

# PLAN

- **Lecture I:** Neutrinos in the SM  
Neutrino masses and mixing: Majorana vs Dirac
- **Lecture II:** Neutrino oscillations and the discovery of neutrino masses and mixings
- **Lecture III:** The quest for the unknowns  
A neutrino look at BSM and the history of the Universe

# SM+3 massive neutrinos

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{PMNS}(\theta_{12}, \theta_{23}, \theta_{13}, \delta, \dots) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

More fundamental parameters in the SM: **3 angles, 3 masses, 1 or 3 phases**

$$U_{PMNS} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \underbrace{\begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1} & 0 \\ 0 & 0 & e^{i\alpha_2} \end{pmatrix}}_{\text{Only Majorana}}$$

$c_{ij} \equiv \cos \theta_{ij} \quad s_{ij} \equiv \sin \theta_{ij}$

Imaginary entries in leptonic mixing matrix imply violation of CP!

# CP violation in oscillations

$$P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$
$$\alpha \neq \beta$$

$$P(\nu_\alpha(\bar{\nu}_\alpha) \rightarrow \nu_\beta(\bar{\nu}_\beta)) = -4 \sum_{i < j} \text{Re}[U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*] \sin^2 \left[ \frac{\Delta m_{ji}^2 L}{4E} \right] \quad \text{CP-even}$$
$$\mp 2 \sum_{i < j} \text{Im}[U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*] \sin \left[ \frac{\Delta m_{ji}^2 L}{4E} \right] \quad \text{CP-odd}$$

# SM+3 massive neutrinos: Global Fits

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{PMNS}(\theta_{12}, \theta_{23}, \theta_{13}, \delta, \dots) \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

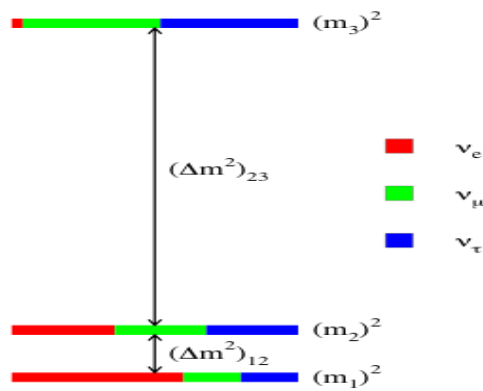
$$\theta_{12} \sim 34^\circ$$

$$\theta_{23} \sim 42^\circ \text{ or } 48^\circ$$

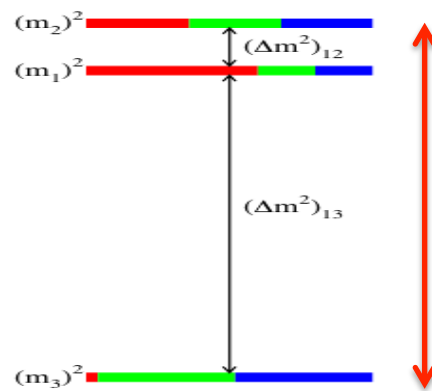
$$\theta_{13} \sim 8.5^\circ$$

$$\delta \sim ?$$

normal hierarchy



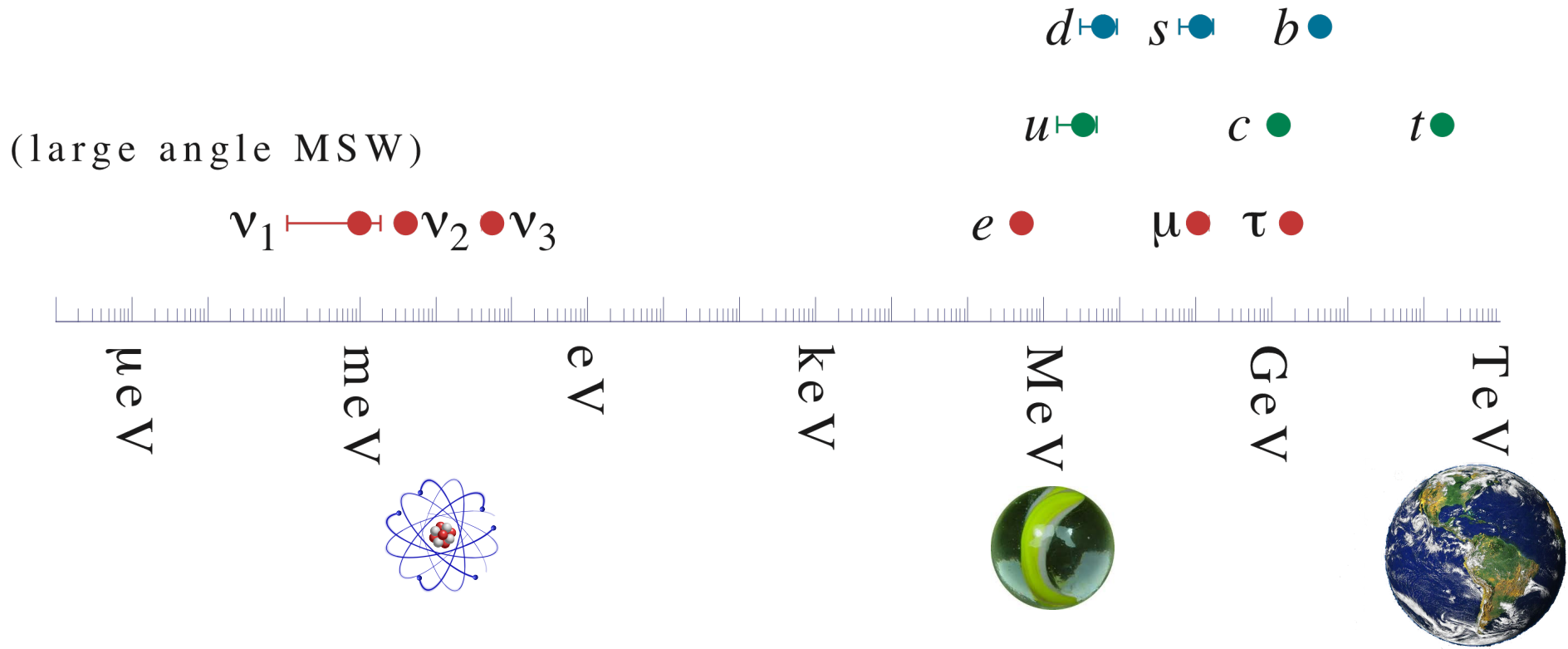
inverted hierarchy



$$\updownarrow 7.5 \cdot 10^{-5} \text{eV}^2$$

$$2.5 \cdot 10^{-3} \text{eV}^2$$

# The flavour puzzle



# Why so different mixing ?

## CKM

$$|V|_{\text{CKM}} = \begin{pmatrix} 0.97427 \pm 0.00015 & 0.22534 \pm 0.0065 & (3.51 \pm 0.15) \times 10^{-3} \\ 0.2252 \pm 0.00065 & 0.97344 \pm 0.00016 & (41.2_{-5}^{+1.1}) \times 10^{-3} \\ (8.67_{-0.31}^{+0.29}) \times 10^{-3} & (40.4_{-0.5}^{+1.1}) \times 10^{-3} & 0.999146_{-0.000046}^{+0.000021} \end{pmatrix}$$

PDG

## PMNS

$$|U|_{3\sigma}^{\text{LID}} = \begin{pmatrix} 0.798 \rightarrow 0.843 & 0.517 \rightarrow 0.584 & 0.137 \rightarrow 0.158 \\ 0.232 \rightarrow 0.520 & 0.445 \rightarrow 0.697 & 0.617 \rightarrow 0.789 \\ 0.249 \rightarrow 0.529 & 0.462 \rightarrow 0.708 & 0.597 \rightarrow 0.773 \end{pmatrix}$$

NuFIT 2016

# Why so different mixing ?

CKM

$$V_{CKM} \simeq \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

PMNS

$$|V_{PMNS}| \simeq \begin{pmatrix} \sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} & 0 \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \\ \sqrt{\frac{1}{6}} & \sqrt{\frac{1}{3}} & \sqrt{\frac{1}{2}} \end{pmatrix}$$

# Where the mixing comes from ?



Discrete or continuous symmetries

Anarchy for leptons ?



# The open questions

Absolute mass scale: minimum  $m_\nu$

What is the **neutrino ordering** normal or inverted ?

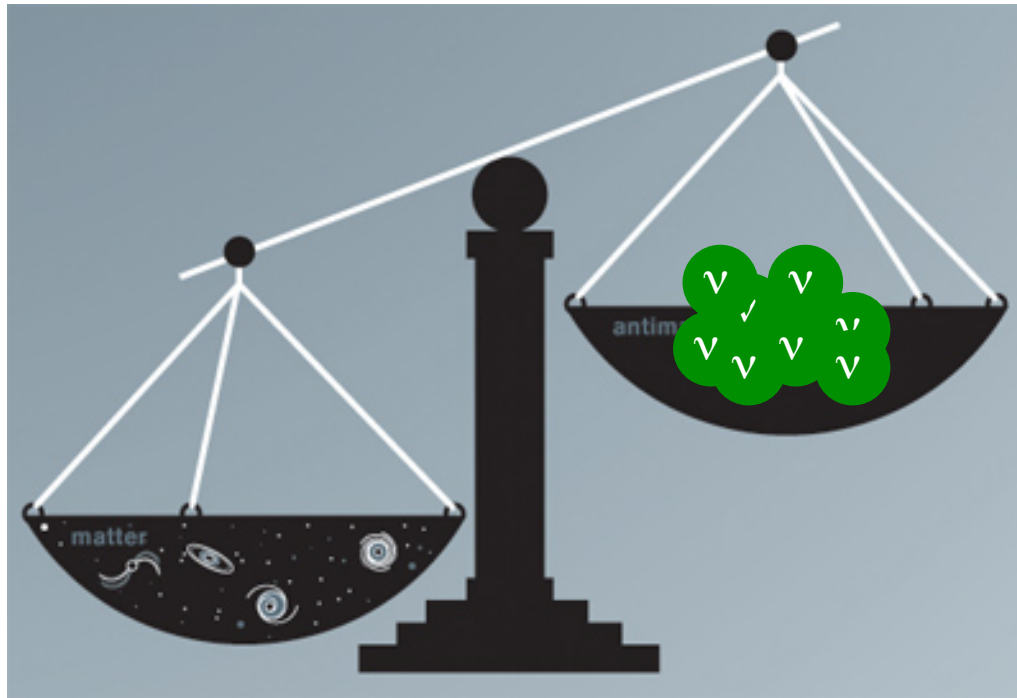
Is there **leptonic CP violation** ?

Are neutrinos **Majorana** and if so, what **new physics** lies behind this fact ?

Can neutrinos explain the **matter-antimatter asymmetry** in the Universe ?

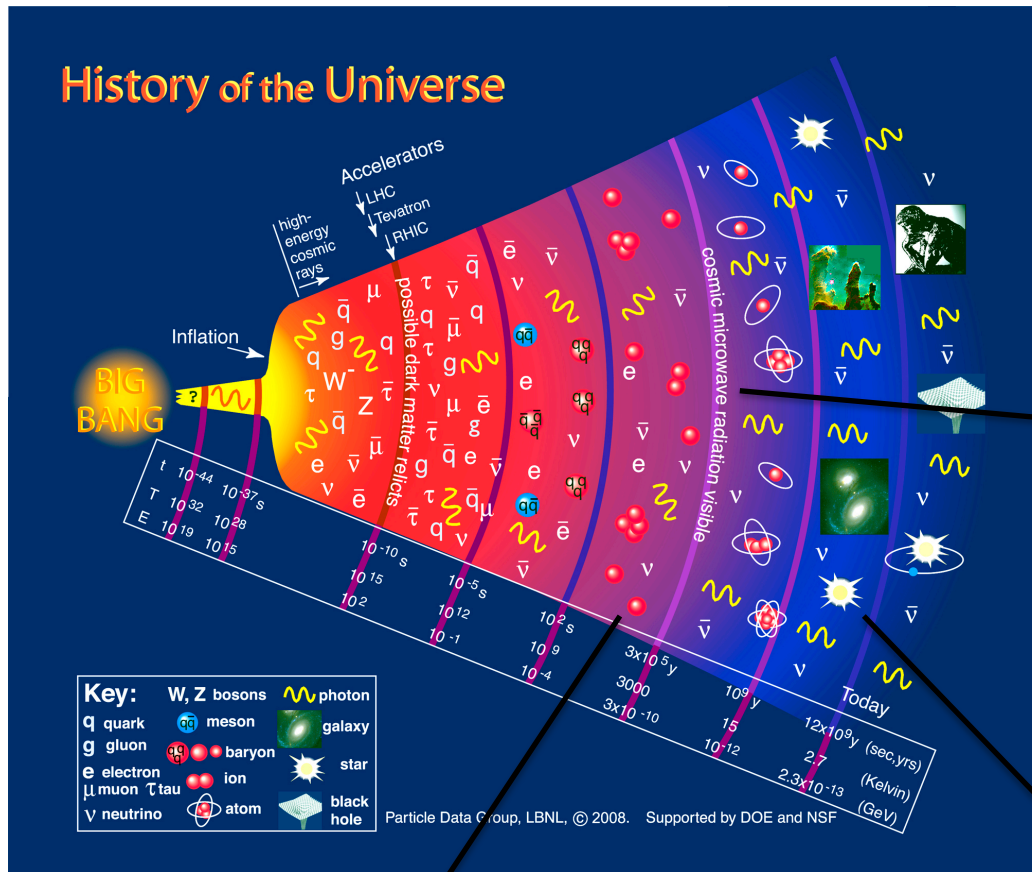
# Absolute mass scale

Best constraints at present from cosmology



# Cosmological neutrinos

Neutrinos have left many traces in the history of the Universe



Cosmic Microwave background

Galaxy distribution (LSS)

Nucleosynthesis (synthesis of light nuclei)

# Neutrinos @ nucleosynthesis (BBN)

Before LEP, the best constraint on  $N_\nu$  came from Big Bang nucleosynthesis!

