APDF Evolution Library

[V. Bertone, et al., Comput. Phys. Commun. 185, 1647 (2014)]

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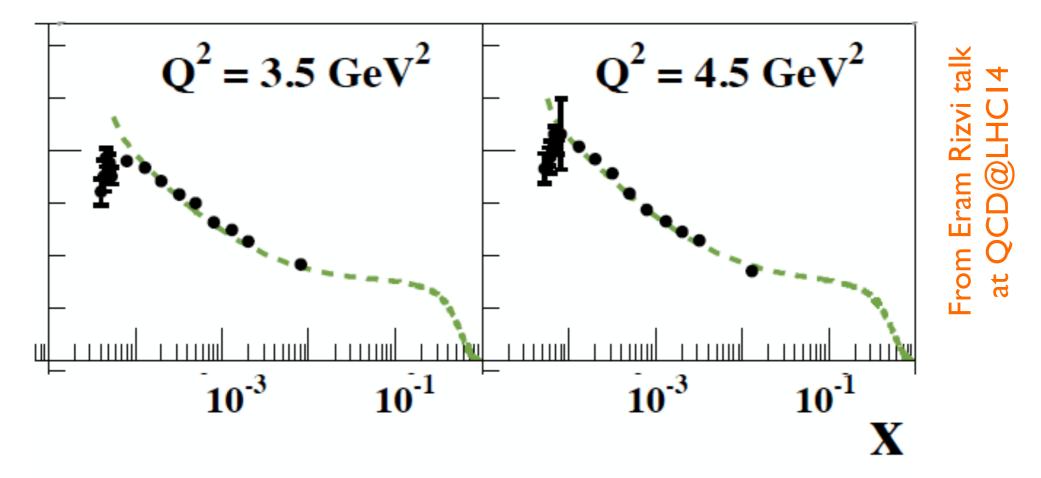
APFEL in a Nutshell

• APFEL is a **public** library for the computation of collinear PDF evolution and DIS structure functions:

- *up* to NNLO in QCD combined to QED corrections up to NLO.
- *FFN* and VFN schemes.
- \checkmark Pole and $\overline{\mathrm{MS}}$ heavy-quark masses.
- fast computation of DIS NC and CC observables in different mass schemes (ZM-VFNS, FFNS and FONLL).
- ✓ Interfaces to FORTRAN, C/C++ and Python.
- Web interface available on http://apfel.mi.infn.it.
- *available from http://apfel.hepforge.org/*.
- Interfaced to **xFitter** (see Ringaile's talk) and **Alpos**.
- ✓ Used for the next generation of the **NNPDF** fits (including FFs).

Small-x Resummation

Tension between fixed-order predictions and data in the small-*x* region reached by HERA:



A similar effect was observed some time ago in the NNPDF framework by F. Caola *et al.* [arXiv:1007.5405].

Suggestion of the need for **small-***x* **resummation**.

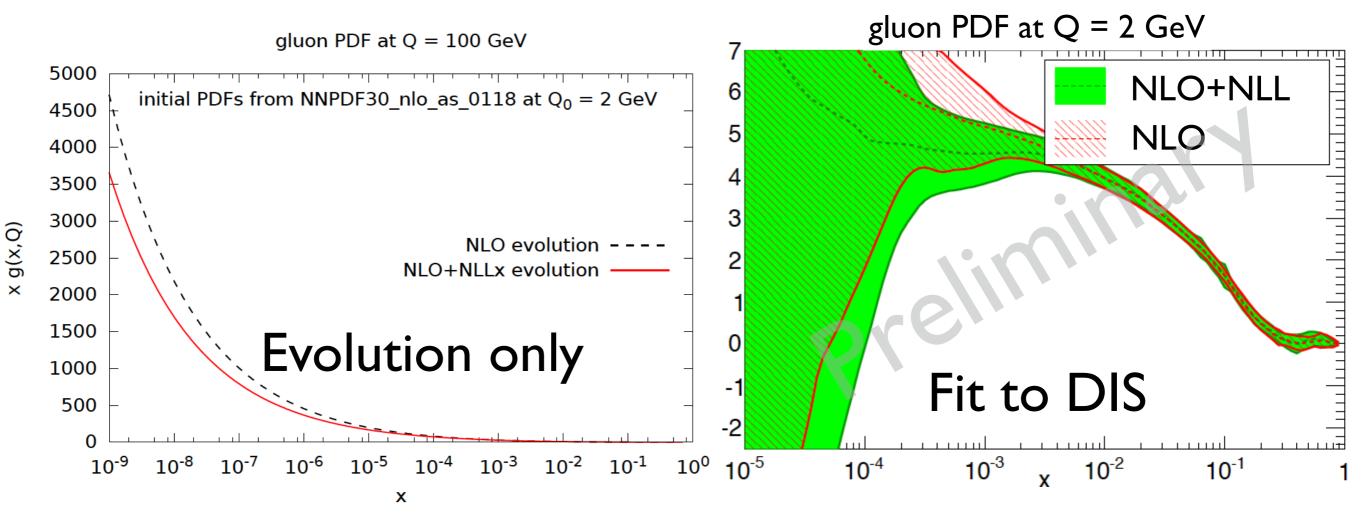
Small-*x***Resummation**

The **HELL** code [arXiv:1607.02153] has been interfaced to APFEL:

- Join the ABF formalism (e.g. see [hep-ph/9501231]).
- **S**mall-*x* **resummed splitting functions** up to **NLL** accuracy,
- **Small-***x* resummed DIS coefficient functions up to NLL:
 - massless (already available from APFEL),
 - \checkmark massive (since very recently \Rightarrow not interfaced to APFEL yet)
- *Kesummed matching conditions (not implemented in HELL, to come).*

In NNPDF we have attempted DIS-only PDF fits with small-x resummation correction at NLL in the evolution and in the ZM sector of the DIS structure functions obtaining **encouraging results**.

Small-*x* **Resummation**



Resummed evolution leads to a **suppression** of the **gluon PDF** at small values of *x* as compared to fixed order.

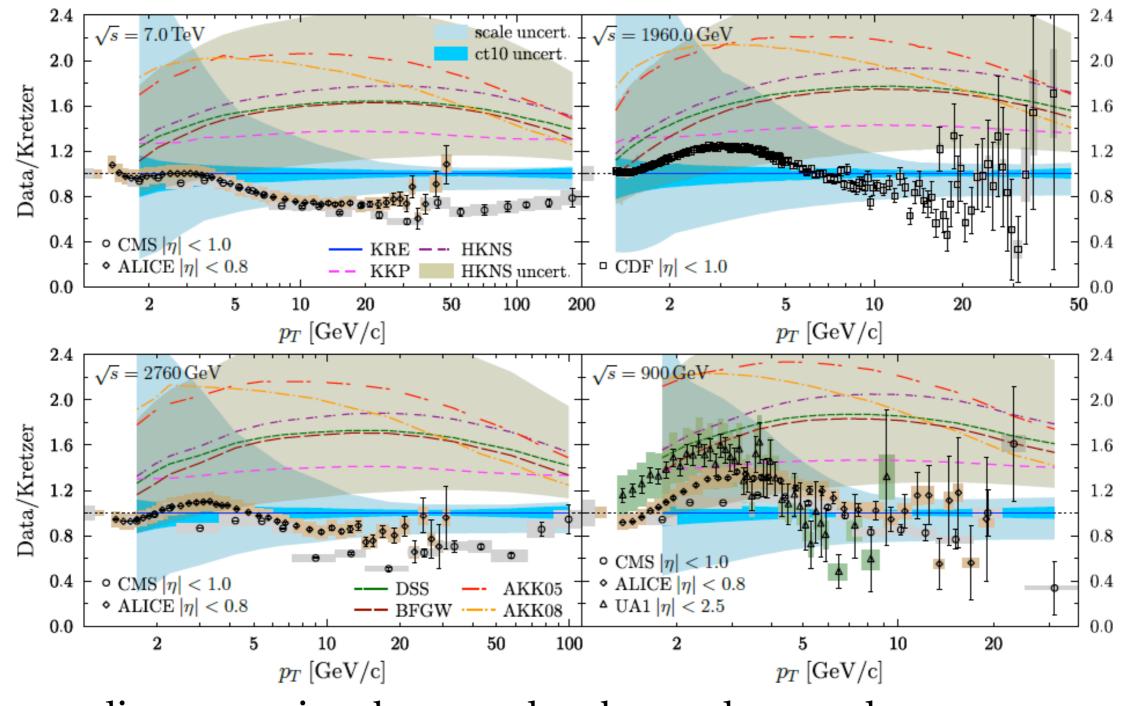
Compensation when also resummed **coefficient functions** are introduced ⇒ effect on the small-*x* gluon PDF at the level of 1- σ .

