

# Some Thoughts on the Relaxion

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with

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# Outline

1. Raising the cutoff
2. Relaxion inflation
3. Other phenomenology (e.g. DM, 750 GeV)

Raising the cutoff

# Reverse Relaxion

Would like to solve hierarchy problem all the way,  
raise relaxion cutoff to scales believed to be fundamental

instead of decreasing Higgs mass<sup>2</sup> we can increase it (from -cutoff<sup>2</sup>)  
already done in some models (e.g. Batell, Giudice, McCullough)  
an example (original non-QCD model):

$$\mathcal{L} \supset m_L L L^c + m_N N N^c + y h L N^c + \tilde{y} h^\dagger L^c N$$

take  $\tilde{y}$  small and  $m_N < m_L$

$$\begin{pmatrix} m_L & yv \\ 0 & m_N \end{pmatrix} \rightarrow \text{seesaw, small eigenvalue is } \begin{matrix} \sim \frac{m_L m_N}{yv} & \text{when } yv \gg m_L \\ \sim m_N & \text{when } yv \ll m_L \end{matrix}$$

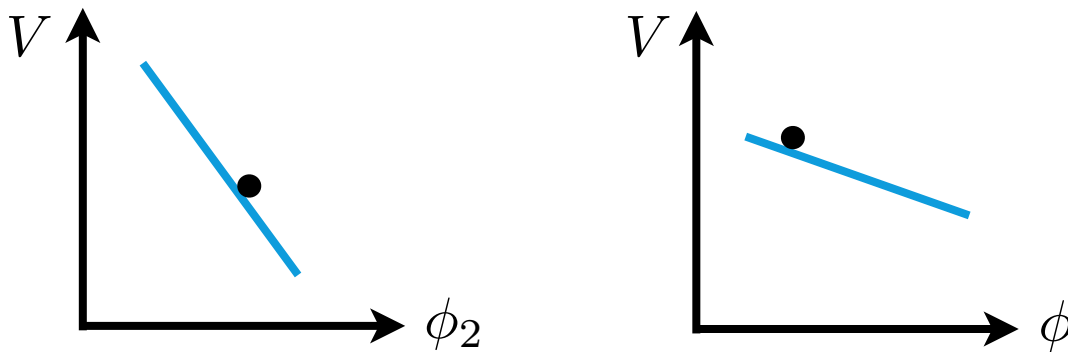
as Higgs vev drops, lightest mass increases  $\rightarrow$  raises barriers, stops vev

# Two Axions

If have two axions coupling to Higgs:

$$\begin{aligned}\mathcal{L} \supset M^2|h|^2 + g\phi|h|^2 + \Lambda^4 \cos(\phi/f) + gM^2\phi + \dots \\ + g_2\phi_2|h|^2 + \Lambda_2^4 \cos(\phi_2/f_2) + g_2M^2\phi_2 + \dots\end{aligned}$$

take scales to be higher for  $\phi_2$  than  $\phi$ , ( $g_2 > g$ ) then both axions roll from start



but 2 rolls faster, stops first  $\rightarrow$  can think of it in stages (though really happens simultaneously)

$\phi_2$  will “go first.” Scans Higgs vev until nearer to zero.

$\phi$  goes next, scans in finer steps, stops when vev even smaller.

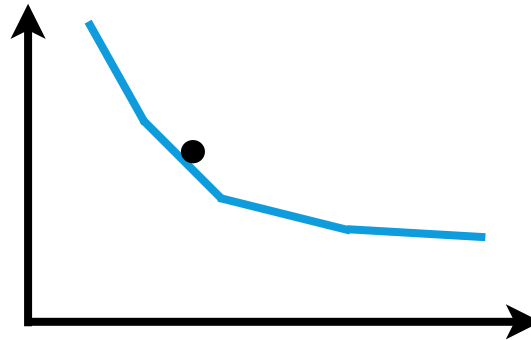
Together do more tuning of Higgs vev than either one alone.

# Multiple Axions

If have multiple axions coupling to Higgs:

$$M^2|h|^2 + \sum_i (g_i \phi_i |h|^2 + \Lambda_i^4 \cos(\phi_i/f_i) + g_i M^2 \phi_i)$$

higher slopes go “first”  
(really simultaneous)  
→ effectively like a potential:



weak scale is special location,  
would be tuned for single field

increases cutoff since steepest at start where “eternal inflation” constraint strictest  $H^3 < V'$   
and shallowest slope at end where barriers must stop it  $V' f < \Lambda^4$

works most easily if these are inflation (CC tuned as usual!)

