

Generalized Einstein-Aether as a dark energy fluid

**Progress on Old and New Themes in cosmology
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Motivation

- Λ CDM is the most popular cosmological model for the universe.
- But there are problems...
- Alternatives? An easy way is to add a new field; scalar (Quintessence...), vector (Einstein-Aether...), tensor (bimetric...)

Generalized Einstein-Aether: the model

$$\mathcal{L}_A = M^2 \mathcal{F}(\mathcal{K}) + \lambda(g_{\mu\nu} A^\mu A^\nu + 1)$$

$$\mathcal{K} = \frac{1}{M^2} K^{\alpha\beta}{}_{\mu\nu} \nabla_\alpha A^\mu \nabla_\beta A^\nu \quad K^{\alpha\beta}{}_{\mu\nu} = c_1 g^{\alpha\beta} g_{\mu\nu} + c_2 \delta_\mu^\alpha \delta_\nu^\beta + c_3 \delta_\nu^\alpha \delta_\mu^\beta + c_4 A^\mu A^\nu g_{\mu\nu}$$

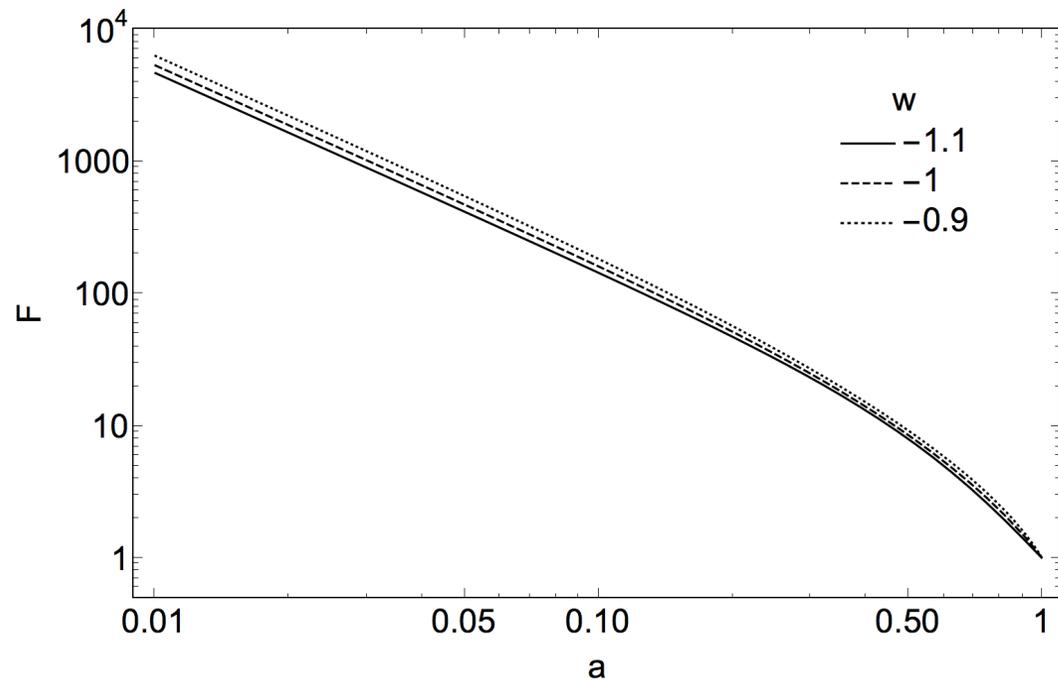
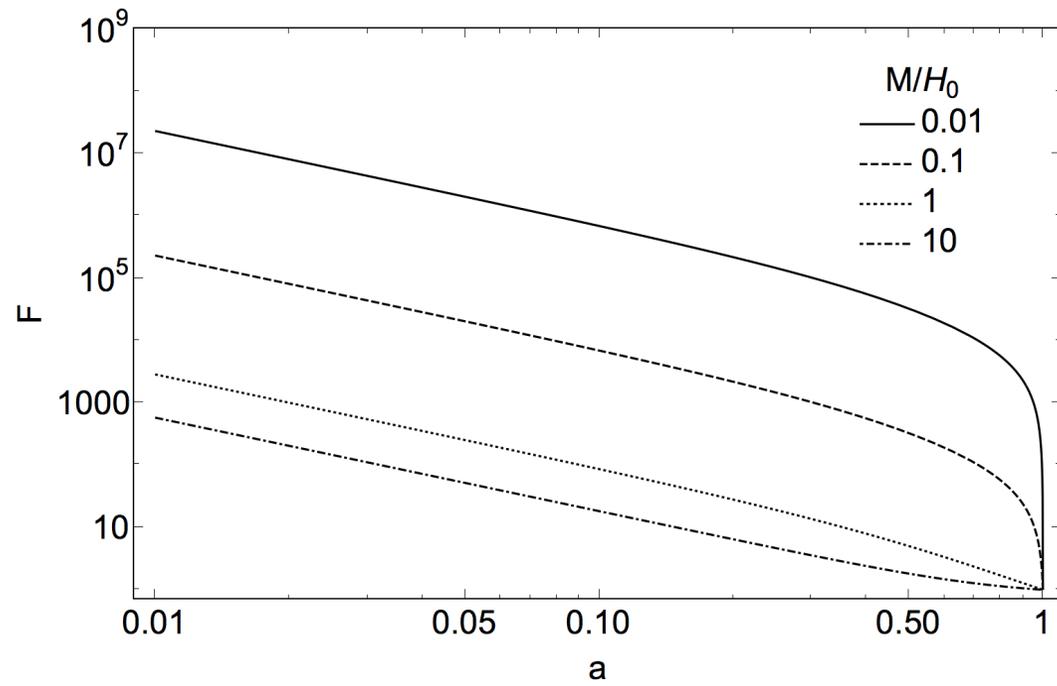
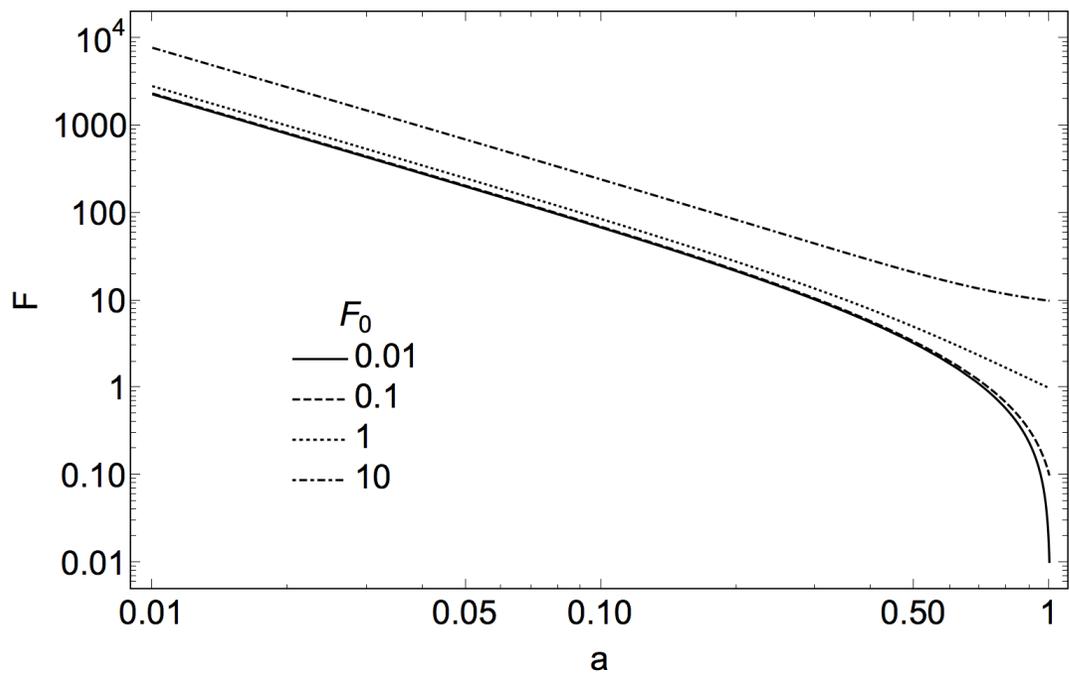
- A new vector field called the Aether.
- Constrained to be of time-like unit norm. This breaks Lorentz invariance.
- Similar to k-essence, the kinetic terms are generalized to a free function.
- “Kinetic terms” + “Potential terms”