• Other PDFs mostly unchanged.

- A faithful determination of fragmentation functions (FFs) is extremely important to study the universality of the QCD factorisation theorems.
- The inclusive hadron measurements at the LHC, sensibly extending the previous kinematical coverage, are particularly useful for studying the FFs.
- Moreover, a good knowledge of FFs is functional to the determination of the **polarised PDFs**.
- The **spread between the different FFs** present on the market is currently very large.
- In addition, none of the existing FF sets describes the recent LHC and Tevatron experimental data.



Inclusive charged-hadron spectrum:

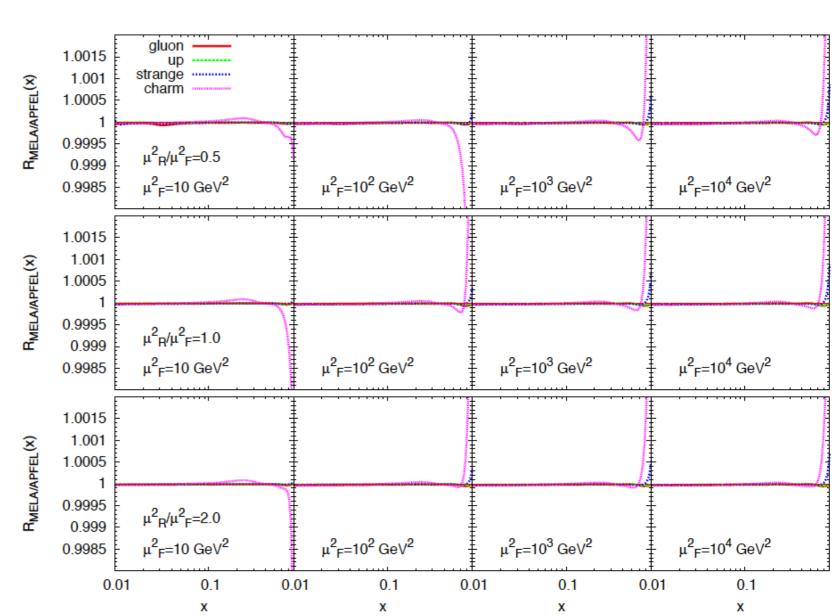


• Large discrepancies that need to be understood.

'Enterria et al. [arXiv:1311.1415]

• APFEL implements the **time-like evolution**:

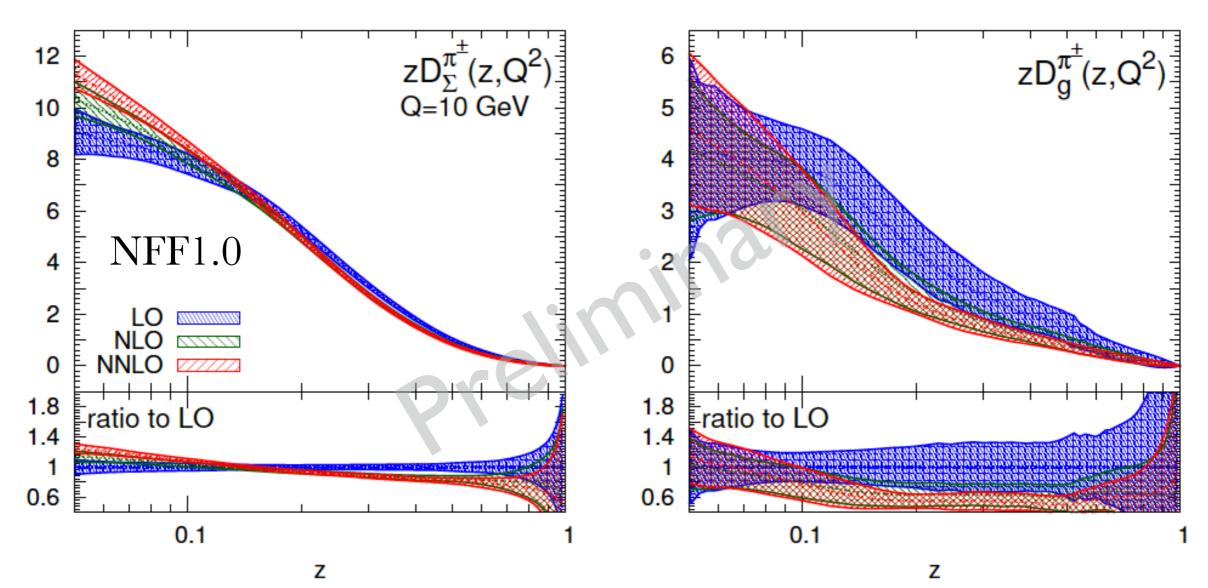
- up to NLO in the VFNS,
- up to NNLO in the FFNS (NNLO matching conditions unknown).
- Careful **benchmark** of the evolution against the MELA code [arXiv:1501.00494].
- Excellent agreement at all perturbative orders.



Single-inclusive e⁺e⁻ annihilation (SIA) structure functions are also implemented in APFEL up to NNLO in QCD:

• partial benchmark against DSS code (thanks to D.P. Anderle).

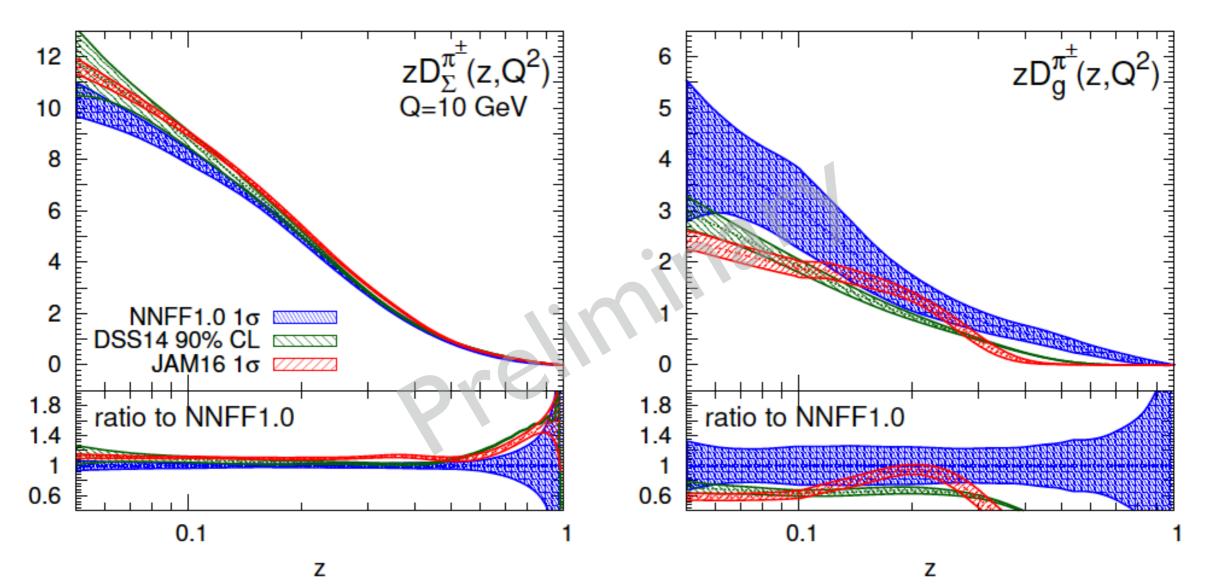
APFEL can now be used to **determine FFs** from SIA data.