Neutrinos decouple from the plasma @  $T_\nu \sim 1\text{MeV}$

$$\text{Rate} \left( \begin{array}{c} \nu_e n \leftrightarrow p e^- \\ e^+ n \leftrightarrow p \bar{\nu}_e \end{array} \right) \simeq \text{Expansion rate}(g_*(T))$$

$g_*(T) \equiv \#$ of relativistic degrees of freedom

Each neutrino species counts like one relativistic d.o.f.:

$g^*$  depends on  $N_\nu \Rightarrow T_\nu(N_\nu)$

# Neutrinos @nucleosynthesis

At  $@T_\nu \sim 1\text{MeV}$  the ratio neutrons/protons freezes and light elements start to form:

$$\frac{N_n}{N_p} = \exp\left(\frac{m_p - m_n}{T_\nu(N_\nu)}\right)$$

The abundance of light nuclei depends strongly on the ratio of n/p

$$Y_{4\text{He}} = \frac{\text{Mass of } ^4\text{He}}{\text{Total Mass}} = \frac{2N_n}{N_p + N_n}$$

# Neutrinos as DM

Neutrino distribution gets frozen at BBN when they are still relativistic

$$N_\nu \simeq N_{\bar{\nu}} \simeq \frac{4}{11} T_\gamma^3$$

$$\Omega_\nu = \frac{\sum_i m_i}{93.5 \text{ eV}} h^{-2} < \Omega_m \rightarrow \sum_i m_i \leq 11.2 \text{ eV}$$

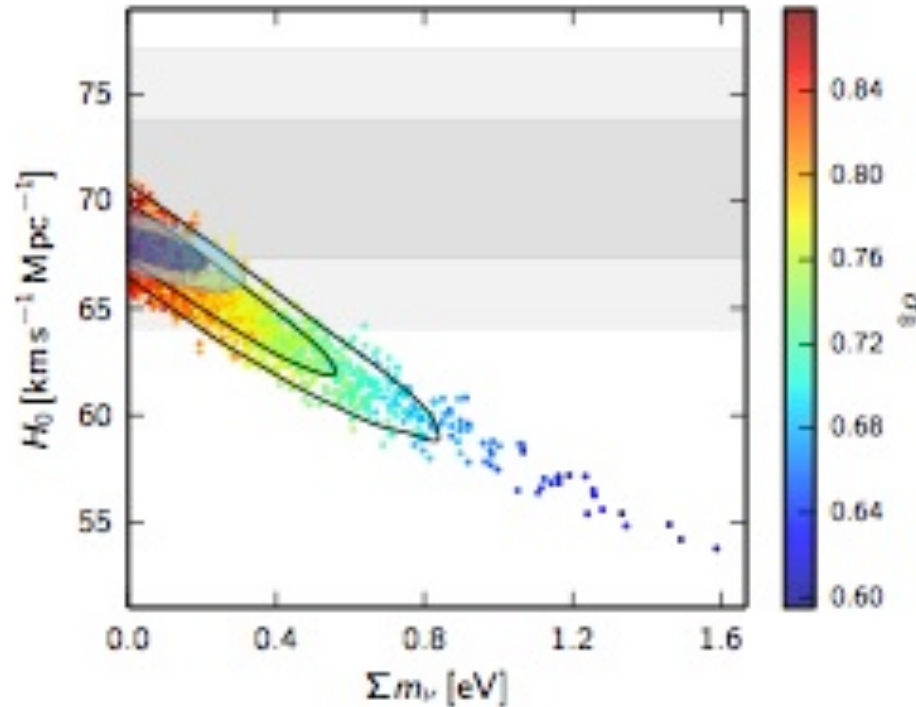
Gershtein, Zeldovich

Massive neutrinos O(eV) contribute significantly to  $\Omega_m$

They tend to produce a Universe with too little structure at small scales

# Absolute mass scale

Neutrinos as light as 0.1-eV modify the large scale structure and CMB



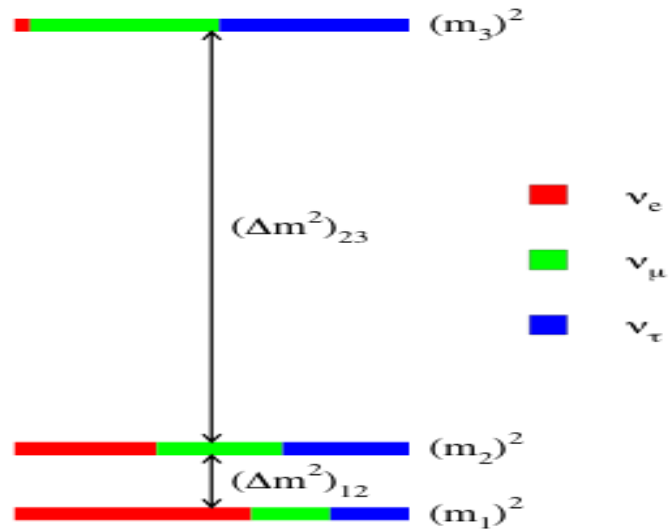
Planck 1502.01589

Conservative limit:

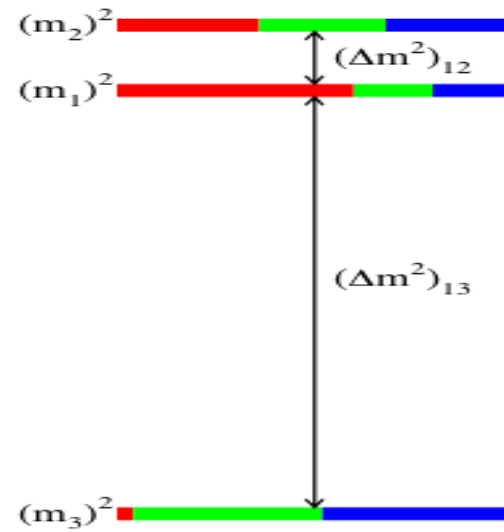
$$\left. \begin{array}{l} \sum m_\nu < 0.23 \text{ eV} \\ \Omega_\nu h^2 < 0.0025 \end{array} \right\} 95\%, \text{ Planck TT+lowP+lensing+ext.}$$

?

normal hierarchy

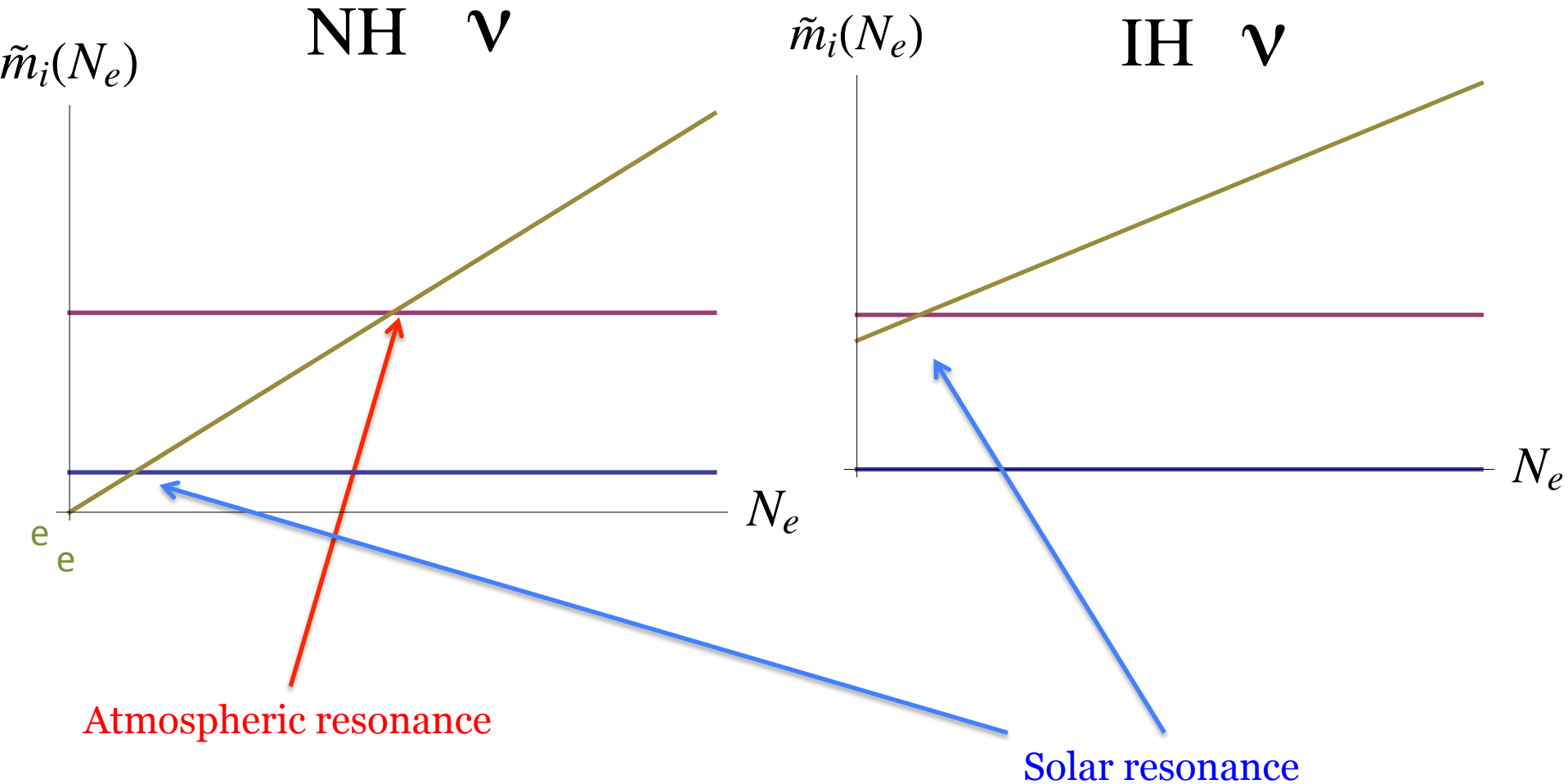


inverted hierarchy





# Neutrino ordering from MSW



$$\Delta m_{23}^2 \cos 2\theta_{13} = \pm 2\sqrt{2}G_F E N_e$$

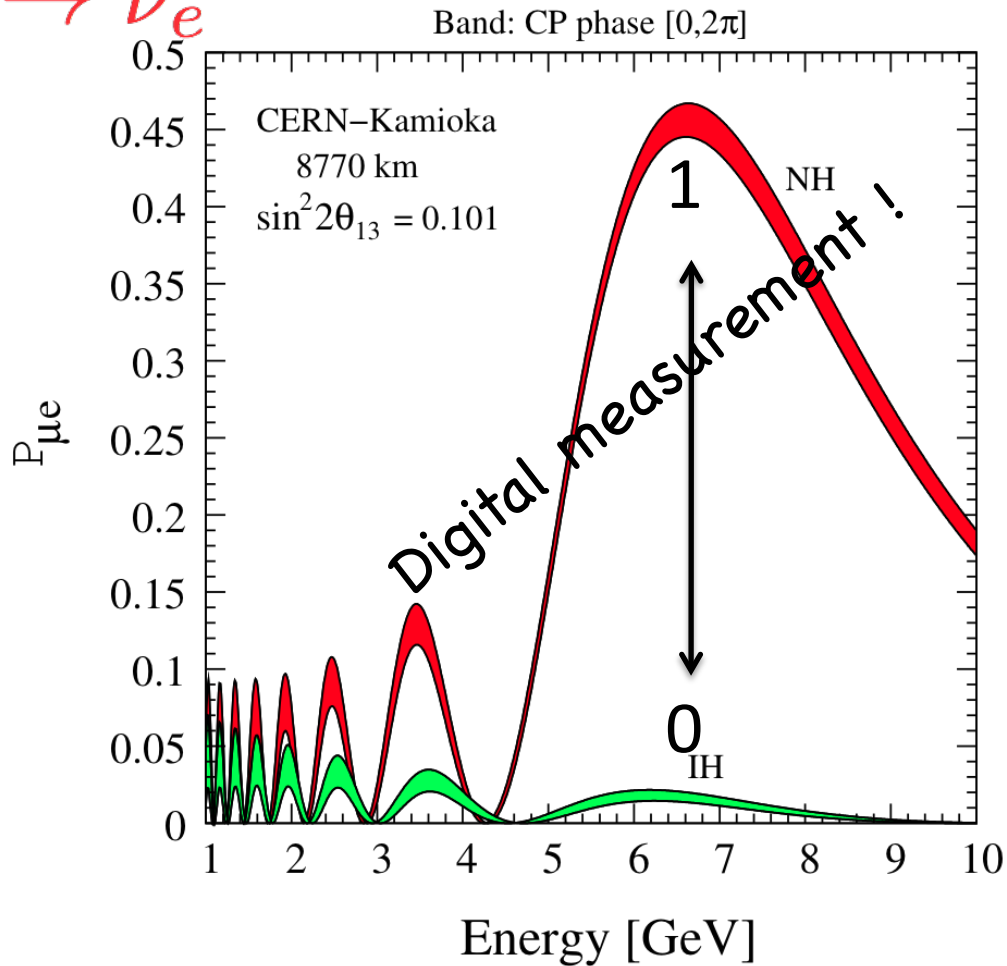
Earth density,  $E_{\text{res}} \sim \text{few GeV} !$