# Multiple Axions

$$M^2|h|^2 + \sum_i (g_i \phi_i |h|^2 + \Lambda_i^4 \cos(\phi_i/f_i) + g_i M^2 \phi_i)$$

Higgs vev

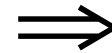
$$\begin{array}{l} \text{---} M_n = M \\ \vdots \\ \text{---} M_2 \\ \text{---} M_1 \\ \text{---} M_w \end{array} \left. \vphantom{\begin{array}{l} \text{---} M_n = M \\ \vdots \\ \text{---} M_2 \\ \text{---} M_1 \\ \text{---} M_w \end{array}} \right\} \phi_1$$

each field will roll over  $\Delta\phi_i \sim \frac{M_i^2}{g_i}$  before stopping

trade Lagrangian parameters for top and bottom of vev ranges:

no eternal inflation:  $H_i^3 < g_i M^2$

barriers form:  $g_i M^2 f_i < \Lambda_i^4 (\sim M_{i-1}^4)$



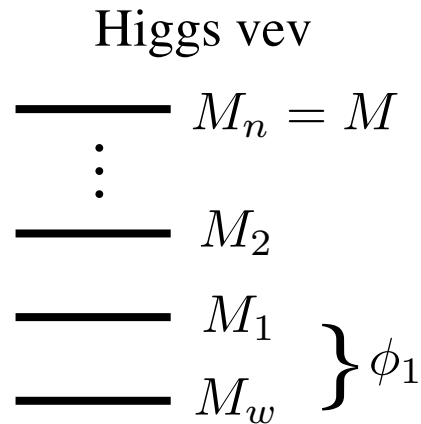
bound on “cutoff” of stage  $M_i$

$$M_i^3 M^4 < M_{i-1}^4 M_{\text{pl}}^3$$

$\nearrow$  was  $M_i$                        $\nwarrow$  was  $M_w$

so  $\{M_i\}$  a series of geometric means

# Increased Cutoff



rearranging  $\rightarrow$  size of each step:  $\left(\frac{M_{i-1}}{M_i}\right)^4 > \frac{M^4}{M_i M_{\text{pl}}^3}$

can keep lowering vev until this = 1

so setting  $M_i = M_w$  find cutoff:  $M^4 \sim M_w M_{\text{pl}}^3$

check classical rolling and quantum fluctuations during any upper stage don't move lower axions too far. Turns out automatically satisfied (weaker than above constraint).

3 axions give cutoff  $\sim 10^{14}$  GeV

asymptotically cutoff  $\sim 10^{16}$  GeV



# Inflation

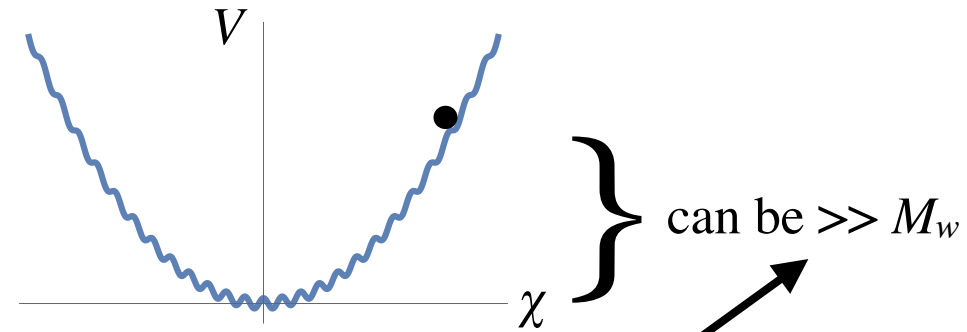
# Two-field Model

*new relaxion model:*

$$\mathcal{L} \supset (-M^2 + g_S S) h^2 + g_S M^2 S + \Lambda_S^4 \cos\left(\frac{S}{f_S}\right) + k S^2 \chi^2 + \Lambda_\chi^4 \cos\left(\frac{\chi}{f_\chi}\right)$$