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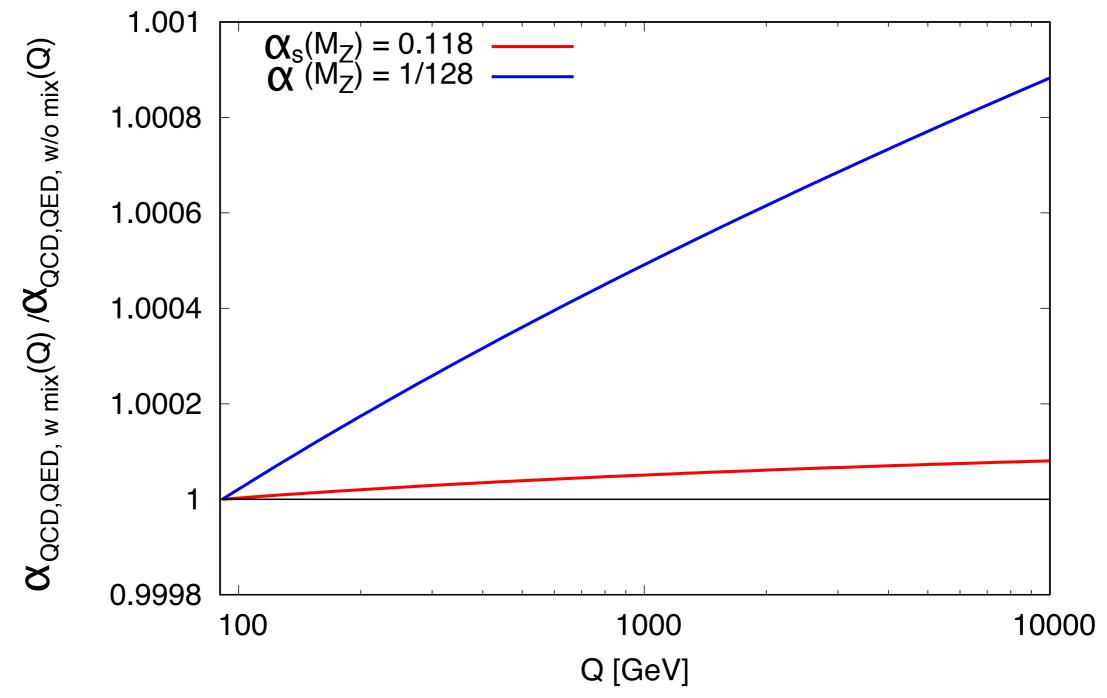


In order to implement the full NLO QCD+QED corrections in the DGLAP evolution two main steps are required:

- 1. Implementing the $O(\alpha_s^2 \alpha)$, $O(\alpha^3)$, $O(\alpha^2 \alpha_s)$ corrections to the β -functions:
 - running of α_s and α is coupled \Rightarrow solve of a coupled ODE,
 - Numerical tests have shown that such terms lead to differences of $O(10^{-4})$ for α_s and $O(10^{-3})$ for $\alpha \Rightarrow$ **unneeded complication**.

NLO QCD+QED Corrections Coupling Evolution

running of the couplings, $N_F = 5$



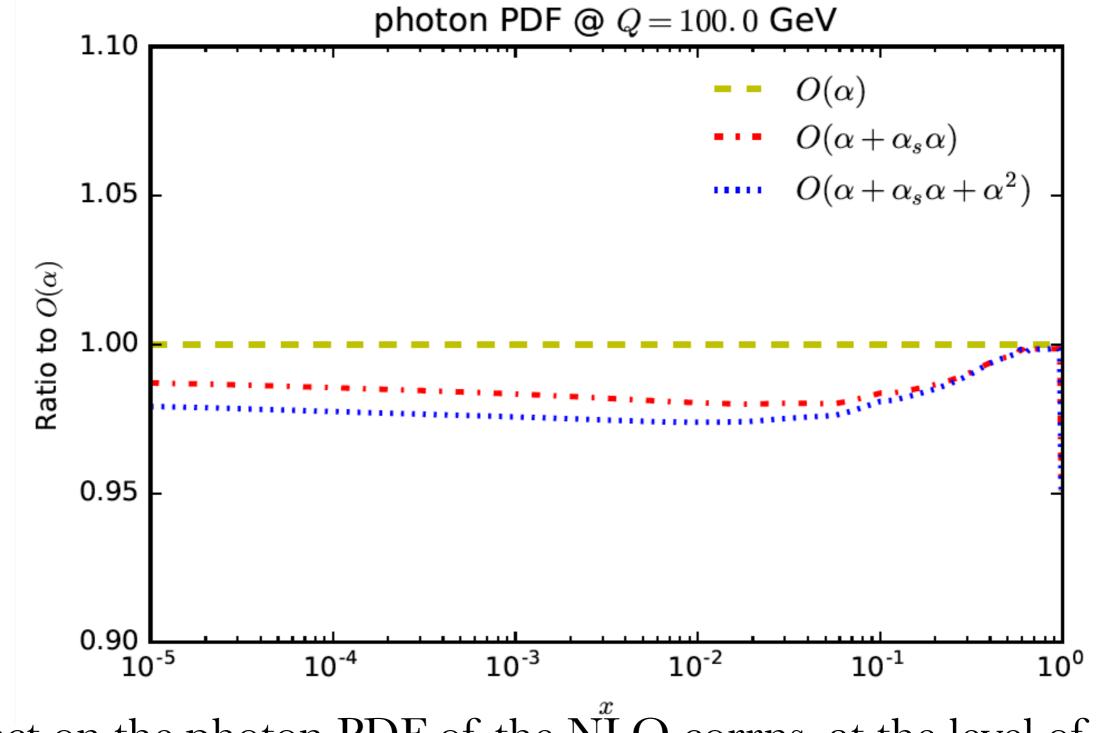
 \mathbf{I} Mixed terms in the β -functions lead to negligible effects.

NLO QCD+QED Corrections *Evolution*

In order to implement the full NLO QCD+QED corrections in the DGLAP evolution two main steps are required:

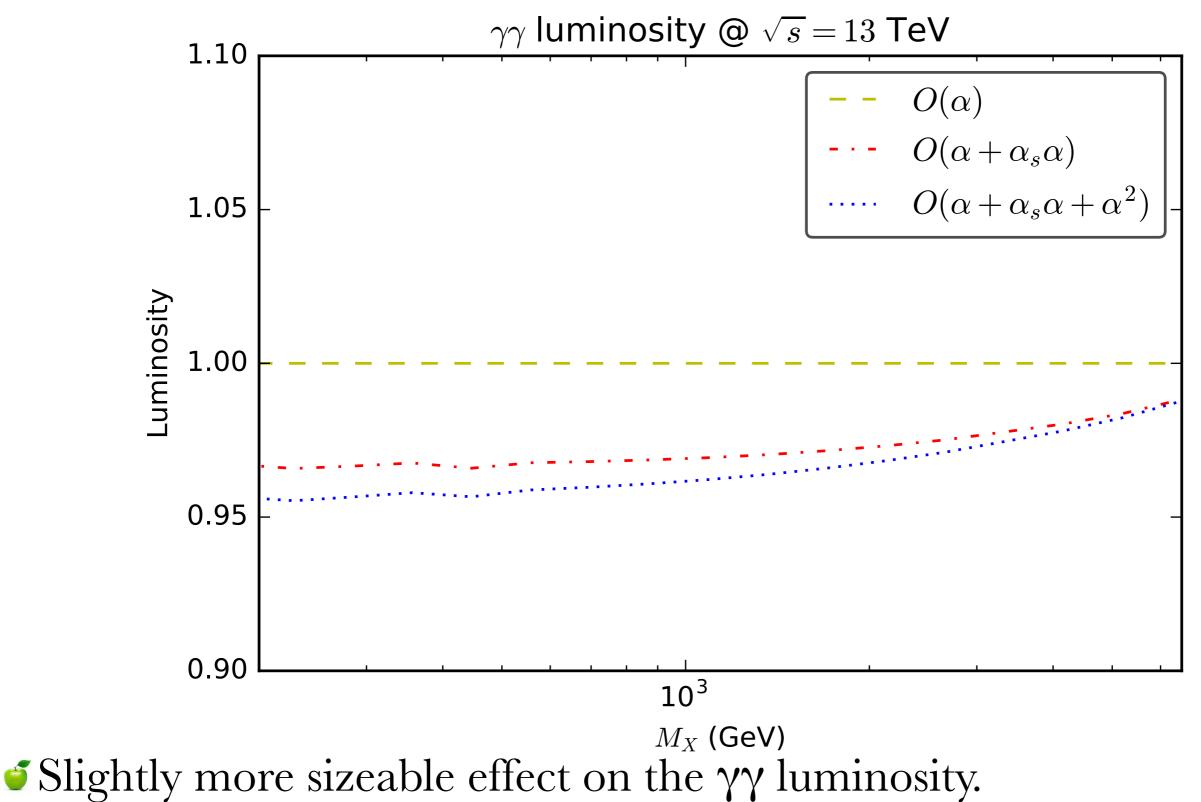
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 - running of α_s and α is coupled \Rightarrow solve of a coupled ODE,
 - Numerical tests have shown that such terms lead to differences of $O(10^{-4})$ for α_s and $O(10^{-3})$ for $\alpha \Rightarrow$ **unneeded complication**.
- 2. Implementing the $O(\alpha_s \alpha)$ and the $O(\alpha^2)$ corrections to the DGLAP **splitting functions** on top of the $O(\alpha)$ ones:
 - complication of the flavour structure due to the presence of terms promotional to e_q^2 and e_q^4 that break the isospin symmetry,
 - need for a more optimal evolution basis as compared to pure QCD.

