

High precision measurement of a_{μ}^{HLO} with a 150 GeV μ beam on e^{-} target at CERN

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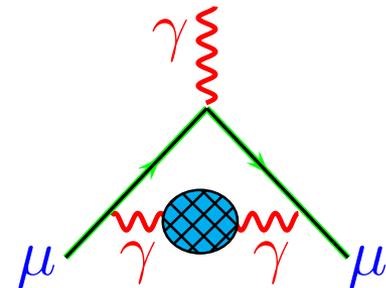
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Outline

- Muon g-2: summary of the present status
- New approach to compute a_{μ}^{HLO} in the space-like region
- Proposal for a measurement a_{μ}^{HLO} with $\mu e \rightarrow \mu e$ at CERN
- Detector considerations
- Conclusions

Muon g-2: summary of the present status

- E821 experiment at BNL has generated enormous interest:

$$a_{\mu}^{E821} = 11659208.9(6.3) \times 10^{-10} \quad (\text{o.54 ppm})$$

- Tantalizing $\sim 3\sigma$ deviation with SM (persistent since >10 years):

$$a_{\mu}^{SM} = 11659180.2(4.9) \times 10^{-10} \quad (DHMZ)$$

M. Davier, A. Hoecker, B. Malaescu and Z. Zhang, Eur. Phys. J. C71 (2011)

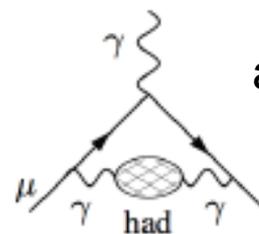
$$a_{\mu}^{E821} - a_{\mu}^{SM} \sim (28 \pm 8) \times 10^{-10}$$

- Current discrepancy limited by:

- Experimental** uncertainty \rightarrow New experiments at FNAL and J-PARC **x4** accuracy
- Theoretical** uncertainty \rightarrow limited by hadronic effects

$$a_{\mu}^{SM} = a_{\mu}^{QED} + \boxed{a_{\mu}^{HAD}} + a_{\mu}^{Weak}$$

Hadronic Vacuum polarization (HLO)



$$a_{\mu}^{HLO} = (692.3 \pm 4.2) 10^{-10}$$

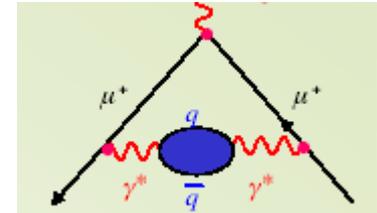
$$\delta a_{\mu} / a_{\mu} \sim 0.6\%$$

a_μ^{HLO} calculation, traditional way: time-like data

$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} \sigma_{e^+e^- \rightarrow \text{hadr}}(s) K(s) ds$$

$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)} \sim \frac{1}{s} \quad \sigma_{e^+e^- \rightarrow \text{hadr}}(s) = \frac{4\pi}{s} \text{Im} \Pi_{\text{had}}(s)$$

$$a_\mu = (g-2)/2$$

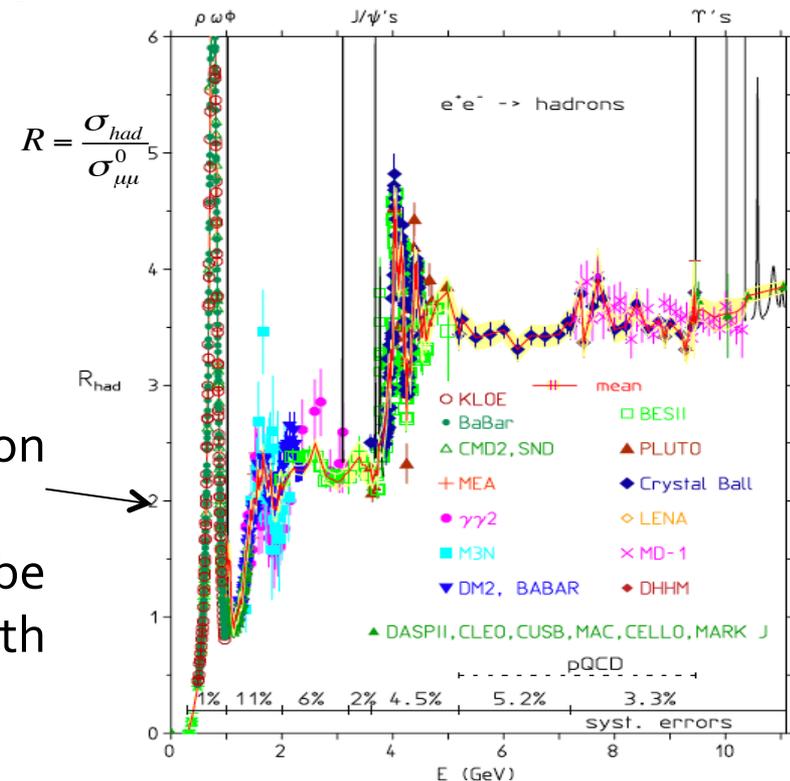


$$2 \text{Im} \left[\text{loop} \right] = \left| \text{cut} \right|^2$$

Traditional way: based on precise experimental (time-like) data:

$$a_\mu^{\text{HLO}} = (692.3 \pm 4.2) 10^{-10} \text{ (DHMZ)}$$

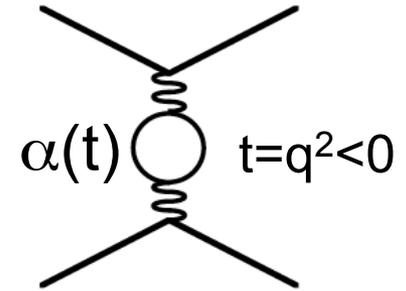
- Main contribution in the low energy region (highly fluctuating!)
- Current precision at 0.6% \rightarrow needs to be reduced by a factor ~ 2 to be competitive with the new $g-2$ experiments



Alternative approach: a_μ^{HLO} from space-like region

[C.M. Carloni Calame, M. Passera, L. Trentadue, G. Venanzoni
Phys.Lett. B746 (2015) 325-32]