$\nwarrow$   
 $\propto \langle h \rangle$

$S$  scans Higgs vev  
 $\chi$  released at  $M_w$   $\rightarrow$  drops  $S$  slope, stops  $S$



main new constraint: slope drop fast enough

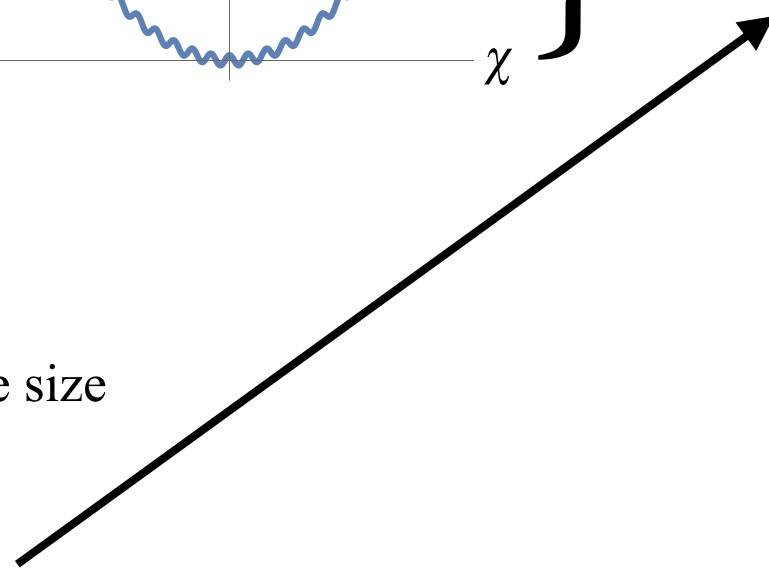
changes relaxion rules

$h$  vev steps can be weak scale size

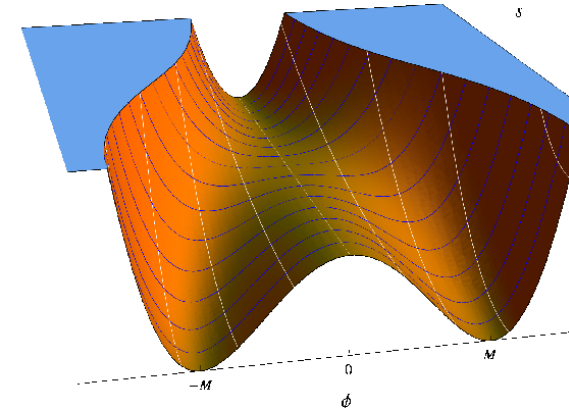
stops in unique vacuum

$\chi$  carries large energy density

$\chi$  can fast roll (satisfies new constraint)



# Inflation



can reheat with energy in  $\chi \rightarrow$  like hybrid inflation ( $\chi$  waterfall)

can solve hierarchy problem and do inflation

to reheat (i.e. fast roll):  $\chi < M_{\text{pl}} \rightarrow$  cutoff  $\sim 100$  TeV

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if barriers unrelated to  $M_w$  then can increase cutoff

if change  $h$  to new “Higgs”  $H_S$

$$\mathcal{L} \supset (-M^2 + g_S S) |H_S|^2 + g_S M^2 S + \Lambda_S^4 \cos\left(\frac{S}{f_S}\right) + k S^2 \chi^2 + \Lambda_\chi^4 \cos\left(\frac{\chi}{f_\chi}\right)$$

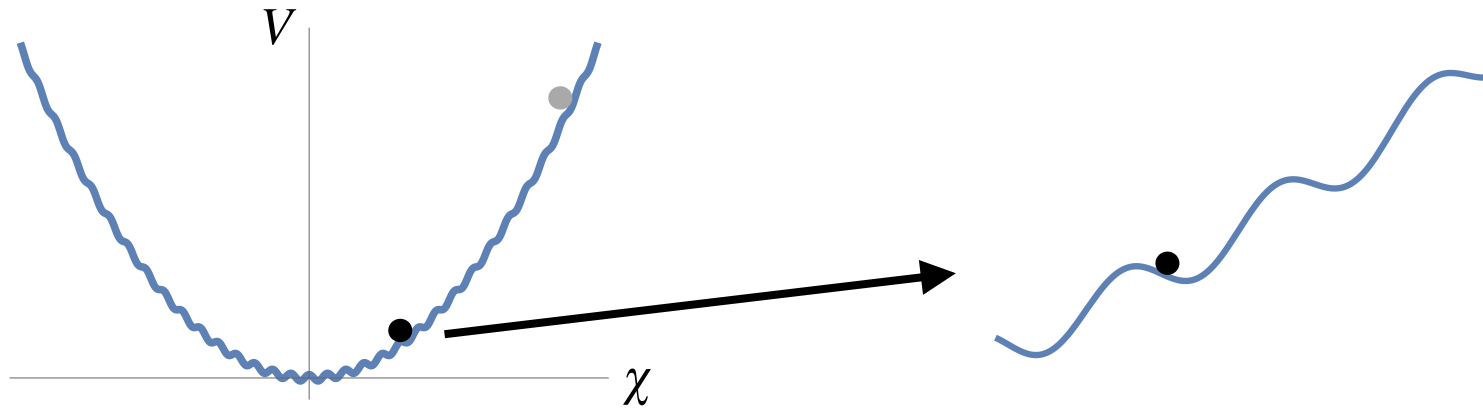
changes relaxation parameters: e.g.  $\Lambda_\chi \rightarrow M$  instead of  $M_w$