$$\Delta m_{12}^2 \cos 2\theta_{12} = 2\sqrt{2}G_F E N_e$$

Solar density,  $E_{\text{res}} \sim \text{few MeV} !$

# Hierarchy through MSW @Earth

$\nu_\mu \rightarrow \nu_e$

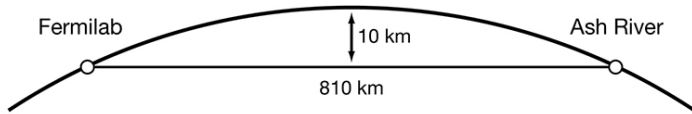


$$E_{\text{res}} \equiv \frac{\Delta m_{31}^2 \cos 2\theta_{13}}{2\sqrt{2}G_F n_e},$$

$$n_e(L)L|_{L_{\text{max}}} = \frac{\pi}{\sqrt{2}G_F \tan 2\theta_{13}}$$

Spectacular MSW effect at O(6GeV) and very long baselines

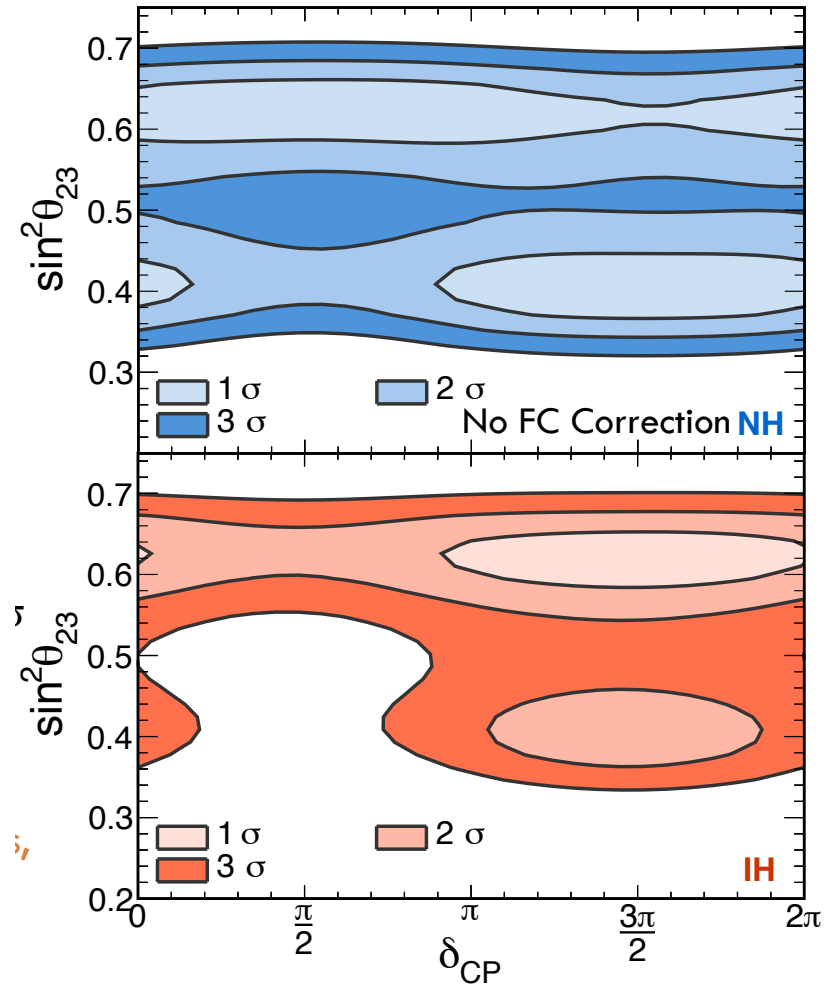
# First attempt at the hierarchy: $\text{NO}\nu\text{A}$



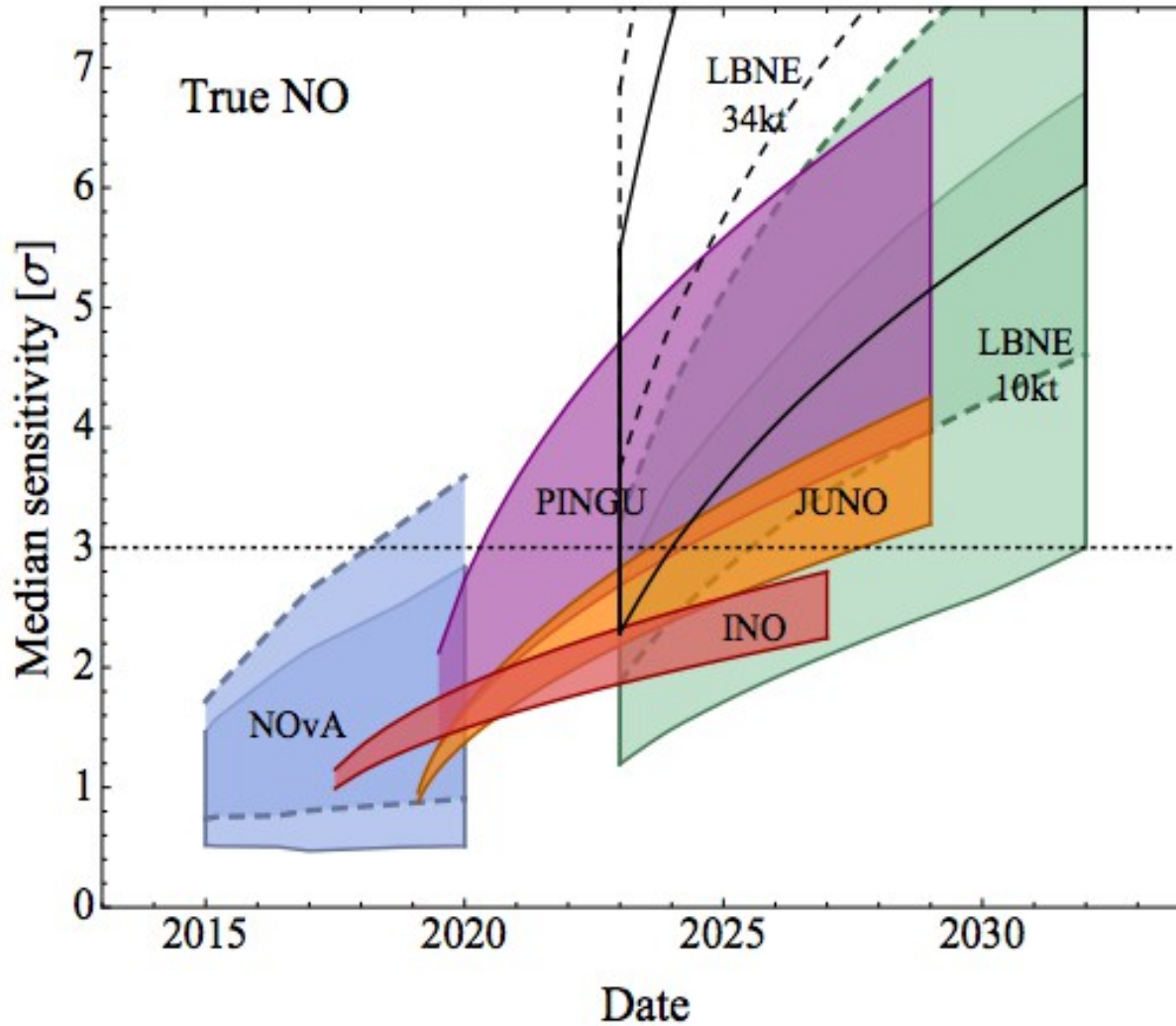
$L=810\text{km}$

$\nu_\mu \rightarrow \nu_e$

NOνA Preliminary



# Hierarchy projects



# Leptonic CP violation

CP violation shows up in a difference between

$$P(\nu_\alpha \rightarrow \nu_\beta) \neq P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \quad \alpha \neq \beta$$

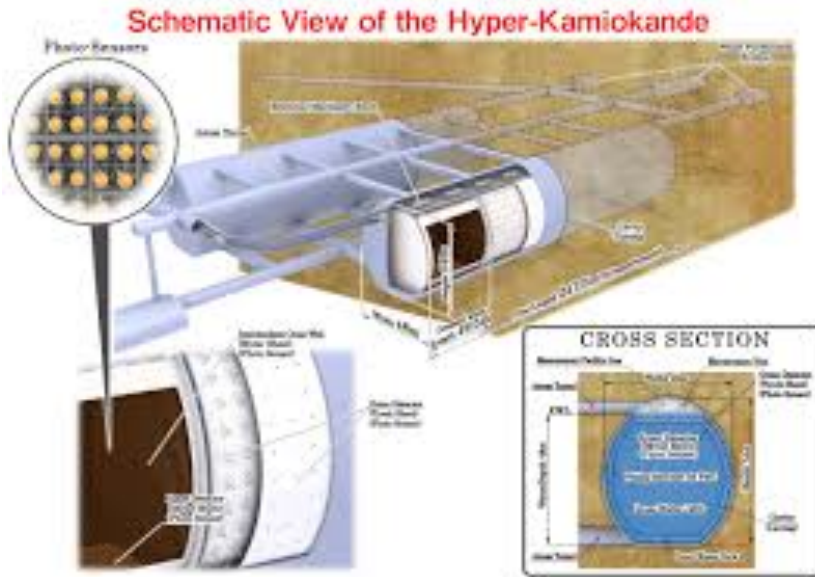
Golden channel:

$$\begin{aligned} P_{\nu_e \nu_\mu}(\bar{\nu}_e \bar{\nu}_\mu) &= s_{23}^2 \sin^2 2\theta_{13} \sin^2 \left( \frac{\Delta_{23} L}{2} \right) \equiv P^{atmos} \\ &+ c_{23}^2 \sin^2 2\theta_{12} \sin^2 \left( \frac{\Delta_{12} L}{2} \right) \equiv P^{solar} \\ &+ \tilde{J} \cos \left( \pm\delta - \frac{\Delta_{23} L}{2} \right) \frac{\Delta_{12} L}{2} \sin \left( \frac{\Delta_{23} L}{2} \right) \equiv P^{inter} \\ &\tilde{J} \equiv c_{13} \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \end{aligned}$$

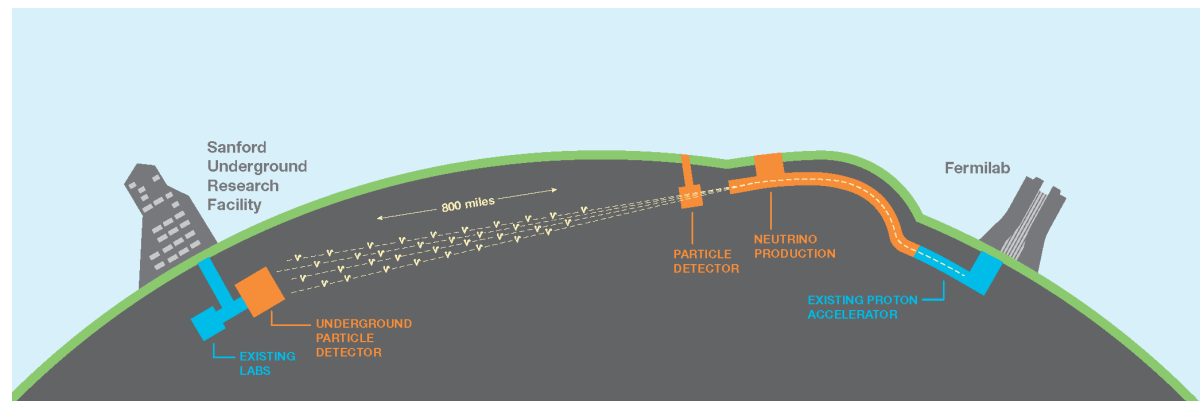
simultaneous sensitivity to both splittings is needed