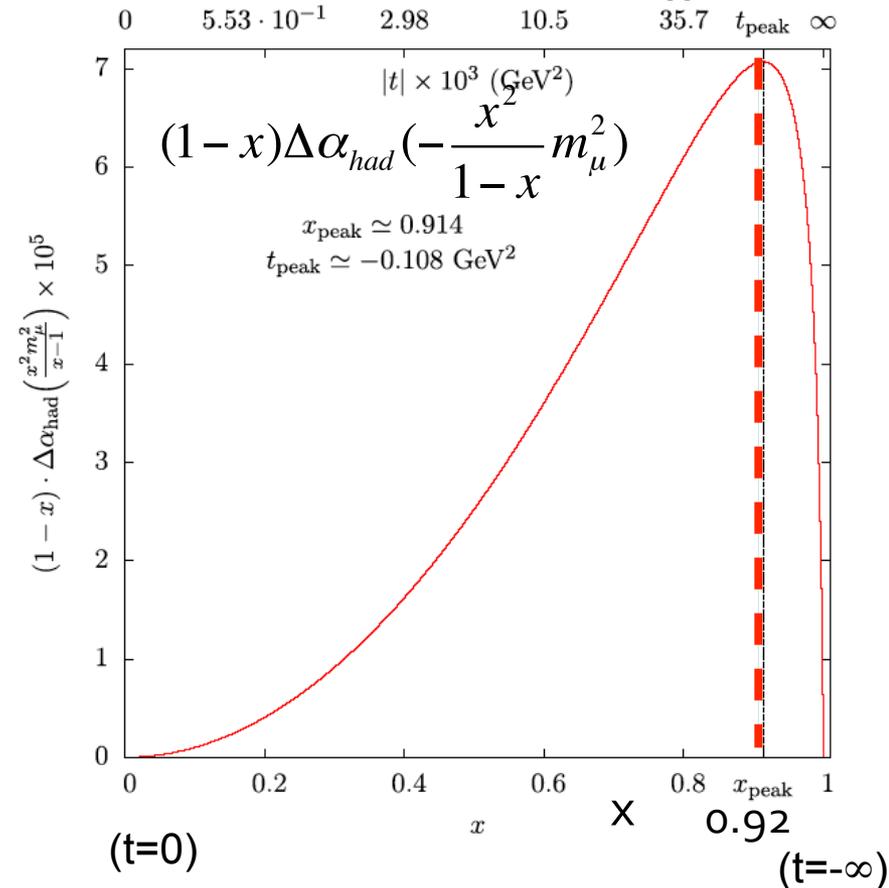
$$a_\mu^{\text{HLO}} = -\frac{\alpha}{\pi} \int_0^1 (1-x) \Delta\alpha_{\text{had}}\left(-\frac{x^2}{1-x} m_\mu^2\right) dx$$



$$t = \frac{x^2 m_\mu^2}{x-1} \quad 0 \leq -t < +\infty$$

$$x = \frac{t}{2m_\mu^2} \left(1 - \sqrt{1 - \frac{4m_\mu^2}{t}}\right); \quad 0 \leq x < 1;$$

$t = -0.11 \text{ GeV}^2$
($\sim 330 \text{ MeV}$)



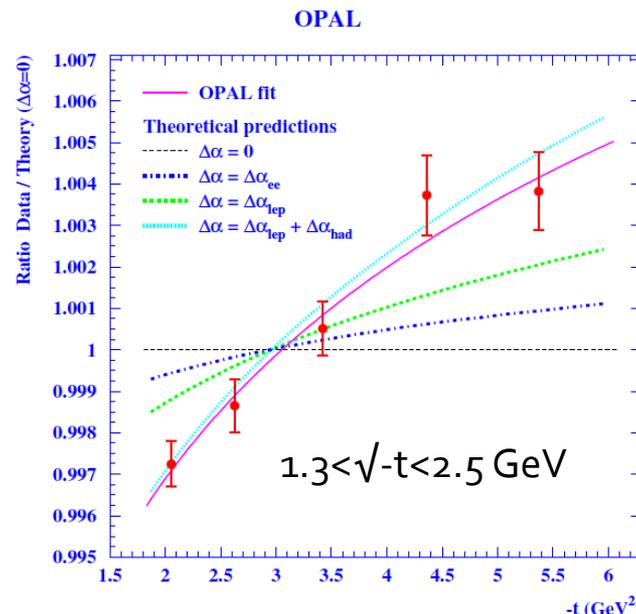
- a_μ^{HLO} is given by the integral of the curve (smooth behaviour)
- It requires a measurement of the hadronic contribution to the effective electromagnetic coupling in the space-like region $\Delta\alpha_{\text{had}}(t)$ ($t=q^2 < 0$)
- It enhances the contribution from low q^2 region (below 0.11 GeV^2)
- Its precision is determined by the uncertainty on $\Delta\alpha_{\text{had}}(t)$ in this region

Measurement of $\Delta\alpha_{\text{had}}(t)$ spacelike at LEP

- $\Delta\alpha_{\text{had}}(t)$ ($t < 0$) has been measured at LEP using small angle Bhabha scattering

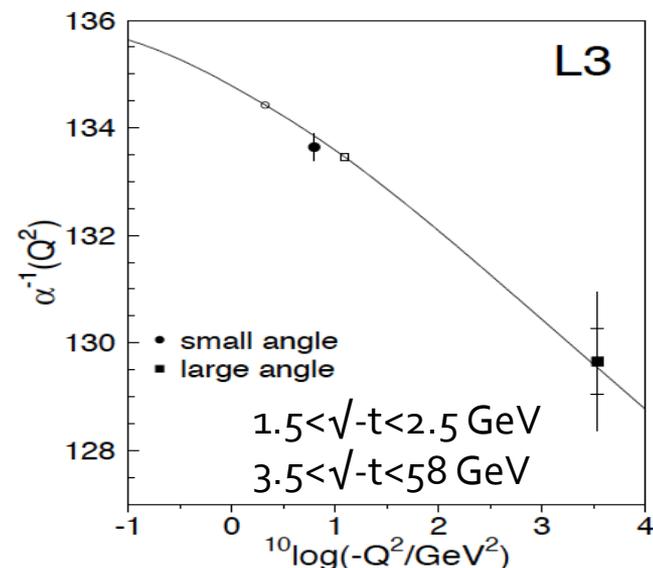
$$f(t) = \frac{N_{\text{data}}(t)}{N_{\text{MC}}^0(t)} \propto \left(\frac{1}{1 - \Delta\alpha(t)} \right)^2.$$

Accuracy at per mill level was achieved!