gives natural inflation model with  $10^{42}$  e-folds and cutoff at  $\sim 10^9$  GeV

$\rightarrow$  can be inflation model for original relaxation

# Reheating

$\chi$  energy redshifts with expansion of universe by  $O(1)$ , turnaround within barriers



greatly increases effective mass of  $\chi$   $m_\chi \sim \frac{M^2}{M_{\text{pl}}} \rightarrow \frac{M^2}{f_\chi}$

rapid decays to SM through  $\frac{\chi}{f_\chi} F \tilde{F}$   $\rightarrow$  can reheat to high scale

can it be end stage of inflation for multiple axions?

Other phenomenology

# Relaxion Baryogenesis?

Can the axion cause baryogenesis?

during electroweak phase transition, axion couples to baryon number current

Servant (2014), ...

sphalerons violate B, axion violates CP

breaking of shift symmetry  $\rightarrow$  relaxion is rolling

$\rightarrow$  non-zero chemical potential for B

$\rightarrow$  Baryon number generated in equilibrium (spontaneous baryogenesis)

goal: single field does hierarchy, strong CP, baryogenesis, and DM

requires some model-building for large enough B

# Summary

Cutoff can go up to  $\sim 10^{16}$  GeV with multiple axions coupled to Higgs

Demonstrated inflation models for relaxion

Can we put these two together?

Interesting phenomenology

- 750 GeV
- dark matter predictions?
- baryogenesis?

A new friction mechanism besides Hubble?

Generic signatures? Other (many) light scalars?





Backup

# Predictions

Dynamics (SUSY, extra dimensions...) → weak-scale particles (e.g. WIMP)

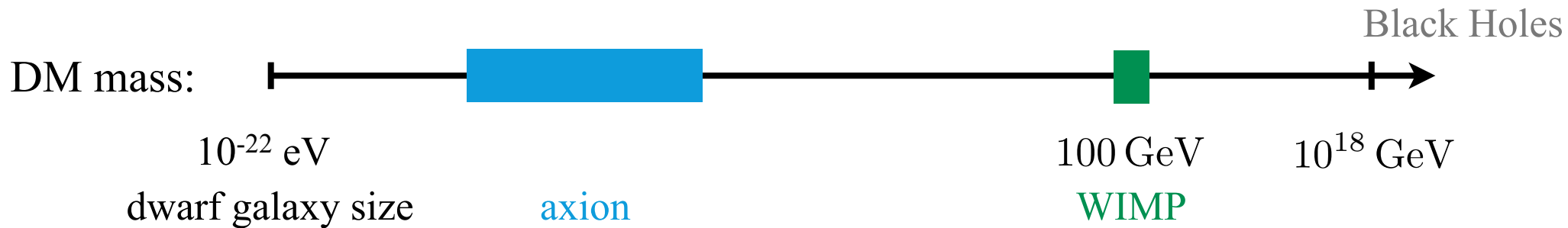
Dynamical Relaxation → light particles (e.g. axion)

- changed predictions for axion DM
- axion DM fluctuates Higgs VEV → oscillates all scales (electron mass...) potentially observable (at low cutoff), would be true proof of mechanism

# Precision Measurement for Dark Matter

# Dark Matter Candidates

What do we know about dark matter?