NLO QCD+QED Corrections DGLAP Evolution



 \checkmark Effect on the photon PDF of the NLO corrns. at the level of 1-2%.

NLO QCD+QED Corrections DGLAP Evolution



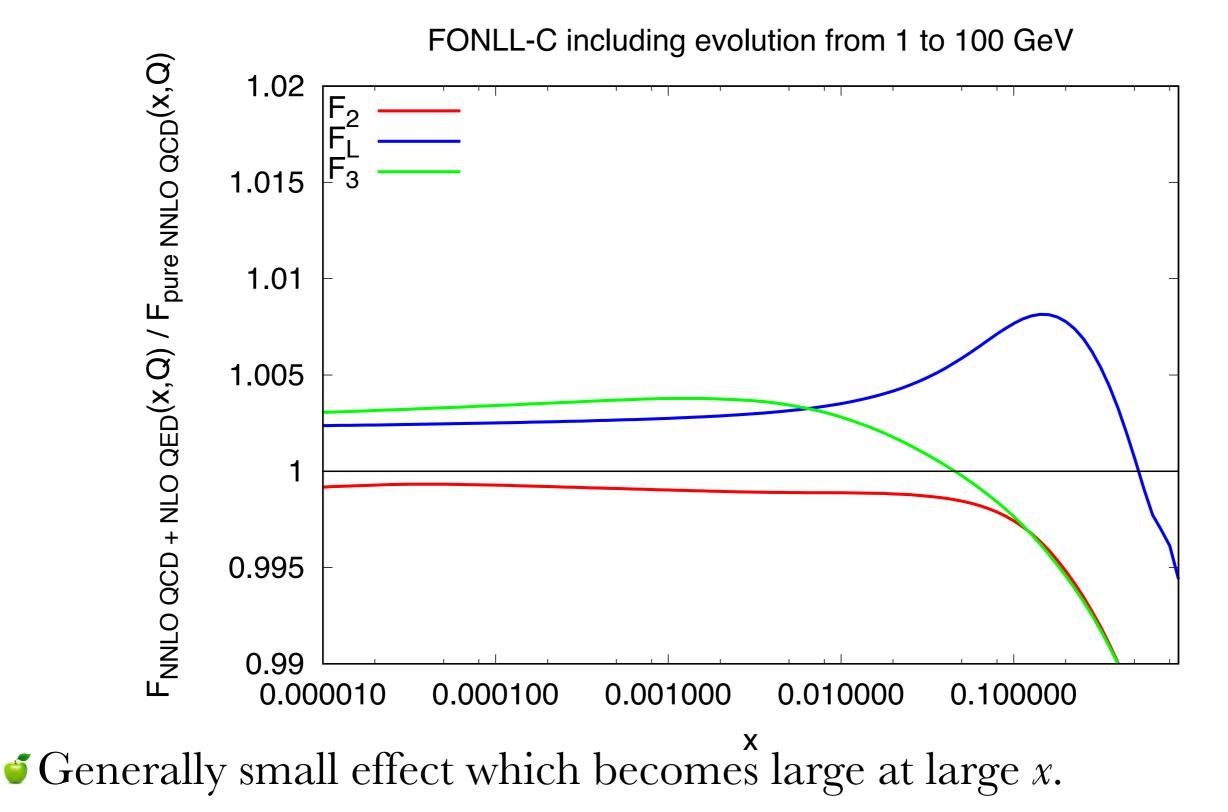
NLO QCD+QED Corrections DIS Structure Functions

While at LO in QED no corrections to the DIS structure functions are required (γ*q → q itself is the LO), at NLO in QED O(α) corrections need to be taken into account:

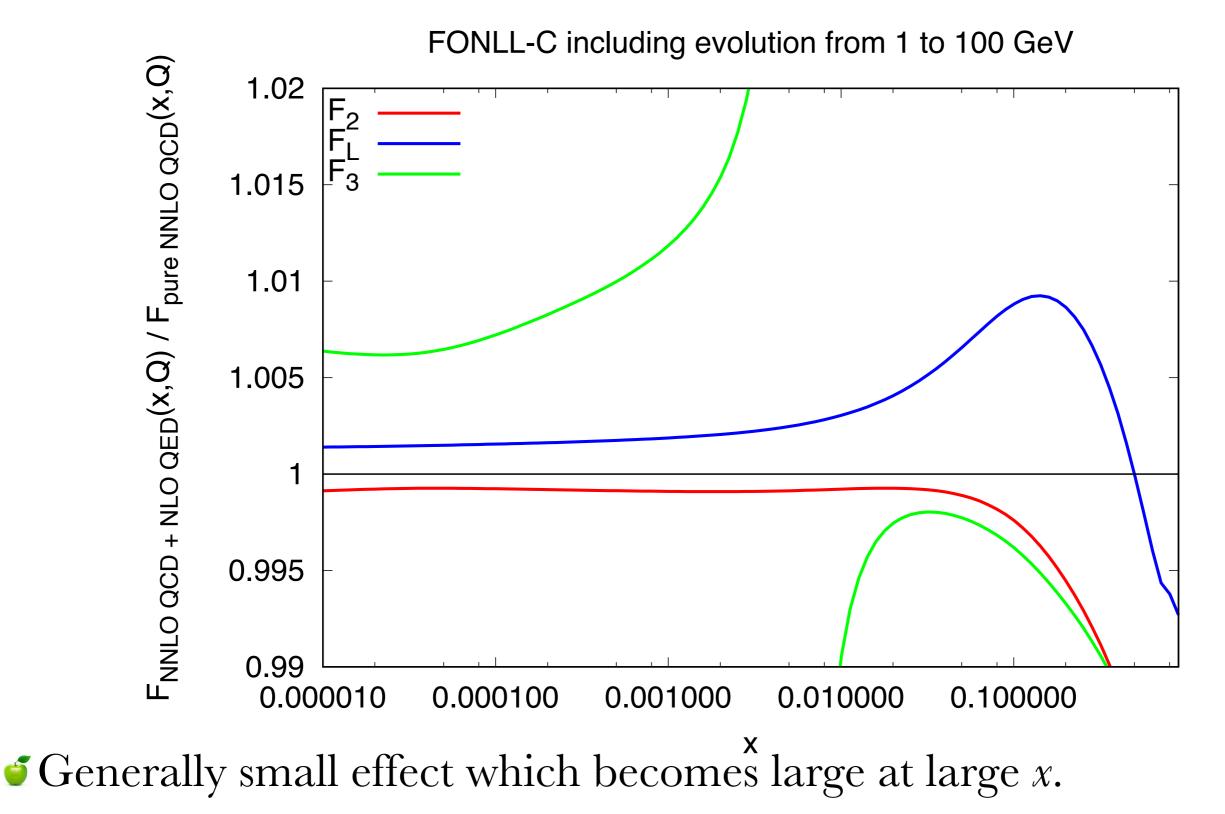
• **new diagrams**: $\gamma^* \gamma \rightarrow q \overline{q}$ and $\gamma^* q \rightarrow q \gamma$,

- easily derivable from the corresponding QCD diagrams.
- The additional diagrams offer a **direct handle** on the photon PDF in DIS observables:
 - at LO in QED the photon PDF was entirely driven by the evolution.
- Small contribution proportional to $\alpha\gamma \sim O(\alpha^2)$ but can be relevant in some kinematic regions:
 - typically at large x and large Q^2 .

NLO QCD+QED Corrections DIS Structure Functions (NC)



NLO QCD+QED Corrections DIS Structure Functions (CC)





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Welcome to APFEL online cluster!

This web-application is a tool designed for High Energy Physics by providing a simple and intuitive interface to plot and compute the most common observables with Parton Distribution Functions (PDFs).