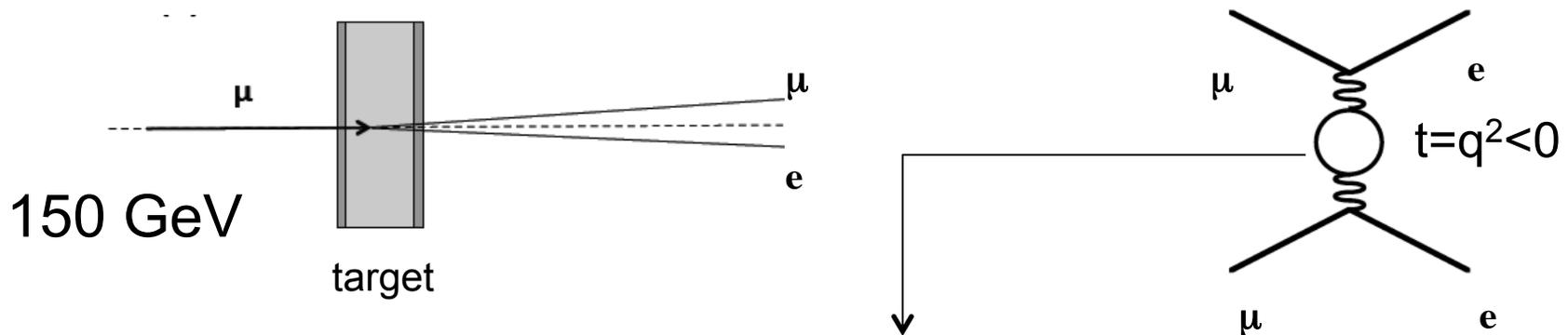
G. Abbiendi et al., Eur. Phys. J. C 45, 1–21 (2006)

- For low t values ($\leq 0.11 \text{ GeV}^2$), like in our case flavour factories are needed!



M. Acciarri et al., Phys. Lett. B476 40–48 (2000)

$\Delta\alpha_{\text{had}}(t)$ can be measured with high energy muon beam ($E \sim 150$ GeV) on electron low-Z target, through the elastic scattering $\mu e \rightarrow \mu e$



$$\left| \frac{\alpha(t)}{\alpha_0} \right|^2 = \left| \frac{1}{1 - \Delta\alpha(t)} \right|^2$$

Why measuring $\Delta\alpha_{had}(t)$ with a 150 GeV μ beam on e^- target ?

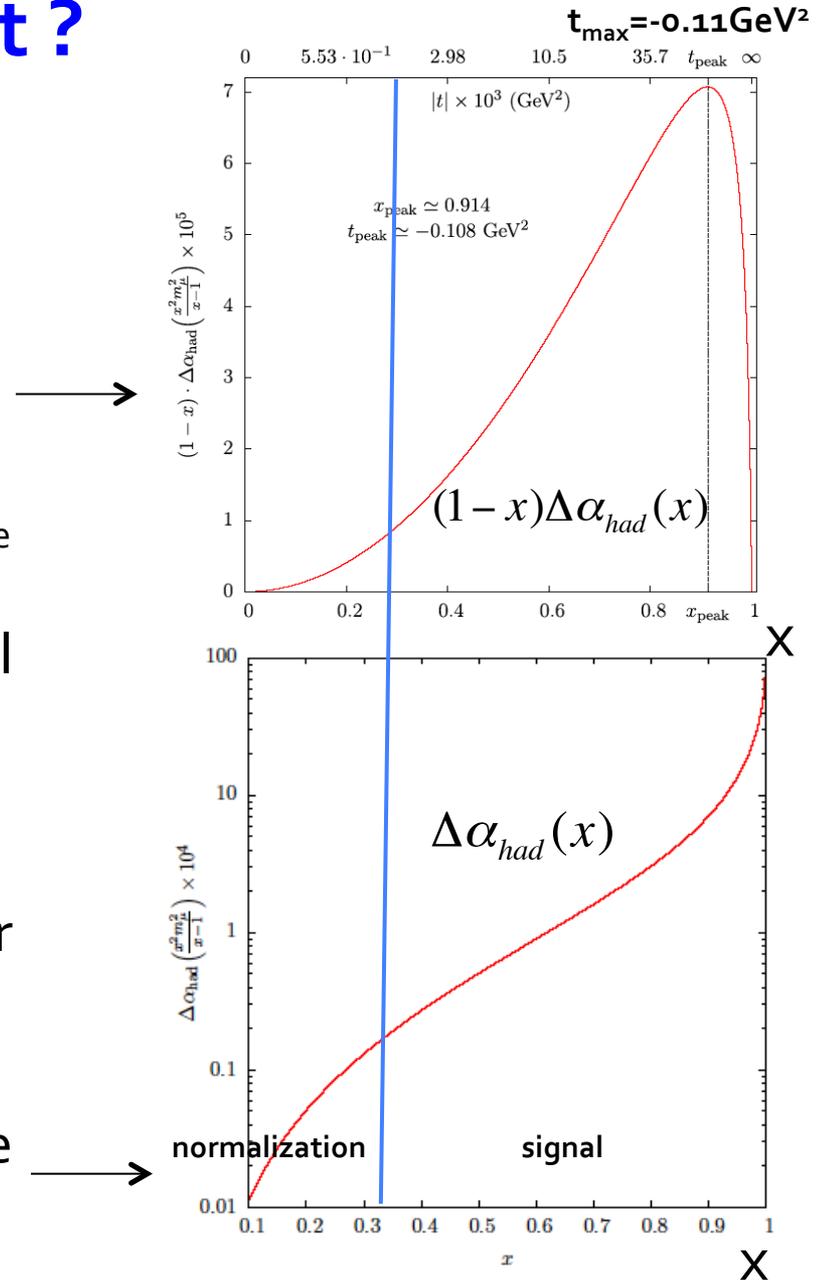
It looks an ideal process!