# Hierarchy + CP in one go... superbeams+superdectors

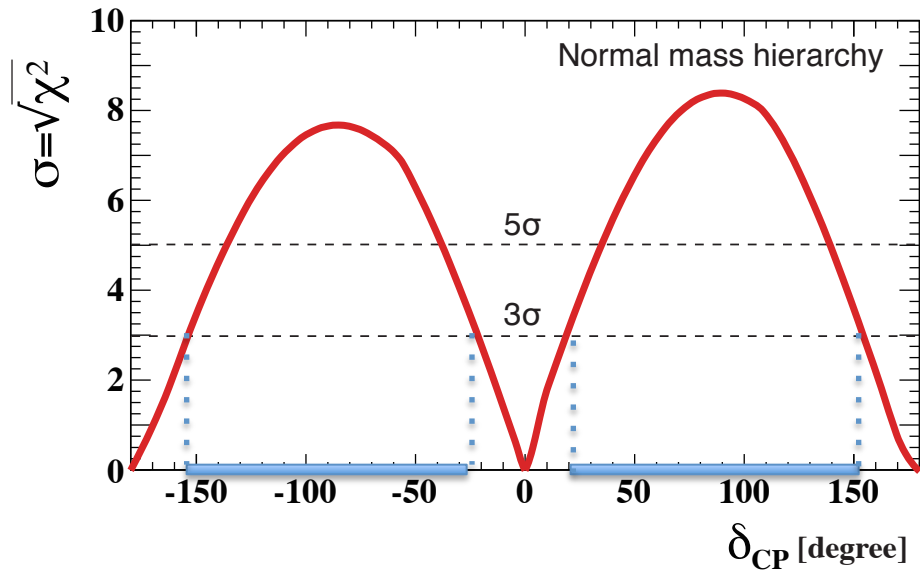
Japan Hyper-Kamiokande: 230km



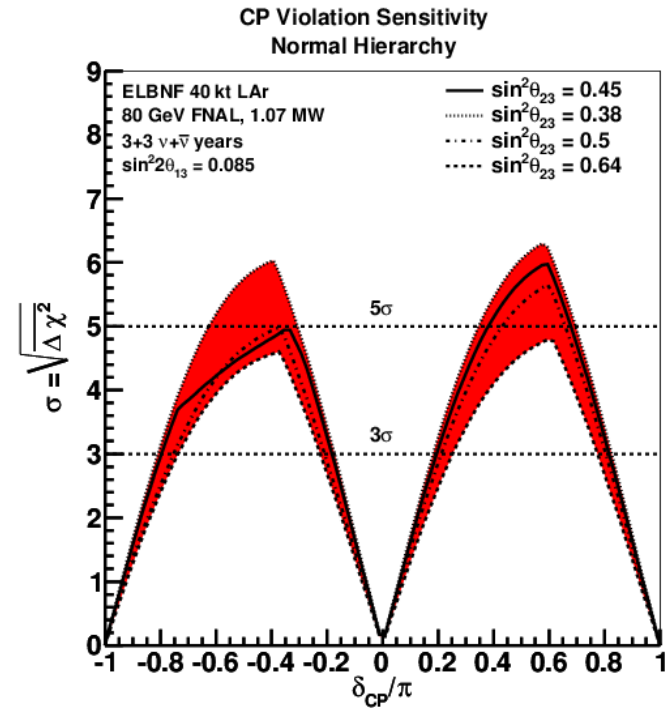
USA DUNE: 1300km



# Leptonic CP violation

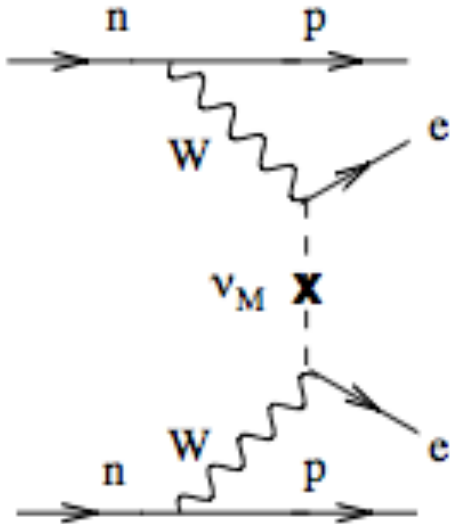


Hyper Kamiokande

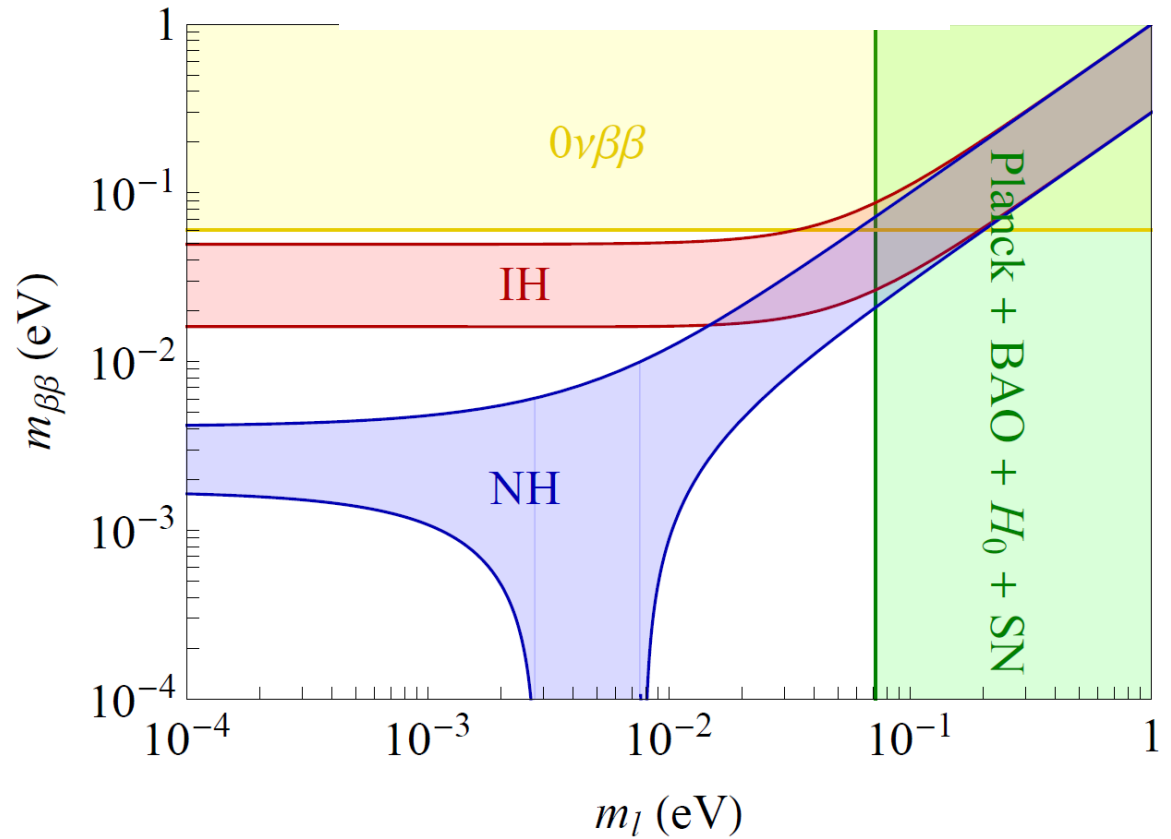


DUNE

# Majorana nature ?

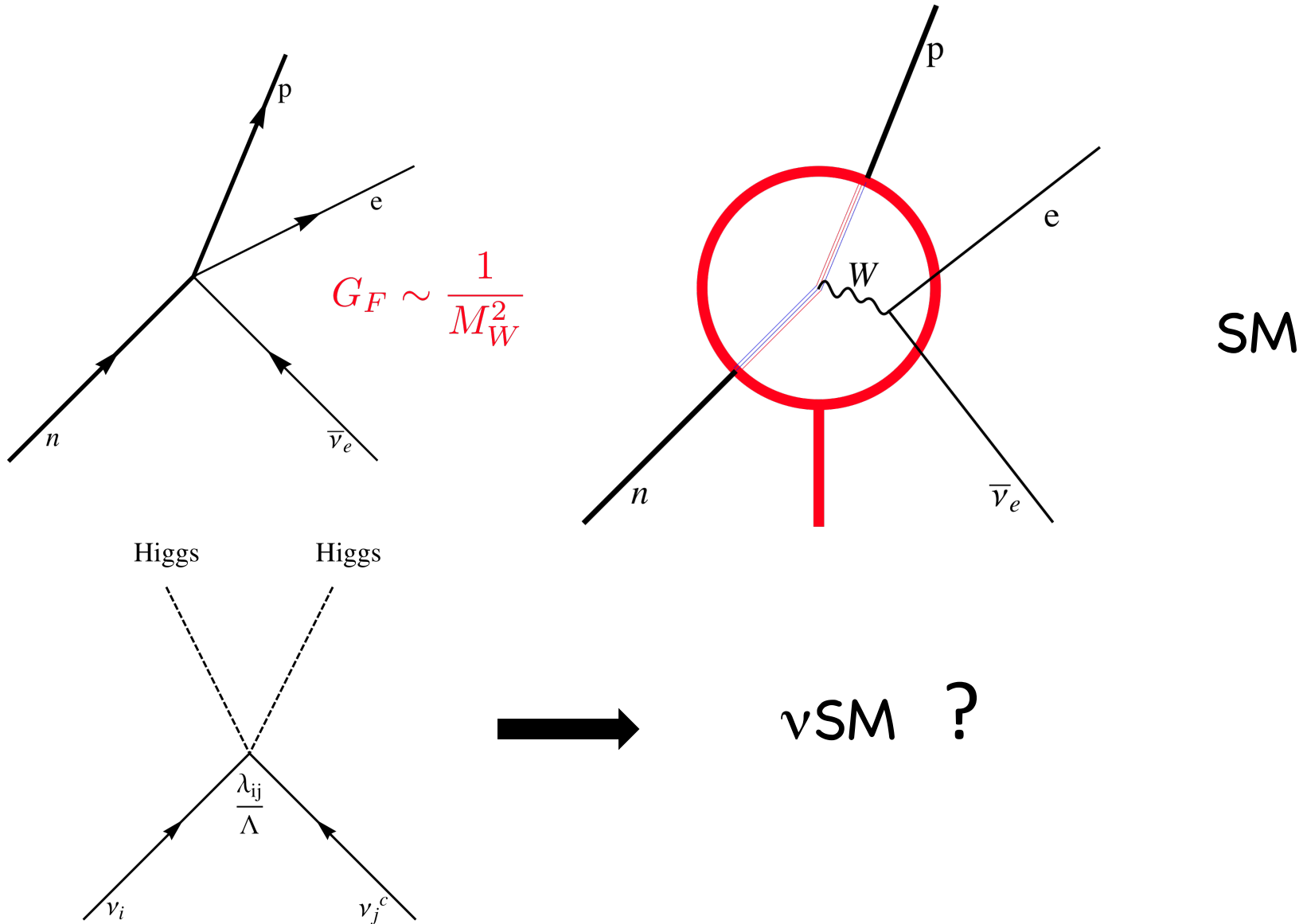


$$m_{\beta\beta} = \sum_{i=1}^3 [(U_{PMNS})_{ei}]^2 m_i$$





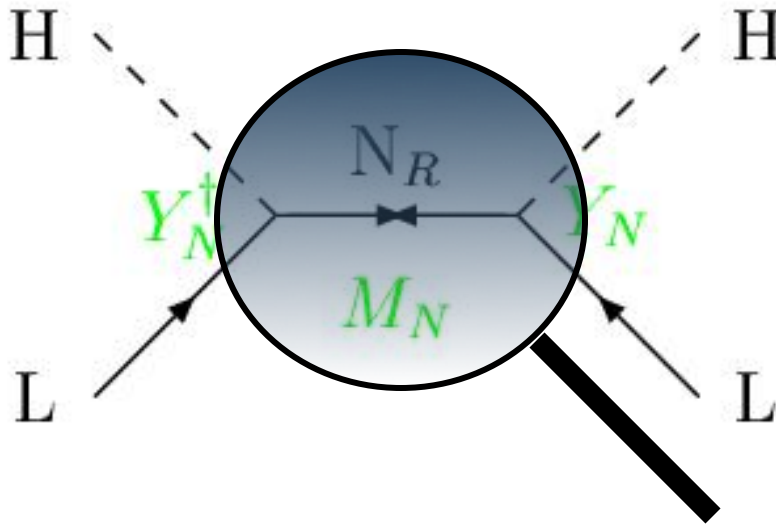
# Majorana neutrinos -> a new physics scale



# Neutrino BSM: the Seesaw Model

SM + heavy singlet fermions

$$\mathcal{L} = \mathcal{L}_{SM} - \sum_{i=1}^{n_R} \bar{l}_L^\alpha Y^{\alpha i} \tilde{\Phi} \nu_R^i - \sum_{i,j=1}^{n_R} \frac{1}{2} \bar{\nu}_R^{ic} M_N^{ij} \nu_R^j + h.c.$$