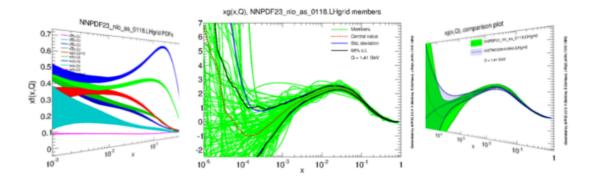
To begin to produce on-line plots, please register and login!

The APFEL library

APFEL, a PDF evolution library, is a computer library specialized in the solution of DGLAP evolution equations up to NNLO in QCD and to LO in QED, both with Pole and $\overline{\rm MS}$ masses. With APFEL you can replace the evolution of LHAPDF sets and check the impact on the choice of evolution parameters. APFEL also computes deep-inelastic scattering processes using multiple schemes.

If you use the APFEL library or the online cluster in a scientific publication, please cite: V. Bertone, S. Carrazza and J. Rojo, "APFEL: A PDF Evolution Library with QED corrections", Comput. Phys. Commun. 185, 1647 (2014), arXiv:1310.1394.

S. Carrazza et al., "APFEL Web: a web-based application for the graphical visualization of parton distribution functions", J. Phys. G: Nucl. Part. Phys. 42 057001, arXiv:1410.5456. Labtalk.



Web developers: D. Palazzo, S. Carrazza, A. Ferrara APFEL developers: V. Bertone, S. Carrazza, J. Rojo. (Contact)

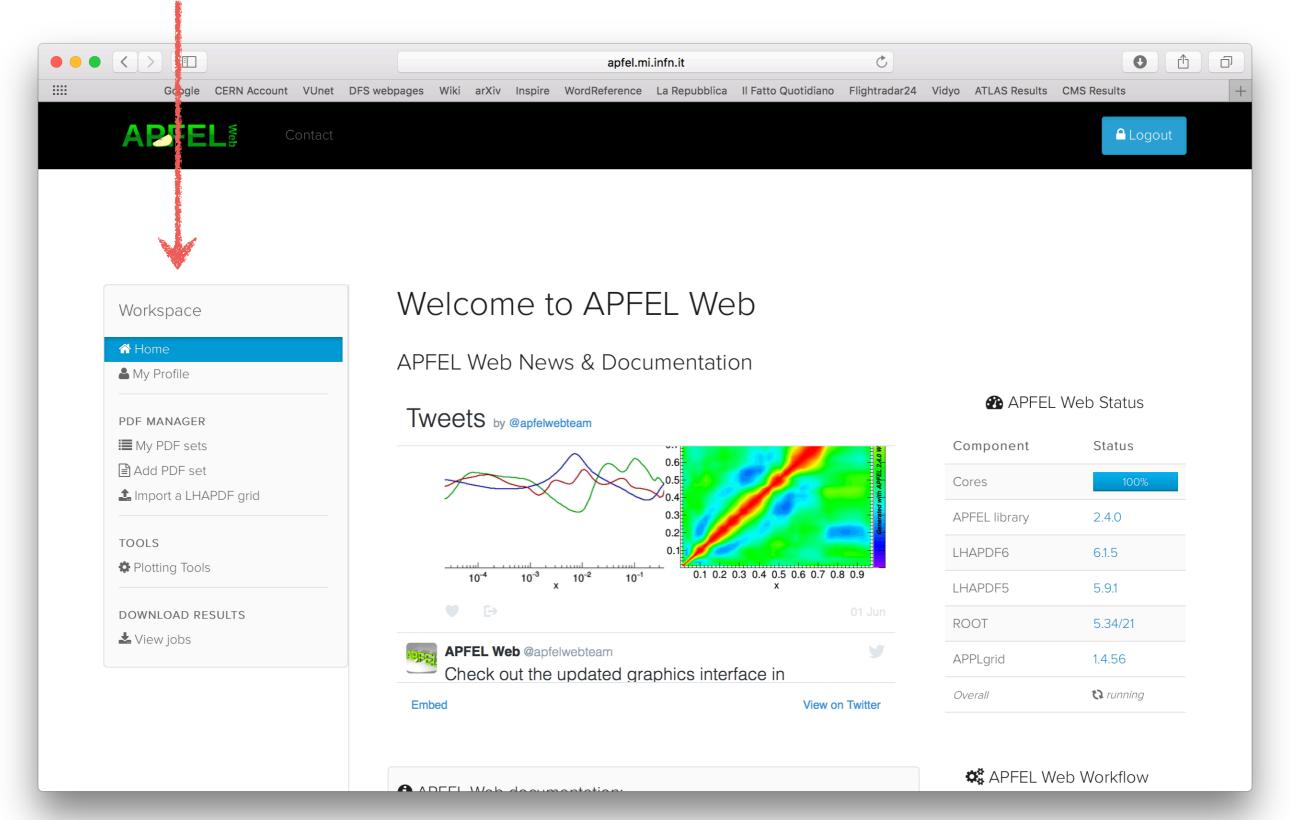


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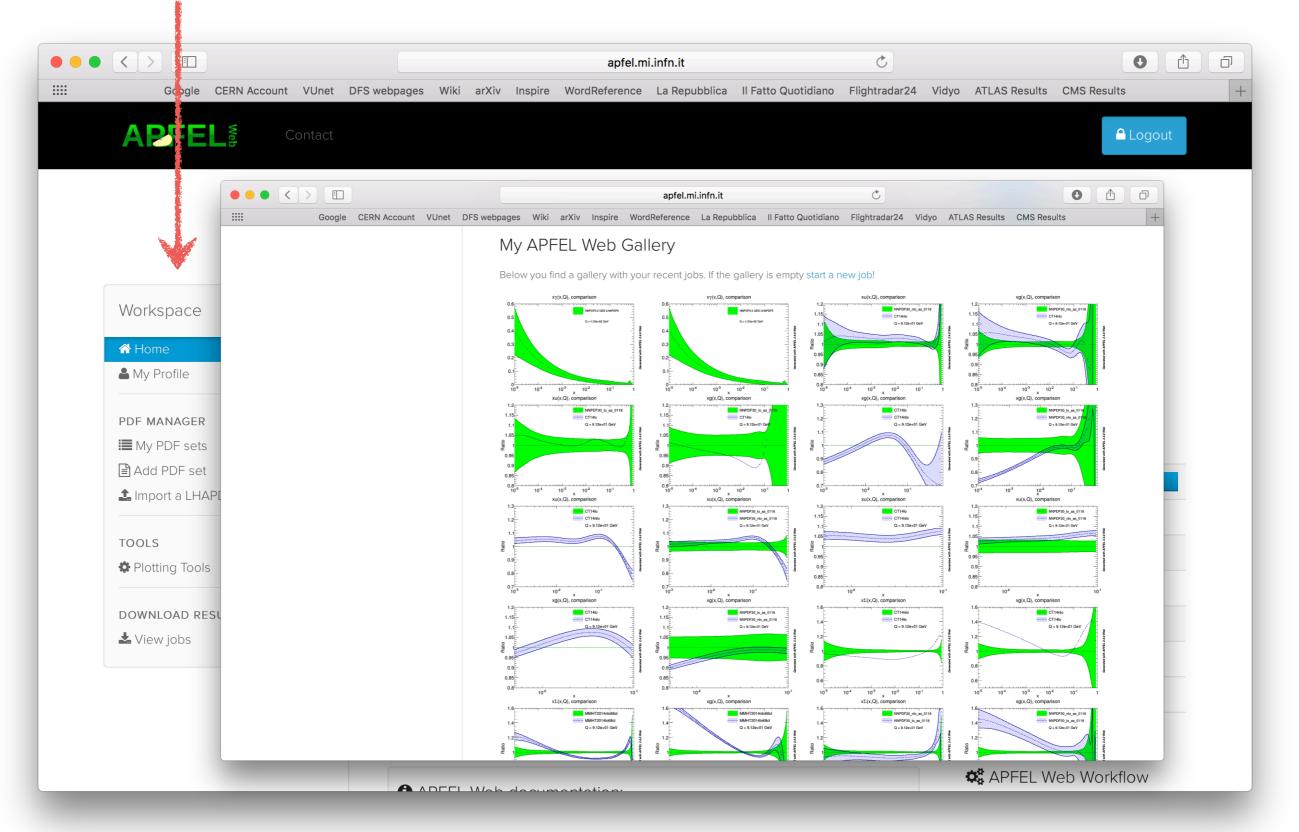


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Workspace





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		m_b , bottom mass in (GeV):	4,18		
		m_t , top mass in (GeV):	173,07		
		$lpha_S$ ($Q_{ m ref}$):	0,118		
		$Q_{ m ref}$ (GeV):	91,2		
		$lpha$ (${\cal Q}_{ m ref}$):	0,007496252		
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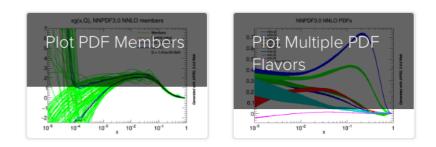
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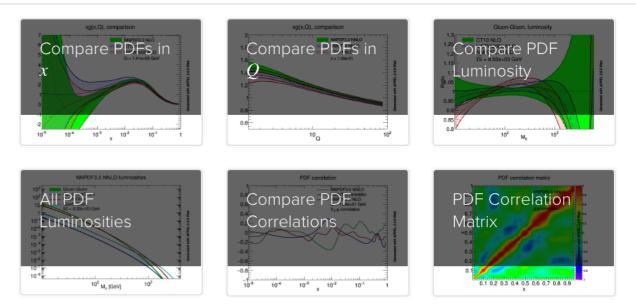
Some jobs, like PDF luminosities, require some time to be finalized. Check the job status at View jobs page.