- $\mu e \rightarrow \mu e$ is pure t-channel (at LO)
- It gives $0 < -t < 0.161 \text{ GeV}^2$ ($0 < x < 0.93$)
- The kinematics is very simple: $t = -2m_e E_e$
- High boosted system gives access to all angles (t) in the cms region

$$\theta_e^{LAB} < 32 \text{ mrad} \quad (E_e > 1 \text{ GeV})$$

$$\theta_\mu^{LAB} < 5 \text{ mrad}$$

- It allows using the same detector for signal and normalization
- Events at $x \sim 0.3$ ($t \sim -10^{-3} \text{ GeV}^2$) can be used as normalization ($\Delta\alpha_{had}(t) < 10^{-5}$)



Detector considerations I

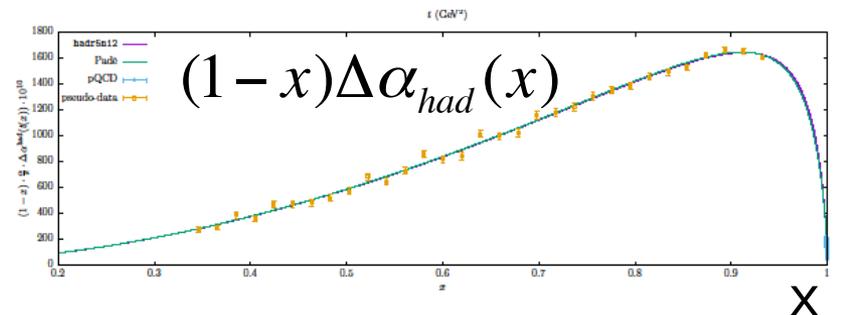
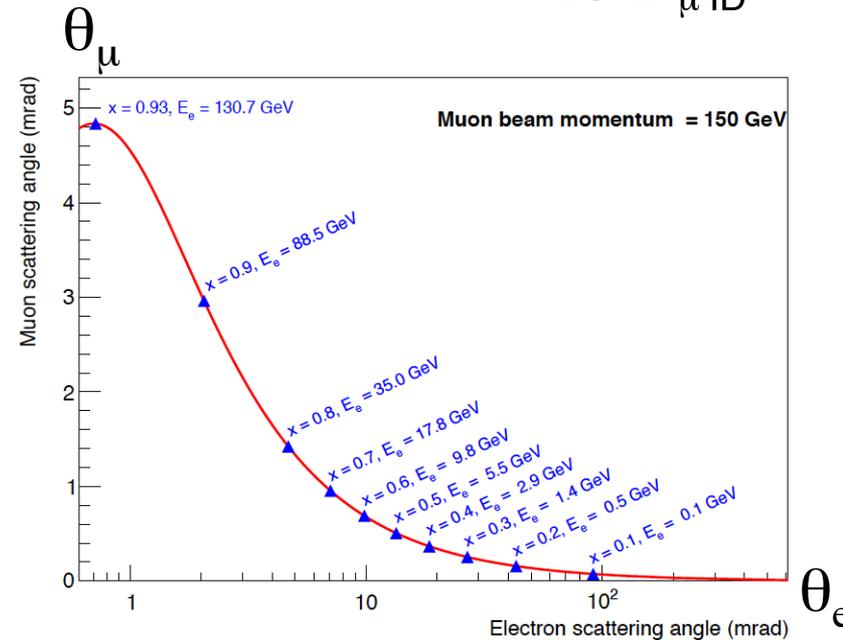
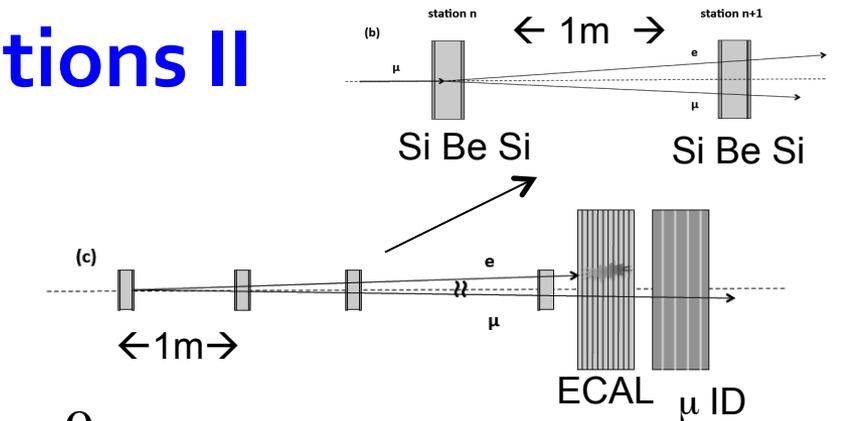
- In order to be competitive with a_{μ}^{HLO} from time-like data (0.6% error) a subpercent uncertainty on a_{μ}^{HLO} is required

$$\delta\Delta\alpha_{\text{had}}(t) \sim 0.5 \sqrt{\left(\frac{\delta N^{\text{data}}(t)}{N^{\text{data}}(t)}\right)^2 + \left(\frac{\delta N^{\text{norm}}(t_0)}{N^{\text{norm}}(t_0)}\right)^2 + \left(\frac{\delta R^{\text{MC}}}{R^{\text{MC}}}\right)^2} + \text{corr. terms}$$
$$R^{\text{MC}} = \frac{d\sigma_0^{\text{MC}}(t)}{d\sigma_0^{\text{MC}}(t_0)}$$

- $\delta\Delta\alpha_{\text{had}}/\Delta\alpha_{\text{had}}$ at 0.5% at peak region ($x=0.92$, $\Delta a_{\text{had}} \sim 10^{-3}$) \rightarrow
 $\delta N(t)/N(t) \sim 10^{-5}$
- Such an accuracy demands **high statistics** keeping **low systematic** errors!
- **Dense** (active) target would provide the required statistics at a price of an unavoidable large multiple scattering and background process (pair production, bremsstrahlung, nuclear interaction)
- Our **choice** goes to **light** Z (Be) target with a modular apparatus which minimizes systematic errors