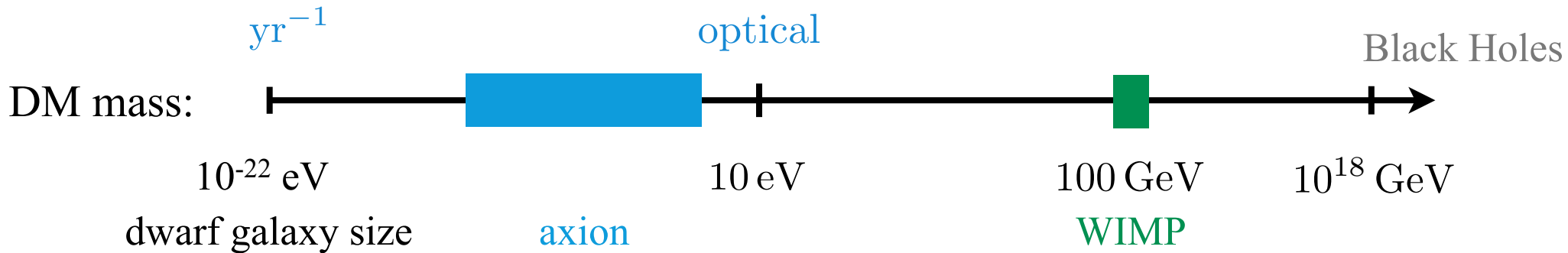
WIMP is well-motivated, significant direct detection effort focused on WIMPs,

Axion is other best-motivated candidate, only a small fraction of parameter space covered

Huge DM parameter space currently unexplored!

# Direct Detection

How can we detect DM?



$$\rho_{\text{DM}} \approx 0.3 \frac{\text{GeV}}{\text{cm}^3} \approx (0.04 \text{ eV})^4 \rightarrow \text{high phase space density if } m \lesssim 10 \text{ eV}$$

field-like (e.g. axion)  
new detectors required

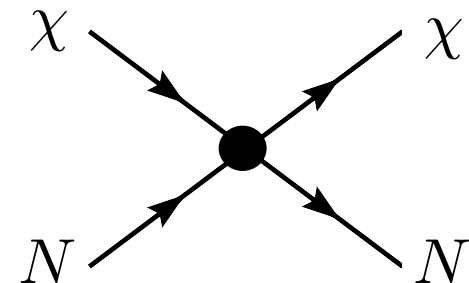
Described as classical field  $a(t,x)$

Detect coherent effects of entire field,  
not single particle scatterings

Frequency range accessible!

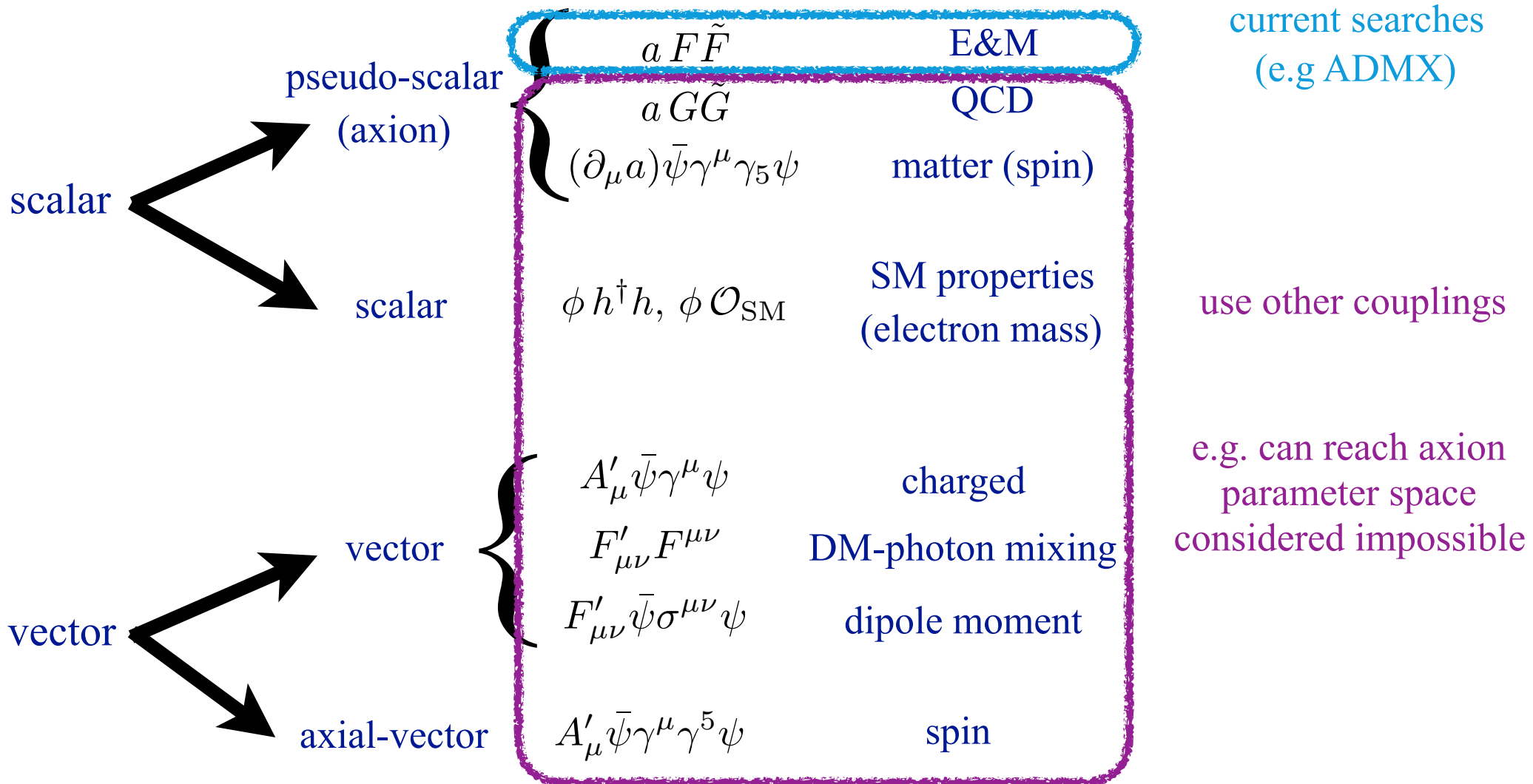
particle-like (e.g. WIMP)  
particle detectors best