The plotting tools can be used for both the LHAPDF libraries: LHAPDF5 and LHAPDF6.

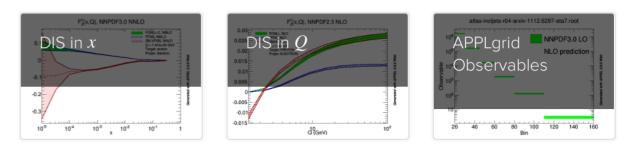
Tools for PDF basic plotting



Tools for PDF analysis & comparisons



Tools for theoretical predictions from PDFs



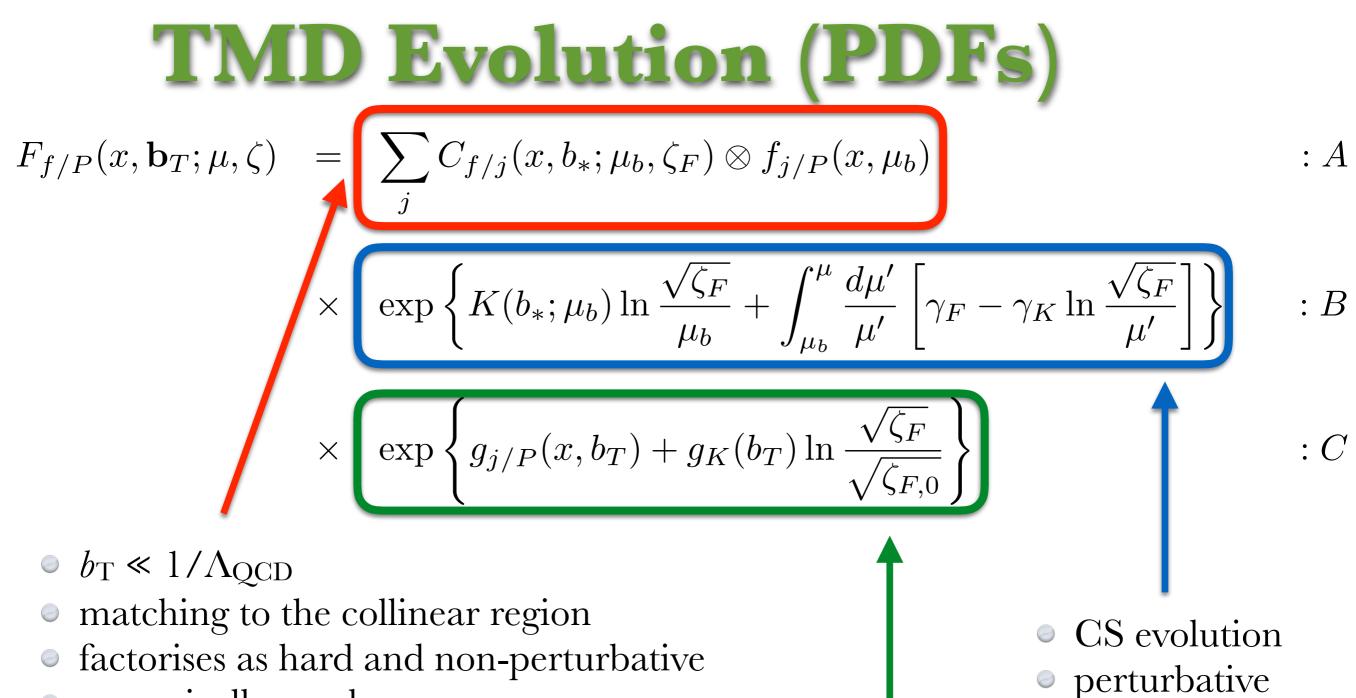
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TMD Evolution (PDFs)

$$F_{f/P}(x, \mathbf{b}_T; \mu, \zeta) = \sum_j C_{f/j}(x, b_*; \mu_b, \zeta_F) \otimes f_{j/P}(x, \mu_b) : A$$

$$\times \exp\left\{K(b_*;\mu_b)\ln\frac{\sqrt{\zeta_F}}{\mu_b} + \int_{\mu_b}^{\mu}\frac{d\mu'}{\mu'}\left[\gamma_F - \gamma_K\ln\frac{\sqrt{\zeta_F}}{\mu'}\right]\right\} : B$$

$$\times \exp\left\{g_{j/P}(x, b_T) + g_K(b_T) \ln \frac{\sqrt{\zeta_F}}{\sqrt{\zeta_{F,0}}}\right\} : C$$



- numerically cumbersome
- precompute using the APFEL technology

- matching between the small and large $b_{\rm T}$
- non perturbative
- parametrised and fitted to data

TMDs in SIDIS

In SIDIS, what enters the computation of the cross sections is:

 $\mathcal{L}_{\text{SIDIS}} = \int \frac{d^2 \mathbf{b}_T}{(2\pi)^2} e^{-i\mathbf{q}_T \cdot \mathbf{b}_T} F_{f/P}(x, \mathbf{b}_T; \mu, \zeta_F) D_{H/f}(x, \mathbf{b}_T; \mu, \zeta_D)$

Fourier transformPDFsFFsThe ingredients are:

- ✓ a set of evolved TMD-PDFs,
- ✓ a set of evolved TMD-FFs,
- the Fourier transform of its product.

Complex set of tasks that have to be performed optimally

- APFEL provides the ideal environment for this computation:
 - fast and accurate interpolation techniques,
 - for precomputation of the time consuming bits.



- Recent developments in APFEL:
 - **small-***x* **resummation** in PDF evolution and structure functions,
 - *framework for the determination of FFs,*
 - NLO QED corrections to evolution and structure functions.
 Other Recent Developments
- **Intrinsic-charm** in DIS *a la* FONLL.
- **Folarised DGLAP evolution** up to NNLO.
- Independent factorisation and renormalisation scale variations both in the DIS structure functions and in the evolution.

In the Pipeline

- Implementation of **TMD evolution** and **SIDIS cross sections**,
- Implementation of the **polarised structure functions**,
- **mass corrections** to SIA structure functions.