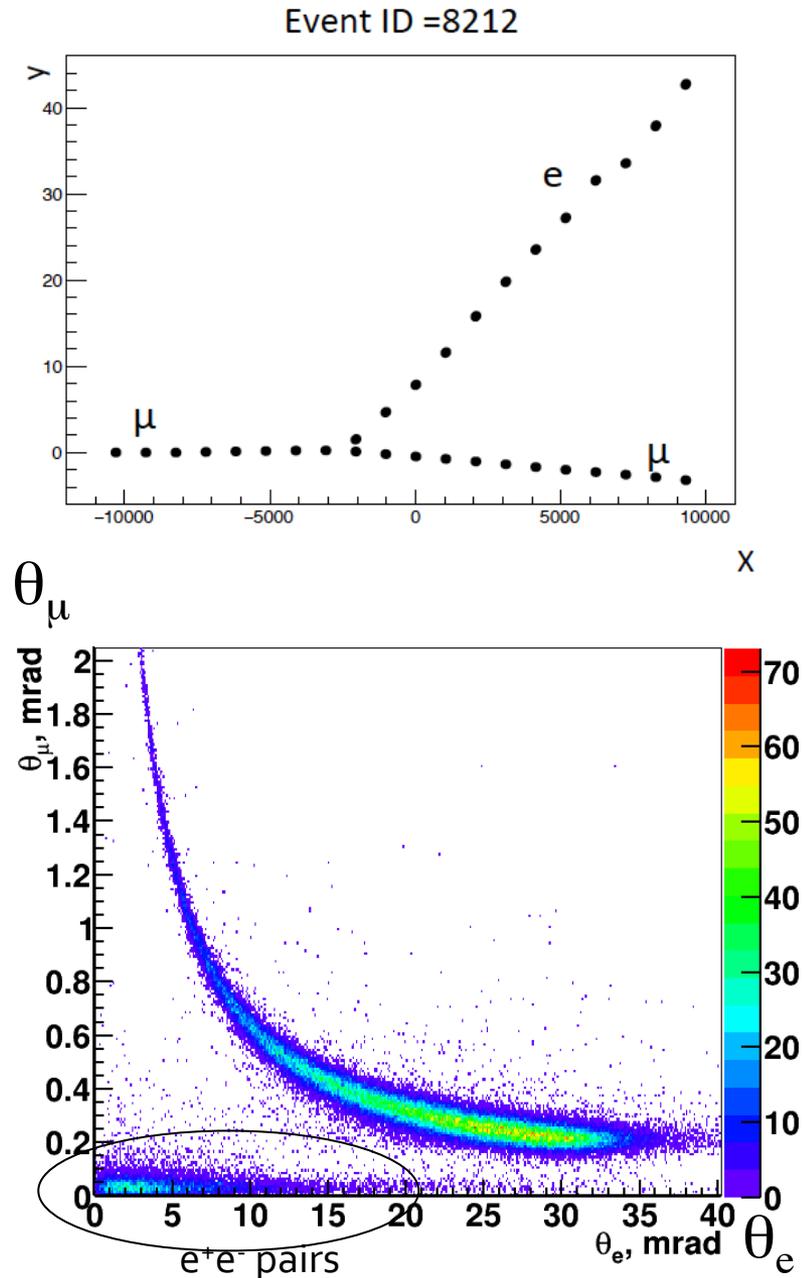
Detector considerations II

- Modular apparatus: 20 layers of 3 cm Be (target), each coupled to 1 m distant Si (0.3 mm) planes. It provides a 0.02 mrad resolution on the scattering angle
- The $t=q^2 < 0$ of the interaction is determined by the electron (or muon) scattering angle (a' la NA7)
- ECAL and μ Detector located downstream to solve PID ambiguity below 5 mrad. Above that, angular measurement gives correct PID
- It provides uniform full acceptance, with the potential to keep the systematic errors at 10^{-5} (main effect is the multiple scattering for normalization which can be studied by data)
- **Statistical considerations show that a 0.3% error can be achieved on a_{μ}^{HLO} in 2 years of data taking with $2 \times 10^7 \mu/s$**



First simulation

- A simulation of the detector based on GEANT-4 has started
- μ -e elastic scattering events have a clear topology
- Background events can be easily identified and rejected in the θ_μ vs θ_e plane
- Multiple scattering can be studied by data as it **breaks** the μ -e two-body angular correlation, moving events out of the kinematic constraint. It also causes **acoplanarity**, while two-body events are planar.
- Simulation will help to optimize the detector (i.e. additional thin layer(s) can be placed for luminosity)



Conclusion I

- We present an alternative method to determine the full contribution to a_{μ}^{HLO} based on the measurement of $\Delta\alpha_{\text{had}}(\mathbf{t})$ in **space-like region**. It consists in measuring the scattering $\mu e \rightarrow \mu e$ using a high energy muon beam ($E \sim 150$ GeV) available in the North Area at CERN on electron target
- A detector concept based on a modular structure of light Z (Be) target and Silicon planes allows to reach a statistical uncertainty of **$\sim 0.3\%$** on a_{μ}^{HLO} in ~ 2 years of data taking and an angular precision of ~ 0.02 mrad on the scattering angle
- The $t=q^2 < 0$ of the events can be obtained by a precise measurement of e and μ scattering angles
- Such a simple and robust technique has the potential to keep systematic effects under control

Conclusion II

- Theory side: high precision MC must be developed to control the systematics. The present knowledge of QED Radiative Corrections is at a few 10^{-4} level; work is in progress to extend MC used at flavour factories (BabaYaga) to μe scattering with expected accuracy at (better than) 10^{-5} on cross section ratios
- In experiments dedicated to high-precision measurements, several systematic effects can be explored within the experiment itself. In this respect the proposed modularity of the apparatus will help. **A test with a single module could provide a proof-of-concept of the proposed method.**

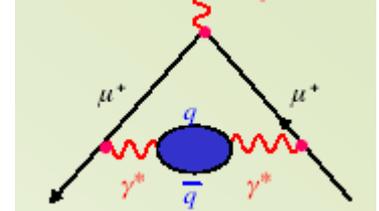
Thanks!

Spare

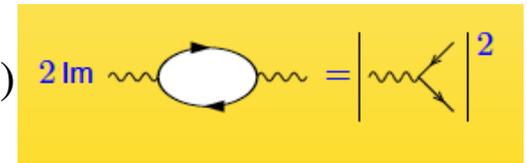
a_μ^{HLO} calculation, traditional way: time-like data

$$a_\mu^{\text{HLO}} = \frac{1}{4\pi^3} \int_{4m_\pi^2}^{\infty} \sigma_{e^+e^- \rightarrow \text{hadr}}(s) K(s) ds$$

$$a_\mu = (g-2)/2$$



$$a_\mu^{\text{HLO}} = \frac{\alpha}{\pi^2} \int_0^{\infty} \frac{ds}{s} K(s) \text{Im} \Pi_{\text{had}}(s) \quad \sigma_{e^+e^- \rightarrow \text{hadr}}(s) = \frac{4\pi}{s} \text{Im} \Pi_{\text{had}}(s)$$

