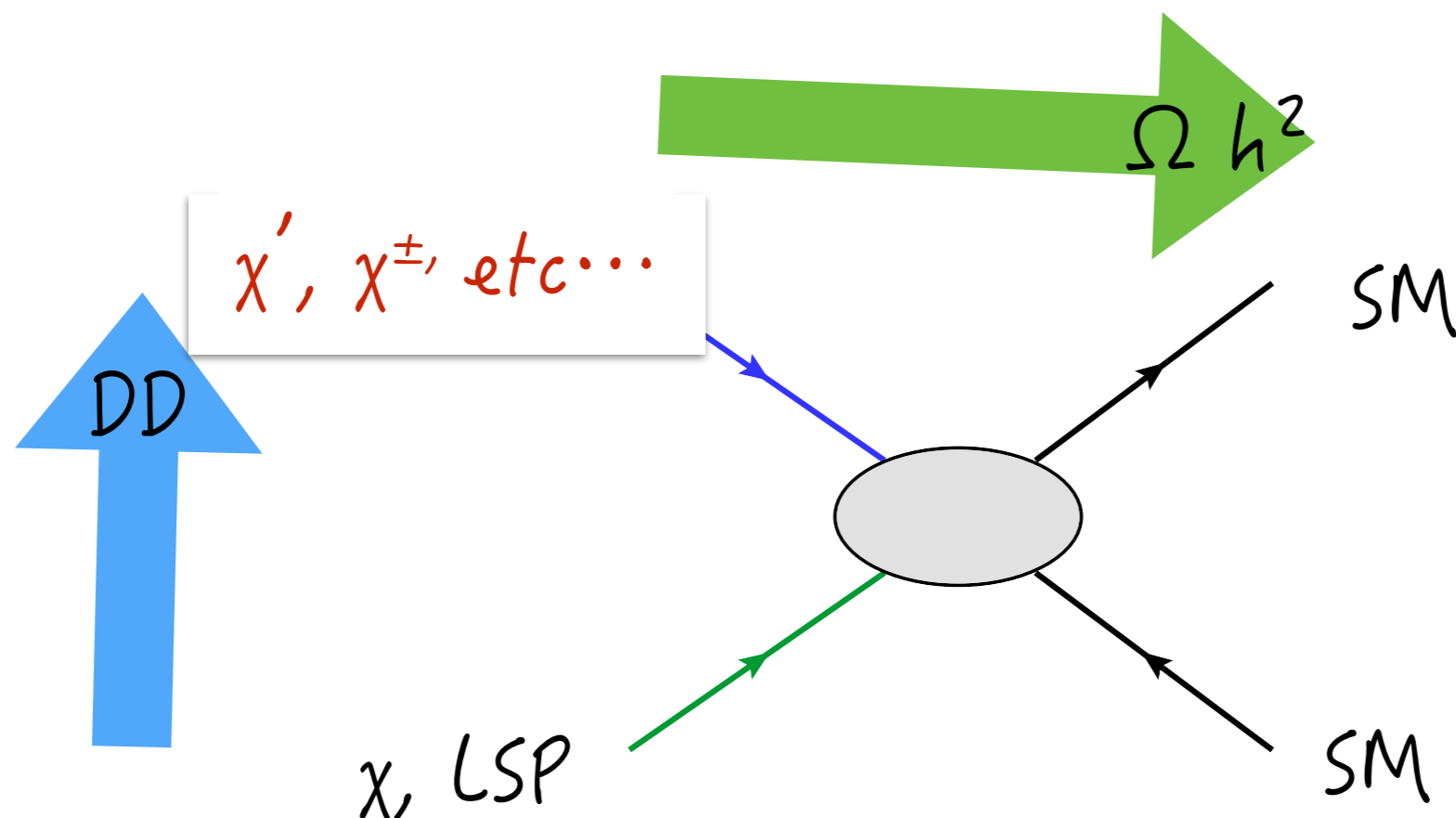
in fact, most current, viable models of EW DM will have a set of states with similar mass



**co-annihilation** fixes the tension, but co-annihilating states must be nearly same mass as DM

Not restricted to SUSY dark matter:

in fact, most current, viable models of EW DM will have a set of states with similar mass



$$\sigma_{eff} \cong \sum_{ij}^N \sigma_{ij} \frac{g_i g_j}{g_{eff}^2} \exp\left(- (m_i + m_j - 2m_{LSP})/T\right)$$

# Summarizing

If supersymmetry is the (or part of) the theory beyond the Standard Model, expect the electroweakinos to be light, fairly **degenerate**

.. this is exactly where they're hard to find

new idea to access these compressed spectra

- use initial radiation to trigger, look for accompanying softer objects ( $\ell^+\ell^-$ , etc.) offline
- should work for non-SUSY setups too!

look good at simulation level, work better for smaller splittings



**THANK YOU!**

have a great Masterclass, and hope to see you again soon!