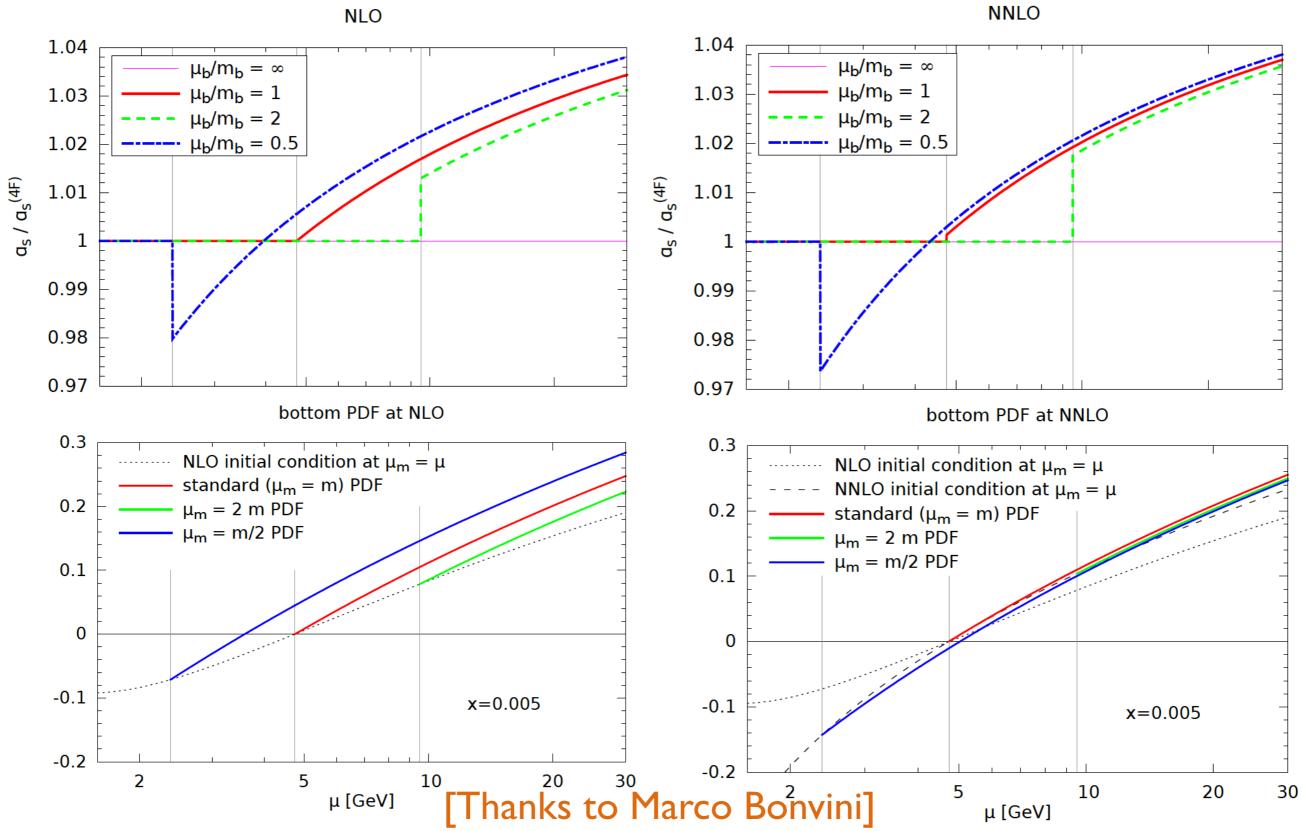
Backup Slides

Displaced Heavy-Quark Thresholds

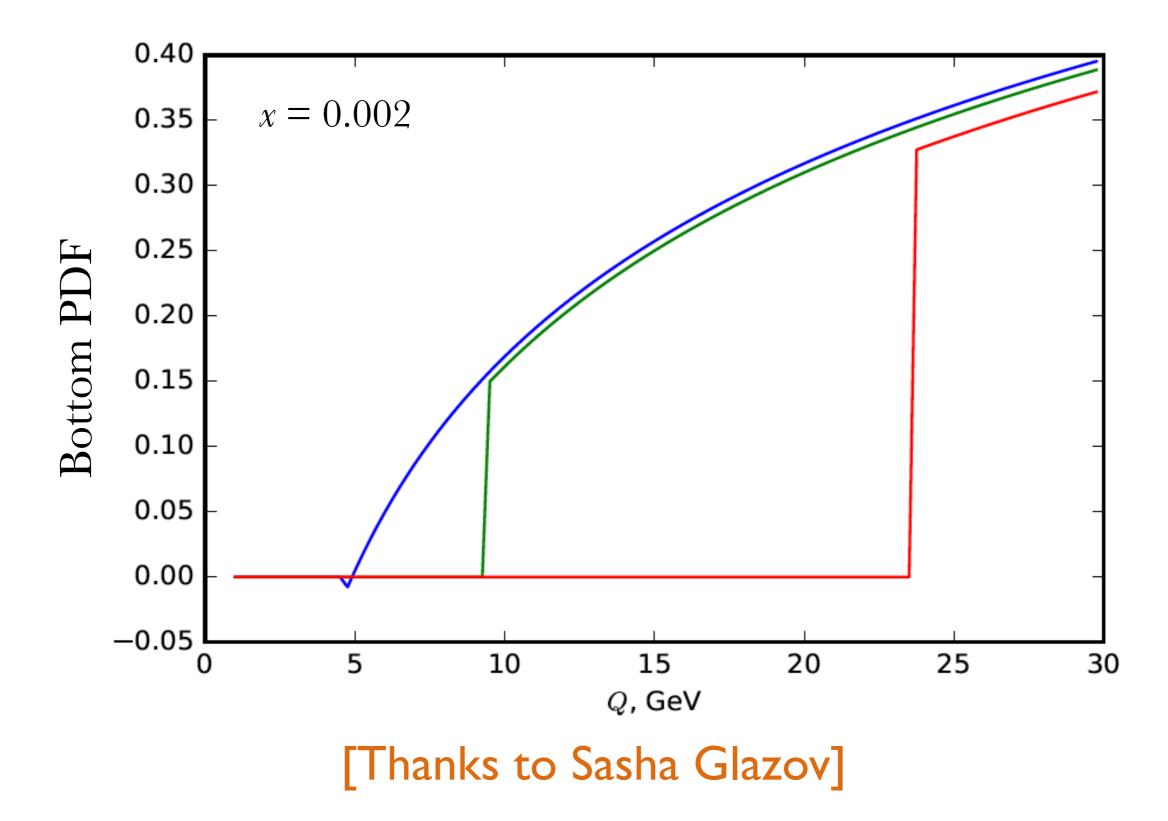
 The implementation of the VFNS evolution both for PDFs and α_s requires matching factorization schemes differing in the number of active flavours:

- the scale at which two consecutive factorization schemes are matched are usually referred to as **heavy-quark thresholds**.
- Given the Heavy-quark thresholds are usually (and for convenience) identified with the heavy quark masses by means of the so-called **matching conditions** presently know up to $O(\alpha_s^2)$ [hep-ph/9612398].
- However, heavy-quark thresholds are actually free parameters and can be chosen **arbitrarily**.
- If masses and thresholds are taken to be different, the matching conditions need to be "generalized" including **logarithmic terms**.
- ✓ APFEL now implements the possibility to set masses and thresholds to different values in a consistent way both in the pole mass and in the MS renormalization schemes.

Displaced Heavy-Quark Thresholds



Displaced Heavy-Quark Thresholds





While being an extremely useful tool, APPLgrid might not be appropriate to be directly employed in a global PDF fit where usually thousands of iterations are needed:

• Need to calculate PDF and α_s evolution in real time.

Solution Not particularly fast convolution.

many tables need to be loaded with the concrete risk of exceeding the memory limit (pretty common on clusters).

✓ In the NNPDF collaboration we developed APFELgrid which, starting from an existing APPLgrid, combines PDF evolution from APFEL to the hard cross sections producing *derived* interpolation tables (FK tables):

Observable	APPLGRID	\mathbf{FK}	optimized FK
W^+ production	$1.03 \mathrm{\ ms}$	0.41 ms (2.5 x)	0.32 ms (3.2 x)
Inclusive jet production	$2.45 \mathrm{\ ms}$	$20.1 \ \mu s \ (120 x)$	$6.57 \ \mu s \ (370 x)$

• APFELgrid will soon be made **public in APFEL.** [thanks to N. Harthland]



In the previous versions of APFEL the DGLAP evolution equations were written in terms of the **evolution operator**:

$$\mu^2 \frac{\partial}{\partial \mu^2} M_{ij}(\mu, \mu_0) = P_{ik}(\mu) \otimes M_{kj}(\mu, \mu_0) \quad \text{with} \quad f_i(\mu) = M_{ij}(\mu, \mu_0) \otimes f_j(\mu_0)$$

