

Unitarization for Vector Boson Scattering at the LHC

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The Standard Model as an Effective Field Theory

Effective Field Theory parameterizes new physics (NP) using the SM information available: working with a *bottom up* description, physics above the SM could be expanded in terms of a cut-off energy (Λ), where the NP is expected to be,

$$\mathcal{L}_{\text{EFT}} = \underbrace{\mathcal{L}_{\text{SM}}}_{\text{dim-4}} + \underbrace{\sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i}_{\text{dim-6}} + \underbrace{\sum_j \frac{f_{S,j}}{\Lambda^4} \mathcal{O}_{S,j} + \sum_k \frac{f_{T,k}}{\Lambda^4} \mathcal{O}_{T,k} + \sum_l \frac{f_{M,l}}{\Lambda^4} \mathcal{O}_{M,l}}_{\text{dim-8}}$$

General Anomalous Quartic Couplings: Dim-8 Operators

Deviations from the SM can be parameterized by:

■ Dim-6 operators

TGCs and QGCs

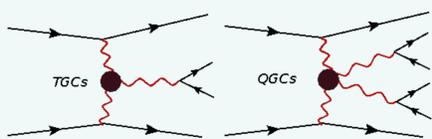
■ Dim-8 operators

QGCs

Dim-6 Anomalous Couplings are related to guarantee gauge invariance: deviations cannot be treated separately and TGCs are strongly constrained; thus, an independent parameterization of QGCs requires **dim-8 operators**.

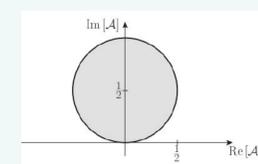
The dim-8 operators are characterized by their field content:

- Covariant derivatives acting on the Higgs doublet, $\mathcal{O}_{S,j}$, $\mathcal{O}_{S,0} = [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\mu \Phi)^\dagger D^\nu \Phi] \rightarrow$ **Longitudinal scattering.**
- Field strength tensors, $\mathcal{O}_{T,j}$, $\mathcal{O}_{T,0} = \text{Tr} [W_{\mu\nu} W^{\mu\nu}] \times \text{Tr} [W_{\alpha\beta} W^{\alpha\beta}] \rightarrow$ **Transverse scattering.**
- Or a mixture of both, $\mathcal{O}_{M,j}$, $\mathcal{O}_{M,0} = \text{Tr} [W_{\mu\nu} W^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi] \rightarrow$ **Mixed polarization.**



Unitarity Problem

Unitarity bounds are derived from on-shell 2→2 scattering:

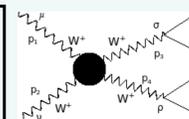


The Argand Circle

Using **partial wave decomposition** to rewrite the scattering amplitude and the **optical theorem**:

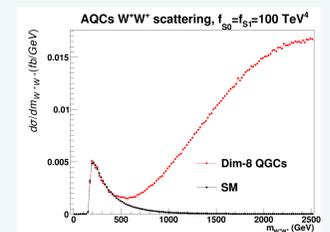
$$|A_j - \frac{i}{2}| \leq \frac{1}{2} \Rightarrow \text{Re}(A_j) < \frac{1}{2}$$

The Scattering amplitudes are not allowed to grow proportional to energy, or the cross section becomes unphysical.



$$\mathcal{A}_{S,0} \propto \frac{f_{S,0}}{\Lambda^4} \epsilon_1 \cdot \epsilon_2 \epsilon_3^* \cdot \epsilon_4^* \Rightarrow \mathcal{A}_{S,0} \propto \frac{s^2}{4}$$

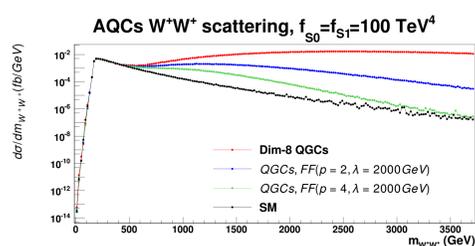
A unitarization scheme needed to obtain a valid prediction in this region.



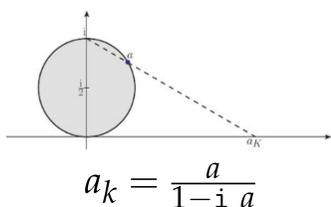
Restoring Unitarity

■ Form Factor Unitarization

In VBFNLO [1], a function $\mathcal{F}(s)$ is multiplied to the couplings, so it suppresses the high energy tail of amplitudes arising from the operators and thus ensures unitarity.



■ On-Shell K-Matrix Unitarization



Real scattering amplitudes are projected onto the Argand circle to restore unitarity, using the Cayley transform of the S-matrix (the K-matrix [2]).

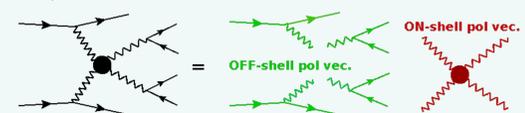
REFERENCES

- (1) J. Baglio et al., *VBFNLO: A parton level Monte Carlo for processes with electroweak bosons* [arxiv:1404.3940]
- (2) W. Killian, T. Ohl, J. Reuter and M. Sekulla, *High-Energy Vector Boson Scattering after the Higgs Discovery*. [arxiv:1408.6207]

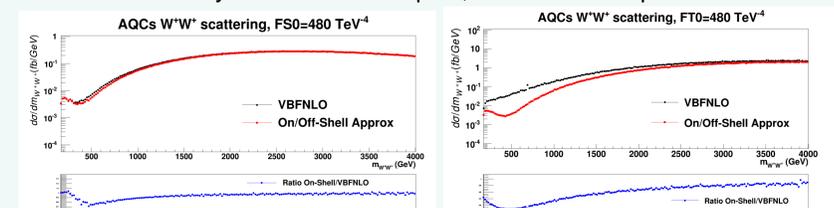
VBS at the LHC

■ On/Off-Shell Approximation

Using the properties of the polarization vectors, the scattering amplitude can be rewrite as,



to restore unitarity for the on-shell part, in an off-shell process.



The on/off-shell approach fails when transverse polarized particles are involved.

■ Analytic Solution

Using the partial wave decomposition for off-shell amplitudes,

$$A_{\text{QC}}(\lambda_1, \lambda_2, \lambda_3, \lambda_4) = 16\pi \sum_{j=0}^2 (2j+1) d_{\lambda_1-\lambda_2, \lambda_3-\lambda_4}^j(\theta) a^j(\lambda_1, \lambda_2, \lambda_3, \lambda_4)$$

we can determine unitarized amplitudes, calculating the analytic K-matrix.