Monte Carlos Herwig and Transverse Momentum

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The Herwig Event Generator



Herwig++ has seen a ten-year development to meet a milestone intended to succeed the FORTRAN HERWIG program.

This milestone evolved over time as the experimental and phenomenological needs did.

On top of its first definition (= at least as good as HERWIG), precision has become the key goal: NLO matched by default.

Herwig++ 3.0 \rightarrow Herwig 7.0

Herwig 7 – Core Features

[Bellm, Gieseke, Grellscheid, Plätzer, M. Rauch, Reuschle, Richardson, Schichtel, Seymour, Siodmok, Wilcock, Fischer, Harrendorf, Nail, Papaefstathiou, D. Rauch – EPJ C 76 (2016) 196]

NLO matched to parton showers as default for the hard process.

- \rightarrow Fully automated, only linking external codes to calculate amplitudes.
- \rightarrow Run in a single program, no event files to move around.
- \rightarrow Subtractive (MC@NLO-type) and multiplicative (POWHEG-type) matching.

Two showers: Angular-ordered and dipole shower.

Spin correlations and QED radiation in angular ordered shower.

Facilities for **parton shower uncertainties.**

Improved kinematics reconstruction.

Vastly improved documentation, usage and installation.

New tunes taking NLO matching into account + much, much more ...

Shower Algorithms

Two shower algorithms available:

"QTilde"

[Gieseke, Stephens, Webber – JHEP 0312 (2003) 045]

- → "Traditional" angular ordered shower: default shower
- \rightarrow QED, spin correlations, shower variations, decays
- → Truncated showering for Powheg-type matching

"Dipole"

[Plätzer, Gieseke – JHEP 1101 (2011) 024]

- \rightarrow Dipole-type evolution, ordered in dipole pt
- \rightarrow Extensive shower variations, decays soon (7.1)
- → Working horse for NLO multijet merging

Outline

Backward Evolution

Kinematics How (Not) To

Uncertainties

Summary

Backward Evolution

Backward evolution with emissions at finite transverse momentum. Phase space bounds link longitudinal & transverse, recoil not unique.

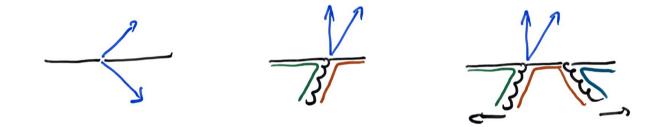
$$\overline{\Pi}(x,q_{L},Q) = \exp\left(-\int_{q_{L}}^{Q} \frac{dh}{h} \int_{z_{L}}^{Q} \frac{dx}{2} \frac{f(x_{l},q_{L},Q)}{f(x_{l},k)} V(q_{l},k,Q)\right)$$

$$\frac{1}{q_{L}} \int_{Q}^{Q} \frac{dh}{h} \int_{z_{L}}^{Q} \frac{dx}{2} \frac{f(x_{l},q_{L},Q)}{f(x_{l},k)} V(q_{l},k,Q)\right)$$

$$\overline{\Pi}(x,q_{L},Q) = \frac{1}{q_{L}} \int_{z_{L}}^{Q} \frac{1}{q_{L}}$$

Kinematics How Not To

Problem in initial dipole-type approaches: Z recoil from first emission only. Spectrum vanishes at zero pt from shower only.



Initial state spectator not changed, beams always aligned.

→ Typical to dipole subtraction terms, so eases matched calculation.

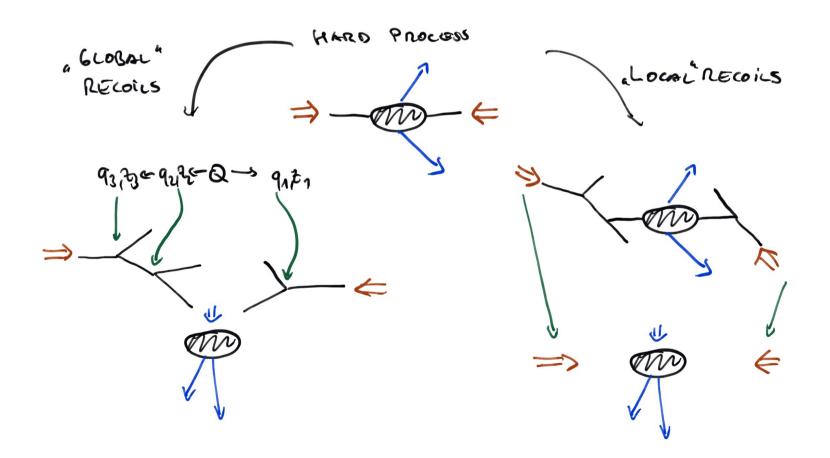
Similar issue in final state?

 \rightarrow Would be visible in EEC, however dipoles provide excellent description.

Kinematics How To

Make multiple emissions contribute to final state pt. Common to angular ordered shower from the beginning, equivalent strategy developed for the dipole shower.