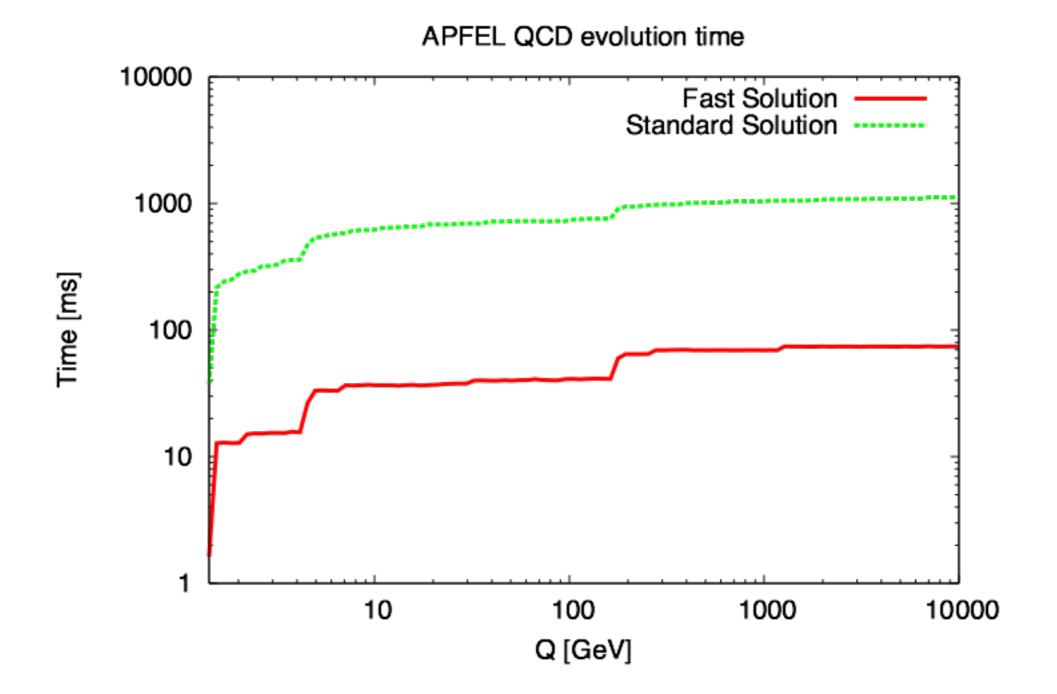
- This may be convenient because the evolution operator can be evaluated once and for all and convoluted with any initial PDF set.
- On the other hand, this requires solving numerically a big coupled system of ODEs, therefore it can be slow.
- Iternatively, one can directly solve the DGLAP equations in terms of **PDFs**:

$$\mu^2 \frac{\partial}{\partial \mu^2} f_i(\mu) = P_{ij}(\mu) \otimes f_j(\mu)$$

This requires the solution of a much smaller system of equations and is consequently much faster.

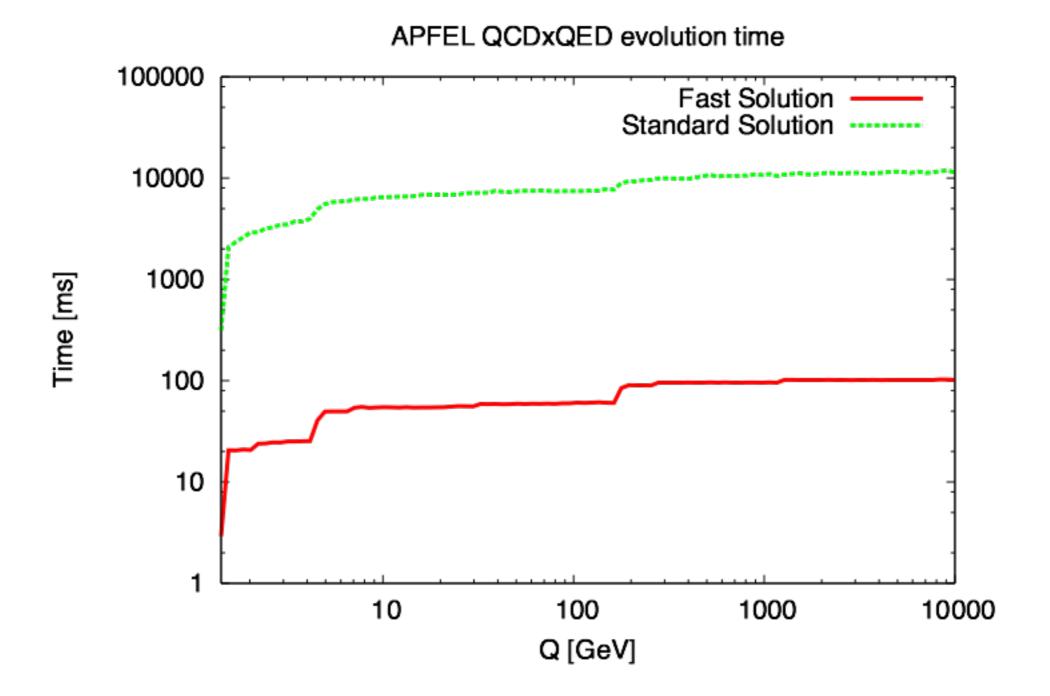


Comparison between old (operatorial) and new (in terms of PDFs) solution for the QCD evolution:





Comparison between old (operatorial) and new (in terms of PDFs) solution for the QCD+QED evolution:



Improvements A New QCD+QED Evolution

In the previous versions of APFEL the QCD+QED evolution was performed by combining the **separate** QCD and QED evolution:

Solution we were due to **subleading terms in** α .

We have now implemented a new evolution basis which allows a simultaneous diagonalization of the QCD+QED evolution matrix:
1) g

2)
$$\gamma$$

3) $\Sigma = \Sigma_u + \Sigma_d$
4) $\Delta_{\Sigma} = \Sigma_u - \Sigma_d$
5) $T_1^u = u^+ - c^+$
6) $T_2^u = u^+ + c^+ - 2t^+$
7) $T_1^d = d^+ - s^+$
8) $T_2^d = d^+ + s^+ - 2b^+$
9) $V = V_u + V_d$
10) $\Delta_V = V_u - V_d$
11) $V_1^u = u^- - c^-$
12) $V_2^u = u^- + c^- - 2t^-$
13) $V_1^d = d^- - s^-$
14) $V_2^d = d^- + s^- - 2b^-$

Improvements A New QCD+QED Evolution

In the previous versions of APFEL the QCD+QED evolution was performed by combining the **separate** QCD and QED evolution:

Solution with the differences, of a few % at most, with the standard implementations which evolve contemporaneously in QCD and QED were due to **subleading terms in** α .

We have now implemented a new evolution basis which allows a **simultaneous diagonalization** of the QCD+QED evolution

matrix:

 $\begin{array}{l} 1) \ g \\ 2) \ \gamma \\ 3) \ \Sigma = \Sigma_u + \Sigma_d \\ 4) \ \Delta_{\Sigma} = \Sigma_u - \Sigma_d \end{array} \qquad \begin{array}{l} 9) \ V = V_u + V_d \\ 10) \ \Delta_V = V_u - V_d \end{aligned} \qquad \begin{array}{l} \mbox{Coupled} \\ 5) \ T_1^u = u^+ - c^+ & 11) \ V_1^u = u^- - c^- \\ 6) \ T_2^u = u^+ + c^+ - 2t^+ & 12) \ V_2^u = u^- + c^- - 2t^- \\ 7) \ T_1^d = d^+ - s^+ & 13) \ V_1^d = d^- - s^- \\ 8) \ T_2^d = d^+ + s^+ - 2b^+ & 14) \ V_2^d = d^- + s^- - 2b^- \end{array} \qquad \begin{array}{l} \mbox{Coupled} \\ \end{array}$

This new basis is also suitable for an easy implementation of the mixed **higher order corrections** to the evolution.

Intrinsic Charm

Introducing an intrinsic charm (IC) component in the context of a GM-VFNS like FONLL (or ACOT, or TR) requires some care:

verturbative generation of heavy quarks at the thresholds,

• take into account **charm-initiated diagrams** both in the **massive** and in the massless sectors [arXiv:1510.00009].

✓A full formulation of the FONLL scheme in the presence of IC has recently been achieved [arXiv:1510.02491]:

interestingly, it has been found that FONLL with IC is equivalent to full ACOT to all orders, while the standard FONLL (w/o IC) is instead equivalent to S-ACOT.

Implemented in APFEL up to NLO both in the NC and CC sector and benchmarked against the public massiveDISsFuntion code (https://www.ge.infn.it/~bonvini/massivedis/).

Intrinsic Charm

Consider realistic models:

Jet BHPS model:

$$f_c^{(3)}(x) = f_{\overline{c}}^{(3)}(x) = Ax^2 \left[6x(1+x)\ln x + (1-x)(1+10x+x^2) \right]$$

SEA model:

$$f_c^{(3)}(x) = f_{\overline{c}}^{(3)}(x) = Ax^{-1.25}(1-x)^3$$

 $\checkmark A$ determined requiring the charm to carry 0.5% of the momentum