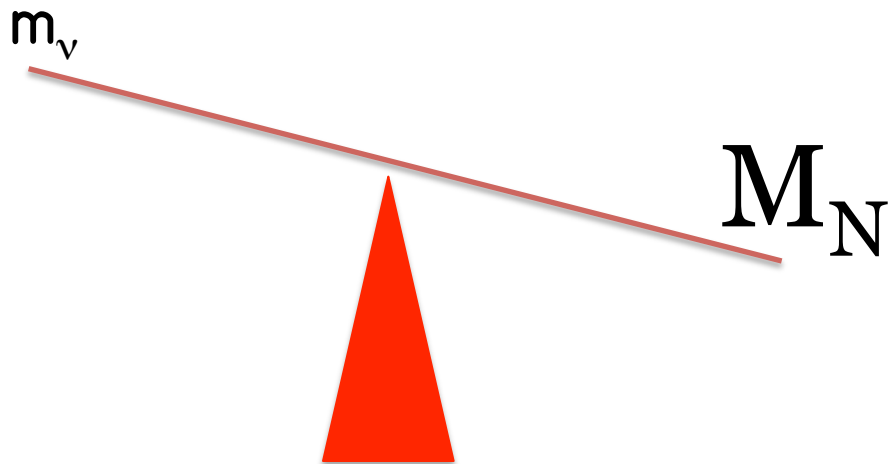


$$\frac{\lambda}{\Lambda} \equiv Y_N^T \frac{1}{M_N} Y_N$$

# Neutrino BSM: the **Seesaw** Model

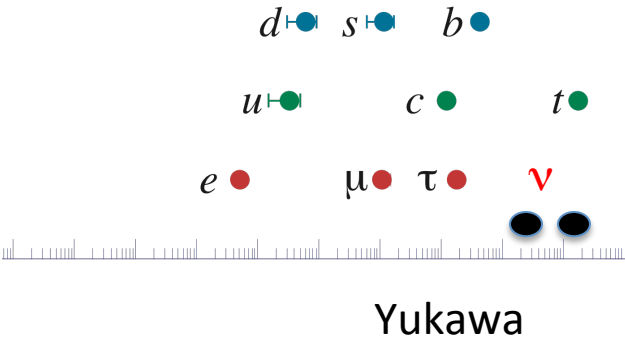
SM + 3 singlet Majorana fermions

$$\mathcal{L} = \mathcal{L}_{SM} - \sum_{i=1}^{n_R} \bar{l}_L^\alpha Y^{\alpha i} \tilde{\Phi} \nu_R^i - \sum_{i,j=1}^{n_R} \frac{1}{2} \bar{\nu}_R^{ic} M_N^{ij} \nu_R^j + h.c.$$

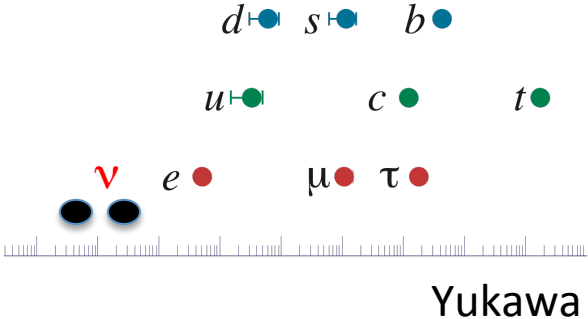

$$m_\nu \propto Y_N^T \frac{v^2}{M_N} Y_N^T$$

# Pinning down the new physics scale: naturalness

$M_N \sim \text{GUT}$

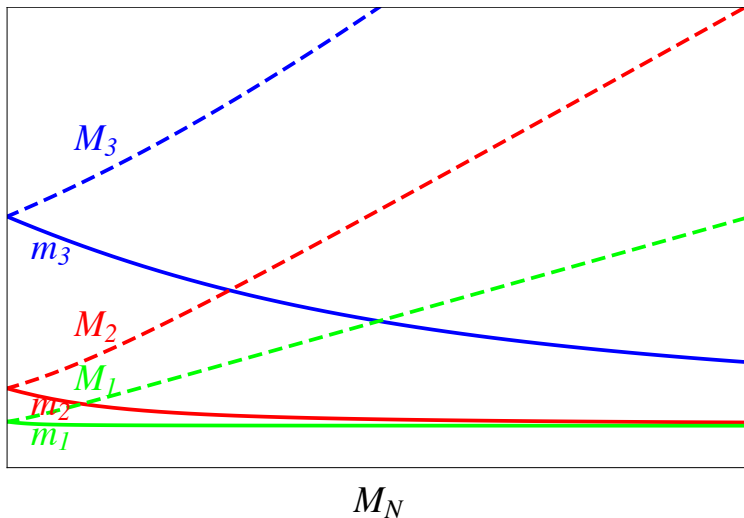


$M_N \sim \text{TeV}$



# Spectra of the seesaw model

$n_R = 3$  : 18 free parameters (6 masses+6 angles+6 phases) out of which we have measured 2 mass differences and 3 angles...



$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = U_{ll} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} + U_{lh} \begin{pmatrix} N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

$$U_{ll} \simeq U_{PMNS} + \mathcal{O}\left(\frac{m_l}{M}\right)$$

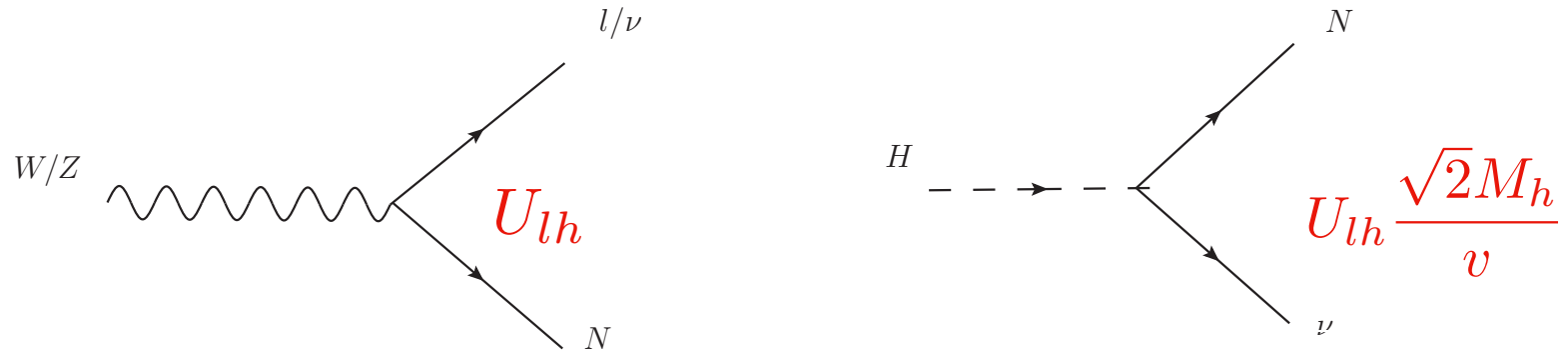
$$U_{lh} \propto \sqrt{\frac{m_l}{M_N}}$$

Unitarity Violations



Strong correlation between active-heavy mixing and neutrino masses

# Spectra of seesaw models



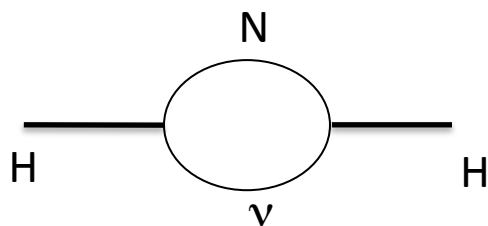
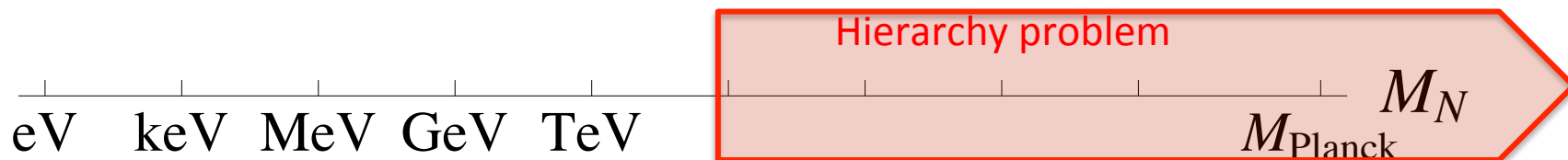
The lower the seesaw scale  $M_N$

- the lower kinematical threshold to produce new states
- the larger mixing with flavour states

# Searching for heavy Majorana neutrinos



# Searching for heavy Majorana neutrinos



$$\delta m_H^2 = \frac{Y^\dagger Y}{4\pi^2} M_N^2 \log \frac{M_N}{\mu}$$

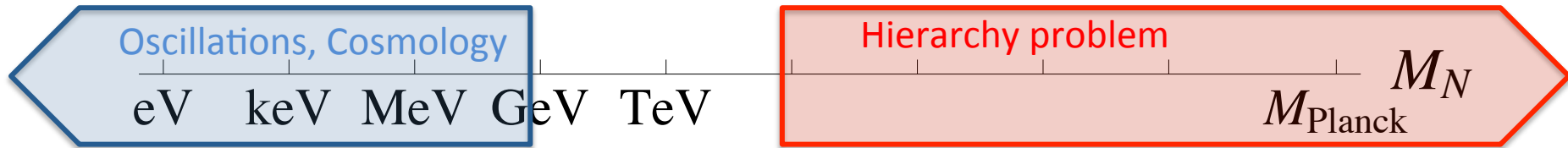
$$M_N \gg m_H$$

Requires a fine-tuning of the Higgs mass in the absence of other physics, like SUSY



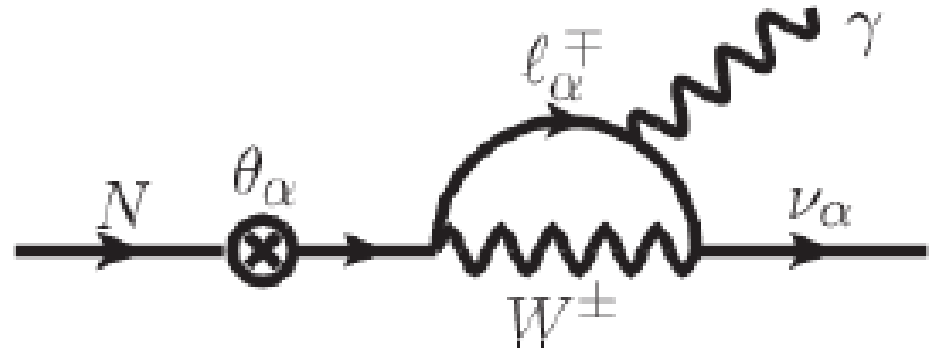
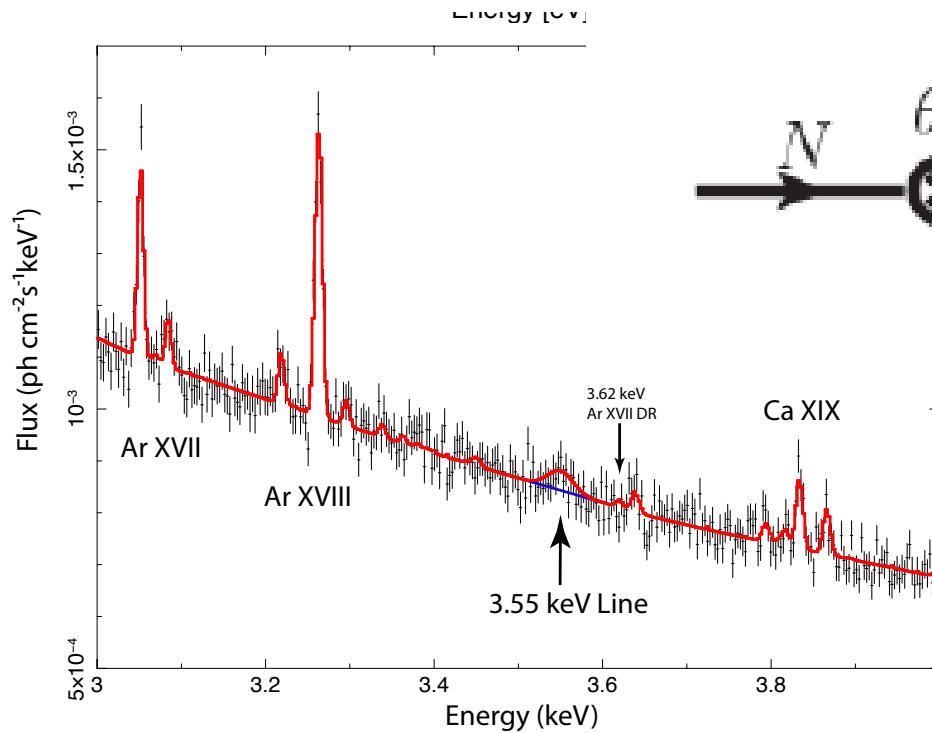
# Searching for heavy Majorana neutrinos

These heavy states can contribute as DM ( $\Omega_{\text{DM}}$ ) and/or extra radiation ( $\Delta N_{\text{eff}}$ )