Search for single, hard particle scattering



# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities (symmetry, effective field theory):



Can cover all these possibilities!

# Cosmic Axion Spin Precession Experiment (CASPER)

Detect axion with NMR and high-precision magnetometry

New field of axion direct detection, similar to early stages of WIMP direct detection

No other way to search for light axions

Would be the discovery of dark matter and glimpse into physics at high energies

Construction beginning at Mainz and BU

Boston University

Alexander Sushkov

Cal State

Derek J. Kimball

JGU Mainz

Dmitry Budker

Peter Blümler

Arne Wickenbrock

Helmholtz Institute Mainz

John Blanchard

Nathan Leifer

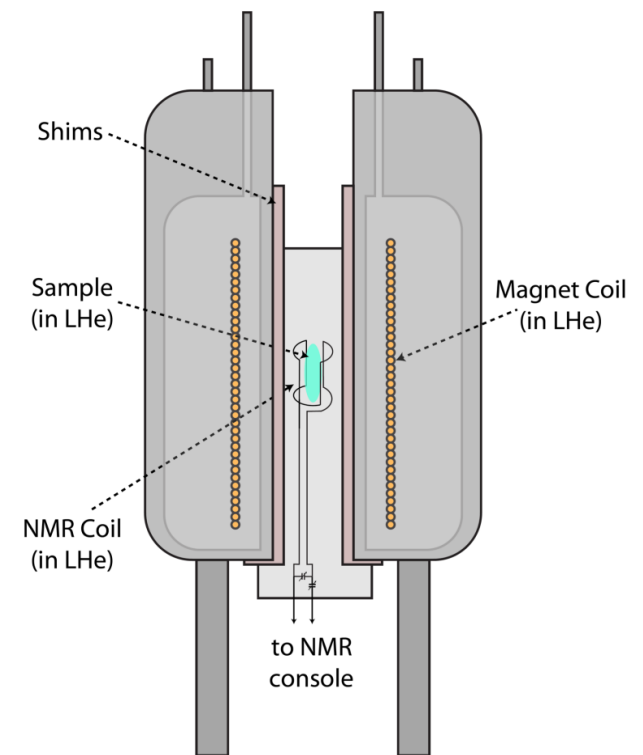
Stanford

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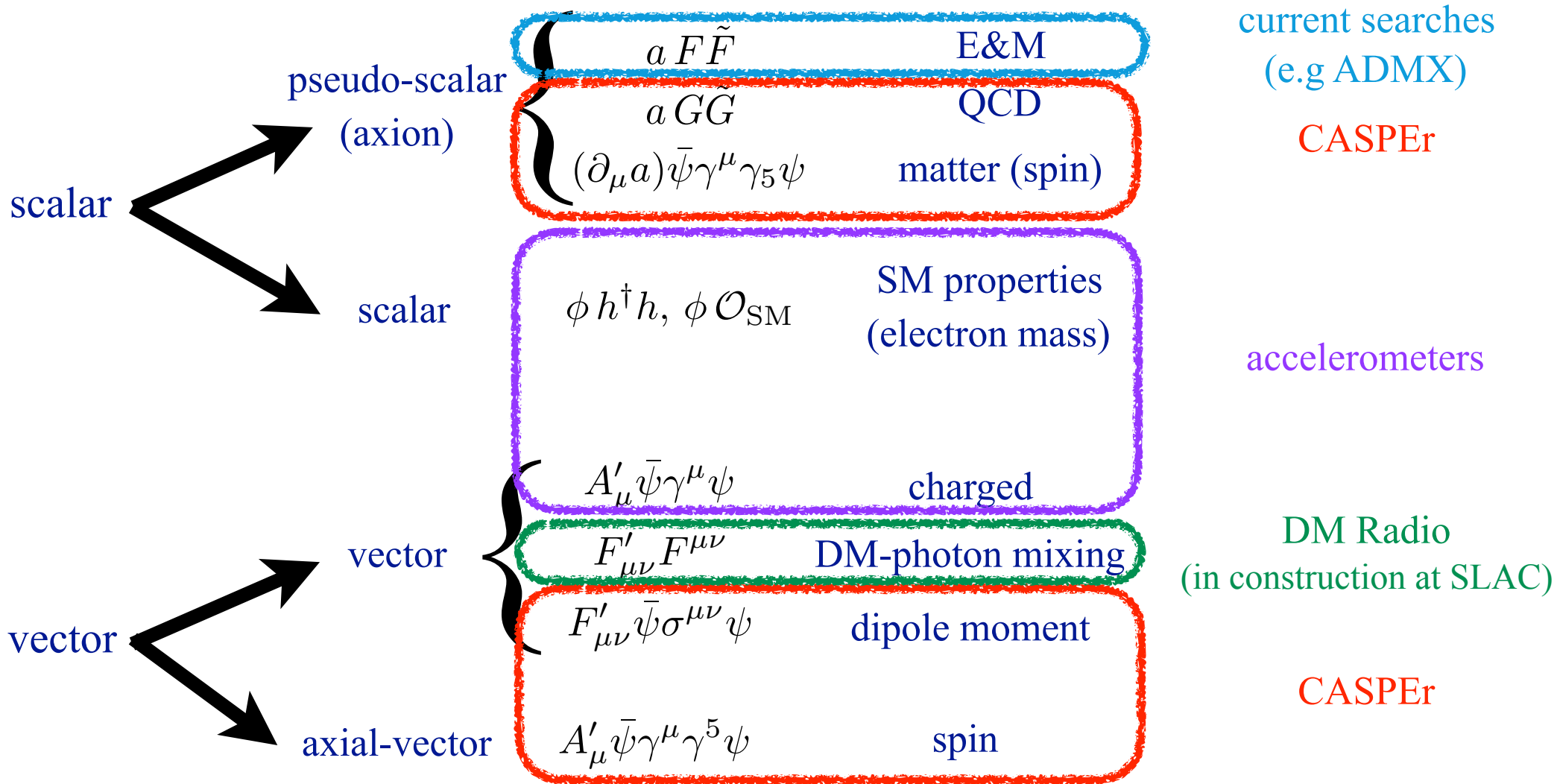
Dmitry Budker

Surjeet Rajendran



# Possibilities for Light Dark Matter

All UV theories summarized by only a few possibilities (symmetry, effective field theory):

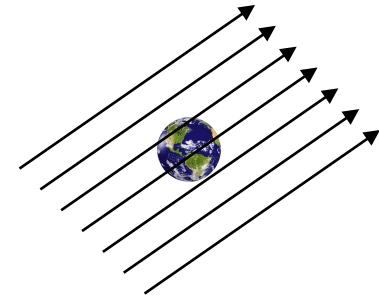


Can cover all these possibilities!