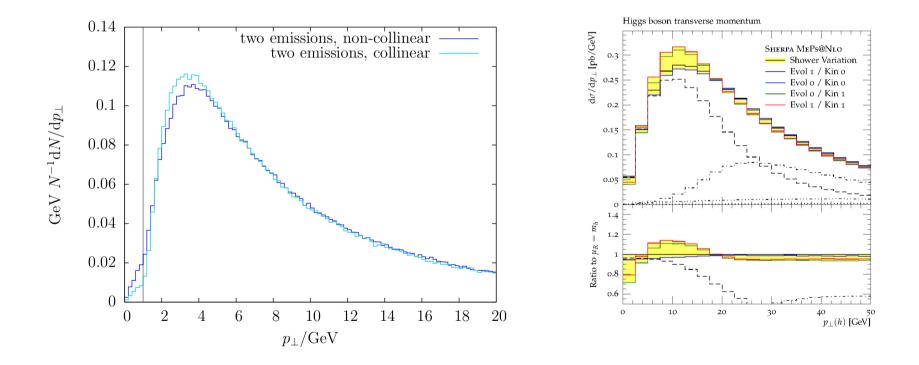
[Plätzer, Gieseke – JHEP 1101 (2011) 024]



Significant effect?

[Plätzer, Gieseke – JHEP 1101 (2011) 024] similar approach in Sherpa, e.g. [Hoeche, Schönherr, Krauss – Phys.Rev. D90 014012]

Look at pT spectra for Drell-Yan or Higgs:

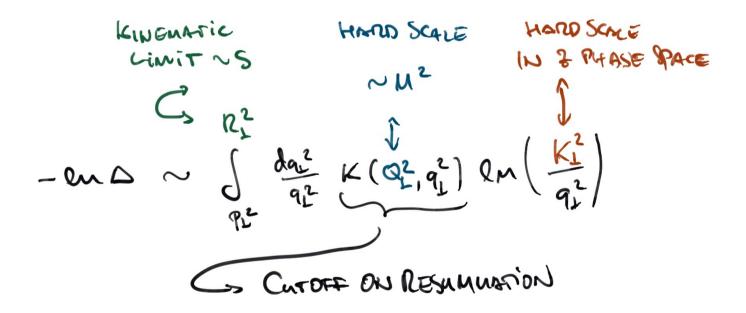


"Algorithmic" uncertainty drowned in "perturbative" uncertainty?

Logarithmic structure → Uncertainties

[Bellm, Nail, Plätzer, Schichtel, Siodmok – arXiv:1605.01338]

Look at generic Sudakov exponent:



AlphaS running on top, also PDF arguments.

Uncertainties

[Bellm, Nail, Plätzer, Schichtel, Siodmok – arXiv:1605.01338]

Aim at evaluating event generator uncertainties in a global prescription

- \rightarrow Need to evaluate uncertainties of building blocks one at a time.
- \rightarrow Then pin down cross feed, making minimal assumptions.

Start with the perturbative part: Parton showers – at leading order! Then check if matching algorithms exhibit the expected improvement.

Shower scale variations not a priori clear to serve as estimating an order one term in the next (logarithmic) order – logarithmic accuracy mostly unclear.

Rather constrain by demanding **controllable** uncertainties:

- → Small/large where showers are expected to be reliable/unreliable.
- → Consistent between two systematically different algorithms.
- \rightarrow Not to mess around with hard process input.

Uncertainty Benchmarks with Herwig 7

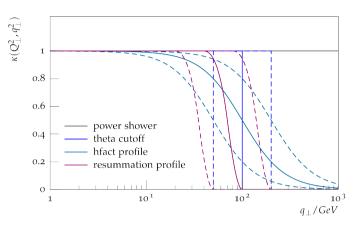
[Bellm, Nail, Plätzer, Schichtel, Siodmok – arXiv:1605.01338]

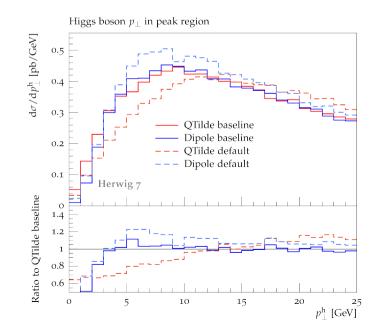
Resummation needs to be cut off at a typical hard scale \rightarrow veto on hard emissions, region to be filled by matching.

Resummation properties are heavily influenced by the way resummation is being switched off.

Study scale variations in angular ordered and Dipole showers at a benchmark setting where we observe absolutely comparable resummation properties:

Hard veto scales, factorization/renormalization scales in the shower and hard process.





Uncertainty Benchmarks with Herwig 7

1st jet) [pb]

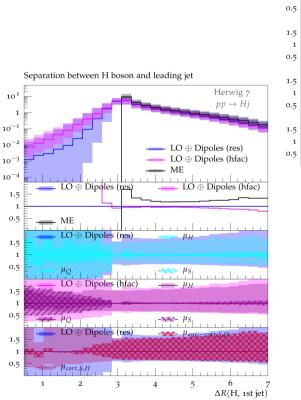
 $d\sigma/d\Delta R(H,$

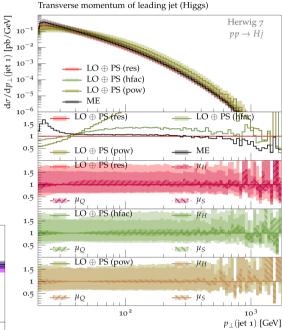
[Bellm, Nail, Plätzer, Schichtel, Siodmok – arXiv:1605.01338]

Choice of the hard veto scale is crucial to reproduce hard process input: typically average transverse momenta of hard objects.

Controllable uncertainties can only be established by narrow, smeared versions of a theta function, confirming simple LL arguments.

We can now check the impact of higher order improvements.







Contemporary shower algorithms generate transverse momenta from many initial state emissions.

Different prescriptions possible, need to be confronted with other sources of uncertainty.

Available shower phase space and "smearing" of pT dependence to cutoff resummation are crucial to variation pattern.

Thank you!