**One exception:** the lighter of the heavy states could be at keV and be the Dark Matter

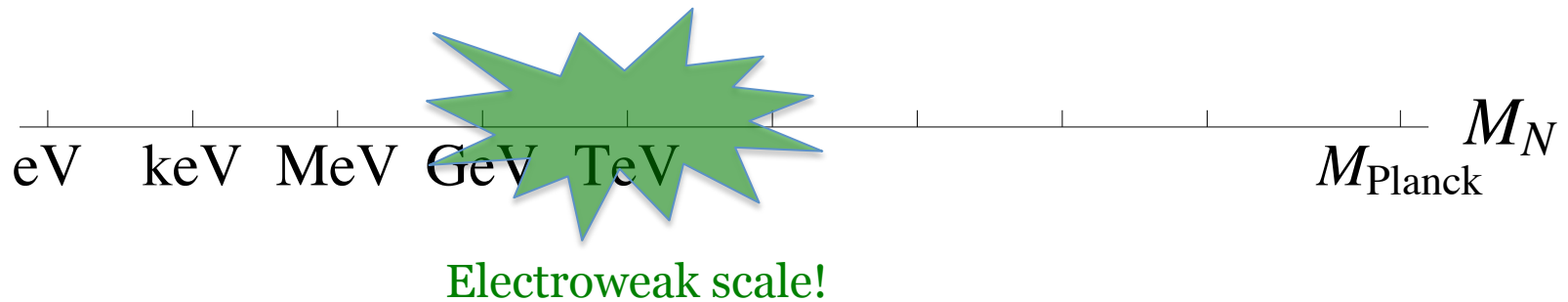
# Neutrino Warm Dark Matter ?



$$E_\gamma = \frac{M_N}{2}$$

$$M_N \simeq 7 \text{ keV}$$

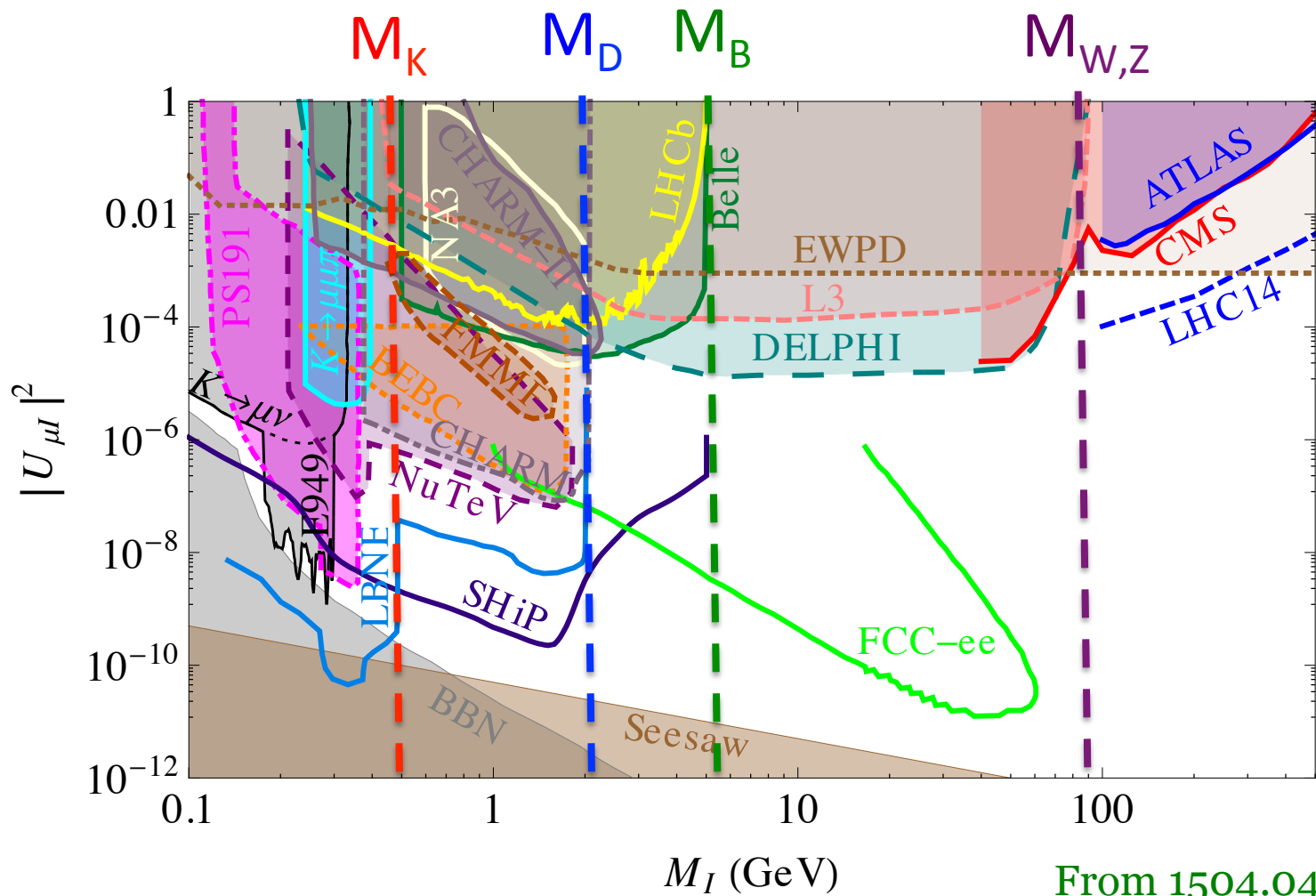
# Searching for heavy Majorana neutrinos



Neutrinos in the range  $\text{GeV} \leq M_N \leq \text{few TeV}$  can modify:

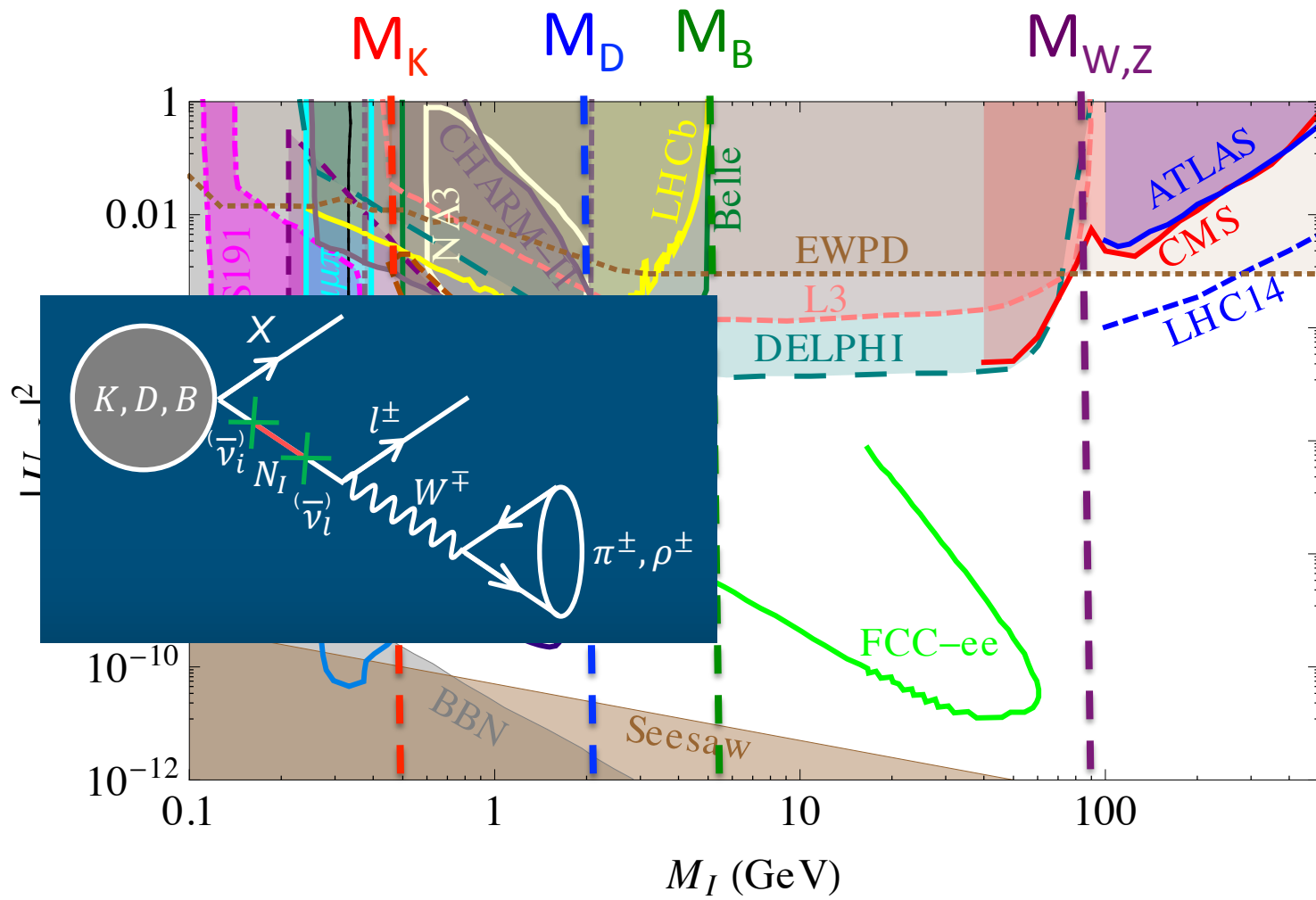
- ✓ Neutrinoless double beta decay
- ✓ Can be directly produced in accelerators: meson decays, W,Z decays
- ✓ Lepton flavour violating processes e.g.  $\mu \rightarrow e\gamma$