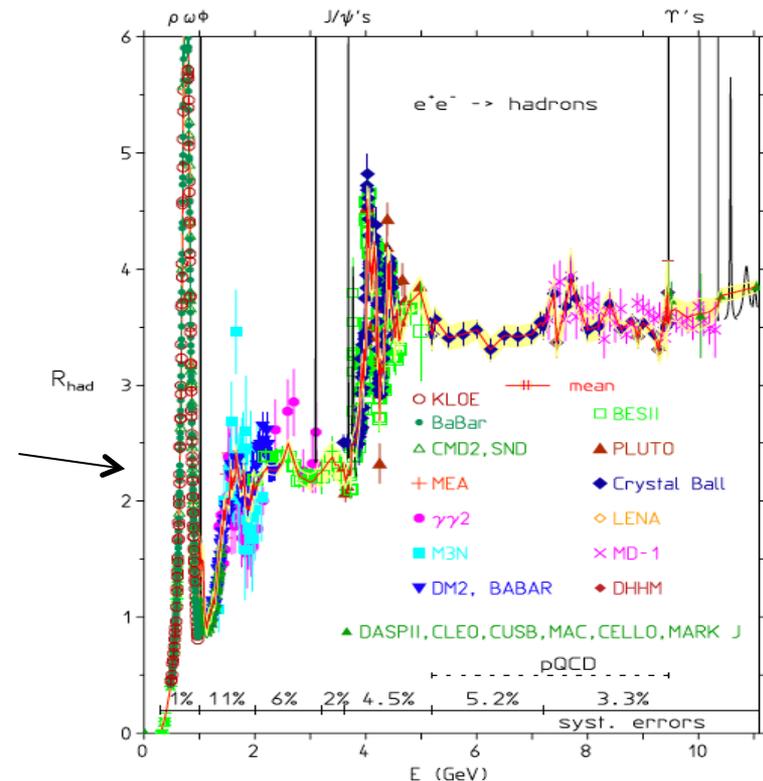


$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)} \sim \frac{1}{s}$$

Traditional way: based on precise experimental (time-like) data:

$$a_\mu^{\text{HLO}} = (692.3 \pm 4.2) 10^{-10} \text{ (DHMZ)}$$

- Main contribution in the low energy region (**highly fluctuating!**)
- Current precision at 0.6% \rightarrow needs to be reduced by a factor **~ 2** to be competitive with the new **E989** experiment



$(g-2)_\mu$: a new experiment at FNAL (E989)

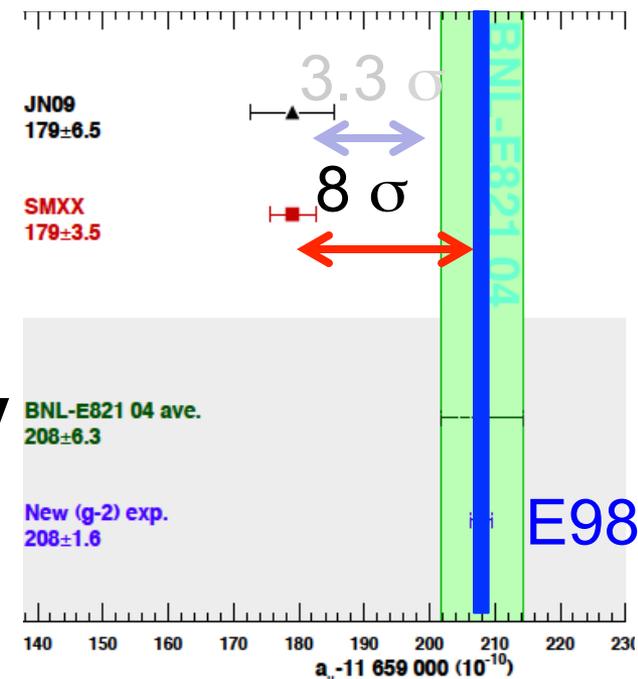
- New experiment at FNAL (E989) at magic momentum, consolidated method. **20 x stat.** w.r.t. E821. Relocate the BNL storage ring to FNAL.

→ $\delta a_\mu \times 4$ improvement (0.14ppm)

If the central value remains the same
 ⇒ 5-8 σ from SM* (enough to claim discovery of **New Physics!**)

***Depending on the progress on Theory**

Thomas Blum; Achim Denig; Ivan Logashenko; Eduardo de Rafael; Lee Roberts, B.; Thomas Teubner; Graziano Venanzoni (2013). "The Muon $(g-2)$ theory Value: Present and Future". [arXiv:1311.2198](https://arxiv.org/abs/1311.2198) [hep-ph].



Complementary proposal at J-PARC in progress

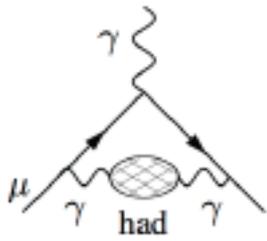
Muon $G-2$: summary of the present status

- E821 experiment at BNL has generated enormous interest
- Tantalizing $\sim 3\sigma$ deviation with SM (persistent since >10 years)
- Current discrepancy limited by:
 - experimental** uncertainty (BNL) \rightarrow New experiments at FNAL and J-PARC
 - Theoretical uncertainty limited by Hadronic Vacuum polarization

$$a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{TH}} \sim (28 \pm 8) \cdot 10^{-10}$$

$$a_{\mu}^{\text{SM}} = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{HAD}} + a_{\mu}^{\text{Weak}}$$