Background cosmology

- What is F ? It would be nice if we could specify some other observational parameters and obtain F from them e.g. designer $F(R)$.

$$(1+w)(2\mathcal{K}\mathcal{F}_{\mathcal{K}} - \mathcal{F}) = (2\mathcal{K}\mathcal{F}_{\mathcal{K}\mathcal{K}} + \mathcal{F}_{\mathcal{K}}) \left[\mathcal{K} + \frac{1}{2}\alpha w (2\mathcal{K}\mathcal{F}_{\mathcal{K}} - \mathcal{F}) \right]$$

- Since the background is forced to match Λ CDM, the perturbations are now important.
- Gives us something to check for (e.g. CMB, weak lensing, ...) without being contaminated by differences in background evolution.



Linear perturbations

- We have the dynamics at background level. Now we want the dynamics of the perturbations.
- We need to specify the metric, the Aether, and the gauge.

$$ds^2 = a^2(\tau) \left[-(1 + 2\Psi)d\tau^2 + (\delta_{ij} + h_{ij})dx^i dx^j \right] \quad \delta A^\mu = \frac{1}{a}(-\Psi, \partial^i V)$$

New degree of freedom

- We can now go ahead and compute all the equations of motion.
- How do we parameterize this theory (and other theories) at the level of linear perturbations?

Equation of State approach

Battye and Peasron, 2013

- Modified gravity or new fluid component?
- The EoS approach treats them as the same. It recasts any theory into an effective fluid.
- Ultimately a matter of interpretation.
- The EoS approach eliminates the new degrees of freedom induced by the modified gravity theory.

Equation of State approach

- At background, need an equation of state to close the equations of motion, $w(a)$.
- What about for linear perturbations?