# Searching for heavy Majorana neutrinos

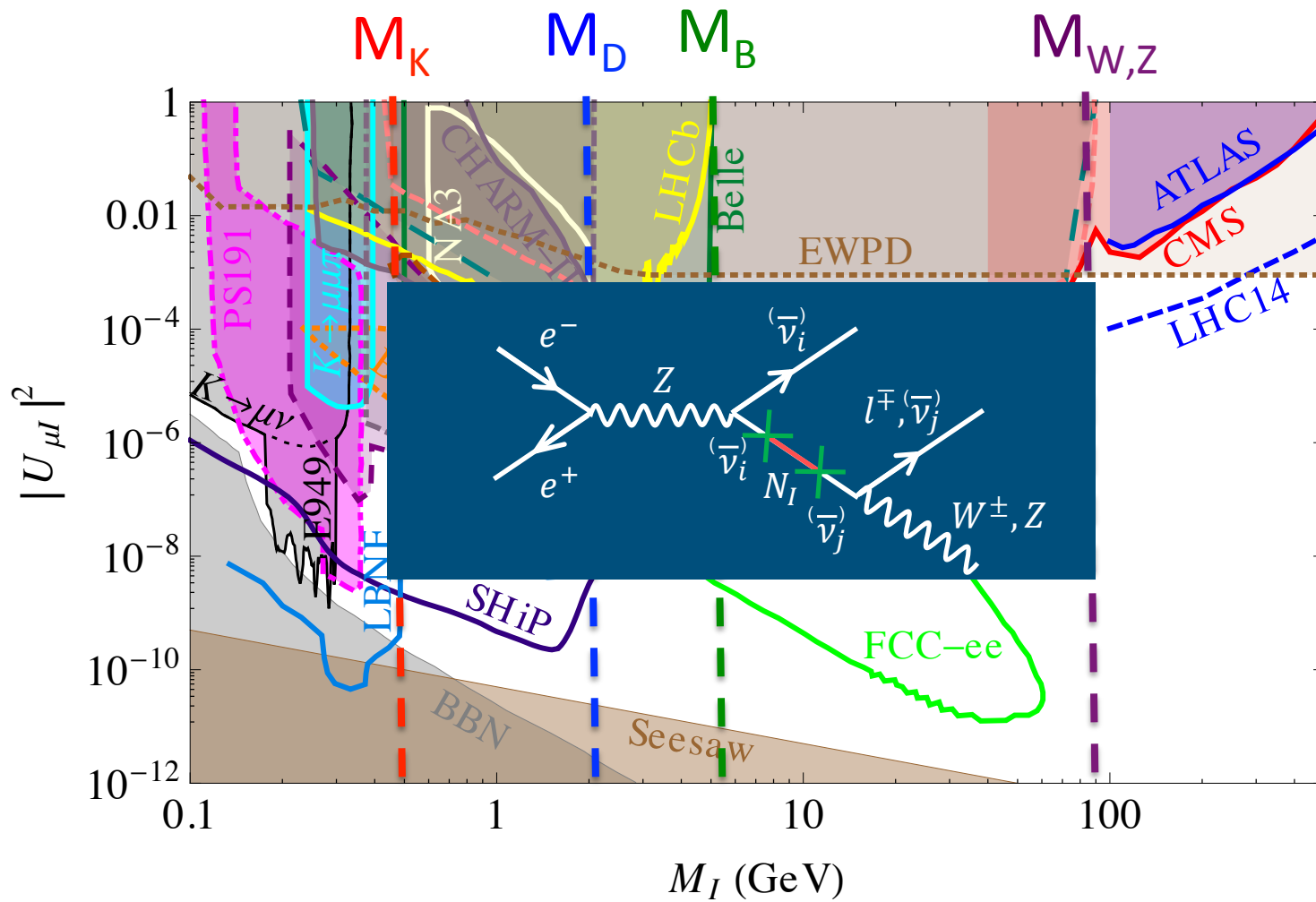


From 1504.04855

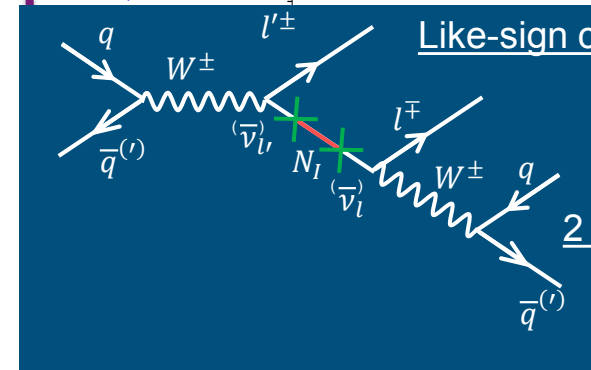
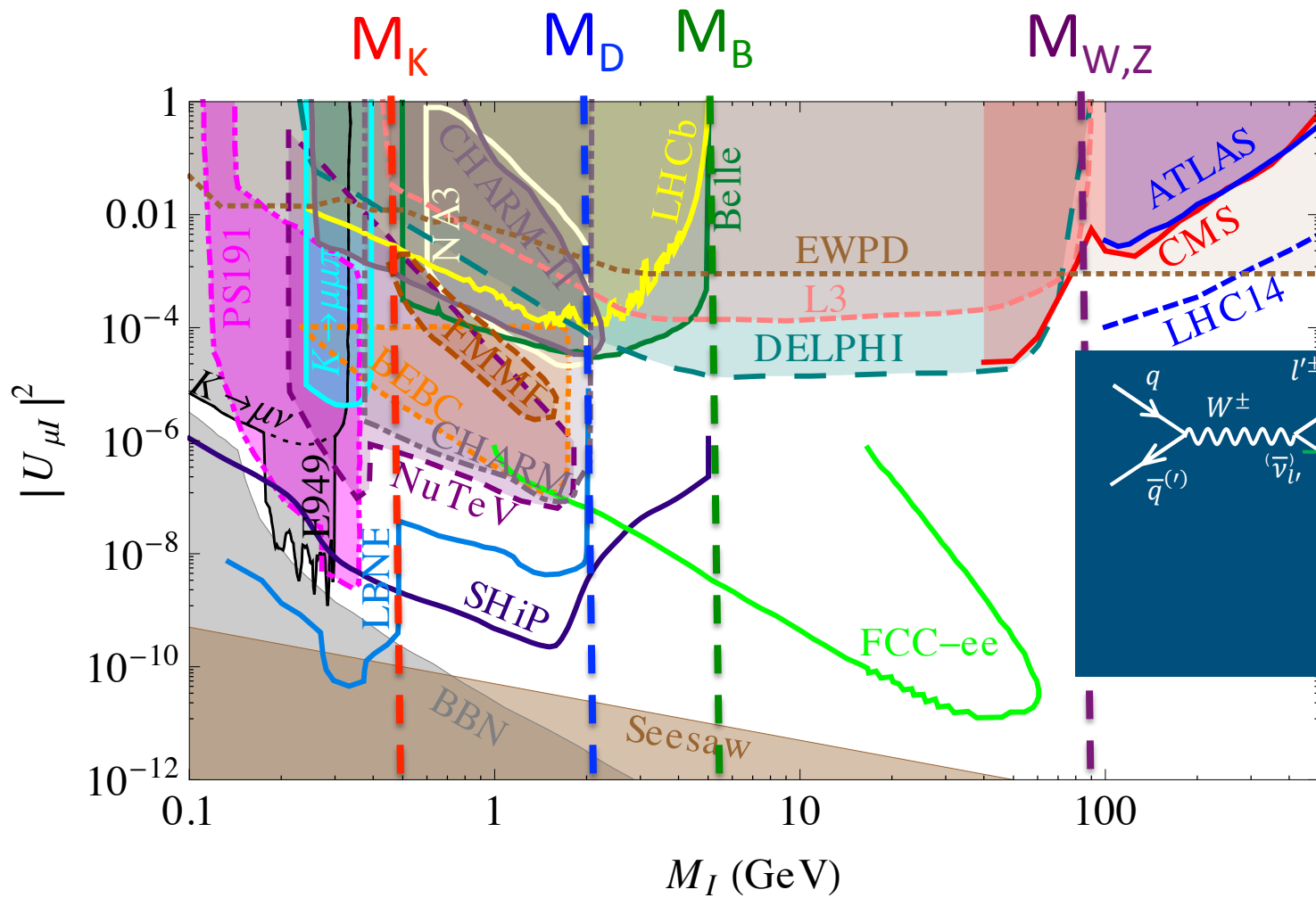
# Searching for heavy Majorana neutrinos



# Searching for heavy Majorana neutrinos

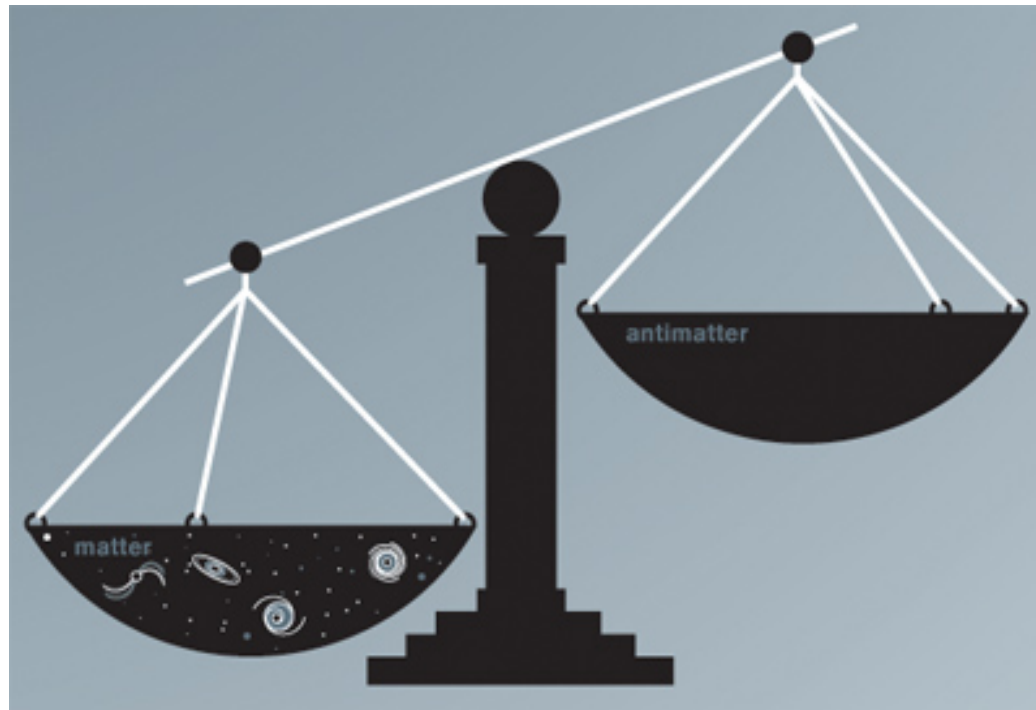


# Searching for heavy Majorana neutrinos



# Matter-antimatter asymmetry

The Universe seems to be made of matter



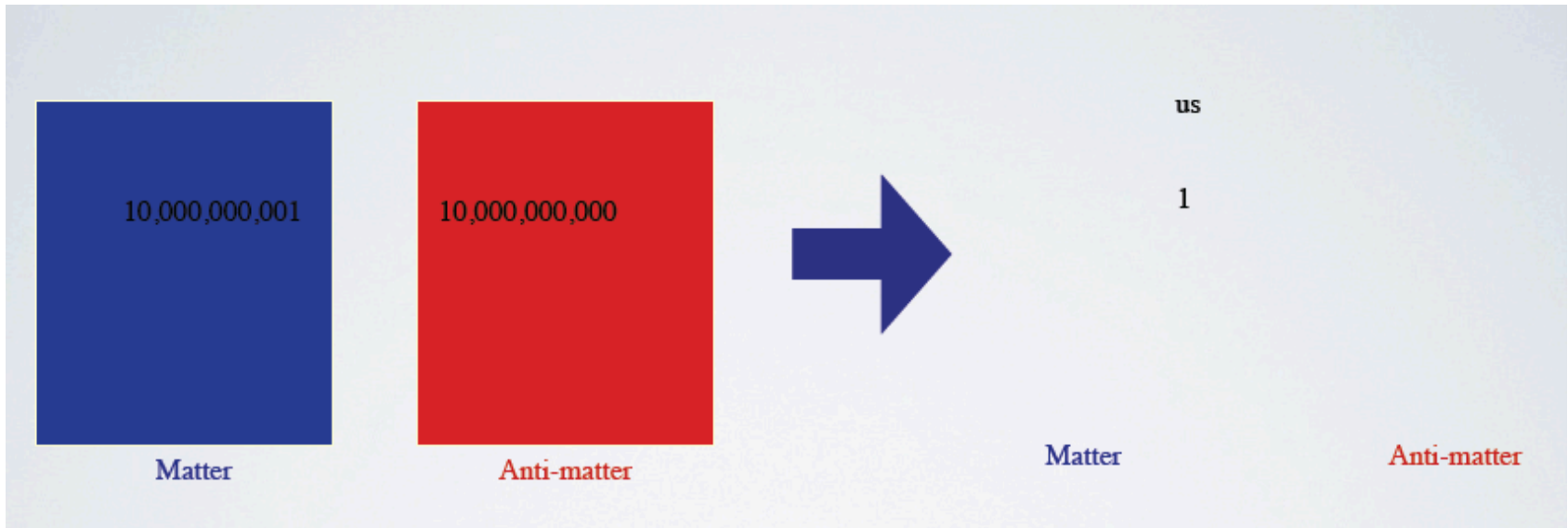
After Planck

$$\eta \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = 6.09(6) \times 10^{-10}$$



# Matter-antimatter asymmetry

The Universe seems to be made of matter



After Planck

$$\eta \equiv \frac{n_B - n_{\bar{B}}}{n_\gamma} = 6.09(6) \times 10^{-10}$$

# Matter-antimatter asymmetry

Can it arise from a symmetric initial condition with same matter & antimatter ?

## Sakharov's necessary conditions for baryogenesis

- ✓ Baryon number violation (B+L violated in the Standard Model)
- ✓ C and CP violation (both violated in the SM)
- ✓ Deviation from thermal equilibrium (at least once: electroweak phase transition)

It does not seem to work in the SM with massless neutrinos ...

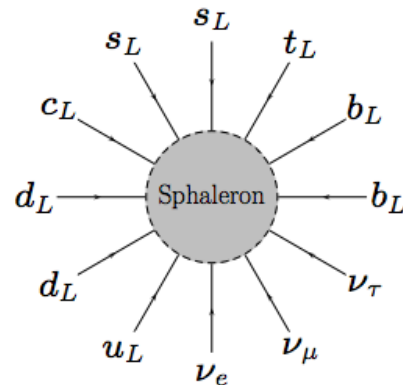
CP violation in quark sector far too small, EW phase transition too weak...

Can we do it with leptons ?

# Baryon number violation

In the SM **there is violation of B+L, preserve B-L**

These processes are strongly suppressed at  $T < T_{EW}$  and in equilibrium at  $T > T_{EW}$

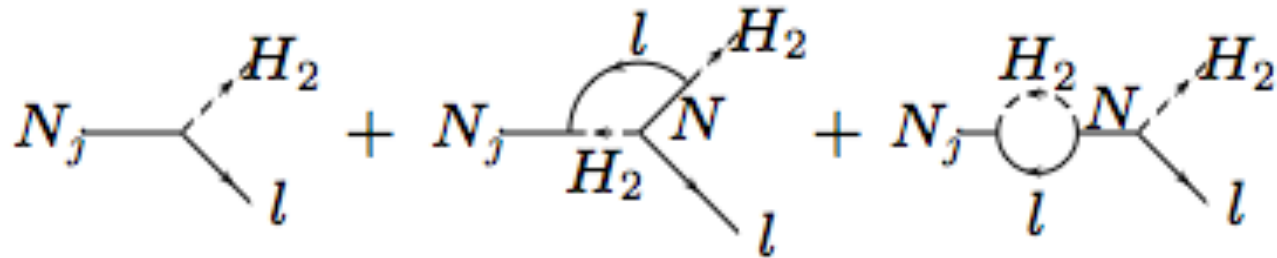


If a net L is generated above  $T_{EW}$  sphalerons produce B

# Leptogenesis: seesaw model $M_N \gg v$

New sources of CP violation and L violation in the neutrino sector can induce CP asymmetries in decays of heavy Majorana  $\nu$

