The LPM effect in sequential gluon bremsstrahlung

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Outline





Overview of the calculation





- At high energies, energy loss is dominated by nearly collinear bremsstrahlung and pair production.
- Naively, the bremsstrahlung rate will be roughly,



where *n* is the density of particles in the medium; σ is the cross section for bremsstrahlung arising from a single collision; and *v* is the relative velocity.

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Complication: The Landau-Pomeranchuk-Migdal effect

- When the formation time of the radiated particle becomes larger than the mean free path in the medium, leading to a significant reduction in the bremsstrahlung rate.
- Consider an *electron* scattering multiple times from the medium and radiating a *photon*.
- The photon cannot resolve details that are smaller than its wavelength. This will create a region of fuzziness, depicted as the blue shaded region.

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Formation time

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Two possibilities interfere quantum mechanically \Rightarrow A reduction in the bremsstrahlung rate from the naive expectation.

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A qualitative difference between QED and QCD formation times

- QED: Softer photon, bigger wavelength \Rightarrow larger formation time.
- QCD: Softer gluon , easier to deflect \Rightarrow shorter formation time.



- LPM effect for QED was figured out in 1950's.
- QCD generalization in 1990's.

So does that solve the problem? Can we just use the LPM result for the rate into a Monte-Carlo-like calculation to study

Image: Image:

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But not when the formation times overlap.



Alternative picture: When we have two splittings within a formation time of each other.

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We would need to calculate interferences of the type



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 α is only moderately small at energy scales reached typically in heavy-ion collisions. Conclusion: Interesting to calculate the effects of overlapping formation times.

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- Previous authors have analyzed this problem in the limiting cases where the radiated gluons have very small energies [Wu, Mehtar-Tani, Blaizot].
- Our calculation goes beyond that approximation and will be valid for general energies of radiated gluons.

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What we do (and what we don't)



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What we do (and what we don't)

• We will focus on the in-medium part of the evolution.



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What we do (and what we don't)

- We will assume the medium to be static, homogeneous and thick, i.e. wider than a typical formation length.
- To simplify the color algebra, we will work in the large *N_C* limit and will focus on the all-gluon case.
- We have calculated the **correction** introduced because of overlapping formation times for general energies of the radiated gluons.

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LPM effect in QCD

In the terms of Feynman diagrams, the LPM effect represents interferences between splittings happening before and after a sequence of elastic scatterings with medium.



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Notation

But for diagrammatic convenience, we will draw



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with interactions with the medium being implicit.

Double Gluon bremsstrahlung:Formalism

• Using this notation, consider one such interference diagram contributing to the double bremsstrahlung rate,



• We will call this the $xy\overline{yx}$ diagram. So the name indicates the order of splittings.

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- For calculations like these, it is helpful to think of these terms as being a single "process".
- In time between the splittings, the number of high energy particles doesn't change.



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• So the problem reduces to sewing QCD splitting matrix elements with the evolution in these three regions.

Complications: 4-particle evolution



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Need to calculate Eigenfrequencies and then, QM HO propagators.

More Complications and the result

- Non-Trivial Helicity contractions: Must use helicity dependent DGLAP splitting functions.
- **Divergences:** Diagrams involving two 3-gluon vertices diverge as △*t* → 0. The divergences cancel when contribution from different diagrams are added together, but we still need to carefully calculate the non-trivial contributions from the poles.

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More Complications and the result

- A lot of diagrams to calculate!.
- Many different permutations and time orderings.



Result = $\int \Delta t$ (Complicated expression)

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• Final Δt integral easy to do numerically.

What we have done

- Completed the calculation of the correction introduced by the considering overlapping formation times.
- P. Arnold and Shahin Iqbal, "The LPM effect in sequential bremsstrahlung", JHEP 04(2015)070
- P. Arnold, H.C. Chang and Shahin Iqbal, "The LPM effect in sequential bremsstrahlung 2: factorization" arXiv:1605.07624 [hep-ph].
- P. Arnold, H.C. Chang and Shahin Iqbal, "The LPM effect in sequential bremsstrahlung: dimensional regularization" arXiv:1606.08853 [hep-ph].
- P. Arnold, H.C. Chang and Shahin Iqbal, "The LPM effect in sequential bremsstrahlung: 4-gluon vertices" arXiv:1608.05718 [hep-ph].
- Can reproduce the results of previous authors in the relevant limiting cases.

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What we find

- The correction is negative only when one of the three daughters is much softer than the the other two.
- In in the limiting case $y \ll x \ll 1$, the correction introduced by considering the effects of overlapping formation times is given by



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Why is the correction negative?

For $y \ll x \ll 1$, x and z are so close to each other that y sees them as if they were a *single* adjoint color particle.



A naive Monte-Carlo calculation treating these splittings as independent would count both possibilities and therefore, double count the chance of *y* emission.

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Work in progress

• Calculating the effects of overlapping formation times in QED.



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3- R. Baier, Y.L. Dokshitzer, A.H. Mueller, S. Peigne and D. Schiff, Radiative energy loss and pT broadening of high-energy partons in nuclei, Nucl. Phys. B 484 (1997) 265 [hep-ph/9608322] [INSPIRE].

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4- J.P. Blaizot and Y. Mehtar-Tani, Renormalization of the jet-quenching parameter, Nucl. Phys. A 929 (2014) 202 [arXiv:1403.2323] [INSPIRE].

5- B. Wu, Radiative energy loss and radiative pt-broadening of high-energy partons in QCD matter, JHEP 12 (2014) 081 [arXiv:1408.5459] [INSPIRE].

Large N limit

- In the 3-particle region, the constraint $T_1 + T_2 + T_3 = 0$ determines all of the T_i . T_j .
- In the 4-particle region, the constraint $T_1 + T_2 + T_3 + T_4 = 0$ cannot determine all $T_i \cdot T_j$.
- Things become much more simplified in the large N limit.



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A Problem

- Integrand $\propto \frac{\text{stuff}}{\Delta t^2} + \frac{\text{stuff}}{\Delta t}$ as $\Delta t \to 0$.
- $\frac{1}{\Delta t^2}$ divergence is canceled by subtracting the vacuum contribution.
- We find that the $\frac{1}{\Delta t}$ divergence was arising in the limit when three of the 4 times t_i approached each other simultaneously.

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• All $\frac{1}{\Delta t}$ divergences get canceled when all six time orderings are added.

A subtlety!

Even though the divergence cancels, one needs to be careful about the *i* e prescriptions.

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• For example,

$$\int_{-\infty}^{+\infty} dt (\frac{1}{t} - \frac{1}{t}) = 0$$

but

$$\int_{-\infty}^{+\infty} dt (\frac{1}{t-i\epsilon} - \frac{1}{t+i\epsilon}) \neq 0$$

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• The different diagrams turn out to get different $i\epsilon$ prescriptions and therefore,we also had to calculate the contributions from these poles.

LPM effect in QED vs QCD



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