2 equations of state to compute

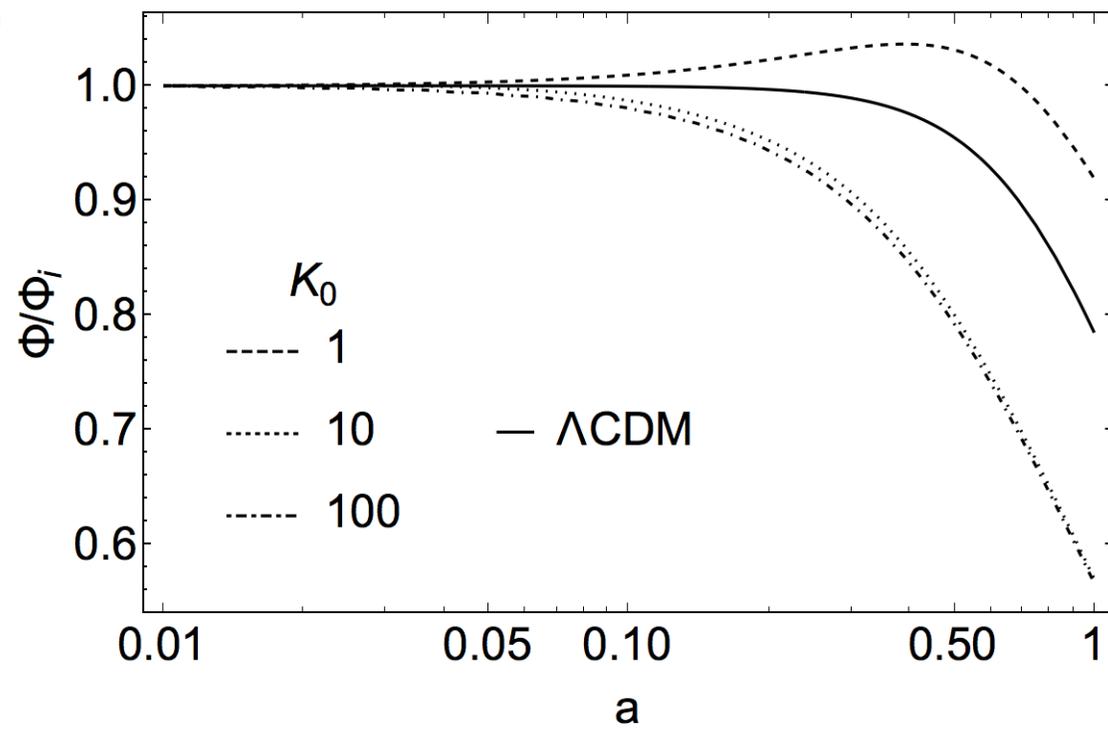
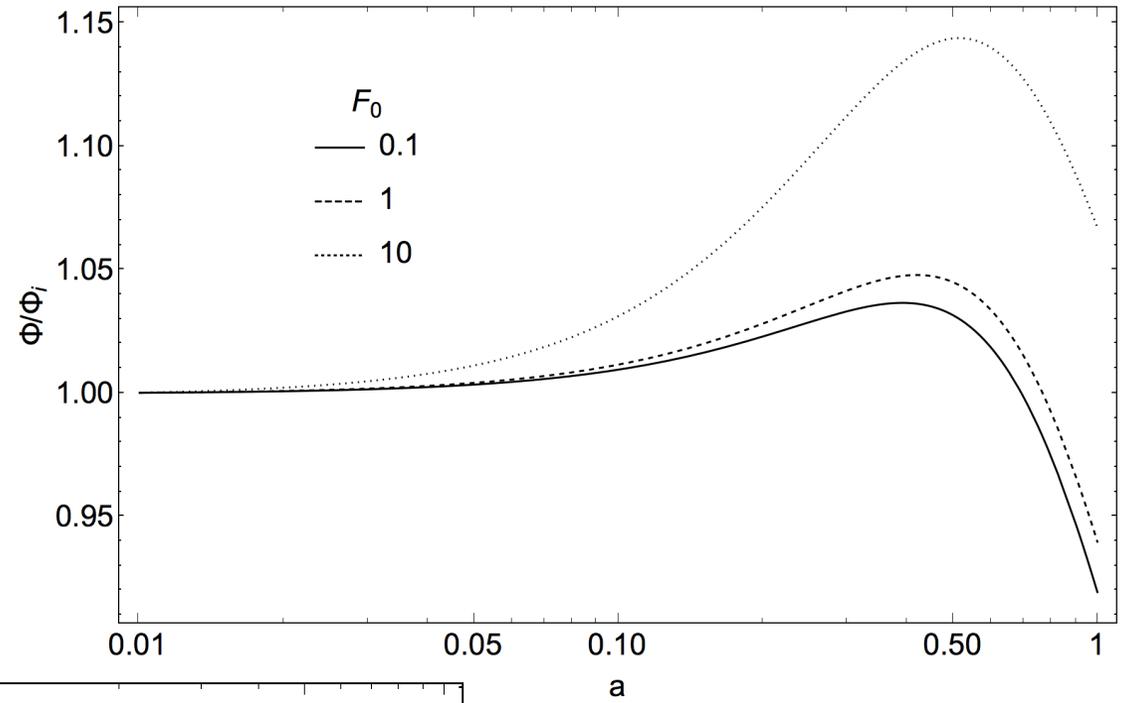
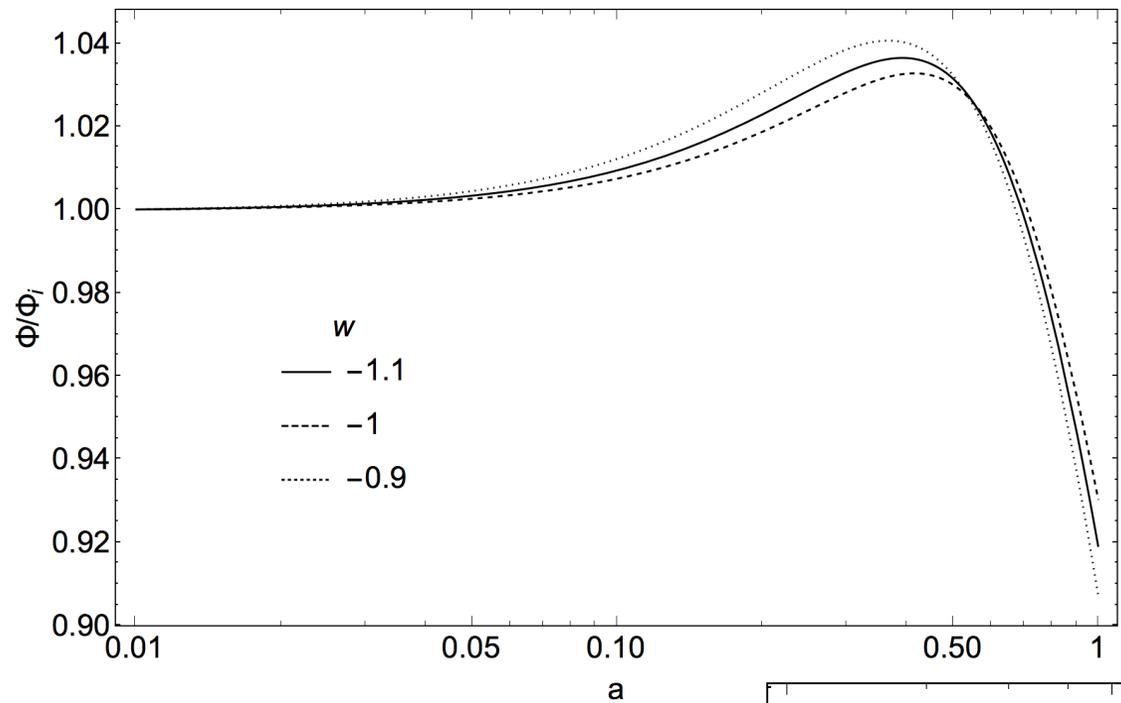
$$\left(\frac{\dot{\delta}}{1+w} \right) = -k^2 \theta^S - \frac{1}{2} \dot{h} - \frac{3\mathcal{H}w}{1+w} \Gamma,$$