Fukuyita, Yanagida



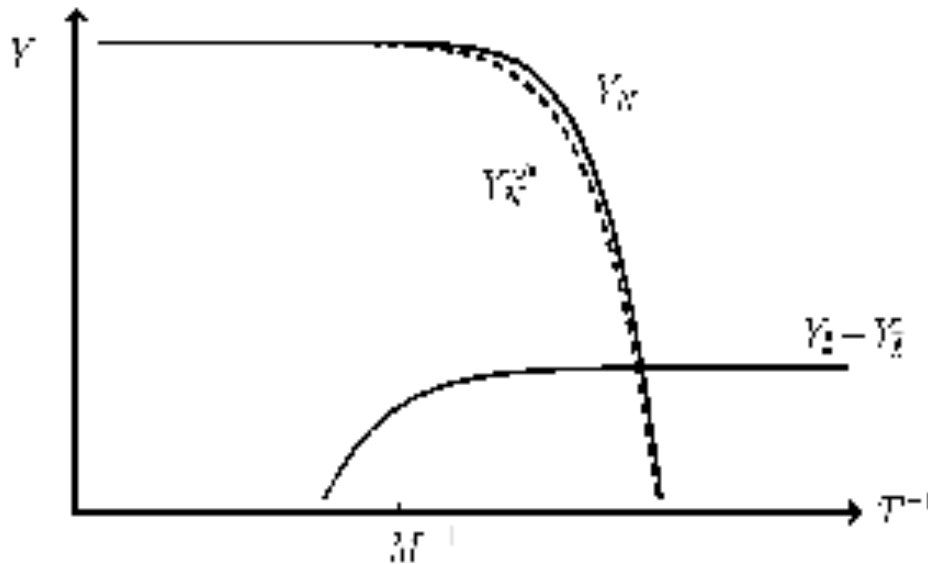
$$\epsilon_1 = \frac{\Gamma(N \rightarrow \Phi l) - \Gamma(N \rightarrow \Phi \bar{l})}{\Gamma(N \rightarrow \Phi l) + \Gamma(N \rightarrow \Phi \bar{l})}$$

A significant asymmetry typically requires  $M_N > 10^7$  GeV

# Out-of-equilibrium

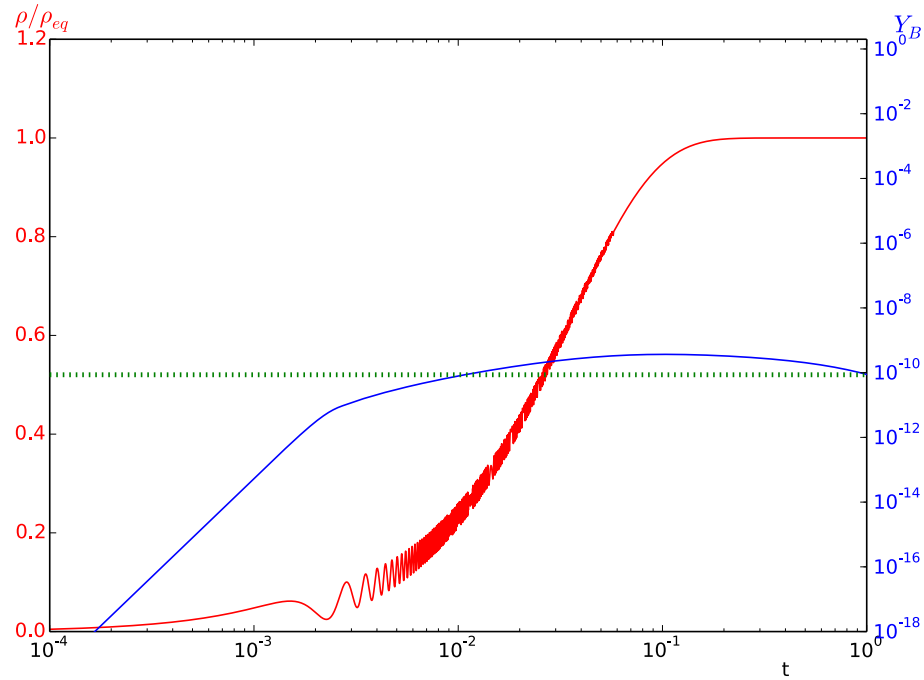
The Majorana neutrinos decay out-of-equilibrium

$$\Gamma_N < \text{expansion rate} \Rightarrow N_N > N_N^{\text{thermal}}$$



# Leptogenesis: seesaw model $M_N < v$

Leptogenesis takes place during the production of heavy states

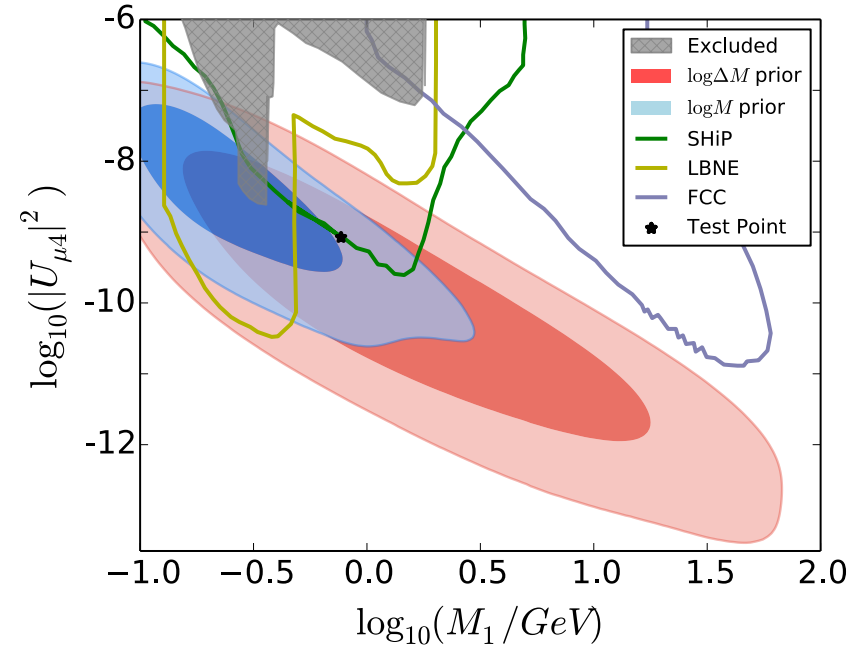
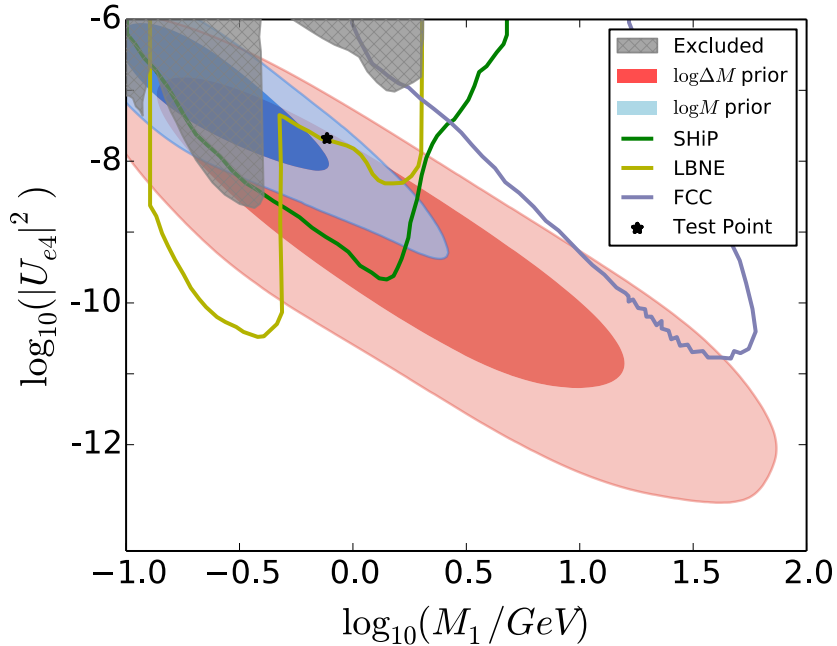


Courtesy of M. Kevic

$T_{EW}$

# Leptogenesis: seesaw model $M_N < v$

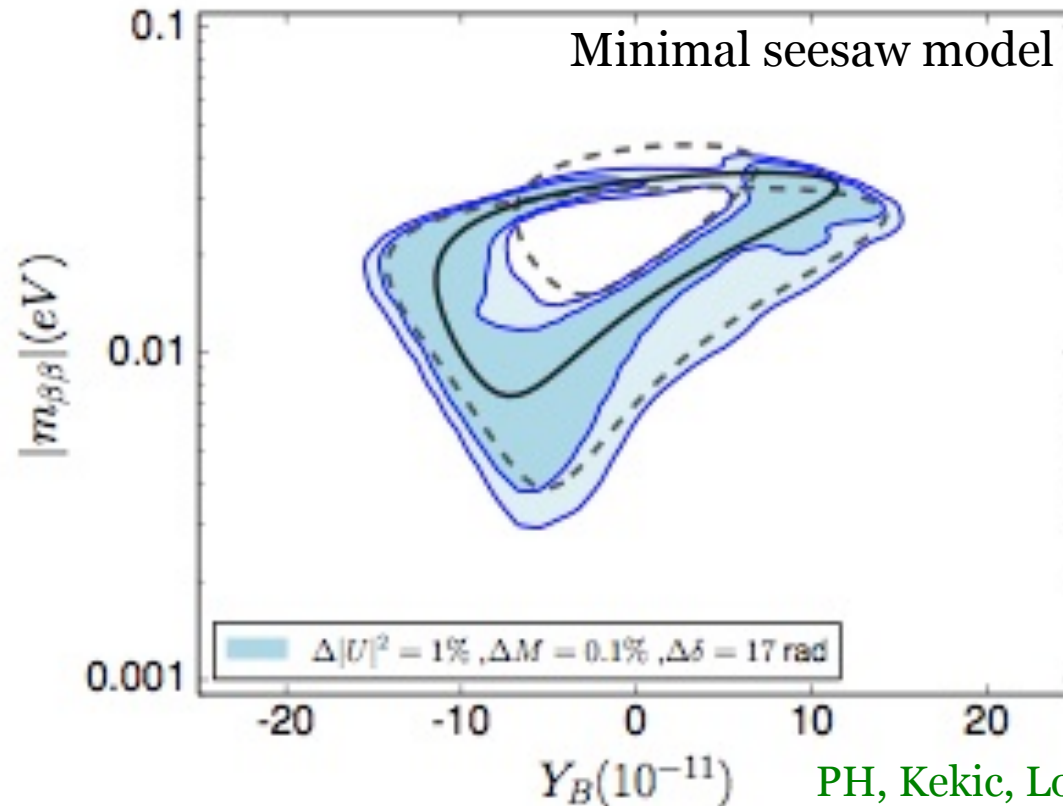
Posterior probabilities mixings vs mass for successful baryon asymmetry



PH, Kekic, Lopez-Pavon, Racker, Salvado

Future searches have a change to discover these states !

One example: sterile neutrinos are discovered in D decays at SHIP and their masses and mixings to e measured with some accuracy + CP phase  $\delta$  is measured in neutrino oscillations



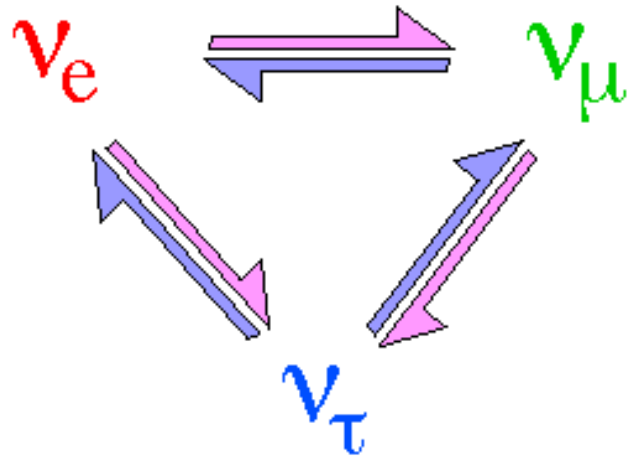
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Neutrinoless double beta decay could be key to predict the baryon asymmetry

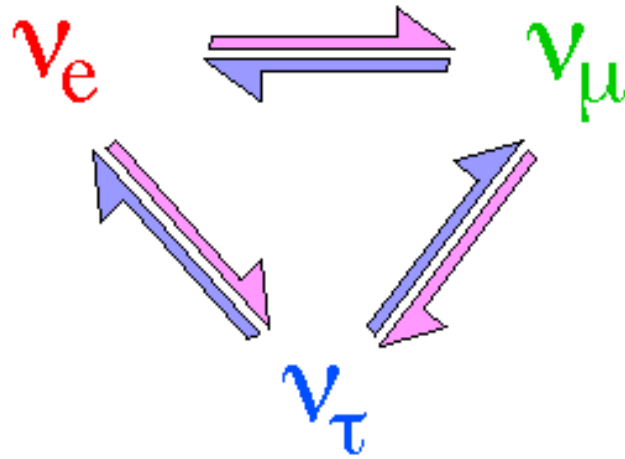


- The results of many beautiful experiments have demonstrated that  $\nu$  are (for the time-being) the less standard of the SM particles
- Many fundamental questions remain to be answered however:  
Majorana nature of neutrinos and scale of new physics? CP violation in the lepton sector? Source of the matter-antimatter asymmetry ?  
Lepton vs quark flavour ?
- A rich experimental programme lies ahead where fundamental physics discoveries are very likely (almost warranted) ...

These elusive pieces of reality have brought many surprises, maybe they will continue with their tradition...



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*ENJOY YOUR SUMMER-STUDENT SUMMER!*