HLO VP



$$a_{\mu}^{\text{HLO}} = (695 \pm 4.3) 10^{-10}$$

$$\delta a_{\mu}^{\text{HLO}} \sim 0.7\%$$

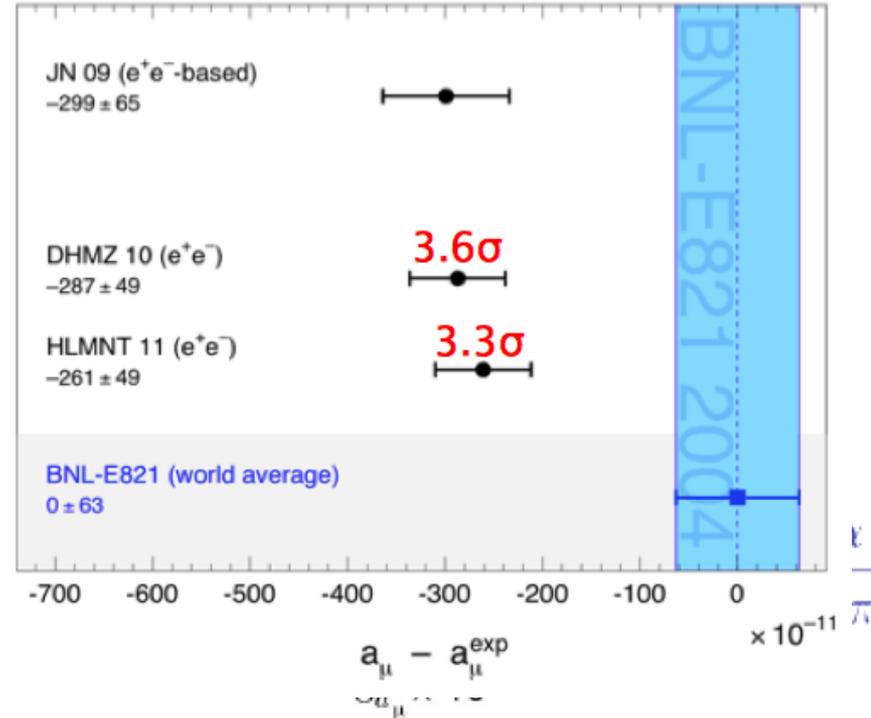
HLbL



$$a_{\mu}^{\text{HLbL}} = (10.5 \pm 2.6) 10^{-10}$$

$$\delta a_{\mu}^{\text{HLbL}} \sim 25\%$$

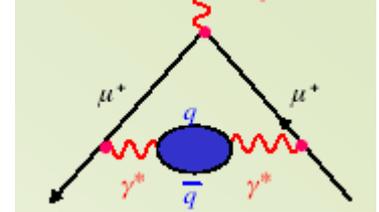
Status: summer 2011 (published results shown only)



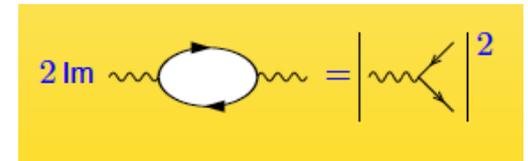
a_μ^{HLO} calculation, traditional way: time-like data

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$$a_\mu = (g-2)/2$$



$$a_\mu^{HLO} = \frac{\alpha}{\pi^2} \int_0^{\infty} \frac{ds}{s} K(s) \text{Im} \Pi_{had}(s) \sigma_{e^+e^- \rightarrow hadr}(s) = \frac{4\pi}{s} \text{Im} \Pi_{had}(s)$$

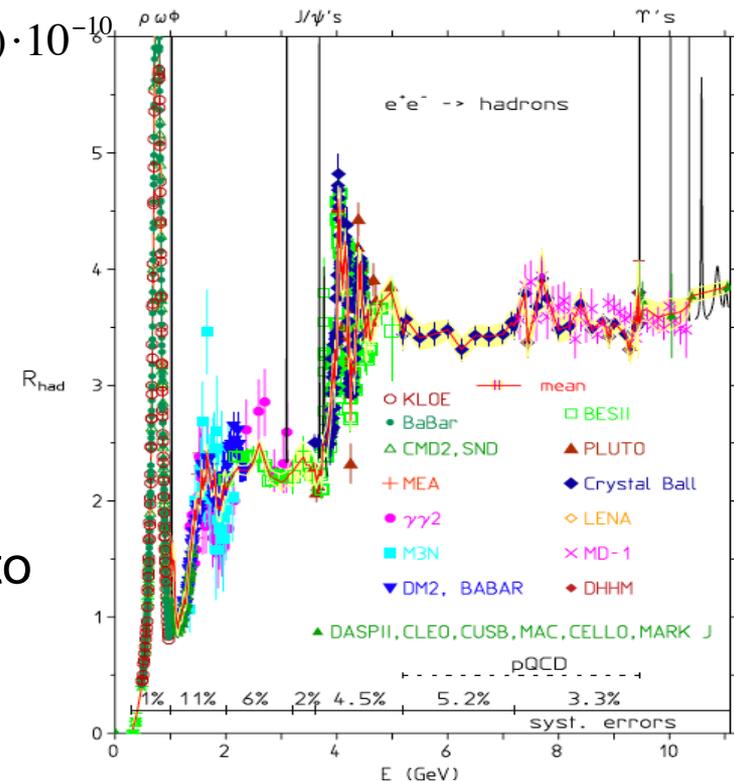


$$K(s) = \int_0^1 dx \frac{x^2(1-x)}{x^2 + (1-x)(s/m^2)} \sim \frac{1}{s} \quad a_\mu^{had} = (689.7 \pm 4.4) \cdot 10^{-10}$$

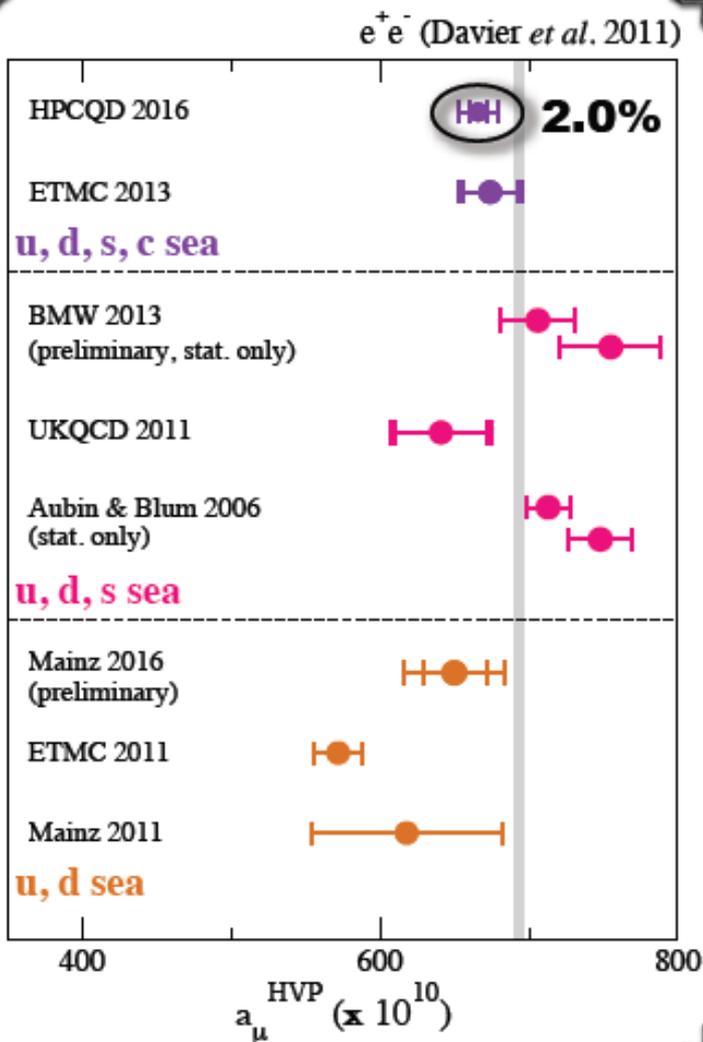
Traditional way: based on precise experimental (time-like) data:

Main contribution in the low energy region (**highly fluctuating!**)

Current precision at 0.7% \rightarrow needs to be reduced by a factor ~ 2 to be competitive with the new E989 experiment: not easy task due to local experimental discrepancies and uncertainties on theoretical ingredients (FSR, Isospin corrections, etc..)



Lattice-QCD progress on a_μ^{HVP}



- Can calculate nonperturbative vacuum polarization function $\Pi(Q^2)$ directly in lattice QCD from simple 2-point correlation function of EM quark current [Blum, PRL 91 (2003) 052001]

- Several ongoing lattice efforts yielding new results since ICHEP 2014 including:

- First calculation of quark-disconnected contribution [RBC/UKQCD, PRL116, 232002 (2016)]

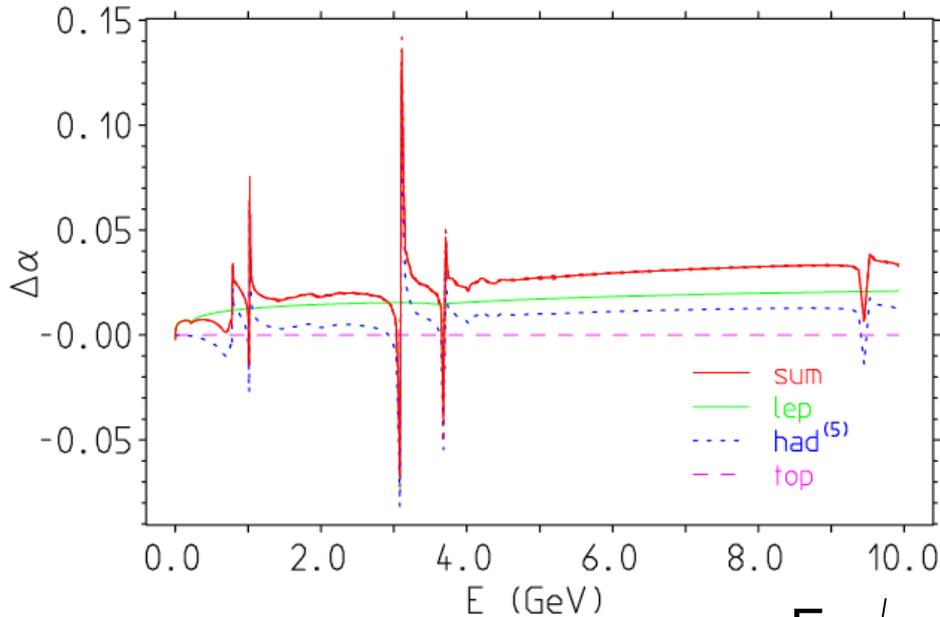
- Second complete calculation of leading-order a_μ^{HVP} [HPQCD, arXiv:1601.03071]

- First to reach precision needed to observe significant deviation from experiment

- ~1% total uncertainty by 2018 possible

- Sub-percent precision will require inclusion of isospin breaking & QED, and hence take longer

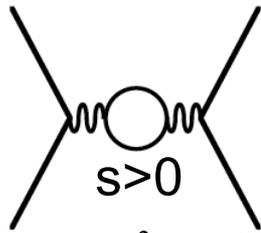
Running of α_{em}



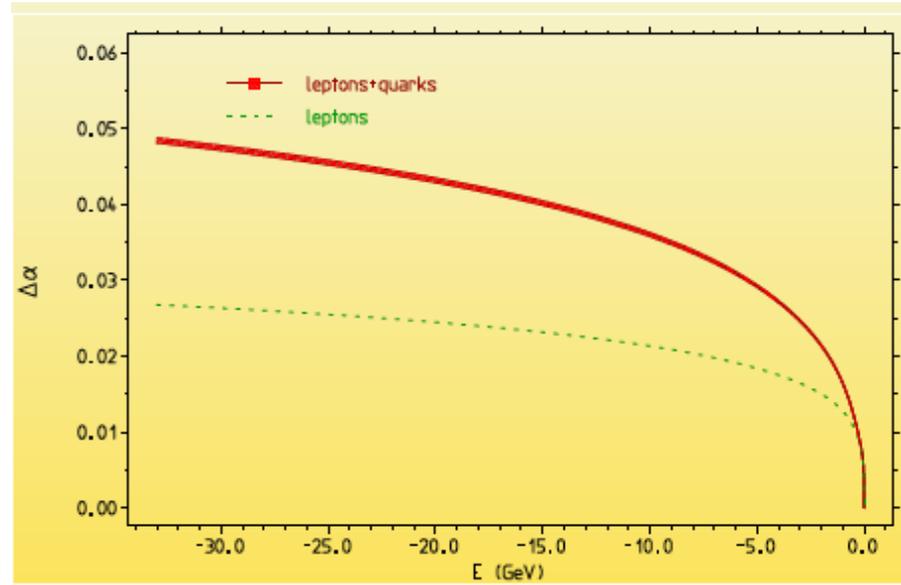
Time-like

$E = \sqrt{s}$

Behaviour characterized by the opening of resonances



$$\Delta\alpha_{had}^{(5)}(M_Z^2) = -\frac{\alpha M_Z^2}{3\pi} \text{Re} \int_{4m_\pi^2}^{\infty} ds \frac{R(s)}{s(s - M_Z^2 - i\epsilon)}$$



Space-like

$E = -\sqrt{-t}$

Very smooth behaviour