$$(1+w)\dot{\theta}^S = \mathcal{H}(1+w) \left(3 \frac{dP}{d\rho} - 1 \right) \theta^S + \frac{dP}{d\rho} \delta + w\Gamma + \frac{2}{3} w \Pi^S.$$

- If we can somehow specify $\Gamma = \Gamma(\delta, \theta^S, \dot{h}, \eta, \dots)$ and $\Pi^S = \Pi^S(\delta, \theta^S, \dot{h}, \eta, \dots)$ only then these equations close.

Equation of State approach

- Eventually will obtain:
$$w\Pi^S = c_{\Pi\Delta_{de}}\Delta_{de} + c_{\Pi\Theta_{de}}\hat{\Theta}_{de} + c_{\Pi\Delta_m}\Delta_m + c_{\Pi\Theta_m}\hat{\Theta}_m$$
$$w\Gamma = c_{\Gamma\Delta_{de}}\Delta_{de} + c_{\Gamma\Theta_{de}}\hat{\Theta}_{de} + c_{\Gamma\Delta_m}\Delta_m + c_{\Gamma\Theta_m}\hat{\Theta}_m$$
- The approach is model independent.
- Any theory dependence is packaged in the c coefficients.
- This has been done for example in $F(R)$ (Battye et al., 2015, and Pace, in prep.), Quintessence and k-essence, Kinetic gravity braiding, and now Generalized Einstein-Aether.



Conclusions

- Generalized Einstein-Aether can give consistent results at background.
- Ultimately, the perturbations will tell us if this gives a viable alternative.
- There are many ways to parameterize theories at the perturbed level.
- We opt for the EoS approach which gives an elegant and straightforward way of solving the equations of motion.
- The two equations of state together with the equations of motion are all we need.