# Force from Dark Matter

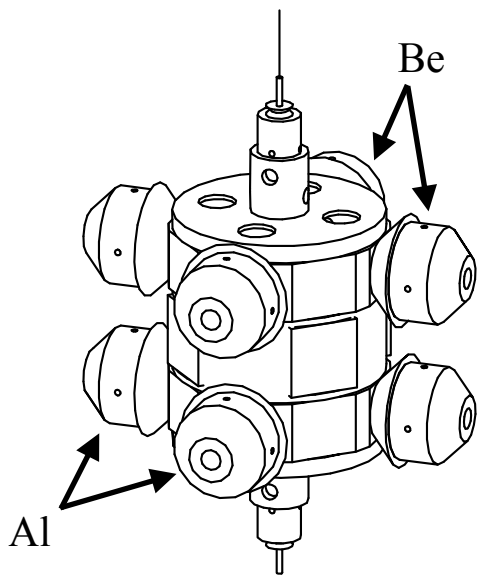
with couplings  $\phi h^\dagger h$ ,  $\phi \mathcal{O}_{\text{SM}}$ ,  $A'_\mu \bar{\psi} \gamma^\mu \psi$  DM acts as a field  
exerts force on matter:  $F \propto g \sqrt{\rho_{\text{DM}}} \cos(m_{\text{DM}} t)$



Force is oscillatory and equivalence-principle violating  
scalar DM would also cause oscillation of “constants” e.g. electron mass

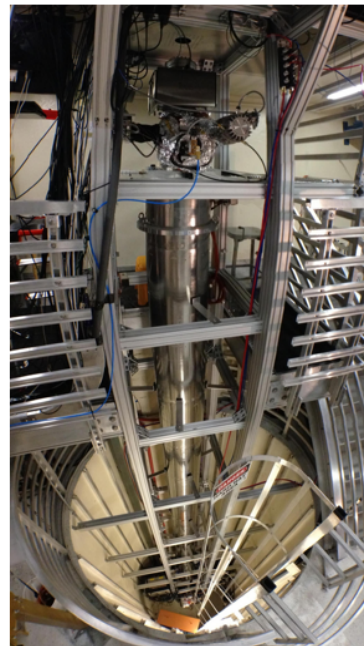
## New Direct Detection Experiments:

### Torsion Balances



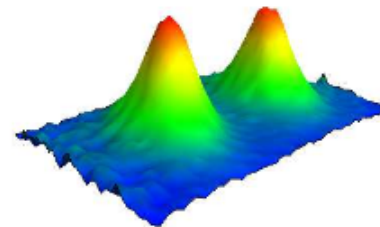
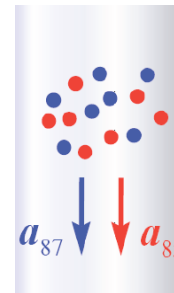
Eot-Wash analysis underway

### Atom Interferometers

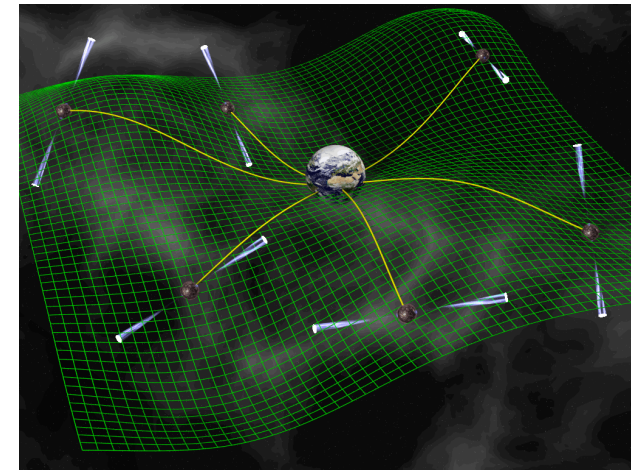


In construction Kasevich/Hogan groups

$^{85}\text{Rb}$ - $^{87}\text{Rb}$



### Pulsar Timing Arrays



Can probe orders of magnitude past current limits

# Summary

Dynamical relaxation provides new class of solutions to hierarchy problem  
physics at weak scale not required

SUSY motivates WIMPs

dynamical relaxation motivates lighter (axion) DM

Precision measurement is a powerful tool for such light fields

new technologies for particle physics beyond traditional particle detectors

1. Cosmic Axion Spin Precession Experiment (CASPER) - in construction at BU and Mainz
2. Accelerometers for DM direct detection - searches by Eot-Wash and Stanford groups
3. DM Radio - in construction at Stanford
4. Atom Interferometry for gravitational wave detection

Many more possibilities...